

Hadron Spectroscopy

Alex Bondar

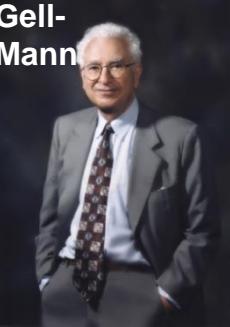
BINP, Novosibirsk

Novosibirsk State University



(Ljubljana, August 18, 2015)

Gell-Mann



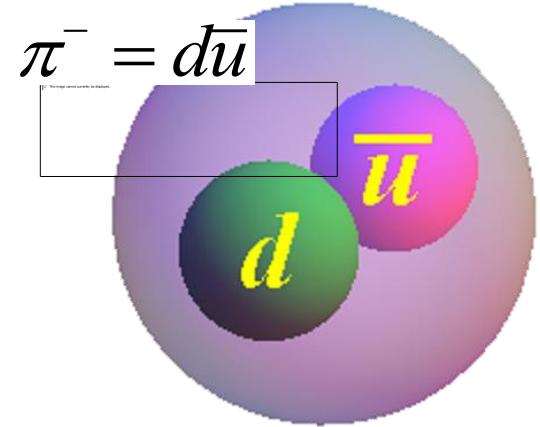
Constituent Quark Model



Mesons are bound states of quark and anti-quark:

$$\pi^+ = u\bar{d} \quad \pi^0 = \frac{1}{\sqrt{2}}(u\bar{u} - d\bar{d}) \quad \pi^- = d\bar{u}$$

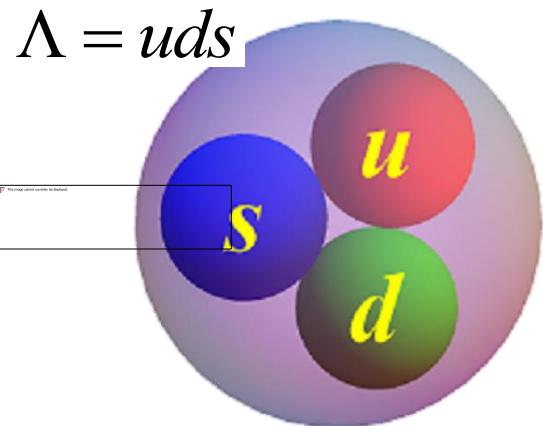
$$K^+ = u\bar{s} \quad K^0 = d\bar{s} \quad \bar{K}^0 = s\bar{d} \quad K^- = s\bar{u}$$



Baryons are bound states of 3 quarks:

$$p = uud \quad n = udd \quad \Lambda = uds$$

$$\bar{p} = \bar{u}\bar{u}\bar{d} \quad \bar{n} = \bar{u}\bar{d}\bar{d} \quad \bar{\Lambda} = \bar{u}\bar{d}\bar{s}$$



Quarkonium Basics

c, b -quarks are heavy:

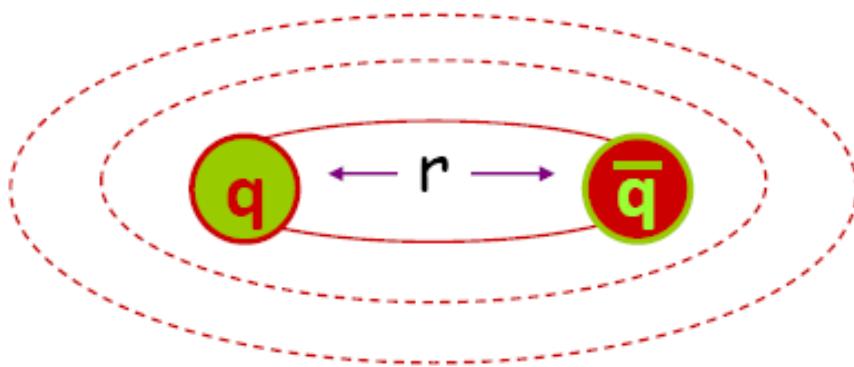
$$m_c \sim 1.5 \text{ GeV} \sim 1.6 m_p ;$$

$$m_b \sim 4.5 \text{ GeV} \sim 4.8 m_p ;$$

velocities are small:

$$v/c \sim 1/4 \text{ (for } b\bar{b}, v/c \sim 0.1)$$

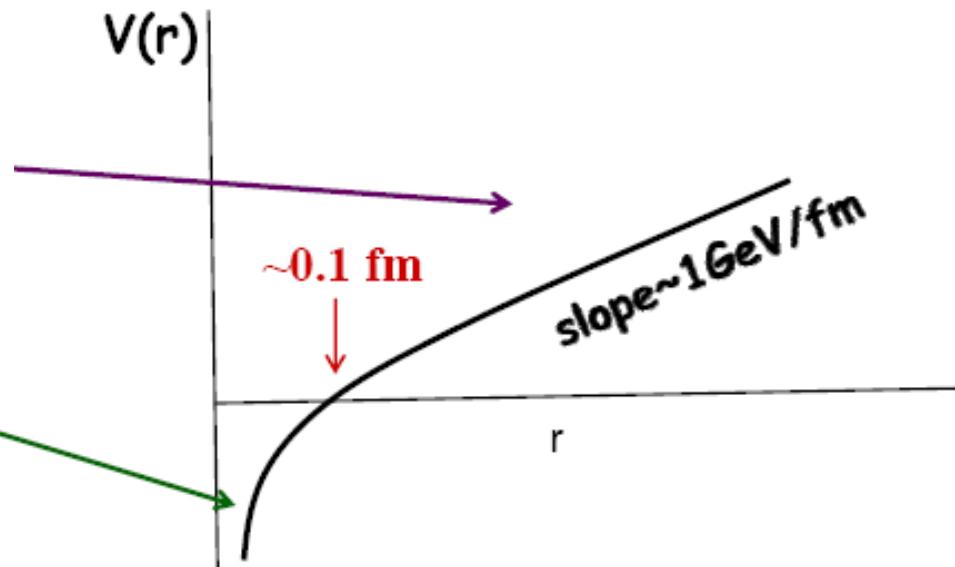
non-relativistic QM applies



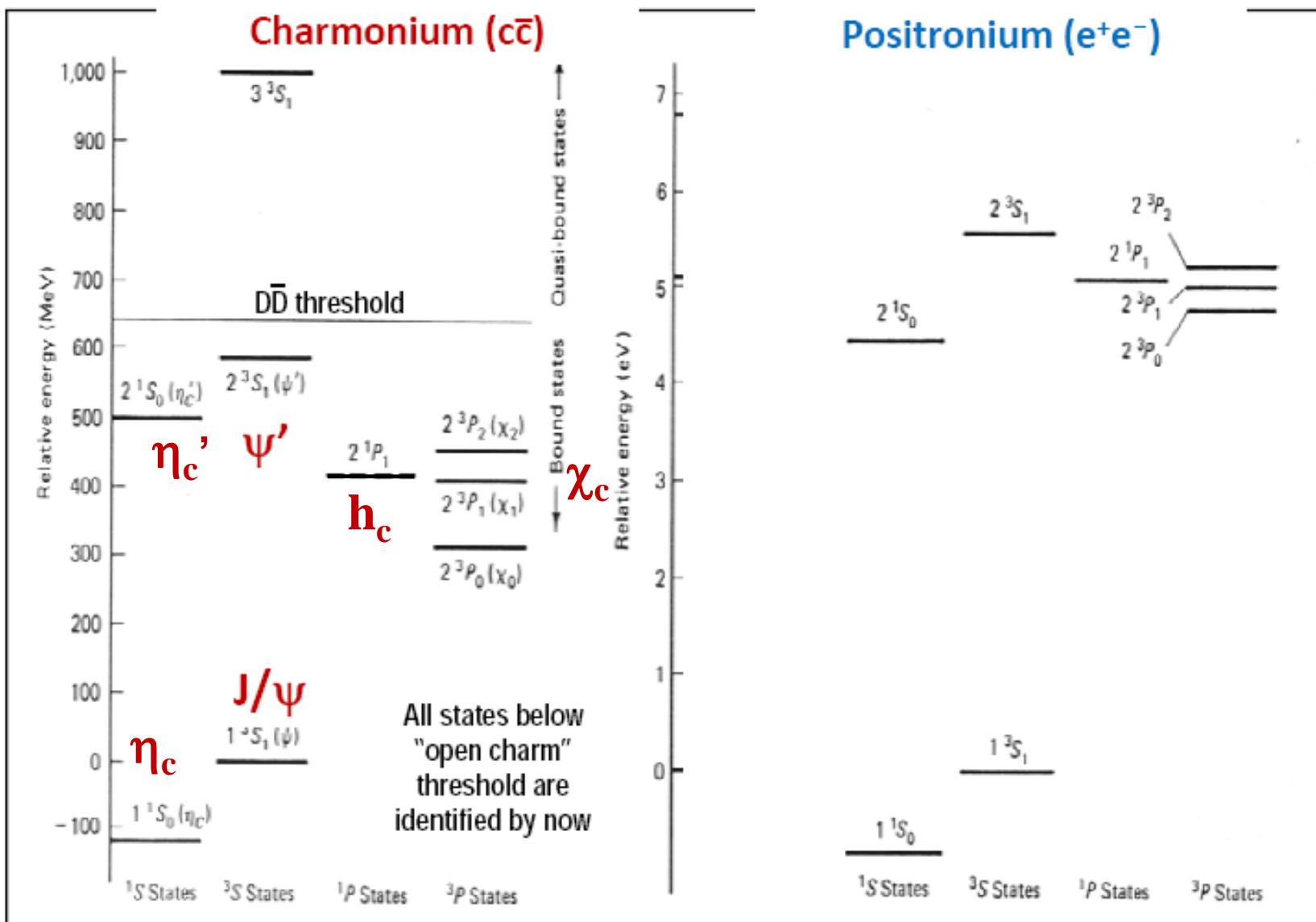
$$-\frac{\hbar^2}{2m_r} \nabla^2 \Psi + V(r) \Psi = E \Psi$$

linear "confinement"
long distance component

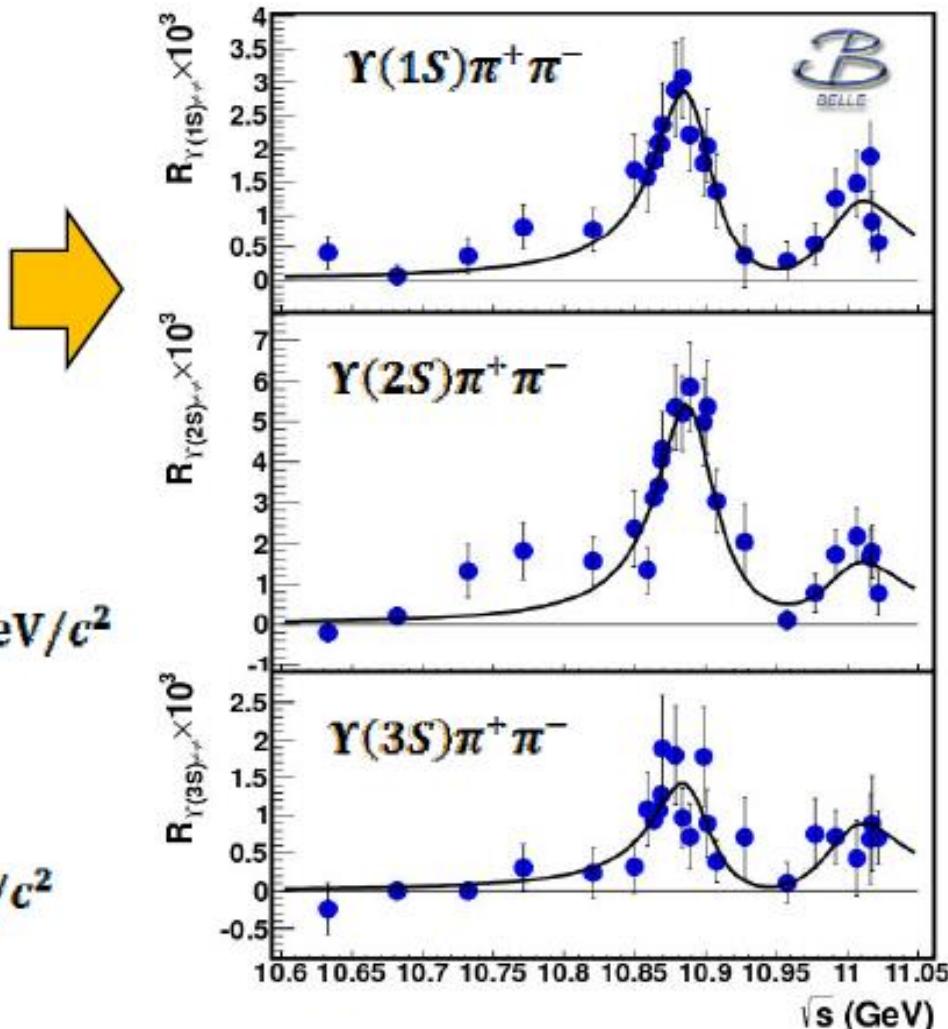
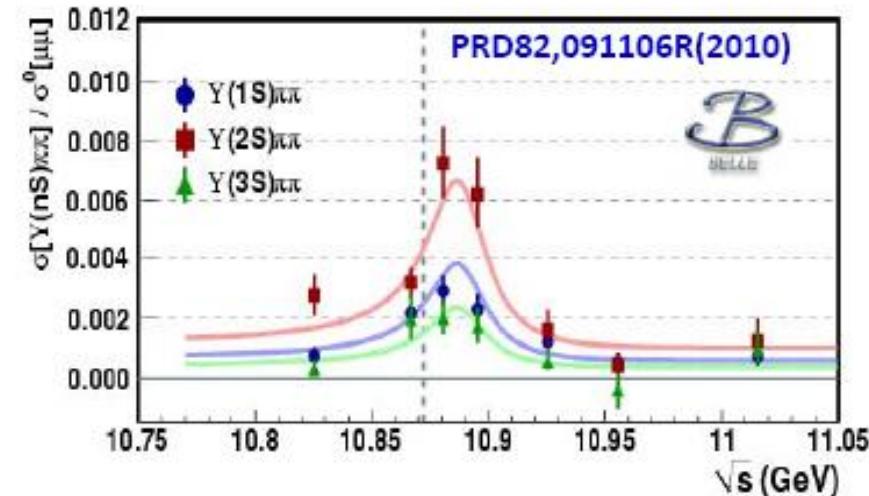
$1/r$ "coulombic"
short distance component



Quarkonium vs Positronium



Anomalous $\Upsilon(5S) \rightarrow \Upsilon(nS)\pi^+\pi^-$ Rates



$\Upsilon(10680)$:

$$M(\Upsilon(10680)) = 10891.1 \pm 3.2^{+0.6}_{-1.5} \text{ MeV}/c^2$$

$$\Gamma(\Upsilon(10680)) = 53.7^{+7.1+0.9}_{-5.6-5.4} \text{ MeV}$$

$\Upsilon(11020)$:

$$M(\Upsilon(11020)) = 10987.5^{+6.4+9.0}_{-2.5-2.1} \text{ MeV}/c^2$$

$$\Gamma(\Upsilon(11020)) = 61^{+9+2}_{-19-20} \text{ MeV}$$

$$\Delta\phi = -1 \pm 0.4^{+1.0}_{-0.1} \text{ rad}$$

hep-ex :1501.01137

Energy dependence of $h_b(1P,2P)\pi^+\pi^-$

Y(10680):

$$M(Y(10680)) = 10884.7 \begin{array}{l} +3.2 \\ -2.9 \end{array} \begin{array}{l} +8.6 \\ -0.6 \end{array} \text{ MeV/c}^2$$

