





Physics at A Fixed-Target ExpeRiment at the LHC: AFTER@LHC

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Part I

Why a new fixed-target experiment for High-Energy Physics now?

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- They exhibit 4 decisive features,
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- which are essential assets to study
 - rare proton fluctuations at large x
 - vector boson production near threshold and other rare processes
 - nuclear dependence in heavy-ion collisions
 - observables involving gluons and the proton spin



 ADVANCE OUR UNDERSTANDING OF THE LARGE-X GLUON, ANTIQUARK AND HEAVY-QUARK CONTENT IN THE NUCLEON & NUCLEUS



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 - · Very large PDF uncertainties for $x \gtrsim 0.5$.

[could be crucial to characterise possible BSM discoveries]

- Proton charm content important to high-energy neutrino & cosmic-rays physics
- EMC effect is an open problem; studying a possible gluon EMC effect is essential
- · Relevance of nuclear PDF to understand the initial state of heavy-ion collisions
- · Search and study rare proton fluctuations



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- · Possible missing contribution to the proton spin: angular momentum
- Test of the QCD factorisation framework
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- · HEAVY-ION COLLISIONS TOWARDS LARGE RAPIDITIES
- Explore the longitudinal expansion of QGP formation with new hard probes
- Test the factorisation of cold nuclear effects from p + A to A + B collisions
- · Test the formation of azimuthal asymmetries: hydrodynamics vs. initial-state radiation

Part II

A fixed-target experiment using the LHC beam(s): AFTER@LHC

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- Bad thing: high multiplicity ⇒ absorber ⇒ physics limitation



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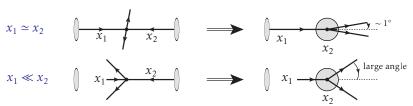
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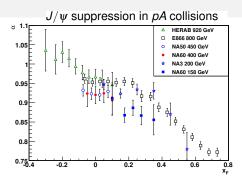
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backward physics = large- x_2 physics

 $(x_F \rightarrow -1)$

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 $(x_F \rightarrow -1)$ J/ψ suppression in pA collisions 1.05 NA60 158 GeV 0.95 0.9 0.85 0.8 0.75

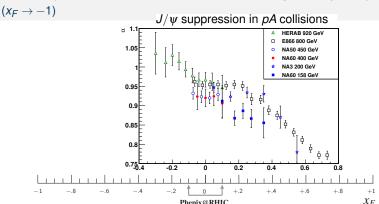
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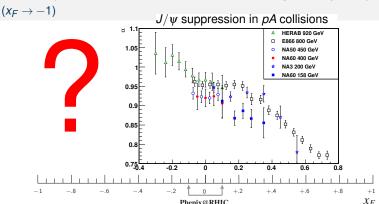
0.2

0.4

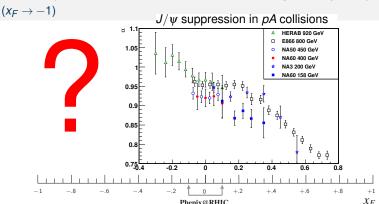
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- If we measure $\Upsilon(b\bar{b})$ at $y_{\rm cms}\simeq -2.5 \ \Rightarrow x_F\simeq {2m_{\Upsilon}\over \sqrt{s}} \sinh(y_{\rm cms})\simeq -1$

Part III

Colliding the LHC beams on fixed targets: 2 options

The extracted-beam option

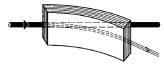
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E. Uggerhøj, U.I Uggerhøj, NIM B 234 (2005) 31, Rev. Mod. Phys. 77 (2005) 1131

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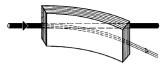
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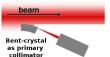
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★ Illustration for collimation





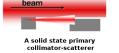
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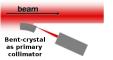
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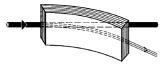
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LUA9 proposal approved by the LHCC

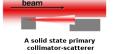
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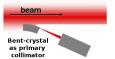
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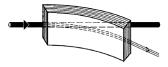
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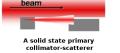
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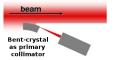
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- ★ 2 crystals and 2 goniometers already installed in the LHC beampipe
- ★ CRYSBEAM: ERC funded project to extract the LHC beams

with a bent crystal

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Target	ρ (g.cm ⁻³)	A	£ (μb-1.s-1)	∫£ (fb-1.yr-1)	
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• For pp and pd collisions : $\mathcal{L}_{H_2/D_2} \simeq 20 \text{ fb}^{-1} y^{-1}$

3 orders of magnitude larger than RHIC



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- Target density: $\frac{\rho}{P} = c = \frac{A}{22400} \text{bar}^{-1} g \, cm^{-3} \Rightarrow \mathscr{L} = \Phi_{beam} \times (\frac{\mathscr{N}_A}{22400} \times P \times \ell)$ [1 mole of a perfect gas occupies 22 400 cm³ at 273 K and 1 bar]

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- Target density: $\frac{\rho}{P} = c = \frac{A}{22400} \text{bar}^{-1} g \, cm^{-3} \Rightarrow \mathcal{L} = \Phi_{beam} \times (\frac{\mathcal{N}_A}{22400} \times P \times \ell)$ [1 mole of a perfect gas occupies 22 400 cm³ at 273 K and 1 bar]
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- Instantaneous Luminosity: $\mathcal{L} = \Phi_{beam} \times N_{target} = N_{beam} \times (\rho \times \ell \times \mathcal{N}_A)/A$
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C. Barschel, P. Lenisa, A. Nass, and E. Steffens, Adv.Hi.En.Phys. (2015) ID:463141



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• A specific gas target could be a competitive alternative to the beam extraction

Part IV

AFTER@LHC: a selection of key measurements

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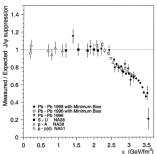


Fig. 7. Measured J/ψ production yields, normalised to the yields expected assuming that the only source of suppression is the ordinary absorption by the nuclear medium. The data is shown as



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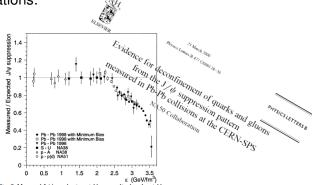


