

# Lattice QCD at non-zero Temperature and Density

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# Lattice QCD/Lattice Field Theory



adapted from C. Davies, LP2013

## *Particle physics*

QCD parameters  
Hadron spectrum  
Hadron structure  
CKM elements  
Glueballs and exotica

## *Nuclear physics*

Nuclear masses  
and properties

## Theories beyond the Standard Model

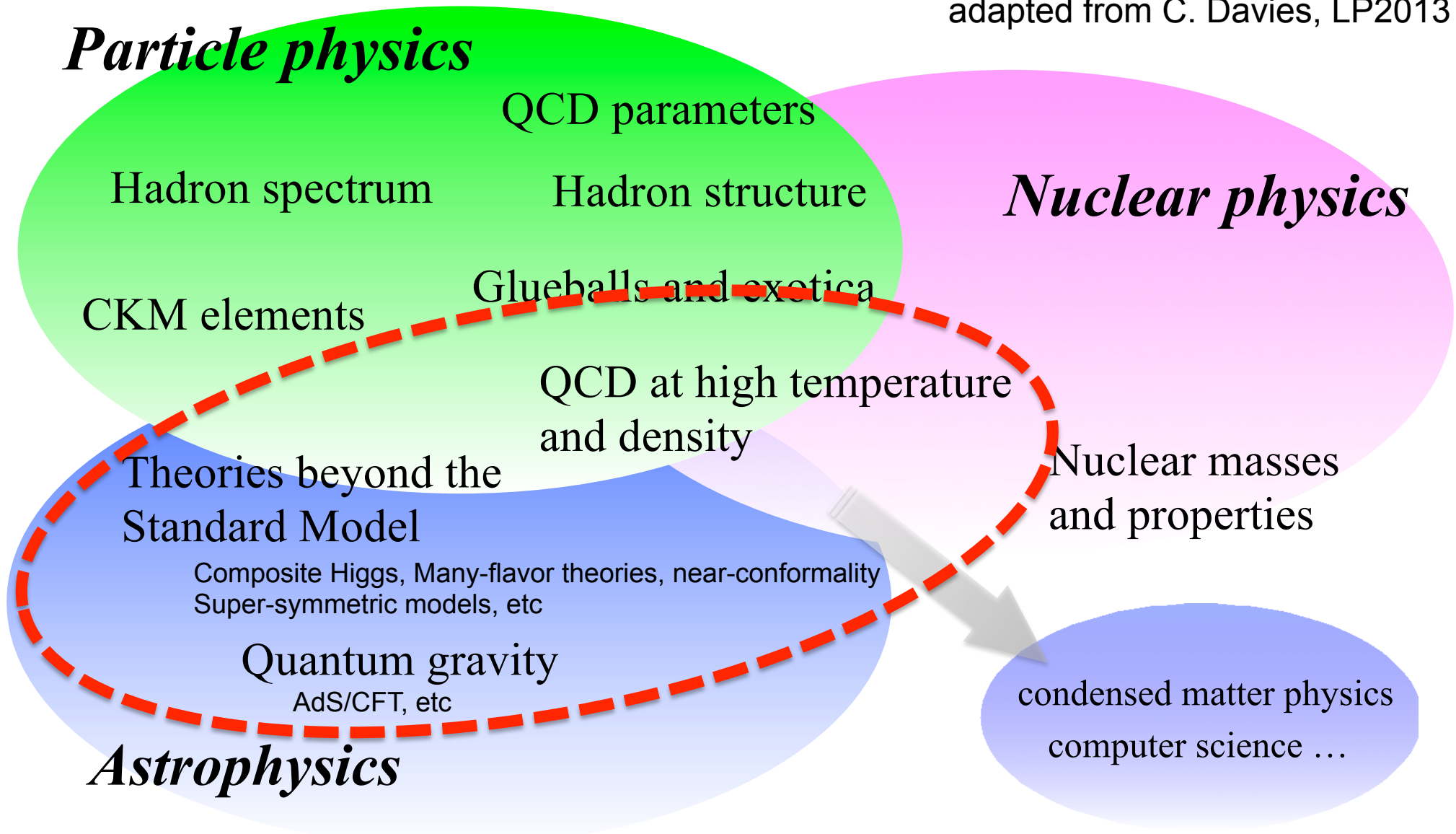
Composite Higgs, Many-flavor theories, near-conformality  
Super-symmetric models, etc

Quantum gravity  
AdS/CFT, etc

## *Astrophysics*

condensed matter physics  
computer science ...

adapted from C. Davies, LP2013





# Plan of talk



- ❖ Non-zero temperature but Zero baryon density
  - ❖ Phase diagram and equation of state (EoS)
  - ❖ Loose ends?
- ❖ Non-zero but small baryon density
  - ❖ Lattice QCD and comparison with phenomenology
- ❖ Conserved charge fluctuations and experiment
- ❖ Search for the critical end point?
- ❖ Summary

# Development since 80's

Titan

- Deeper theoretical understandings
- Better numerical algorithms
- $O(10^{10})$  increase of FLOPS power

2015

K computer



Sequoia



VAX 11/780



1980

1<sup>st</sup> Monte Carlo of lattice gauge theory by M. Creutz, L. Jacobs, C. Rebbi

1<sup>st</sup> hadron mass calculations by D. Weingarten, H. Hamber & G. Parisi

- includes all light quarks (u, d, s), and even charm

$$N_f = 2(\text{degenerate u \& d}) + 1(\text{s}) \quad \text{or} \quad N_f = 2(\text{u \& d}) + 1(\text{s}) + 1(\text{c})$$

- physical values for the quark masses

$$(m_u + m_d) / 2 \approx 3.5 \text{ MeV}, \quad m_s \approx 95 \text{ MeV}, \quad m_c \approx 1.3 \text{ GeV} \quad (\overline{MS} \text{ at } 2 \text{ GeV})$$

- results are extrapolated to *the continuum limit*

$$a(\text{lattice spacing}) \approx 0.1 \text{ fm} - 0.05 \text{ fm} \Rightarrow 0(\text{continuum})$$

- A word of caution: thermodynamic calculations still use

- mostly “staggered quark action” with some unwanted properties (continuum limit mandatory)
- Full-fledged results with chirally symmetric actions (e.g., domain-wall) needs more FLOPS to catch up.



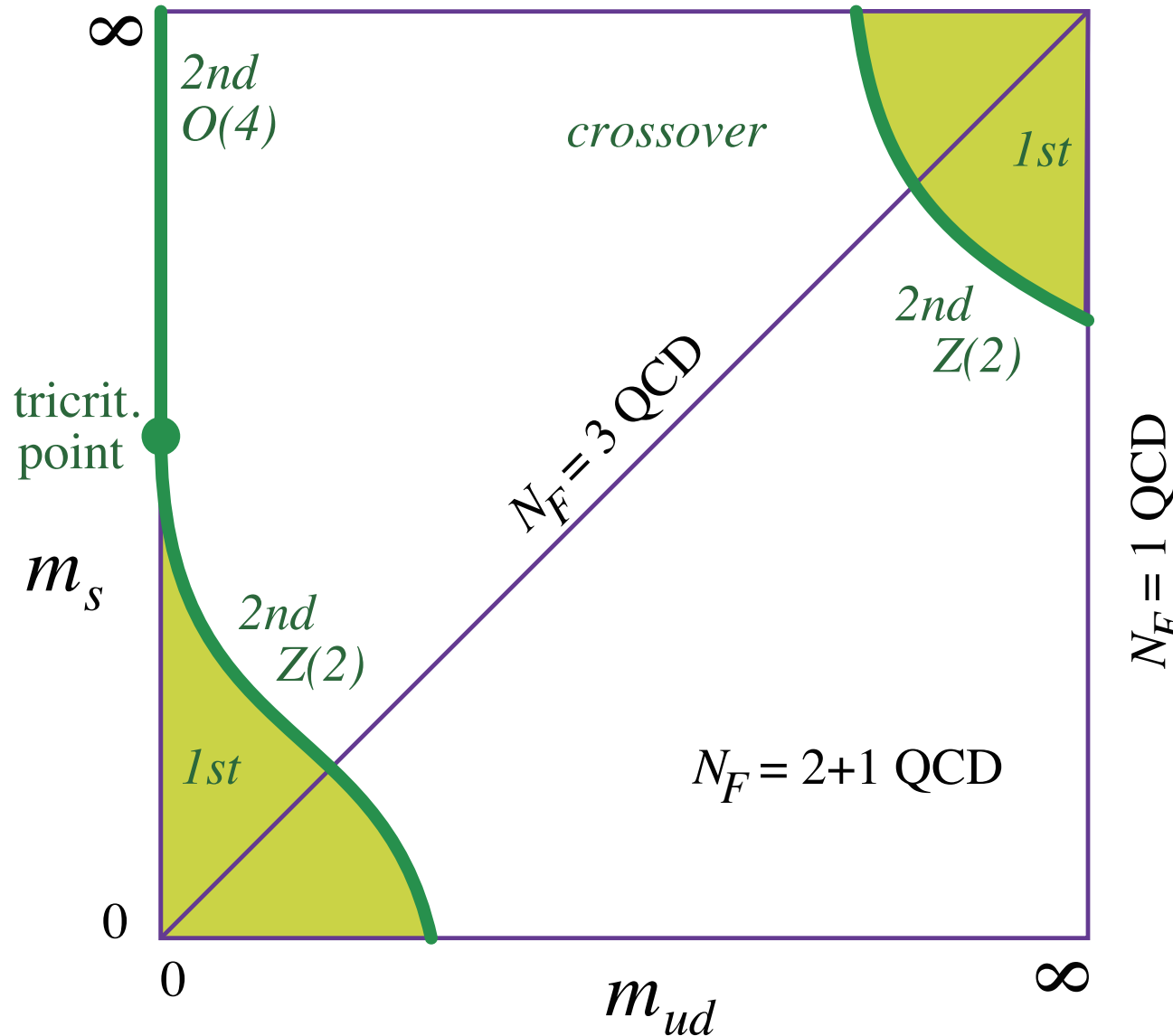
Non-zero temperature  
but Zero baryon density

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# “Standard” picture of the QCD phases

$N_F = 2$  QCD

SU(3) YM : pure gluon theory

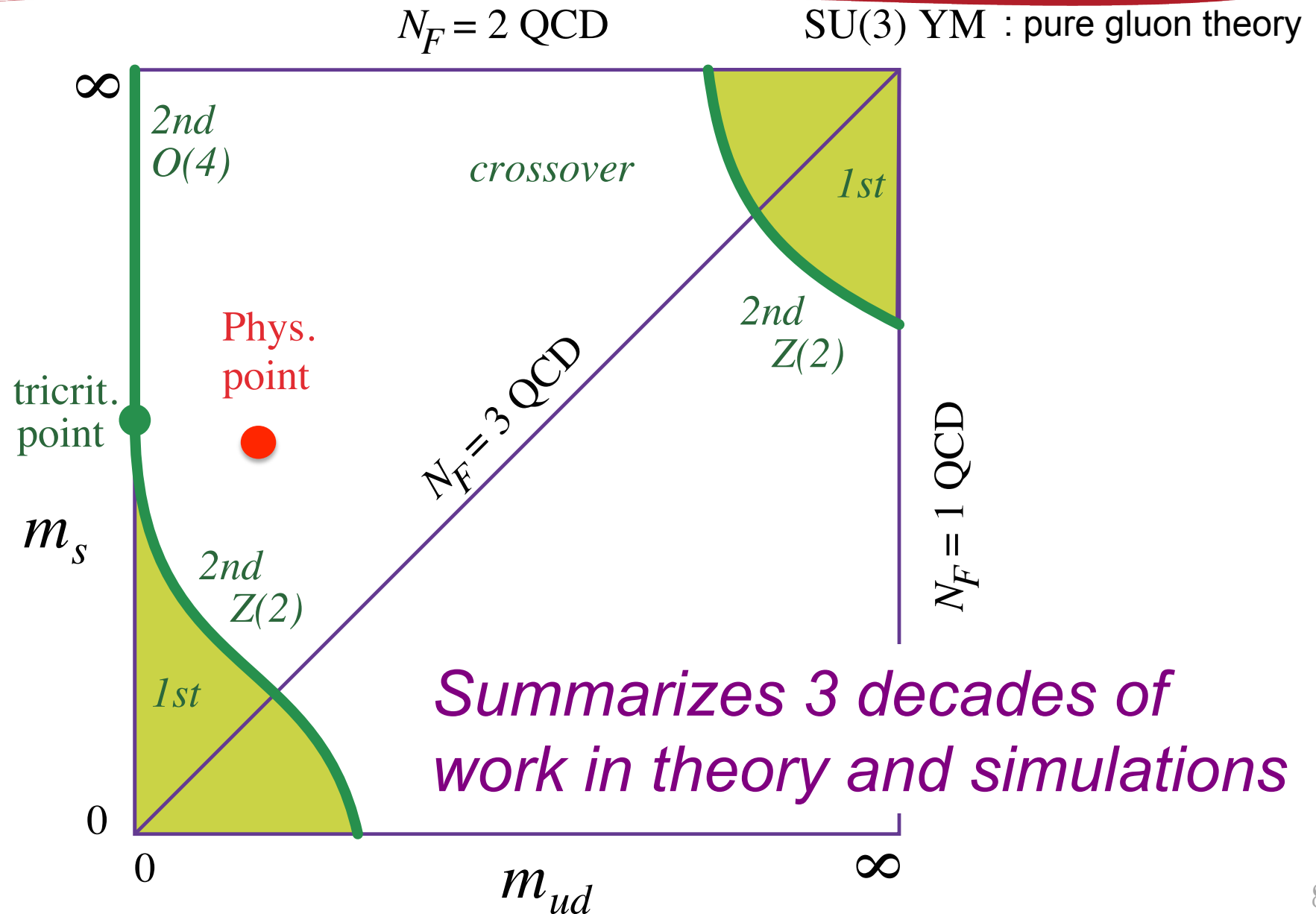


“Columbia Plot”