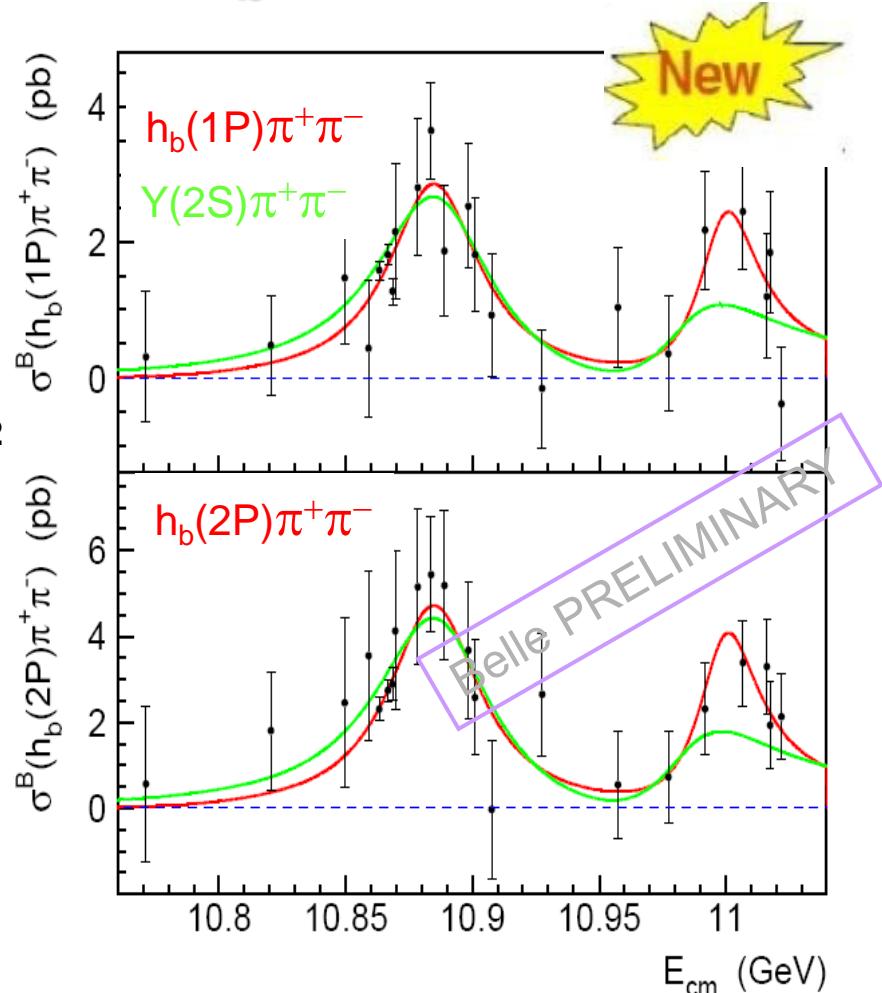
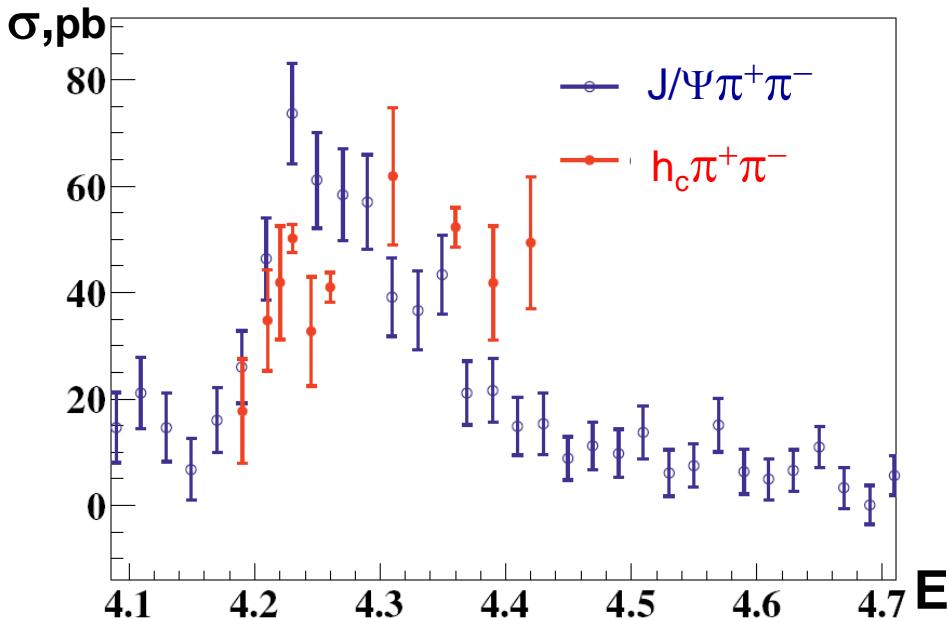
$$\Gamma(Y(10680)) = 44.2 \begin{array}{l} +11.9 \\ -7.8 \end{array} \begin{array}{l} +2.2 \\ -15.8 \end{array} \text{ MeV}$$

Y(11020):

$$M(Y(11020)) = 10998.6 \pm 6.1^{+16.1}_{-1.1} \text{ MeV/c}^2$$

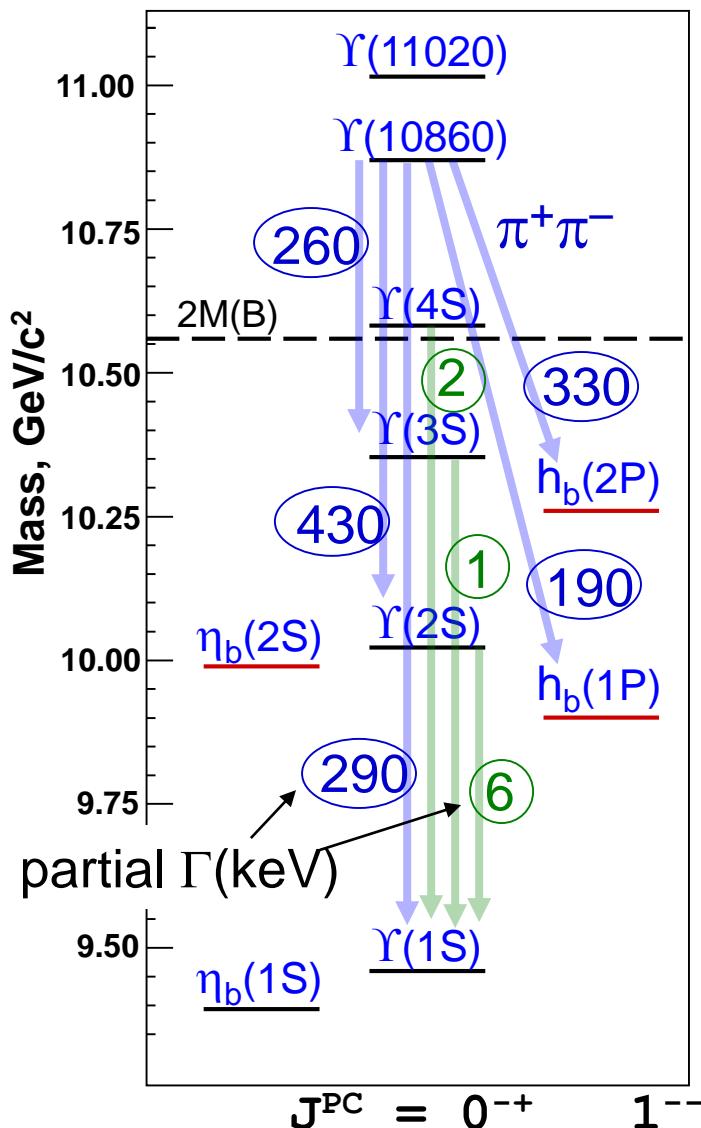
$$\Gamma(Y(11020)) = 29 \begin{array}{l} +20 \\ -12 \end{array} \begin{array}{l} +2 \\ -7 \end{array} \text{ MeV}$$

$$\Delta\phi = 0.3 \begin{array}{l} +1.0 \\ -1.5 \end{array} \text{ rad.}$$



$$\frac{\sigma(h_c\pi\pi)}{\sigma(h_b\pi\pi)} = \frac{m_{Y(5S)}^2 q_c^2 \Gamma_{Y(5S)}}{m_{Y(4260)}^2 q_b^2 \Gamma_{Y(4260)}} \sim 7$$

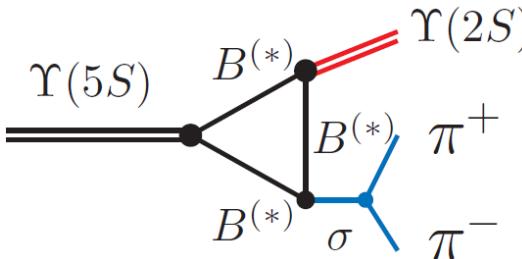
Anomalies in $\Upsilon(5S) \rightarrow (b\bar{b})\pi^+\pi^-$ transitions



Belle: PRL100, 112001 (2008)

$\Gamma[\Upsilon(5S) \rightarrow \Upsilon(1,2,3S)\pi^+\pi^-] \gg \Gamma[\Upsilon(4,3,2S) \rightarrow \Upsilon(1S)\pi^+\pi^-]$

⇐ Rescattering of on-shell $B^{(*)}\bar{B}^{(*)}$?



Belle: PRL108, 032001 (2012)

$\Upsilon(5S) \rightarrow h_b(1,2P)\pi^+\pi^-$ are not suppressed



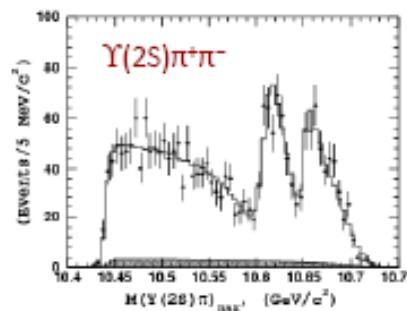
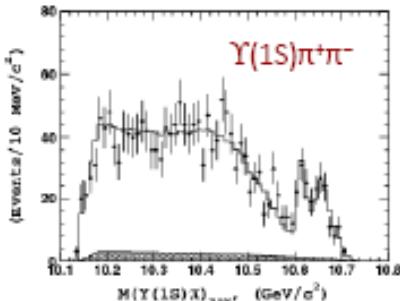
1^{-+}

expect suppression $\sim \Lambda_{\text{QCD}}/m_b$
~~Heavy Quark Symmetry~~

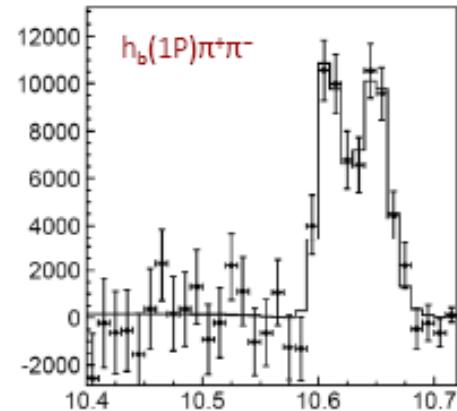
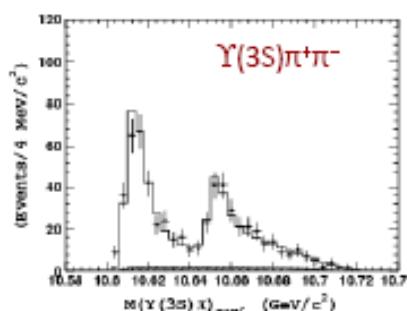
h_b production mechanism? ⇒ Study resonant structure in $h_b(mP)\pi^+\pi^-$

Properties of $Z_b(10610)$ & $Z_b(10650)$

Z_b^\pm Observed in five different modes:



PRL 108, 122001(2012)



$Z_b(10610)$

$Z_b(10650)$

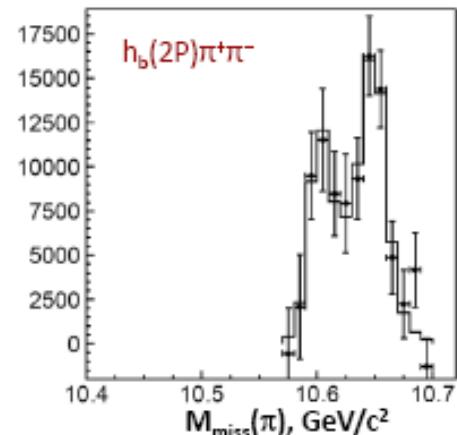
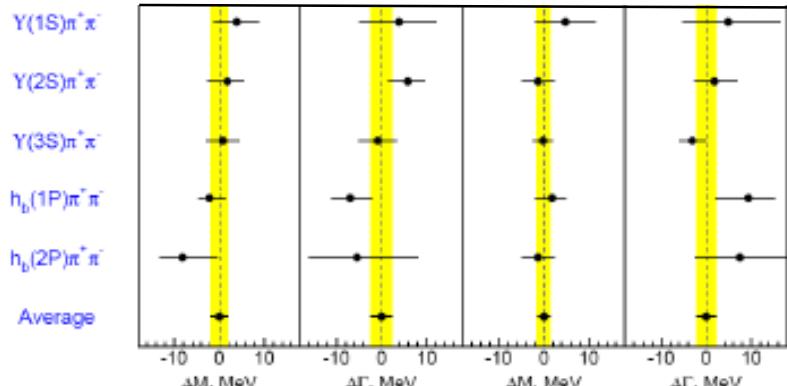
Average for Z_b^\pm :

$$\langle M_1 \rangle = 10607.2 \pm 2.0 \text{ MeV}$$

$$\langle \Gamma_1 \rangle = 18.4 \pm 2.4 \text{ MeV}$$

$$\langle M_2 \rangle = 10652.2 \pm 1.5 \text{ MeV}$$

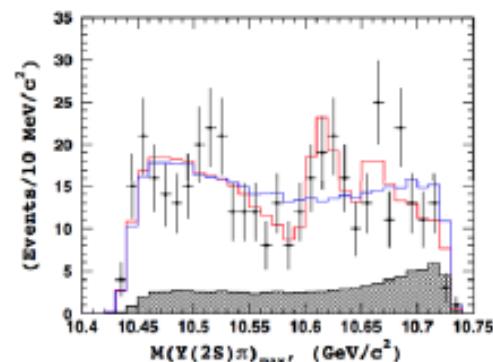
$$\langle \Gamma_2 \rangle = 11.5 \pm 2.2 \text{ MeV}$$



