

# Exact $SU(5)$ Yukawa matrix unification in the General Flavour Violating MSSM

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based on M. Kowalska, JHEP 1504 (2015) 120

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**INNOVATIVE ECONOMY**  
NATIONAL COHESION STRATEGY



**EUROPEAN UNION**  
EUROPEAN REGIONAL  
DEVELOPMENT FUND



1. SU(5) Yukawa matrix unification
2. Minimal Supersymmetric Standard Model
3. chirally-enhanced SUSY threshold corrections
4. off-diagonal soft terms help  $\rightarrow$  General Flavour Violating MSSM
5. Phenomenology of Yukawa unification in the GFV MSSM:
  - ▶ 2nd + 3rd generation
  - ▶ 1st + 2nd + 3rd generation

# Unification - SU(5) model: matter & Higgs sector

Georgi, Glashow, 1974

$$\underbrace{(\bar{3}, 1, \frac{1}{3})}_{d_R^*} \oplus \underbrace{(1, 2, -\frac{1}{2})}_l = \underbrace{\bar{5}}_{\Psi_{\bar{5}}}$$
$$\underbrace{(3, 2, \frac{1}{6})}_q \oplus \underbrace{(\bar{3}, 1, -\frac{2}{3})}_{u_R^*} \oplus \underbrace{(1, 1, 1)}_{e_R^*} = \underbrace{10}_{\Psi_{10}},$$

$$W \ni \Psi_{10} \mathbf{Y}^{de} \Psi_{\bar{5}} H_{\bar{5}} + \Psi_{10} \mathbf{Y}^u \Psi_{10} H_5$$

# Unification - SU(5) model: matter & Higgs sector

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$$W \ni \Psi_{10} \mathbf{Y}^{de} \Psi_{\bar{5}} H_{\bar{5}} + \Psi_{10} \mathbf{Y}^u \Psi_{10} H_5$$

$$Y_{ii}^{d, MSSM} = Y_{ii}^{e, MSSM}$$

# Gauge coupling unification

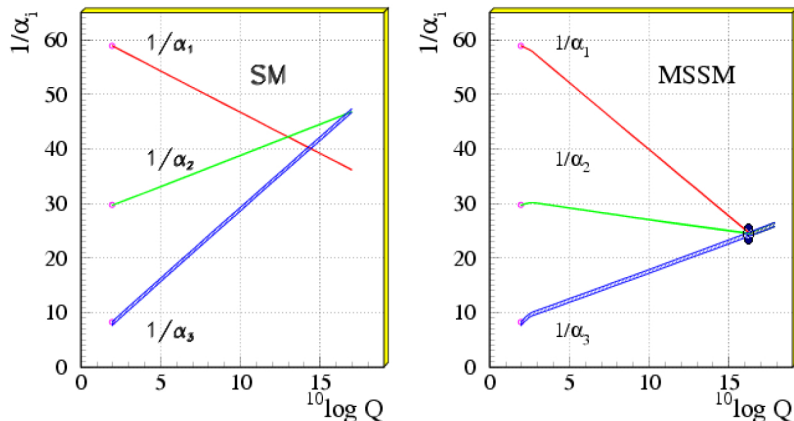
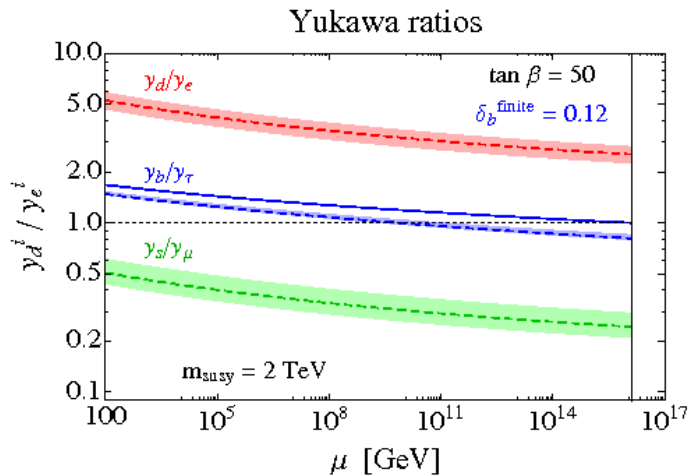


Figure : Gauge coupling unification in non-SUSY GUTs on the left vs. SUSY GUTs on the right using the LEP data (1991)  
arXiv: hep-ph/0012288

# Yukawa couplings at the GUT scale



Elor, Hall,  
Pinner,  
Ruderman,  
JHEP 1210  
(2012) 111,  
arXiv:1206.5301

2nd generation:  $Y_\mu(M_{\text{GUT}}) \approx 3Y_s(M_{\text{GUT}})$   
1st generation:  $Y_e(M_{\text{GUT}}) \approx 1/3Y_d(M_{\text{GUT}})$

Change the boundary condition at the high scale

- ▶ additional Higgs fields, e.g.

*H. Georgi and C. Jarlskog, Phys. Lett. B86 (1979) 297*

$$H_5, \quad H_{\bar{5}}, \quad (H_{45}) \rightarrow Y_\mu = 3Y_s, \quad Y_e = 1/3Y_d$$

- ▶ correction  $O(1)$  from higher-dim. operators

*D. Emmanuel-Costa and S. Wiesenfeldt, Nucl. Phys. B 661 (2003) 62*

*S. Antusch and M. Spinrath, Phys. Rev. D 79 (2009) 095004*

*S. Antusch, S.F. King and M. Spinrath, Phys. Rev. D 89 (2014) 055027*

$$W = W_Y + W_{HO} \rightarrow \begin{aligned} Y_d^{ij} &= Y_{de}^{ij} + \frac{1}{M_{\text{PL}}} F^{ij} \\ Y_e^{ij} &= Y_{de}^{ij} + \frac{1}{M_{\text{PL}}} G^{ij} \end{aligned}$$

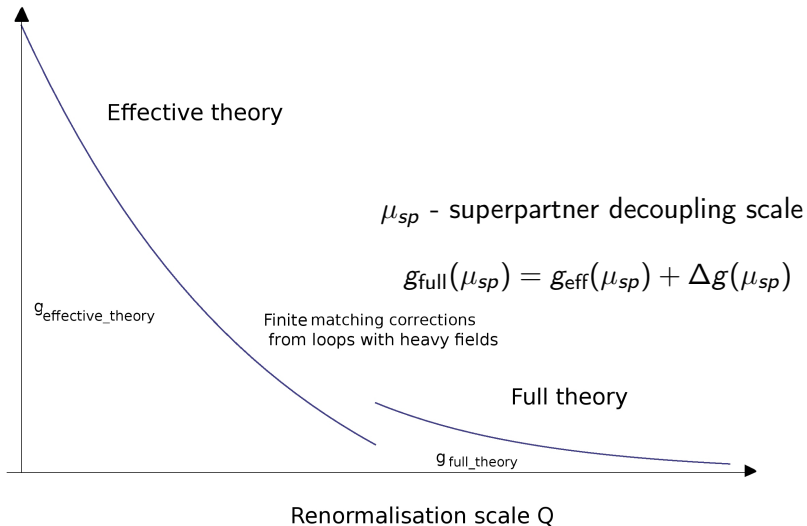
Manipulate the boundary condition between SM and MSSM - play with threshold corrections

