

# Edge-TCT characterization of irradiated HV-CMOSv3 sensors

Christian Gallrapp, Michael Moll, Hannes Neugebauer, Constantin Weisser  
CERN PH-DT-SSD

Marcos Fernández García  
IFCA-CSIC-Universidad de Cantabria,  
visiting scientist at CERN PH-DT-SSD

Daniel Muenstermann  
University of Geneva

**25th RD50 Workshop (CERN)**

19-21 November 2014 CERN

# Outline

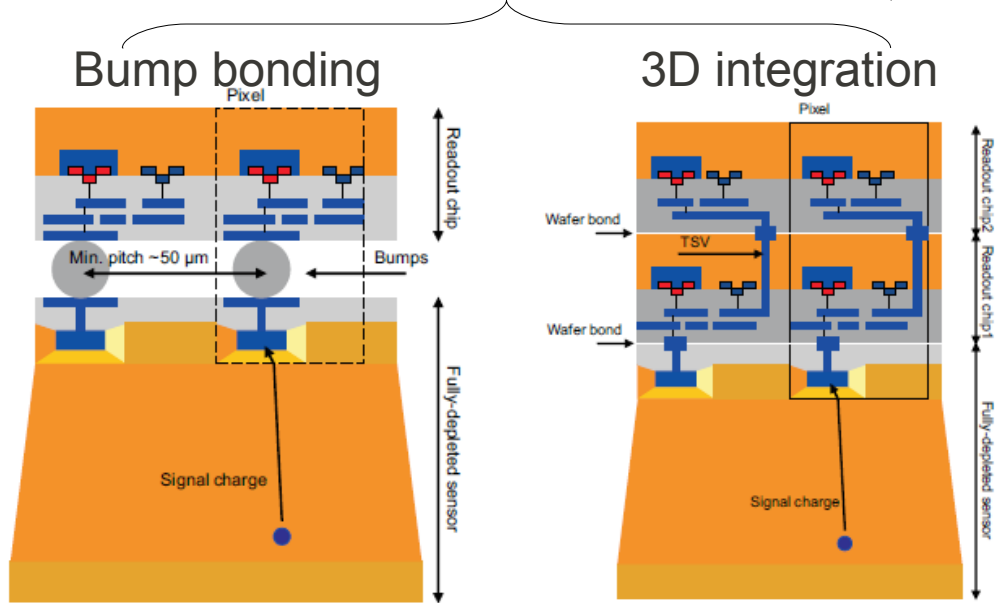
**Introduction and background information**

**Edge-TCT measurements on HVCMOSv3**

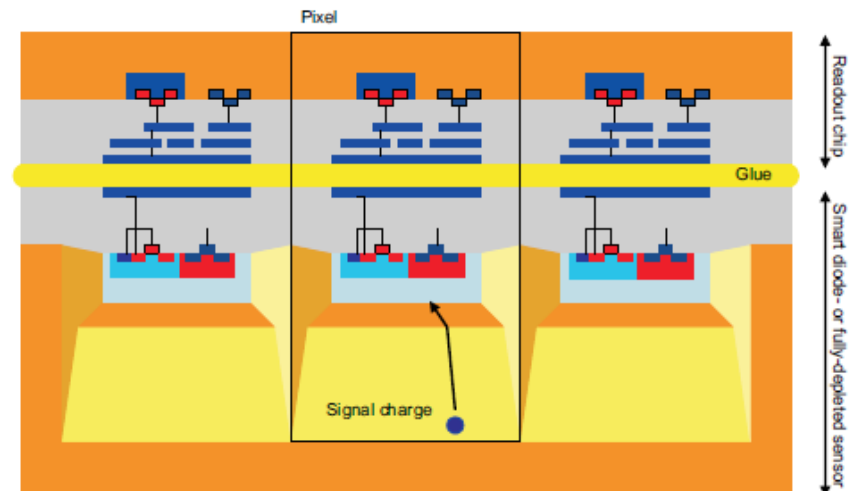
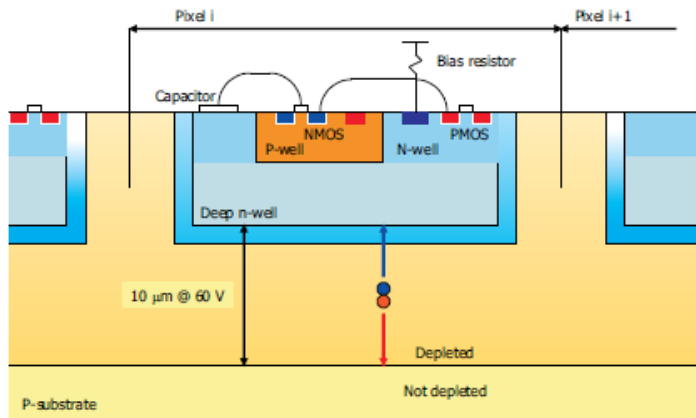
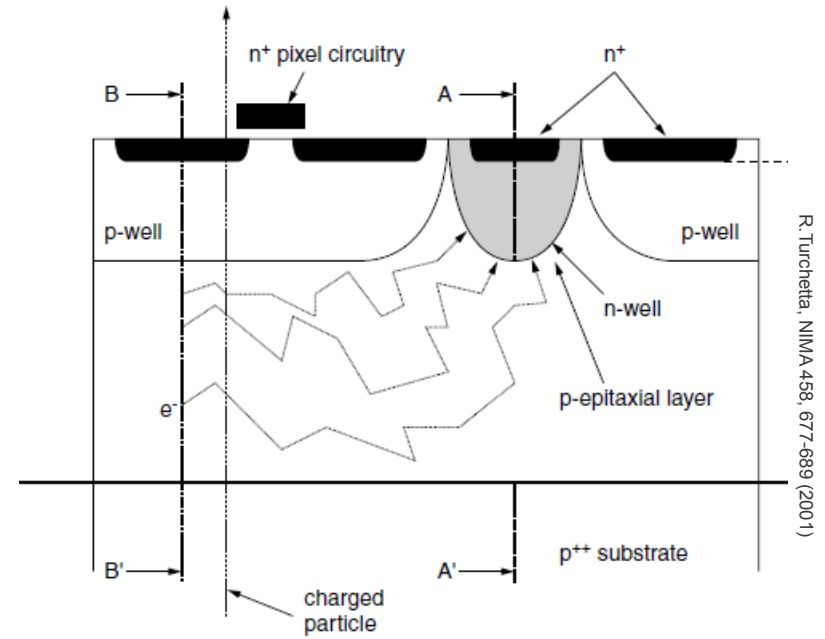
**Radiation hardness up to  $2 \times 10^{16} n_{eq}/cm^2$**