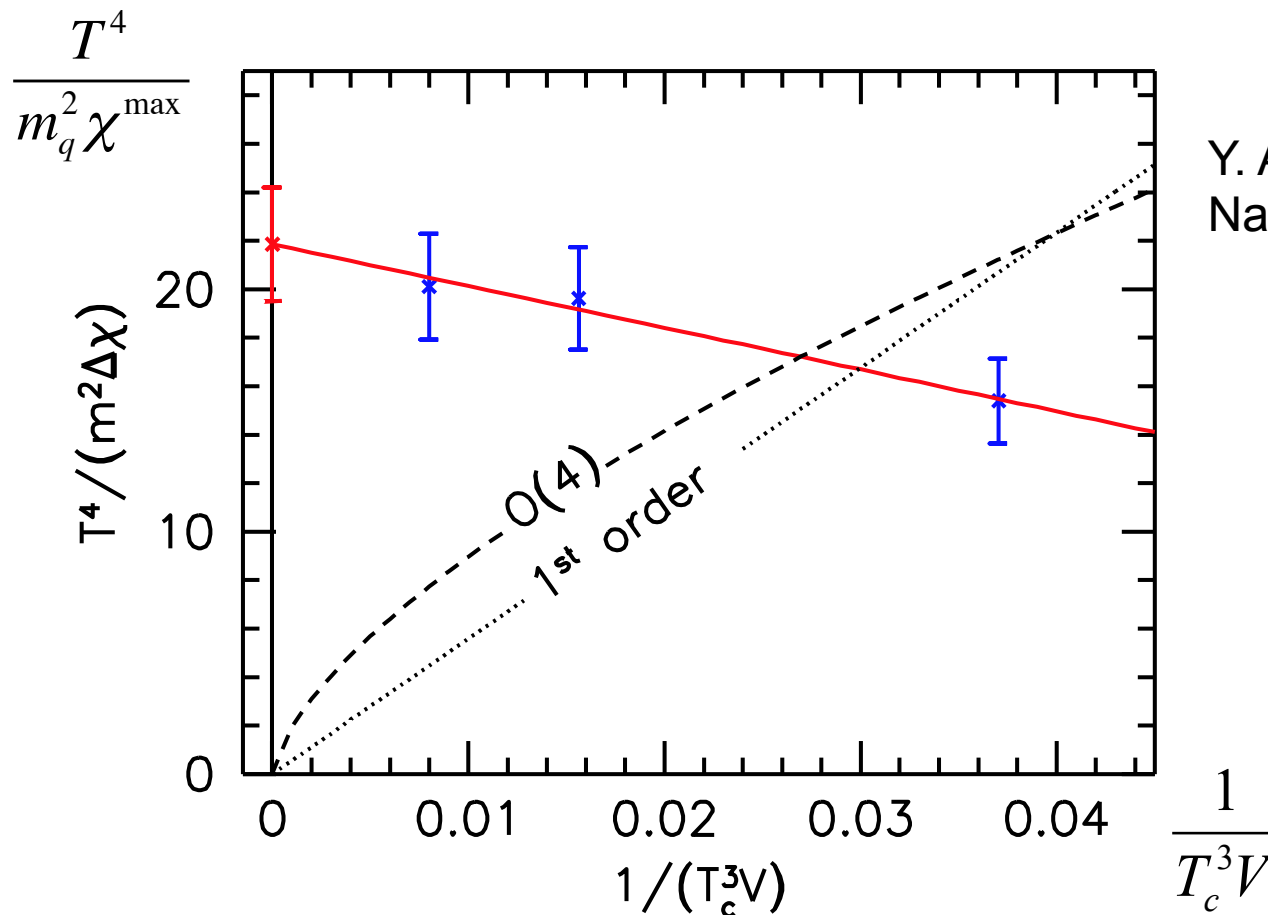
F. R. Brown et al,  
PRL 65, 2491 (1990)



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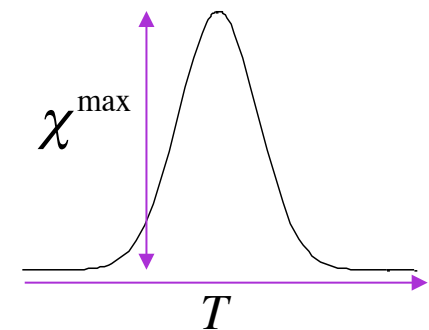


- *Continuous crossover* as a function of temperature  
i.e., no jumps or divergences in the observables

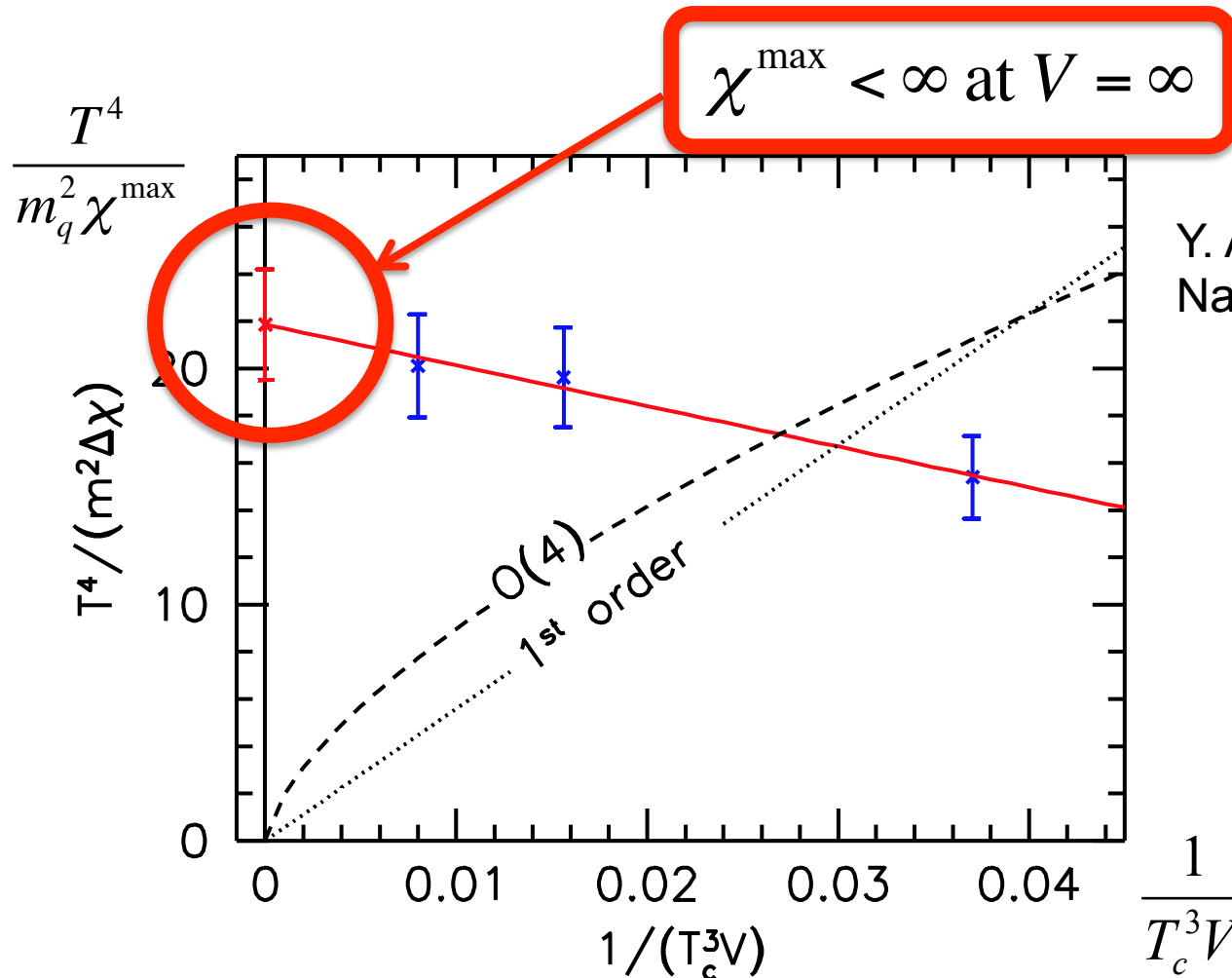


Y. Aoki et al,  
Nature 443:675 (2006)

$$\chi = \int dx (\langle \bar{q}q(x) \bar{q}q(0) \rangle - \langle \bar{q}q(x) \rangle \langle \bar{q}q(0) \rangle)$$

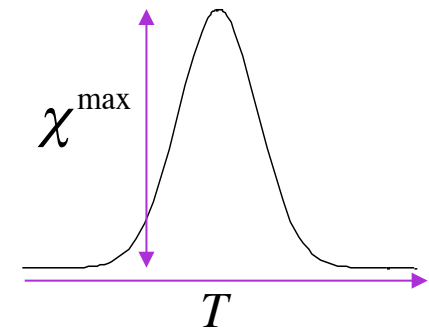


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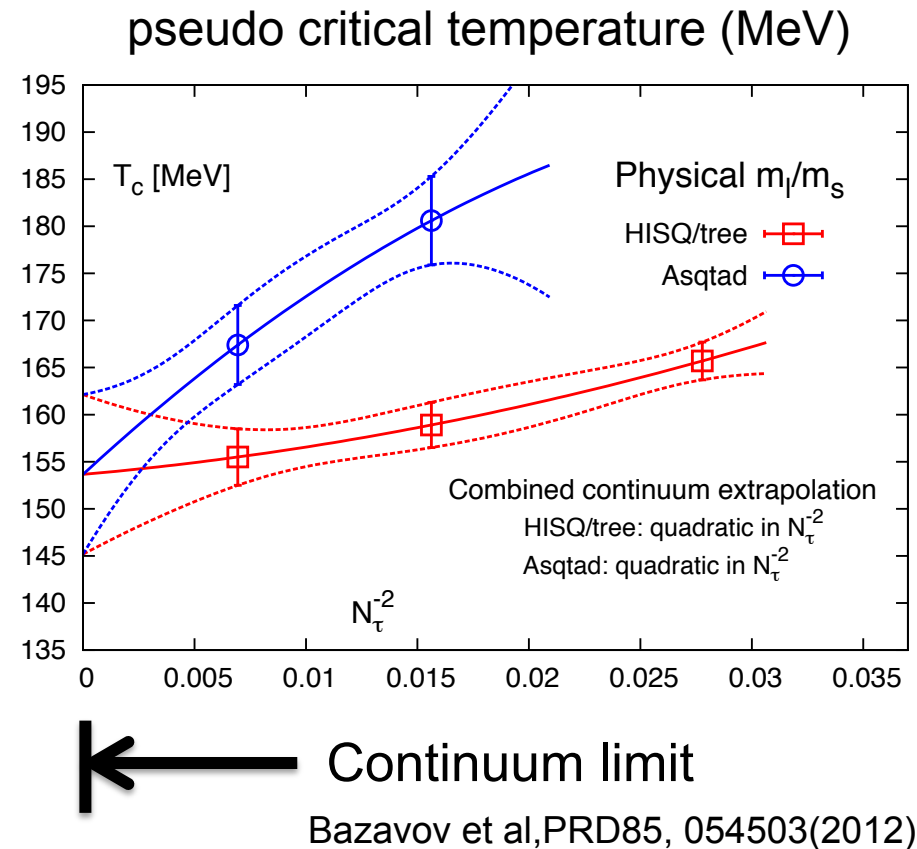
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- Checked with 3 types of staggered actions (stout, Asqtad, HISQ)
- Significant difference at finite lattice spacings, but consistent results in the continuum limit
- Pseudo critical temperature:

$$T_{pc} = \begin{cases} 155 \pm 6 \text{ MeV} & \text{WB} \\ 154 \pm 9 \text{ MeV} & \text{HotQCD} \end{cases}$$