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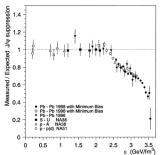
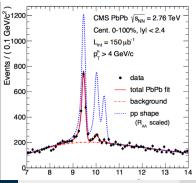


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- Enough stat to perform the same study as CMS at low energy



PRL **109**, 222301 (2012)

Selected for a Viewpoint in Physics
PHYSICAL REVIEW LETTERS

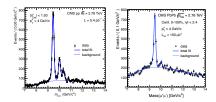
week ending 30 NOVEMBER 2012



Observation of Sequential Y Suppression in PbPb Collisions

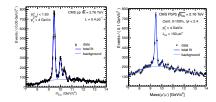
S. Chatrchyan *et al.** (CMS Collaboration)

CMS PRL 109 222301 (2012), JHEP04(2014)103



$\frac{[\Upsilon(nS)/\Upsilon(1S)]_{ij}}{[\Upsilon(nS)/\Upsilon(1S)]_{pp}}$	28	3 <i>S</i>
PbPb	$0.21 \pm 0.07 (stat.) \pm 0.02 (syst.)$	$0.06 \pm 0.06 \text{ (stat.)} \pm 0.06 \text{ (syst.)}$

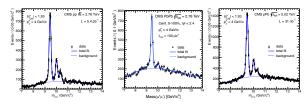
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In addition to QGP formation, differences between quarkonium production yields in PbPb and pp collisions can also arise from cold-nuclear-matter effects [21]. However, such effects should have a small impact on the double ratios reported here. Initial-state nuclear effects are expected to affect similarly each of the three Υ states, thereby canceling out in the ratio. Final-state "nuclear absorption" becomes weaker with increasing energy [22] and is expected to be negligible at the LHC [23].

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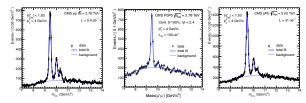
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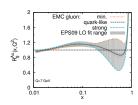
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If the effects responsible for the relative nS/1S suppression in pPb collisions factorise, they could be responsible for half of the PbPb relative suppression !!!

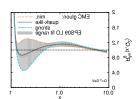


Large-x gluon nPDF: unknown

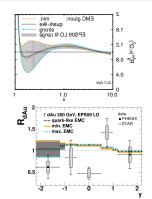
Gluon EMC effect: unknown



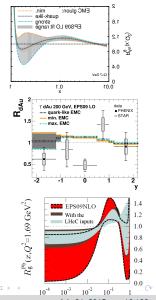
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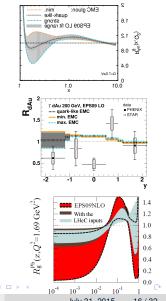


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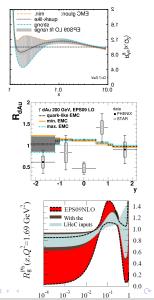
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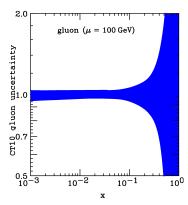
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- AFTER allows for extensive studies of gluon sensitive probes in pA
- Unique potential for gluons at x > 0.1



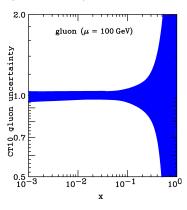
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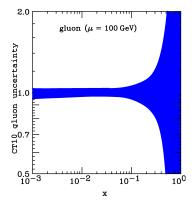
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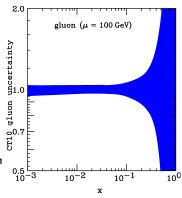


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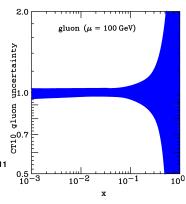
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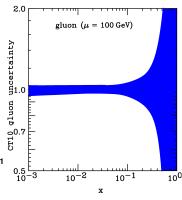
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Multiple probes needed to check factorisation



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1.25 NNPDP23 NNLO
1.25 CT10 NNLO
MSTW2008 NNLO
1.60.95
0.85 C.0.95

LHC 8 TeV - Ratio to NNPDF2.3 NNLO - α_e = 0.118

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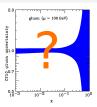
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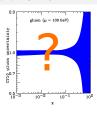
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Large-*x* gluons: important to characterise some possible BSM findings at the LHC



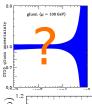


Gluon PDF for the neutron unknown



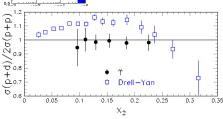
Gluon PDF for the neutron unknown possible experimental probes

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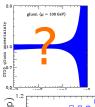
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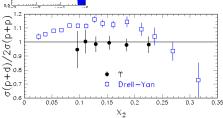
Pioneer measurement by E866

- using $\Upsilon \rightarrow Q^2 \simeq 100 \text{ GeV}^2$
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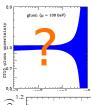


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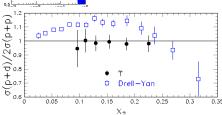
could be extended with AFTER

- using J/ψ , ..., C = +1 onia, ...
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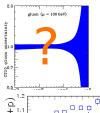
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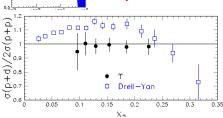
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1m Liq. H ₂	20 fb ⁻¹	4.0×10^{8}	9.0 × 10 ⁵
1m Liq. D_2	$24 \; \text{fb}^{-1}$	9.6×10^8	1.9×10^6



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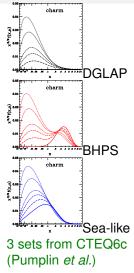
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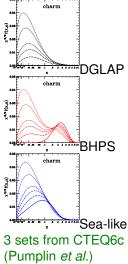
If $g_n(x) - g_p(x)$ is too small, this measurement would anyhow be sensitive to the EMC and Fermi-motion effects in the deuteron

• Heavy-quark distributions (at high x)

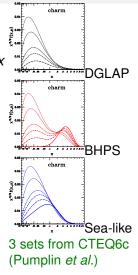
- Heavy-quark distributions (at high *x*)
 - Pin down intrinsic charm, ... at last



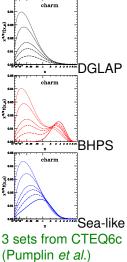
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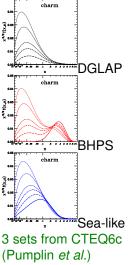


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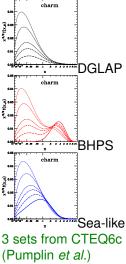
(Pumplin et al.)