# Introduction to HVCMOS

## Hybrid pixels



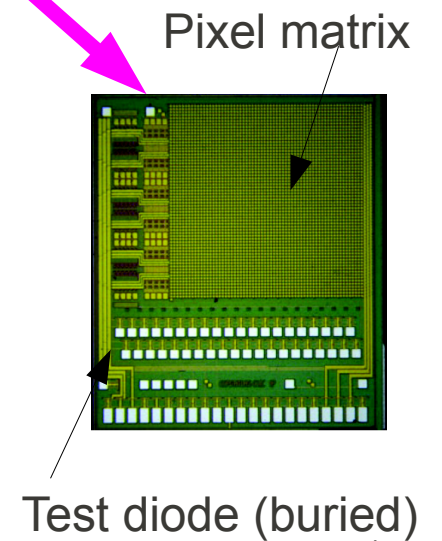
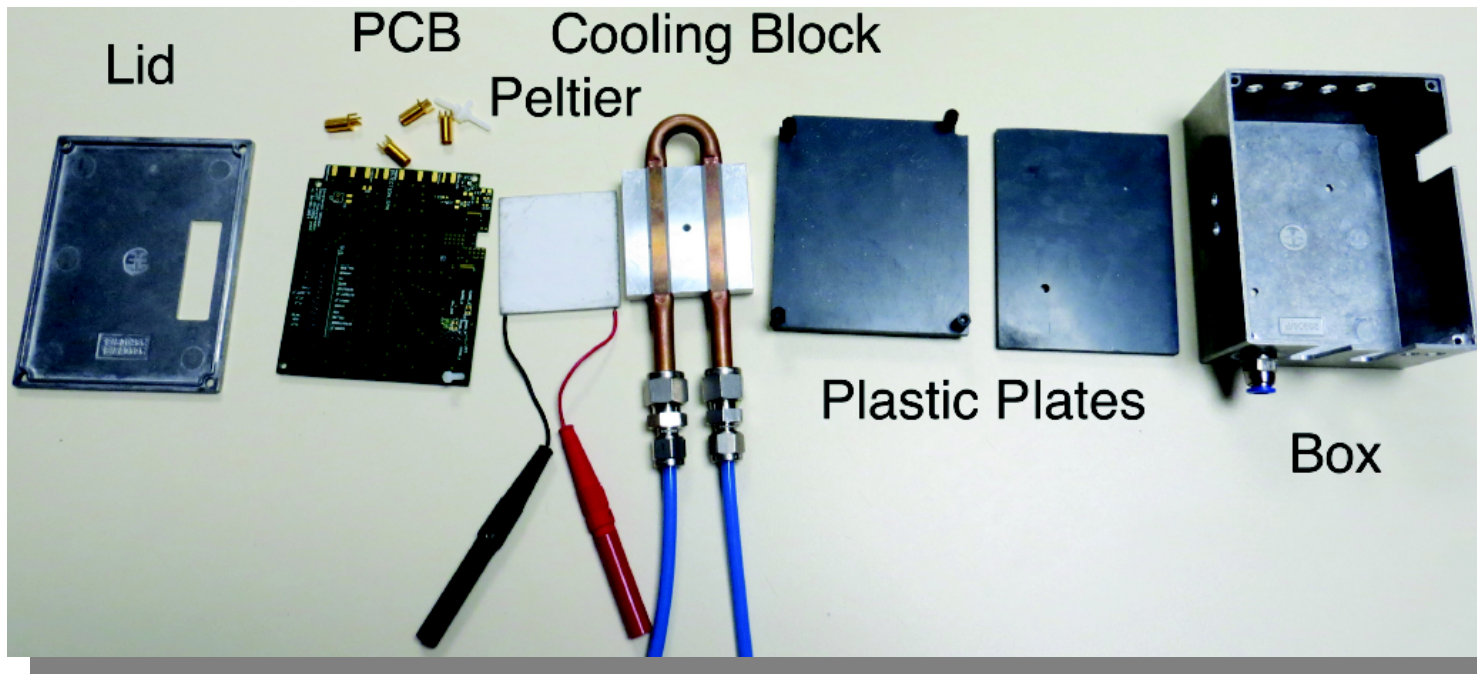
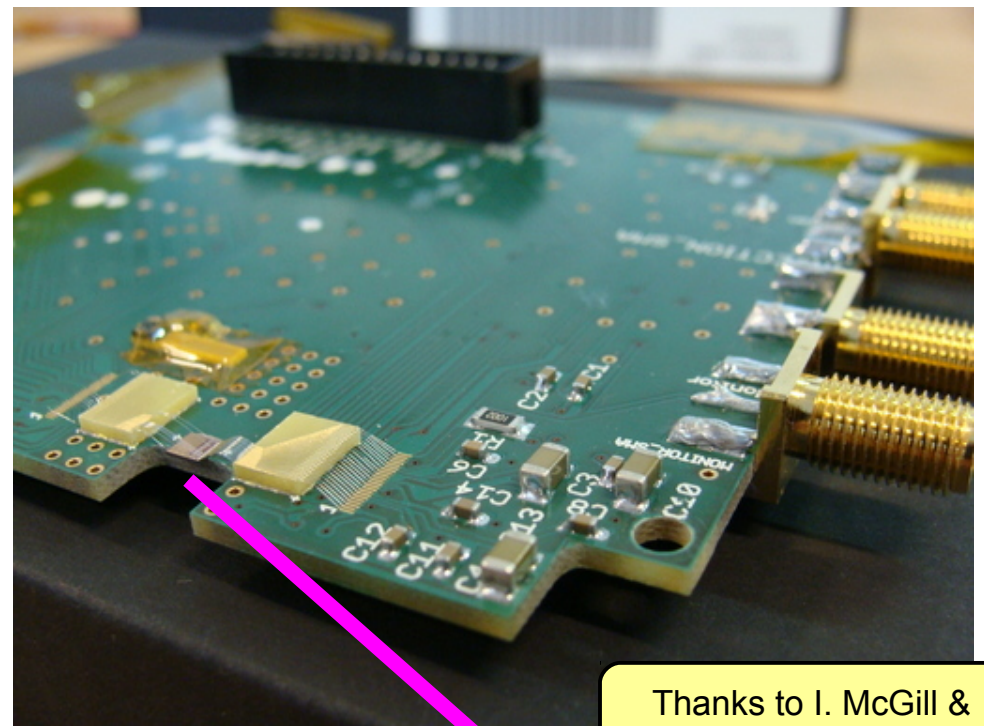
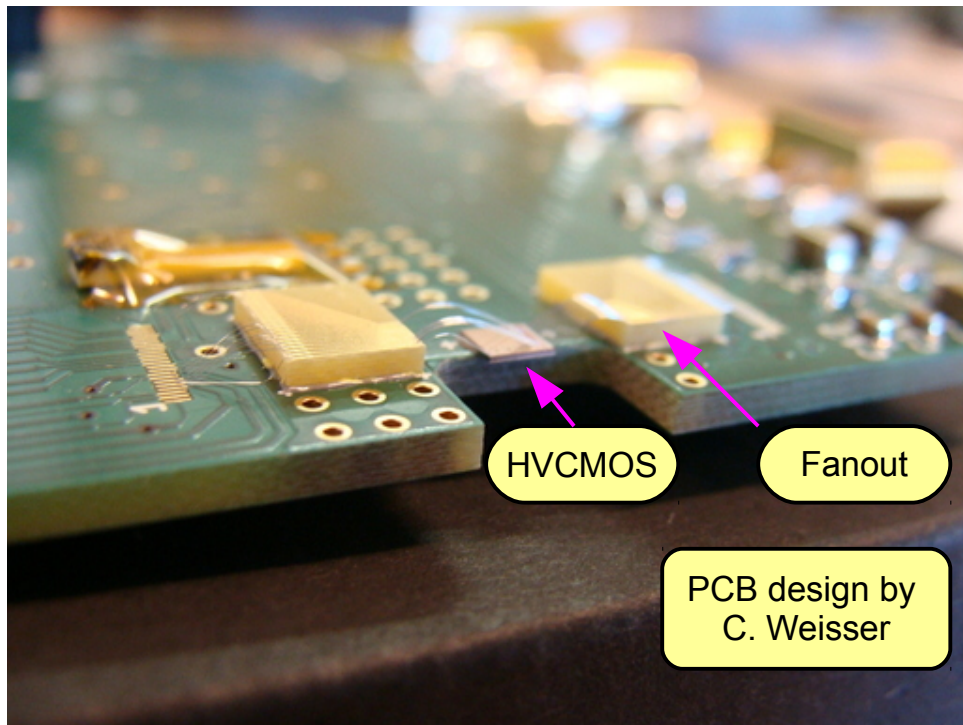
## Monolithic CMOS



