# **RPV from Discrete R Symmetries**

(and Recent Nucleon Decay Searches in Super-K)

# Volodymyr Takhistov (UCI)



SUSY 2015

Lake Tahoe – August 25, 2015

*Primarily based on:* - M.-C. Chen, M. Ratz and V. Takhistov [NPB 819 (2015) 322; hep-ph/1410.3474]

#### **Overview**

- Problems of the MSSM
- "Fixing" the MSSM → **discrete symmetries**
- Properties of discrete symmetries
- Equivalent symmetries
- Algorithm for finding maximal discrete symmetry
- Survey of discrete symmetries for R-parity violating and conserving MSSM
- What else can we do with these symmetries ... ?

+ (if time permits) a word about recent nucleon decay searches at Super-Kamiokande

#### **Problems of the MSSM**

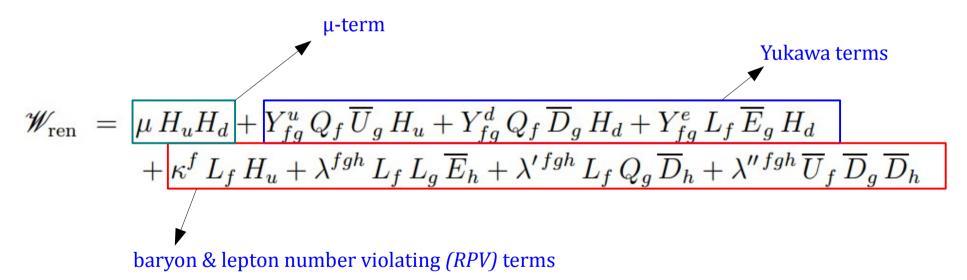
- SUSY well motivated
  - $\rightarrow$  top down / "formally" (Haag's theorem, local SUSY  $\rightarrow$  gravity, strings)
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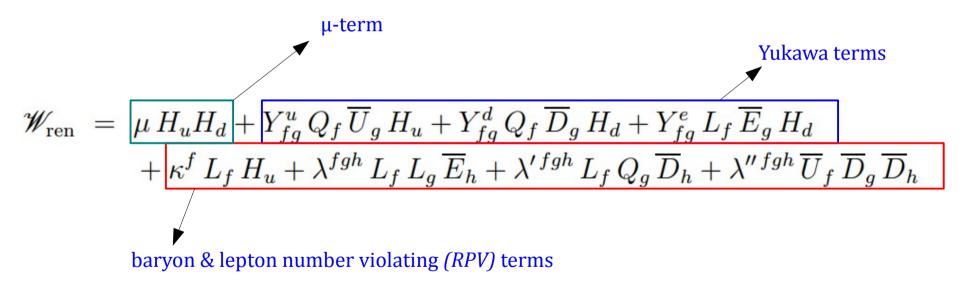


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- Problems of the MSSM:
  - $\rightarrow$  rapid proton decay (*RPV terms*) and  $\mu$  problem

• Phenomenological constraints:

μ-term (EWSB):  $\mu \sim 10^{-16} M_{Pl}$ 

Neutrino mass:  $\kappa^i \leq 10^{-21} M_{Pl}$ 

**Proton stability** 

$$\lambda' \lambda'' \leq 10^{-27}$$
$$\lambda' \lambda_3 \leq 10^{-10}$$
$$\lambda_1 \leq 10^{-8}, \lambda_2 \leq 10^{-8}$$

#### + higher order RPV terms

 $\begin{array}{rcl}
\mathcal{O}_{1} &=& \left[Q \, Q \, Q \, L\right]_{F} , & \mathcal{O}_{2} &=& \left[\overline{U} \, \overline{U} \, \overline{D} \, \overline{E}\right]_{F} , \\
\mathcal{O}_{3} &=& \left[Q \, Q \, Q \, H_{d}\right]_{F} , & \mathcal{O}_{4} &=& \left[Q \, \overline{U} \, \overline{E} \, H_{d}\right]_{F} , \\
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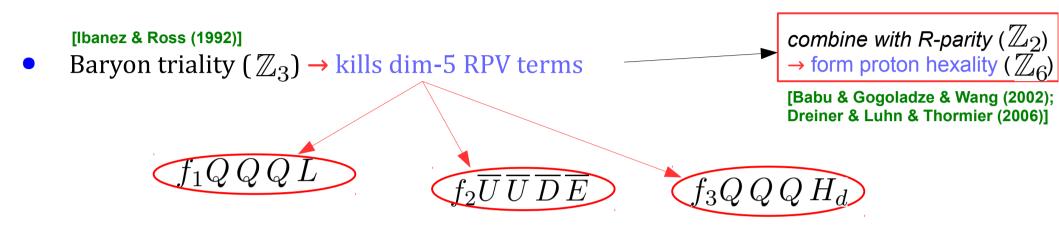
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• Typically impose **R-parity** ("matter parity") [Farrar & Bayet (1978); Dimopoulos, Raby & Wilczek (1981)]  $\rightarrow$  discrete  $\mathbb{Z}_2$  symmetry, kills  $L H_u$  term and dim-4 RPV terms