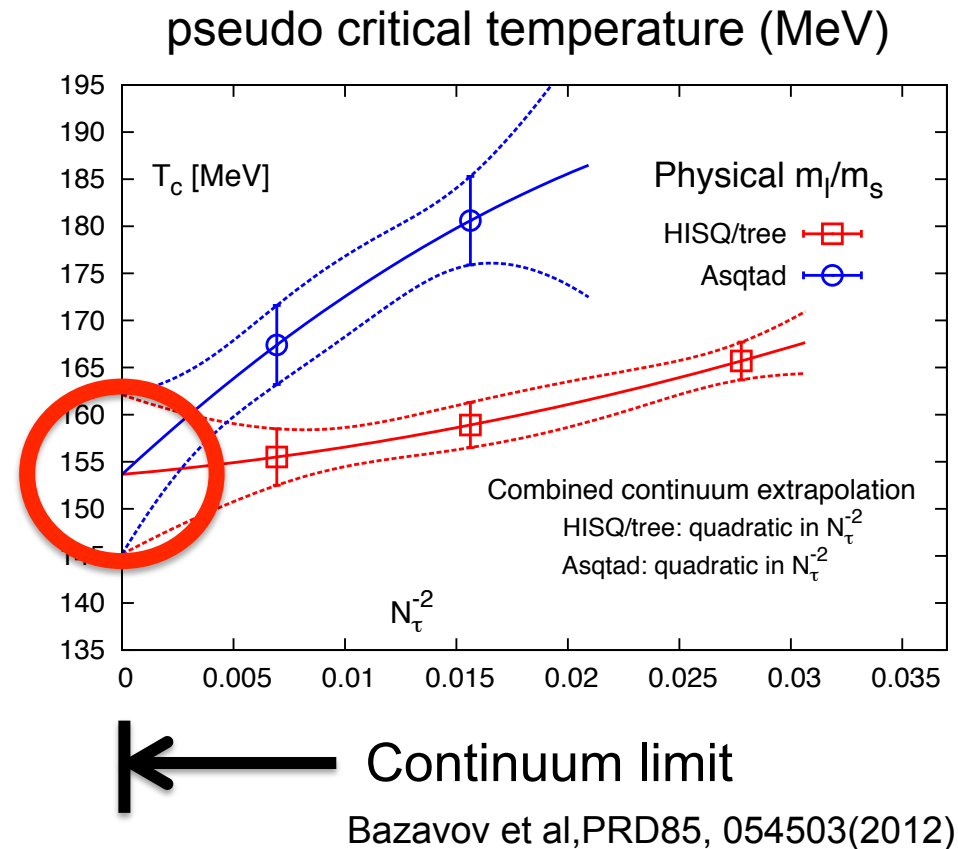
- Wuppertal-Butapest Collaboration
  - Y. Aoki et al, Nature 443: 675 (2006); S. Borsanyi et al, JHEP 2010:73 (2010)
- HotQCD Collaboration
  - A. Bazavov et al, PRD 85:054503 (2012)

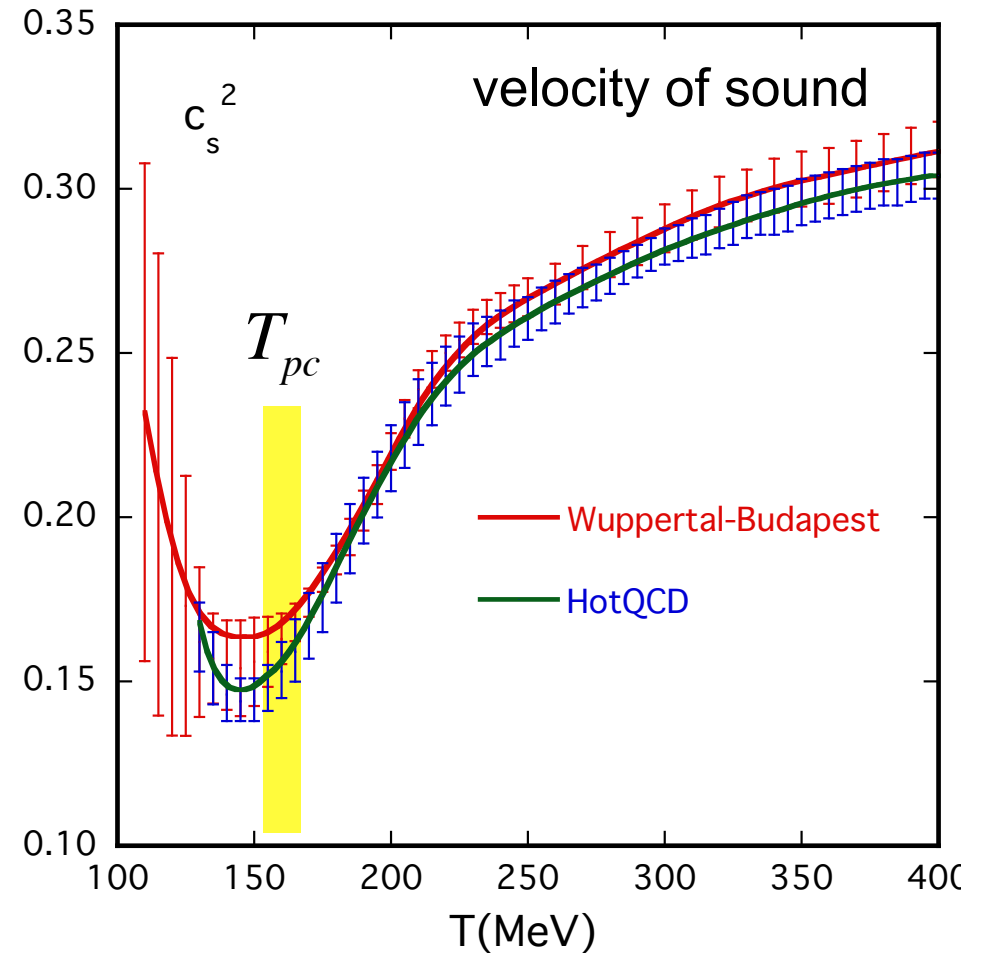
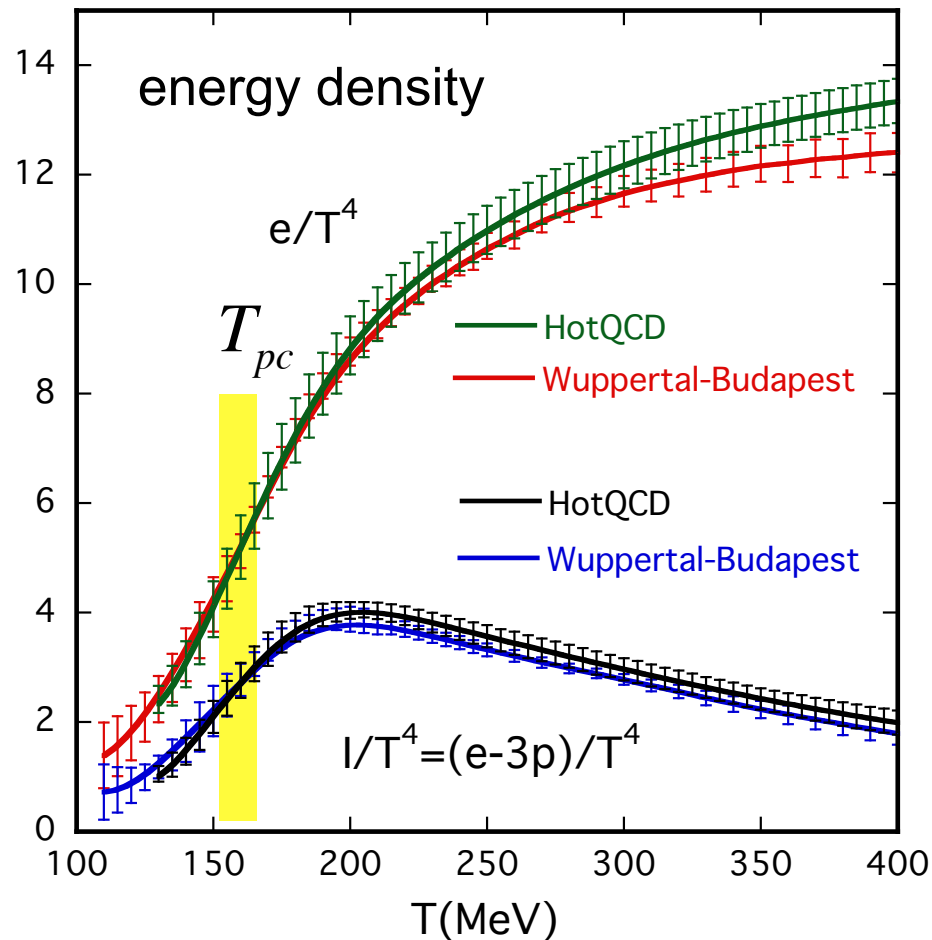


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- Wuppertal-Budapest: S. Borsanyi et al, Phys. Lett. B730, 99 (2014)
- HotQCD: A. Bazavov et al, PRD90, 094503(2014)

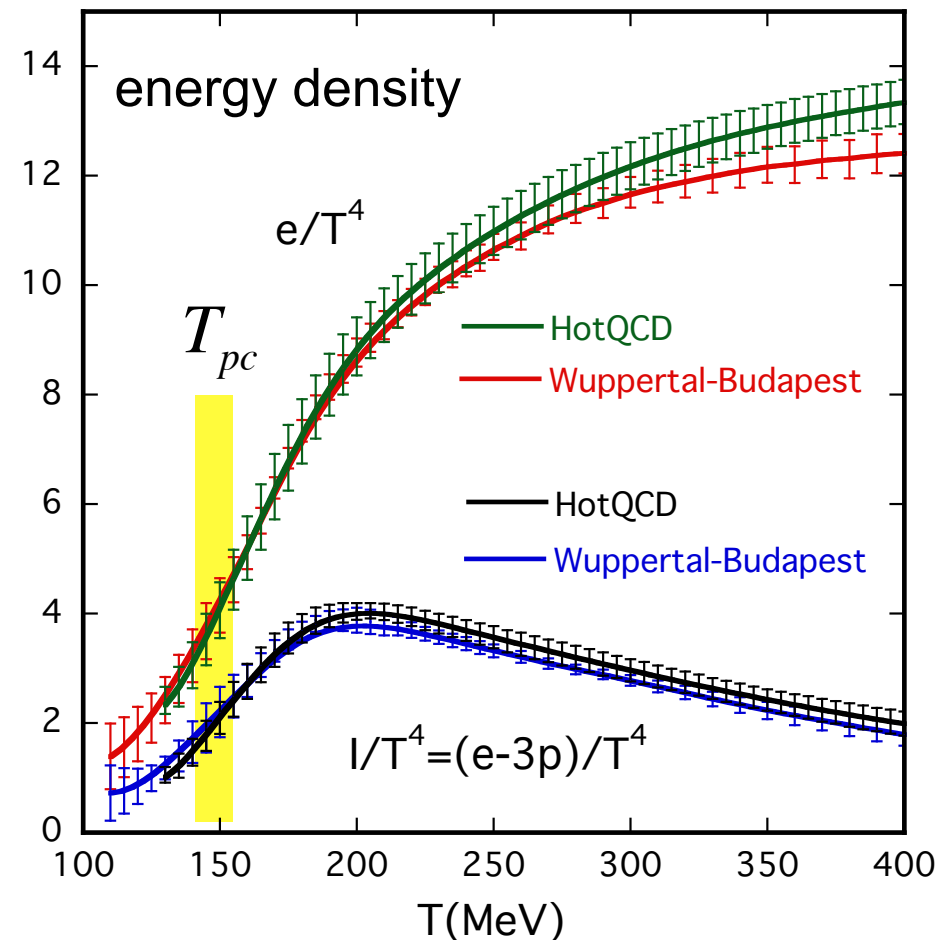
- Direct photon transverse distribution as a probe of temperature in heavy ion collisions

$$n_\gamma \propto \exp\left(-\frac{p_T}{T_\gamma}\right)$$

$$T_\gamma = \begin{cases} 221 \pm 27 \text{ MeV} & \text{RHIC } \sqrt{s} = 200 \text{ GeV} \\ 304 \pm 51 \text{ MeV} & \text{LHC } \sqrt{s} = 2.7 \text{ TeV} \end{cases}$$