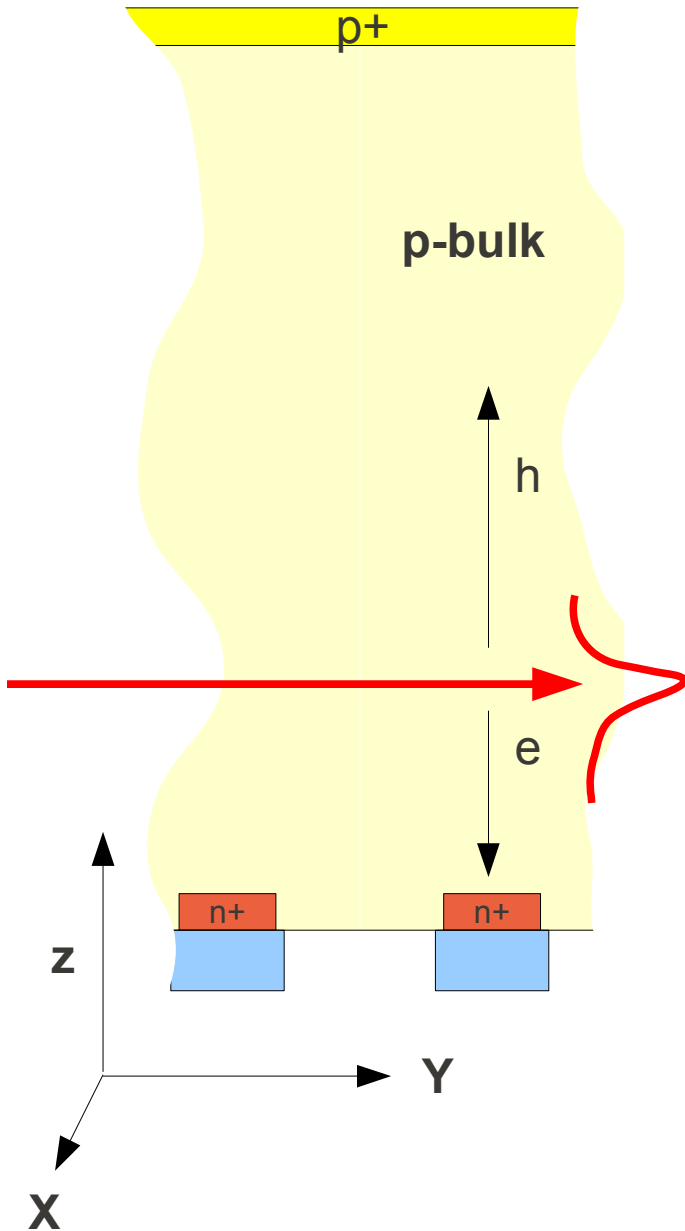
Unconfirmed guesses:  
 DNwell  $\sim 1 \times 10^{20} \text{ cm}^{-3}$  ? , 5  $\mu\text{m}$  depth (n-type)  
 Bulk  $10 \Omega \cdot \text{cm} = 1.4 \times 10^{15} \text{ cm}^{-3}$  (p-type)

“PowerPoint  
 HVCMOS”  
 (next slides)