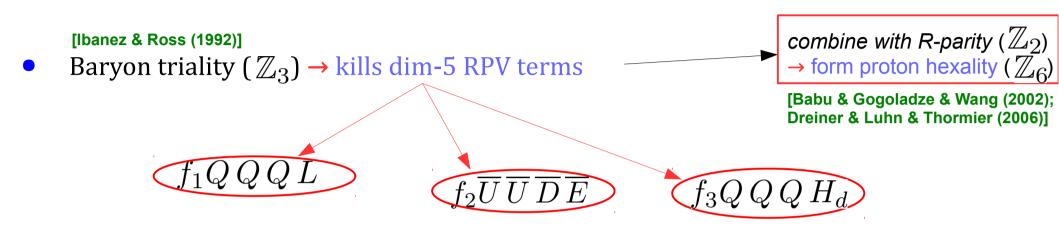
$$\mathcal{W}_{\text{ren}} = \mu H_u H_d + Y_{fg}^u Q_f \overline{U}_g H_u + Y_{fg}^d Q_f \overline{D}_g H_d + Y_{fg}^e L_f \overline{E}_g H_d$$
$$+ \kappa^f L_f H_u + \lambda^{fgh} L_f L_g \overline{E}_h + \lambda'^{fgh} L_f Q_g \overline{D}_h + \lambda''^{fgh} \overline{U}_f \overline{D}_g \overline{D}_h$$

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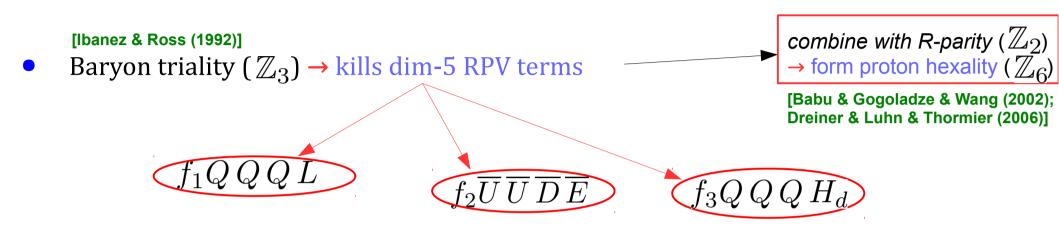
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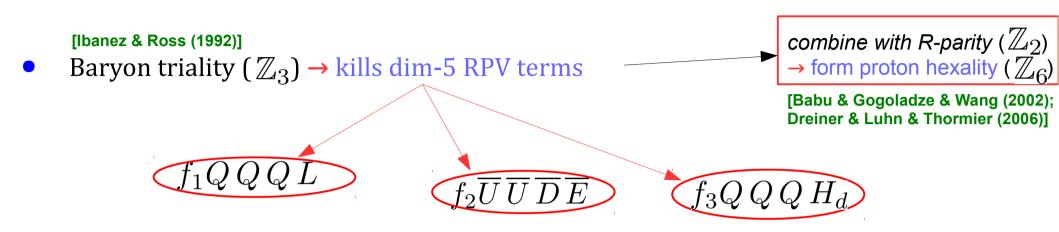
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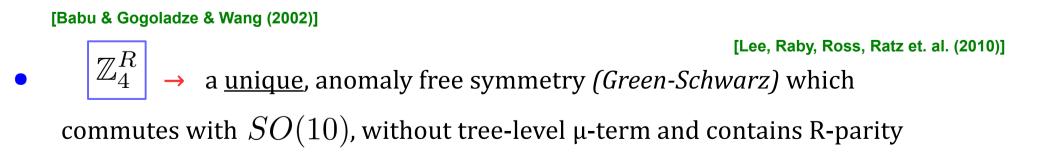


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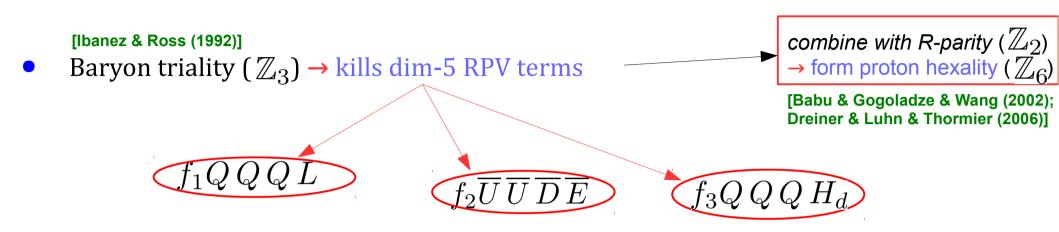


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# [Babu & Gogoladze & Wang (2002)] $\mathbb{Z}_{4}^{R} \rightarrow a \text{ unique, anomaly free symmetry (Green-Schwarz) which}$ commutes with SO(10), without tree-level $\mu$ -term and contains R-parity (also found w/o GS but with adding extra fields [Kurosawa, Maru, Yanagida (2001)])

#### 08/25/15

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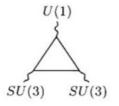
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• **Anomalies** (failure of maintaining the symmetry at quantum level)

OR



non-invariant path integral measure (Fujikawa method)



anom. coeff. for this diagram is 0 in SM

QM measure vs. classical measure  

$$\mathcal{D}\psi'\mathcal{D}\overline{\psi'} = (\det J)^{-1} (\det \overline{J})^{-1} \mathcal{D}\psi\mathcal{D}\overline{\psi}.$$
  
 $\equiv \exp\left(i\int d^4x \alpha(x)\mathcal{A}(x)\right)$   
 $\mathcal{A}(x) = \frac{1}{16\pi^2} \epsilon_{\mu\nu\rho\sigma} F_a^{\mu\nu}(x) F_b^{\rho\sigma}(x) \operatorname{tr}[tT^aT^b]$ 

anomaly coefficient

 $\rightarrow$  usually set anomaly coefficient to 0 (SM)

