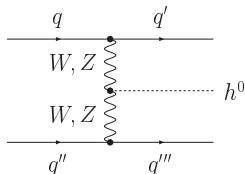


New-physics effects in (VH and) VBF-Higgs production

Michael Rauch | LHC HXSWG 8th meeting, Jan 2015

INSTITUTE FOR THEORETICAL PHYSICS





- mediated by HWW and HZZ couplings

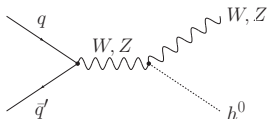
VBF: [Han, Valencia, Willenbrock; Figy, Oleari, Zeppenfeld; Campbell, Ellis, Berger]

- clear signature due to two tagging jets
- QCD corrections relatively small $\sim 5\%$

[Figy, Oleari, Zeppenfeld; Harlander et al.; Bolzoni et al.]

- EW corrections of same size

[Ciccolinni, Denner, Dittmaier; Figy, Palmer, Weiglein]



VH:

- similar matrix elements
- boosted H in exp. analyses

Two dedicated tools for NLO QCD + anomalous couplings:

- VBFNLO

VBF in VBF approx., WH(j)

[Zeppenfeld, MR et al.]

- HAWK

VBF, VH including interferences

[Denner, Dittmaier, Kallweit, Mück]

NLO EW in both cases treated additively (no EW corrections to anomalous couplings)

Effective Field Theories

General structure of $HV_1^\mu V_2^\nu$ vertex ($V = W, Z, \gamma, g$)

$$T^{\mu\nu} = a_1(p_1, p_2)g^{\mu\nu} + a_2(p_1, p_2)[p_1 \cdot p_2 g^{\mu\nu} - p_2^\mu p_1^\nu] + a_3(p_1, p_2)\varepsilon^{\mu\nu\rho\sigma} p_{1\rho} p_{2\sigma}$$

SM (LO): $HWW, HZZ \propto a_1, a_2 = a_3 = 0$

$$H\gamma\gamma, Hgg \propto a_2, a_1 = a_3 = 0$$

Alternative parametrization given by dimension-6 operators of effective Lagrangian

$$\mathcal{L}_{eff} = \frac{f_{WW}}{\Lambda_6^2} \mathcal{O}_{WW} + \frac{f_{BB}}{\Lambda_6^2} \mathcal{O}_{BB} + \frac{f_W}{\Lambda_6^2} \mathcal{O}_W + \frac{f_B}{\Lambda_6^2} \mathcal{O}_B + \text{CP-odd part}$$

$$= \frac{eM_W}{\Lambda_6^2} \left(\frac{1}{2 \cos^2 \theta_W} (\cos^2 \theta_W f_W + \sin^2 \theta_W f_B) Z_{\mu\nu} Z^\mu \partial^\nu H - \frac{1}{2 \cos^2 \theta_W} (\sin^4 \theta_W f_{BB} + \cos^4 \theta_W f_{WW}) HZ_{\mu\nu} Z^{\mu\nu} + \frac{1}{2} f_W (W_{\mu\nu}^+ W_-^\mu + W_{\mu\nu}^- W_+^\mu) \partial^\nu H - f_{WW} H W_{\mu\nu}^+ W_-^{\mu\nu} \right) + H\gamma\gamma, HZ\gamma \text{ terms, CP-odd part}$$

Direct correspondence to a_i

$$a_2^{HWW} = \frac{2eM_W}{\sin \theta_W \Lambda_6^2} \left(f_{WW} - \frac{1}{2} f_W \right)$$

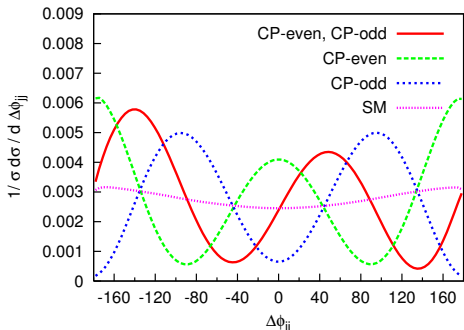
$$a_2^{HZZ} = \frac{2eM_W}{\sin \theta_W \Lambda_6^2} \left(\cos^2 \theta_W f_{WW} + \left(\frac{\sin^2 \theta_W}{\cos^2 \theta_W} - \sin^2 \theta_W \right) f_{BB} - \frac{1}{2} f_W - \frac{\sin^2 \theta_W}{2 \cos^2 \theta_W} f_B \right)$$

(only subset shown here, see <https://www.itp.kit.edu/~vbfnlweb/wiki/doku.php?id=documentation:details:anomaloushiggs> for all relations)

→ for these dimension-6 operators no momentum dependence of a_i

Particular sensitivity given by azimuthal angle difference of the two tagging jets $\Delta\phi_{jj}$

[Hankele, Klämke, Zeppenfeld, Figy]



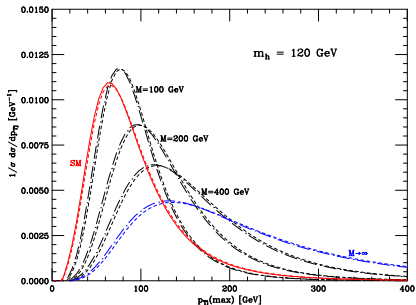
SM:	$a_1 \neq 0$
CP-even:	$a_2 \neq 0$
CP-odd:	$a_3 \neq 0$
CP-even, CP-odd:	$a_2 = a_3 \neq 0$

- almost flat distribution for SM = a_1 -Term
- $\cos(2\Delta\phi_{jj})$ -behaviour for a_2 and a_3 terms
- admixture of a_2 and a_3 can be determined via phase shift of distribution

Physical Observables – $p_T(j)$

New-physics effects can also modify transverse-momentum distributions

$a_1 = 0, a_2 = a_3 = \text{const.}$ [Figy, Zeppenfeld]



- Huge c.s. increase for large p_T
- \rightarrow dominated by (D6-op.)²-term
- \leftrightarrow violates EFT expansion: D8-op. \times SM of same order
- \rightarrow eventually unitarity violation
- \Rightarrow higher-order operators become equally important
- effective parametrization by form factor, e.g.
$$a_i(q_1, q_2) = a_i(0, 0) \left[\left(1 - \frac{q_1^2}{M^2} \right) \left(1 - \frac{q_2^2}{M^2} \right) \right]^{-1}$$
- $\rightarrow p_T$ distrib. can become SM-like

\Rightarrow New-physics effects can modify p_T distributions

\Leftrightarrow Non-observation of modified p_T distribution not sufficient

\Rightarrow need to check other observables as well

\rightarrow When using D6-operators, need to ensure that exclusion is not based purely on large- p_T region