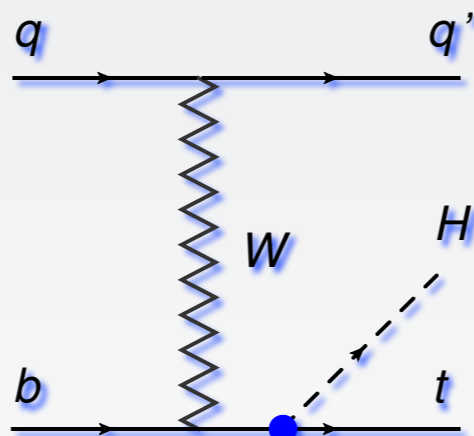
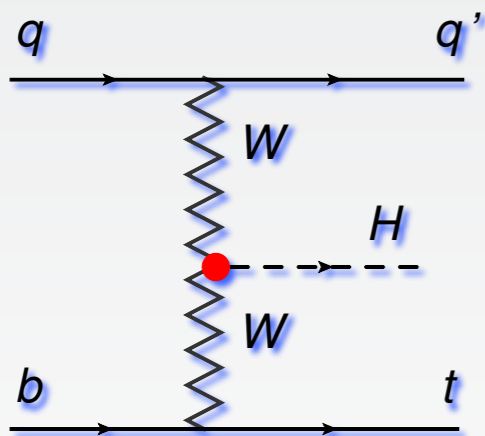
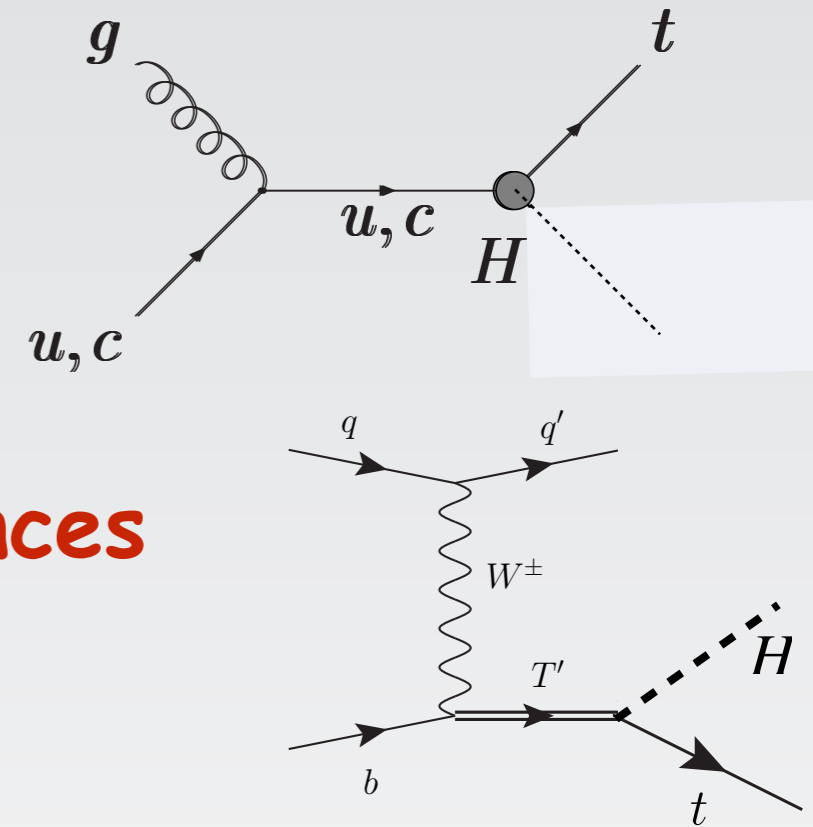


Single top + Higgs (theory)



Why to look at $t+H$ final states ?

- probe FCN couplings tuH , tcH
- search for singly-produced tH resonances
- test structure of g_{Htt} coupling
(beyond what can be achieved by $gg \rightarrow H$, ttH)
-



in this talk focus on g_{ttH} tests
(including g_{ttH} CP properties !)

ttH interaction beyond the SM

$$\mathcal{L}_t^{SM} = -\frac{1}{\sqrt{2}} Y_t^{SM} \bar{t} t H \quad \text{Yukawa interaction}$$

$$Y_t^{SM} = \frac{\sqrt{2} m_t}{v} \simeq \sqrt{2} \frac{173}{246} \simeq 0.995 \simeq 1 \quad !!!$$

→ special role of top in EWSB ?

● beyond the SM (scalar + pseudo-scalar coupling)

$$\mathcal{L}_t = -\frac{1}{\sqrt{2}} Y_t^{SM} (\kappa_t \bar{t} t + i \tilde{\kappa}_t \bar{t} \gamma_5 t) H$$

Part 1:

real $g_{t\bar{t}H}$ (\rightarrow CP conserving)

$$Y_t^{SM} \bar{t} t H \rightarrow Y_t \bar{t} t H$$

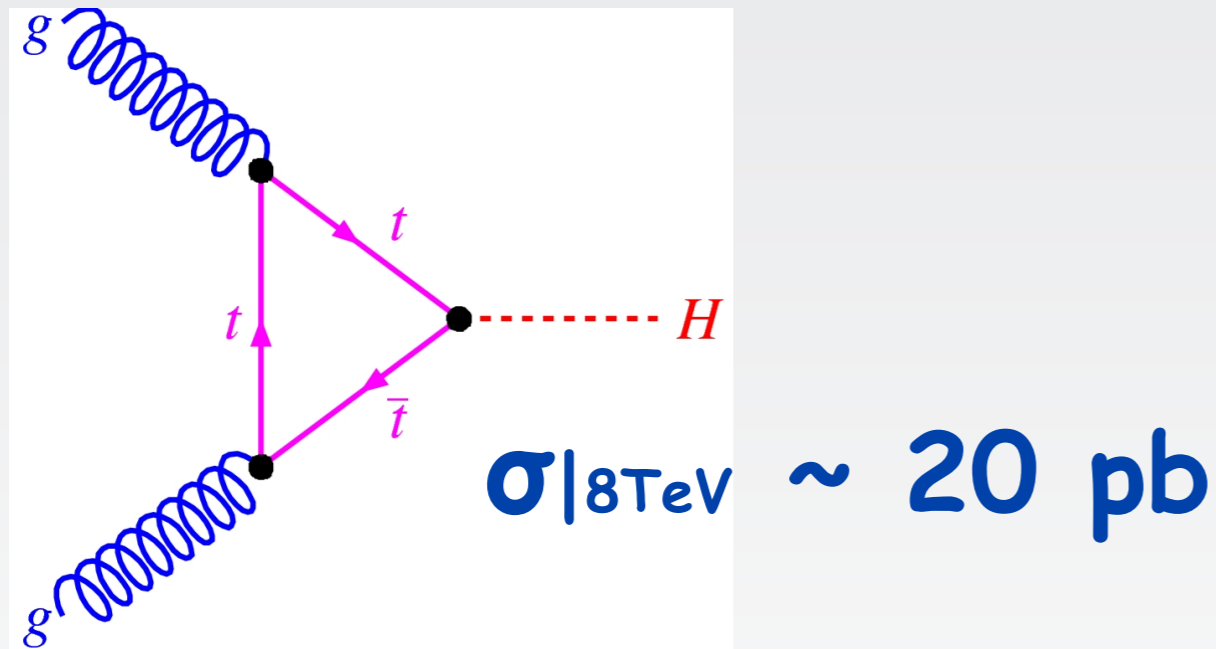
notation :

$$Y_f = C_f Y_f^{SM}$$

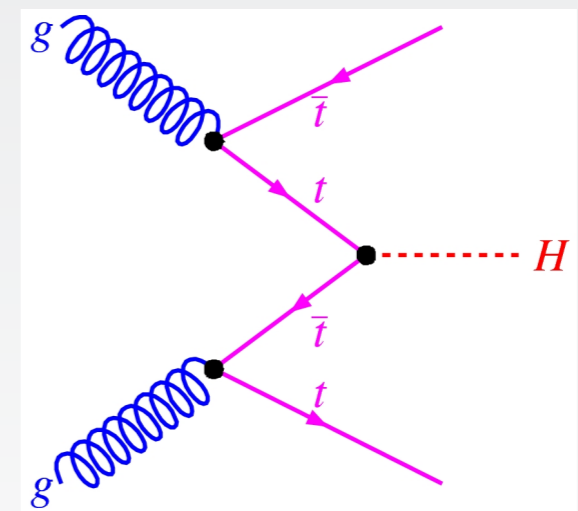
C_f modifier of Hff couplings

main sensitivity to g_{Htt} ($\sim C_t$)

- dominant but indirect (top quark not observed $\rightarrow C_t$ effects could be faked by New Virtual contributions)



- direct (top quark observed) but not yet quite at reach



$$\sigma|_{8\text{TeV}} \sim 130 \text{ fb}$$

both σ 's depend on $C_t^2 \rightarrow$
no sensitivity to g_{Htt} sign!

BSM theories can predict $C_f < 0$!

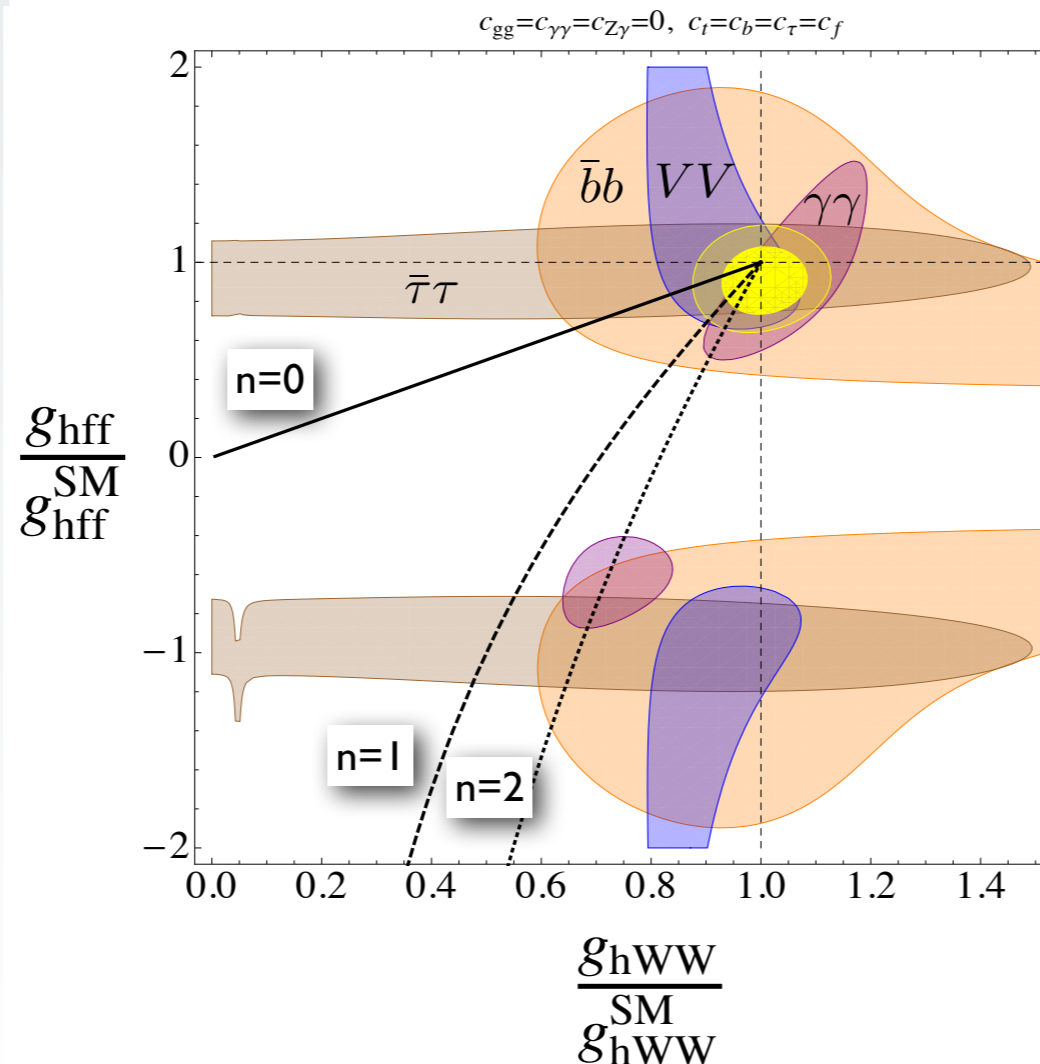
example : Minimal Composite Higgs Models $SO(5) \rightarrow SO(4)$
 global symmetry in a strong sector broken at a scale $f_H > v$

Agashe, Contino, Pomarol, [hep-ph/0412089]

$$\xi = \frac{v^2}{f_H^2}$$

$$\frac{g_{hWW}}{g_{hWW}^{\text{SM}}} = \sqrt{1 - \xi} \quad \frac{g_{hff}}{g_{hff}^{\text{SM}}} = \frac{1 - (1 + n)\xi}{\sqrt{1 - \xi}} \quad n = 0, 1, 2, \dots$$

Depending on realization



■ MCHM limit

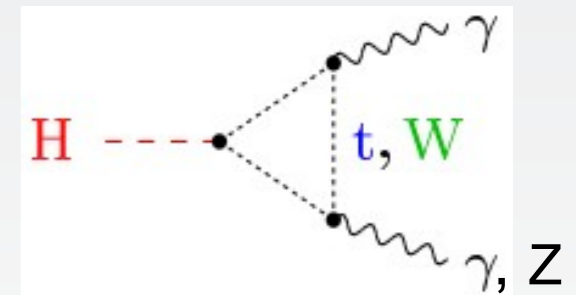
$$\rightarrow f_H \gtrsim 700 \text{ GeV}$$

Falkowski, Riva, Urbano, [1303.1812]

Looking for some way to experimentally discriminate the C_f sign

- linear terms in C_f needed \rightarrow look for interferences in squared amplitudes

• in decays: mainly in loop channels



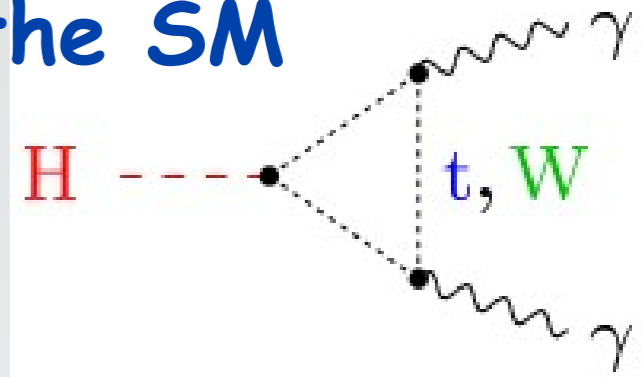
• what about production mechanisms ???