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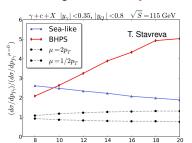
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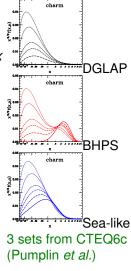
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(Pumplin et al.)

- Heavy-quark distributions (at high x)
 - Pin down intrinsic charm, ... at last
 - c(x) & b(x) & the 5-flavour scheme at large x for BSM studies F.Maltoni,..., JHEP 1207 (2012) 022 requires
 - several complementary measurements
 - good coverage in the target-rapidity region
 - high luminosity to reach large x





S. J. Brodsky, et al., Adv.Hi.En.Phys. (2015) ID:231547.

Gluon Sivers effect
 ← correlation between
 the gluon transverse momentum & the proton spin

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 - Transverse single spin asymetries

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A. Bacchetta, et al., PRL 99 (2007) 212002 J.W. Qiu, et al., PRL 107 (2011) 062001

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- the target-rapidity region corresponds to high x^{\uparrow} where the k_T -spin correlation is the largest
- In general, one can carry out an extensive spin-physics program
- Even w/o target polarisation via the Boer-Mulders effect [backup slides] of the Boer-Mulders of the Boer

Part V

First simulation results

First simulation: is the boost an issue?

B. Trzeciak, L. Massacrier et al., 1504.05145 [hep-ex], to appear in Adv.Hi.En.Phys.



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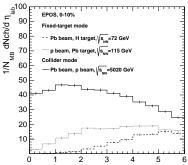
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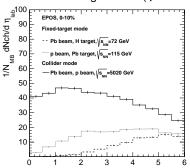


• Despite the boost, the multiplicity in the LHCb acceptance [forward η] is lower in the fixed mode than in the collider mode (at higher \sqrt{s})

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- Despite the boost, the multiplicity in the LHCb acceptance [forward η] is lower in the fixed mode than in the collider mode (at higher \sqrt{s})
- Simulation backed-up with a comparison of the number-of-track distribution between simulations at the detector level and data
 Z. Yang, private comm.

FAST SIMULATIONS FOR QUARKONIA (pp \sqrt{s} = 115 GeV) USING LHCB RECONSTRUCTION PARAMETERS

- ☐ Simulations with Pythia 8.185
- □ LHCb detector is NOT simulated but LHCb reconstruction parameters are introduced in the fast simulation (resolution, analysis cuts, efficiencies...)

Requirements

Momentum resolution : $\Delta p/p = 0.5\%$ Muon identification efficiency: 98%

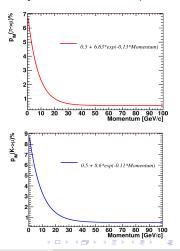
Cuts at the single muon level

 $2 < \eta_{\mu} < 5$ $p_{\tau}^{\mu} > 0.7 \text{ GeV/c}$

Muon misidentification

If π and K decay before the calorimeters (12m), they are rejected by the tracking Else a misidentification probability is applied

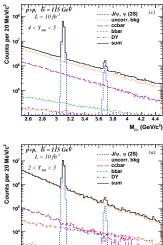
Performance of the muon identification at LHCb, F. Achilli et al, arXiv:1306.0249

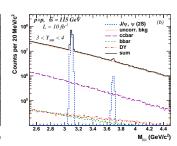


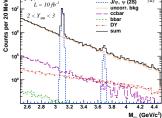
Trzeciak, L. Massacrier et al., 1504.05145 [hep-ex

Charmonium background & its rapidity dependence

B. Trzeciak, L. Massacrier et al., 1504.05145 [hep-ex], to appear in Adv.Hi.En.Phys.

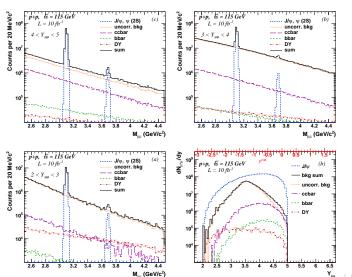






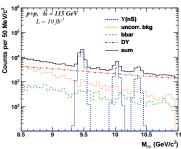
Charmonium background & its rapidity dependence

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Bottomonium background & signal reach

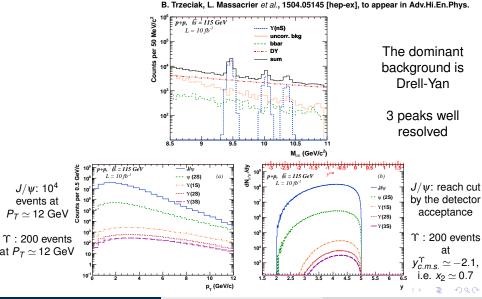
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The dominant background is Drell-Yan

3 peaks well resolved

Bottomonium background & signal reach



Part VI

Further readings

Heavy-Ion Physics

- Gluon shadowing effects on J/ψ and ↑ production in p+Pb collisions at √s_{NN} = 115 GeV and Pb+p collisions at √s_{NN} = 72 GeV at AFTER@LHC by R. Vogt. Adv.Hi.En.Phys. (2015) ID:492302.
- Prospects for open heavy flavor measurements in heavy-ion and p+A collisions in a fixed-target experiment at the LHC by D. Kikola. Adv.Hi.En.Phys. (2015) ID:783134
- Quarkonium suppression from coherent energy loss in fixed-target experiments using LHC beams by F. Arleo, S.Peigné. [arXiv:1504.07428 [hep-ph]]. Adv.Hi.En.Phys. (2015) ID:961951
- Anti-shadowing Effect on Charmonium Production at a Fixed-target Experiment Using LHC Beams by K. Zhou, Z. Chen, P. Zhuang. arXiv:1507.05413 [nucl-th].
- Lepton-pair production in ultraperipheral collisions at AFTER@LHC
 By J.P. Lansberg, L. Szymanowski, J. Wagner. arXiv:1504.02733 [hep-ph].
- Quarkonium Physics at a Fixed-Target Experiment using the LHC Beams. By J.P. Lansberg, S.J. Brodsky, F. Fleuret, C. Hadjidakis. [arXiv:1204.5793 [hep-ph]]. Few Body Syst. 53 (2012) 11.

