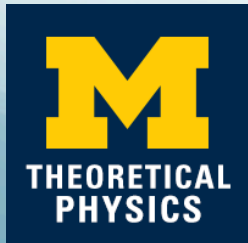


# Scattering Amplitudes

Henriette Elvang  
University of Michigan – Ann Arbor

**SUSY 2015**  
Lake Tahoe  
August 25, 2015



# Gluon scattering amplitudes

Using Feynman diagrams, the tree-level gluon amplitude requires

$$g + g \rightarrow g + g \quad 4 \text{ diagrams}$$

$$g + g \rightarrow g + g + g \quad 25 \text{ diagrams}$$

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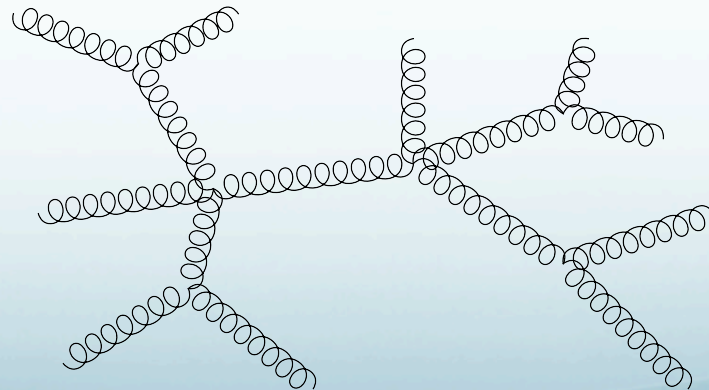
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Each diagram gets increasingly complicated as  $n$  grows



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However, the result for the amplitudes can be written much more compactly, e.g.

$$A_n[1^+ \dots i^- \dots j^- \dots n^+] = \frac{\langle ij \rangle^4}{\langle 12 \rangle \langle 23 \rangle \dots \langle n1 \rangle}.$$

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However, the result for the amplitudes can be written much more compactly, e.g.

**This is a  
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simplification!**



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# Scattering Amplitudes

Why are the results so simple?

Is there a better way to calculate amplitudes?

That is what this talk is about!

# What makes gauge theory amplitudes difficult?

Feynman rules (and hence the diagrams) depend on

- Choice of gauge
- Field redefinitions

because the Lagrangian is off-shell.

But the *on-shell amplitudes are independent* of both.

*If we can work only with on-shell invariant input, then these complications are avoided.*

There is a way to do this: *on-shell recursion relations*

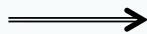
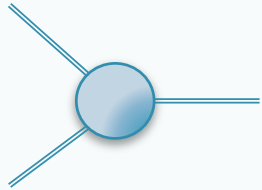
[Britto, Cachazo, Feng, Witten (2005)]

# Better ways to calculate

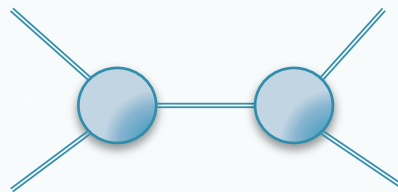
## On-shell recursion relations

*Idea:* build up n-point amplitudes from lower-point *on-shell* amplitudes

3-particle amplitude



4-particle amplitude



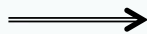
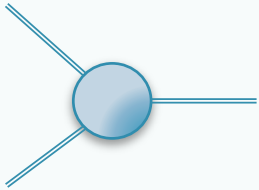


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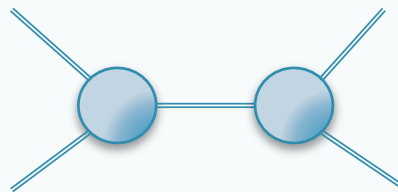
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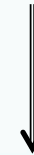
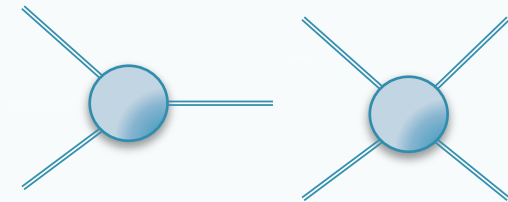
3-particle amplitude