- ▶ Diaz-Cruz, Murayama, Pierce, Phys.Rev.D 65:075011, 2002  
(particular ansatz using  $A$ -terms for unification)
- ▶ Ts. Enkhbat, arXiv:0909.5597  
(general diagonal  $A$ -terms)
- ▶ MI, Eur.Phys.J. C75 (2015) 51  
(update - new exp results, broader  $\tan \beta$  range, weaker impact on flavour observables)



# Threshold corrections

Renormalised constant  $g$



*A. Crivellin, L. Hofer, J. Rosiek, JHEP 1107 (2011) 017*

$$v_f Y_{ii}^{f MSSM} = v_f Y_{ii}^{f SM} - \Sigma_{ii}^f(Y_j^{f'}, \dots).$$

A. Crivellin, L. Hofer, J. Rosiek, *JHEP* 1107 (2011) 017

$$v_f Y_{ii}^{f MSSM} = v_f Y_{ii}^{f SM} - \sum_{ii}^f (Y_j^{f'}, \dots).$$

$$m_i^{d(\ell) SM} - v_d Y_{ii}^{d(\ell) MSSM} = \sum_{ii}^{d(\ell) LR} + \epsilon_i^{d(\ell)} v_u Y_{ii}^{d(\ell)(0)} + O\left(\frac{v^2}{M_{SUSY}}\right),$$

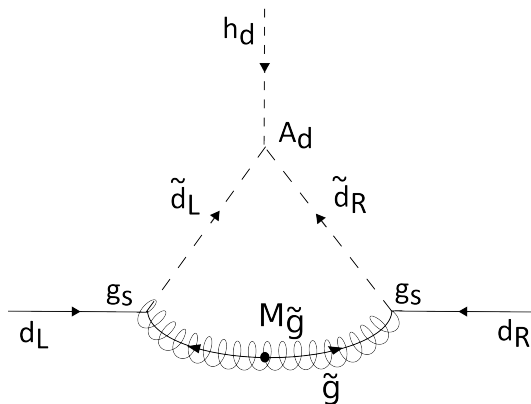
A. Crivellin, L. Hofer, J. Rosiek, *JHEP* 1107 (2011) 017

$$v_f Y_{ii}^{f MSSM} = v_f Y_{ii}^{f SM} - \sum_{ii}^f (Y_j^{f'}, \dots).$$

$$m_i^{d(\ell) SM} - v_d Y_{ii}^{d(\ell) MSSM} = \sum_{ii}^{d(\ell) LR} \mathcal{X} + \epsilon_i^{d(\ell)} v_u Y_{ii}^{d(\ell)(0)} + O\left(\frac{v^2}{M_{SUSY}}\right),$$

$$Y_{ii}^{d(\ell) MSSM} = \frac{m_i^{d(\ell) SM} - \sum_{ii}^{d(\ell) LR} \mathcal{X}}{v_d (1 + \tan \beta \cdot \epsilon_i^{d(\ell)})}.$$

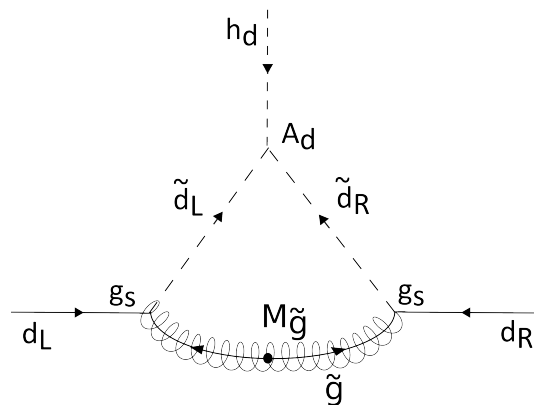
# Threshold corrections - example diagrams



- ▶ Diaz-Cruz, Murayama, Pierce, Phys.Rev.D 65:075011, 2002
- ▶ Ts. Enkhbat, arXiv:0909.5597
- ▶ Ml, Eur.Phys.J. C75 (2015) 51

$$(\Sigma_{ii}^d)^{\tilde{g}} \sim \alpha_S m_{\tilde{g}} (v_d A_{ii}^d - v_d Y_{ii}^d \mu \tan \beta)$$

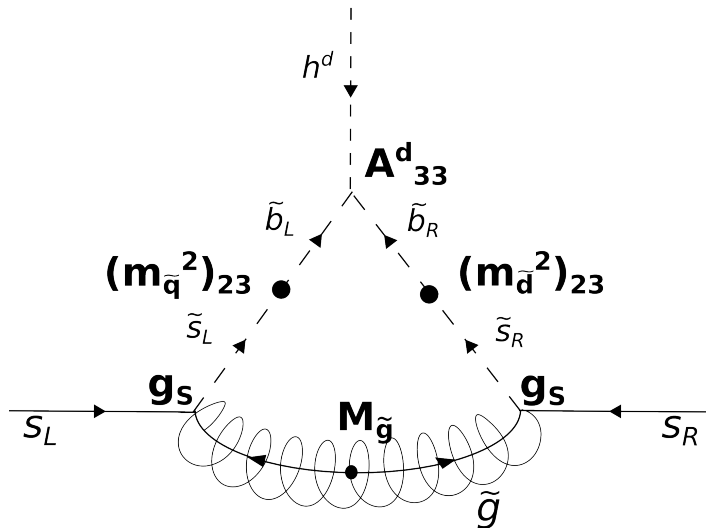
# Threshold corrections - example diagrams



- ▶ Diaz-Cruz, Murayama, Pierce, Phys.Rev.D 65:075011, 2002
- ▶ Ts. Enkhbat, arXiv:0909.5597
- ▶ Ml, Eur.Phys.J. C75 (2015) 51

$A_s \sim m_{\tilde{s}}$  required for strange-muon unification  
 $\Rightarrow$  MSSM vacuum metastable