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$$\mathcal{L} \supset -\frac{a}{8}F^a \tilde{F^a}$$

can absorb path integral shift from anomaly

$$\Delta \mathcal{L}_{\text{anomaly}} = \sum_{G} \frac{\alpha}{32\pi^2} F^a \tilde{F^a} A_{G-G-\mathbb{Z}_N}$$

into Lagrangian shift after axion transforms under the symmetry  $a 
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• <u>Anomaly universality</u>: GS can cancel anomalies separately, but coupling unification and/or GUT requires universality  $A_{G_i-G_i-\mathbb{Z}_N^R} = \rho \mod N$ [Chen, Ratz, Staud, Vaudrevange (2012)]  $(\rho = 0 \text{ without GS})$ 

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## **Surveying Discrete Symmetries**

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#### We will provide:

- general criteria for identical symmetries  $\rightarrow$  identify redundancies
- novel algorithm for identifying maximal discrete symmetry
- models consistent with Pati—Salam group, some favoring Dirac neutrino mass
- minimal solutions of both RPV and R-parity conserving surveys
- a counter-example to statement in the literature regarding L-viol. symmetries

• Discrete symmetries may be equivalent  $\rightarrow$  redundancy

Discrete symmetries may be equivalent → redundancy

#### **Criteria of equivalence for symmetry of order** *N*:

- <u>Common divisors</u>: if symmetry order N and all the charges have a common divisor M,  $\mathbb{Z}_N^R$  is equivalent to  $\mathbb{Z}_{N/M}^R$  with its charges divided by M - <u>Non-trivial centers</u>: in the presence of an SU(N) gauge factor, acting with the center of SU(N), the  $Z_{SU(M)} \simeq \mathbb{Z}_M$ , is an equivalent symmetry - <u>Hypercharge-shifts</u>: can add multiples of the hypercharge to respective field charges to obtain equivalent symmetries

- <u>**Coprime factors</u>**: multiplying all charges by a factor co-prime with the symmetry order *N* leads to the same symmetry</u>

less general criteria can be found in [Petersen, Ratz, Schieren (2009)]

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- With charges as variables, using *Smith form* can "diagonalize" constraint matrix

- $d_{n_q}$  is the maximal meaningful order N
- <u>Bonus</u>: can also apply to inequality constraints

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- Under discrete symmetry, the field transforms  $\phi \stackrel{\mathbb{Z}_N}{\mapsto} \mathrm{e}^{2\pi \,\mathrm{i}\,q/N} \phi$
- Can parametrize the set of constraints as  $\sum_{j=1}^{k} a_{ij} q_j = 0 \mod N$   $\forall 1 \le i \le n_c$ if system is overconstrainted  $\rightarrow$  no U(1) solution, but  $\mathbb{Z}_N^R$  solution is possible
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 $\blacktriangleright \quad U \cdot A \cdot V = D$ U. V – unimodular  $A \cdot q = 0 \mod N$ matrices and  $q = V \cdot \begin{pmatrix} k_1 \frac{d_{n_q}}{d_1} \\ \vdots \\ k_{n_q} \frac{d_{n_q}}{d_1} \end{pmatrix} \mod N$ where  $D = \begin{pmatrix} a_1 \\ \ddots \\ d_{n_q} \end{pmatrix}$ applying the method has all other  $d_i$  as divisors - imposing SU(5) + Weinberg op.  $\rightarrow$  maximal order is 24  $d_{n_q}$  is the maximal meaningful order N- imposing SO(10) + Weinberg op.  $\rightarrow$  maximal order is 4 Bonus: can also apply to inequality constraints (agreement with literature)

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- RPV in general not compatible with SU(5), SO(10)
- (  $\overline{U}\overline{D}\overline{D}$  and  $LL\overline{E}$  are allowed or forbidden simultaneously)
- partial unification of Pati—Salam  $G_{PS} = SU(4) \times SU(2)_L \times SU(2)_R$  is allowed (not considered in previous surveys)

- don't enforce anomaly universality for Pati—Salam, since no single unifying gauge group  $\rightarrow$  PS doesn't predict coupling unification in general

## **Survey of the Solutions**

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- many RPV symmetries consistent with Pati—Salam
- found models (eg.  $\mathbb{Z}_{12}^R$ ) "effectively" R-parity conserving and even order, but without R-parity
- confirmed through scan the unique R-parity conserving solution  $\mathbb{Z}_4^R$
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- Can identify some interesting features just from operators:

*Example*: for B-viol. found

 $\left. \begin{array}{c} \text{assume Pati-Salam compatibility} \\ \text{allow } \overline{U} \, \overline{D} \, \overline{D} \\ \text{forbid } L \, H_u \end{array} \right\} \curvearrowright \text{Forbid Weinberg operator } L \, H_u \, L \, H_u \\ \end{array} \right\}$ 

→ B-viol. models consistent with pheno + Pati—Salam prefer Dirac neutrino masses

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- "non-perturbative"  $\mathbb{Z}_3^R$ 
  - non-universal anomalies (can prove no solution with uni. anom.)

field	Q	$\overline{U}$	$\overline{D}$	L	$\overline{E}$	$H_u$	$H_d$	$\theta$
$\mathbb{Z}_3^R$	1	1	1	1	1	0	0	1