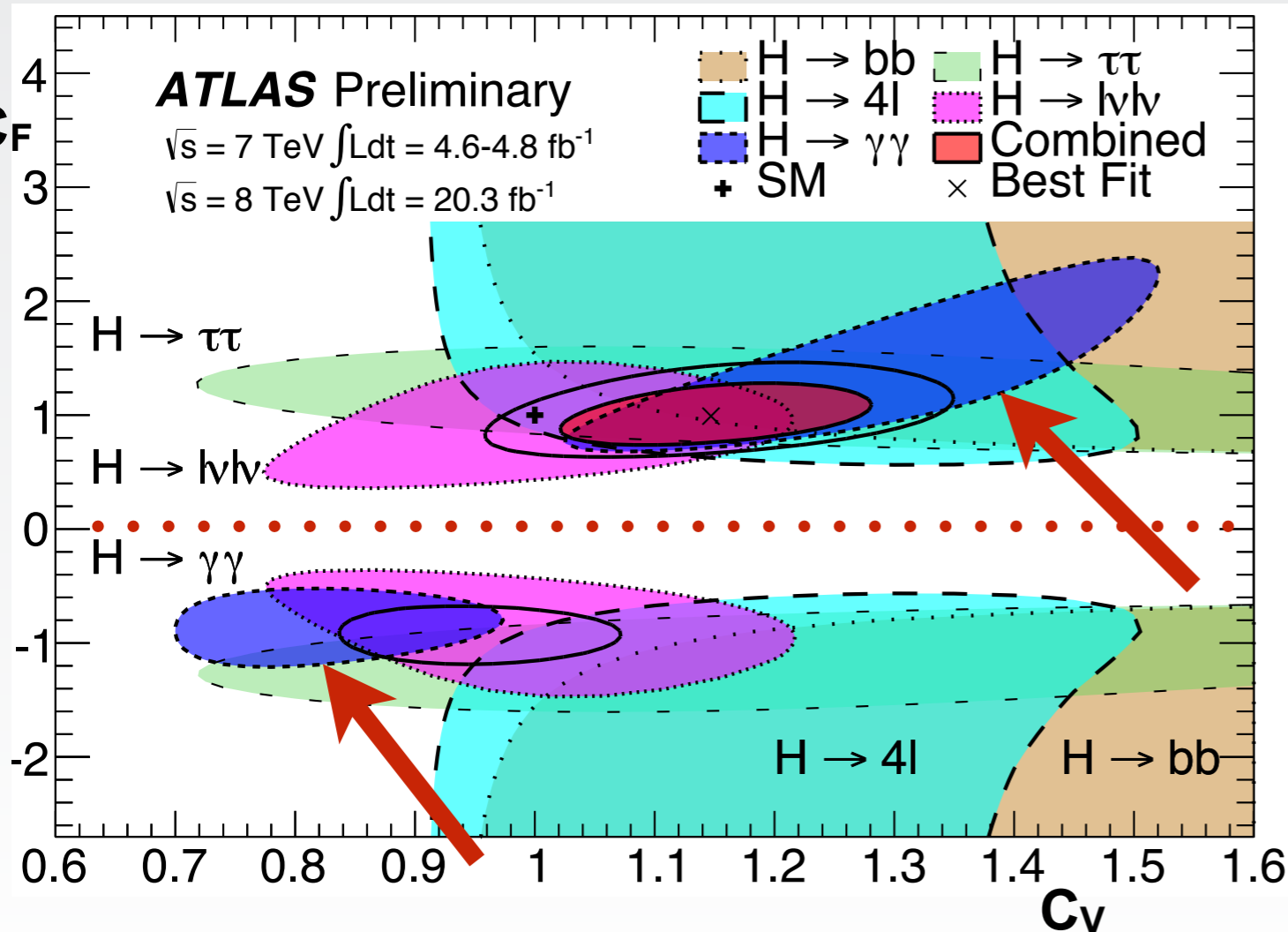
$H \rightarrow \gamma\gamma$ breaks $C_t \rightarrow -C_t$ degeneracy

W and top loops interfere destructively in the SM

$C_t \sim +1$ (SM) $\rightarrow C_t \sim -1$ enhances $BR_{\gamma\gamma}$



ATLAS-CONF-2014-009



$$\sigma(H \rightarrow \gamma\gamma) \sim (5C_V - C_t)^2$$

\rightarrow gives asymmetric constraints for $C_t \rightarrow -C_t$

enhanced $\sigma^{\gamma\gamma}$ rates favor $C_t < 0$ ranges

note !

two solutions for any $\sigma^{\gamma\gamma}$

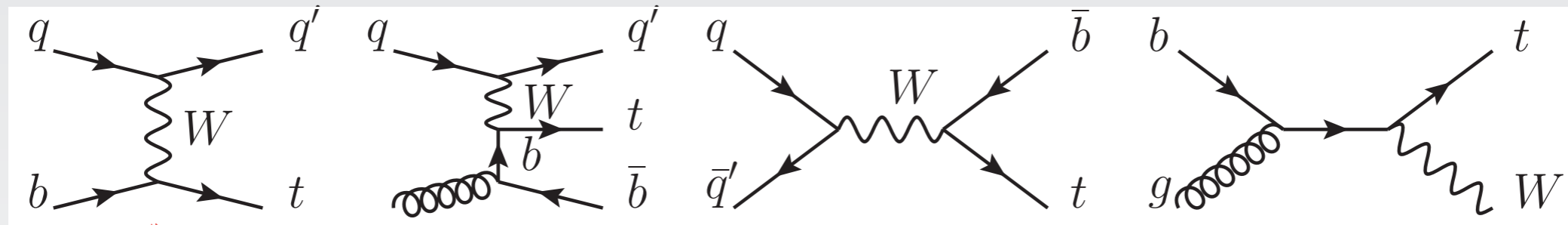
$$\sigma_{SM}^{\gamma\gamma} = \sigma^{\gamma\gamma}(C_V = C_t = 1)$$

$$\simeq \sigma^{\gamma\gamma}(C_V = -C_t = 0.66)$$

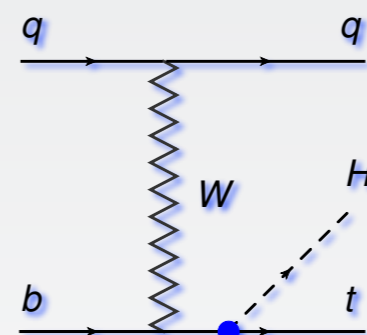
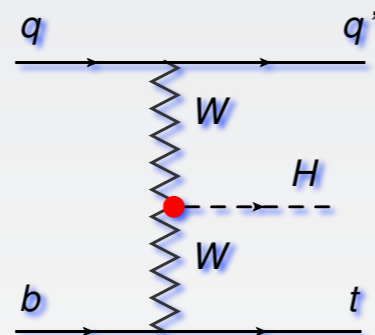
Single top + Higgs production

$$p p \rightarrow t H x \quad (x = q, b, W)$$

- ask for an extra Higgs in single-top production



(t-channel)



(same for tb and tW channels)

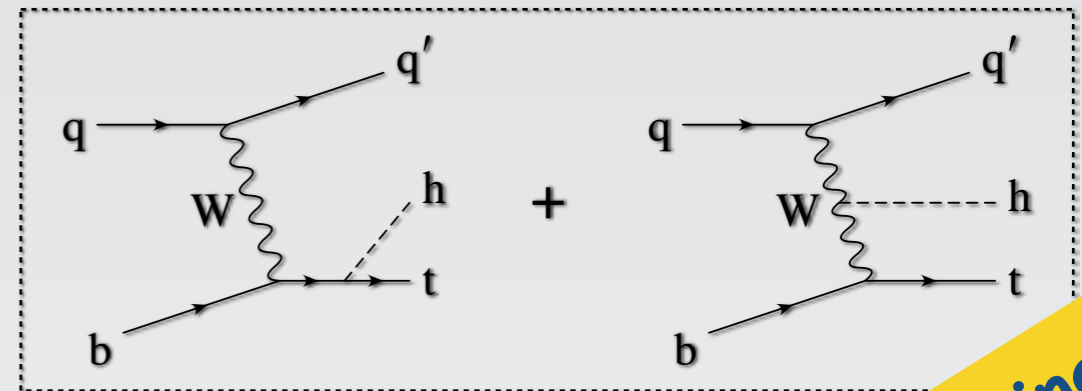
- EW process where Higgs emission from a W interferes with Higgs emission from a top-quark :

$$\sigma \sim a C_t^2 + b C_W^2 + c C_W C_t$$

interference term in σ_{tot} sensitive to C_t sign !

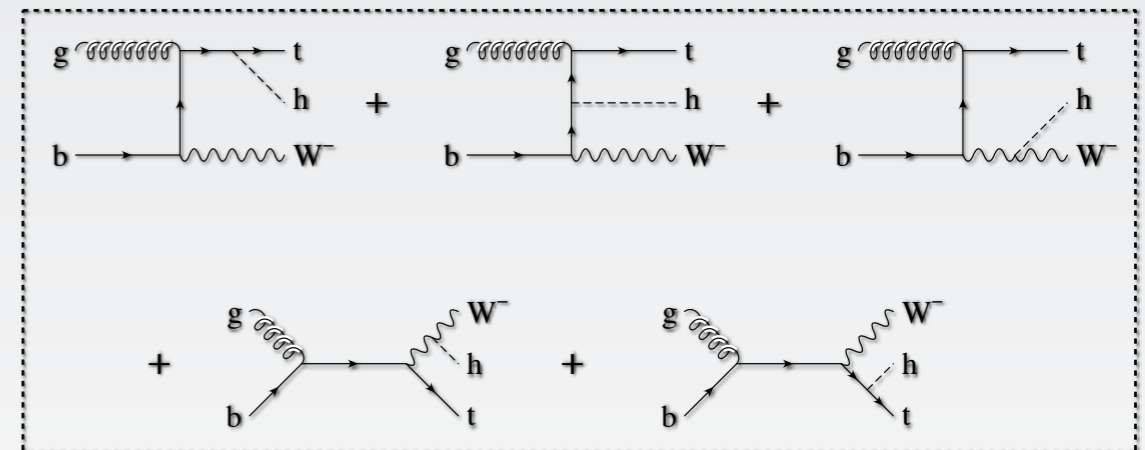
$\sigma(pp \rightarrow t H + X)$ in the SM: 3 channels

- sensitive to g_{Htt} and g_{HWW}
- $\sigma(\text{t-channel}) \sim 1/10 \sigma(\text{ttH})$

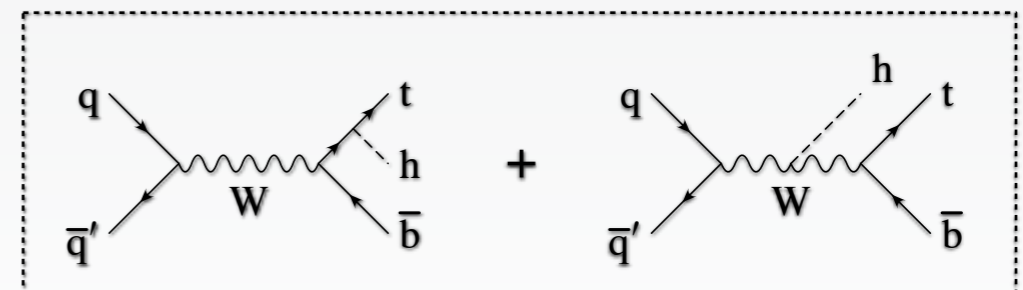


(t-channel)

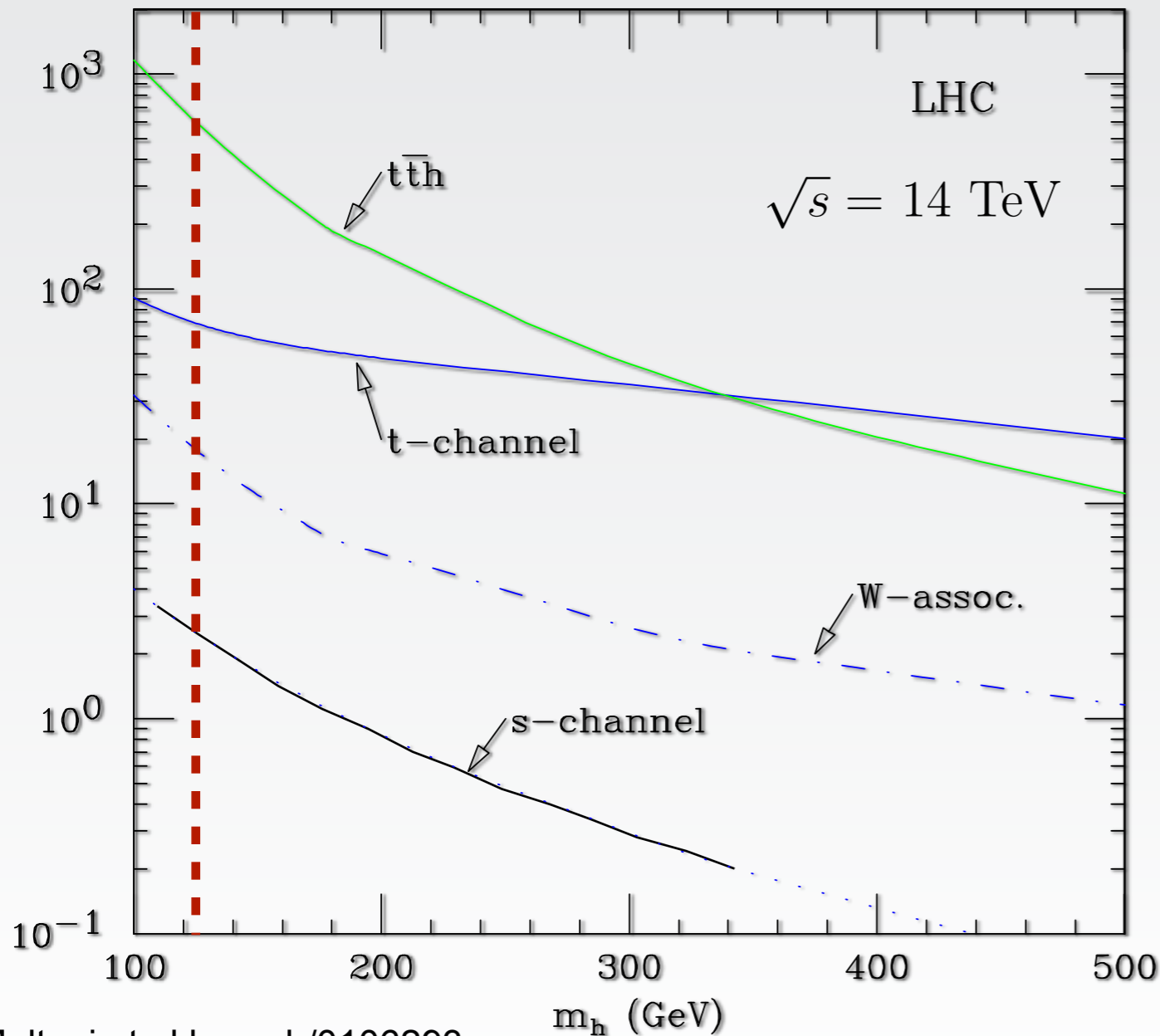
(dominant !)



(W-assoc.)