# Threshold corrections - example diagrams



$$(\Sigma_{22}^d)^{\tilde{g}} \sim \alpha_S M_{\tilde{g}} v_d (A_{33}^d - Y_b \mu \tan \beta) (m_{\tilde{q}}^2)_{23} (m_{\tilde{d}}^2)_{23}$$

$$(m_{\tilde{l}}^2)_{ij} = (m_{\tilde{d}}^2)_{ij} \equiv (m_{\text{dl}}^2)_{ij}$$

$$(m_{\tilde{q}}^2)_{ij} = (m_{\tilde{u}}^2)_{ij} = (m_{\tilde{e}}^2)_{ij} \equiv (m_{\text{ue}}^2)_{ij}$$

$$A_{ij}^d = A_{ij}^e \equiv \mathbf{A}_{ij}^{\text{de}}$$

$$A_{ij}^u$$

$$M_1 = M_2 = M_3 \equiv \mathbf{M}_{1/2},$$

$$\tan \beta = \frac{v_u}{v_d}$$

$$m_{\text{H}_u}^2, \quad m_{\text{H}_d}^2$$



**BayesFITSv.3.2** A. Fowlie, M. Kazana, K. Kowalska, S. Munir, L. Roszkowski, E. M. Sessolo, S. Trojanowski, Y. L. S. Tsai [arXiv:1206.0264], K. Kowalska [arXiv:1406.0710]

## MultiNest v2.7

F. Feroz, M. P. Hobson and M. Bridges, [arXiv:0809.3437]

### SUSY\_Flavor v2.10

A. Crivellin, J. Rosiek,  
P. H. Chankowski, A. Dedes,  
S. Jaeger and P. Tanedo  
[arXiv:1203.5023]

### SPheno v3.3.3

W. Porod and F. Staub [arXiv:1104.1573]

### HIGGSBounds v4.0.0 HIGGSSIGNALS v1.0.0

P. Bechtle et al. [arXiv:0811.4169]  
[arXiv:1102.1898], [arXiv:1311.0055],  
[arXiv:1305.1933]

### DarkSUSY v5.0.6

P. Gondolo, J. Edsjo,  
P. Ullio, L. Bergstrom,  
M. Schelke, E. A. Baltz,  
[astro-ph/0406204]

# Ranges of input parameters

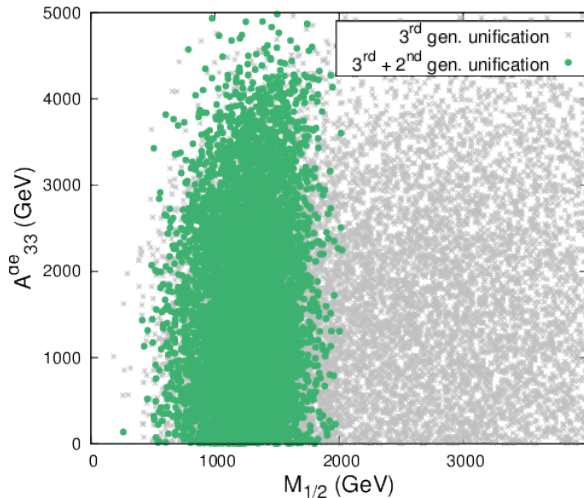
$m_{ii}^{dl}, i = 1, 2, 3$	[100, 7000] GeV		
$m_{23}^{dl}/m_{33}^{dl}$	[0, 1]	$M_{1/2}$	[100, 4000] GeV
$m_{13}^{dl}/m_{33}^{dl}$	[0, 1]	$m_{H_u}$	[100, 8000] GeV
$m_{12}^{dl}/m_{33}^{dl}$	[0, 1]	$m_{H_d}$	[100, 8000] GeV
$m_{ii}^{ue}, i = 1, 2, 3$	[100, 7000] GeV	$\tan \beta$	[3, 45]

$A_{33}^{de}$	[0, 5000] GeV
$A_{33}^u$	[-9000, 9000] GeV
$A_{11}^{de}/A_{33}^{de}$	[-0.00028, 0.00028]
$A_{22}^{de}/A_{33}^{de}$	[-0.065, 0.065]
$A_{22}^u/A_{33}^u$	[-0.005, 0.005]
$A_{ij}^{de}/A_{33}^{de}, i \neq j$	[-0.5, 0.5]

$$m_{ij}^{dl} \equiv \sqrt{(m_{dl}^2)_{ij}}, \quad m_{ij}^{ue} \equiv \sqrt{(m_{ue}^2)_{ij}}.$$

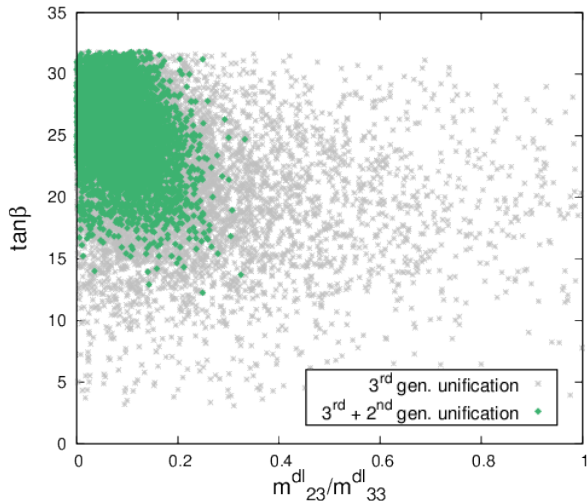
# 3rd + 2nd family Yukawa unification

relevant GFV parameter:  $m_{23}^{dl}$



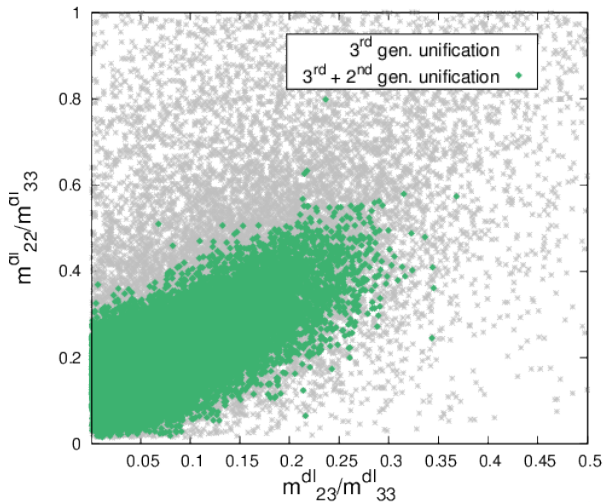
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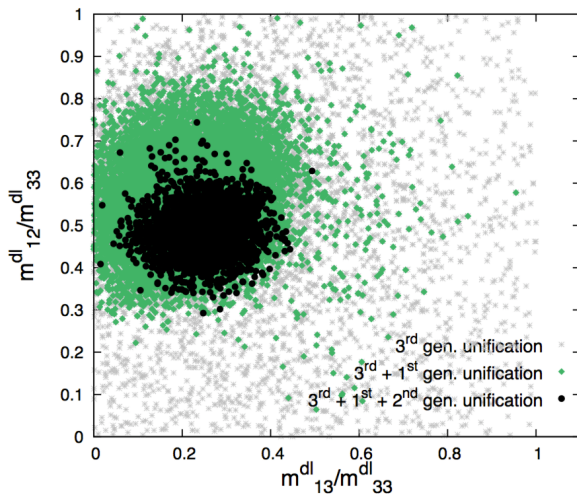
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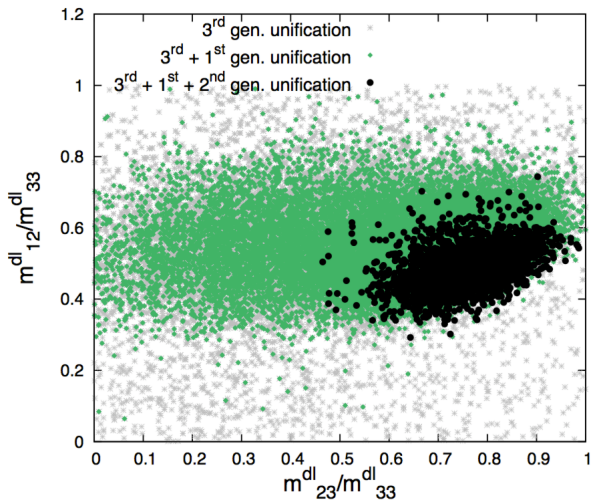
# 3rd + 2nd + 1st family Yukawa unification