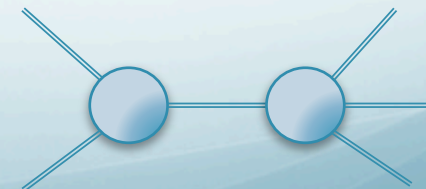
4-particle amplitude



3- and 4-particle amplitudes



5-particle amplitude

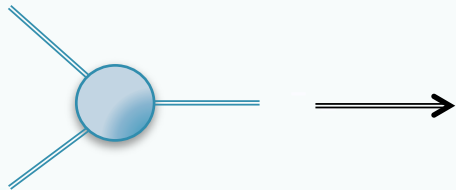


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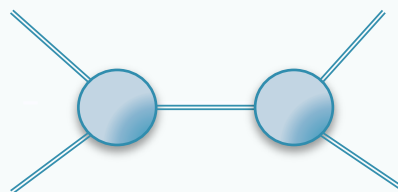
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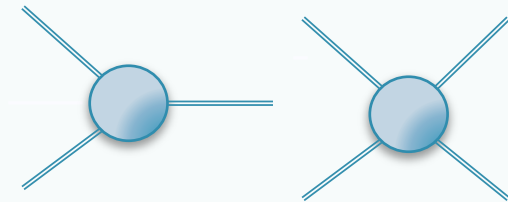
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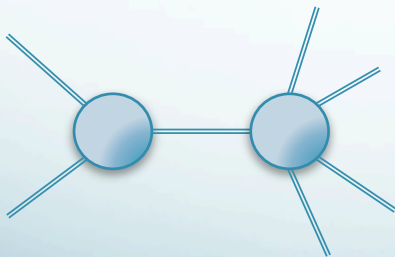
4-particle amplitude



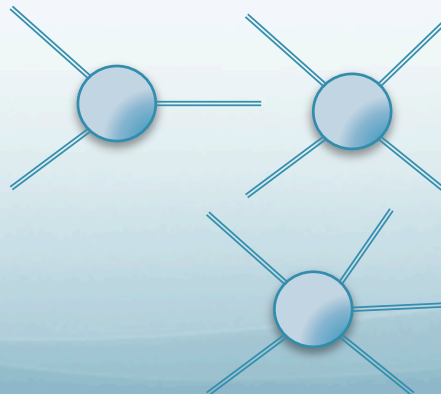
3- and 4-particle amplitudes



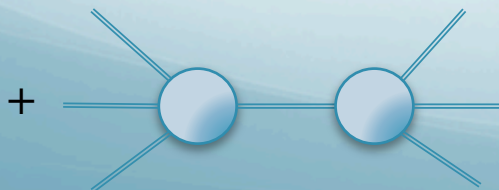
6-particle amplitude



3-, 4- and 5-particle amplitudes



5-particle amplitude



# Better ways to calculate

## On-shell recursion relations

*Idea:* build up n-point amplitudes from lower-point *on-shell* amplitudes

Mathematically

$$A_n = \sum_{\text{diagrams } I} \hat{A}_L(z_I) \frac{1}{P_I^2} \hat{A}_R(z_I)$$

“hat” indicates momentum shift

lower-point

The *derivation* exploits complex analysis and Cauchy’s theorem in a very simple way: uses knowledge of physical poles in amplitudes and factorization.

# Example 1 gluon + gluon $\rightarrow$ squark + squark

$$A_4 = \text{[Diagram 1]} + \text{[Diagram 2]} + \text{[Diagram 3]} + \text{[Diagram 4]}$$

As if four diagrams weren't bad enough.... having to square this  $|A_4|^2$  and sum over colors & helicities to get the cross-section just makes it worse.