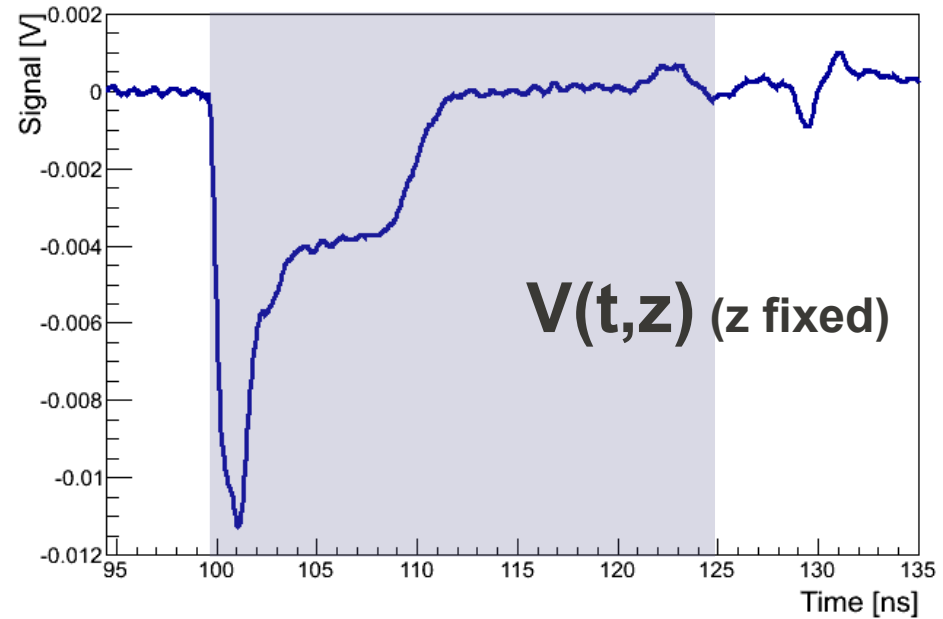
HVCMOS sketches from I. Peric



# Some eTCT jargon



Transient – current waveform



**Signal maps:**  
2D scan

$$I(t; z) \text{ } z \text{ variable}$$

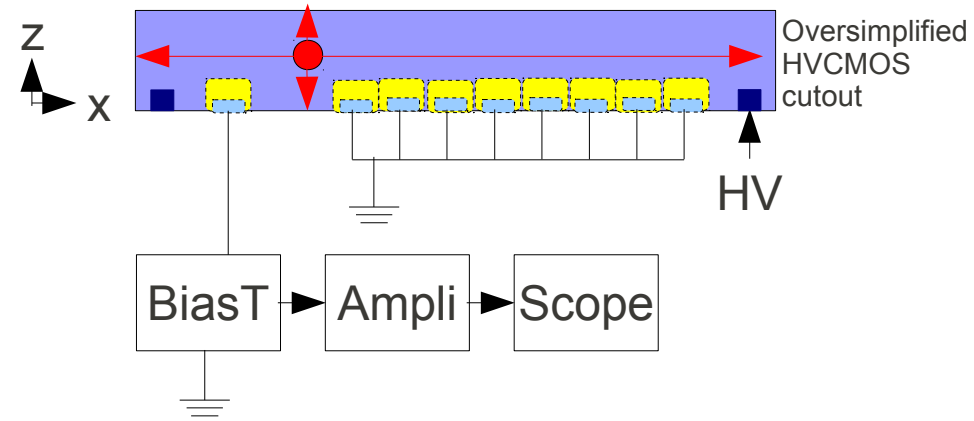
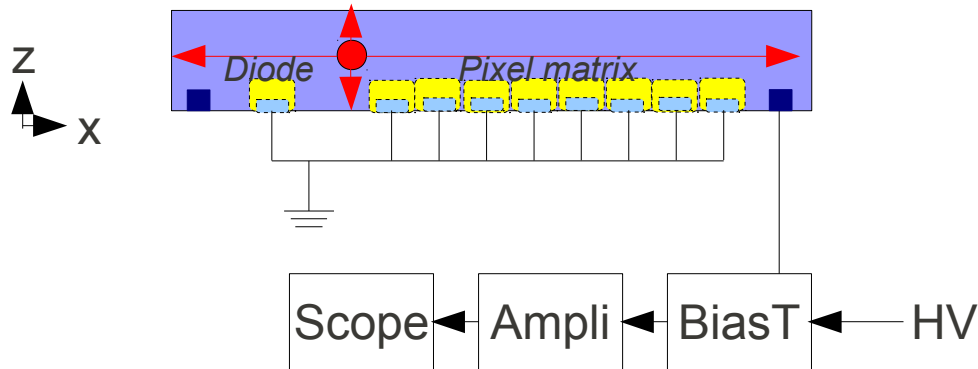
**Charge profile**  
1D scan

$$Q(z) = \int_0^{25 \text{ ns}} I(t, z) dt$$

**Charge maps:**  
2D scan

$$Q(x, z) = \int_0^{25 \text{ ns}} I(t; x, z) dt$$

# HVCMOSv3: 2 biasing configurations

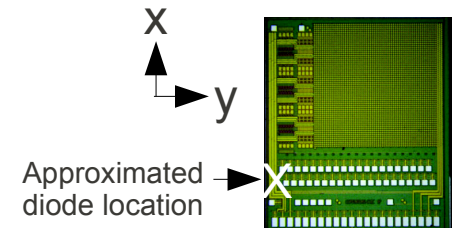


## Full chip readout, *a.k.a.* “configuration 1”

Good to “find” the detector [~2 mm wide]

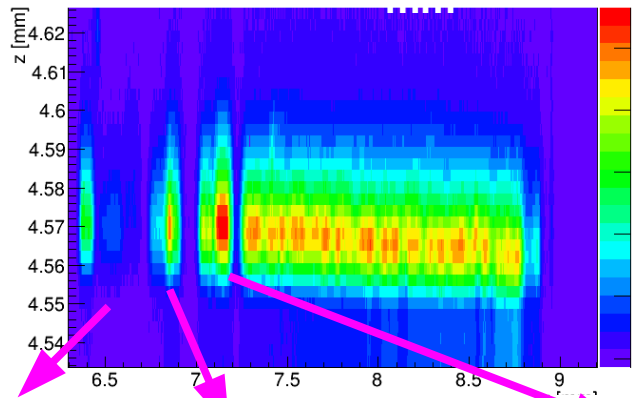
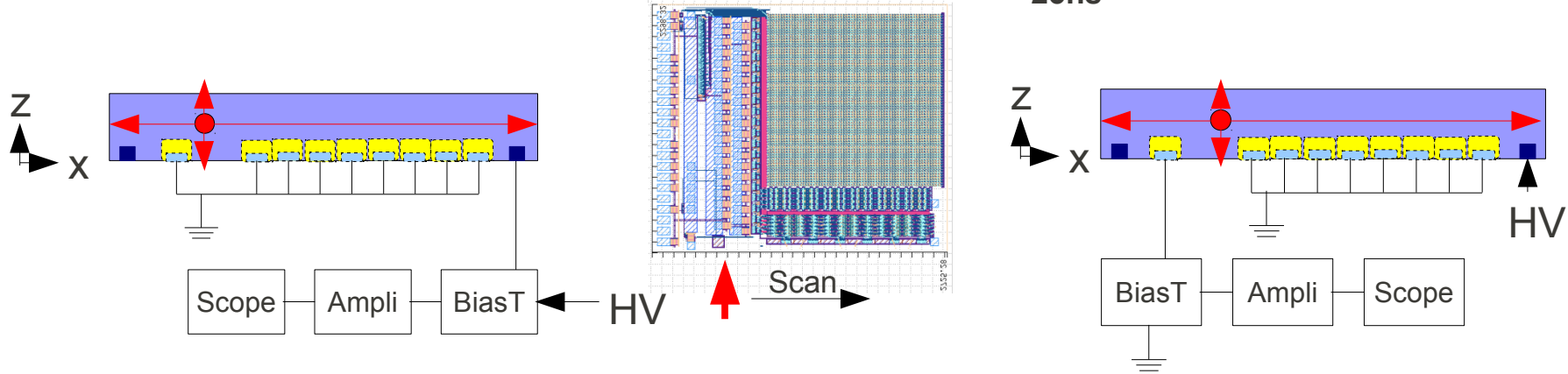
### Comments:

- 1) **PCB** gives **mechanical support** and allows, if needed, to **access the pixel matrix**. This feature was not used in these measurements.
- 2) In both configurations, we use the detector as a “monolithic **passive pixel**”, (we do not use the active part (preamp) at all). Analog waveform is picked at the PCB, amplified and displayed in the scope.
- 3) In edge-TCT **illumination** is produced **uniformly** along the light path (over several mm). Carriers produced before/after the measured “feature” can still be picked via **capacitive coupling**. In low res., they **will appear as diffusion**.