(s-channel)



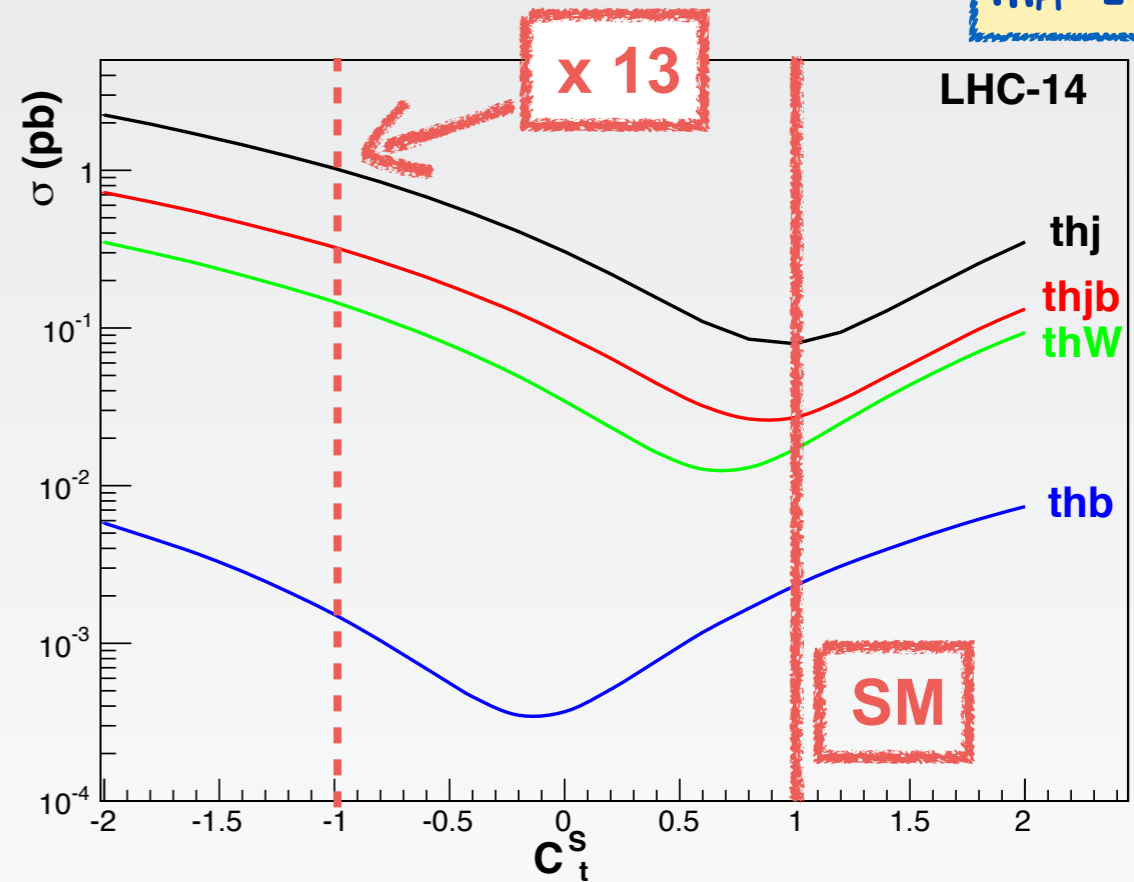
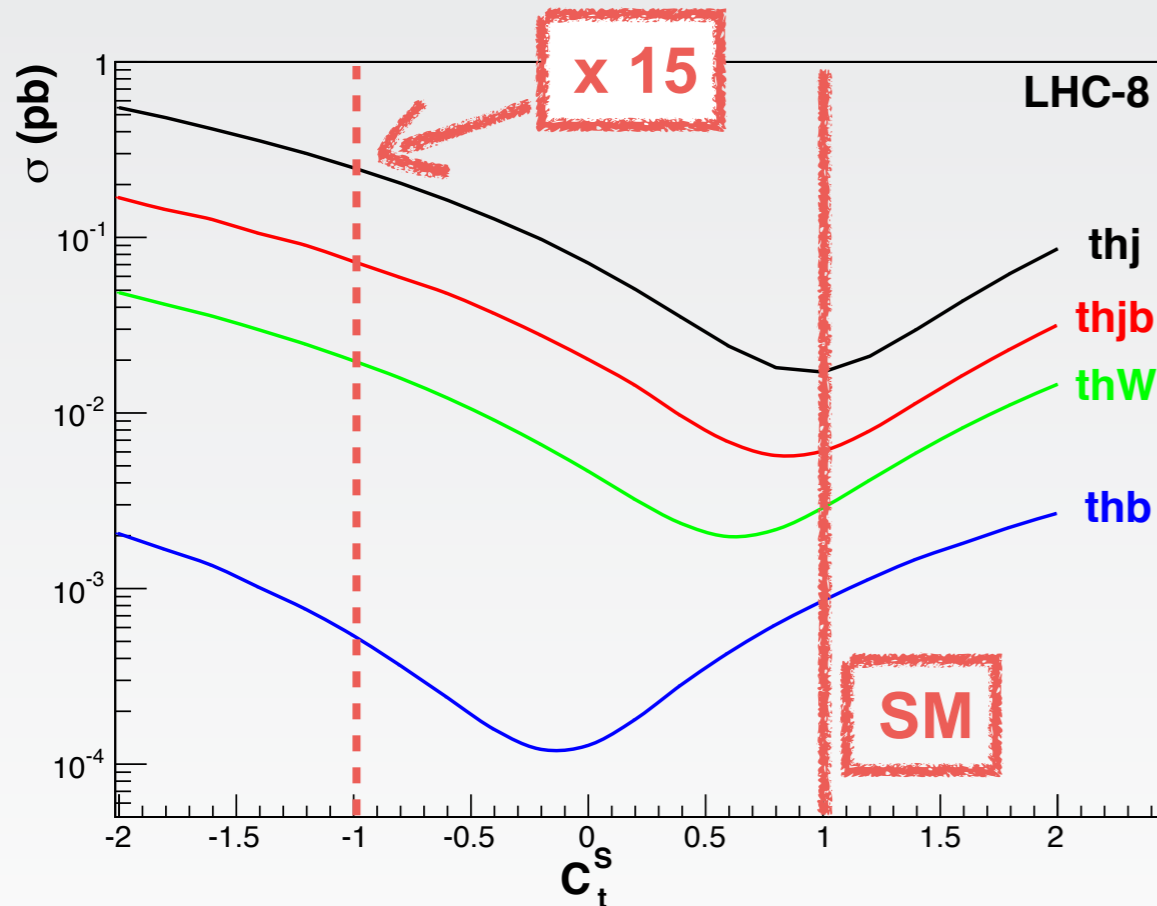
Maltoni et al, hep-ph/0106293

$\sigma(pp \rightarrow t H X)$ [fb] vs C_t at 14 TeV (8 TeV)

	$X = j$	$X = j + b$	$X = W$	$X = b$
$C_t^S = +1$ (SM)	79.4 (17.1)	27.1 (5.95)	17.0 (2.89)	2.32(0.833)
$C_t^S = 0$	305 (71.4)	90.0 (19.8)	34.4 (4.66)	0.368 (0.126)
$C_t^S = -1$	1030 (249)	325 (72.8)	146 (19.8)	1.52 (0.536)

Chang et al, arXiv:1403.2053

$m_H = 125$ GeV



destructive (constructive) interference
in tHj, tHW (tHb) channels for $C_t = 1$

prompted dedicated
studies $\rightarrow \rightarrow \rightarrow$

a few references on tHq sensitivity to C_t at LHC

Tait and Yuan, *Phys.Rev.* D63 (2000) 014018 [hep-ph/0007298]

Maltoni et al, *Phys.Rev.* D64 (2001) 094023 [hep-ph/0106293]

Barger et al, *Phys.Rev.* D81 (2010) 034020 [arXiv:0911.1556]

Biswas et al , *JHEP* 01 (2013) 088 [arXiv:1211.0499] ($H \rightarrow \gamma\gamma + (\text{had}) \text{ top}$)

Farina et al , *JHEP* 05 (2013) 022 [arXiv:1211.3736] ($H \rightarrow bb + (\text{lep}) \text{ top}$)

Biswas et al , *JHEP* 07 (2013) 073 [arXiv:1304.1822] ($H \rightarrow \gamma\gamma, WW, \tau\tau$)

Ellis et al , *JHEP* 1404 (2014) 004 [arXiv:1312.5736]

Englert and Re, *Phys.Rev.* D89 (2014) 073020 [arXiv:1402.0445]

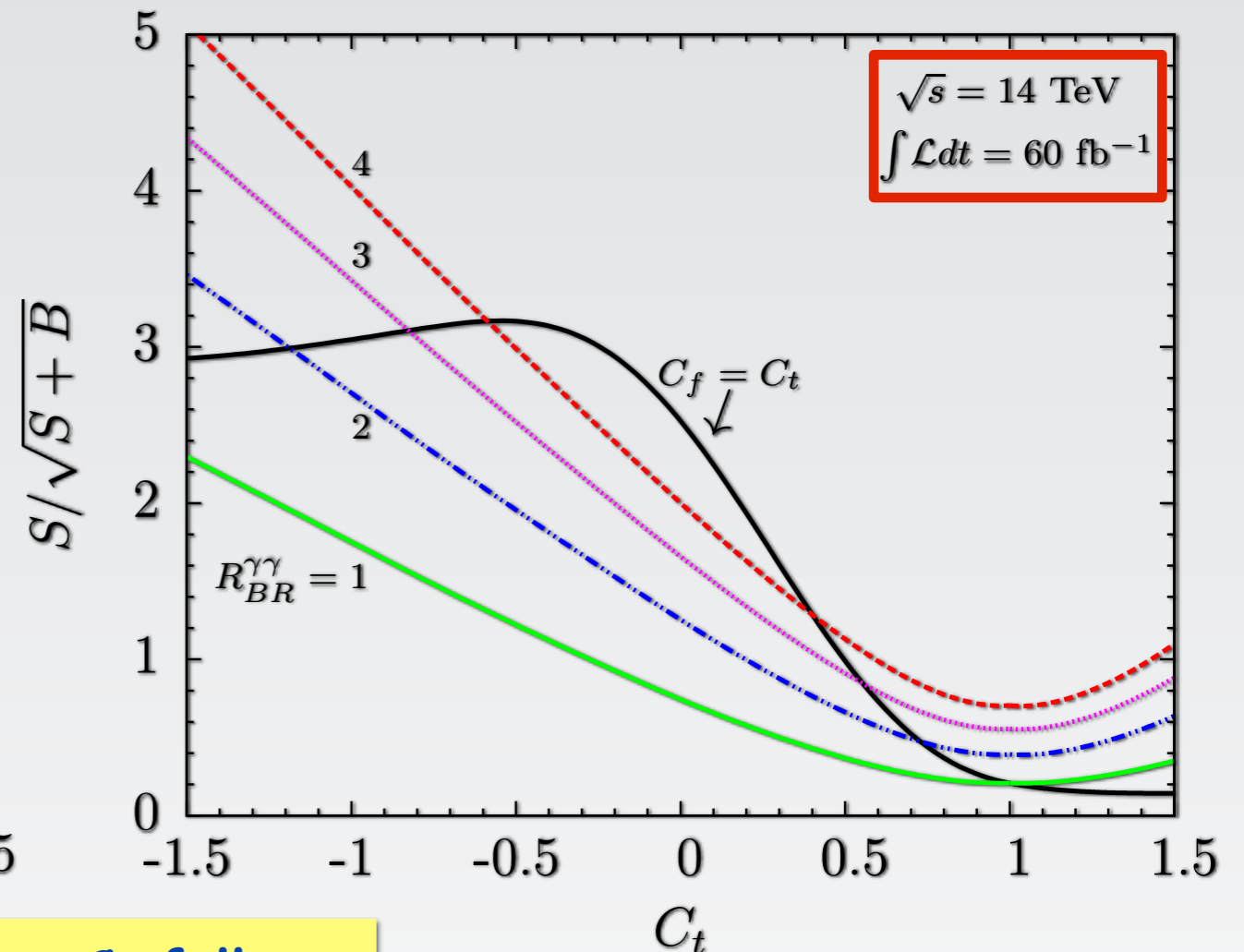
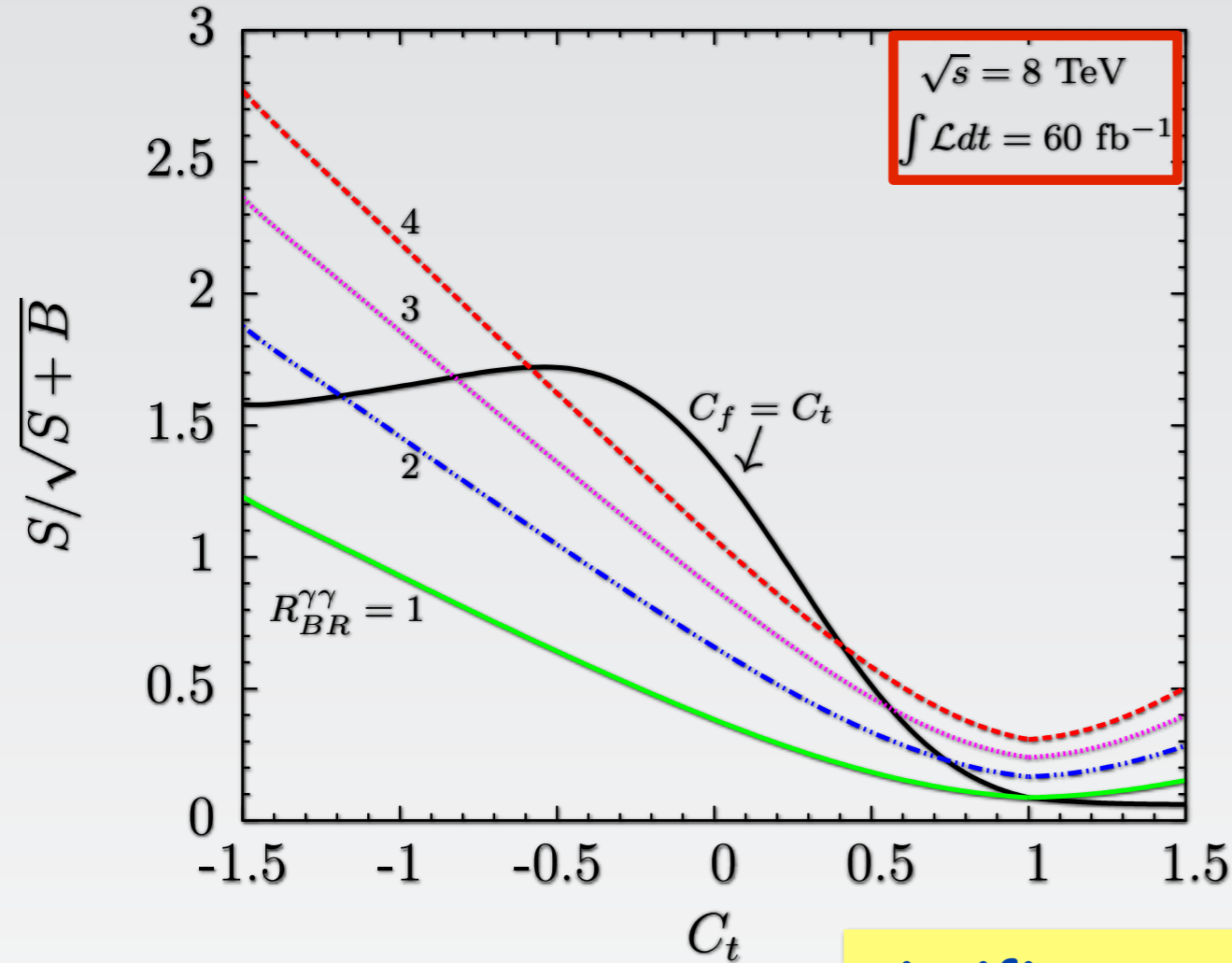
Chang et al, *JHEP* 1405 (2014) 062 [arXiv:1403.2053]