PHENIX collaboration, PRL104, 132301(2010)

M. Wilde for ALICE collaboration, arXiv:1210.5958 (2012)



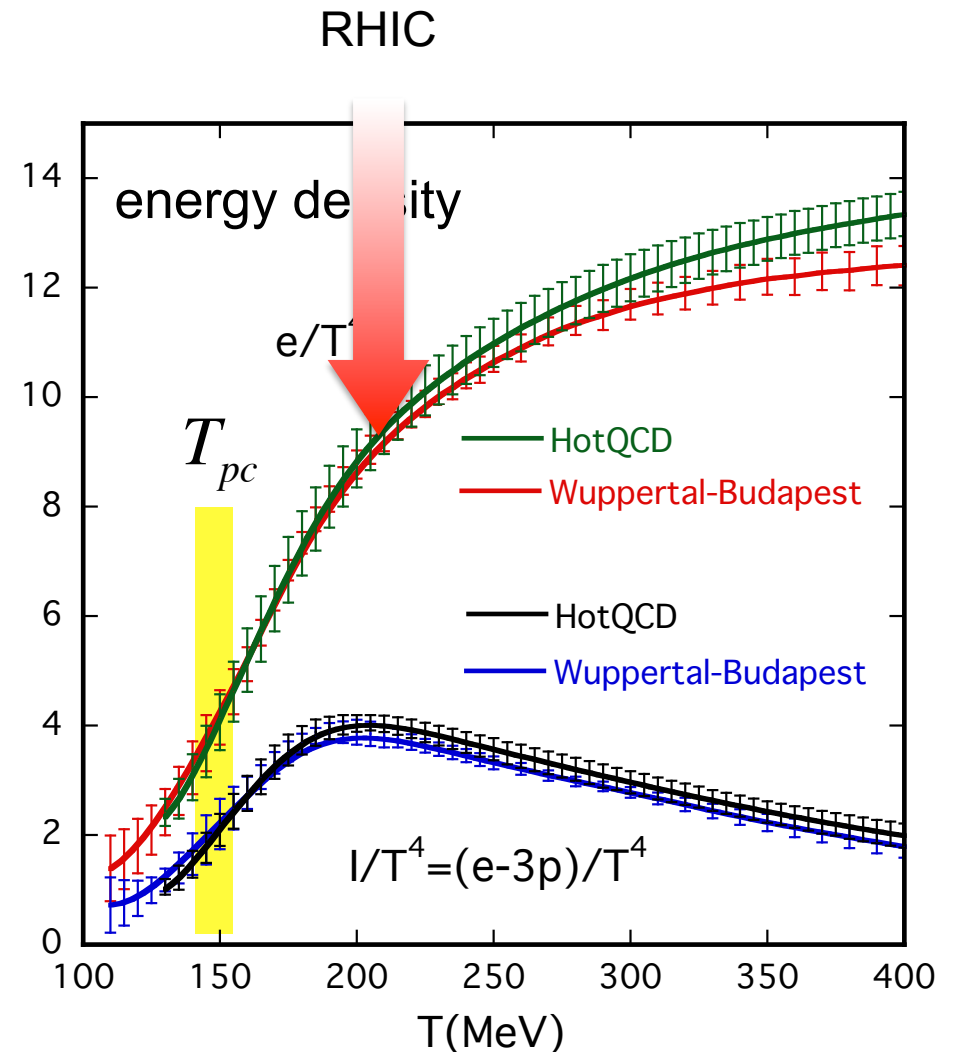
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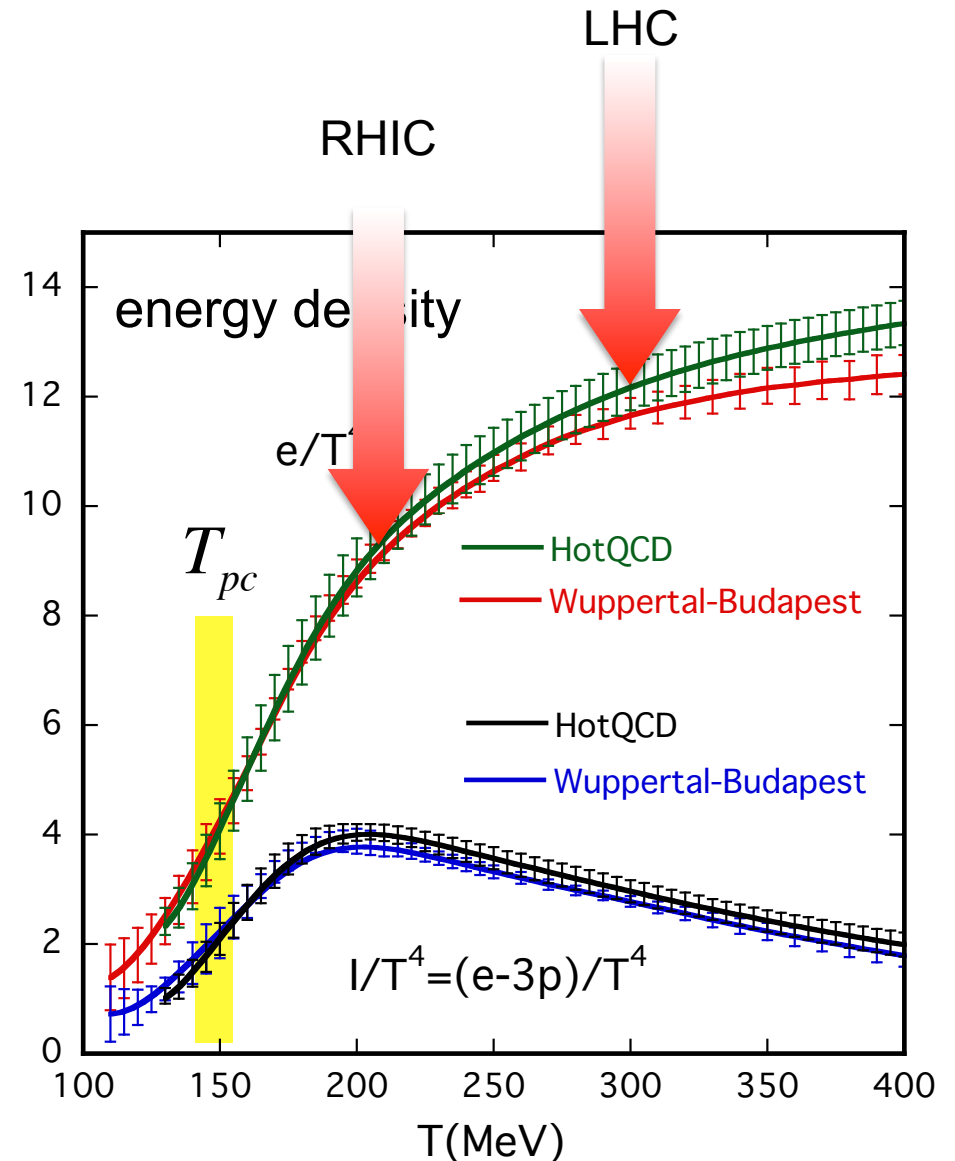
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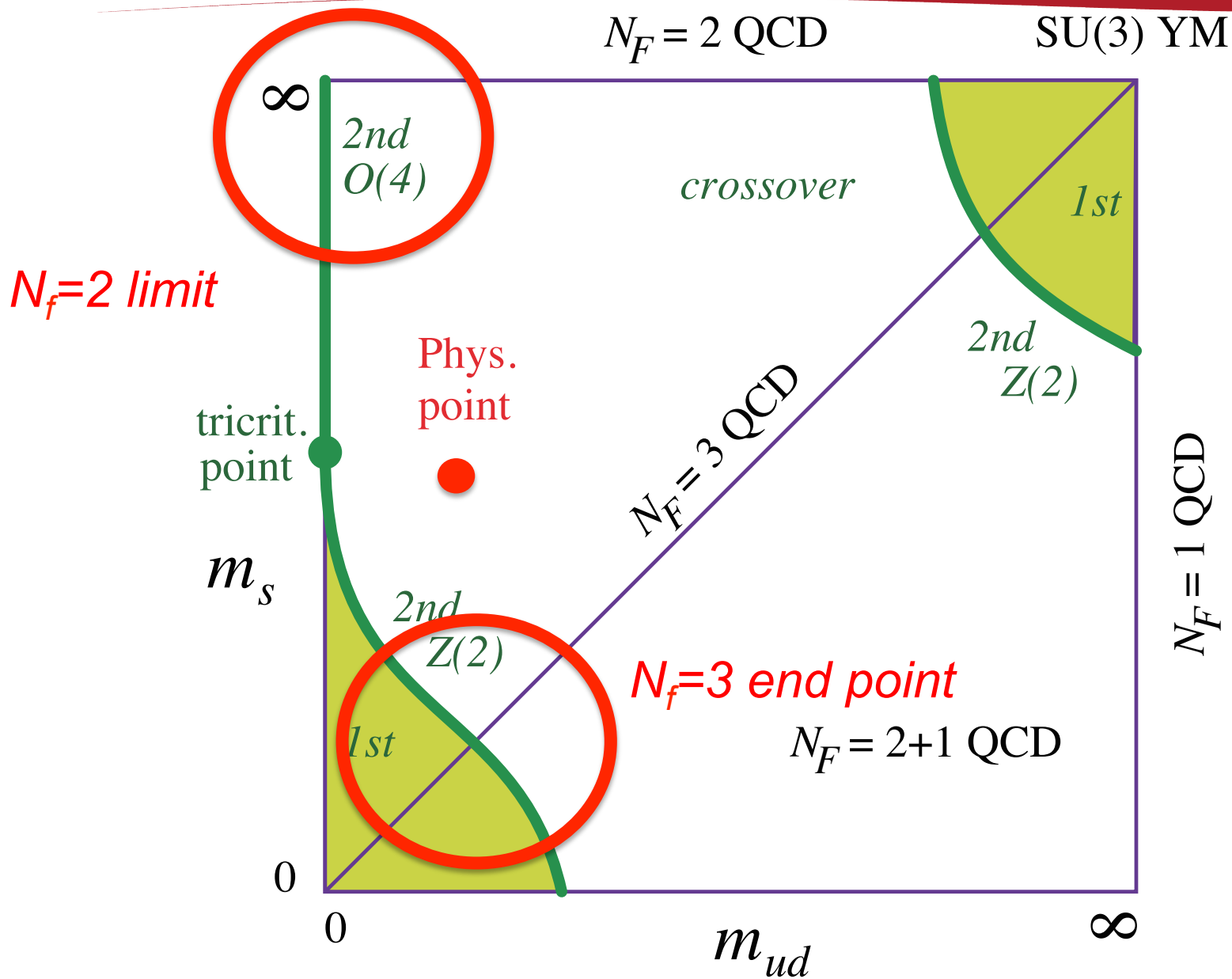
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# Loose ends?