relevant GFV parameters:  $m_{23}^{dl}$ ,  $m_{13}^{dl}$ ,  $m_{12}^{dl}$ ,  $A_{12}^{de}$



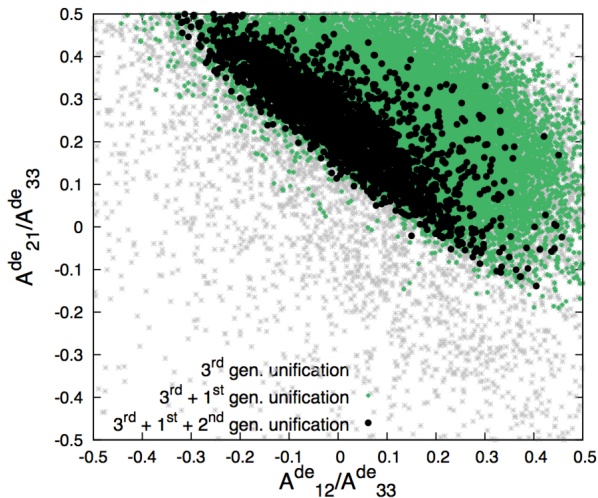
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relevant GFV parameters:  $m_{23}^{dl}$ ,  $m_{13}^{dl}$ ,  $m_{12}^{dl}$ ,  $A_{12}^{de}$



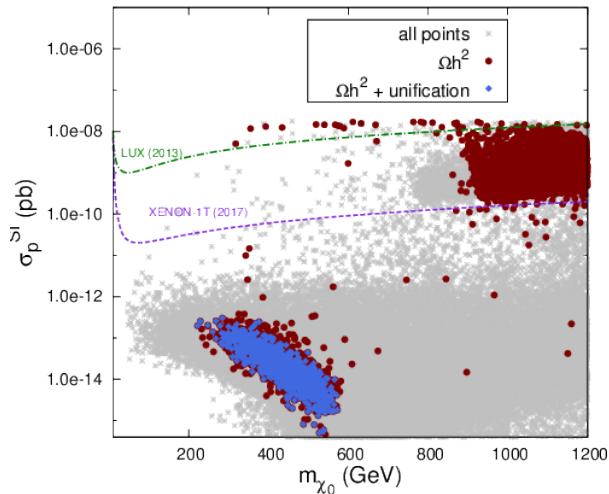


# Experimental constraints

Measurement	Mean or range	Error [ exp., th.]
$\Omega_\chi h^2$	0.1199	[0.0027, 10%]
$m_h$ (by CMS)	125.7 GeV	[0.4, 3.0] GeV
$\sin^2 \theta_{\text{eff}}$	0.23155	[0.00012, 0.00015]
$M_W$	80.385 GeV	[0.015, 0.015] GeV
$\text{BR}(\bar{B} \rightarrow X_s \gamma) \times 10^4$	3.43	[0.22, 0.23]
$\text{BR}(B_s \rightarrow \mu^+ \mu^-) \times 10^9$	2.8	[0.7, 0.23]
$\text{BR}(B_d \rightarrow \mu^+ \mu^-) \times 10^{10}$	3.9	[1.6, 0.2]
$\Delta M_{B_s} \times 10^{11}$	1.1691 GeV	[0.0014, 0.1580] GeV
$\Delta M_{B_d} \times 10^{13}$	3.357 GeV	[0.033, 0.340] GeV
$\Delta M_{B_d} / \Delta M_{B_s} \times 10^2$	2.87	[0.02, 0.14]
$\sin(2\beta)_{\text{exp}}$	0.682	[0.019, 0.003]
$\text{BR}(B_u \rightarrow \tau \nu) \times 10^4$	1.14	[0.27, 0.07]
$\text{BR}(K^+ \rightarrow \pi^+ \nu \bar{\nu}) \times 10^{10}$	1.73	[1.15, 0.04]
$ d_n  \times 10^{26}$	$< 2.9$ e cm	[0, 30%]
$\epsilon_K \times 10^3$	2.228	[0.011, 0.17]

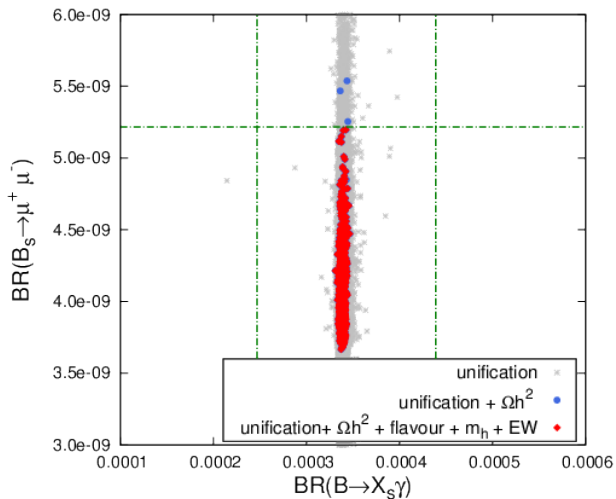
$$\begin{array}{l|l} \text{BR}(\mu^+ \rightarrow e^+ \gamma) \times 10^{13} & < 5.7 \\ \text{BR}(\tau^\pm \rightarrow e^\pm \gamma) \times 10^8 & < 3.3 \\ \text{BR}(\tau^\pm \rightarrow \mu^\pm \gamma) \times 10^8 & < 4.4 \\ \text{BR}(\mu^+ \rightarrow e^+ e^+ e^-) \times 10^{12} & < 1.0 \\ \text{BR}(\tau^\pm \rightarrow e^\pm e^+ e^-) \times 10^8 & < 2.7 \\ \text{BR}(\tau^\pm \rightarrow \mu^\pm \mu^+ \mu^-) \times 10^8 & < 2.1 \end{array}$$