Kobakhidze et al, arXiv:1406.1961

Yue, arXiv:1410.2701

bckgr studies : $H \rightarrow \gamma\gamma + \text{top (had)}$

Biswas et al, arXiv:1211.0499



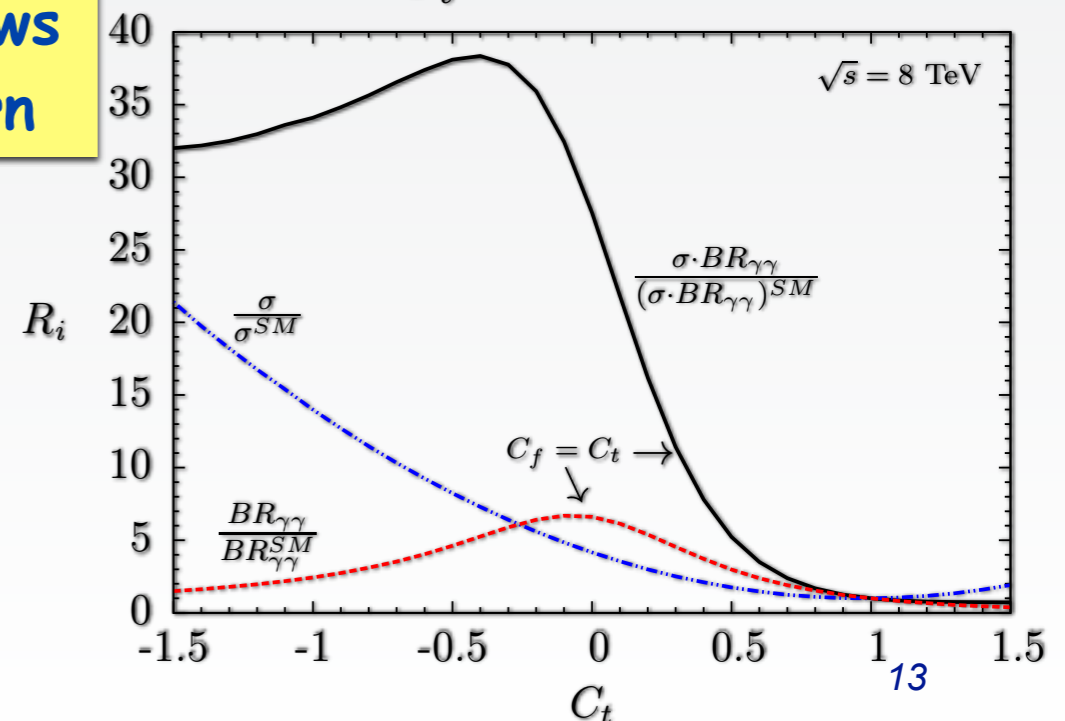
significance vs C_t follows signal strength pattern

$$\frac{BR_{\gamma\gamma}(C_t = -1)}{BR_{\gamma\gamma}^{SM}} \simeq 2.4$$

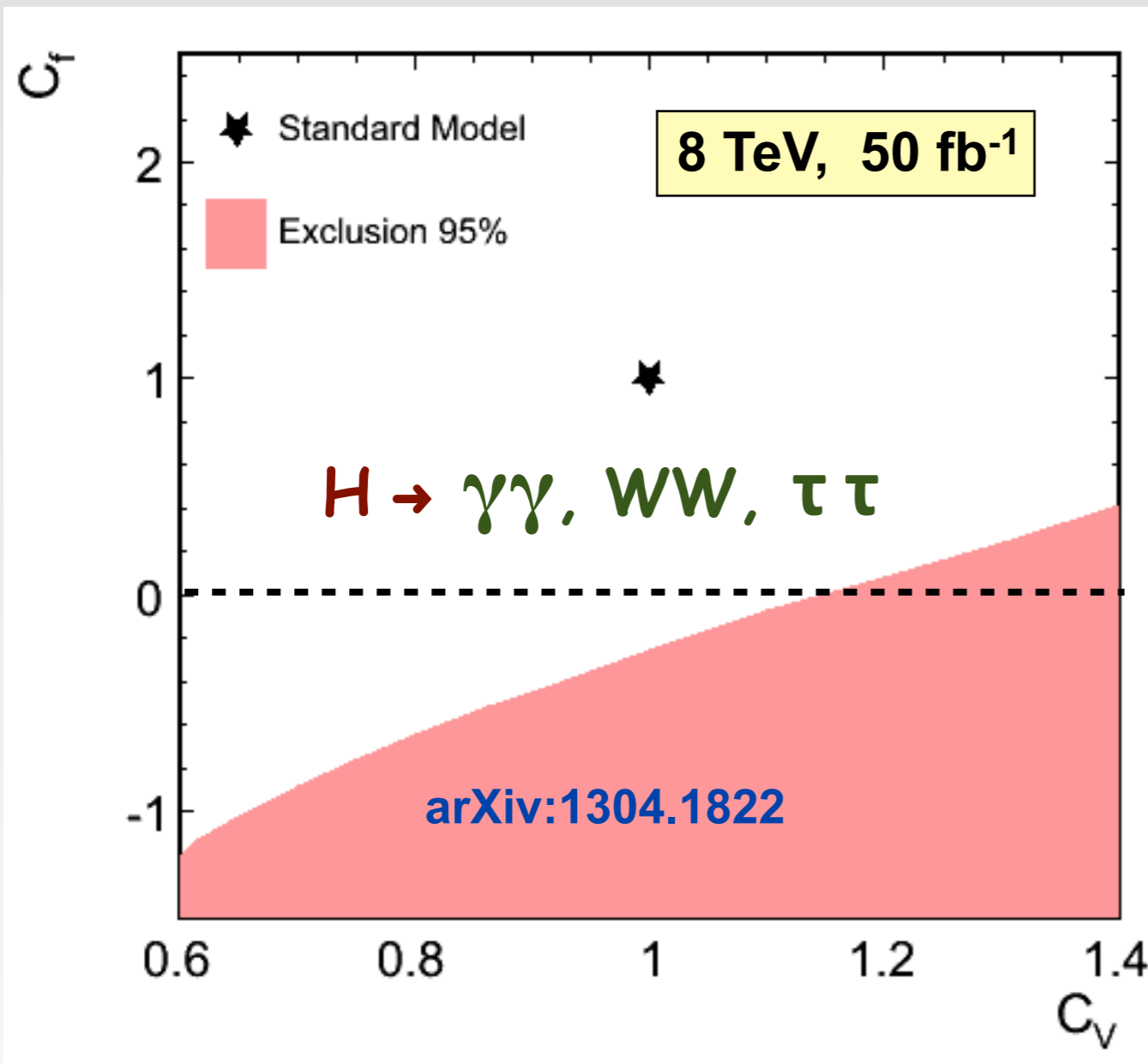
2 benchmark scenarios :

● universal $C_f (=C_t)$ \longrightarrow

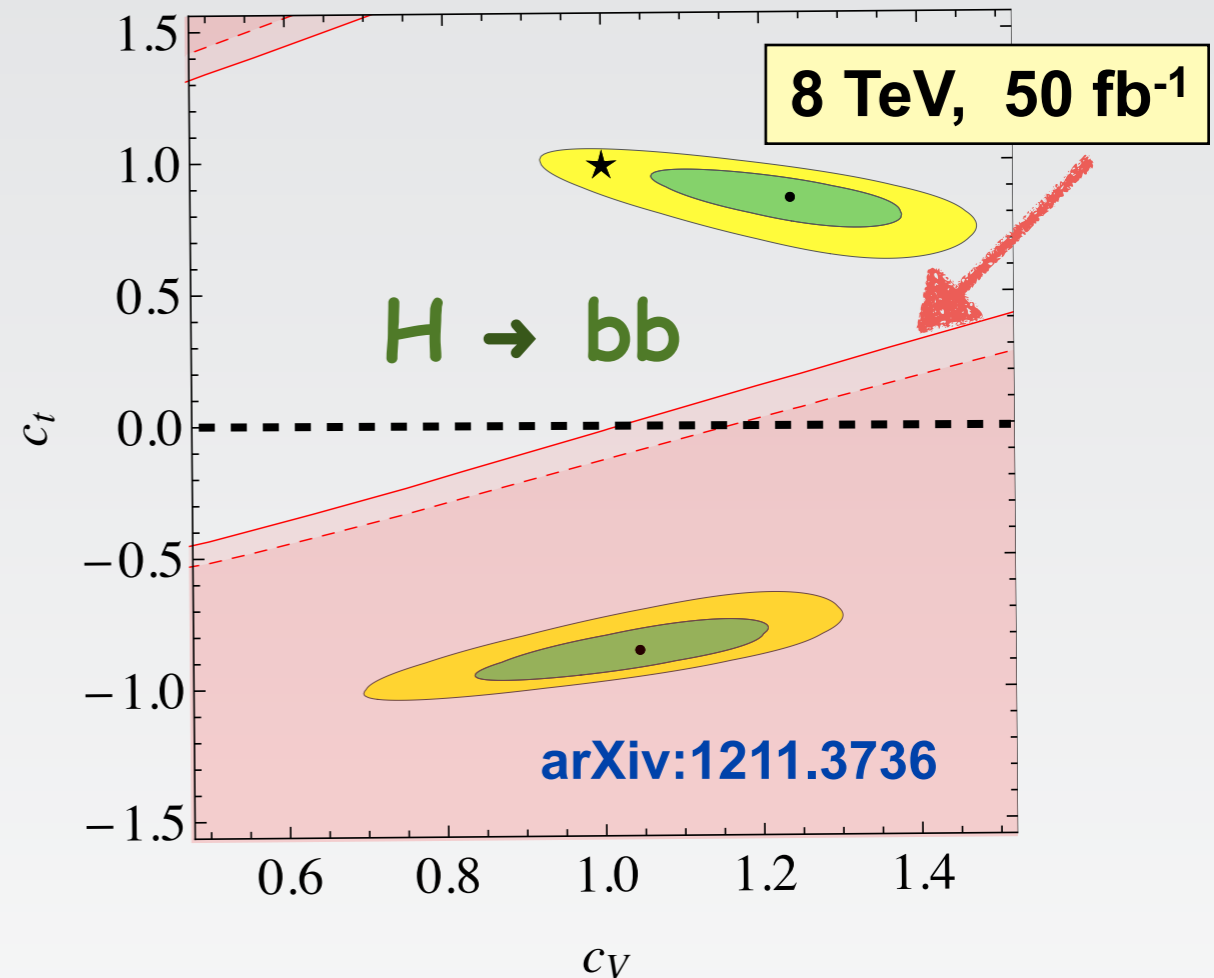
● free C_t , and SM-like $C_b=C_c=C_t=1$



exclusion potential in (C_V, C_T) plane at 8 TeV



free C_T (with $C_{f\neq t} = 1$) scenario

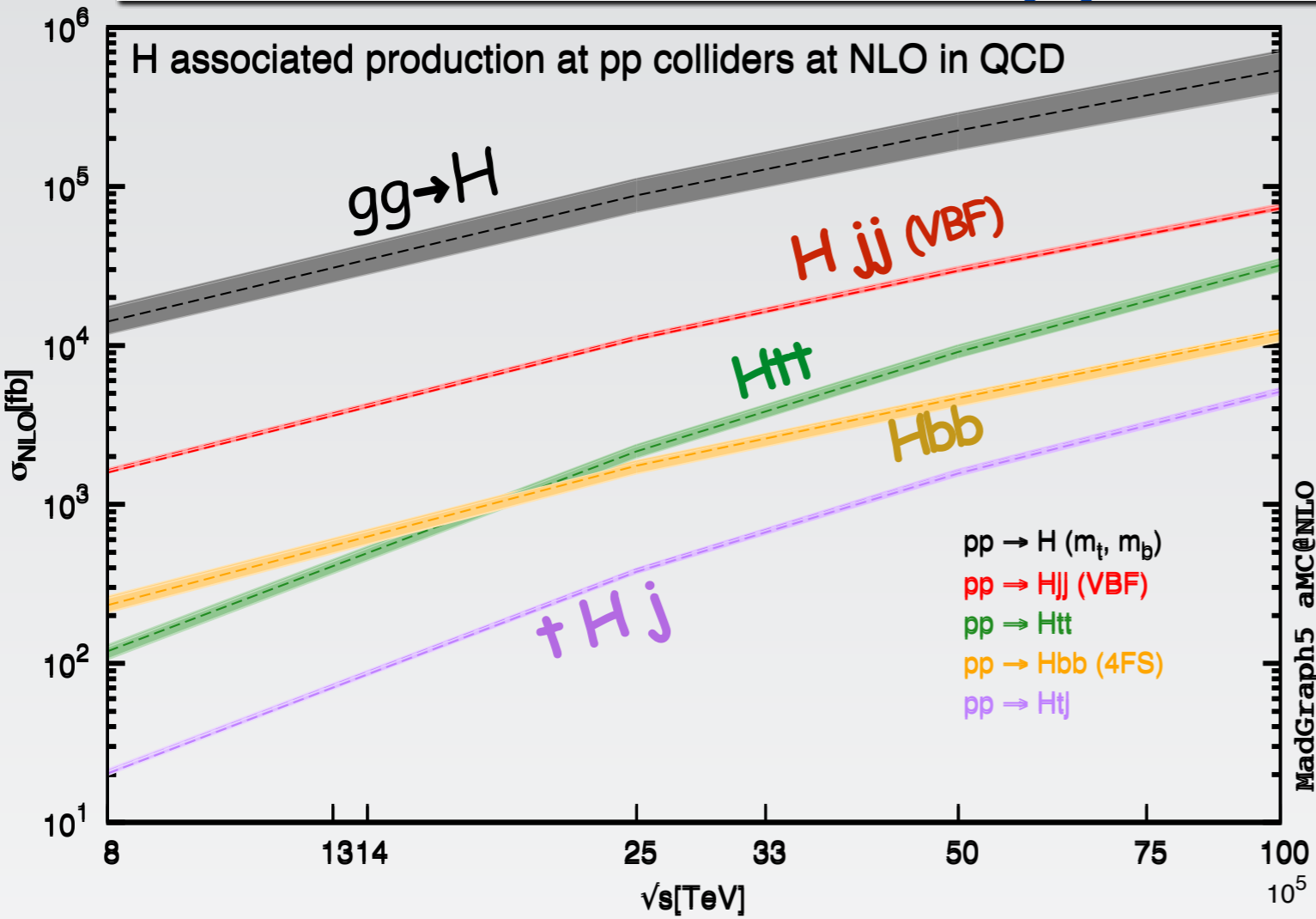


can now be compared with real analyzes ! (see next talk...)