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- $U_A(1)$  symmetry is broken by anomaly  $\partial_\mu J_\mu^{U_A(1)} = \frac{1}{8\pi^2} \text{Tr}(F_{\mu\nu} \tilde{F}_{\mu\nu})$  but its effect maybe negligible around the transition
- This may affect the *order and universality class* of the  $N_f=2$  transition at  $m_u=m_d=0$
- Long debate, not yet completely settled

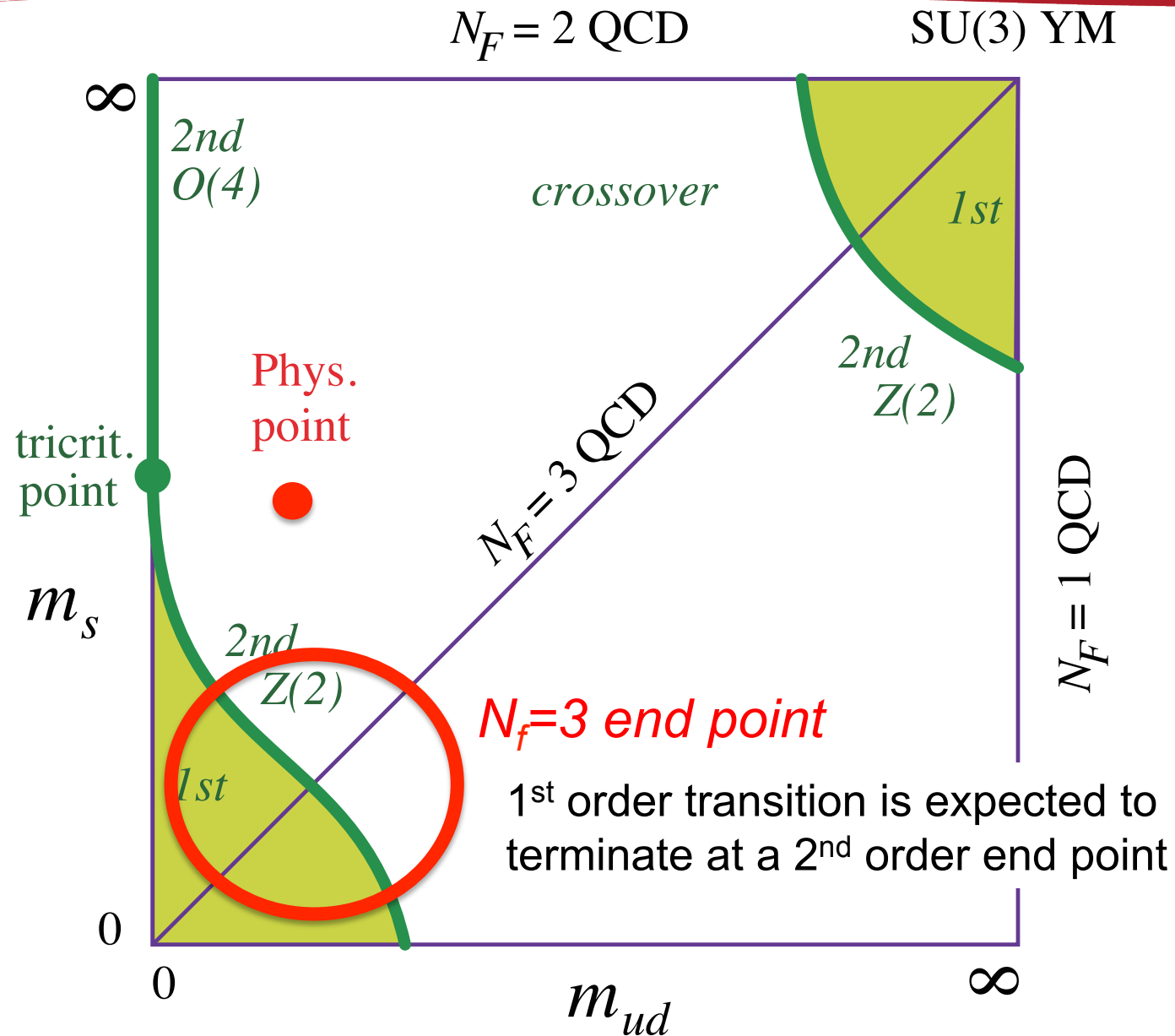
	$U_A(1)$ broken	$U_A(1)$ restored
symmetry	$SU(2) \otimes SU(2) \approx O(4)$	$SU(2) \otimes SU(2) \otimes U(1) \approx O(4) \otimes O(2)$
Pisarski-Wilczek(1984)	<i>2nd order</i>	<i>1st order</i> (no RG fixed point at 1-loop)
Pellissetto et al(2000)		<i>2nd order</i> (RG fixed point at 6-loop)
Aoki et al(2012)	<i>cannot be <math>O(4)</math></i>	<i>not visible in certain correlators</i>
Nakayama et al(2014)		<i>2nd order</i> (conformal bootstrap)



# My assessment



- No strong evidence of 1<sup>st</sup> order; 2<sup>nd</sup> order likely
- Theoretically important to settle
- But, quantitative effects likely limited at the physical point, since
  - $O(4)$  and  $O(4) \times O(2)$  critical exponents are similar  
(only about 10% difference)
  - Finite quark mass effects further softens the difference



- Pion mass marking the end point is very uncertain:

- For the staggered action, small values, *even consistent with zero?*

*Karsch et al (p4) '03*

*Forcrand-Philipsen (unimproved) '07*

*Endrodi et al (stout) '07*

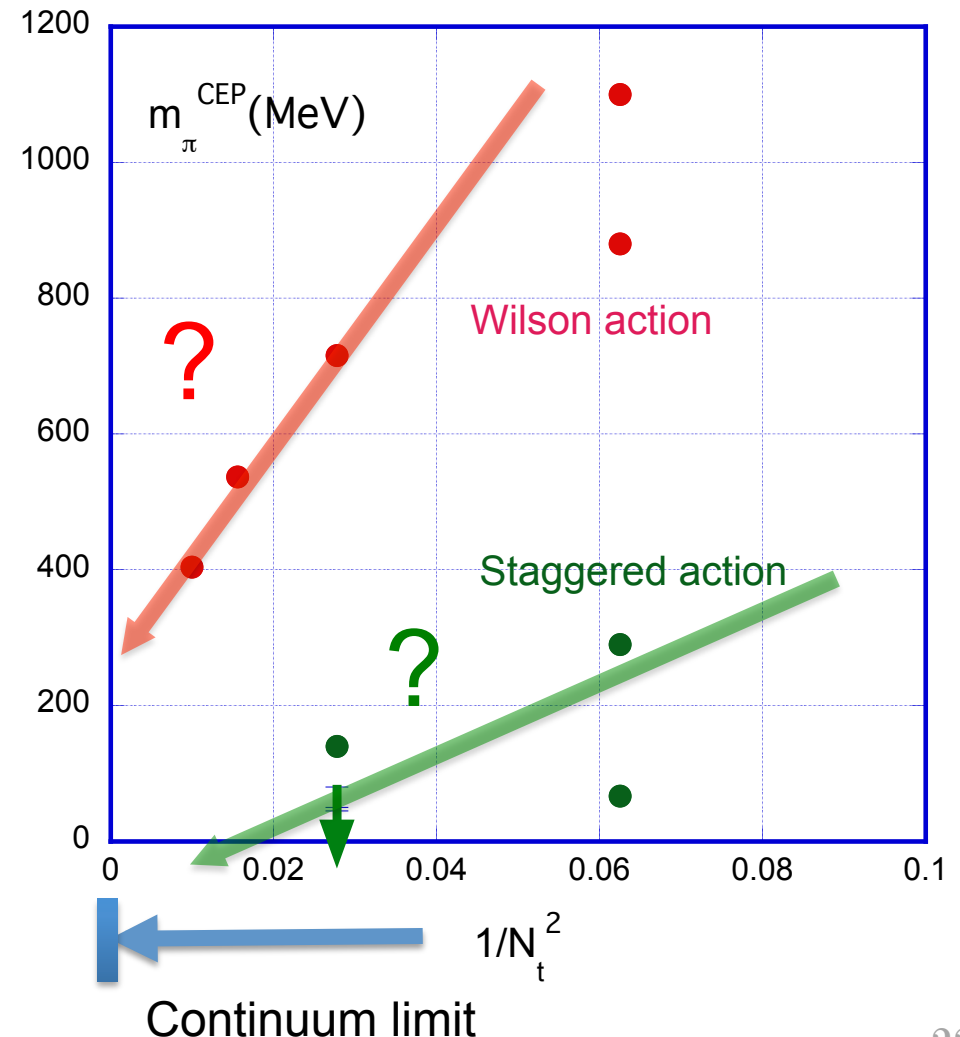
*Ding et al (HISQ) '11*

- For the Wilson action, large scaling violation, and *sizable difference from the staggered values*

*Iwasaki et al '96*

*Nakamura et al '14*

pion mass at the end point



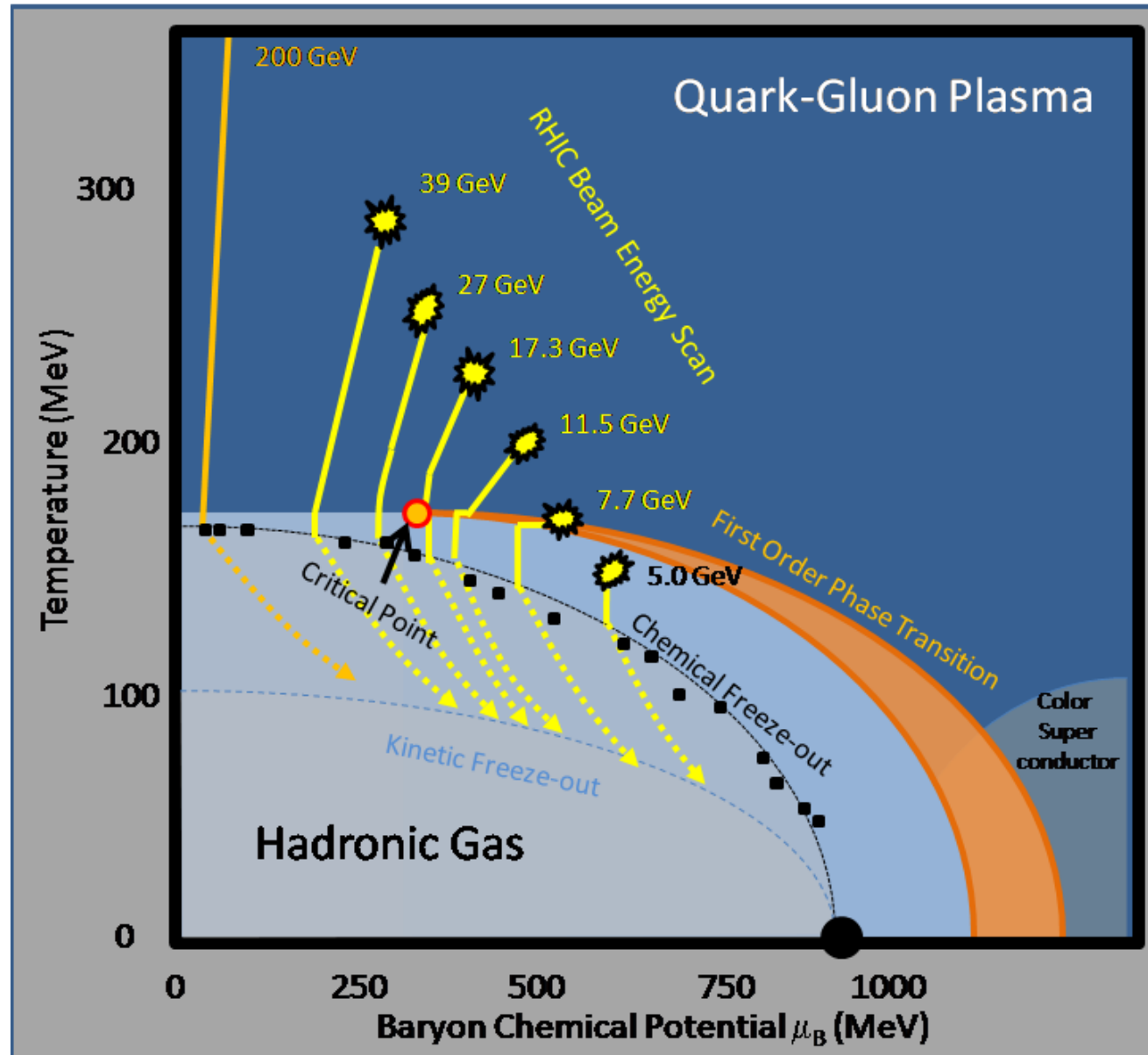


Non-zero but small  
baryon density

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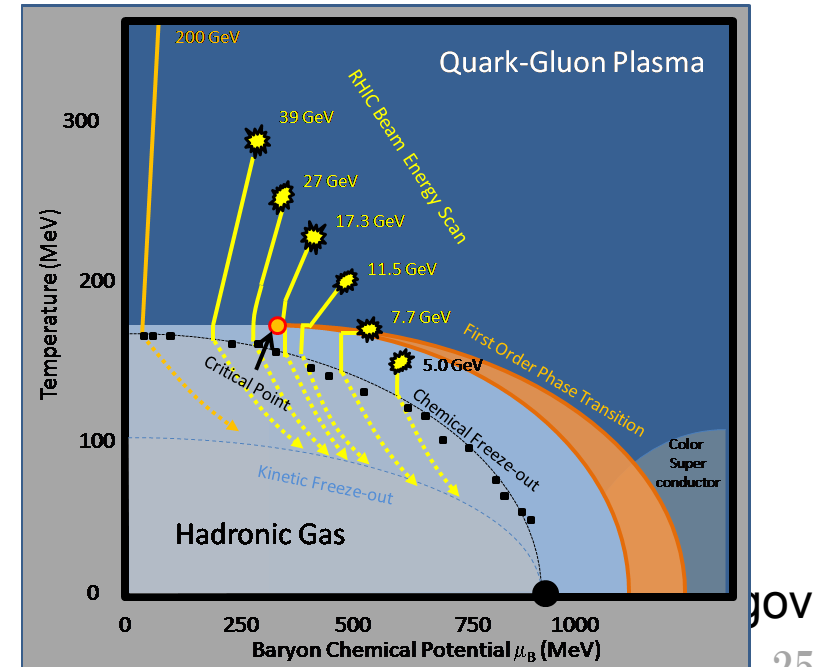


# How much do we really know?



From BNL@gov

- How does the pseudo critical temperature extend to non-zero density?
- How does it compare with phenomenological “freeze-out” temperatures?
- Is there a critical end point, and if yes, where?