Spin physics

- Transverse single-spin asymmetries in proton-proton collisions at the AFTER@LHC experiment by K. Kanazawa, Y. Koike, Andreas Metz, and D. Pitonyak. [arXiv:1502.04021 [hep-ph]. Adv.Hi.En.Phys. (2015) ID:257934.
- Transverse single-spin asymmetries in proton-proton collisions at the AFTER@LHC experiment in a TMD factorisation scheme
 by M. Anselmino, U. D'Alesio, and S. Melis. [arXiv:1504.03791 [hep-ph]]. Adv.Hi.En.Phys. (2015) ID:475040.
- The gluon Sivers distribution: status and future prospects
 by D. Boer, C. Lorcé, C. Pisano, and J. Zhou. [arXiv:1504.04332 [hep-ph]].
 Adv.Hi.En.Phys. (2015) ID:371396
- Azimuthal asymmetries in lepton-pair production at a fixed-target experiment using the LHC beams (AFTER)
 By T. Liu, B.Q. Ma. [arXiv:1203.5579 [hep-ph]]. Eur.Phys.J. C72 (2012) 2037.
- Polarized gluon studies with charmonium and bottomonium at LHCb and AFTER
 By D. Boer, C. Pisano. [arXiv:1208.3642 [hep-ph]]. Phys.Rev. D86 (2012) 094007.



Hadron structure

- Double-quarkonium production at a fixed-target experiment at the LHC (AFTER@LHC).
 by J.P. Lansberg, H.S. Shao. [arXiv:1504.06531 [hep-ph]].
- Next-To-Leading Order Differential Cross-Sections for Jpsi, psi(2S) and Upsilon Production in Proton-Proton Collisions at a Fixed-Target Experiment using the LHC Beams (AFTER@LHC)
 by Y. Feng, and J.X. Wang. Adv.Hi.En.Phys. (2015) ID:726393, in press.
- η_c production in photon-induced interactions at a fixed target experiment at LHC as a probe of the odderon
 By V.P. Goncalves, W.K. Sauter. arXiv:1503.05112 [hep-ph].Phys.Rev. D91 (2015) 9, 094014.
- A review of the intrinsic heavy quark content of the nucleon by S. J. Brodsky, A. Kusina, F. Lyonnet, I. Schienbein, H. Spiesberger, and R. Vogt. Adv.Hi.En.Phys. (2015) ID:231547, in press.
- Hadronic production of ≡_{cc} at a fixed-target experiment at the LHC
 By G. Chen et al.. [arXiv:1401.6269 [hep-ph]]. Phys.Rev. D89 (2014) 074020.



Feasibility study and technical ideas

- Feasibility studies for quarkonium production at a fixed-target experiment using the LHC proton and lead beams (AFTER@LHC) by L. Massacrier, B. Trzeciak, F. Fleuret, C. Hadjidakis, D. Kikola, J.P.Lansberg, and H.S. Shao arXiv:1504.05145 [hep-ex]. Adv.Hi.En.Phys. (2015) ID:986348
- A Gas Target Internal to the LHC for the Study of pp Single-Spin Asymmetries and Heavy Ion Collisions by C. Barschel, P. Lenisa, A. Nass, and E. Steffens. Adv.Hi.En.Phys. (2015) ID:463141
- Quarkonium production and proposal of the new experiments on fixed target at LHC by N.S. Topilskaya, and A.B. Kurepin. Adv.Hi.En.Phys. (2015) ID:760840

Generalities

Physics Opportunities of a Fixed-Target Experiment using the LHC Beams
 By S.J. Brodsky, F. Fleuret, C. Hadjidakis, J.P. Lansberg. [arXiv:1202.6585 [hep-ph]].
 Phys.Rept. 522 (2013) 239.

Part VII

Conclusion and outlooks

 Both p and Pb LHC beams can be extracted without disturbing the other experiments

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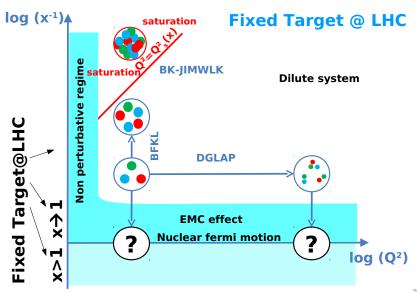


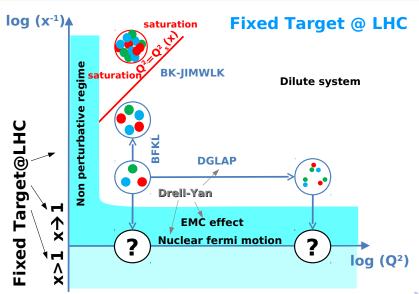
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- A wealth of possible measurements:

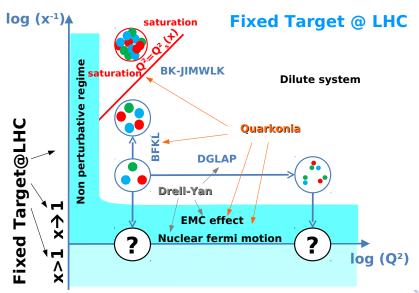
DY, Open b/c, jet correlation, UPC... (not mentioning secondary beams) 32/32