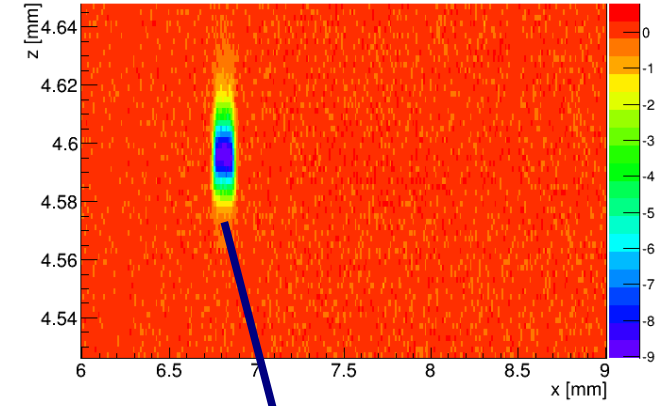


4) Measurements shown in this talk were taken at the same laser power ( $\pm 0.5 \mu\text{W}$ )

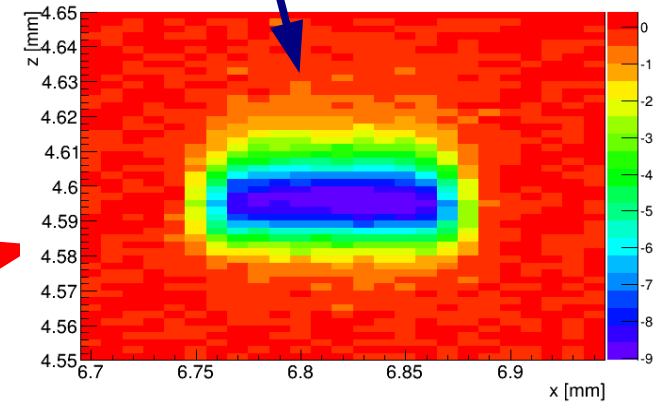
# Unirradiated HVCMOS – Charge maps: $Q_{25ns}(x,z)$



Sensor board is tilted



Diode readout in conf. 2



$(\Delta x=10, \Delta z=5)$   $\mu\text{m}$  resolution

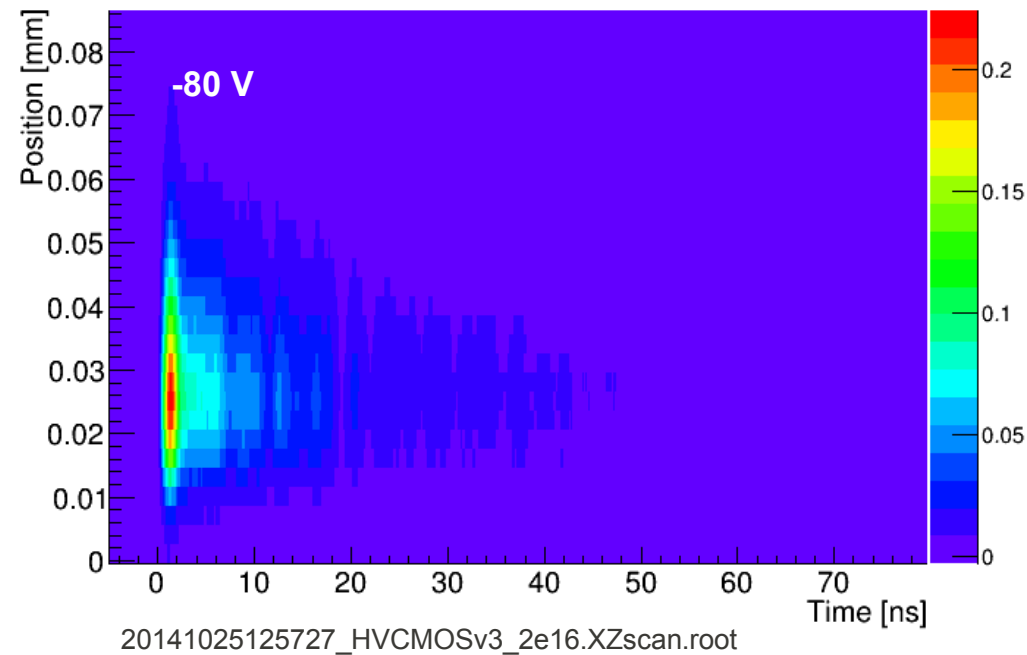
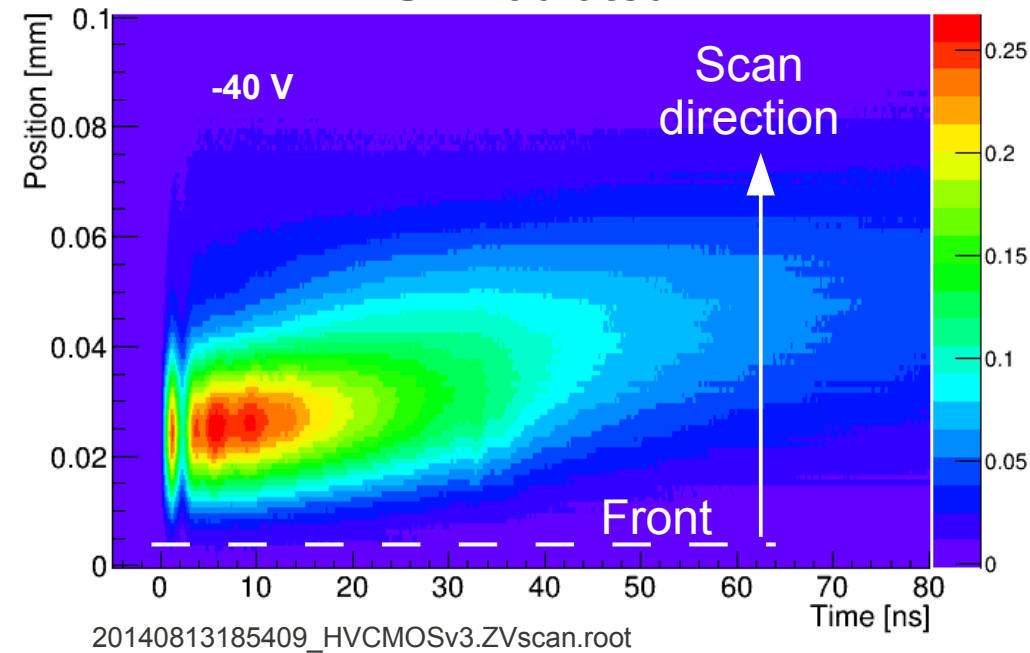
# Signal maps for non-irradiated and $2 \times 10^{16}$

## Effect of trapping on diffusion

- Comparison of **full chip readout** (conf 1) **non-irradiated** and  $2 \times 10^{16}$
- Showing signal maps:  $V(\text{waveform time, laser position})$  [ $V \equiv \text{scope signal}$ ]