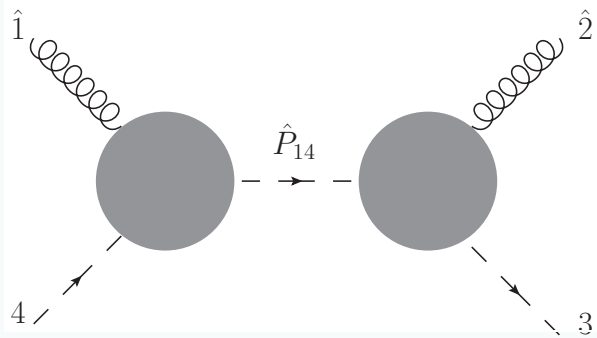
Recursion relations makes this much simpler!

I'll illustrate it first in the limit of  $m_{\text{squark}} = 0$

# Example 1 gluon + gluon $\rightarrow$ squark + squark

$$m_{\text{squark}} = 0$$

Recursion relations: only one diagram



$$A_4[g_1^\pm g_2^\pm \tilde{q}_3 \tilde{q}_4^*] = 0$$

$$A_4[g_1^- g_2^+ \tilde{q}_3 \tilde{q}_4^*] = \frac{[23]^2 [24]^2}{[12][23][34][41]}$$

$$A_4[g_1^+ g_2^- \tilde{q}_3 \tilde{q}_4^*] = ([ij] \leftrightarrow \langle ij \rangle)$$

$$\langle ij \rangle = \sqrt{2p_i \cdot p_j} e^{i\text{phase}}$$

$$[ij] = \sqrt{2p_i \cdot p_j} e^{-i\text{phase}}$$

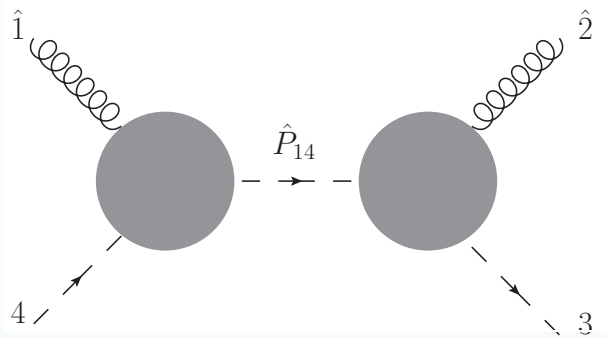
$$\langle 12 \rangle [12] = (p_1 + p_2)^2 = -s, \quad \text{etc}$$

$$\left| A_4[g_1^- g_2^+ \tilde{q}_3 \tilde{q}_4^*] \right|^2 = \frac{s_{23}^2 s_{24}^2}{s_{12} s_{23} s_{34} s_{41}} = \frac{t^2}{s^2}$$

# Example 1 gluon + gluon $\rightarrow$ squark + squark

Recursion relations: only one diagram

$m_{\text{squark}}$  nonzero



$$A_4[g_1^- g_2^- \tilde{q}_3 \tilde{q}_4^*] = \frac{1}{[12]^2} \frac{sm_{\tilde{q}}^2}{u_1}$$

$$A_4[g_1^- g_2^+ \tilde{q}_3 \tilde{q}_4^*] = -\frac{1}{u_1 s} \langle 1|4|2][2|3|1\rangle,$$