Z_b^0 Results:

$$\langle M_1 \rangle = 10609 \pm 7 \pm 6 \text{ MeV}$$

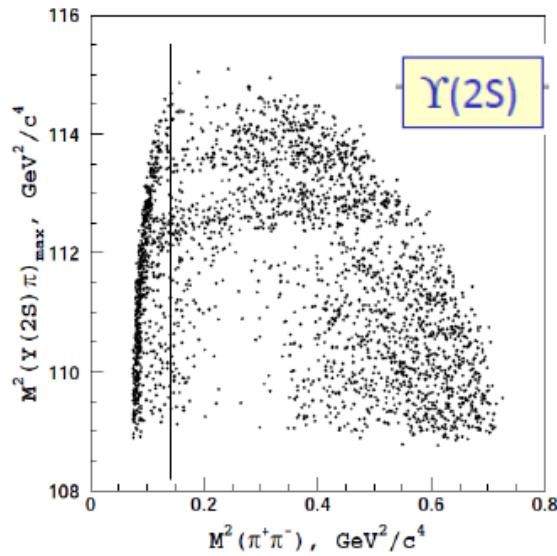
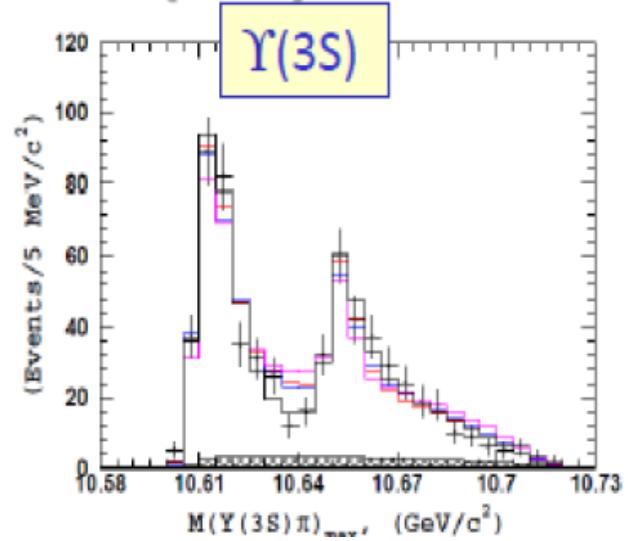
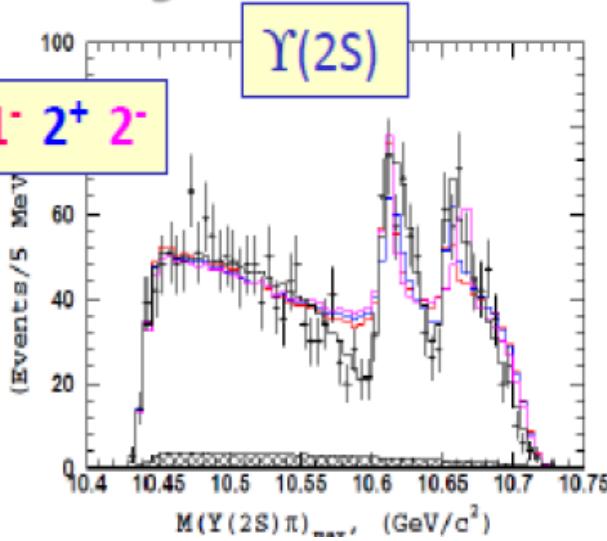
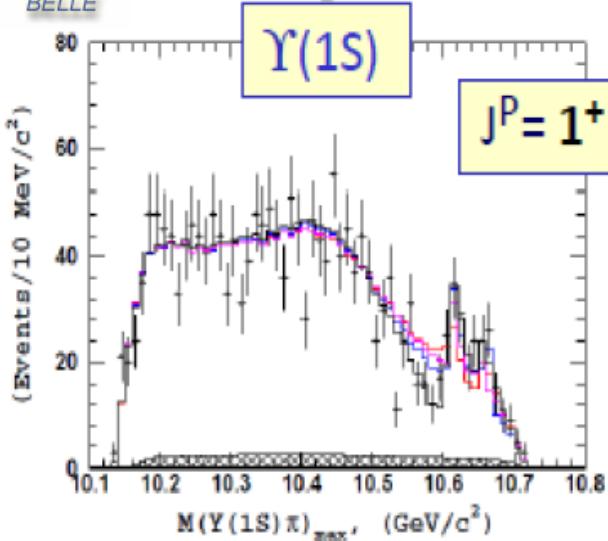
Consistent with Z_b^\pm



$$M_1 - M_B - M_{B^*} = 2.36 \pm 2.1 \text{ MeV}$$

$$M_2 - M_{B^*} - M_{B^*} = 1.8 \pm 1.75 \text{ MeV}$$

Amplitude analysis of $e^+e^- \rightarrow \Upsilon(nS)\pi^+\pi^-$



	$Z_b(10650)$	1^+	1^-	2^+	2^-
$Z_b(10610)$					
1^+		0 (0)	60 (33)	42 (33)	77 (63)
1^-		226 (47)	264 (73)	224 (68)	277 (106)
2^+		205 (33)	235 (104)	207 (87)	223 (128)
2^-		289 (99)	319 (111)	321 (110)	304 (125)

Spin-parity of $Z_b(10610)$ and $Z_b(10650)$ is 1^+ .
 All other $J^P < 3$ are excluded.

Molecule?

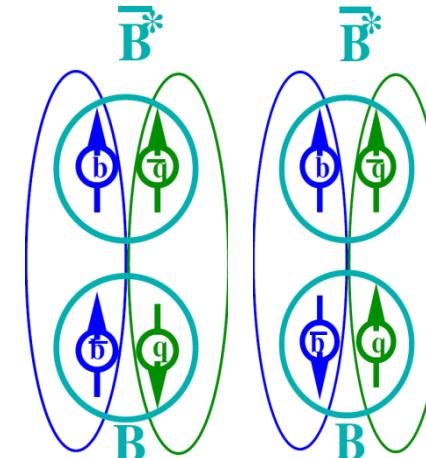
Heavy quark structure in Z_b

A.B.,A.Garmash,A.Milstein,R.Mizuk,M.Voloshin PRD84 054010 (arXiv:1105.4473)

Wave func. at large distance – $B(*)B^*$

$$|Z'_b\rangle = \frac{1}{\sqrt{2}} \mathbf{0}_{bb}^- \otimes \mathbf{1}_{Qq}^- - \frac{1}{\sqrt{2}} \mathbf{1}_{bb}^- \otimes \mathbf{0}_{Qq}^-$$

$$|Z_b\rangle = \frac{1}{\sqrt{2}} \mathbf{0}_{bb}^- \otimes \mathbf{1}_{Qq}^- + \frac{1}{\sqrt{2}} \mathbf{1}_{bb}^- \otimes \mathbf{0}_{Qq}^-$$



Explains

- Why $h_b\pi\pi$ is unsuppressed relative to $\Upsilon\pi\pi$
- Relative phase ~ 0 for Υ and $\sim 180^\circ$ for h_b
- Production rates of $Z_b(10610)$ and $Z_b(10650)$ are similar
- Widths –”–
- Dominant decays to $B(*)B^*$

Other Possible Explanations

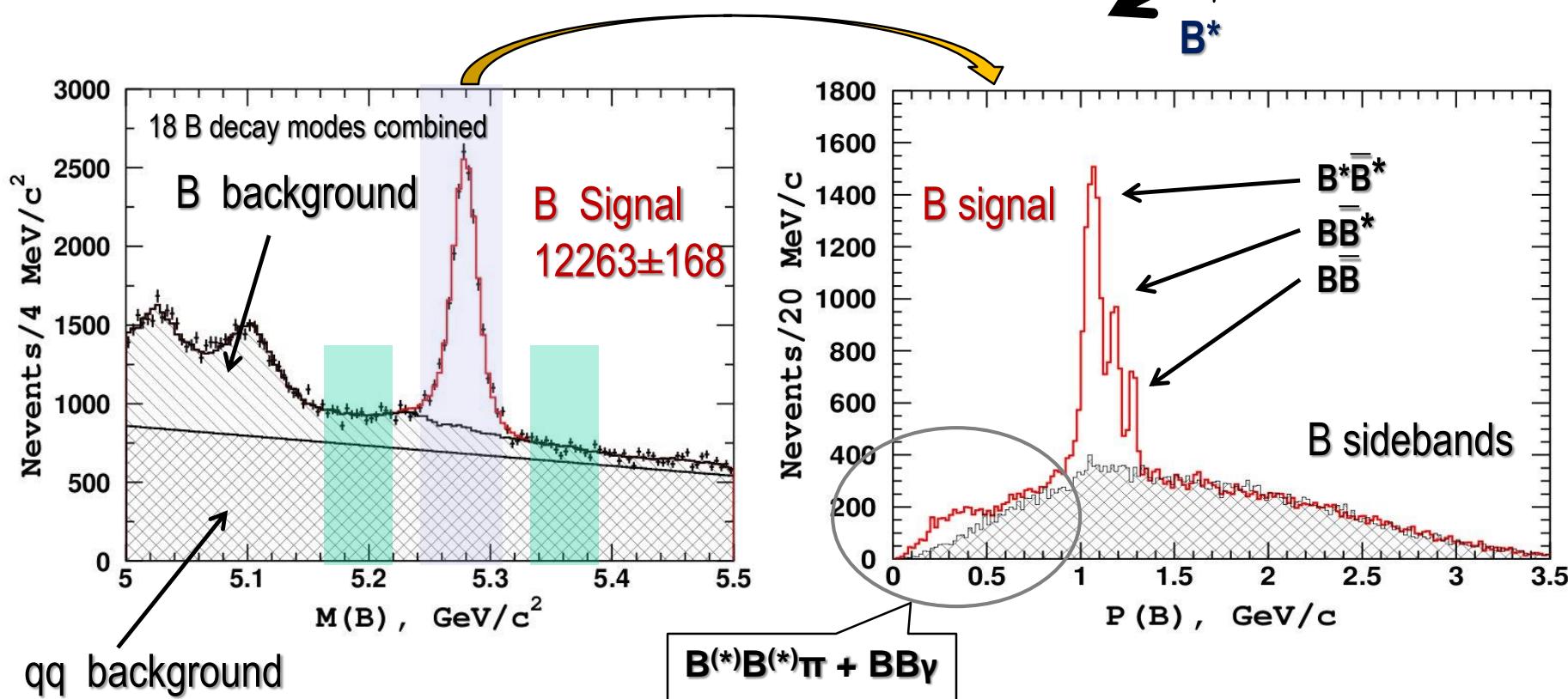
- Coupled channel resonances (I.V.Danilkin et al, arXiv:1106.1552)
- Cusp (D.Bugg Europhys.Lett.96 (2011),arXiv:1105.5492)
- Tetraquark (M.Karlener, H.Lipkin, arXiv:0802.0649)

$\Upsilon(5S) \rightarrow Z_b \pi^\pm$

Masses of the observed Z_b resonances are close to the BB^* and B^*B^* thresholds, respectively: branching fractions $Z_b \rightarrow B(*)B^*$ might be large (dominant).

Analysis strategy:

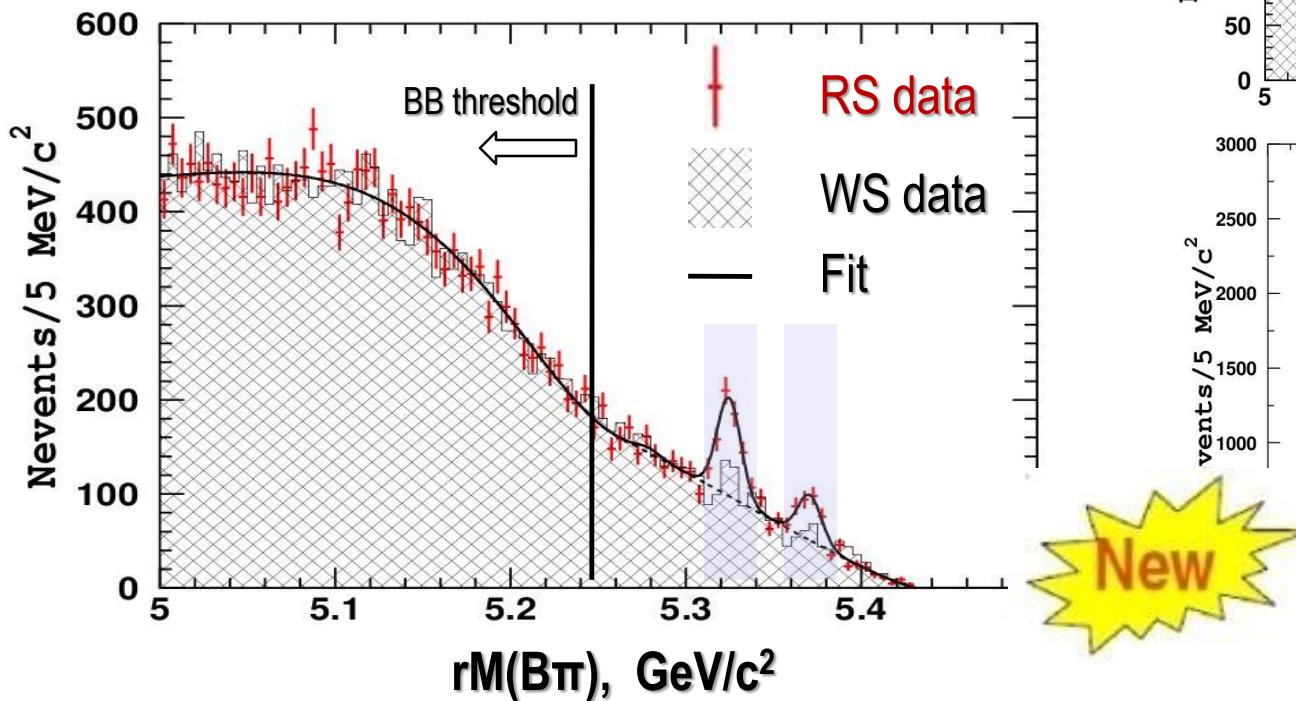
- Full reconstruction of one B meson



$\Upsilon(5S) \rightarrow Z_b \pi^\pm$

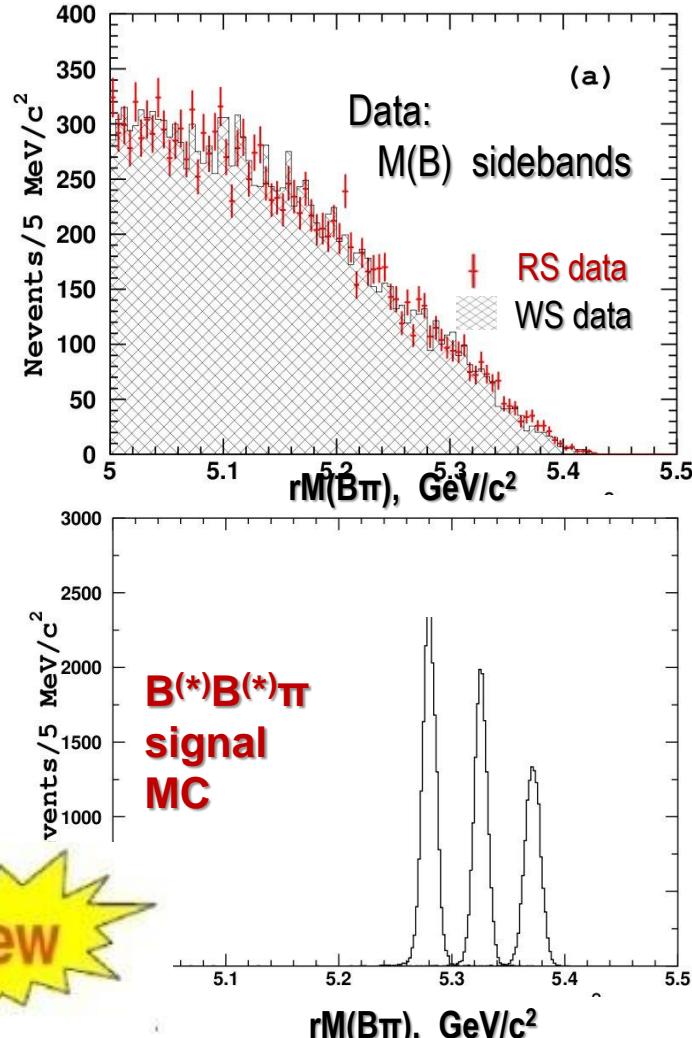
Analysis strategy:

- Combine the B with a charged pion from the rest of the event \rightarrow calculate recoil mass against the $B\pi$ system



$$\begin{aligned} N(BB\pi) &= 13 \pm 25 \\ N(BB^{(*)}\pi) &= 357 \pm 30 \\ N(B^{(*)}B^{(*)}\pi) &= 161 \pm 21 \end{aligned}$$