Unirradiated

$2e16$



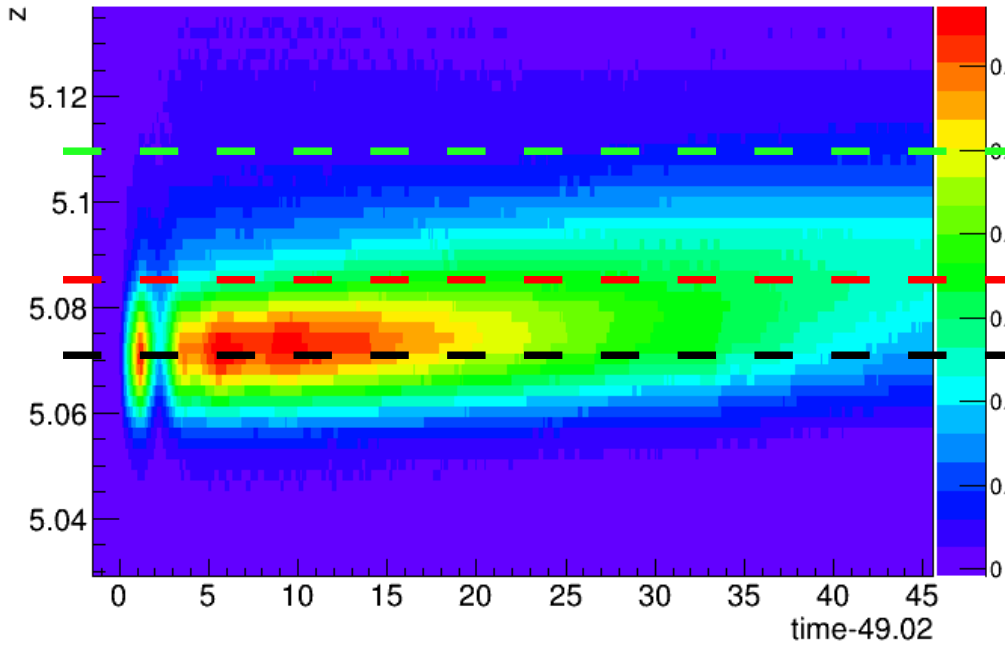
- Drift (time < 2 ns spot) and diffusion (time > 2 ns) clearly seen.
- Drift and diffusion amplitude decrease as beam moves away from the surface. After  $\sim 50 \mu\text{m}$  drift disappears while diffusion still survives, though amplitude reduces and shifts to longer times.

- Diffusion contribution is much smaller. Fast contribution still present.

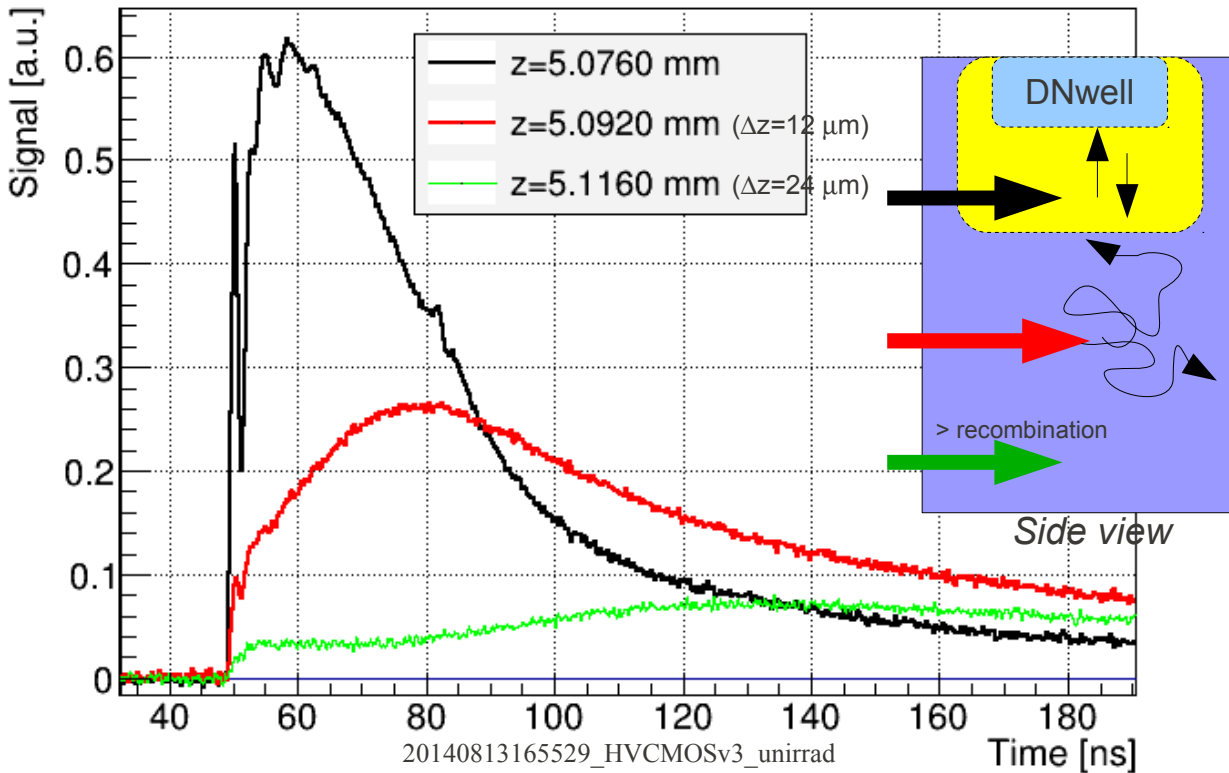
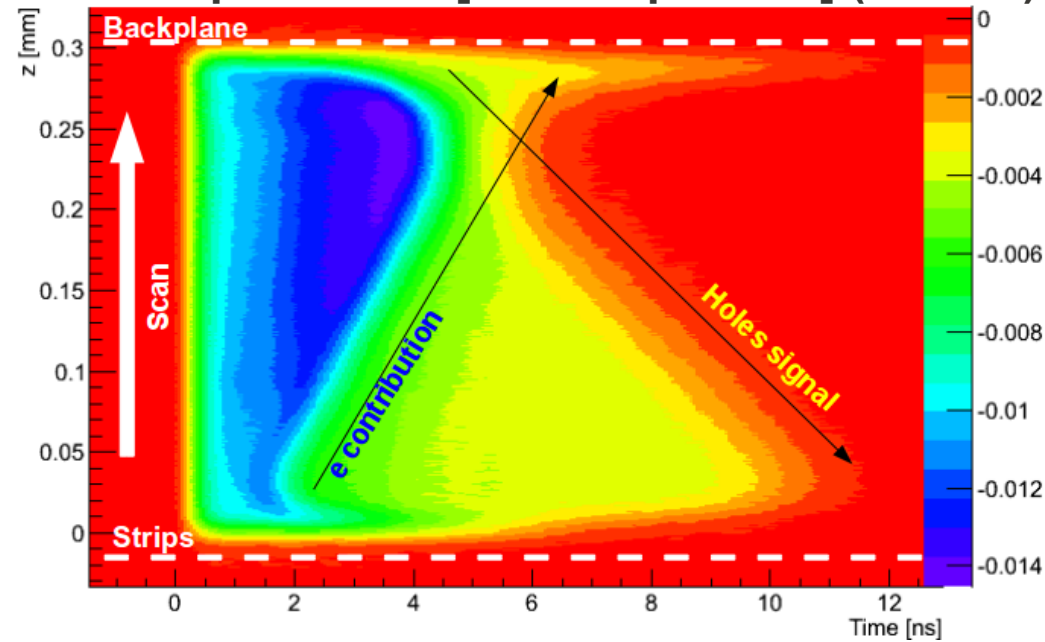