# 3rd + 2nd family unification: Dark matter



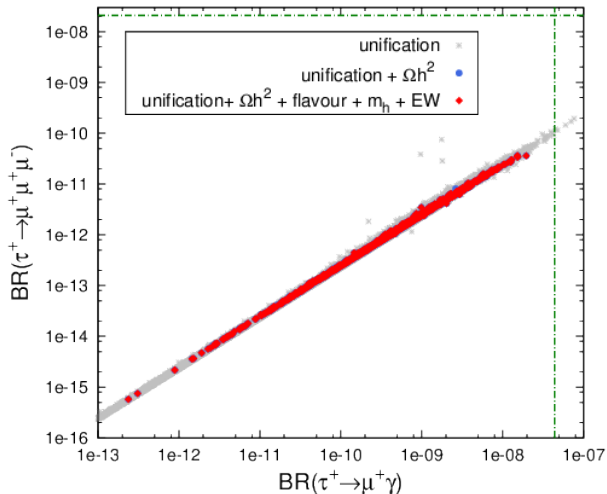
only bino DM

# 3rd + 2nd family unification: Flavour observables



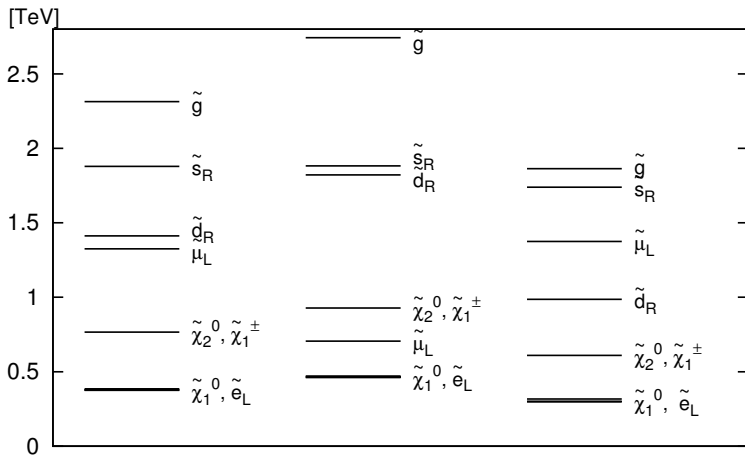
dashed lines -  $3\sigma$  experimental limits

# 3rd + 2nd family unification: Flavour observables

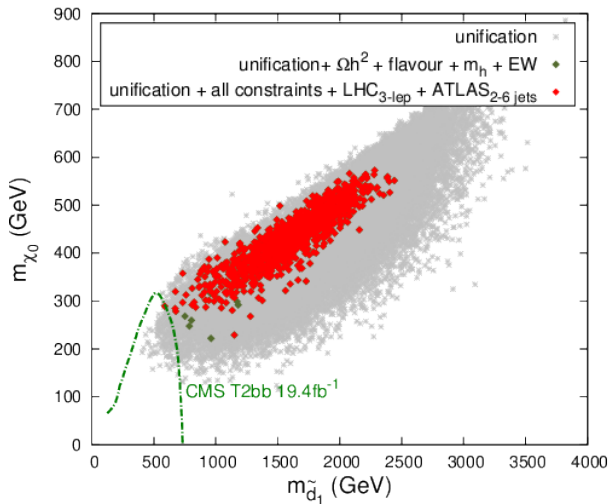


dashed lines -  $3\sigma$  experimental limits

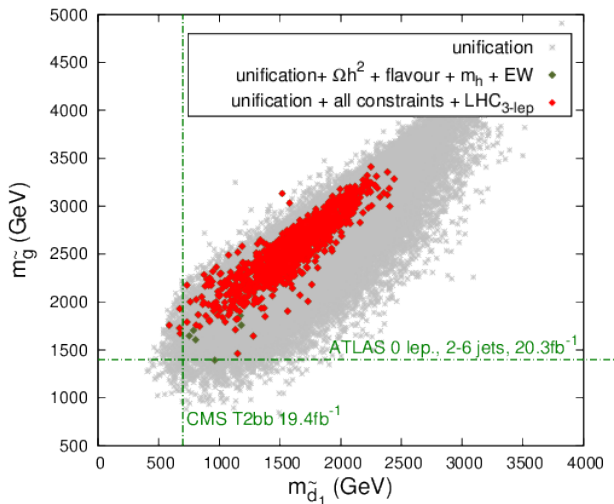
# 3rd + 2nd family unification: typical spectra



# 3rd + 2nd family unification: LHC SUSY searches



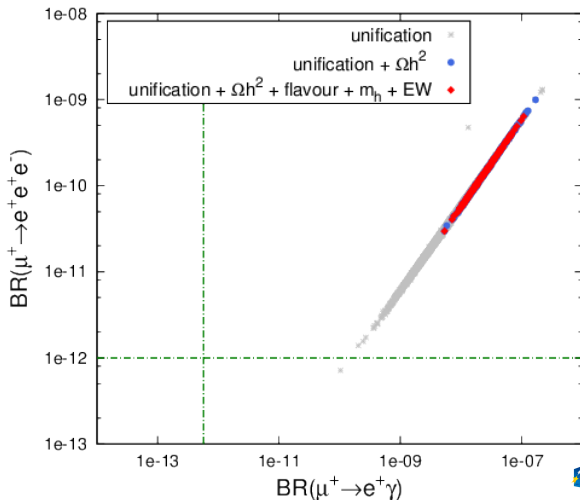
# 3rd + 2nd family unification: LHC SUSY searches





# 3rd + 2nd + 1st family unification: LFV

- ▶ consistent with quark flavour observables
- ▶ **strongly disfavoured** by the Lepton Flavour Violating observables



- ▶ Are there other regions consistent with Yukawa unification?
- ▶ Could the exclusion of  $GFV_{123}$  Yukawa unification be avoided?  
e.g. much higher SUSY masses,  
an  $SU(5)$  GUT scenario with  $m_{\tilde{t}_1} \neq m_{\tilde{d}}$
- ▶ Could two-loop threshold corrections be any relevant?
- ▶  $Y_d = Y_e$  in a  $GFV_{23}$ -like scenario without vacuum metastability?