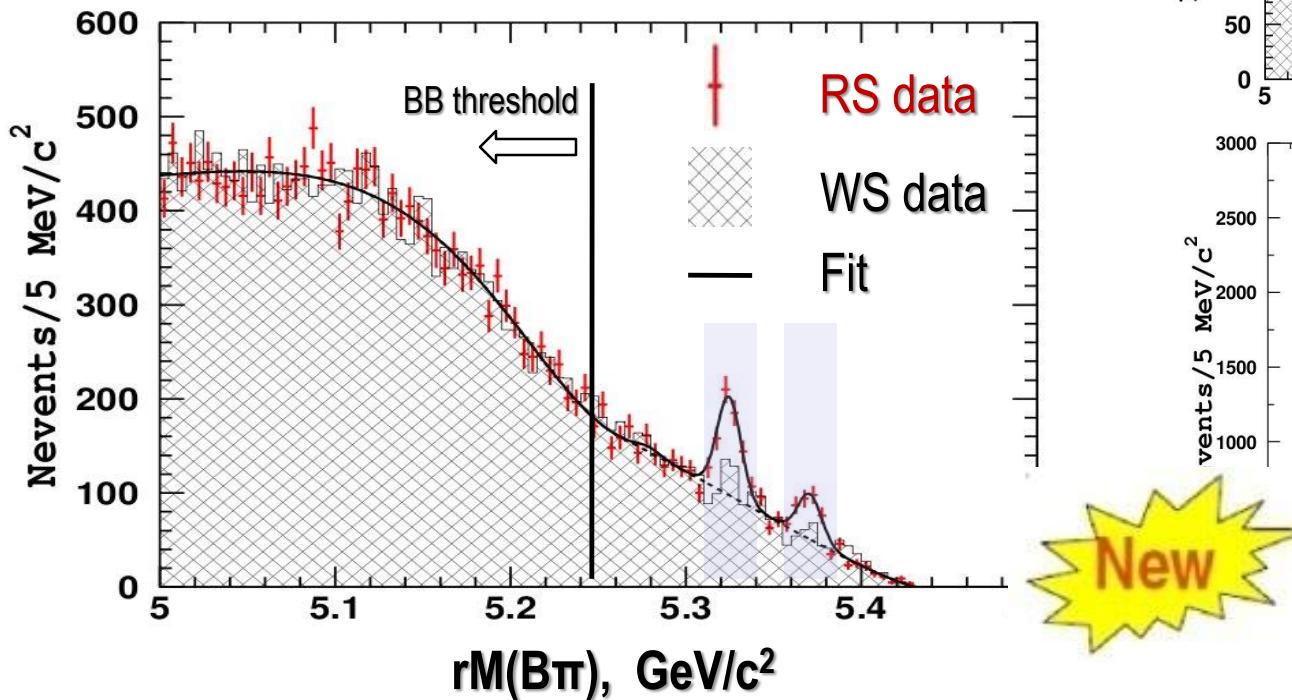
$$\begin{aligned} \sigma(e^+e^- \rightarrow BB\pi) &< \\ \sigma(e^+e^- \rightarrow BB^{(*)}\pi) &= \\ \sigma(e^+e^- \rightarrow B^{(*)}B^{(*)}\pi) &= \end{aligned}$$



$\Upsilon(5S) \rightarrow Z_b \pi^\pm$

Analysis strategy:

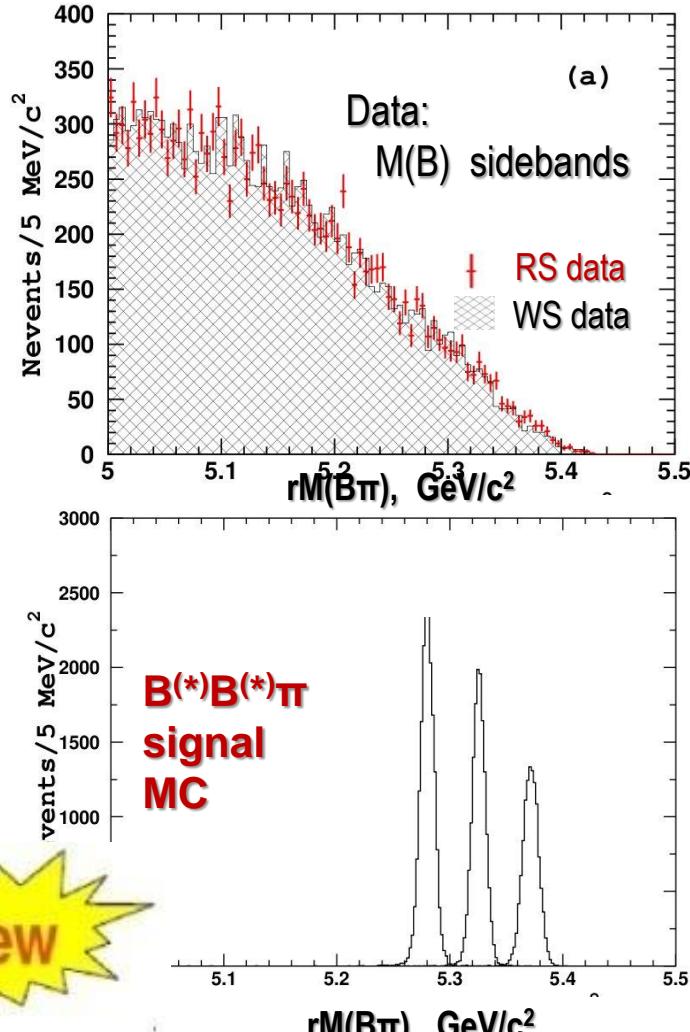
- Combine the B with a charged pion from the rest of the event → calculate recoil mass against the $B\pi$ system



$$\begin{aligned} N(BB\pi) &= 13 \pm 25 \\ N(BB^{(*)}\pi) &= 357 \pm 30 \\ N(B^{(*)}B^{(*)}\pi) &= 161 \pm 21 \end{aligned}$$

$$\begin{aligned} \sigma(e^+e^- \rightarrow BB\pi) \\ \sigma(e^+e^- \rightarrow BB^{(*)}\pi) \\ \sigma(e^+e^- \rightarrow B^{(*)}B^{(*)}\pi) \end{aligned}$$

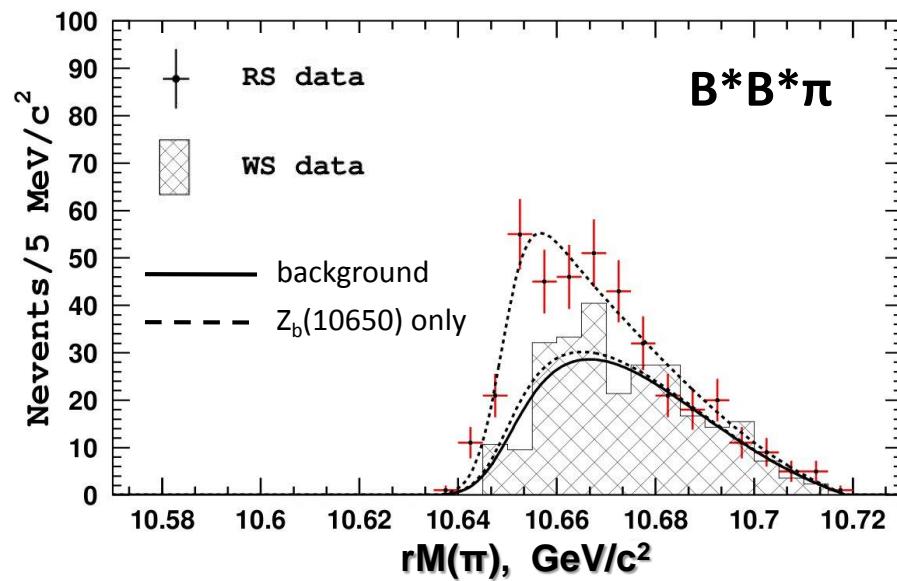
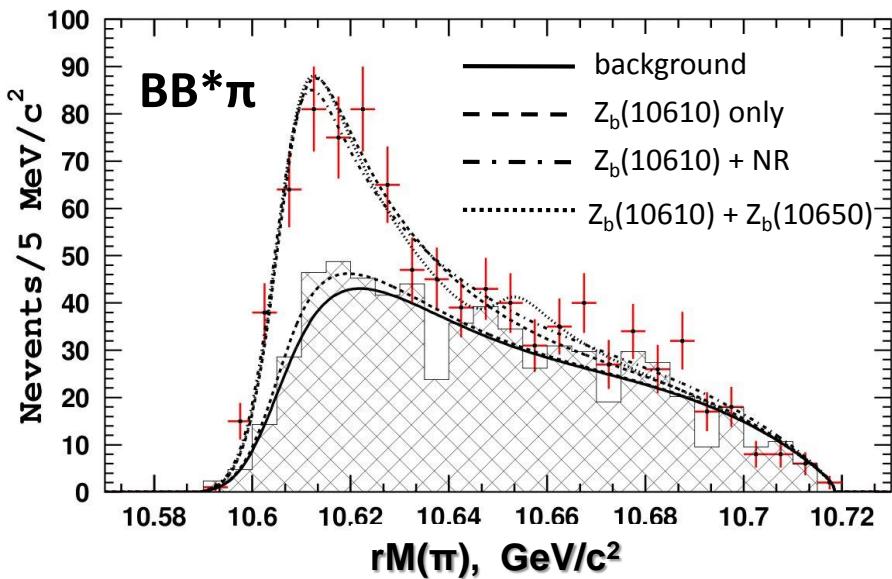
Oops... Sorry, Belle Collaboration
is not ready yet to show official
numbers.



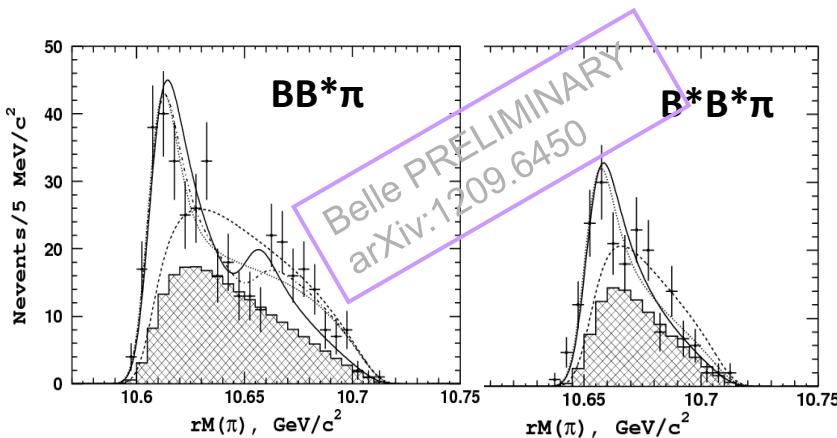
$\Upsilon(5S) \rightarrow Z_b \pi^\pm$



For events from 3-body sig. region \rightarrow recoil mass against primary π^-



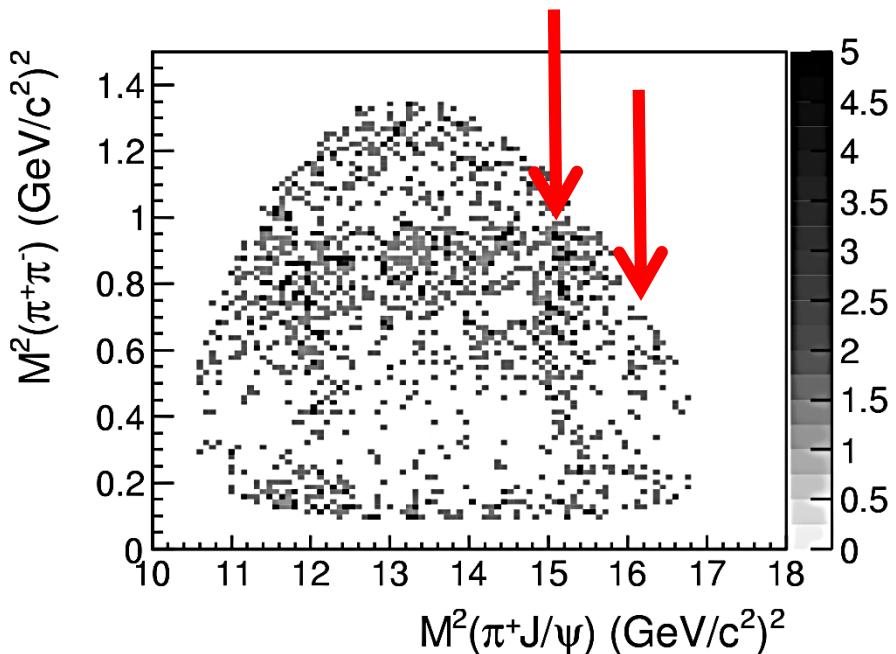
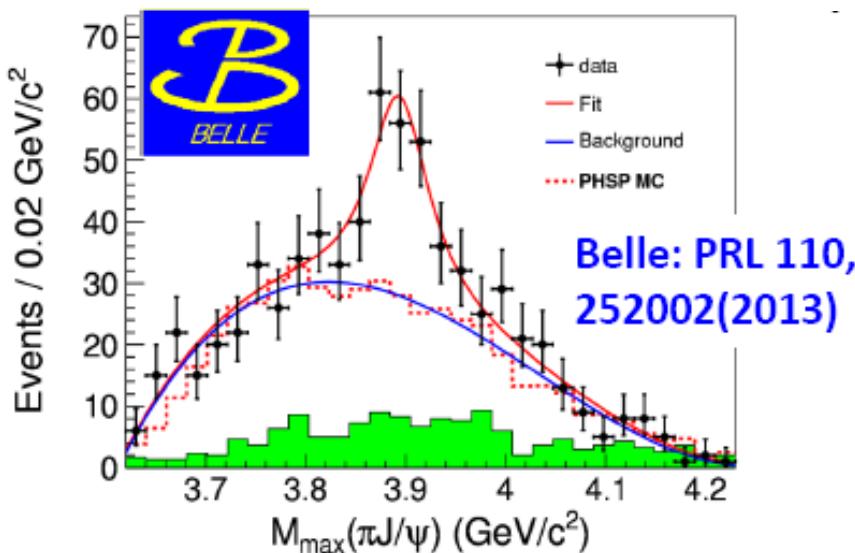
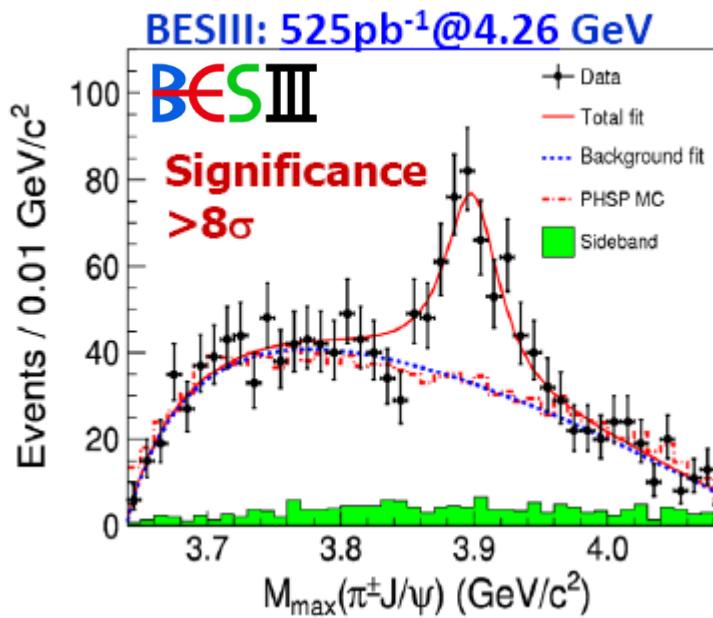
BB* π and B*B* π data fits well to just $Z_b(10610)$ and $Z_b(10650)$ signal, respectively



Assuming Z_b decays are saturated by already observed channels

B(*)B* channels dominate the Z_b decays

Observation of $Z_c(3900)$ at BESIII & Belle



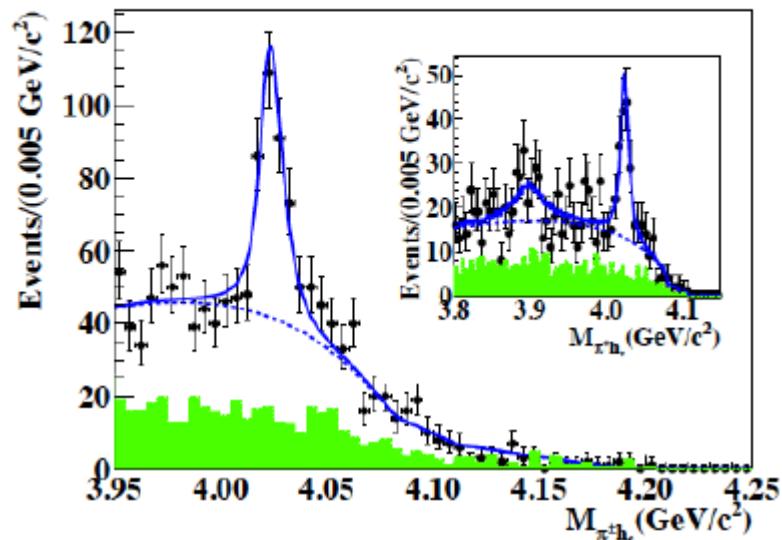
BESIII: PRL110, 252001 (2013)

- $M = 3899.0 \pm 3.6 \pm 4.9$ MeV
- $\Gamma = 46 \pm 10 \pm 20$ MeV
- 307 ± 48 events

The mass position is 24 MeV away from DD* threshold!