- Standard Monte Carlo simulation methods not applicable
- Two alternative methods in use for small  $\mu_B / T$ 
  - Taylor expansion in  $\mu_B / T$

$$O(T, \mu_B) = O_0(T) + O_1(T) \frac{\mu_B}{T} + O_2(T) \left( \frac{\mu_B}{T} \right)^2 + \dots$$

expansion coefficients  $O_n(T)$  are calculable with simulations at  $\mu_B = 0$

- Analytic continuation from imaginary potential  $\mu_B / T = i\rho / T$

$$O(T, \mu_B = i\rho)$$

calculable by Monte Carlo, hence analytically continued to real potential

## ■ Numerical methods for small chemical potential:

■ Taylor expansion in  $\mu_B / T$

O. Kacmarek et al, Phys.Rev.D83:014504('11)

G. Endrodi et al, JHEP04, 001('11)

■ Imaginary chemical potential  $\mu_B / T = i\rho / T$

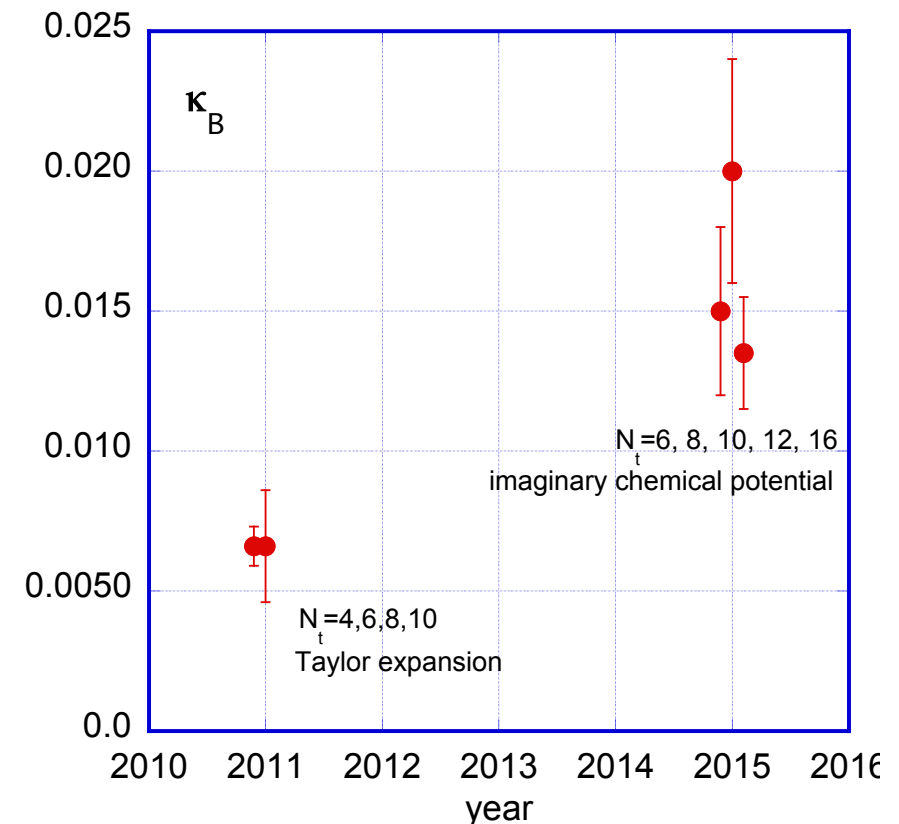
$$\frac{T_{pc}(\mu_B)}{T_{pc}(0)} = 1 - K_B \left( \frac{\mu_B}{T_{pc}(\mu_B)} \right)^2 + \dots$$

■ Recent continuum results sizably larger than the pioneering numbers

L. Cosmai et al, Lattice 2015

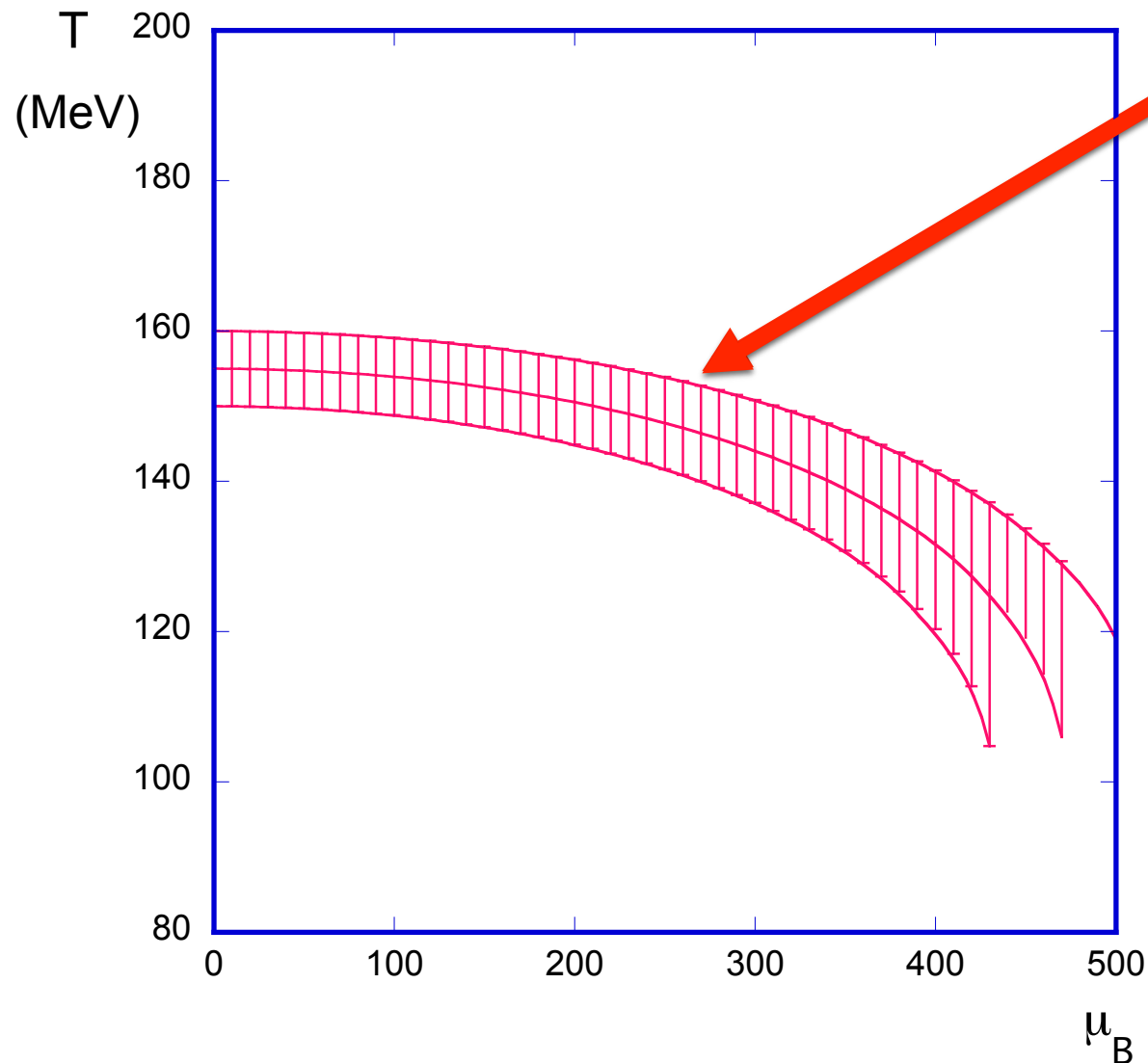
R. Bellwied et al, Lattice 2015

C. Bonatti et al, Lattice 2015



$$K_B \approx 0.0162 \pm 0.0017$$

My eyeball value



$$\frac{T_{pc}(\mu_B)}{T_{pc}(0)} = 1 - \kappa_B \left( \frac{\mu_B}{T_{pc}(\mu_B)} \right)^2 + \dots$$