# Full chip readout: signal maps

## HVCMOSv3



## Microstrip detector [for comparison] (Micron)



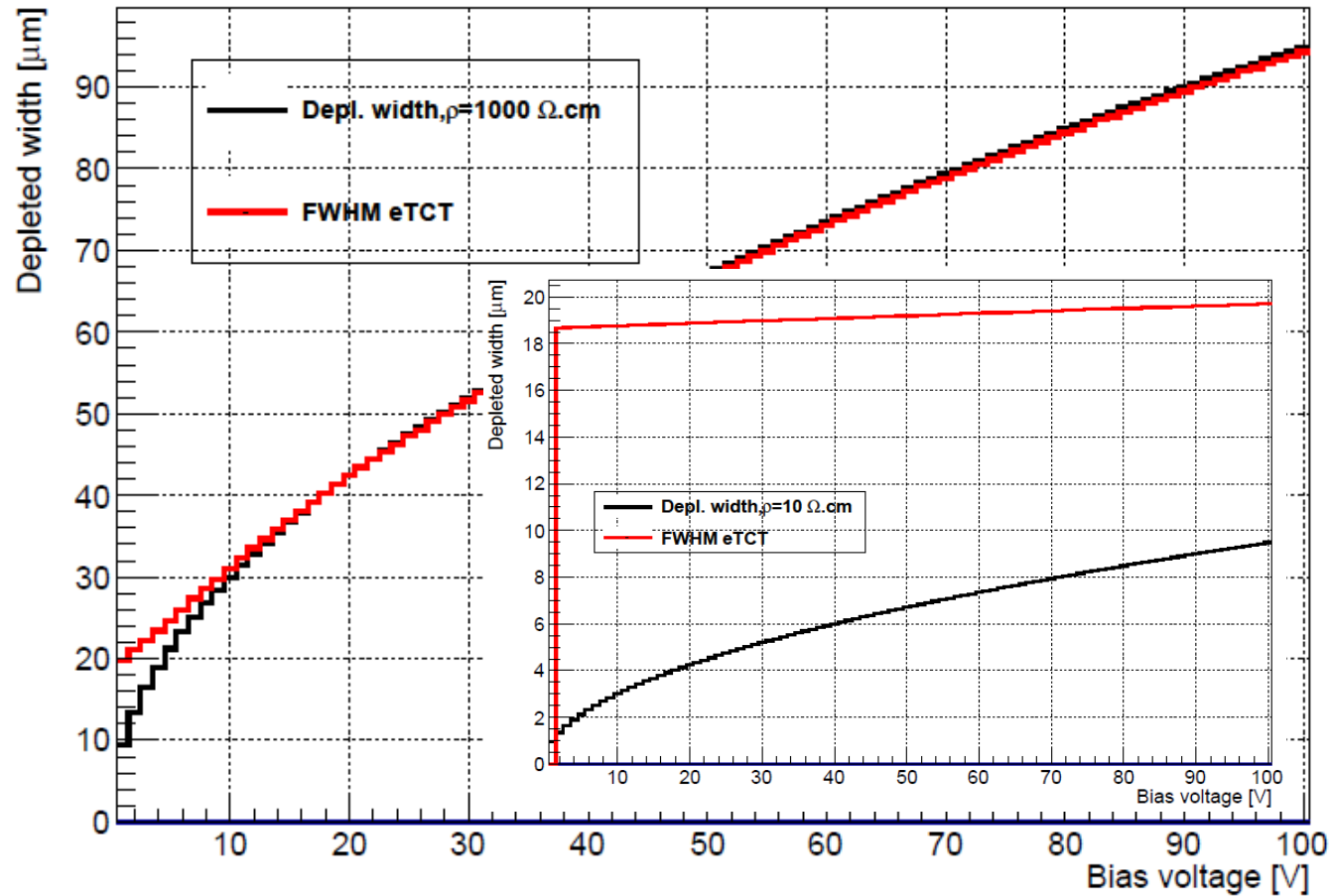
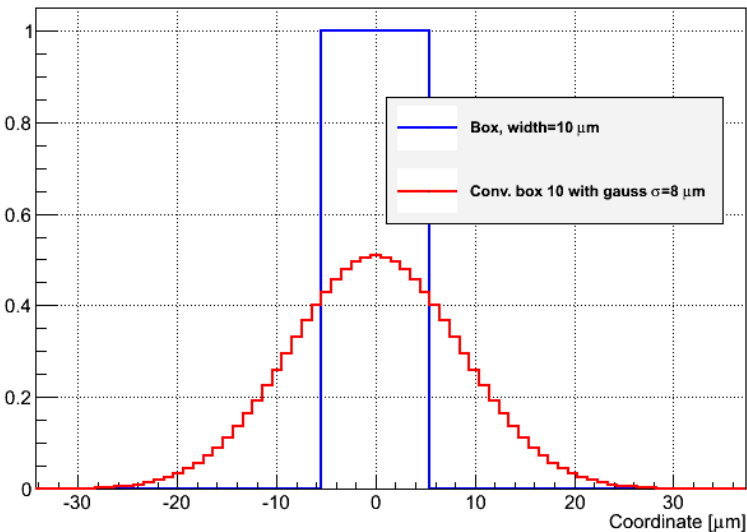
- Fast peak will be assimilated to drift, slower to diffusion.
- Diffusion maximum decreases in amplitude and shifts towards longer times as injection moves away from depleted region. Laser tails overlapping the depleted region may still produce direct drift.
- Drift peak inside the depleted bulk does not move

# Limiting spatial resolution in edge-TCT

Spatial resolution in edge-TCT is limited by laser beam gaussian width  $\sim 8 \mu\text{m}$ .

Assume a **real depleted region**  $10 \mu\text{m}$  wide.

**Measured depleted region** is the result of the convolution of the beam with the real **depleted "box"**.

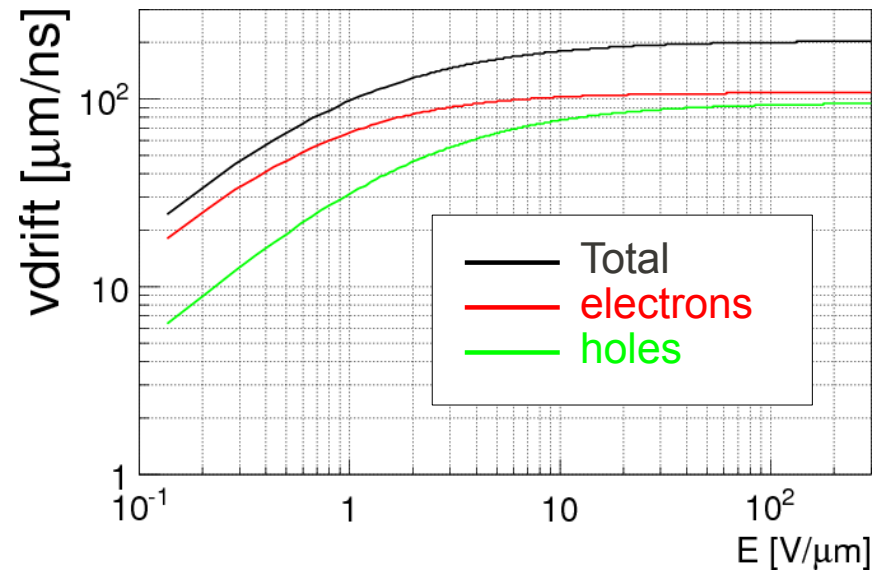


At 90V, for  $\rho=10 \Omega\cdot\text{cm}$ , a depletion region of  $\sim 10 \mu\text{m}$  (see inset) is expected. The apparent width measured will be FWHM  $\sim 20 \mu\text{m}$  ( $4\sigma=35 \mu\text{m}$ ).