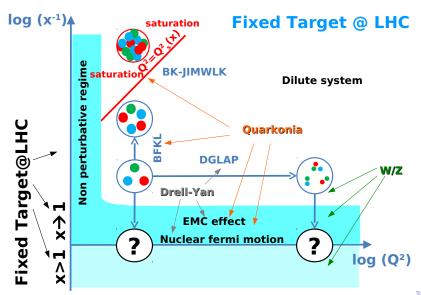
Part VIII

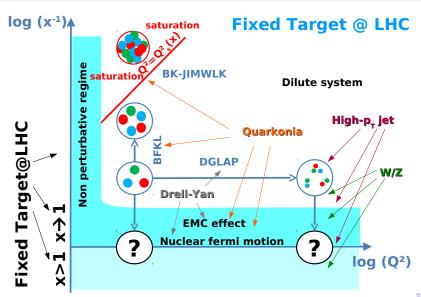
Backup slides











Gas target

C. Barschel, P. Lenisa, A. Nass, and E. Steffens, Adv.Hi.En.Phys. (2015) ID:463141

TABLE 1: Comparison of gas targets in storage rings with a hypothetical target for the proposed AFTER@LHC initiative [1, 2]. The target gas ¹H, ²D, or ³He is assumed to be spin polarized.

Storage ring	Particle	E _{max} [GeV]	Target type	L [m]	T [K]	L _{max} [1/cm ² s]	Remarks	Reference
HERA-e DESY (term. 2007)	e [±] pol.	27.6	Cell ¹ H, ² D, ³ He	0.4	100 25	$2.5 \cdot 10^{31} \\ 2.5 \cdot 10^{32}$	HERMES exp. 1995–2007	[9]
RHIC-p BNL	p pol.	250	Jet	_	_	$1.7 \cdot 10^{30}$	Absolute p polarimeter	[10]
COSY FZ Jülich	p, d pol.	3.77 T = 49.3 MeV	Cell ¹ H, ² D Cell ¹ H	0.4	300	$10^{29} \\ 2.75 \cdot 10^{29}$	ANKE exp. PAX exp.	[4, 5] [11]
LHC CERN (proposed)	p unpol. heavy ions	7,000 2,760 · A	Cell ^{1}H , ^{2}D Xe $M \approx 131$	1.0	100 ≥100	$10^{33} \\ 10^{27} - 10^{28}$	Based on techn. of HERMES target	this paper

 \rightarrow beam lifetime with $\mathscr{L}_{pp}=10^{33} cm^{-2} s^{-1}\text{=}~10~nb^{-1} s^{-1} of~2\times 10^6~s$ (or 23 days).

Accessing the large x glue with quarkonia:

PYTHIA simulation $\sigma(y) / \sigma(y=0.4)$ statistics for one month 5% acceptance considered

Statistical relative uncertainty Large statistics allow to access very backward region

Gluon uncertainty from MSTWPDF

- only for the gluon content of the target
- assuming

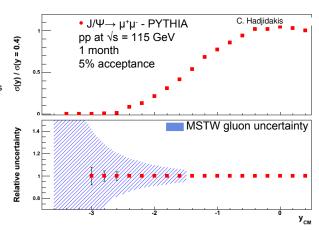
$$x_g = M_{J/\Psi}/\sqrt{s} e^{-yCM}$$

J/Ψ

$$y_{CM} \sim 0 \rightarrow x_g = 0.03$$

 $y_{CM} \sim -3.6 \rightarrow x_g = 1$

Y: larger x_g for same y_{CM} $y_{CM} \sim 0 \rightarrow x_g = 0.08$ $y_{CM} \sim -2.4 \rightarrow x_g = 1$



⇒ Backward measurements allow to access large x gluon pdf

Assuming that we understand the quarkonium-production mechanisms

Distribution of linearly polarised gluons in unpolarised protons



Distribution of linearly polarised gluons in unpolarised protons

PHYSICAL REVIEW D 86, 094007 (2012)

Polarized gluon studies with charmonium and bottomonium at LHCb and AFTER

Daniël Boer*

Theory Group, KVI, University of Groningen, Zemikelaan 25, NL-9747 AA Groningen, The Netherlands

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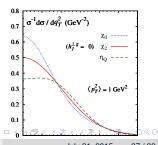
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- Affect the low P_T spectra:

$$\frac{1}{\sigma} \frac{d\sigma(\eta_O)}{d\mathbf{q}_T^2} \propto 1 - R(\mathbf{q}_T^2) & \frac{1}{\sigma} \frac{d\sigma(\chi_{0,O})}{d\mathbf{q}_T^2} \propto 1 + R(\mathbf{q}_T^2)$$

(R involves $f_1^g(x, k_T, \mu)$ and $h_1^{\perp g}(x, k_T, \mu)$)



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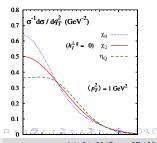
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(*R* involves $f_1^g(x, k_T, \mu)$ and $h_1^{\perp g}(x, k_T, \mu)$)

 The boost is of great help to access low P_T P-wave quarkonia



PHYSICAL REVIEW D 86, 094007 (2012)

Polarized gluon studies with charmonium and bottomonium at LHCb and AFTER

Daniël Boer*

Theory Group, KVI, University of Groningen, Zemikelaan 25, NL-9747 AA Groningen. The Netherlands

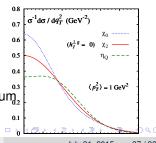
Cristian Pisano

Istituto Nazionale di Fisica Nucleare, Sezione di Cagliari, C.P. 170, I-09042 Monserrato (CA), Italy

- Low P_T C-even quarkonium production is a good probe of the distribution of linearly polarised gluons in unpolarised protons: $h_1^{\perp g}$
- Affect the low P_T spectra:

$$\frac{1}{\sigma} \frac{d\sigma(\eta_O)}{d\mathbf{q}_T^2} \propto 1 - R(\mathbf{q}_T^2) \& \frac{1}{\sigma} \frac{d\sigma(\chi_{0,O})}{d\mathbf{q}_T^2} \propto 1 + R(\mathbf{q}_T^2)$$
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- The boost is of great help to access low P_T P-wave quarkonia
- $h_{1}^{\perp g}$ is connected to the Higgs transverse-momentum, distribution D. Boer, et al. PRL 108 (2012) 032002





PRL 112, 212001 (2014) PHYSICAL REVIEW LETTERS

week ending 30 MAY 2014

Accessing the Transverse Dynamics and Polarization of Gluons inside the Proton at the LHC

Wilco J. den Dunnen, ^{1,2} Jean-Philippe Lansberg, ^{2,2} Cristian Pisano, ^{3,2} and Marc Schlegel ^{1,4}
¹Institute for Theoretical Physics, Universität Tülningen, Auf der Morgenstelle 14, D-720'6 Tülningen, Germany
²PRO, Universitär Paris-Sud, (CNSRVP2), F-91406, Ornay, France
Nikhef and Department of Physics and Astronomy, VU University Austerdam,
De Boelsdum (818, IA-108) HV, Austerdam, The Meherlands







• Gluon B-M can also be accessed via back-to-back $\psi/\Upsilon + \gamma$ associated production at the LHC. Also true at AFTER@LHC!