$$u_1 = u - m_{\tilde{q}}^2$$

$$t_1 = t - m_{\tilde{q}}^2$$



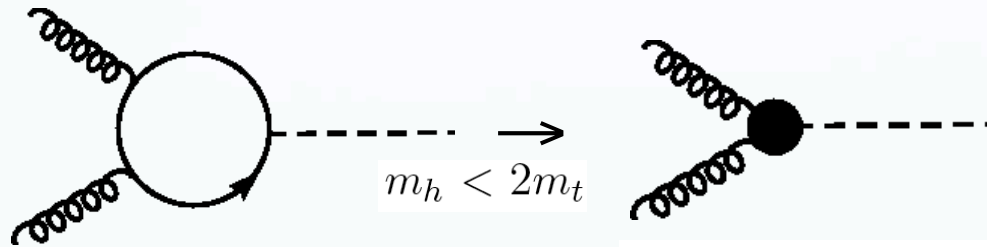
$$\sum |A_4(g_1 g_2 \tilde{q}_3 \tilde{q}_4^*)|^2 = 2g^4 \left\{ N(N^2 - 1) \left( 1 - 2\frac{u_1 t_1}{s^2} \right) - \frac{N^2 - 1}{N} \right\} \left[ 1 - 2\frac{m_{\tilde{q}}^2 s}{u_1 t_1} \left( 1 - \frac{m_{\tilde{q}}^2 s}{u_1 t_1} \right) \right]$$

This is the well-known result from the gluon to squark cross-section in the literature. Here derived via recursion relations by a Michigan undergrad, Filipe Rudriguez.



## Example 2: Gluon fusion

Gluon-Higgs fusion:



$g g \rightarrow \text{Higgs}$

[Dixon, Glover, Khoze' 04;  
Berger, Del Duca, Dixon' 06]

Applied  
on-shell recursion  
to gluon-Higgs fusion

$\Rightarrow$  Higgs+ $n$  gluon  
amplitudes

$h \text{Tr} F_{\mu\nu} F^{\mu\nu}$   
dim-5 operator

1-loop  
 $\Rightarrow$  tree-level

2-loop  
 $\Rightarrow$  1-loop

etc

Proof of validity  
[Cohen, H.E., Kiermaier (2010)]

# Study of Scattering Amplitudes

## Two pillars

*Practical application*  
to phenomenologically  
relevant processes

Uncover the  
*mathematical structure*  
of the amplitudes

TRUTH



BEAUTY



# Three Major Research Directions

**1) Try to push loop calc as far as possible in a very controlled simple theory: “(planar) N=4 Super Yang Mills Theory” (SYM) (gluons, gluinos, scalars - all massless)**

*Goals: ‘Solve’ theory at all loop order.  
Compact expressions? Understand why!*

**2) Adapt lessons from N=4 SYM to phenomenologically relevant theories to find new computational methods.**

*Goals: application to analysis of data from  
LHC and future particle experiments. New physics insights?*

**3) Use new methods to explore perturbative quantum gravity.**

*Goals: Point-particle quantum gravity perturbatively sensible?  
Gravity as (gauge theory)<sup>2</sup>  
Structure of string theory amplitudes*