Observation of $Z_c(4020)$ in $e^+e^- \rightarrow h_c\pi^+\pi^-$

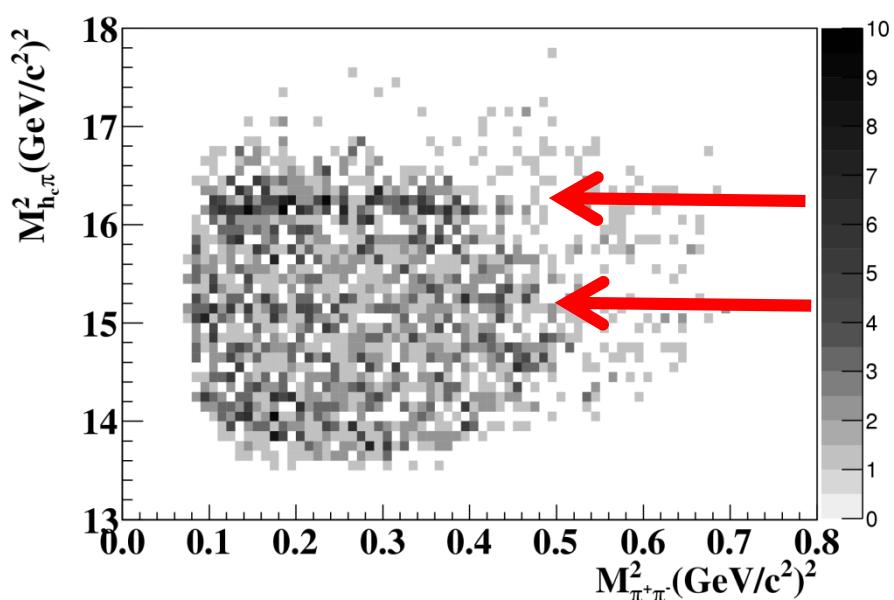
BESIII: PRL 111, 242001 (2013)



Simultaneous fit to
4.23/4.26/4.36 GeV data, 16 η_c
decay modes.

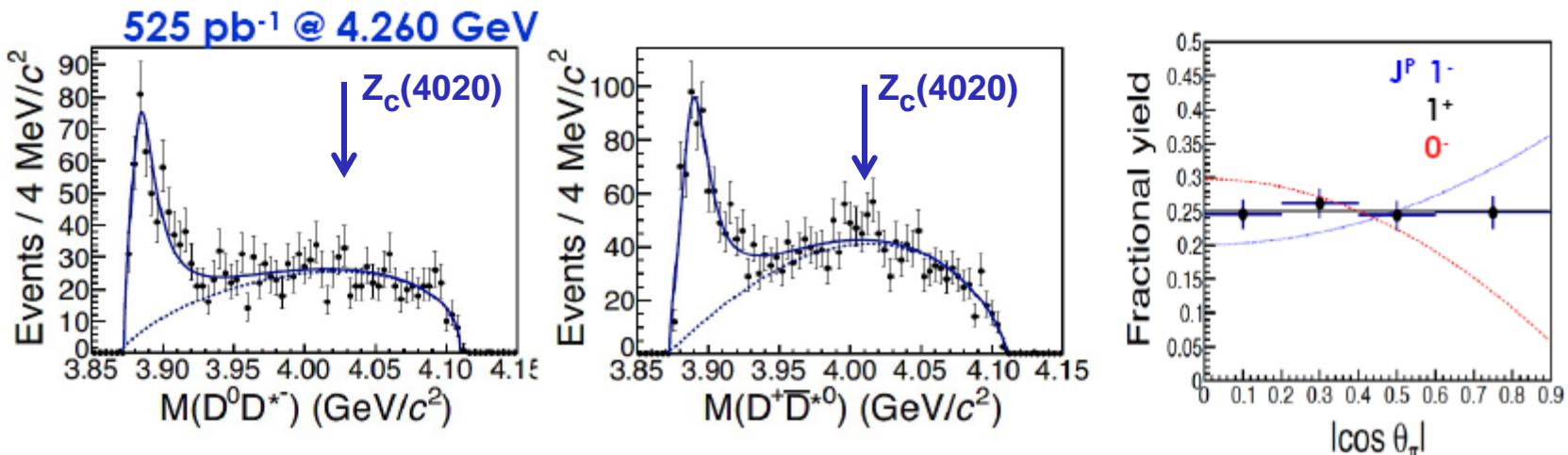
$$M = 4022.9 \pm 0.8 \pm 2.7 \text{ MeV}/c^2$$

$$\Gamma = 7.9 \pm 2.7 \pm 2.6 \text{ MeV}$$



Significance: 8.9σ ($Z_c(4020)$)
No significant $Z_c(3900)$ (2.1σ)

BESIII: PRL 112, 022001 (2014)



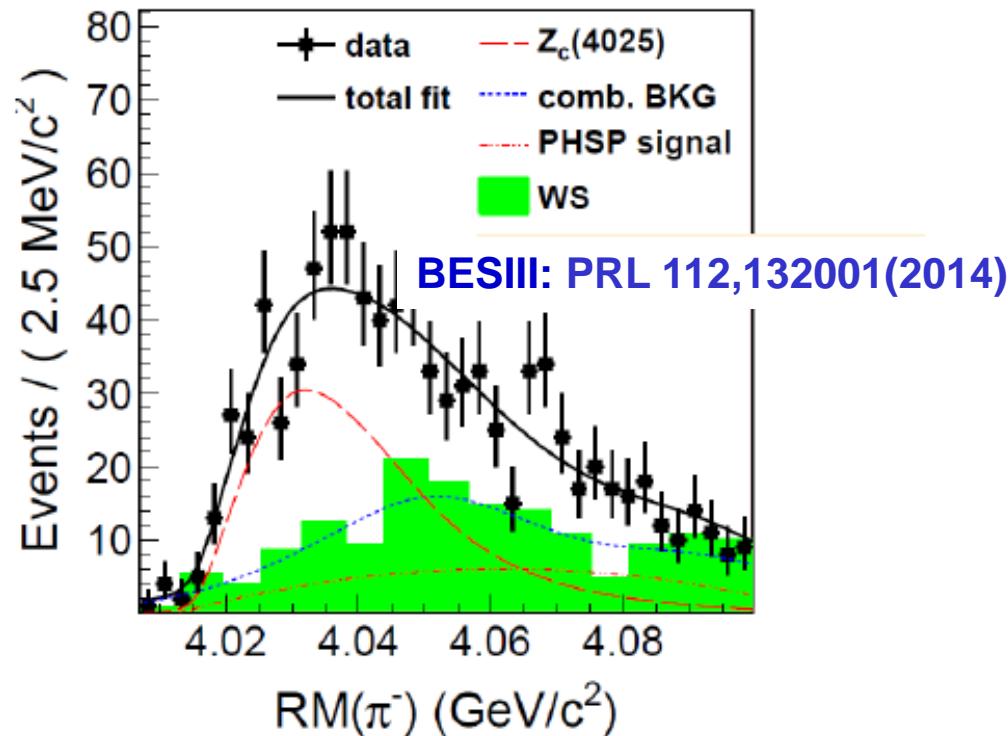
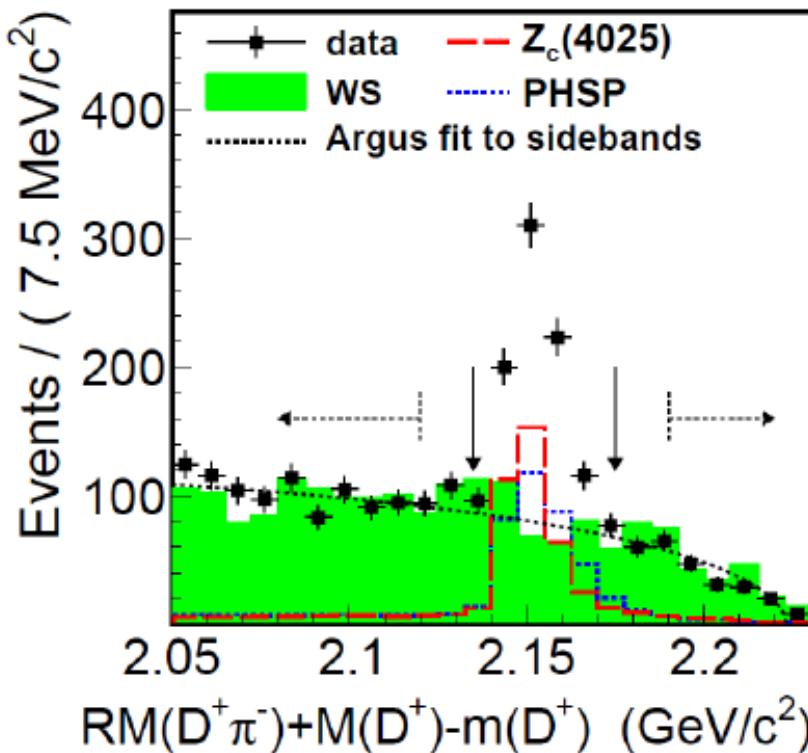
- $M = 3883.9 \pm 1.5 \pm 4.2 \text{ MeV}$; $\Gamma = 24.8 \pm 3.3 \pm 11.0 \text{ MeV}$
- $\sigma \times B = 85.3 \pm 6.6 \pm 22.0 \text{ pb}$
- **fits favor 1^+ distribution assumption**

fit with mass-dependent-width BW with phase space and efficiency correction

$$\frac{\Gamma(Z_c(3885) \rightarrow D\bar{D}^*)}{\Gamma(Z_c(3900) \rightarrow \pi J/\psi)} = 6.2 \pm 1.1 \pm 2.7$$

Assuming $Z_c(3885)$ due to $Z_c(3900)$

BESIII Observation of $Z_c(4025)$ in $e^+e^- \rightarrow \pi^\pm(D^*D^*)^\mp$



Fit to π^\pm recoil mass yields 401 ± 47 $Z_c(4025)$ events. $>10\sigma$

$$M(Z_c(4025)) = 4026.3 \pm 2.6 \pm 3.7 \text{ MeV}; \quad \Gamma(Z_c(4025)) = 24.8 \pm 5.6 \pm 7.7 \text{ MeV}$$

$$R = \frac{\sigma(e^+e^- \rightarrow \pi^\pm Z_c^\mp(4025) \rightarrow \pi^\pm (D^* \bar{D}^*)^\mp)}{\sigma(e^+e^- \rightarrow \pi^\pm (D^* \bar{D}^*)^\mp)} = (65 \pm 9 \pm 6)\%$$

$$\sigma(e^+e^- \rightarrow \pi^\pm (D^* \bar{D}^*)^\mp) = (137 \pm 9 \pm 15) \text{ pb}$$

Summary of the Z_c states

Channel	Mass (MeV/c ²)	Width (MeV)	
$\pi^\pm J/\psi$	$3899.0 \pm 3.6 \pm 4.9$	$46 \pm 10 \pm 20$	Close to D \bar{D}^* threshold (3875 MeV)
$(D \bar{D}^*)^\pm$	$3883.9 \pm 1.5 \pm 4.2$	$24.8 \pm 3.3 \pm 11.0$	
2 σ difference		1 σ difference	
$\pi^\pm h_c$	$4022.9 \pm 0.8 \pm 2.7$	$7.9 \pm 2.7 \pm 2.6$	Close to $D^* \bar{D}^*$ threshold (4017 MeV)
$(D^* \bar{D}^*)^\pm$	$4026.3 \pm 2.6 \pm 3.7$	$24.8 \pm 5.6 \pm 7.7$	
1 σ difference		2 σ difference	

- Near threshold, but all the measurements are the results of 1D fits without interference, to be improved by amplitude analysis;
- Couple to $D^*D^{(*)}$ final states stronger than to final states with charmonium;
- Necessary to measure the energy dependence of the $Z_c\pi^+$ production cross-sections;
- As well as $D^{(*)}D^*\pi$;

Bottomonium-like vs Charmonium-like states

- Z_b are very close to $\bar{B}B^*$, $B^*\bar{B}^*$ threshold
- $I^G J^{P(C)} = 1^+ 1^+ (-)$
- Observed both in the hidden-bottom modes: $\pi Y(1S, 2S, 3S)$, $\pi h_b(1P, 2P)$ and open-bottom modes: $\bar{B}B^*$, $B^*\bar{B}^*$
- $B(^*)\bar{B}^*$ dominate Z_b decays with the branching ratio 86% and 73%
- Energy behaviors $h_b(mP)\pi\pi$ and $Y(nS)\pi\pi$ are similar
- Z_c are very close to $\bar{D}D^*$, $D^*\bar{D}^*$ threshold
- $I^G J^{P(C)} = 1^+ 1^+ (-)$
- Observed both in the hidden-charm modes: $\pi J/\psi$, πh_c and open-charm modes: $\bar{D}D^*$, $D^*\bar{D}^*$
- $\bar{D}D^*$ dominates $Z_c(3900)$ decay; \bar{D}^*D^* dominates $Z_c(4020)$
- no $Z_c(3900)$ decay to πh_c ?
- no $Z_c(4020)$ decay to $\pi J/\psi$?
- Energy behaviors $h_c\pi\pi$ and $J/\psi\pi\pi$ are different

Z_b properties agree well with $B(^*)\bar{B}^*$ molecular structure, while in case of Z_c the situation is still unclear (masses too far above thresholds, no $J/\psi\pi^+$ at D^*D^* threshold)

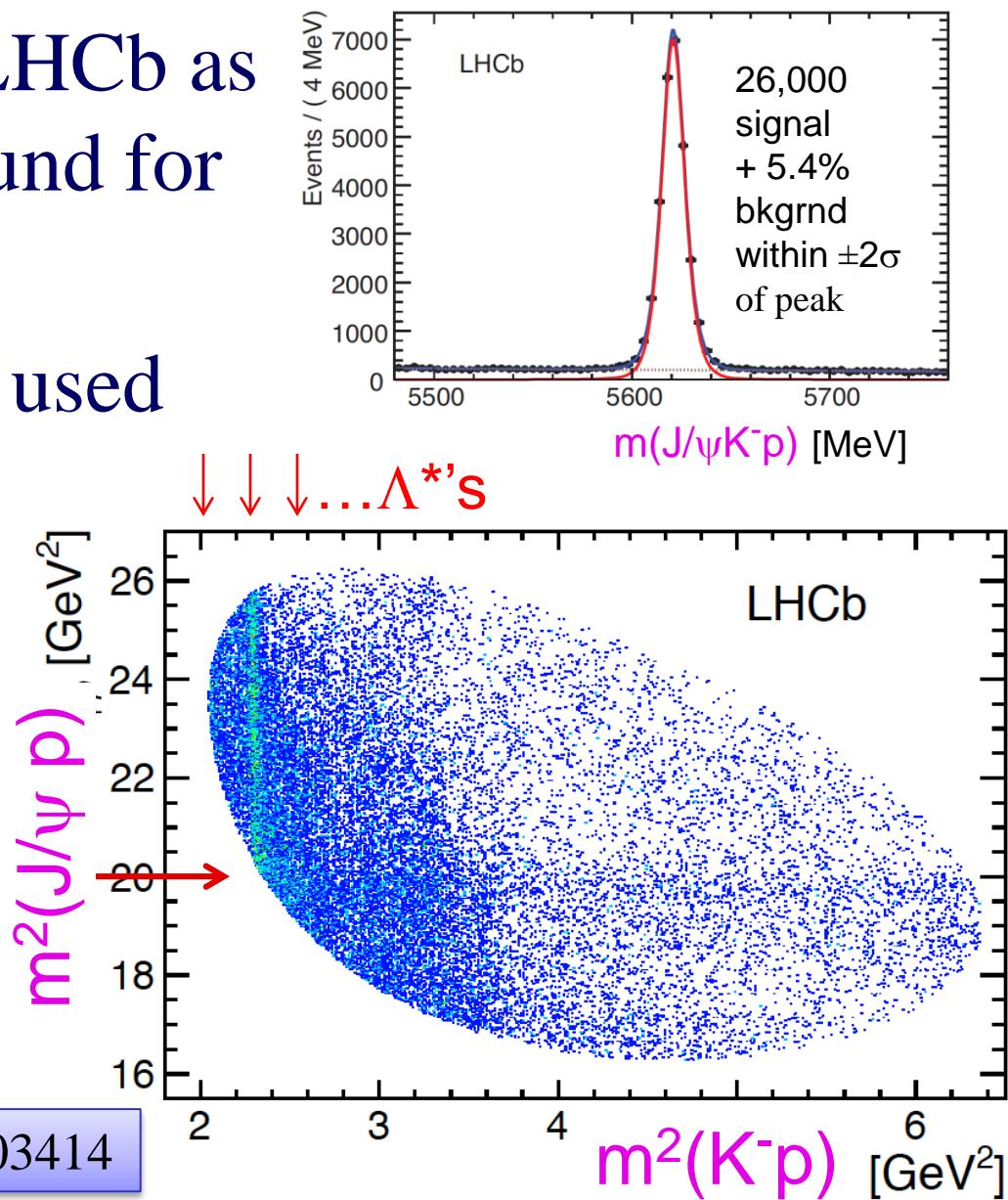
Pentaquark?