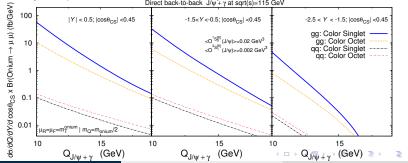


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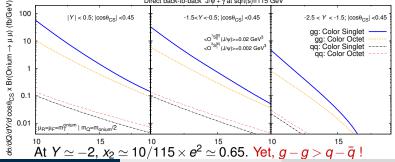
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SSA in Drell-Yan studies with AFTER@LHC

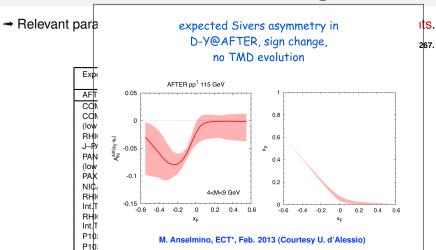
→ Relevant parameters for the future proposed polarized DY experiments.

S.J. Brodsky, F. Fleuret, C. Hadjidakis, JPL, Phys. Rep. 522 (2013) 239 V. Barone, F. Bradamante, A. Martin, Prog. Part. Nucl. Phys. 65 (2010) 267.

Experiment	particles	energy (GeV)	\sqrt{s} (GeV)	x_p^{\uparrow}	\mathcal{L} (nb ⁻¹ s ⁻¹)
AFTER	$p+p^{\uparrow}$	7000	115	$0.01 \div 0.9$	1
COMPASS	$\pi^{\pm} + p^{\uparrow}$	160	17.4	$0.2 \div 0.3$	2
COMPASS	$\pi^{\pm} + p^{\uparrow}$	160	17.4	~ 0.05	2
(low mass)					
RHIC	$p^{\uparrow} + p$	collider	500	$0.05 \div 0.1$	0.2
J-PARC	$p^{\uparrow} + p$	50	10	$0.5 \div 0.9$	1000
PANDA	$\bar{p} + p^{\uparrow}$	15	5.5	$0.2 \div 0.4$	0.2
(low mass)					
PAX	$p^{\uparrow} + \bar{p}$	collider	14	$0.1 \div 0.9$	0.002
NICA	$p^{\uparrow} + p$	collider	20	$0.1 \div 0.8$	0.001
RHIC	$p^{\uparrow} + p$	250	22	$0.2 \div 0.5$	2
Int.Target 1					
RHIC	$p^{\uparrow} + p$	250	22	$0.2 \div 0.5$	60
Int.Target 2					
P1027	$p^{\uparrow} + p$	120	15	$0.35 \div 0.85$	400-1000
P1039	$p + p^{\uparrow}$	120	15	$0.1 \div 0.3$	400-1000

- → For AFTER, the numbers correspond to a 50 cm polarized *H* target.
- \rightarrow $\ell^+\ell^-$ angular distribution: separation Sivers vs. Boer-Mulders effects

SSA in Drell-Yan studies with AFTER@LHC



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Instantaneous Luminosity:



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$$\mathcal{L} = \Phi_{beam} \times N_{target} = N_{beam} \times (\rho \times \ell \times \mathcal{N}_A)/A$$

 $\Phi_{beam} = 2 \times 10^5 \text{ Pb s}^{-1}, \quad \ell = 1 \text{ cm (target thickness)}$

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Target	ρ (g.cm ⁻³)	A	$\mathcal{L} (mb^{-1}.s^{-1}) = \int \mathcal{L} (nb^{-1}.yr^{-1})$
1m Liq. H ₂	0.07	1	800
1m Liq. D ₂	0.16	2	1000
1cm Be	1.85	9	25
1cm Cu	8.96	64	17
1cm W	19.1	185	13
1cm Pb	11.35	207	7



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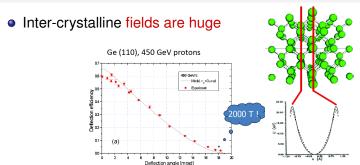
- Planned lumi for PHENIX Run15AuAu 2.8 nb⁻¹ (0.13 nb⁻¹ at 62 GeV)
- Nominal LHC lumi for PbPb 0.5 nb⁻¹



Inter-crystalline fields are huge

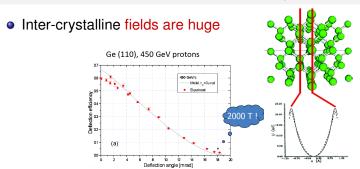
Ge (110), 450 GeV protons





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Ge (110), 450 GeV protons

Ge (100), 450 GeV protons

- The channeling efficiency is high for a deflection of a few mrad
- One can extract a significant part of the beam loss $(10^9 p^+ s^{-1})$
- Simple and robust way to extract the most energetic beam ever:





The beam extraction: news

[S. Montesano, Physics at AFTER using LHC beams, ECT* Trento, Feb. 2013]

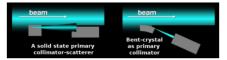
Goal : assess the possibility to use bent crystals as primary collimators in hadronic accelerators and colliders



UA9 installation in the SPS

Prototype crystal collimation system at SPS:

- local beam loss reduction (5÷20x reduction for proton beam)
- beam loss map show average loss reduction in the entire SPS ring
- halo extraction efficiency 70÷80% for protons (50÷70% for Pb)



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Towards an installation in the LHC: propose and install during LSI a min. number of devices