$$\kappa_B \approx 0.0162 \pm 0.0017$$

$$T_{pc} \approx 155 \pm 5 \text{ MeV}$$

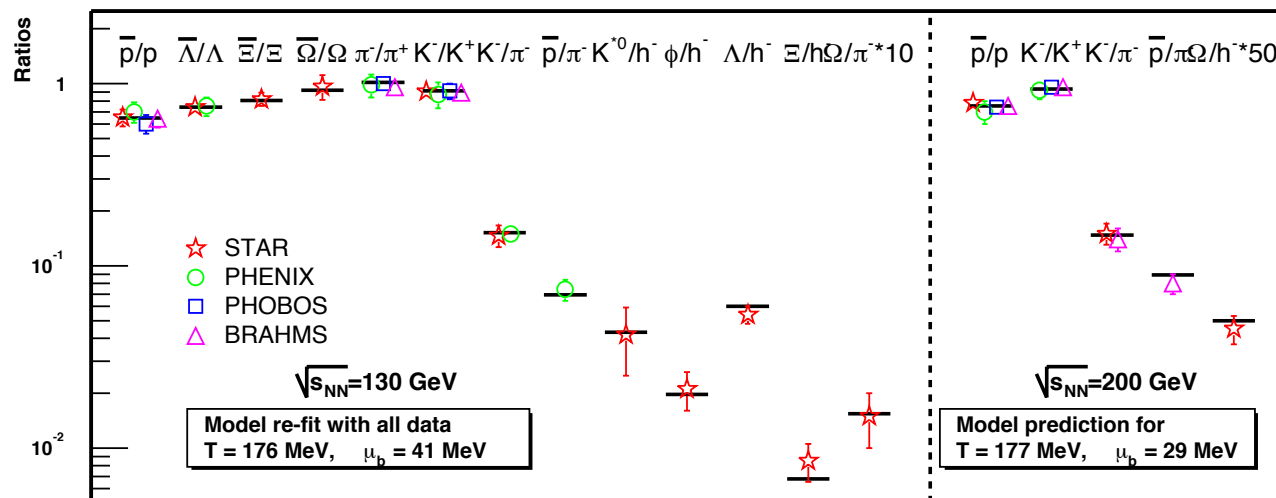
## ■ Statistical hadronization model

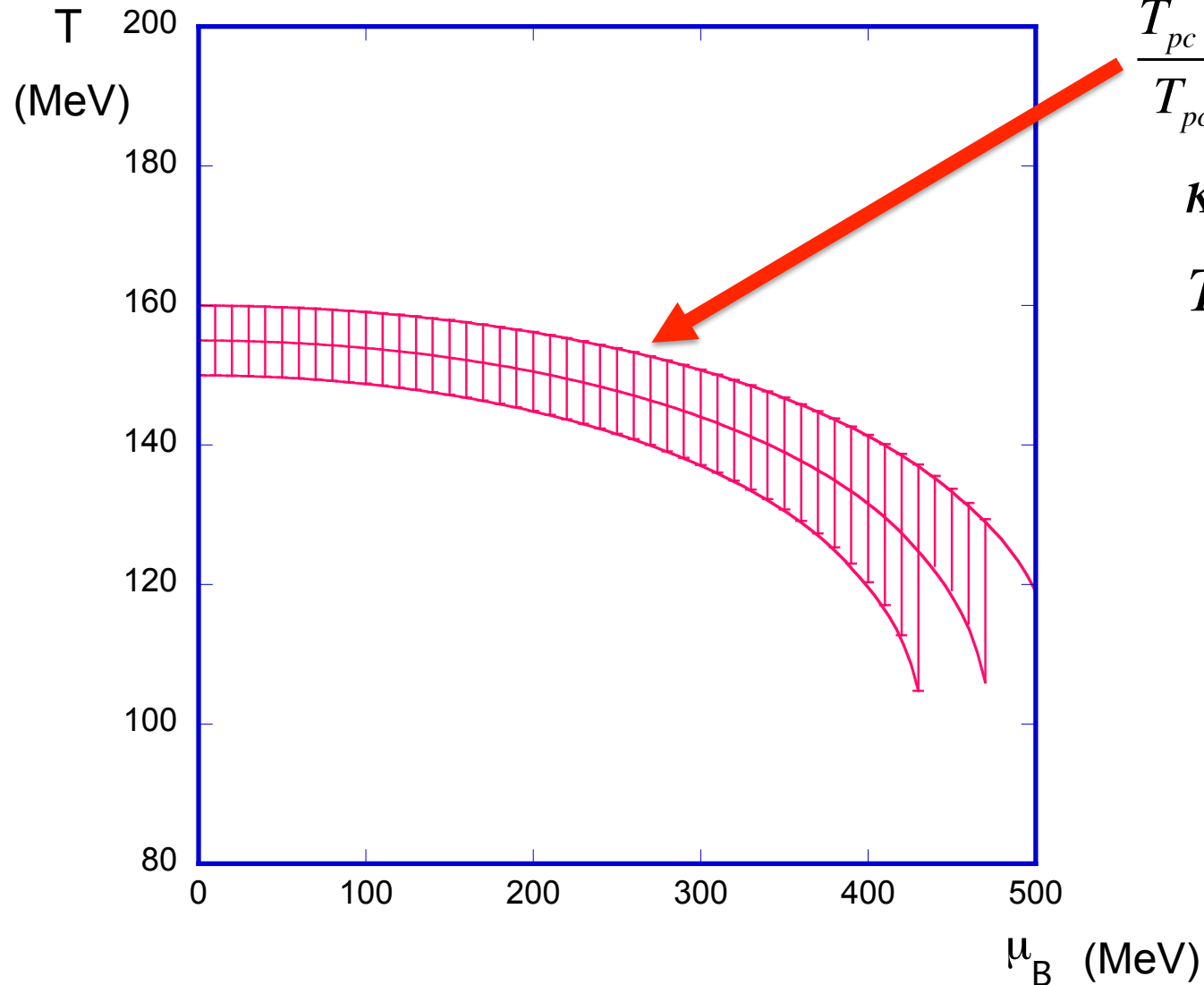
e.g., review by P. Braun-Munzinger et al, Nucl-th/0304013(2003)

- Statistical equilibrium at the time of hadronization (freeze out)
- Boltzmann distribution for hadron yields

$$N_h(T, \mu_B, V) = V \int \frac{d^3 p}{(2\pi)^3} \frac{1}{\exp\left(\frac{\epsilon_h(p) - B_h \mu_B}{T}\right) \pm 1}$$

- Fits to experimental yields provides an estimate of  $T, \mu_B$

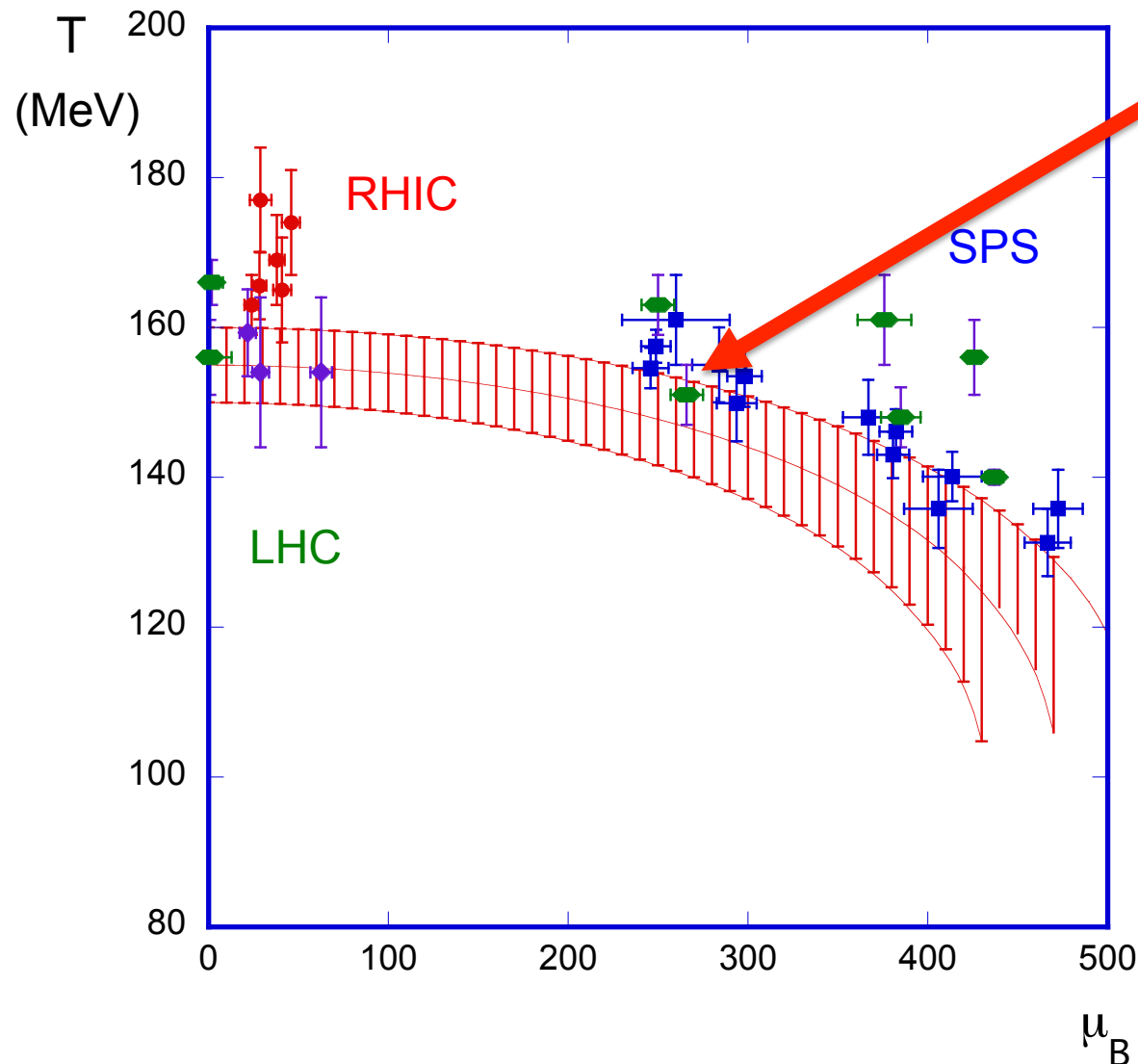




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$$T_{pc} \approx 155 \pm 5 \text{ MeV}$$

RICH, SPS

Cleymans et al, PRC.73.034905 (2006)

RICH Star collaboration

Abelev et al, PRC79, 034909 (2009)

LHC, SPS

Bocattini et al, PRL 111, 082302 (2012)



- Qualitative agreement but higher value from phenomenology?
- “Precise” comparison rather difficult since
  - Lattice QCD side
    - Pseudo critical temperature has inherent ambiguity, e.g., definition, choice of observable etc
    - Statistical and systematic errors of Monte Carlo calculation
  - Heavy ion phenomenology side
    - Inherent ambiguity in the extraction of  $T, \mu_B$   
i.e., at what stage do hadrons materialize?



# Conserved charge fluctuations and experiment

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- Moments of conserved quantities, e.g., electric charge  $Q$

$$\chi_1^Q(T, \mu_B) = \frac{1}{VT^3} \langle N_Q \rangle$$

$$\chi_2^Q(T, \mu_B) = \frac{1}{VT^3} \langle (\delta N_Q)^2 \rangle, \quad \delta N_Q = N_Q - \langle N_Q \rangle$$

$$\chi_3^Q(T, \mu_B) = \frac{1}{VT^3} \langle (\delta N_Q)^3 \rangle$$

$$\chi_4^Q(T, \mu_B) = \frac{1}{VT^3} \left[ \langle (\delta N_Q)^4 \rangle - \langle (\delta N_Q)^2 \rangle^2 \right]$$

$$M_Q = \langle N_Q \rangle, \quad \sigma_Q^2 = \langle (\delta N_Q)^2 \rangle$$

$$S_Q = \frac{\langle (\delta N_Q)^3 \rangle}{\sigma_Q^3}, \quad \kappa_Q = \frac{\langle (\delta N_Q)^4 \rangle}{\sigma_Q^4} - 3$$

or

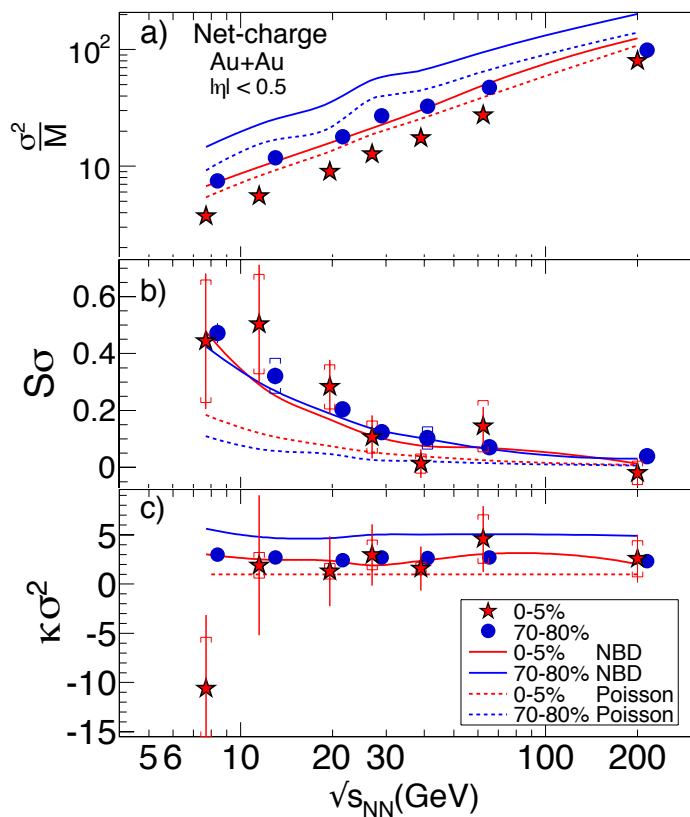
$$\frac{M_Q}{\sigma_Q^2} = \frac{\chi_1^Q}{\chi_2^Q} \equiv R_{12}^Q, \quad \frac{S_Q \sigma_Q^3}{M_Q} = \frac{\chi_3^Q}{\chi_1^Q} \equiv R_{31}^Q$$

- Measurable in experiments for given beam energy  $\sqrt{s}$
- Measurable in lattice QCD as functions of  $T, \mu_B$
- Hence, provide a *direct thermometer* of QGP, e.g.,

$$\frac{M_Q}{\sigma_Q^2}(\sqrt{s}) = R_{12}^Q(T, \mu_B)$$

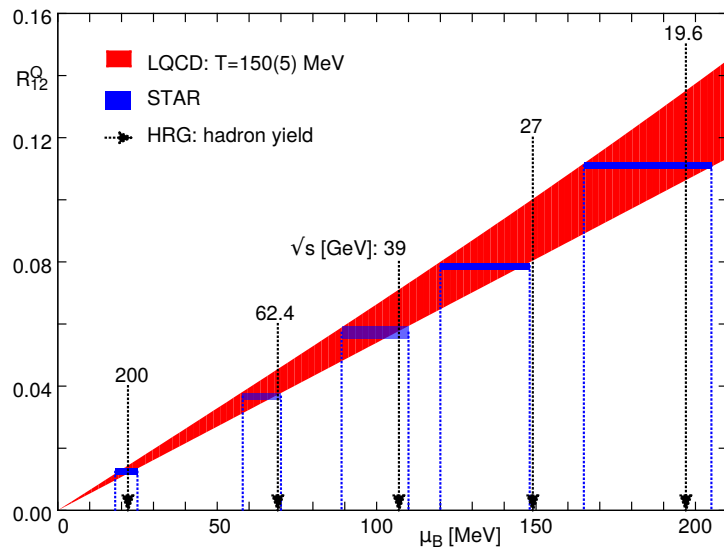
A. Bazavov et al, PRL 109, 192302 (2012)