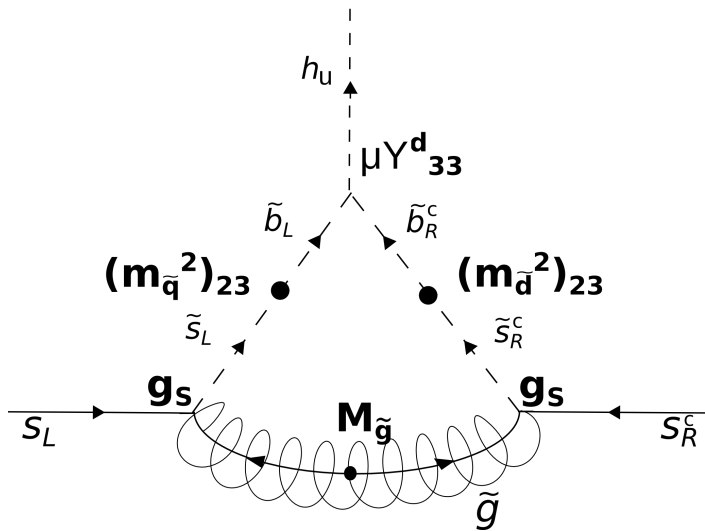
Non-trivial flavour structure of the MSSM

can facilitate the SU(5) Yukawa matrix unification

- ▶ Unification of the 2nd and 3rd generation phenomenologically allowed (relevant parameter:  $(m_{dl}^2)_{23}$ )
- ▶ Full unification of all three generations is strongly disfavoured by the limits on LFV (problems with:  $(m_{dl}^2)_{12}$ ,  $A_{12/21}^{de}$ )

# Supplementary slides

# Threshold corrections - example diagrams



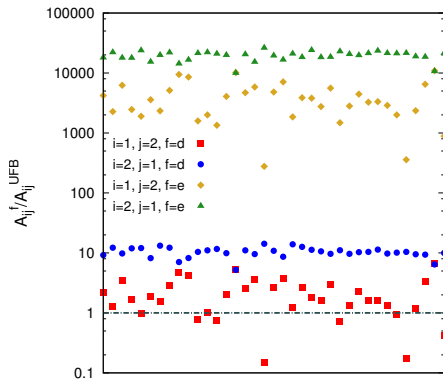
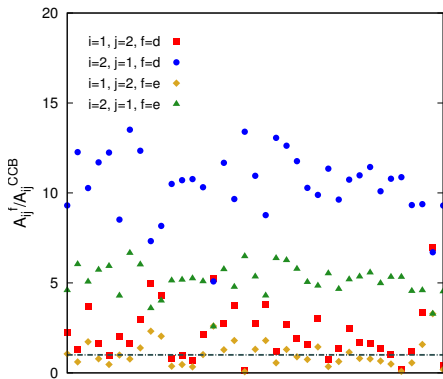
In the down-squark sector, Tree-level formulae for the CCB and UFB bounds in the down-squark sector:

$$(v_d/\sqrt{2})A_{ij}^d \leq m_k^d [(m_{\tilde{q}}^2)_{ii} + (m_{\tilde{d}}^2)_{jj} + m_{H_d}^2 + \mu^2]^{1/2}, \quad k = \text{Max}(i, j)$$

$$(v_d/\sqrt{2})A_{ij}^d \leq m_k^d [(m_{\tilde{q}}^2)_{ii} + (m_{\tilde{d}}^2)_{jj} + (m_{\tilde{l}}^2)_{ii} + (m_{\tilde{e}}^2)_{jj}]^{1/2}$$

J. A. Casas and S. Dimopoulos, [hep-ph/9606237]

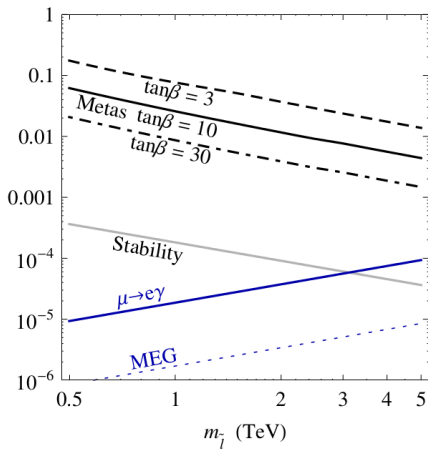
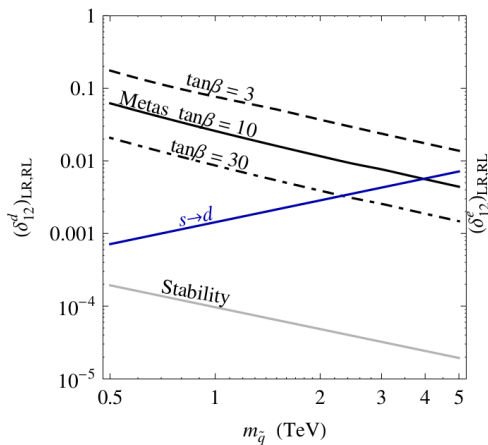
# EW vacuum stability



EW vacuum CCB (a) and UFB (b) upper bounds (dashed) on the elements  $A_{12/21}^{d,e}$

# EW vacuum stability

J. h. Park, [arXiv:1011.4939]:



metastability bounds are 2-3 orders of magnitude weaker.



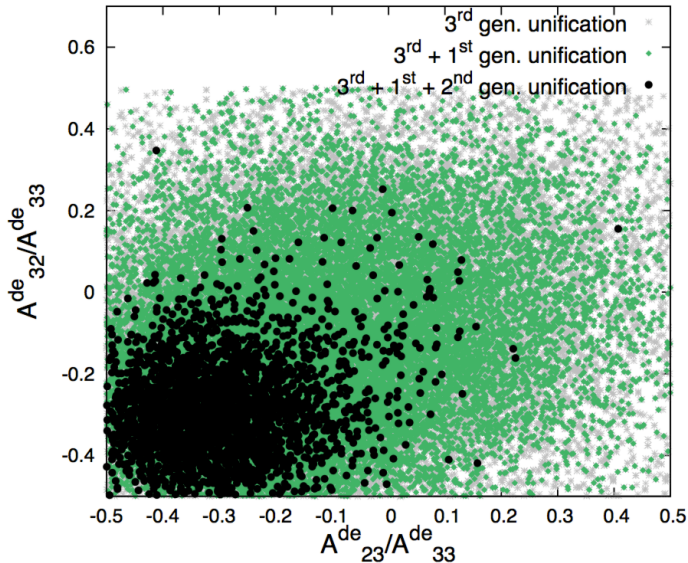
# Constants values

we scanned over ( $m_t^{\text{pole}}$ ,  $m_b^{\overline{MS}}(m_b)$ ,  $\alpha_{\text{em}}^{-1}(M_Z)$  and  $\alpha_s^{\overline{MS}}(M_Z)$ ) ( $\bar{\rho}$ ,  $\bar{\eta}$ ,  $A$ ,  $\lambda$ )