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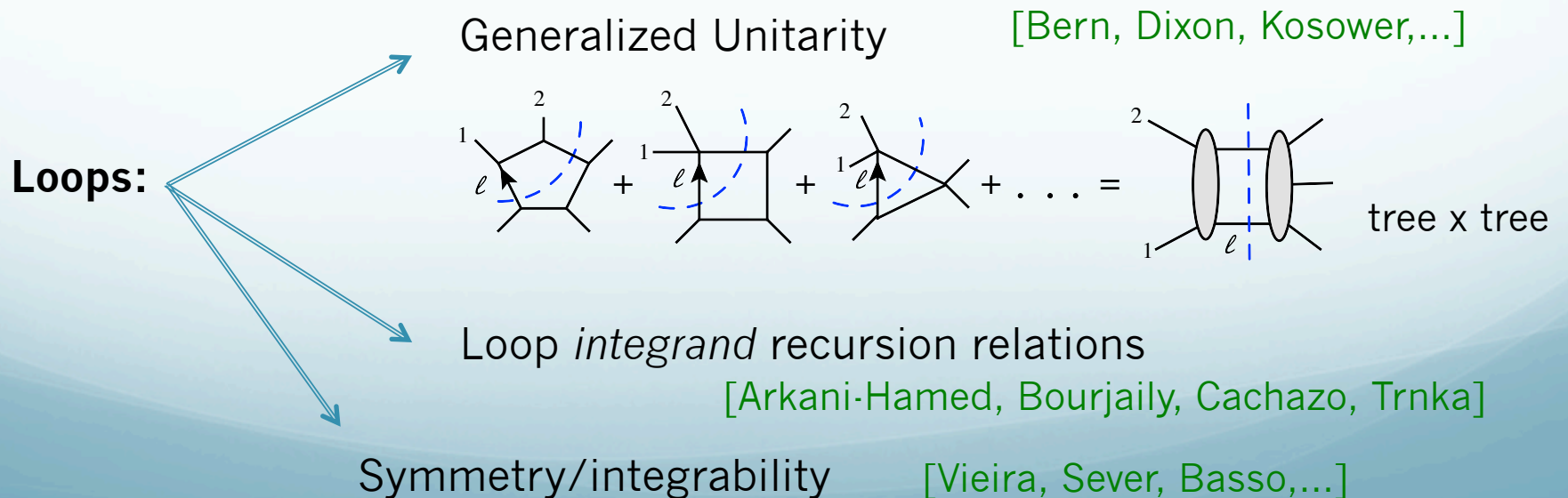
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**Tree-level:** all tree amplitudes in N=4 SYM solved via recursion.



# Planar N=4 Super Yang Mills Theory and GEOMETRY

It turns out that amplitudes have a geometric interpretation.

Recursion relations  $\Rightarrow$  amplitude = sum of terms

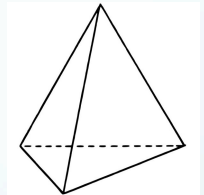
Each of these terms can be interpreted as the volume of a 4-simplex

0-simplex = a point

1-simplex = line segment

2-simplex = triangle

3-simplex = tetrahedron.



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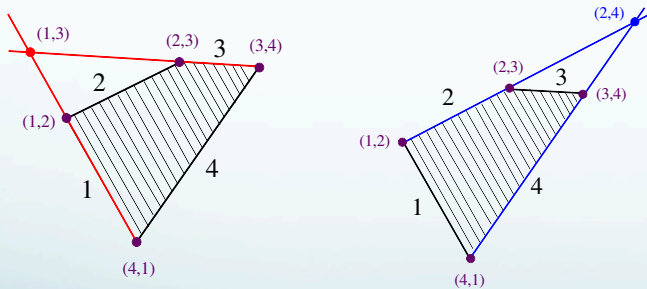
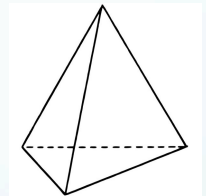
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In this way, we can interpret the amplitude as the volume of polytope in a higher-dimensional space!!