See for details Tomasz Skwarnicki presentation

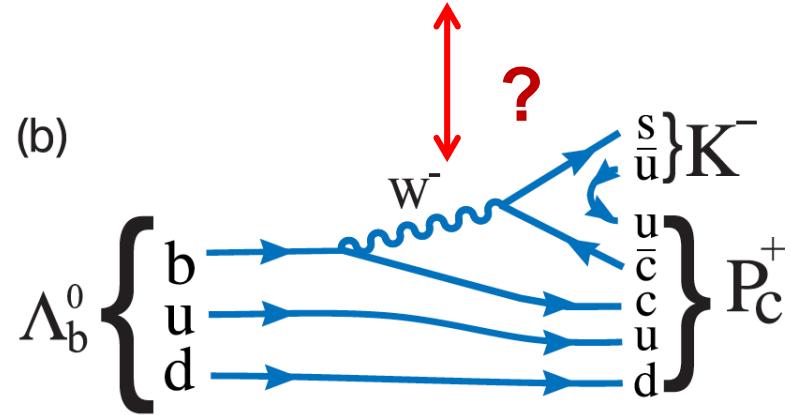
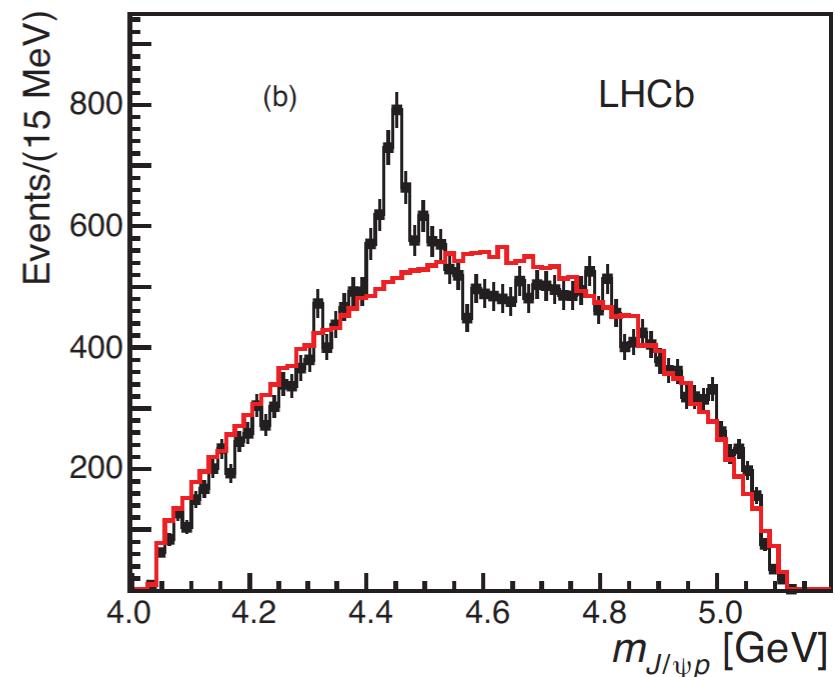
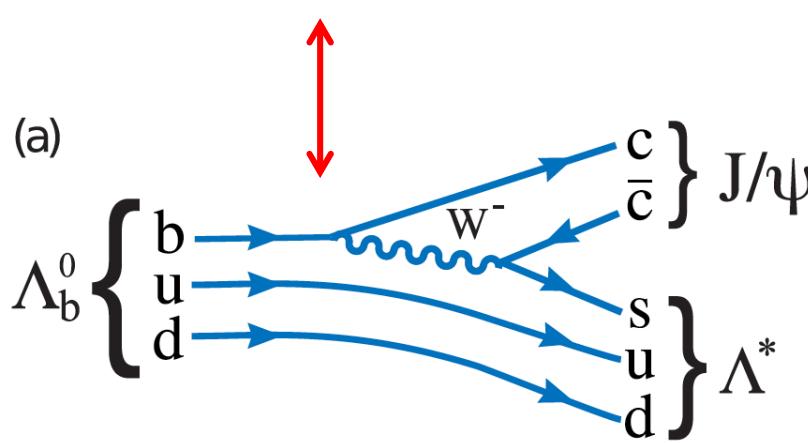
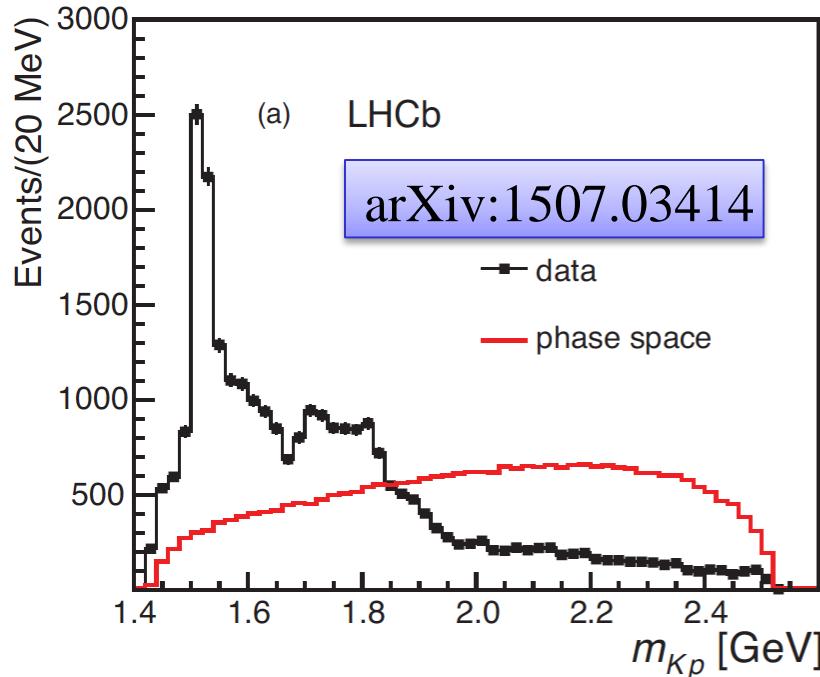
$\Lambda_b \rightarrow J/\psi K^- p$

- First looked for at LHCb as a potential background for $B^0 \rightarrow J/\psi K^+ K^-$
- Large signal found, used for Λ_b lifetime

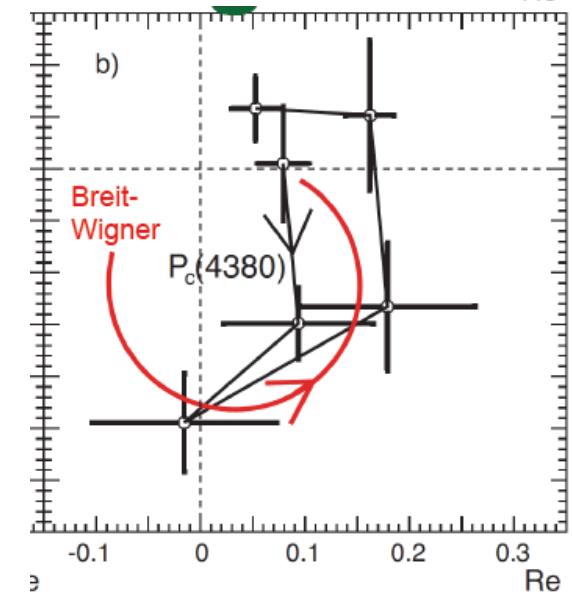
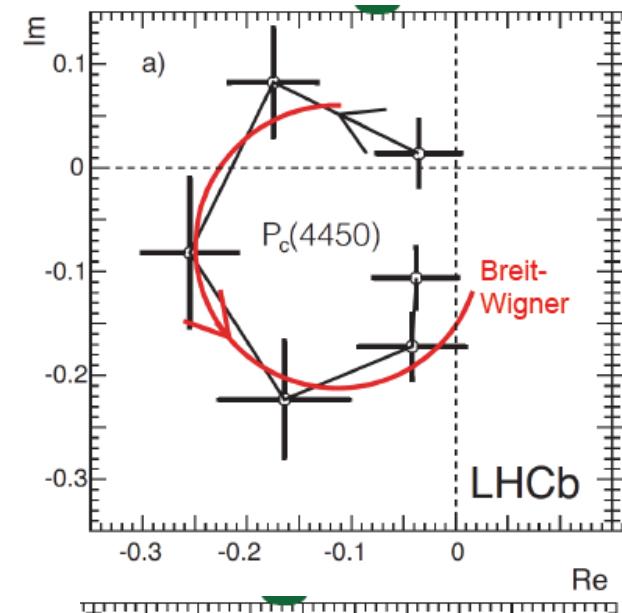
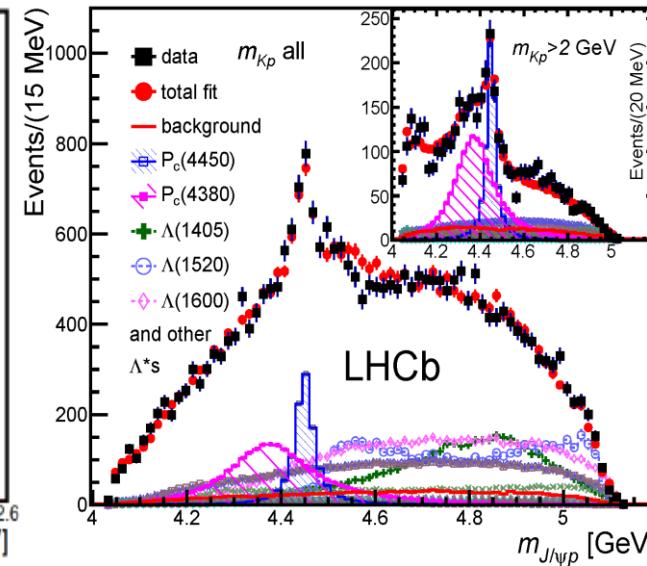
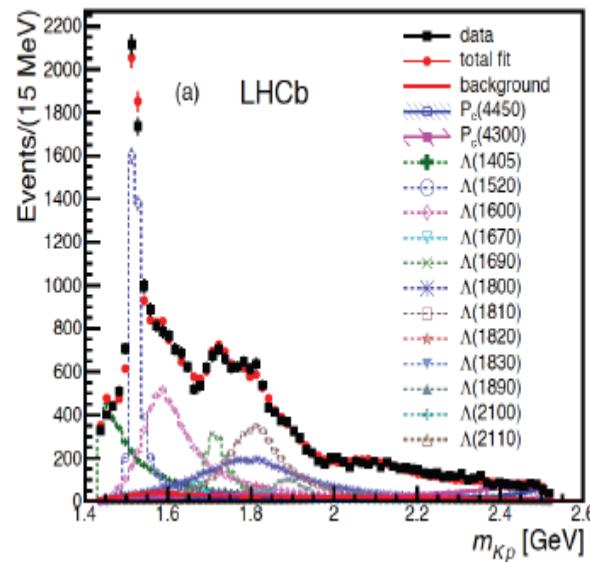
Dalitz plot showed an unusual feature



Projections



Mass (MeV)	Width (MeV)	Fit fraction (%)
$4380 \pm 8 \pm 29$	$205 \pm 18 \pm 86$	$8.4 \pm 0.7 \pm 4.2$
$4449.8 \pm 1.7 \pm 2.5$	$39 \pm 5 \pm 19$	$4.1 \pm 0.5 \pm 1.1$
$\Lambda(1405)$		$15 \pm 1 \pm 6$
$\Lambda(1520)$		$19 \pm 1 \pm 4$



Models

Tightly bound quarks

Two colored diquarks plus the anti-quark, L.Maiani, et. al, [arXiv:1507.04980], ibid [PRD20(1979) 748]

Colored diquark + colored triquark, R. Lebed [arXiv:1507.05867]

Hadroquarkonium

$\chi_{c1} p$ resonance ,U. Meißner,J. Oller[arXiv:1507.07478]

Rescattering

The narrow peak in pJ/ψ spectra at 4450 MeV is antiquark diquark-diquark state, while the broad bump is the result of rescatterings in the (pJ/ψ)-channel, V.V.Anisovich et.al [arXiv:1507.07652]

Molecule

Baryon-meson molecule generally with meson exchange for binding ,

Törnqvist [Z.Phys. C61 (1994) 525]

π exchange dominance, Σ_c (2445) D^* S-wave molecule,

M.Karliner/J.Rosner [arXiv:1506.06386]

What can we do more?

- Search for other possible P_c decay channels ($J/\Psi p\pi, \chi_{c1}p, \Sigma_c^{(*)} D^{(*)}, \Lambda_c D^{(*)} \dots$)
- Search in the inclusive production at LHC
- Search in the inclusive production at e^+e^- factories
- Similar structures in the $\Xi_b \rightarrow J/\Psi \Lambda K$?
- Photoproduction in γ -nucleon interactions

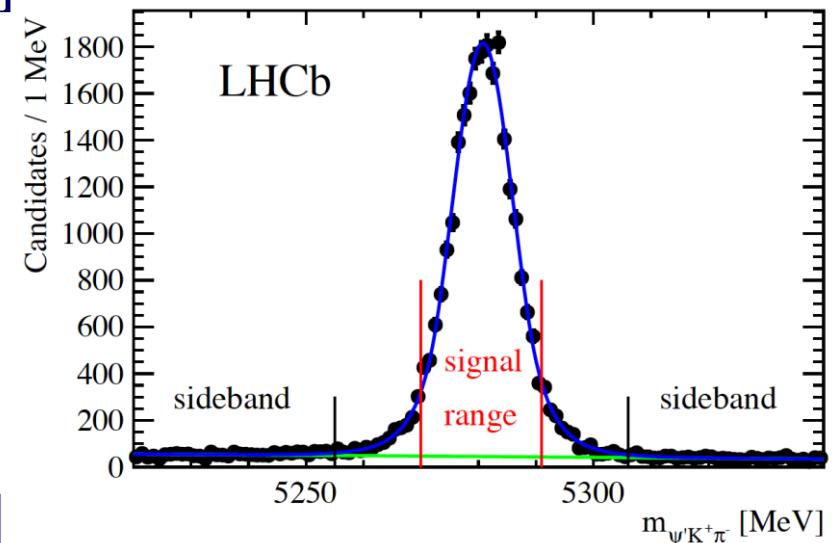
Qian Wang, Xiao-Hai Liu, Qiang Zhao [arXiv:1508.00339], V.Kubarovsky, M.Voloshin [1508.00888], M.Karliner/J.Rosner [arXiv:1508.01496]

Peak cross section for $\gamma + p \rightarrow P_c \rightarrow J/\psi + p$, proportional to $[Br(P_c \rightarrow J/\psi + p)]^2$, can reach tens of nanobarns or more, if $Br(P_c \rightarrow J/\psi + p) \sim 10\%$ ($E\gamma \sim 10\text{GeV}$)

Tetraquark?