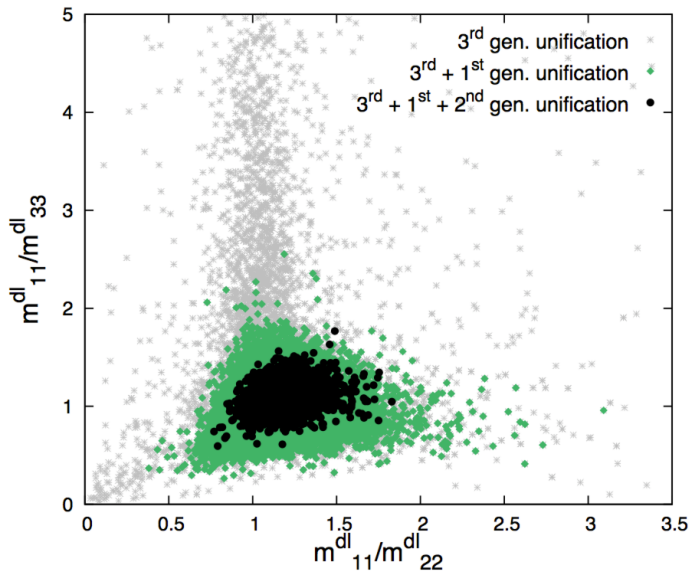
$m_t^{\text{pole}}$		$m_b^{\overline{MS}}(m_b)$		$\alpha_s^{\overline{MS}}(M_Z)$		$\alpha_{\text{em}}^{-1}(M_Z)$	
173.34 ± 0.76 GeV		4.18 ± 0.03 GeV		0.1184 ± 0.0007		127.944 ± 0.015	
$m_u^{\overline{MS}}$	$m_d^{\overline{MS}}$	$m_s^{\overline{MS}}$	$m_c^{\overline{MS}}(m_c)$	$m_e^{\text{pole}}$	$m_\mu^{\text{pole}}$	$m_\tau^{\text{pole}}$	$M_Z^{\text{pole}}$
2.3 MeV	4.8 MeV	95 MeV	1.275 GeV	511 keV	106 MeV	1.777 GeV	91.19
$\bar{\rho}$		$\bar{\eta}$		$A$		$\lambda$	
0.159 ± 0.045		0.363 ± 0.049		0.802 ± 0.020		0.22535 ± 0.0006	

**Table :** Standard Model parameters (PDG 2014) used in our numerical calculations. The light ( $u$ ,  $d$ ,  $s$ ) quark masses are  $\overline{MS}$ -renormalized at 2 GeV.

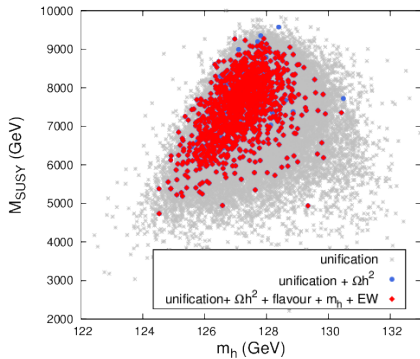
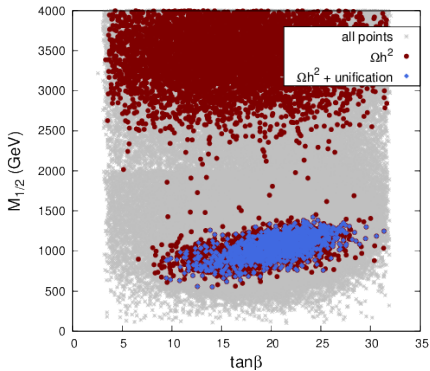
# Yukawa unification



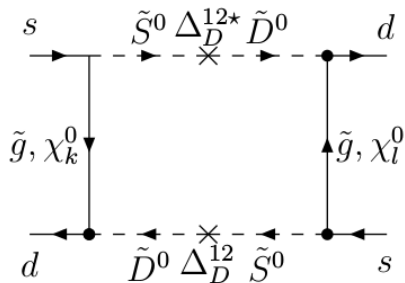
# Yukawa unification



# Dark matter & Higgs mass



# Kaon and B mixing



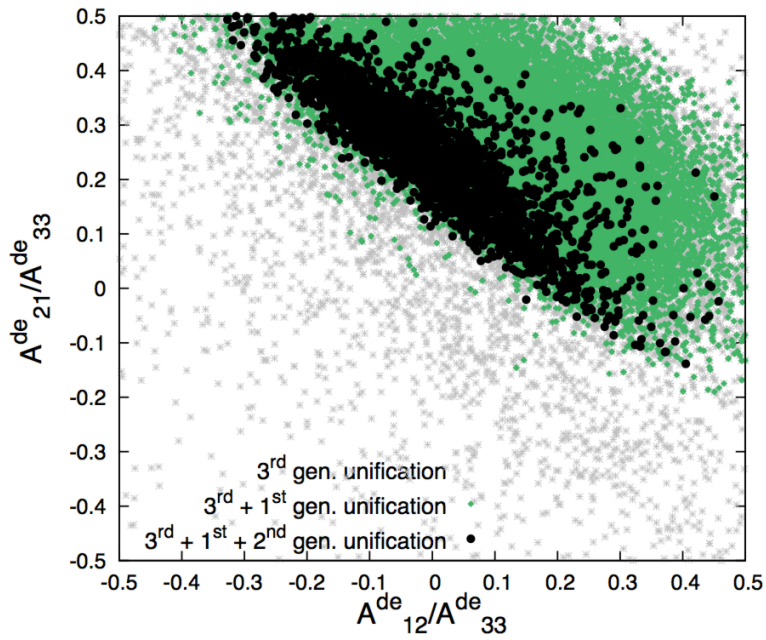
$$\Delta M_{B_{d(s)}} = 2 \left| \langle \bar{B}_{d(s)}^0 | H_{\text{eff}}^{\Delta B=2} | B_{d(s)}^0 \rangle \right|$$

$$\varepsilon_K = \frac{\exp(i\pi/4)}{\sqrt{2}\Delta M_K} \Im \langle \bar{K}^0 | H_{\text{eff}}^{\Delta S=2} | K^0 \rangle$$