- no B and L-viol. @ renormalizable level, only "non-perturbatively"
- $\rightarrow$  assuming R symmetry breaking is of order  $m_{3/2}$

$$\mathscr{W}_{\text{eff}} \supset \frac{m_{3/2}}{M_{\text{P}}} L L \overline{E} + \frac{m_{3/2}}{M_{\text{P}}} Q L \overline{D} + \frac{m_{3/2}}{M_{\text{P}}} \overline{U} \overline{D} \overline{D}$$

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- $LH_u$  is suppressed by  $\,m_{3/2}^2/M_P\,$  , but the  $\mu$  term is of order  $\,m_{3/2}\,$
- $\rightarrow$  <u>counter-example to statement</u>: in L-viol. RPV they are <u>always</u> of same size

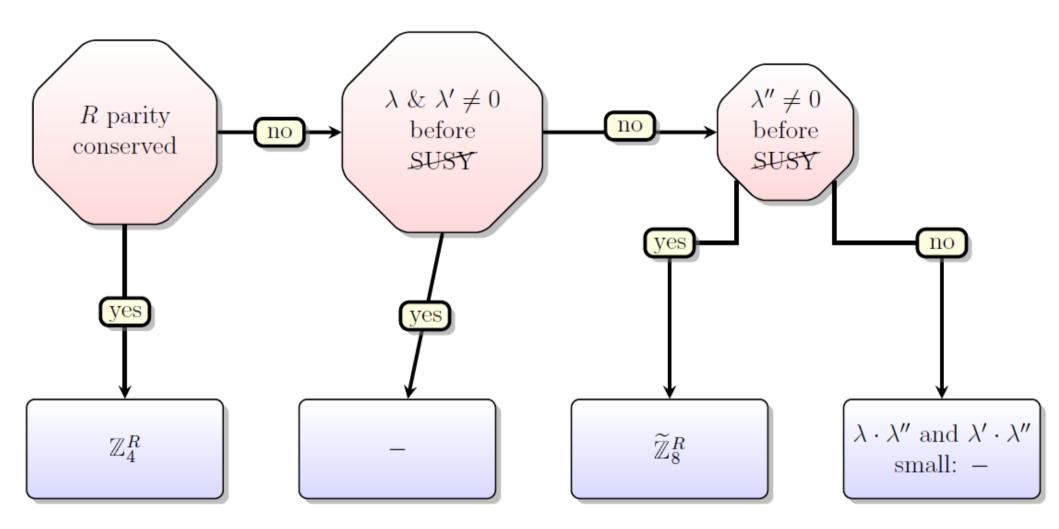


Figure 1: Summary of our results. We present the simplest discrete R symmetries with universal anomalies and the specified properties. The symbol "–" indicates the absence of a solution.

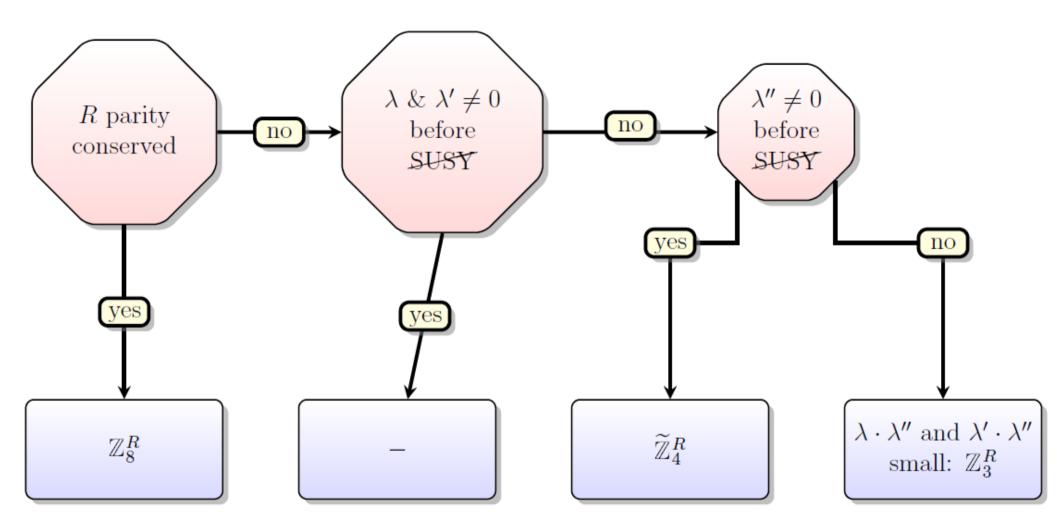


Figure 2: Summary of our results. We present the simplest discrete R symmetries with non–universal anomalies and the specified properties. The symbol "–" indicates the absence of a solution.

### With flavor symmetry:

- anomalous U(1) always present after compactification
- long history of using it as flavor symmetry [Binetruy, Ramond (1994)]

(introduce flavon(s), U(1) charges determine how many flavons couple to

different Yukawas, after flavon gets vev different Yukawa "textures" emerge)

- combine baryon triality  $\mathbb{Z}_3$  and U(1) flavor:  $U(1)_R \to \mathbb{Z}_3$  [Dreiner, Luhn, Murayama, Thormier (2008)]
- combine nice  $\mathbb{Z}_4^R$  and U(1) flavor:  $U(1)_R o \mathbb{Z}_4^R$  [Dreiner, Opferkuch, Luhn (2014)]
- use  $\mathbb{Z}_N^R$  directly as flavor symmetry [Babu, Gogoladze, Wang (2002)]

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### With baryogenesis:

– decay of heavy baryon scalar which dominates before nucleosynthesis is a source of both cold DM and baryon asymmetry, model given by  $\mathbb{Z}_9 \times \mathbb{Z}_2$ [Kitano, Murayama, Ratz (2002)]

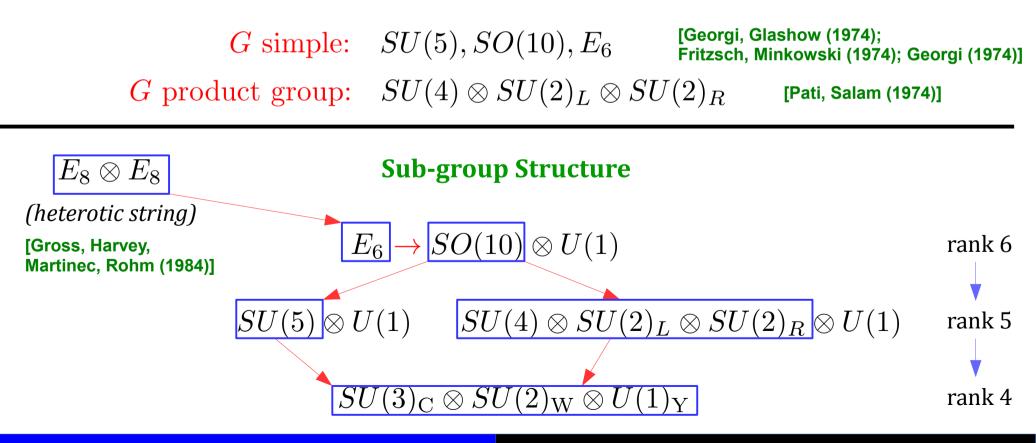
(... some work in progress)

Seen how constraining proton decay to model building is, so

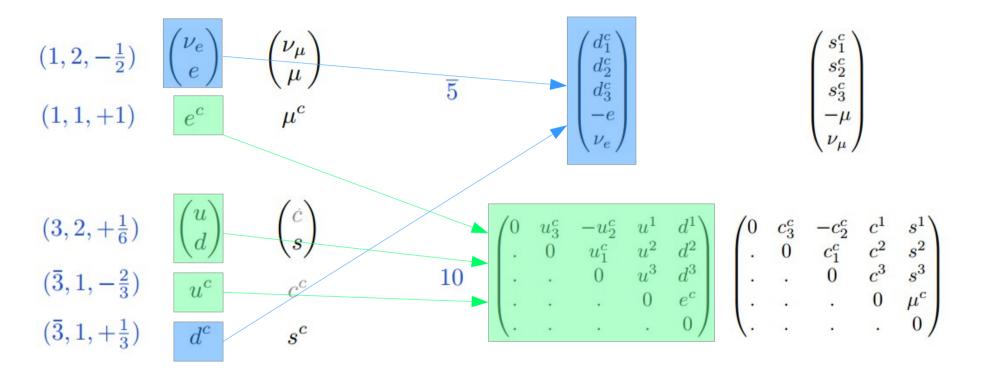
... a brief word about

recent nucleon decay searches at Super-K

- Can unify 3 seeming unrelated forces of the SM into 1 simple Lie group  $SU(3)_{
  m C}\otimes SU(2)_{
  m W}\otimes U(1)_{
  m Y} o G$
- **Many hints**: gauge coupling "unification" at high energies, charge is quantized, quark and lepton mixing patterns seem to have some structure, etc. *(inflation?)*
- Most promising candidates (anomaly free, rank ≥ 4, contain SM as subgroup):



• Particle content of the SM fits into  $\overline{5}$ , 10 of SU(5)



- Gauge and Higgs sectors fit into  $24, \overline{5}$  of SU(5) (new g. bosons X, Y and H. triplet T)
- Even better with SO(10) 16 of SO(10) contains all SM particles + rh.  $\nu$
- Many great features

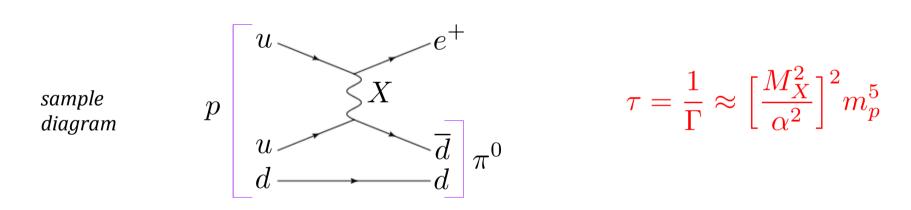
...but GUT scale ~10^(16) GeV, colliders can't reach

...but proton decay can!

# Proton Decay: non-SUSY GUTs

- In non-SUSY models, such as minimal SU(5), proton decay originates from dim-6 operators:  $\frac{QQQL}{\Lambda^2}, \frac{\overline{U}\overline{U}\overline{E}\overline{D}}{\Lambda^2}, \frac{\overline{U}\overline{E}QQ}{\Lambda^2}, \frac{\overline{D}\overline{U}QL}{\Lambda^2}$
- The typical dominant non-SUSY decay channel is  $p 
  ightarrow e^+ \pi^0$

→ mediated by GUT gauge bosons X and Y, can also be by GUT color triplet Higgs T (*limit on this leads to the doublet-triplet splitting problem*)