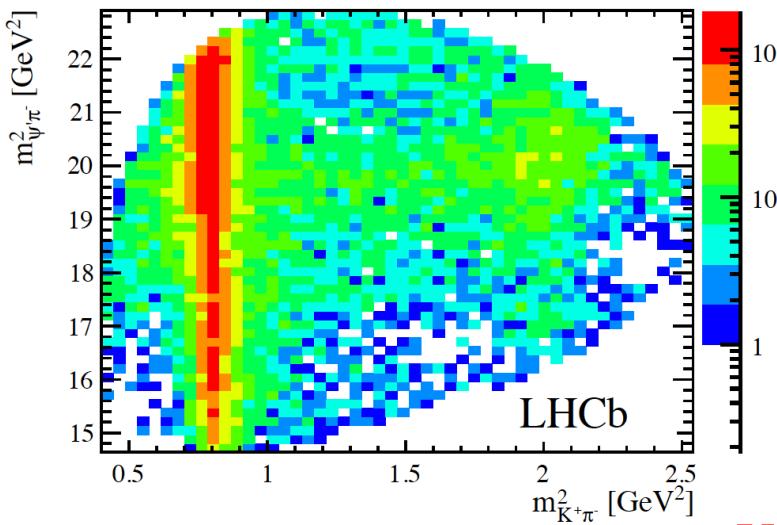
Z(4430)⁺

- $B^0 \rightarrow \psi' \pi^- K^+$, peak in $m(\psi' \pi^-)$, charged charmonium state must be exotic, not $q\bar{q}$
 - First observed by Belle $M=4433 \pm 5$ MeV, $\Gamma=45$ MeV
 - Challenged by BaBar: explanation in terms of K^* 's
 - Belle reanalysis using full amplitude fit:
 $M=4485 \pm 22^{+28}_{-11}$ MeV, $\Gamma=200$ MeV, 1^+ preferred
[Belle PRD 88, 074026 (2013)]
- LHCb analysis also uses full amplitude fit
 - $M=4475 \pm 7^{+15}_{-25}$ MeV
 - $\Gamma=172$ MeV[LHCb PRL 112, 222002 (2014)]

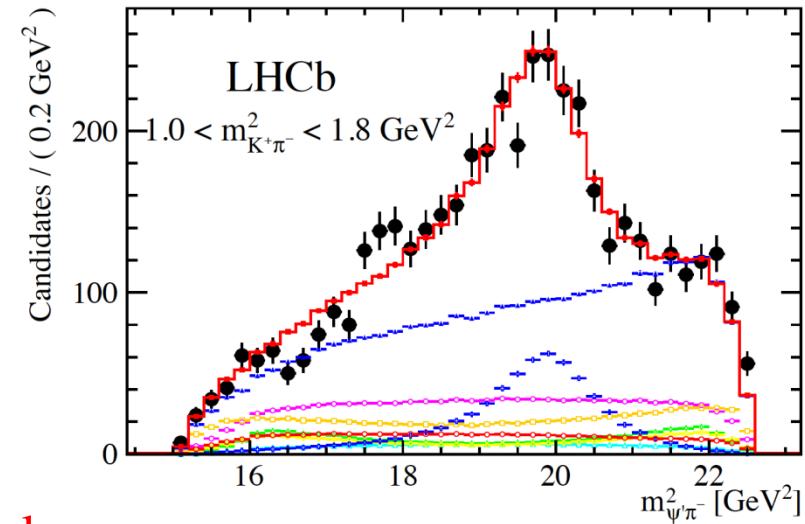


LHCb Amplitude analysis

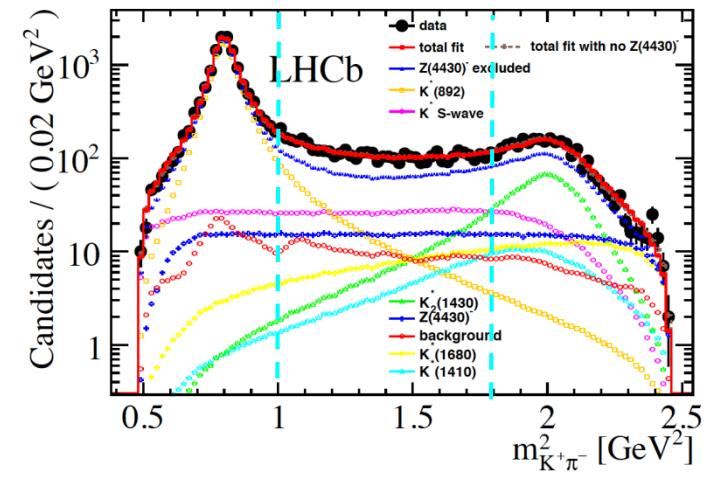
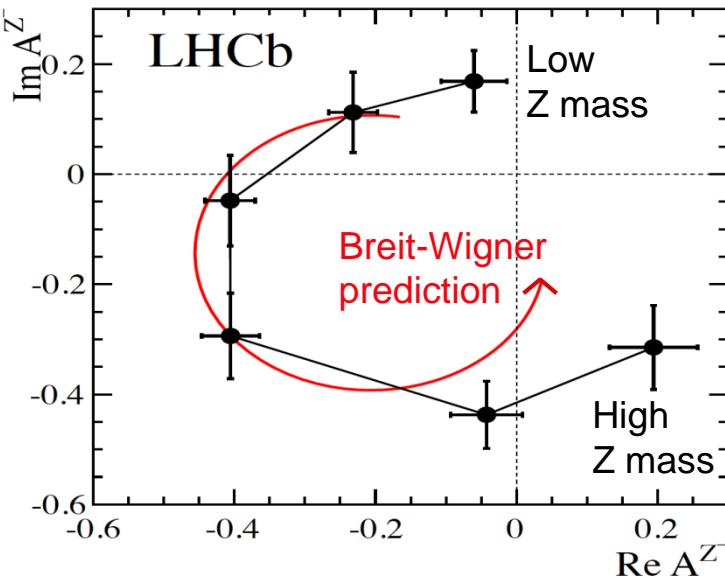
- Full 4D fit to both $K^* \rightarrow K^- \pi^+$ & $Z \rightarrow \psi' \pi^-$ states



$\text{JP} = 1^+$



Unambiguously



Hadron Spectroscopy enters the new region – Physics of Highly Excited Quarkonium or/and Chemistry of Heavy Flavor

We can expect much more from Super B factory, BESIII and Run2 at LHC

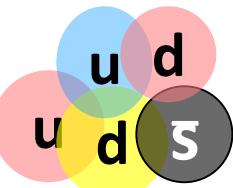
We need new Theoretical instruments for analysis

What about other color-singlet combinations?

Other possible “white” combinations of quarks & gluons:

Pentaquark:

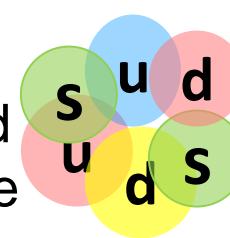
S=+1 Baryon



H-diBaryon

Glueball

tightly bound
6-quark state



Color-singlet multi-gluon bound state

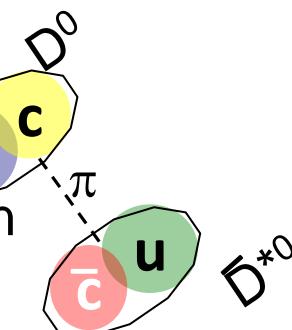


Tetraquark mesons

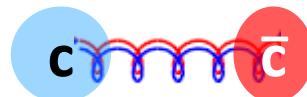
tightly bound
diquark-dantiquark



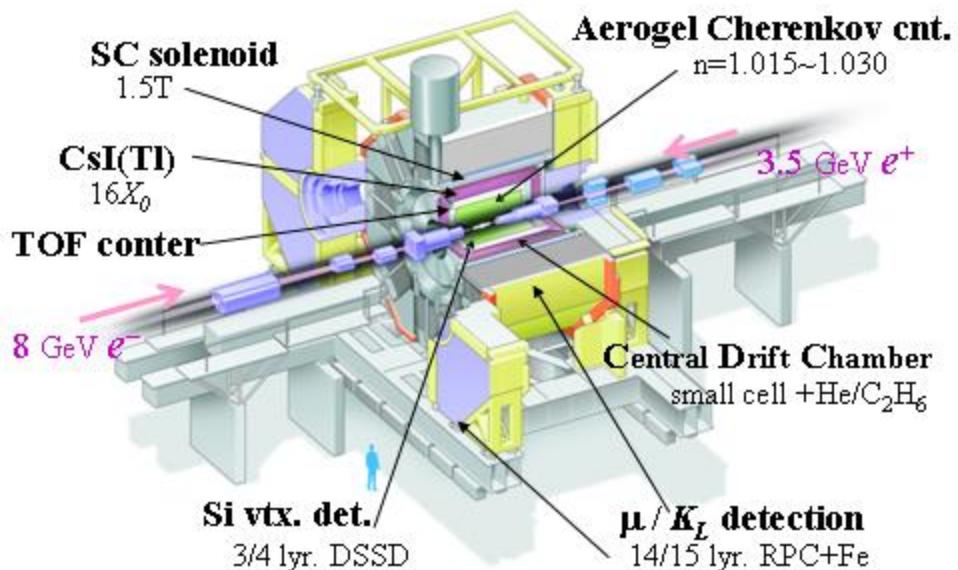
loosely bound
meson-antimeson
“molecule”



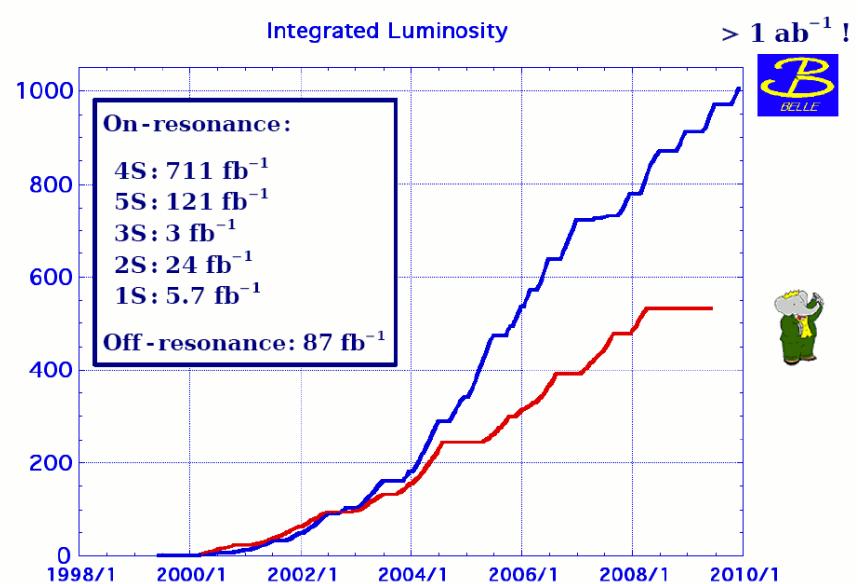
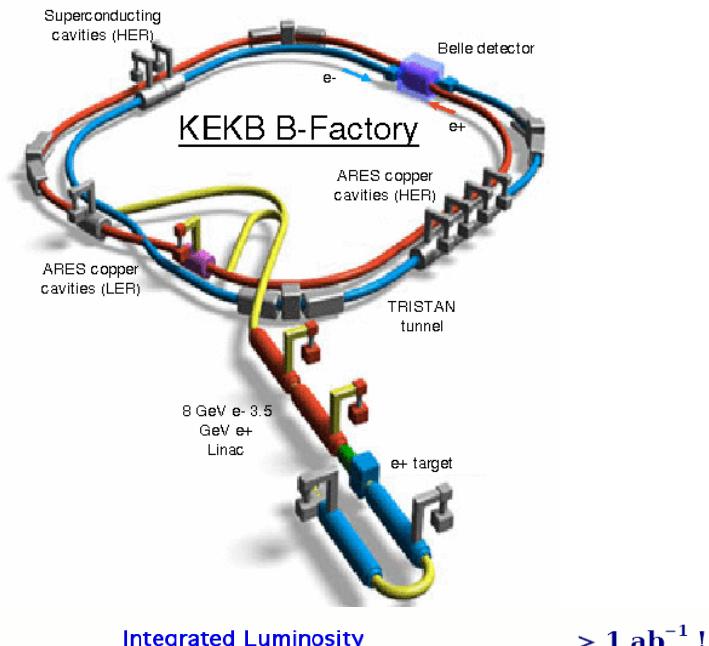
q-q-bar-gluon hybrid mesons



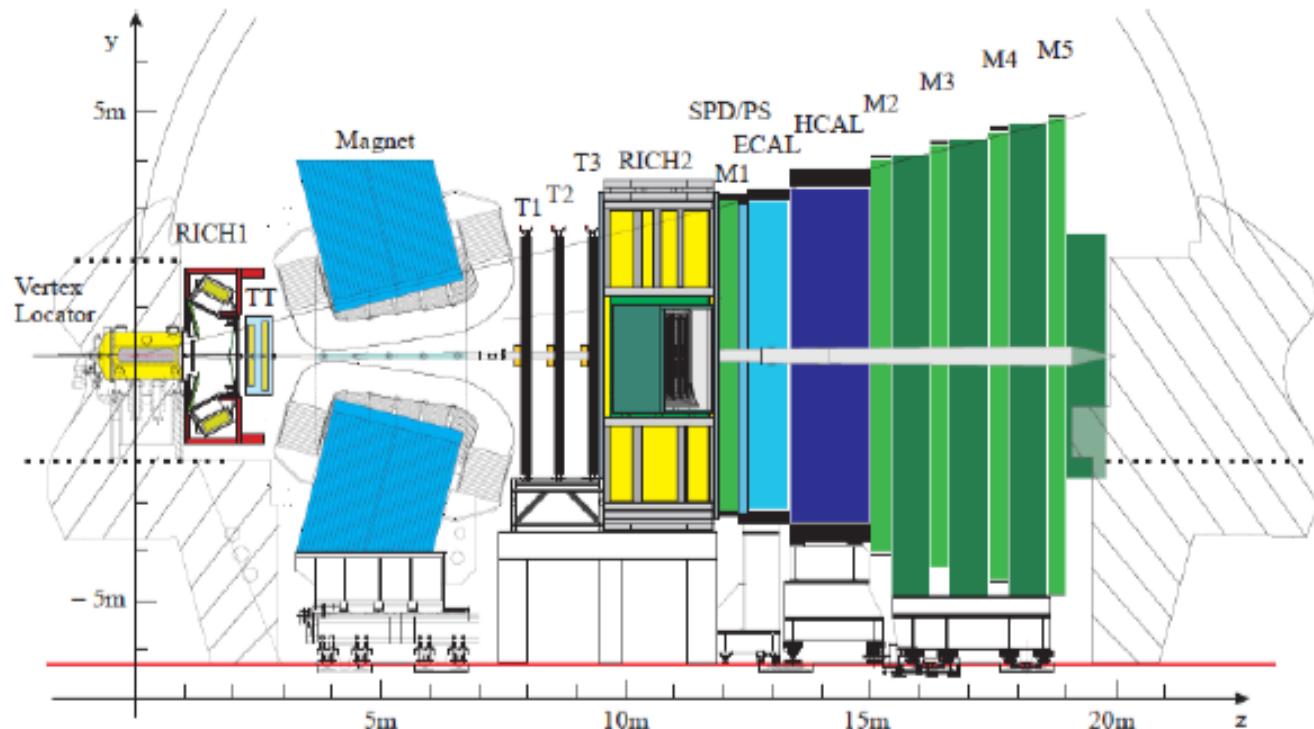
Belle Detector



- $3.5 \text{ GeV } e^+ \times 8.0 \text{ GeV } e^-$.
- $\mathcal{L}_{\max} = 2.1 \times 10^{34} \text{ cm}^{-2} s^{-1}$
- Continuous injection
→ $1.1 \text{ fb}^{-1}/\text{day}$.
- $\int \mathcal{L} dt \approx 1 \text{ ab}^{-1}$