• 2 crystals

Long term plan is ambitious: propose a collimation system based on bent crystals for the upgrade of the current LHC collimation system

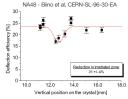
42 / 32

Simone Montesano - February 11th, 2013 - Physics at AFTER using the LHC beams

Crystal resistance to irradiation

- IHEP U-70 (Biryukov et al, NIMB 234, 23-30):
 - 70 GeV protons, 50 ms spills of 10¹⁴ protons every 9.6 s, several minutes irradiation
 - equivalent to 2 nominal LHC bunches for 500 turns every 10 s
 - 5 mm silicon crystal, channeling efficiency unchanged
- SPS North Area NA48 (Biino et al, CERN-SL-96-30-EA):
 - 450 GeV protons, 2.4 s spill of 5 x 10¹² protons every 14.4 s, one year irradiation, 2.4 x 10²⁰ protons/cm² in total,
 - · equivalent to several year of operation for a primary collimator in LHC
 - 10 x 50 x 0.9 mm³ silicon crystal, 0.8 x 0.3 mm² area irradiated, channeling efficiency reduced by 30%.
- HRMT16-UA9CRY (HiRadMat facility, November 2012):
 - 440 GeV protons, up to 288 bunches in 7.2 $\mu s,$ 1.1 x 10 11 protons per bunch (3 x 10 13 protons in total)
 - · energy deposition comparable to an asynchronous beam dump in LHC
 - 3 mm long silicon crystal, no damage to the crystal after accurate visual inspection, more tests planned to assess possible crystal lattice damage
 - · accurate FLUKA simulation of energy deposition and residual dose







S. Montesano (CERN - EN/STI) @ ECT* Trento workshop. Physics at AFTER using the LHC beams (Feb. 2013)

■ Beam loss: 10⁹ p⁺s⁻¹

• Extracted intensity: $5 \times 10^8 p^+ s^{-1}$ (1/2 the beam loss) E. Uggerhøj, U.I Uggerhøj, NIM B 234 (2005) 31



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 - one extracts $5.10^8/3.10^7 \simeq 15p^+$ from each bunch at each pass
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- This means $1.8 \times 10^{13}/3.2 \times 10^{14} \simeq 5.6\%$ of the p^+ in the beam These protons are lost anyway!

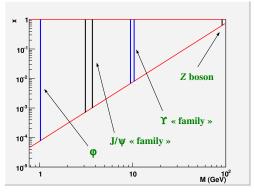
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- similar figures for the Pb-beam extraction



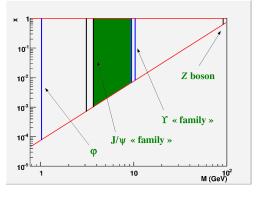
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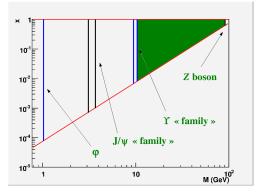
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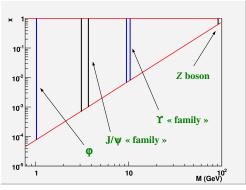
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- \rightarrow Above $c\bar{c}$: $\dot{x} \in [10^{-3}, 1]$
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Note: $x_{target} (\equiv x_2) > x_{projectile} (\equiv x_1)$ "backward" region

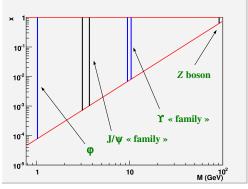


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- at large(est) x: backward ("easy")
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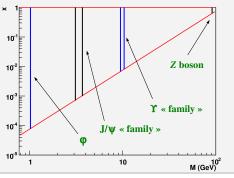
AFTER@LHC: A dilepton observatory?

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To do: to look at the rates to see how competitive this will be



Target	∫£ (fb ⁻¹ .yr ⁻¹)	$N(J/\Psi)$ yr ⁻¹ = ALBσ _Ψ	N(Υ) yr ⁻¹ =A <i>L</i> Bσ _Υ
1 m Liq. H ₂	20	4.0 10 ⁸	8.0 10 ⁵
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LHC pp 14 Tev (low pT)	0.05 (ALICE) 2 LHCb	3.6 10 ⁷ 1.4 10 ⁹	1.8 10 ⁵ 7.2 10 ⁶
RHIC pp 200GeV	1.2 10 ⁻²	4.8 10 ⁵	1.2 10 ³

Interpolating the world data set:

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- Probe of the (very) large x in the target



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 - but also pp collisions in gg-fusion process
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PHYSICAL REVIEW D

VOLUME 37, NUMBER 5

1 MARCH 1988

Structure-function analysis and ψ , jet, W, and Z production: Determining the gluon distribution

A. D. Martin

Department of Physics, University of Durham, Durham, England

R. G. Roberts Rutherford Appleton Laboratory, Didcot, Oxon, England

W. J. Stirling

Department of Physics, University of Durham, Durham, England (Received 27 July 1987)

We perform a next-to-leading-order structure-function analysis of deep-inelastic μN and νN scattering data and find acceptable fits for a range of input gluon distributions. We show three equally acceptable sets of parton distributions which correspond to gluon distributions which are (1) "soft," (2) "hard," and (3) which behave as $xG(x) \sim 1/\sqrt{x}$ at small x. J/ψ and prompt photon hadroproduction data are used to discriminate between the three sets. Set 1, with the "soft"-gluon distribution, is favored. W, Z, and jet production data from the CERN collider are well described but do not distinguish between the sets of structure functions. The precision of the predictions for σ_W and σ_{τ} allow the collider measurements to yield information on the number of light neutrinos and the mass of the top quark. Finally we discuss how the gluon distribution at very small x may be directly measured at DESY HERA.

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We perform a next-to-leading-order structure-function analysis of deep-inelastic μN and νN scattering data and find acceptable fits for a range of input gluon distributions. We show three equally acceptable sets of parton distributions which correspond to gluon distributions which are (1) "soft," (2) "hard," and (3) which behave as $xG(x) \sim 1/\sqrt{x}$ at small x. J/ψ and prompt photon hadroproduction data are used to discriminate between the three sets. Set 1, with the "soft"-gluon distribution, is favored. W. Z. and jet production data from the CERN collider are well described but do not distinguish between the sets of structure functions. The precision of the predictions for σ_W and σ_{τ} allow the collider measurements to yield information on the number of light neutrinos and the mass of the top quark. Finally we discuss how the gluon distribution at very small x may be directly measured at DESY HERA.