- $H \rightarrow \gamma\gamma$ CMS-PAS-HIG-14-001; (ATLAS) CERN-PH-EP-2014-179
- $H \rightarrow bb$ CMS-PAS-HIG-14-015

present data-set not yet sensitive to $C_T \sim -1$

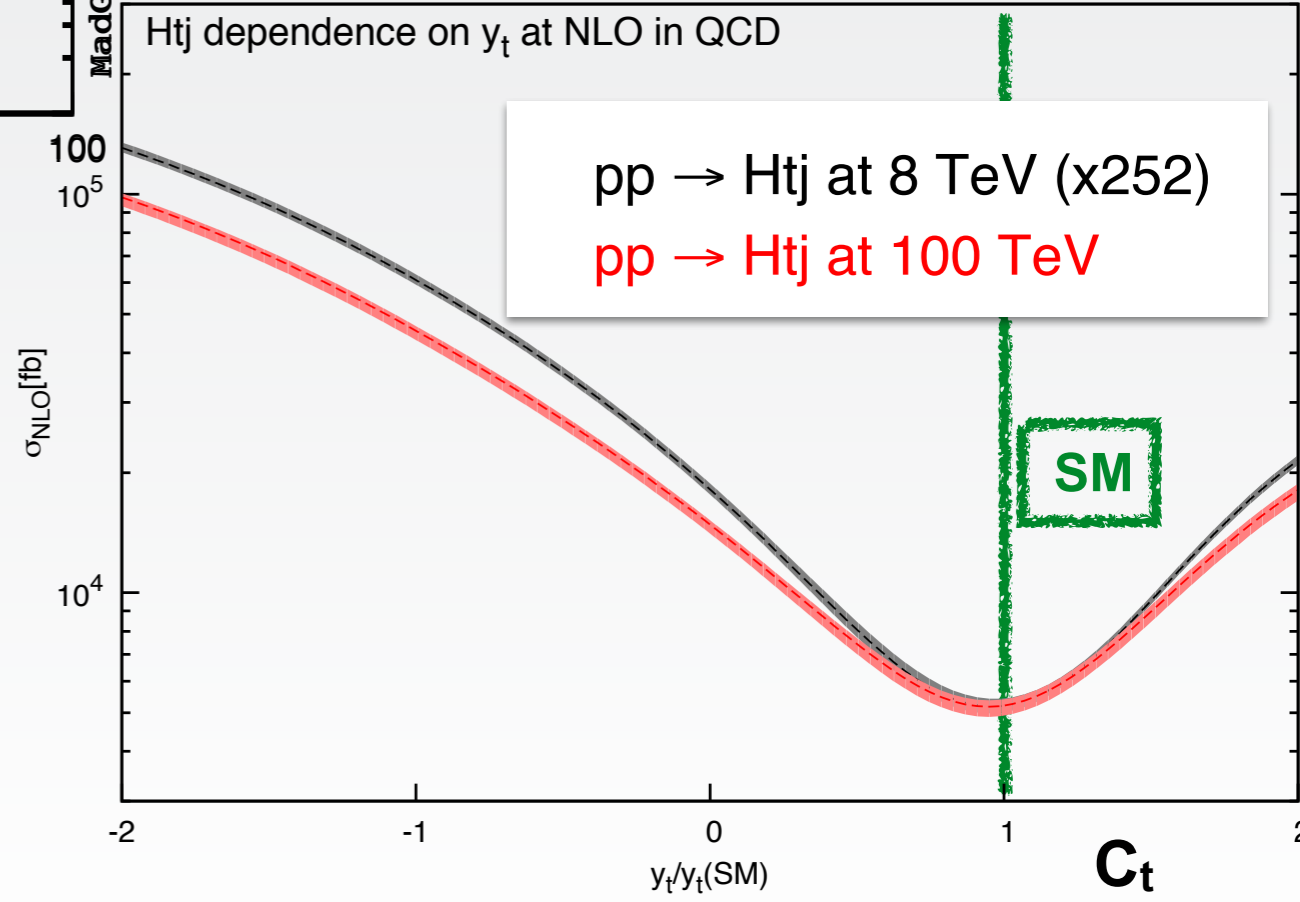
(NLO) SM $\sigma(pp \rightarrow t H j)$ vs \sqrt{s}



$\sigma(H tt)$ and $\sigma(t H j)$ have similar (fast) increase vs E_{cm}

MadGraph5_aMC@NLO

tHJ sensitivity to C_t



(NLO) SM $\sigma(pp \rightarrow t H j)$ vs \sqrt{S}

$\sigma(100\text{TeV}) / \sigma(8\text{TeV})$

Process	$\sigma_{\text{NLO}}(8 \text{ TeV})$ [fb]	$\sigma_{\text{NLO}}(100 \text{ TeV})$ [fb]	ρ
$pp \rightarrow H (m_t, m_b)$	$1.44 \cdot 10^4$ $\begin{matrix} +20\% & +1\% \\ -16\% & -2\% \end{matrix}$	$5.46 \cdot 10^5$ $\begin{matrix} +28\% & +2\% \\ -27\% & -2\% \end{matrix}$	38
$pp \rightarrow H j j$ (VBF)	$1.61 \cdot 10^3$ $\begin{matrix} +1\% & +2\% \\ -0\% & -2\% \end{matrix}$	$7.40 \cdot 10^4$ $\begin{matrix} +3\% & +2\% \\ -2\% & -1\% \end{matrix}$	46
$pp \rightarrow H t \bar{t}$	$1.21 \cdot 10^2$ $\begin{matrix} +5\% & +3\% \\ -9\% & -3\% \end{matrix}$	$3.25 \cdot 10^4$ $\begin{matrix} +7\% & +1\% \\ -8\% & -1\% \end{matrix}$	269
$pp \rightarrow H b \bar{b}$ (4FS)	$2.37 \cdot 10^2$ $\begin{matrix} +9\% & +2\% \\ -9\% & -2\% \end{matrix}$	$1.21 \cdot 10^4$ $\begin{matrix} +2\% & +2\% \\ -10\% & -2\% \end{matrix}$	51
$pp \rightarrow H t j$	$2.07 \cdot 10^1$ $\begin{matrix} +2\% & +2\% \\ -1\% & -2\% \end{matrix}$	$5.21 \cdot 10^3$ $\begin{matrix} +3\% & +1\% \\ -5\% & -1\% \end{matrix}$	252
$pp \rightarrow H W^\pm$	$7.31 \cdot 10^2$ $\begin{matrix} +2\% & +2\% \\ -1\% & -2\% \end{matrix}$	$1.54 \cdot 10^4$ $\begin{matrix} +5\% & +2\% \\ -8\% & -2\% \end{matrix}$	21
$pp \rightarrow H Z$	$3.87 \cdot 10^2$ $\begin{matrix} +2\% & +2\% \\ -1\% & -2\% \end{matrix}$	$8.82 \cdot 10^3$ $\begin{matrix} +4\% & +2\% \\ -8\% & -2\% \end{matrix}$	23

Torrielli, arXiv:1407.1623

$$\sqrt{S} = 8 \text{ TeV} \rightarrow \sigma(t H j) \sim 10^{-3} \sigma(H)$$

$$\sqrt{S} = 100 \text{ TeV} \rightarrow \sigma(t H j) \sim 10^{-2} \sigma(H)$$

Part 2:

complex g_{tH} (\rightarrow CP violating)

$$\mathcal{L}_t = -\frac{m_t}{v} (\kappa_t \bar{t}t + i\tilde{\kappa}_t \bar{t}\gamma_5 t) H$$

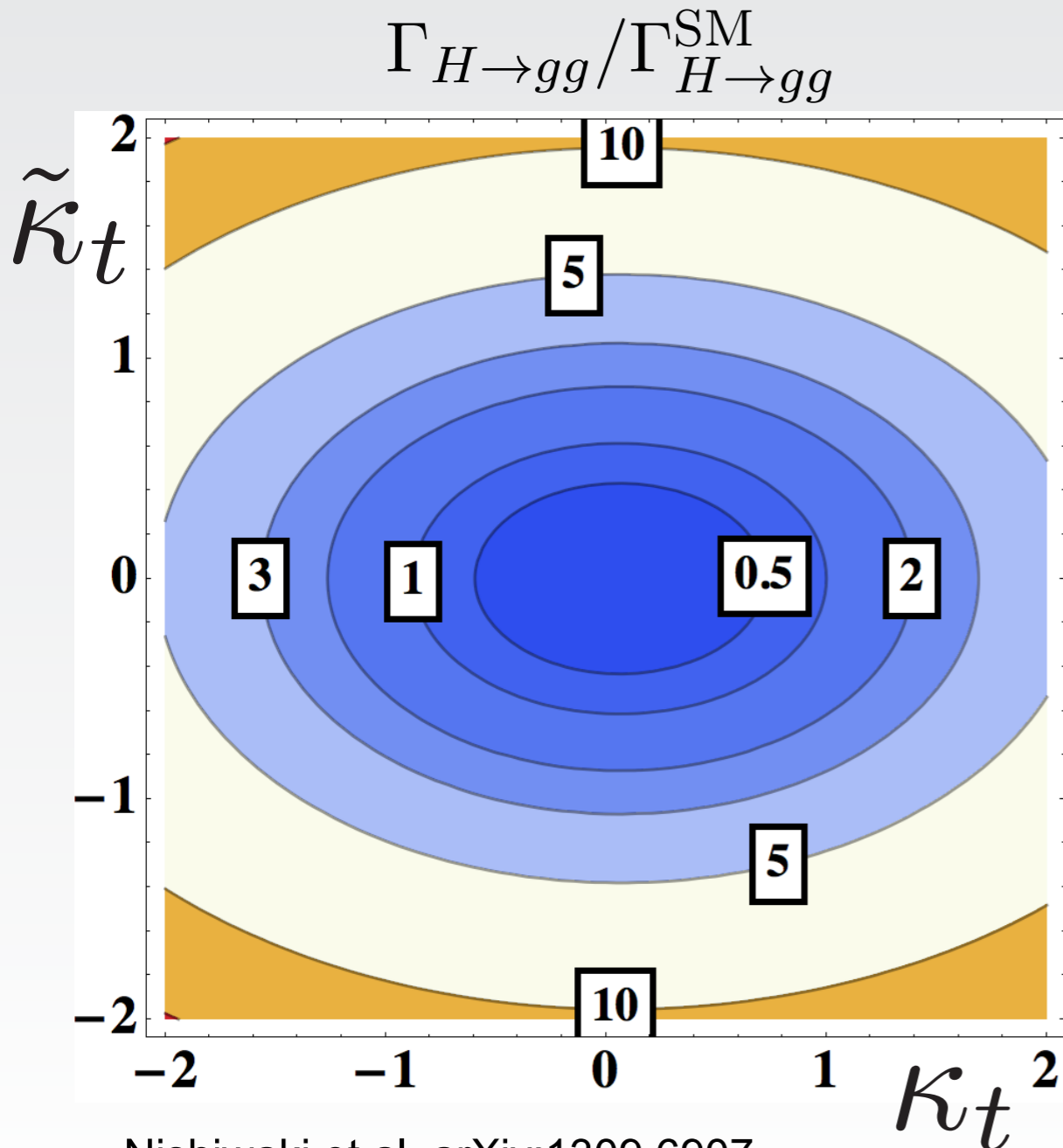
$$\kappa_t \equiv C_t \quad (\tilde{\kappa}_t = 0)$$

great sensitivity in tHj features also
to pseudo-scalar coupling $\tilde{\kappa}_t$!

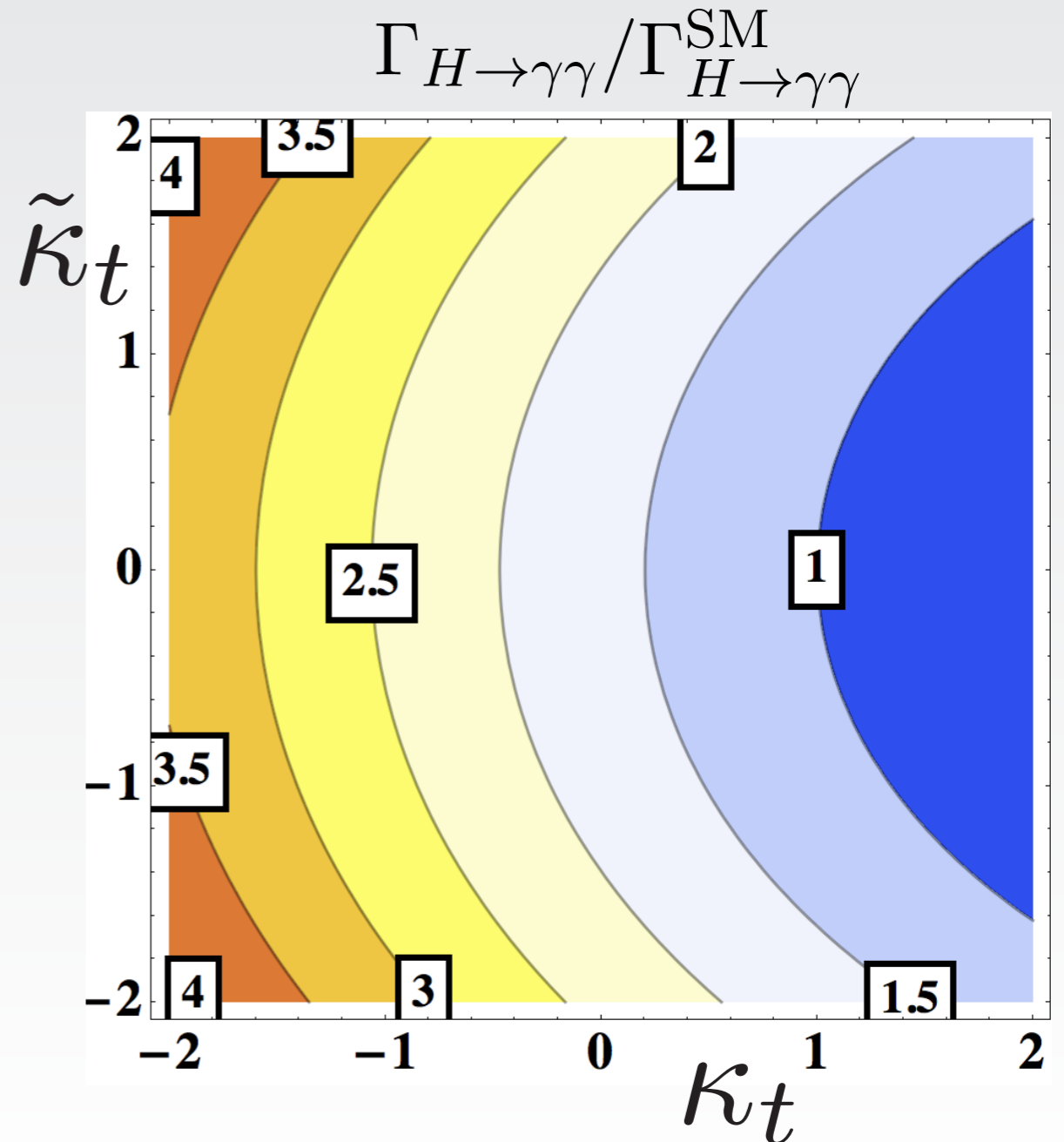
indirect constraints on $(\kappa_t, \tilde{\kappa}_t)$ from C_g and C_γ

$$C_g^2 = \mu_{gg} \simeq \kappa_t^2 + 2.6\tilde{\kappa}_t^2 + 0.11\kappa_t(\kappa_t - 1)$$

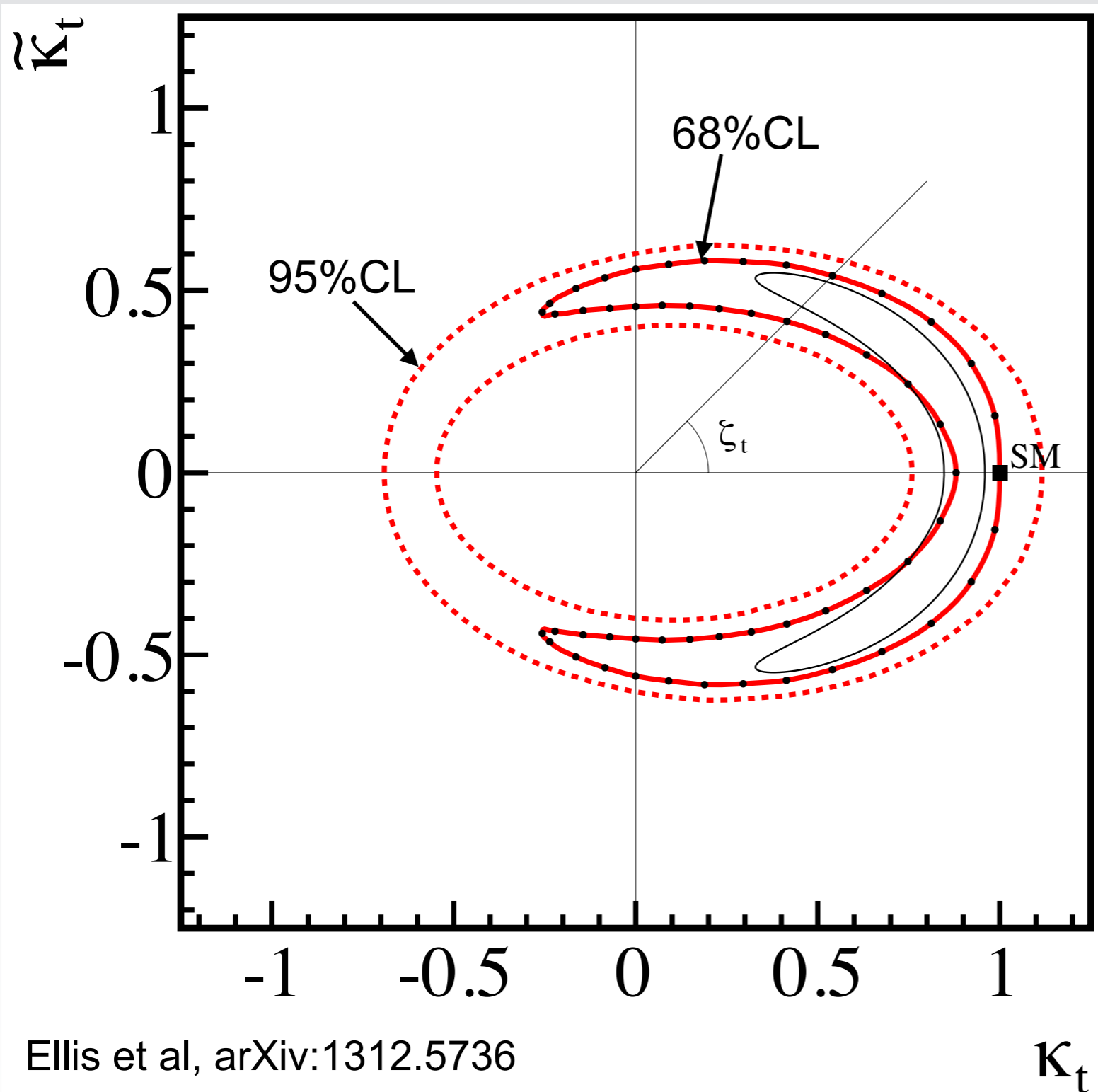
$$C_\gamma^2 = \mu_{\gamma\gamma} \simeq (1.28 - 0.28\kappa_t)^2 + (0.43\tilde{\kappa}_t)^2$$