LHCb detector



Impact parameter:

$$\sigma_{IP} = 20 \mu\text{m}$$

Proper time:

$$\sigma_\tau = 45 \text{ fs for } B_s^0 \rightarrow J/\psi \phi \text{ or } D_s^+ \pi^-$$

Momentum:

$$\Delta p/p = 0.4 \sim 0.6\% \text{ (5 - 100 GeV/c)}$$

Mass :

$$\sigma_m = 8 \text{ MeV}/c^2 \text{ for } B \rightarrow J/\psi X \text{ (constrained } m_{J/\psi})$$

RICH $K - \pi$ separation:

$$\epsilon(K \rightarrow K) \sim 95\% \text{ mis-ID } \epsilon(\pi \rightarrow K) \sim 5\%$$

Muon ID:

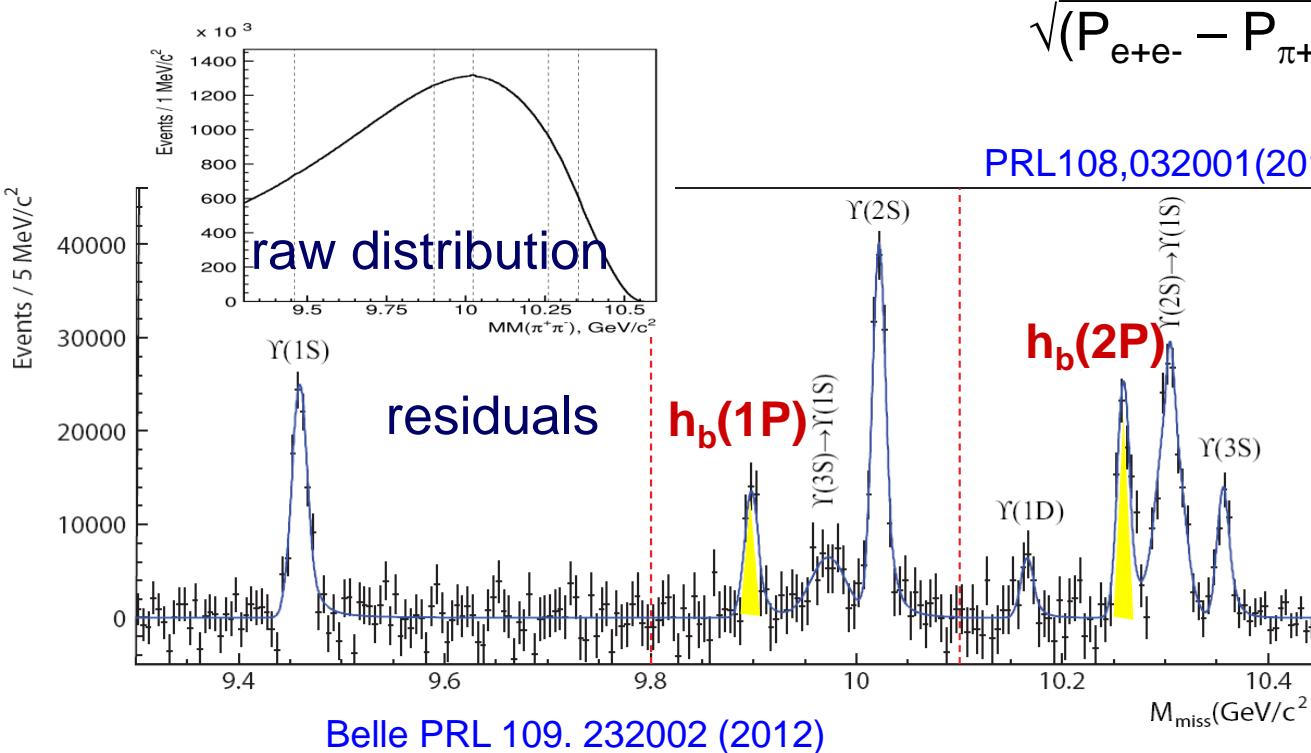
$$\epsilon(\mu \rightarrow \mu) \sim 97\% \text{ mis-ID } \epsilon(\pi \rightarrow \mu) \sim 1 - 3\%$$

ECAL:

$$\Delta E/E = 1 \oplus 10\%/\sqrt{E(\text{GeV})}$$

Observation of $h_b(1P,2P)$

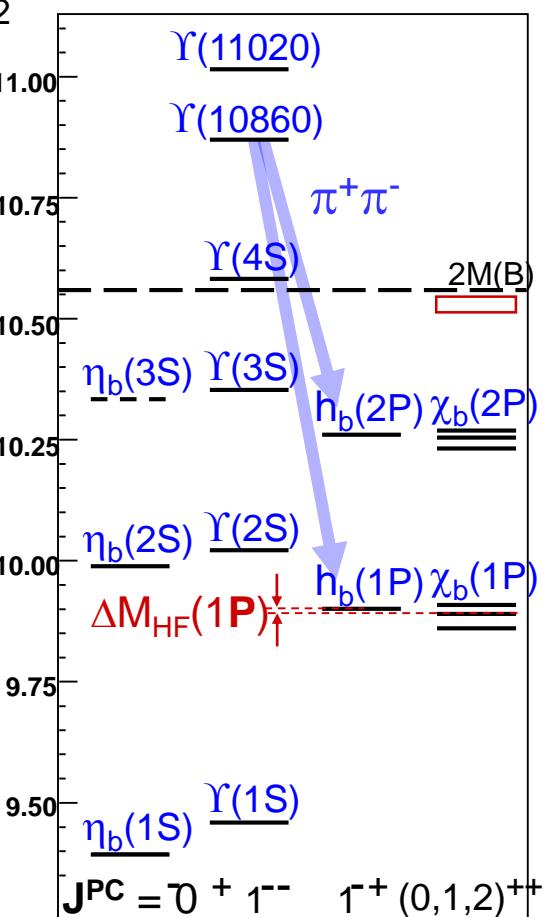
$e^+e^- \rightarrow \gamma(5S) \rightarrow h_b(nP) \pi^+\pi^-$ \leftarrow reconstructed, use $M_{\text{miss}}(\pi^+\pi^-)$



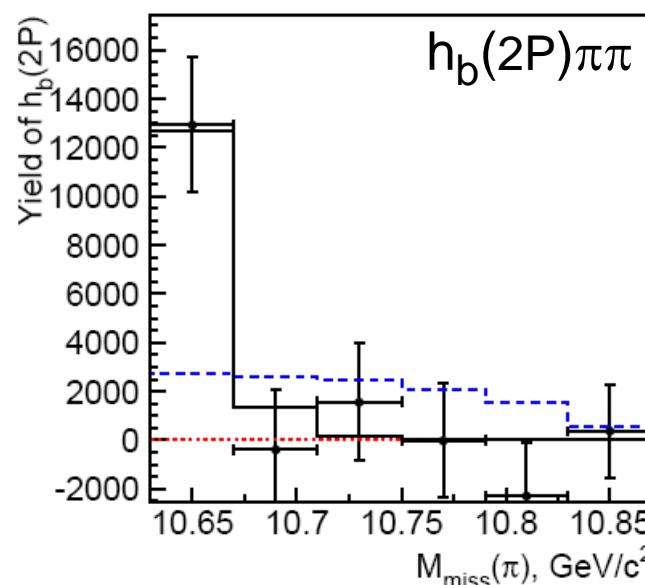
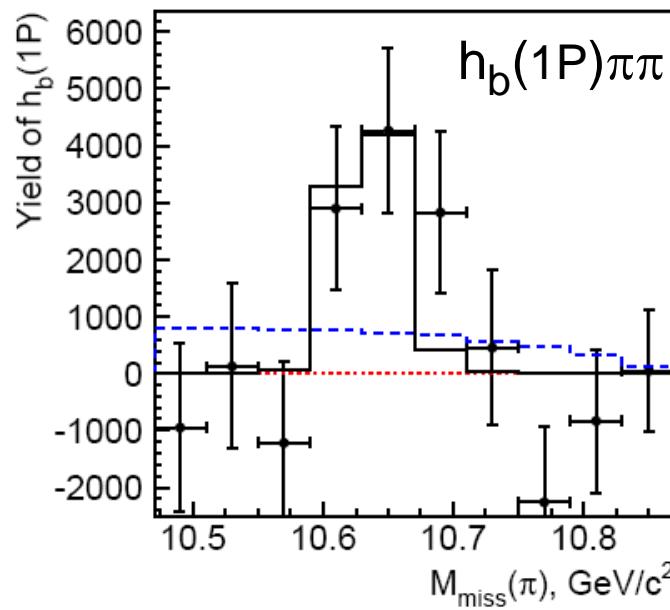
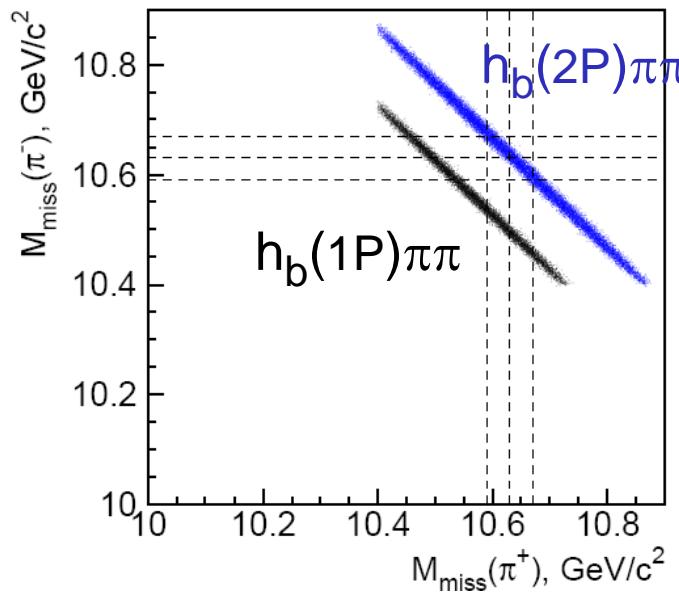
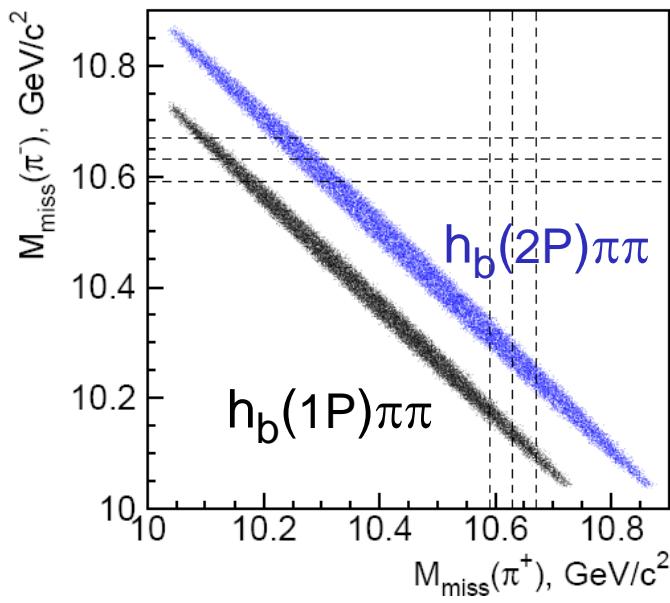
$$\Delta M_{HF}(1P) = +0.8 \pm 1.1 \text{ MeV}$$

$$\Delta M_{HF}(2P) = +0.5 \pm 1.2 \text{ MeV}$$

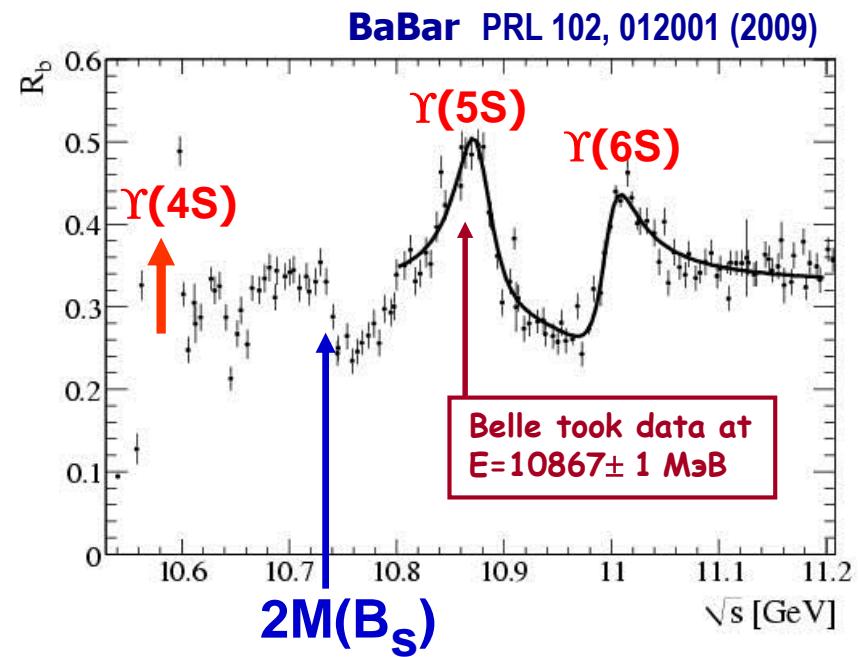
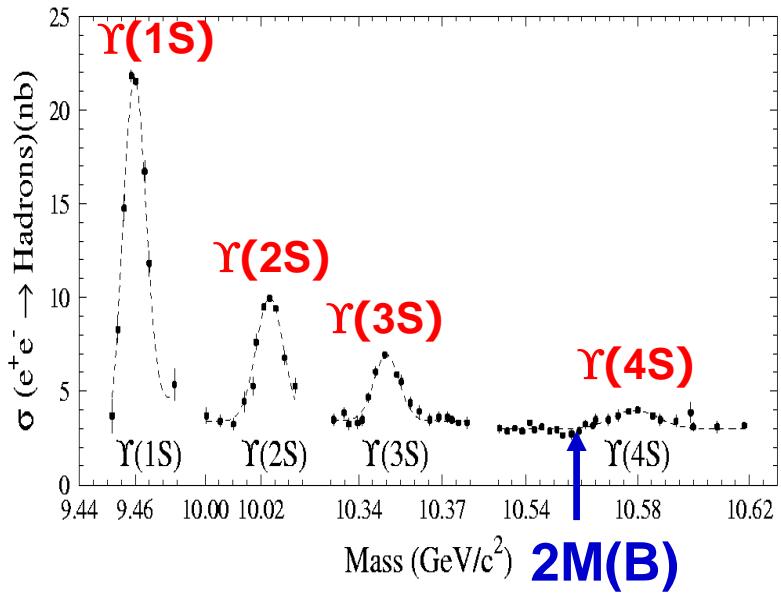
consistent with zero,
as expected



Large $h_b(1,2P)$ production rates
c.f. CLEO/BESIII $e^+e^- \rightarrow h_c \pi^+\pi^-$

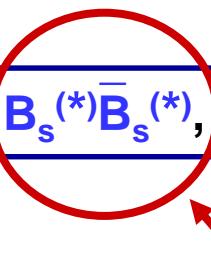


e^+e^- hadronic cross-section



$e^+ e^- \rightarrow \Upsilon(4S) \rightarrow B\bar{B}$, where B is B^+ or B^0

$e^+ e^- \rightarrow b\bar{b} (\Upsilon(5S)) \rightarrow B^{(*)}\bar{B}^{(*)}, B^{(*)}\bar{B}^{(*)}\pi, B\bar{B}\pi\pi, B_s^{(*)}\bar{B}_s^{(*)}, \Upsilon(1S)\pi\pi, \Upsilon X \dots$



main motivation
for taking data at $\Upsilon(5S)$