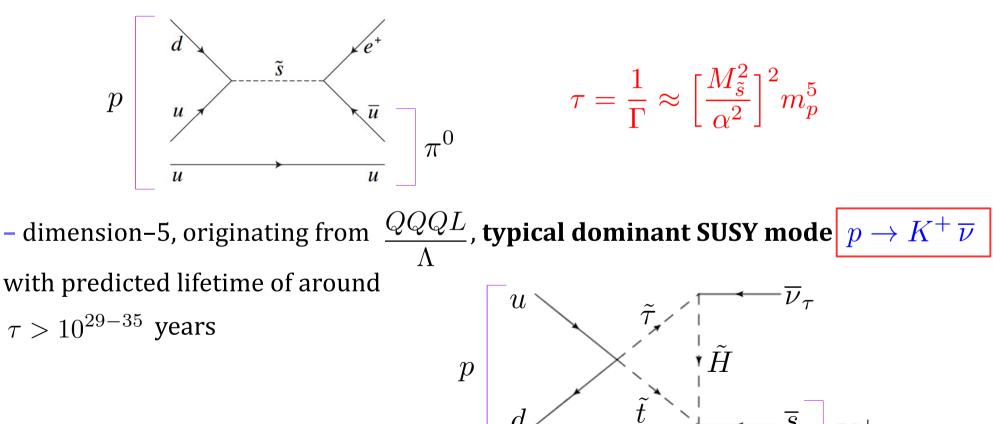


• For minimal *SU(5)*, predicted lifetime is  $\tau > 10^{29\pm2}$  years

ightarrow ruled out by experiment with  $au > 10^{34}$  years [Nishino et. al. (Super-K) (2012)]

# Proton Decay: SUSY GUTs

- SUSY pushes up unification scale  $\rightarrow \tau(p \rightarrow e^+\pi^0) > 10^{35-38}$  years
- Sparticles present → new decays start to dominate
  - dimension–4, originating from operators  $LQ\overline{D}, \overline{U}\overline{D}\overline{D}$ 
    - $\rightarrow \tau \sim 1s$  if squark mass ~ TeV, forbidden by *R* parity



# Searching for Proton Decay: Water Cherenkov Detectors

 To see if proton lives longer than 10<sup>31</sup> years, can either look at 1 proton for 10<sup>32</sup> years ... OR ... look at 10<sup>33</sup> protons (~10 kiloton) for 1 month
 → large underground detectors → water Cherenkov detectors, cheap + large

- (1979) Irvine–Michigan–Brookhaven (IMB)
  - no proton found, limit  $au(p 
    ightarrow e^+ \pi^0) > 10^{32}$  years (1990)
  - saw SN1987A neutrinos and atmospheric neutrino "anomaly" (later oscillations)
- (1980~) Kamiokande

- saw SN1987A neutrinos and atmospheric
 neutrino "anomaly", solar neutrinos
 *Koshiba's Nobel Prize (2002)*



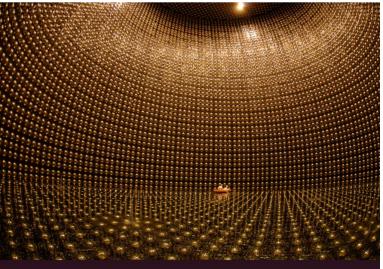
IMB experiment

# Searching for Proton Decay: Water Cherenkov Detectors

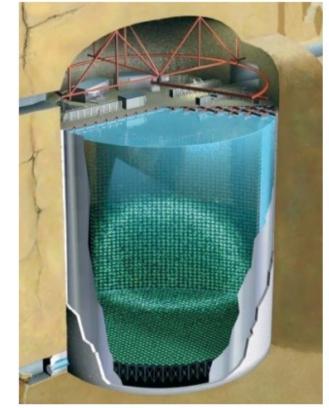
- (1996 ) Super–Kamiokande (SK)
  - largest water-C. detector, successor of Kamiokande
  - discovery of neutrino oscillations (1998)
  - → neutrinos have mass
  - lifetime of proton