S. Borsanyi et al, PRL 111, 062005 (2013)



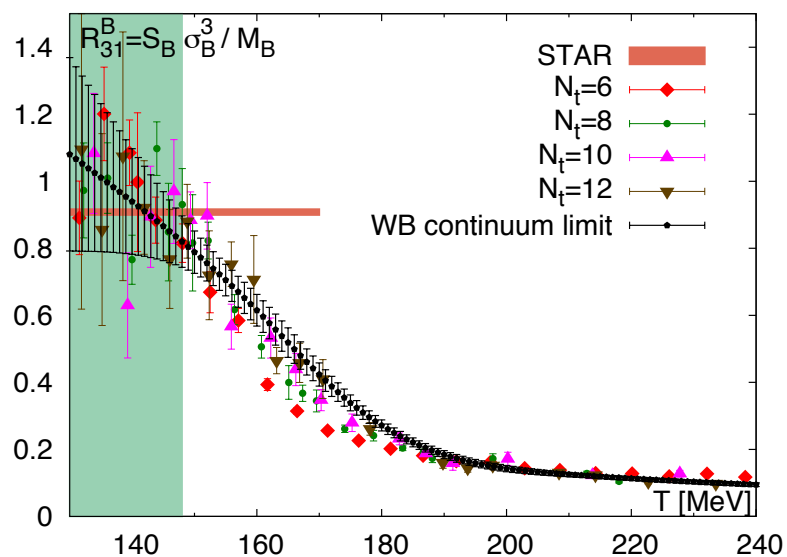
RHIC STAR measurement

L. Adamczyk et al, PRL 113, 092301(2014)



$\mu_B$  for given  $T$

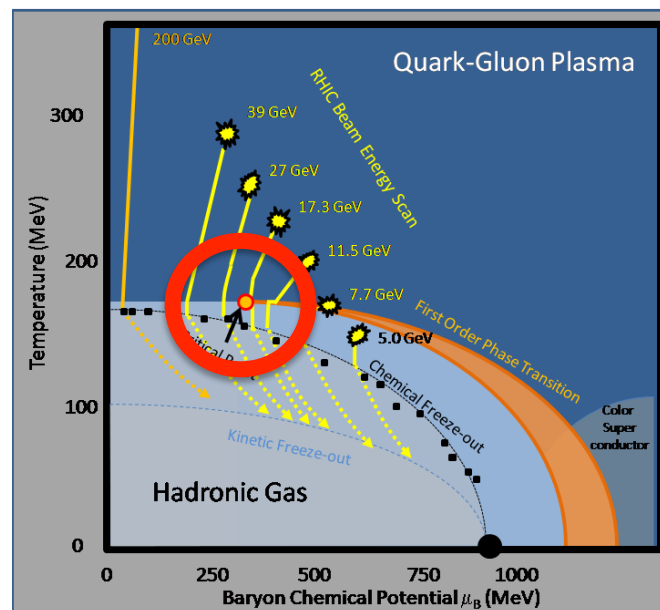
S. Mukherjee,  
PoS CPOD2014, 005

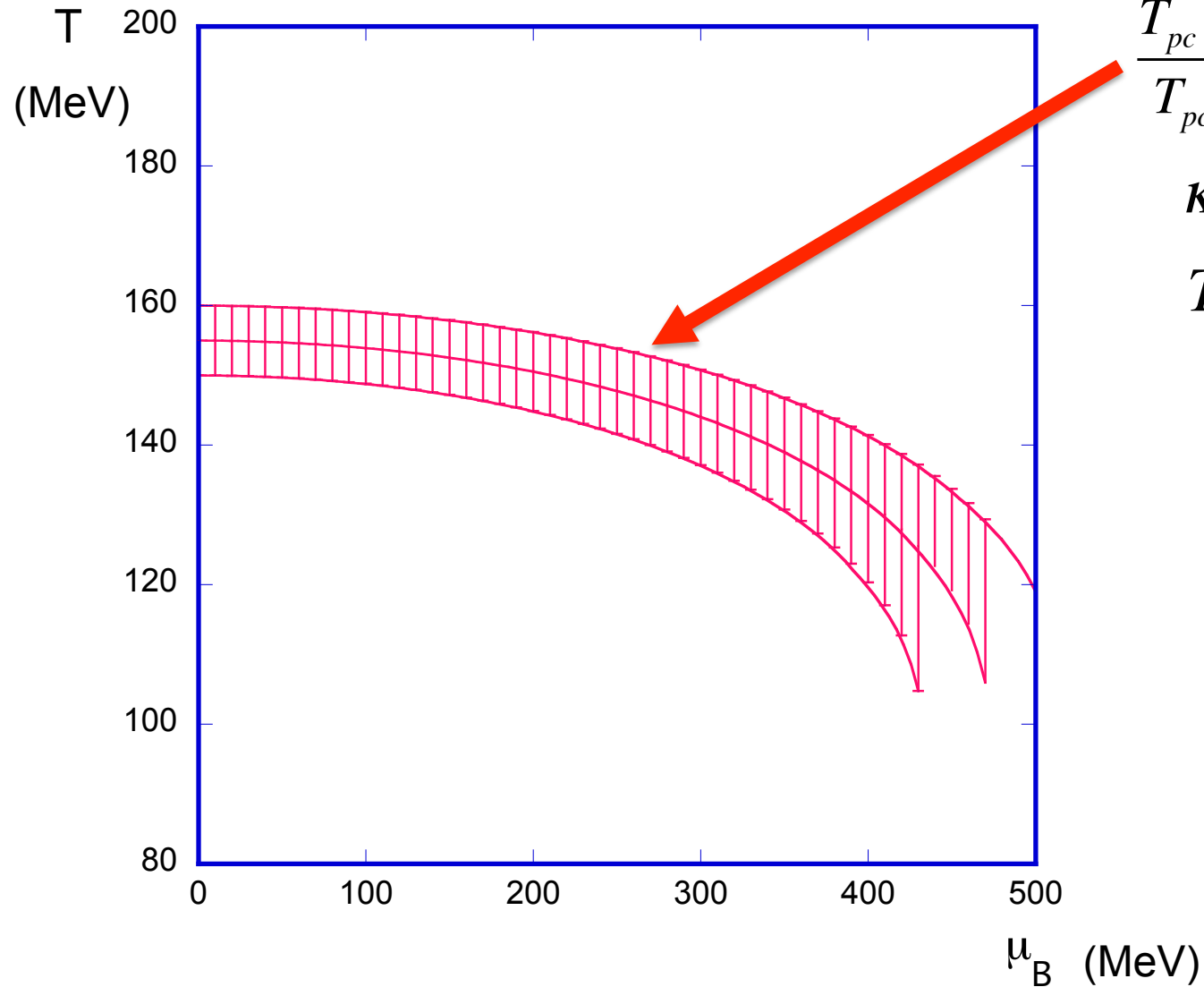


$T$  for  $\mu_B \approx 0$

S. Borsanyi et al,  
PRL 113, 052301(2014)

# Search for the critical end point?



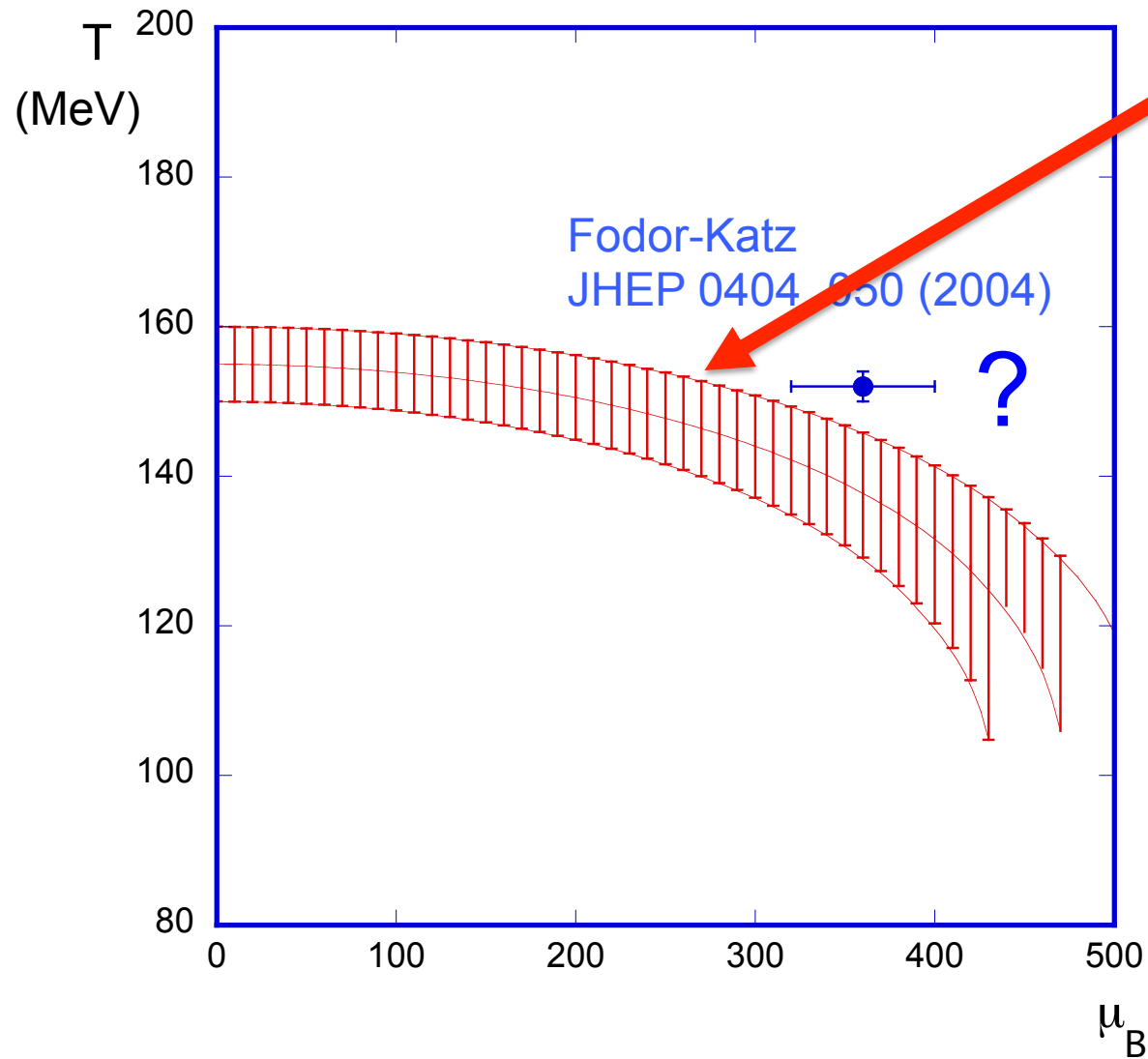


$$\frac{T_{pc}(\mu_B)}{T_{pc}(0)} = 1 - \kappa_B \left( \frac{\mu_B}{T_{pc}(\mu_B)} \right)^2 + \dots$$

$$\kappa_B \approx 0.0162 \pm 0.0017$$

$$T_{pc} \approx 155 \pm 5 \text{ MeV}$$

# Critical end point?



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$$\kappa_B \approx 0.0162 \pm 0.0017$$

$$T_{pc} \approx 155 \pm 5 \text{ MeV}$$

No new results in spite of many attempts, e.g.,  
reweighting,  
complex Langevin,  
Lefschetz thimble, ...



# *Summary*

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- *1<sup>st</sup> review of lattice QCD thermodynamics at LP/ICHEP conference series in 10 years!*
- *Significant progress in both lattice QCD and experiment (RHIC, LHC) in the mean time.*
- *Direct comparison of theory and experiment is now beginning to be possible.*
- *The location of critical end point remains elusive; hope for development in near future.*