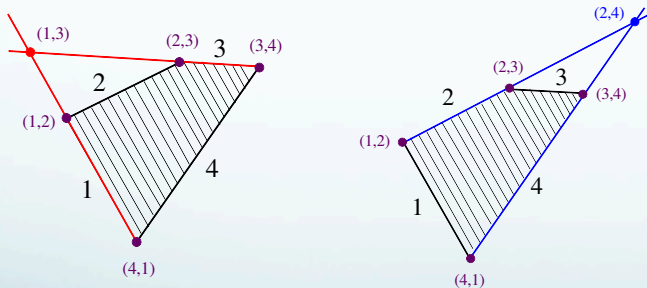
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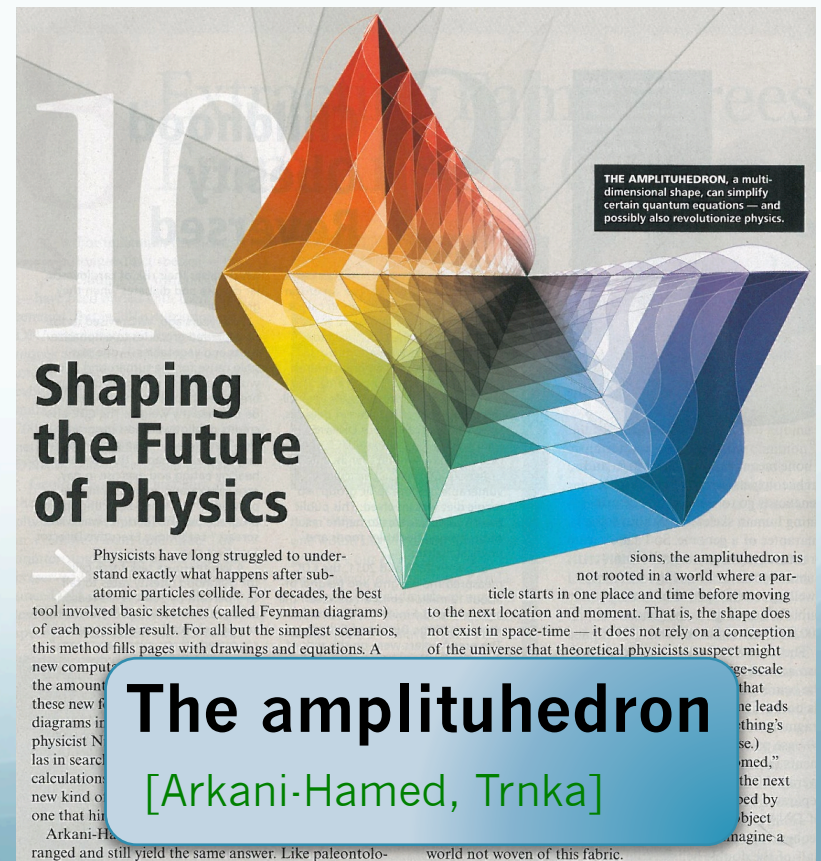
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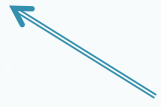
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So not a good theory of quantum gravity.



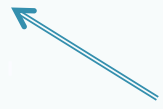
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Z. Bern<sup>a</sup>, L. J. Dixon<sup>b</sup>, R. Roiban<sup>c</sup> (2007)

$\mathcal{N}=8$  supergravity in 4d:

< 2007: 4-graviton 1,2,3-loop finite

2009: 4-graviton 4-loop finite

[Bern, Carrasco, Dixon, Johansson,  
Kosower, Roiban (2007)+(2009)]

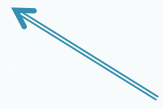


Done using *generalized unitarity* AND *gravity as (gauge theory)<sup>2</sup>*

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2010: all  $n$ -graviton amplitudes finite  $L < 7$   
guaranteed by symmetries.

Used amplitude techniques to assess possible UV counterterms. No loops needed.



[Freedman, Kiermaier, HE; Kiermaier, HE;  
Beisert, Kiermaier, Freedman, Morales, Stieberger (2010)]

Perturbative finiteness still an open question.

# Summary

Modern on-shell methods for scattering amplitudes are incredibly powerful.

Applications in a wide range of problems:

- Pheno amplitudes
- Formal developments (mathematical structure, geometry)
- Quantum gravity
- Studies of formal aspects of QFT
- Non-renormalization theorems

and much more!

# Textbook

Published by  
Cambridge University Press  
(2015)

(preview: [arXiv:1308.1697](https://arxiv.org/abs/1308.1697))

*Includes introductions to:*

Spinor helicity formalism  
for Feynman diagrams.

On-shell recursion relations.

Supersymmetry applications.

Unitarity cuts and loops.

Advanced subjects  
(Grassmannians, on-shell diagrams, polytopes, supergravity,  
**supergravity** = (SYM)<sup>2</sup>...)