$$\tau(p \rightarrow e^+ \pi^0) > 10^{34} \ \ \, {\rm years}$$

#### SK experiment



from Super-K Webpage / Ed Kearns, NEPPRS 09



### Super-Kamiokande

22.5 kton fiducial volume 7.5×10<sup>33</sup> p + 6×10<sup>33</sup> n

SK-I: 1996 - 2001 11146 50-cm inner PMTs , 40% coverage 1885 20-cm outer PMTs

#### SK-II: Jan 2003 - Oct 2005

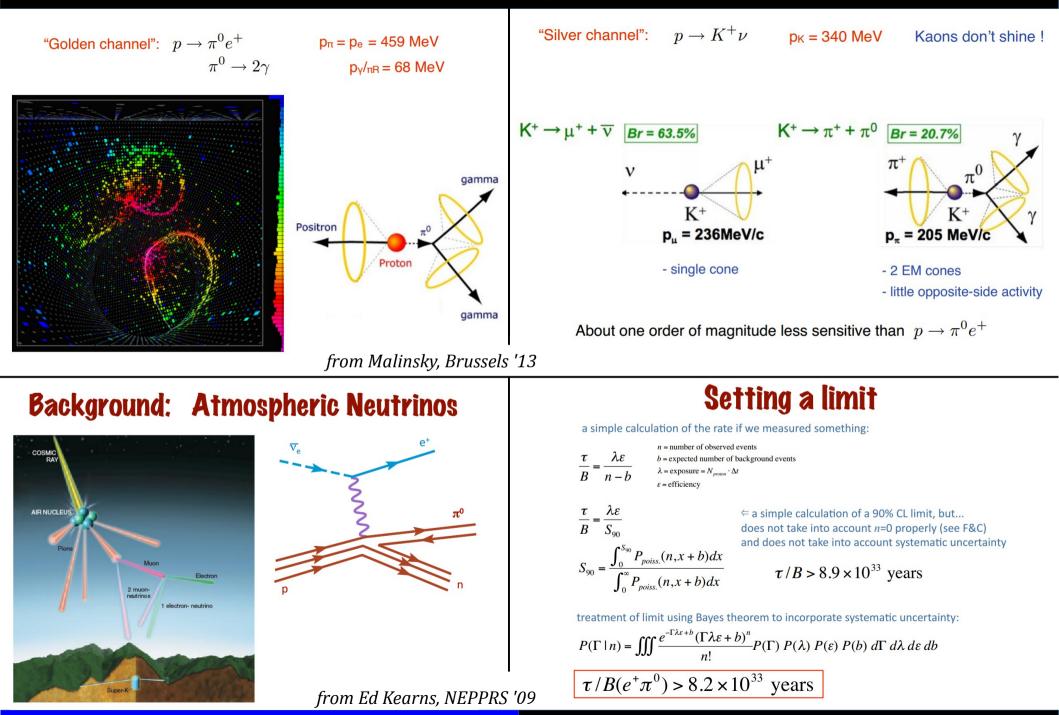
Recovery from accident 5182 50-cm inner PMTs Acrylic + FRP protective Outer detector fully restored



SK-III: May 2006 - August 2008 Restored 40% coverage Outer detector segmented (top | barrel | bottom)

SK-IV: September 2008 -SK-IV Replace all electronics – 2008 T2K beam – late 2009 Add gadolinium - 201?

# Proton Decay at Super-Kamiokande



# **Story is Actually More Complex ...**

- Proton decay hasn't been found in  $\sim$ 30 years, why continue looking?
- Many models beyond simple *SU(5)*, proton decay can rule some of them out ...

Model	Ref.	Modes	$\tau_N$ (years)
Minimal $SU(5)$	Georgi, Glashow [2]	$p \rightarrow e^+ \pi^0$	$10^{30} - 10^{31}$
Minimal SUSY $SU(5)$	Dimopoulos, Georgi [11], Sakai [12]	$p \rightarrow \bar{\nu}K^+$	
	Lifetime Calculations: Hisano,	$n \rightarrow \bar{\nu} K^0$	$10^{28} - 10^{32}$
	Murayama, Yanagida [13]		
SUGRA $SU(5)$	Nath, Arnowitt [14, 15]	$p \rightarrow \bar{\nu}K^+$	$10^{32} - 10^{34}$
SUSY $SO(10)$	Shafi, Tavartkiladze [16]	$p \rightarrow \bar{\nu}K^+$	
with anomalous		$n \rightarrow \bar{\nu} K^0$	$10^{32} - 10^{35}$
flavor $U(1)$		$p \rightarrow \mu^+ K^0$	
SUSY $SO(10)$	Lucas, Raby [17], Pati [18]	$p \rightarrow \bar{\nu}K^+$	$10^{33} - 10^{34}$
MSSM (std. $d = 5$ )		$n \rightarrow \bar{\nu} K^0$	$10^{32} - 10^{33}$
SUSY $SO(10)$	Pati [18]	$p \rightarrow \bar{\nu}K^+$	$10^{33} - 10^{34}$
ESSM (std. $d = 5$ )			$\lesssim 10^{35}$
SUSY $SO(10)/G(224)$	Babu, Pati, Wilczek [19, 20, 21],	$p \rightarrow \bar{\nu}K^+$	$\lesssim 2 \cdot 10^{34}$
MSSM or ESSM	Pati [18]	$p \rightarrow \mu^+ K^0$	
$(new \ d = 5)$		B	$\sim (1 - 50)\%$
SUSY $SU(5)$ or $SO(10)$	Pati [18]	$B \cdot p \rightarrow e^+ \pi^0$	$\sim 10^{34.9\pm1}$
MSSM $(d = 6)$			
Flipped $SU(5)$ in CMSSM	Ellis, Nanopoulos and Wlaker[22]	$p \rightarrow e/\mu^+ \pi^0$	$10^{35} - 10^{36}$
Split $SU(5)$ SUSY	Arkani-Hamed, et. al. [23]	$p \rightarrow e^+ \pi^0$	$10^{35} - 10^{37}$
SU(5) in 5 dimensions	Hebecker, March-Russell[24]	$p \rightarrow \mu^+ K^0$	$10^{34} - 10^{35}$
		$p \rightarrow e^+ \pi^0$	
SU(5) in 5 dimensions	Alciati et.al.[25]	$p \rightarrow \bar{\nu}K^+$	$10^{36} - 10^{39}$
option II			
GUT-like models from	Klebanov, Witten[26]	$p \rightarrow e^+ \pi^0$	$\sim 10^{36}$
Type IIA string with D6-branes			

#### [Bueno (2007)]

TABLE I: Summary of the expected nucleon lifetime in different theoretical models.

### ..... many predicted modes

B + L
$\Delta B = 2$ , TeV < scale < GUT
$\lambda'_{\rm uds} < 10^{-8}$
6 dimensions
invisible
radiative from Ed Kearns, NEPPRS '09

Many predictions in the  $\tau \approx 10^{34-36}$  year range

... are we on the verge of discovery (*Super–K*, or near future Hyper–K)?