Nishiwaki et al, arXiv:1309.6907



by combining LHC constrains on C_g and C_γ :



CP-violation
phase in $g_{t\bar{t}H}$:

$$\zeta_t \equiv \arctan \left(\frac{\tilde{\kappa}_t}{\kappa_t} \right)$$

low-energy bounds on
electric dipole moments

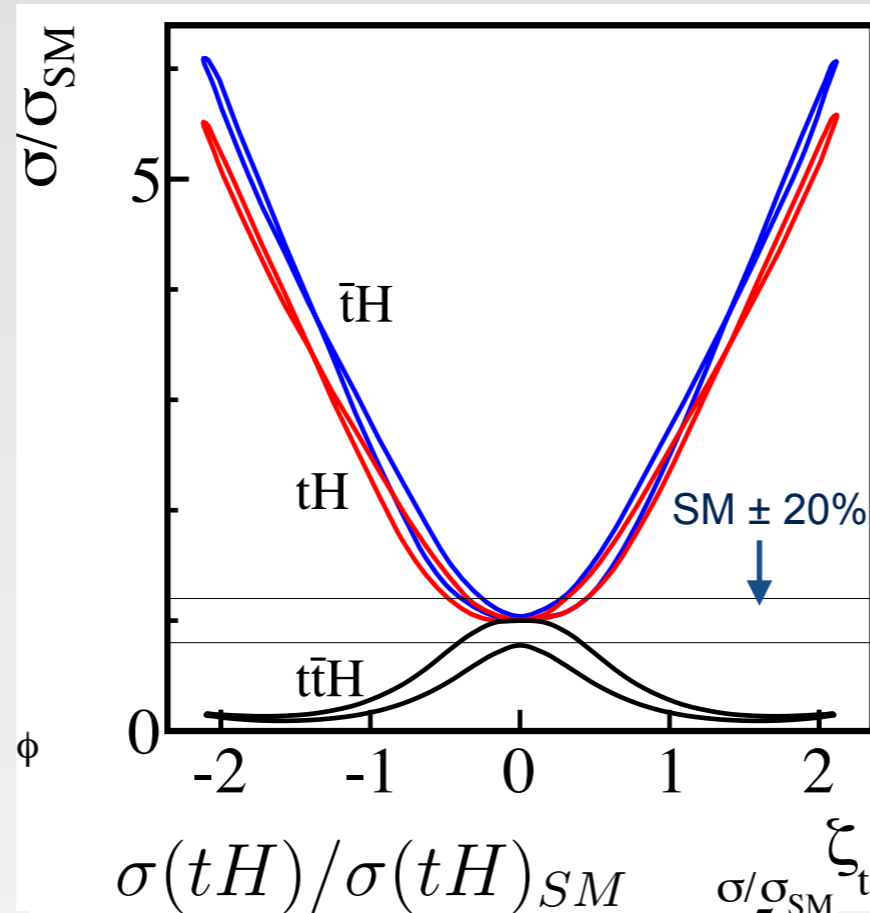
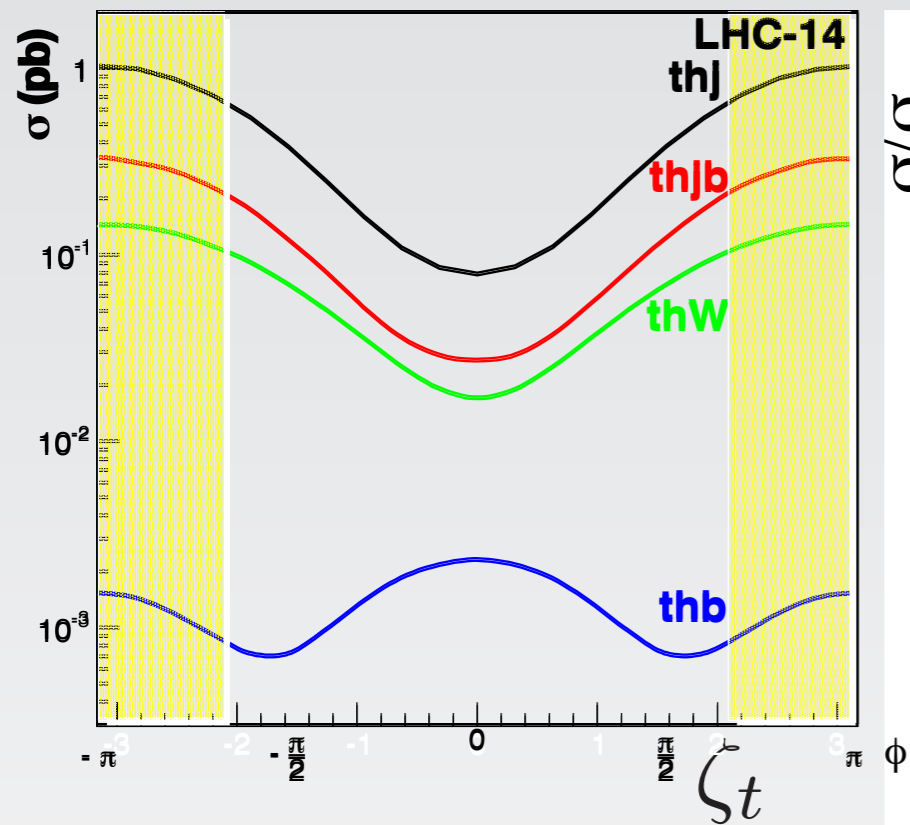
$$|\tilde{\kappa}_t| < 0.01 \quad \text{if} \quad \kappa_e = 1$$

(model dependent !)

Brod et al, arXiv:1310.1385

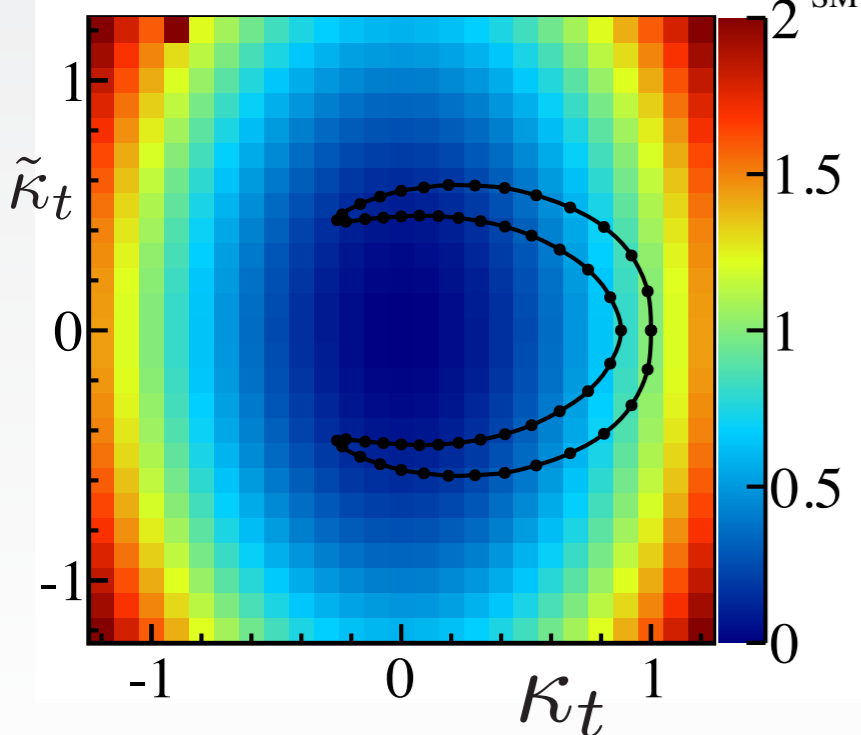
tHj and Htt cross sections versus $(\kappa_t, \tilde{\kappa}_t, \zeta_t)$

Ellis et al, arXiv:1312.5736

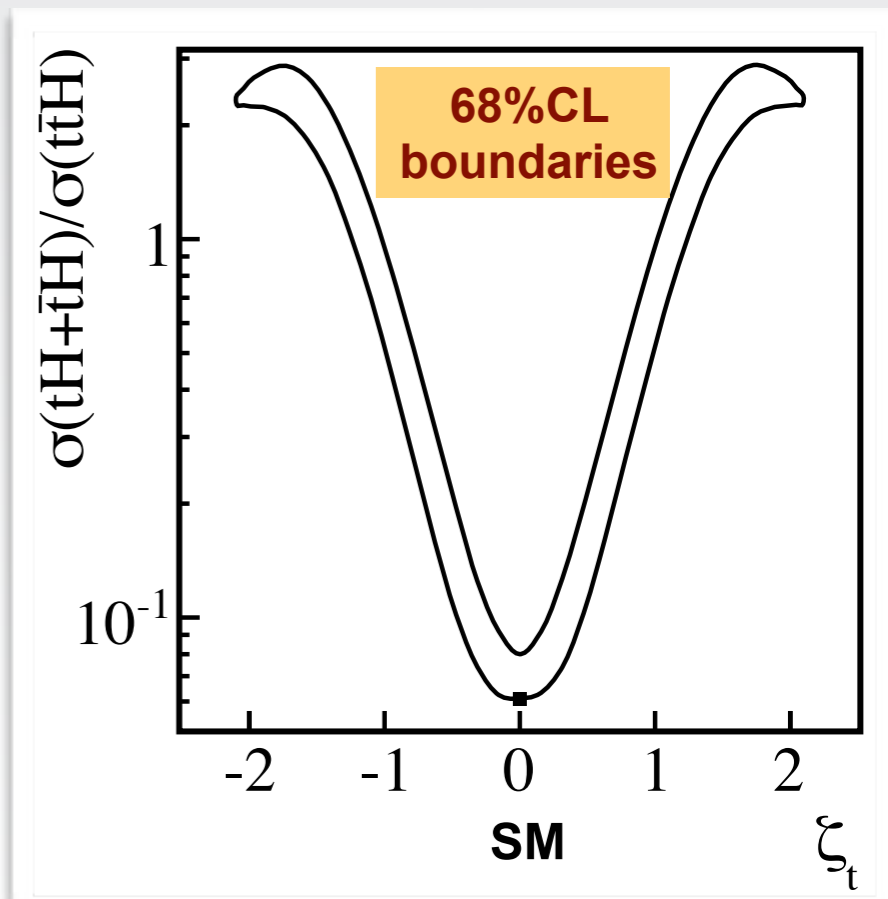
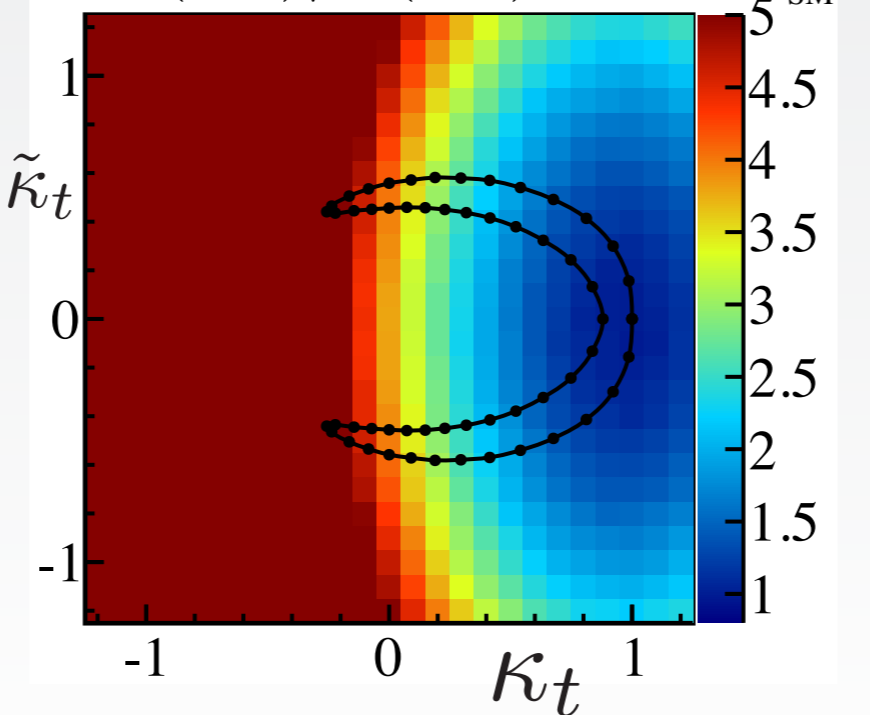