$\Delta_D^{12} = m_{12}^d$  in super-CKM basis

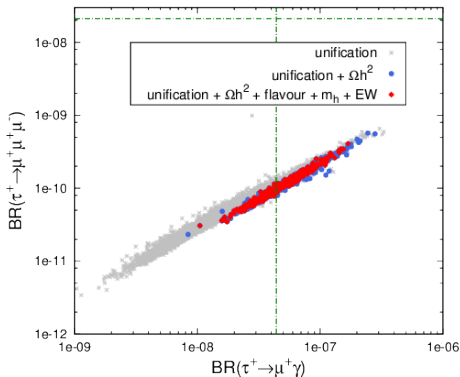
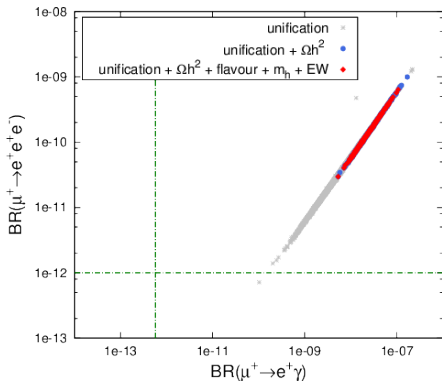
Misiak, Pokorski, Rosiek,

hep-ph/9703442



# 3rd + 2nd + 1st family unification: LFV

- ▶ consistent with quark flavour observables
- ▶ **strongly disfavoured** by the Lepton Flavour Violating observables



Parameter	Scanning Range	Parameter	Scanning Range
$M_{1/2}$	[100, 4000] GeV	$m_{ii}^{dl}, i = 1, 2, 3$	[100, 7000] GeV
$m_{H_u}$	[100, 8000] GeV	$m_{23}^{dl}/m_{33}^{dl}$	[0, 1]
$m_{H_d}$	[100, 8000] GeV	$m_{13}^{dl}/m_{33}^{dl}$	[0, 1]
$\tan \beta$	[3, 45]	$m_{12}^{dl}/m_{33}^{dl}$	[0, 1]
$\text{sgn } \mu$	-1	$m_{ii}^{ue}, i = 1, 2, 3$	[100, 7000] GeV
$A_{33}^{de}$	[0, 5000] GeV	<p>Table : Ranges of the input SUSY parameters used in our initial scan. Several omitted soft SUSY-breaking parameters at the GUT scale (namely <math>A_{11}^u</math> as well as <math>A_{ij}^u</math> and <math>m_{ij}^{ue}</math> for <math>i \neq j</math>) have been set to zero.</p>	
$A_{33}^u$	[-9000, 9000] GeV		
$A_{11}^{de}/A_{33}^{de}$	[-0.00028, 0.00028]		
$A_{22}^{de}/A_{33}^{de}$	[-0.065, 0.065]		
$A_{22}^u/A_{33}^u$	[-0.005, 0.005]		
$A_{ij}^{de}/A_{33}^{de}, i \neq j$	[-0.5, 0.5]		



# Minimal Supersymmetric Standard Model

Superfields	Fermions	Scalars
$Q = \begin{pmatrix} U_L \\ D_L \end{pmatrix}$	$q = \begin{pmatrix} u_L \\ d_L \end{pmatrix}$	$\tilde{q} = \begin{pmatrix} \tilde{u}_L \\ \tilde{d}_L \end{pmatrix}$
$U_R$	$u_R$	$\tilde{u}_R$
$D_R$	$d_R$	$\tilde{d}_R$
$L = \begin{pmatrix} N \\ E_L \end{pmatrix}$	$l = \begin{pmatrix} \nu \\ e_L \end{pmatrix}$	$\tilde{l} = \begin{pmatrix} \tilde{\nu} \\ \tilde{e}_L \end{pmatrix}$
$E_R$	$e_R$	$\tilde{e}_R$
$H_d$	$\tilde{h}_d = \begin{pmatrix} \tilde{h}_d^0 \\ \tilde{h}_d^- \end{pmatrix}$	$h_d = \begin{pmatrix} h_d^0 \\ h_d^- \end{pmatrix}$
$H_u$	$\tilde{h}_u = \begin{pmatrix} \tilde{h}_u^+ \\ \tilde{h}_u^0 \end{pmatrix}$	$h_u = \begin{pmatrix} h_u^+ \\ h_u^0 \end{pmatrix}$



Yukawa interactions in the superpotential of the minimal  $SU(5)$  SUSY GUT:

$$\mathcal{W} \ni \psi_{10} \mathbf{Y}^{\text{de}} \psi_5 \bar{H}_5 + \psi_{10} \mathbf{Y}^{\text{u}} \psi_{10} H_5 \quad (0.1)$$

Here  $\bar{H}_5$  and  $H_5$  are two Higgs superfields that couple to model's matter fields. The masses of known fermions are thus given by only two independent  $3 \times 3$  matrices  $\mathbf{Y}_{\text{de}}$  and  $\mathbf{Y}_{\text{u}}$

the superpotential of MSSM:

$$\mathcal{W}_{MSSM} = QY^u U_R H_u + QY^d D_R H_d + LY^e E_R H_d + \mu H_d H_u.$$

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$$\mathcal{W}_{MSSM} = QY^u U_R H_u + QY^d D_R H_d + LY^e E_R H_d + \mu H_d H_u.$$

the soft supersymmetry-breaking terms:

$$\begin{aligned} \mathcal{L}_{soft}^{MSSM} = & -\frac{1}{2}[m_{\tilde{g}}(\tilde{G}^a)^T C \tilde{G}^a + m_{\tilde{W}}(\tilde{W}^I)^T C \tilde{W}^I + m_{\tilde{B}}\tilde{B}^T C \tilde{B} + h.c.] - m_{h_d}^2 h_d^\dagger \\ & - \tilde{q}^\dagger(\mathbf{m}_{\tilde{q}}^2)\tilde{q} - (\tilde{u}_R)^\dagger(\mathbf{m}_{\tilde{u}}^2)(\tilde{u}_R) - (\tilde{d}_R)^\dagger(\mathbf{m}_{\tilde{d}}^2)(\tilde{d}_R) - \tilde{l}^\dagger(\mathbf{m}_{\tilde{l}}^2)\tilde{l} - (\tilde{e}_R)^\dagger(\mathbf{m}_{\tilde{e}}^2) \\ & + \tilde{q}\mathbf{A}^u \tilde{u}_R h_u + \tilde{q}\mathbf{A}^d \tilde{d}_R h_d + \tilde{l}\mathbf{A}^e \tilde{e}_R h_d + B\mu h_d h_u + h.c. \end{aligned}$$

# Problem's anatomy in SU(5)

In SM and MSSM the fermion masses are independent parameters and are given by 3 Yukawa matrices:

$$Y^u \rightarrow m_u, m_c, m_t$$

$$Y^d \rightarrow m_d, m_s, m_b$$

$$Y^e \rightarrow m_e, m_\mu, m_\tau$$

In the minimal SU(5) Grand Unified Theory the symmetry requires:

$$Y_d = Y_e, Y_s = Y_\mu, Y_b = Y_\tau$$

flavour mixing (CKM matrix can be included in)  $\mathbf{Y}_u$