Production puzzle → quarkonium not used anymore in global fits



- Many hopes were put in quarkonium studies to extract gluon PDF
 - in photo/lepto production (DIS)
 - but also pp collisions in gg-fusion process
 - mainly because of the presence of a natural "hard" scale: m_O
 - and the good detectability of a dimuon pair

PHYSICAL REVIEW D

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Structure-function analysis and ψ , jet, W, and Z production: Determining the gluon distribution

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We perform a next-to-leading-order structure-function analysis of deep-inelastic μN and νN scattering data and find acceptable fits for a range of input gluon distributions. We show three equally acceptable sets of parton distributions which correspond to gluon distributions which are (1) "soft," $2^{\rm D}$ + and," and (3) which behave as $\kappa G(\nu) - 1/\nu X$ at small κ . J/ ν and prompt photon hadroproduction data are used to discriminate between the three sets. Set 1, with the "soft"-gluon distribution, is favored, νN , κ and growth production data from the CERN collider are well described but do not distinguish between the sets of structure functions. The precision of the predictions for $\nu \mu$ and $\nu \alpha$ gluon distribution is a soft only distribution and $\nu \alpha$ gluon that the mass of the top quark. Finally we discuss how the gluon distribution at very small $\nu \alpha$ may be directly measured at DESY HERA.

- Production puzzle → quarkonium not used anymore in global fits
- With systematic studies, one would restore its status as gluon probe

Target	A	∫£ (fb-¹.yr-¹)	N(J/Ψ) yr ⁻¹ = A£βσ _Ψ	N(Υ) yr ⁻¹ =A <i>£</i> βσ _Υ
1cm Be	9	0.62	1.1 10 ⁸	2.2 10 ⁵
1cm Cu	64	0.42	5.3 10 ⁸	1.1 10 ⁶
1cm W	185	0.31	1.1 10°	2.3 10 ⁶
1cm Pb	207	0.16	6.7 10 ⁸	1.3 10 ⁶
LHC pPb 8.8 TeV	207	10-4	1.0 107	7.5 10 ⁴
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 - not to mention ratio with open charm, Drell-Yan, etc ...



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- One should be careful with factorization breaking effects:
 - This calls for multiple measurements to (in)validate factorisation

Luminosities and yields with the extracted 2.76 TeV Pb beam

 $(\sqrt{s_{NN}} = 72 \text{ GeV})$

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The same picture also holds for open heavy flavour



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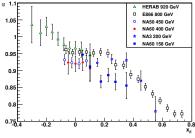
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 - χ_c never studied in AA collisions
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- the possibilities for cc recombination
 - Open charm studies are difficult where recombination matters most
 i.e. at low P_T
 - Only indirect indications –from the y and P_T dependence of R_{AA} that recombination may be at work
 - CNM effects may show a non-trivial y and P_T dependence ...



SPS and Hera-B

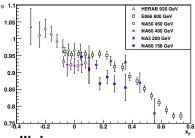
$-J/\psi$ data in pA collisions



NA60 Phys.Lett. B 706 (2012) 263 NA 50 Eur.Phys.J. C48 (2006) 329 NA 3 Z.Phys. C20 (1983) HERA-B Eur.Phys.J. C60 (2009) 525

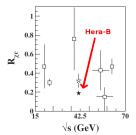
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HERA-B PRD 79 (2009) 012001, and ref. therein



LHB

Our idea is not completely new

Nuclear Instruments and Methods in Physics Research A 333 (1993) 125-135 North-Holland NUCLEAR
INSTRUMENTS
& METHODS
IN PHYSICS
RESEARCH
Section A

LHB, a fixed target experiment at LHC to measure CP violation in B mesons

Flavio Costantini

University of Pisa and INFN, Italy

A fixed target experiment at LHC to measure CP violation in B mesons is presented. A description of the proposed apparatus is given together with its sensitivity on the CP violation asymmetry measurement for the two benchmark decay channels $B^0 \to J/\psi + K_s^0$, $B^0 \to \pi^+\pi^-$. The possibility of obtaining an extracted LHC beam hinges on channeling in a bent silicon crystal. Recent results on beam extraction efficiencies measured at CERN SPS based on this technique are presented.

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This paper presents a fixed target experiment to measure CP violation in the B system based on the possibility of extracting the 8 TeV LHC proton beam using a bent silicon crystal [4]. A 10% extraction efficiency of the LHC beam halo will give an extracted beam intensity of about 10^8 protons/s allowing the production of as many as 10^{10} BB pairs per year, i.e. about two orders of magnitude more than what could be produced by an e⁺e⁻ asymmetric B factory with 10^{34} cm⁻²s⁻¹ luminosity [5].



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- After a year, one simply moves the crystal by less than one mm ...



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C.H. Chang, J.X. Wang, X.G. Wu. Comput. Phys. Commun. 177 (2007) 467



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• they should also be calculated for $x_F \rightarrow -1$

where IQ could dominate

Isolated-γ in p(7 TeV)-p(rest): √s ~ 115 GeV

■ p-p photon kinematics at fixed-target LHC (central rapidities): To access x > 0.3 one needs isolated- γ at: p_T = x_T√s/2 > 20 GeV/c

JETPHOX NLO (preliminary) pQCD calculations: p-p at √s=115 GeV |y| < 0.5, p₋>20 GeV/c Isolation: R=0.4, E_Thad<5 GeV 10⁻³ ~1 count 10-4 \mathcal{L} (10 cm H₂-target) ~ 2 • 10³ pb⁻¹/year p_ (GeV/c) PDF: CT10 52 eigenval. (90% CL) Scales: $\mu_i = p_{\tau}$ FF = BFG-II x-section uncertainties(*) of ±150% (*) (68%CL)/(90% CL) ~ 1.65 p_ (GeV/c)