$\sigma(tH)/\sigma(ttH)$
increased
up to ~ 20

$\sigma(\bar{t}tH)/\sigma(\bar{t}tH)_{SM}$



$\sigma(tH)/\sigma(tH)_{SM}$



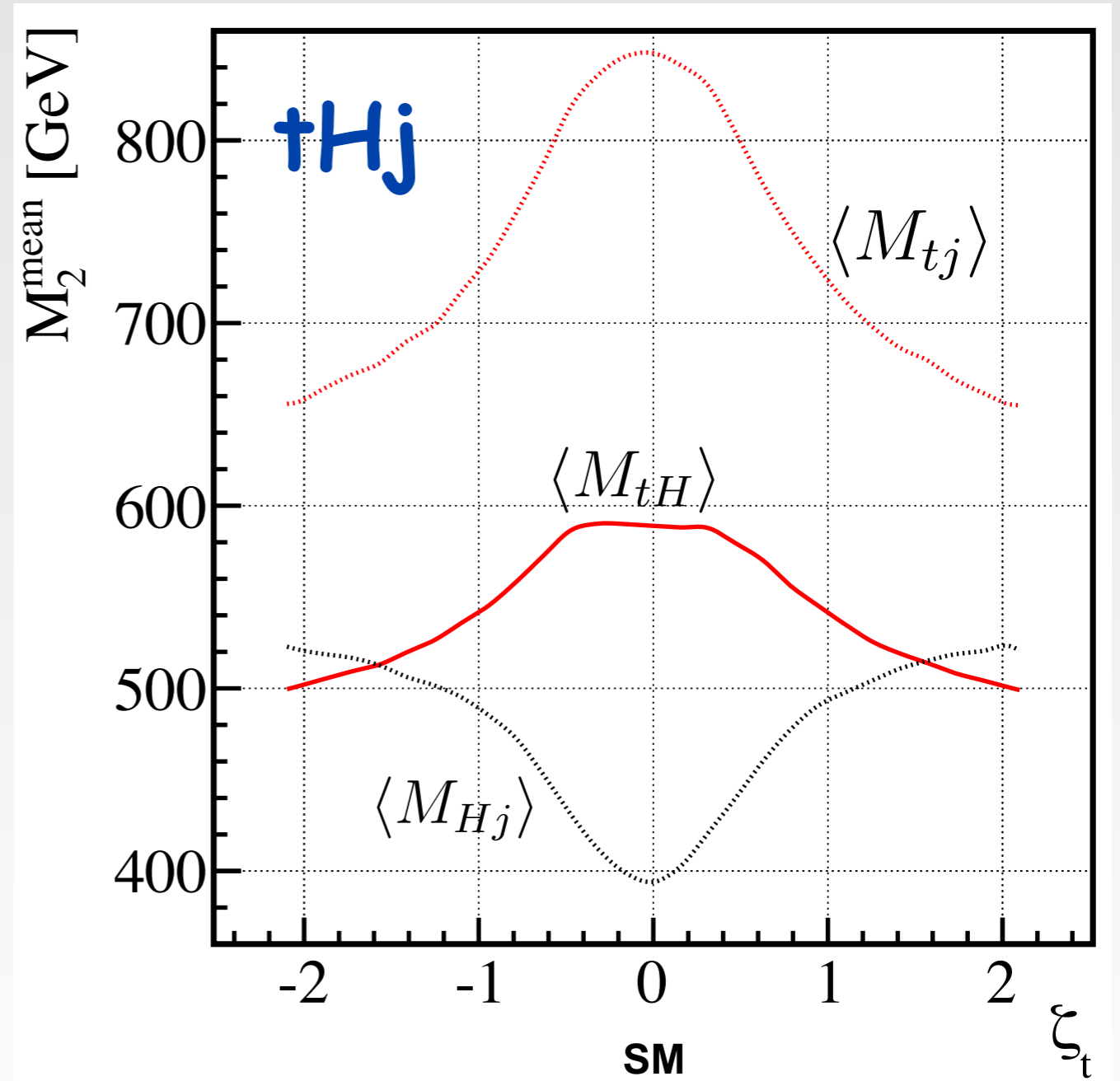
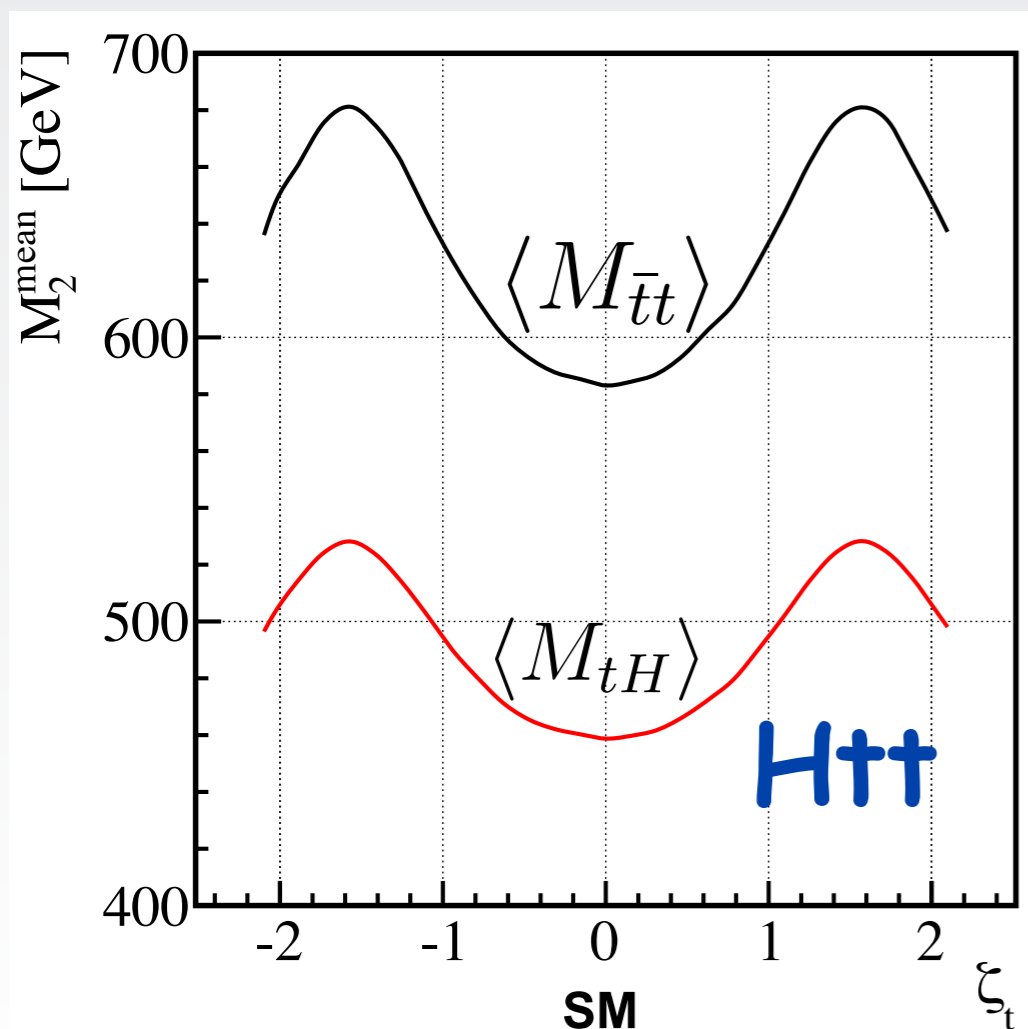
$\sigma(tH)/\sigma(ttH)|_{SM} \sim 0.06$

(see also Khatibi Najafabadi,
arXiv:1409.6553)

tHj and Htt invariant-mass distributions

(along a contour passing through the middle of 68%CL crescent-shape allowed region)

sensitive variables
for all sub-systems !



lepton angle in $[t(\rightarrow \ell vb) H j]$ and $|\tilde{\kappa}_t/\kappa_t|$

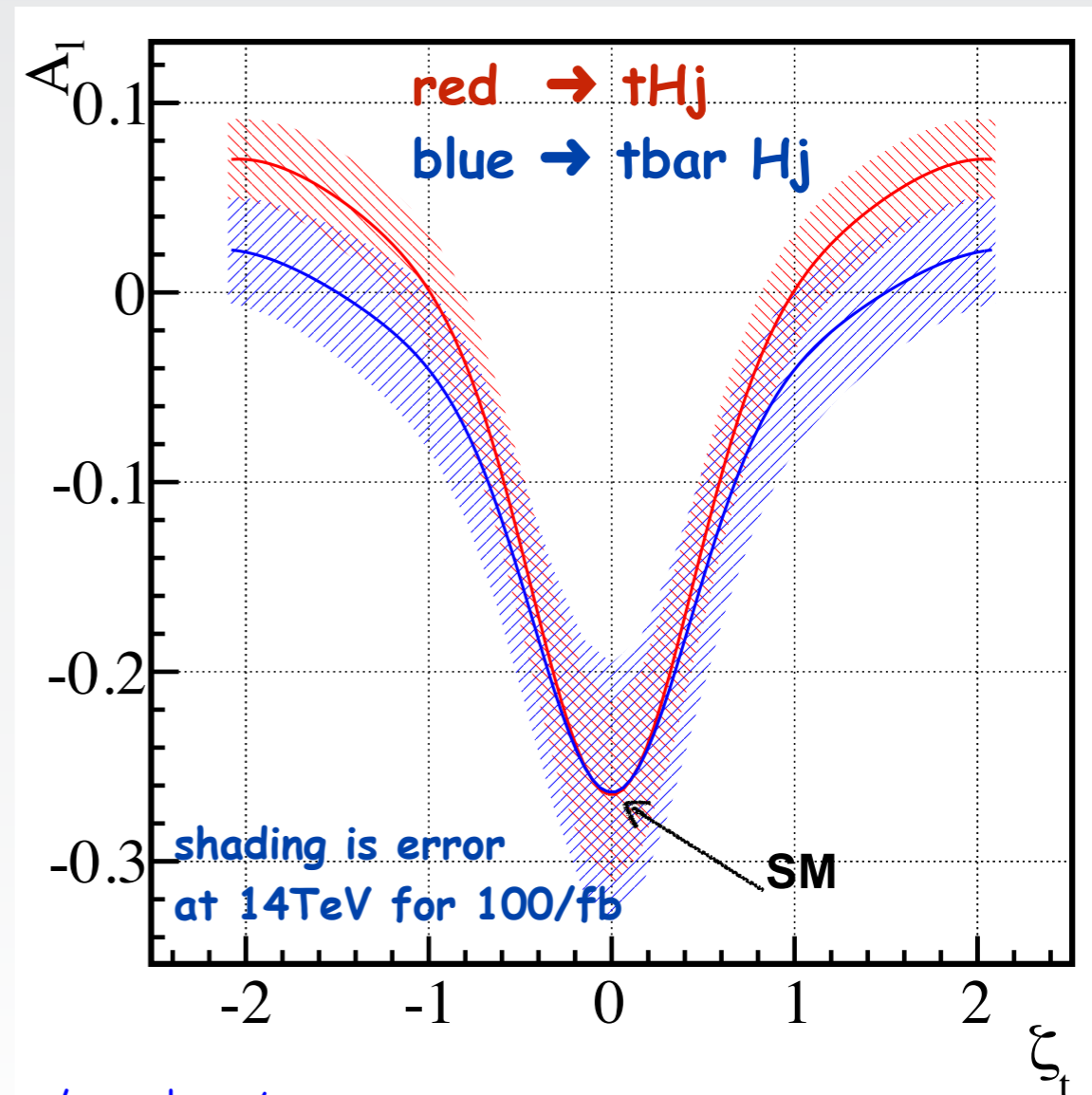
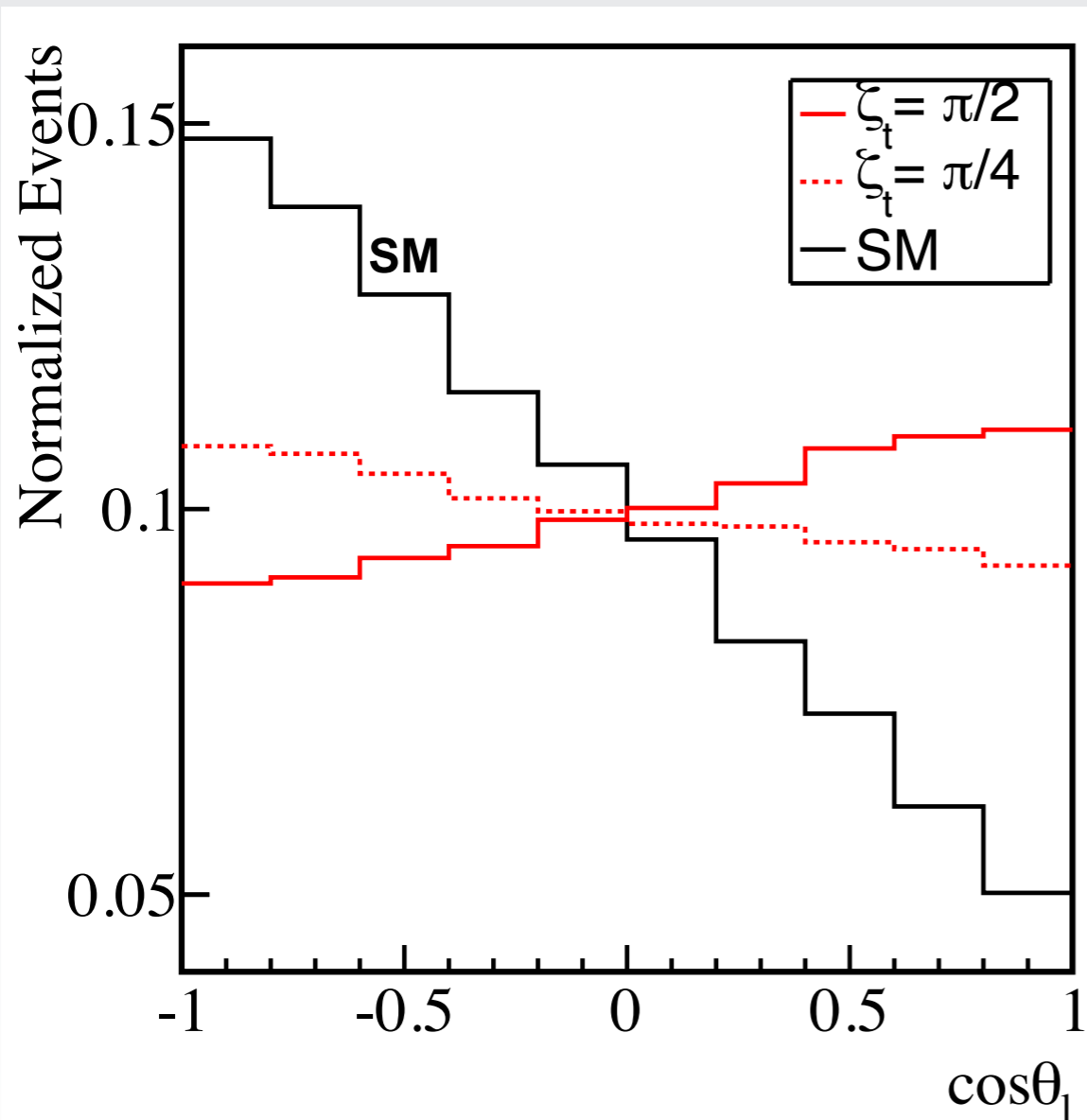
$$\frac{1}{\Gamma_f} \frac{d\Gamma_f}{d\cos\theta_f} = \frac{1}{2}(1 + \omega_f P_t \cos\theta_f)$$

top spin quantization axis
in top rest frame
(angle wrt top boost direction)

top in tHj highly polarized in the SM

large $P_t = \frac{N(\uparrow) - N(\downarrow)}{N(\uparrow) + N(\downarrow)} \rightarrow$ large asymmetry

$$A_\ell = \frac{N(\cos\theta_\ell > 0) - N(\cos\theta_\ell < 0)}{N(\cos\theta_\ell > 0) + N(\cos\theta_\ell < 0)}$$