### **Some Novel SK Searches**

- Minimal *SU(5)* and SUSY *SU(5)* ~ *ruled out by experiment*
- Can consider larger group, eg.  $SO(10) \rightarrow Pati-Salam$
- In certain variations trilepton modes can be significant
  - $\rightarrow$  maybe also useful for baryogenesis [Gu and Sarkar (2012)]:

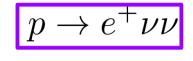
 $\rightarrow$  predicted lifetimes  $\,\tau\approx 10^{31-33}$  years  $\,$  [Pati (1984), Gu and Sarkar (2012)]:

• First 3-body decay search in SK

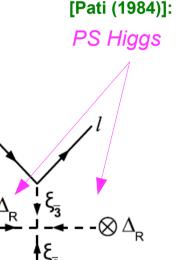
• In SK can't see neutrinos, only spectra from e+,  $\mu$ + (can use  $\mu \rightarrow \text{evv}$  spectra to describe the above) Chen, Takhistov [PRD (2014)]

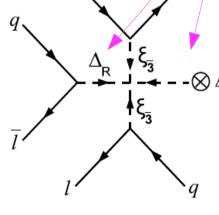
- Spectral fit analysis
  - $\rightarrow$  set limit ~ 10^32 years (>1 order improvement)

Takhistov et. al. (Super-K Collab.) [PRL (2014)]



$$p \to \mu^+ \nu \nu$$





### **Some Novel SK Searches**

**<u>Other modes</u>** which can be similarly analyzed (*spectral fit to momenta*):

- $p \to e^+ X \mid p \to \mu^+ X$  inclusive decays (X is invisible particle)
- $n 
  ightarrow \gamma 
  u$  radiative mode [Nath and Perez (2007)]:
- $np \to e^+ \nu$   $np \to \mu^+ \nu$   $np \to \tau^+ \nu$  dinucleon decays, which can arise in

models with extended Higgs sector, may be connected to baryogenesis

### Results for SK search of nucleon decay modes with charged lepton + inv.:

#### Takhistov et. al. (Super-K Collab.) (accepted to PRL)

Mode	SK I-IV Sensitivity	(years)	SK I-IV Limit (years)	PDG Limit (years)
$p \rightarrow e^+ X$	$7.9 \cdot 10^{32}$		$7.9 \cdot 10^{32}$	—
$p \to \mu^+ X$	$7.7 \cdot 10^{32}$		$4.1 \cdot 10^{32}$	—
$n  ightarrow \nu \gamma$	$5.8 \cdot 10^{32}$	in	$5.5 \cdot 10^{32}$	$2.8 \cdot 10^{31}$
$np \rightarrow e^+ \nu$	$9.9 \cdot 10^{31}$	arer	$2.6 \cdot 10^{32}$	$2.8 \cdot 10^{30}$
$np \to \mu^+ \nu$	$1.1 \cdot 10^{32}$		$2.2 \cdot 10^{32}$	$1.6 \cdot 10^{30}$
$np \to \tau^+ \nu$	$1.1 \cdot 10^{31}$		$2.9 \cdot 10^{31}$	
C				

not in PDG, first ever search

08/25/15

[Arnellos Marciano (1982), Arnold, Fornal, Wise (2013); Bryman (2014)]:

# **Other Recent SK Searches**

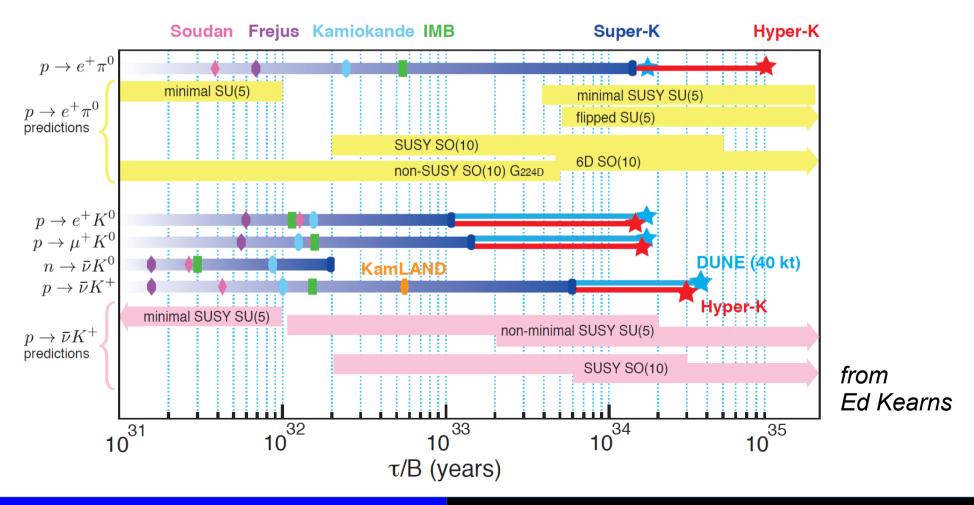
$$pp 
ightarrow \pi^+ \pi^+$$
 [Gustafson et. al. (2015)]:  
 $n 
ightarrow 
u \pi^0$  [Abe et. al. (2014)]:  
 $p 
ightarrow 
u \pi^+$  [Abe et. al. (2014)]:  
 $pp 
ightarrow K^+ K^+$  [Litos et. al. (2014)]:

....

### Future

- HyperK is bigger version of SuperK (20 x SK size)
- aside mass hierarchy and CP violation also improved proton decay search ....

### (1<sup>st</sup> proto-collaboration meeting June 2015)



# **Thank You!**