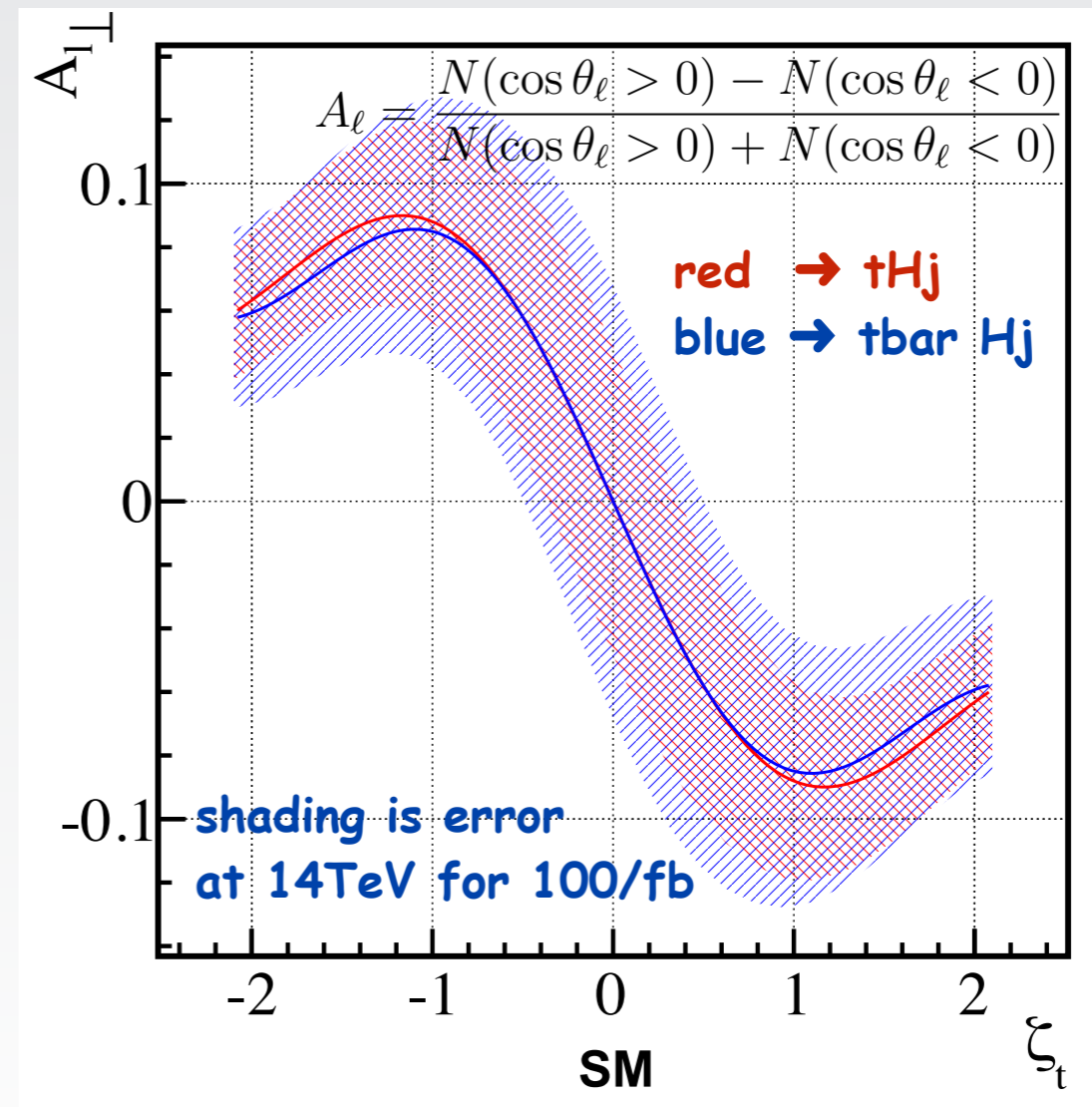
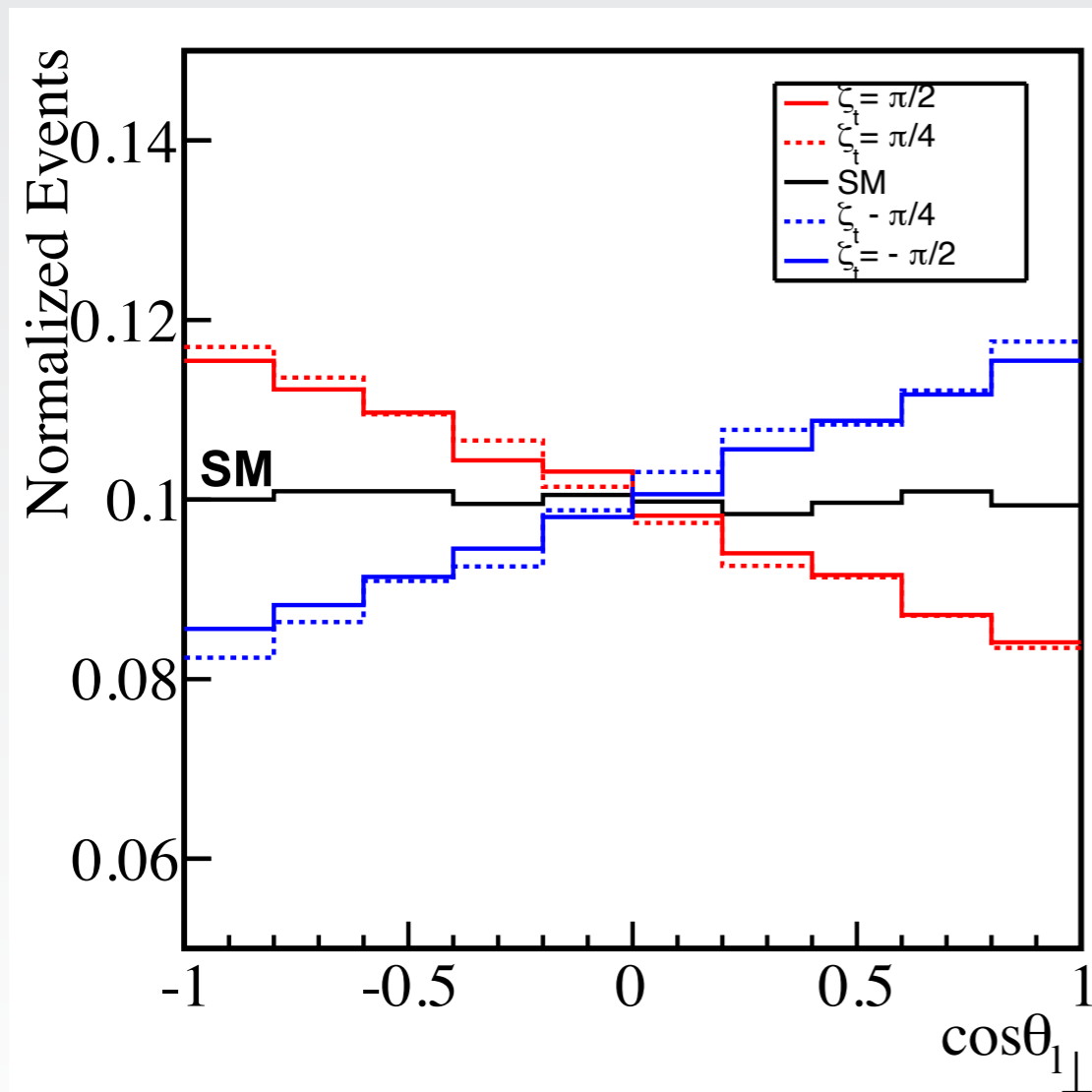
$|\tilde{\kappa}_t/\kappa_t| \neq 0$ decreases Asymmetry!

$t(\rightarrow \ell vb) H j$ and sign of $\tilde{\kappa}_t/\kappa_t$

lepton decay angle out of tHj production plane sensitive to sign of $\tilde{\kappa}_t/\kappa_t$

$\vec{p}_j \times \vec{p}_H \rightarrow$ top spin quantization axis perpendicular to production plane in top rest frame

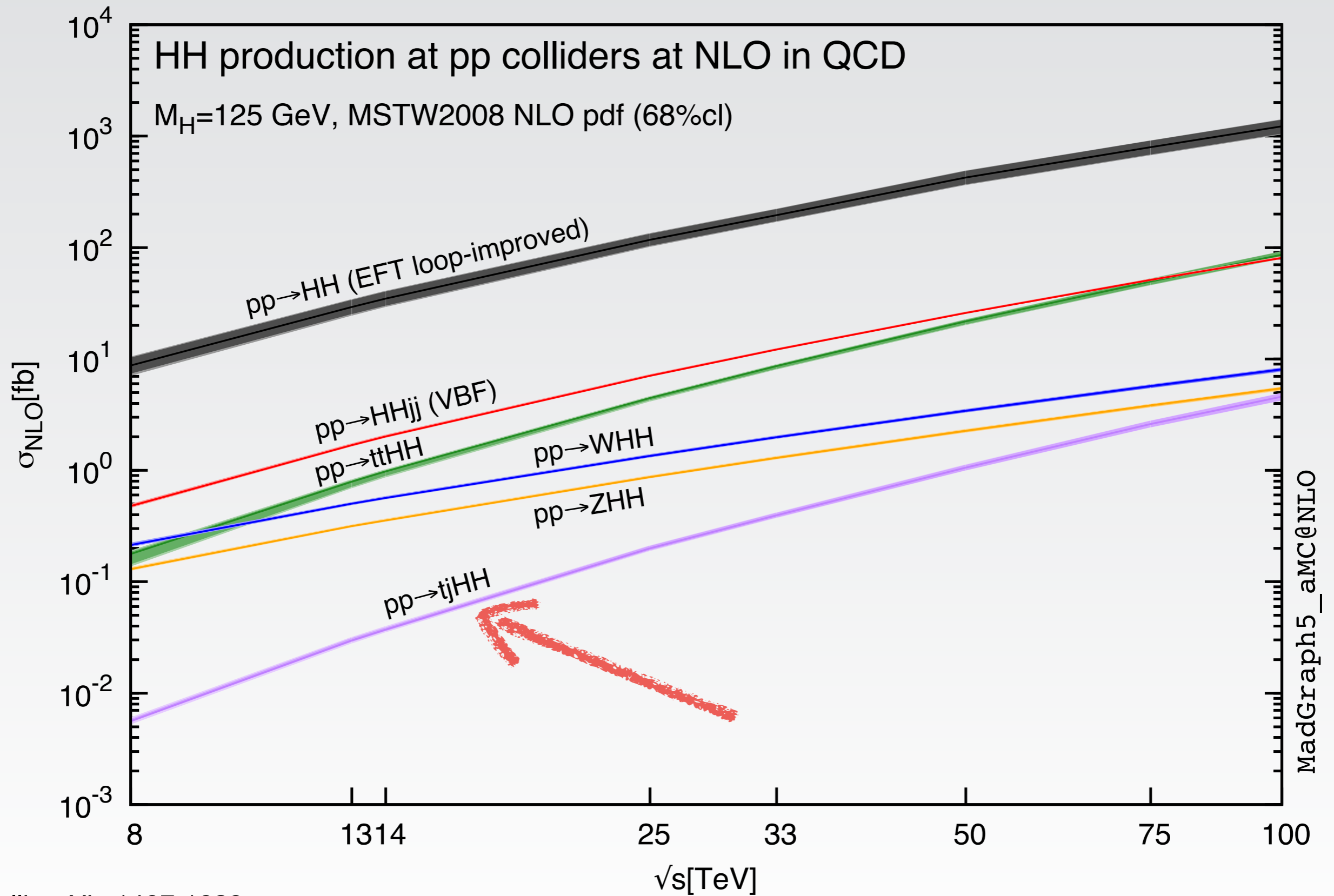
$$\frac{1}{\Gamma_f} \frac{d\Gamma_f}{d \cos \theta_f} = \frac{1}{2} (1 + \omega_f P_t \cos \theta_f)$$



Ellis et al, arXiv:1312.5736

just a look to the far future of $t+H$...

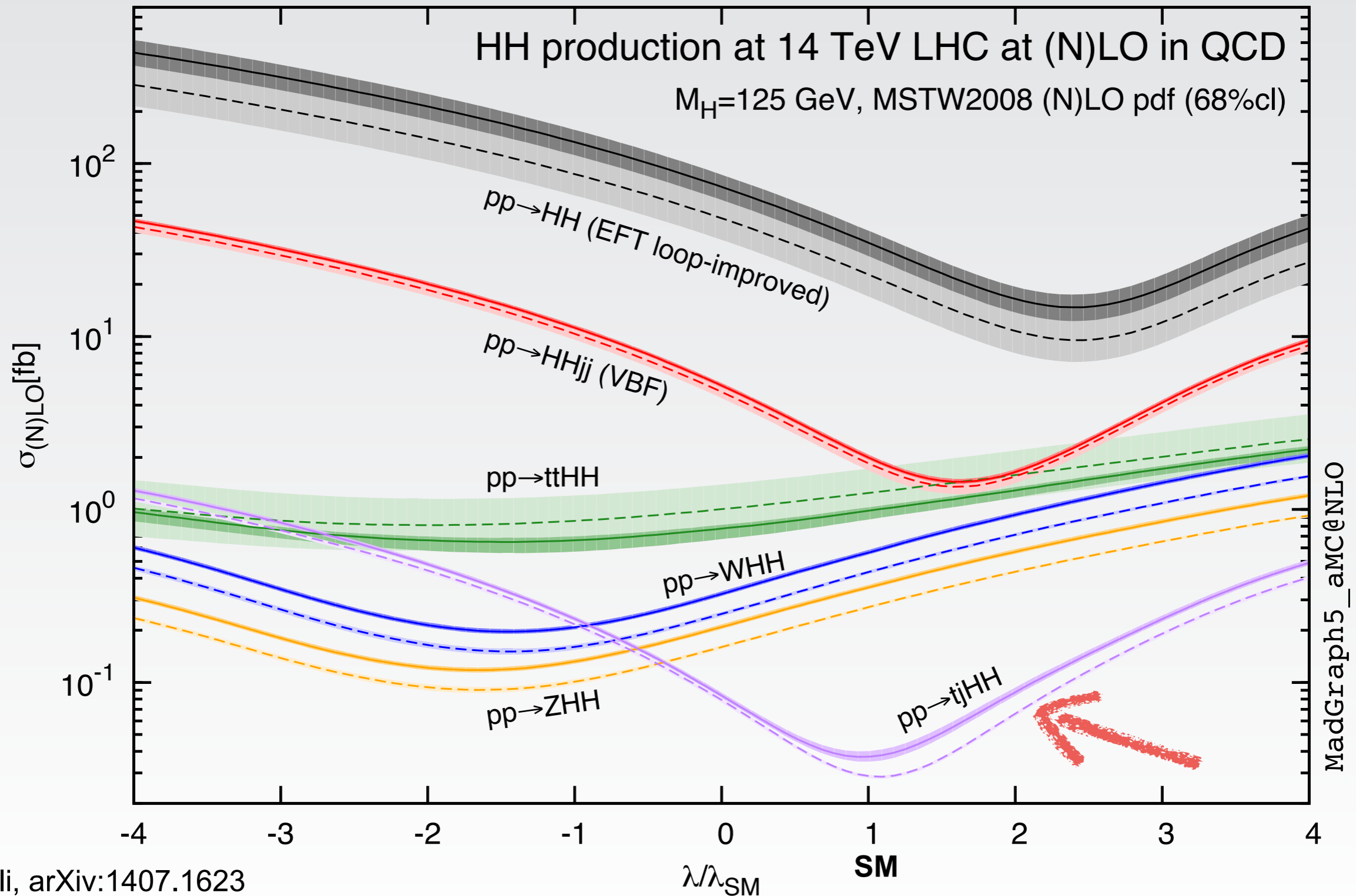
$t+HH$ production



Torrielli, arXiv:1407.1623

nice sensitivity to trilinear coupling λH^3

$$[\lambda^{\text{SM}} = (m_H/v)^2/2 = 0.13]$$



Torrielli, arXiv:1407.1623

Outlook

- $g_{H\bar{t}t}$ coupling tightly connected to **EWSB** mechanism
- crucial to probe it in a **direct** way
- $p p \rightarrow H \bar{t} t$ major role, but **not sensitive to $g_{H\bar{t}t}$ sign**
- $p p \rightarrow \bar{t} H q$ production excellent test of **$g_{H\bar{t}t}$ sign**
- $p p \rightarrow \bar{t} H q$ also sensitive to **pseudo-scalar $g_{H\bar{t}t}$ components** (in a complementary way to **$H\bar{t}t$ production**);
kinematical distributions and asymmetries crucial !
- **SM tHq cross sections** hard to probe in near future
- **13-TeV run** expected to reach sensitivity needed for testing the **$C_t \sim -1$** hypothesis through $p p \rightarrow \bar{t} H q$!