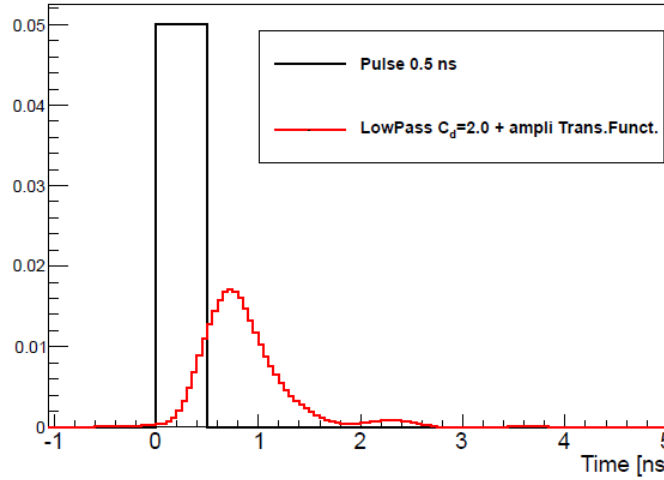
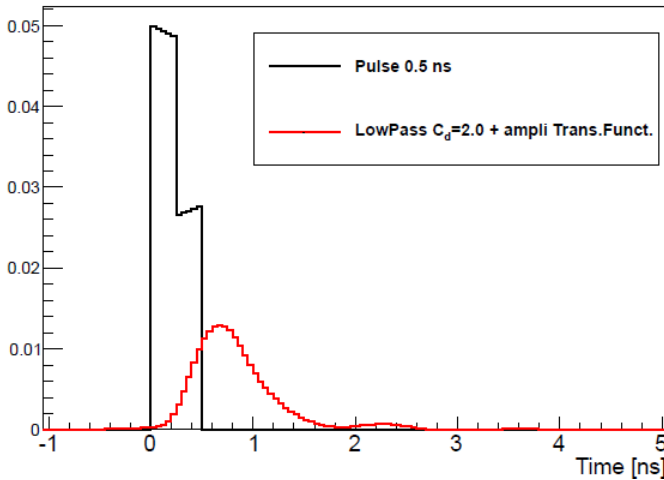
Real depleted width and edge-TCT measured ( $\sigma=8 \mu\text{m}$ ) coincide for real thickness  $> 10 \mu\text{m}$ .

# Drift velocity saturation in HVCMOS

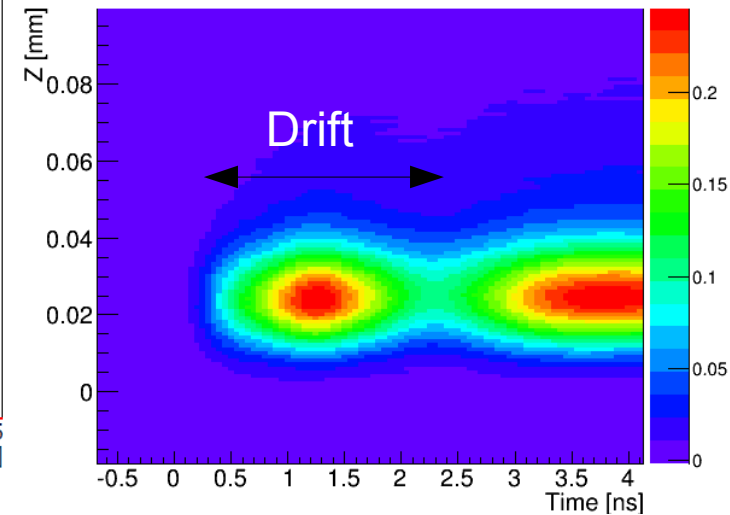
90V bias over 10  $\mu\text{m}$   $\rightarrow$  drift velocity saturation. Depleted region is crossed in (much) less than 1 ns, even at 10 V. Assumed, in this talk, that drift velocity in low resistivity silicon same as in high res.



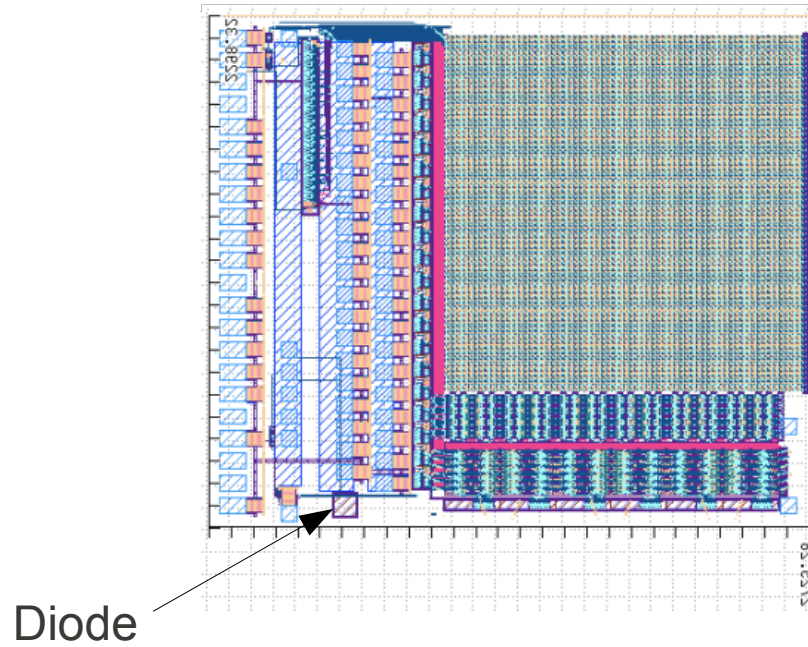
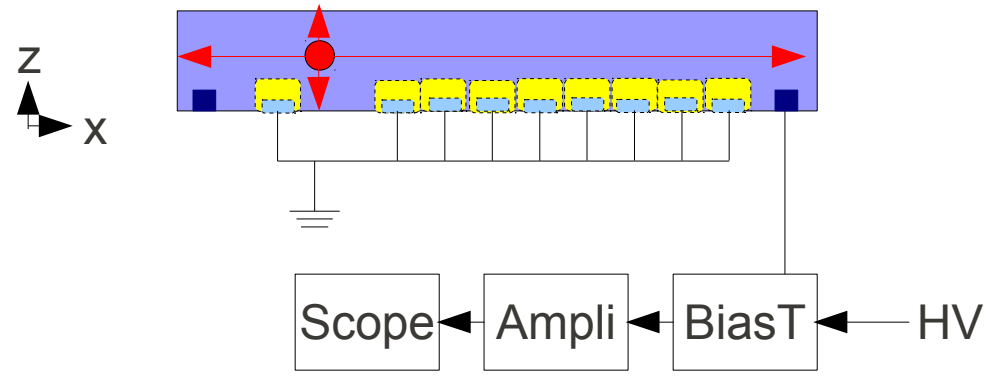
**Calculated** time response of Low Pass filtered (2pF) and amplified to various assumed HVCMOS-like signal. Output signal does not resolve electron/holes drift features (as we are used to see in thick sensors)



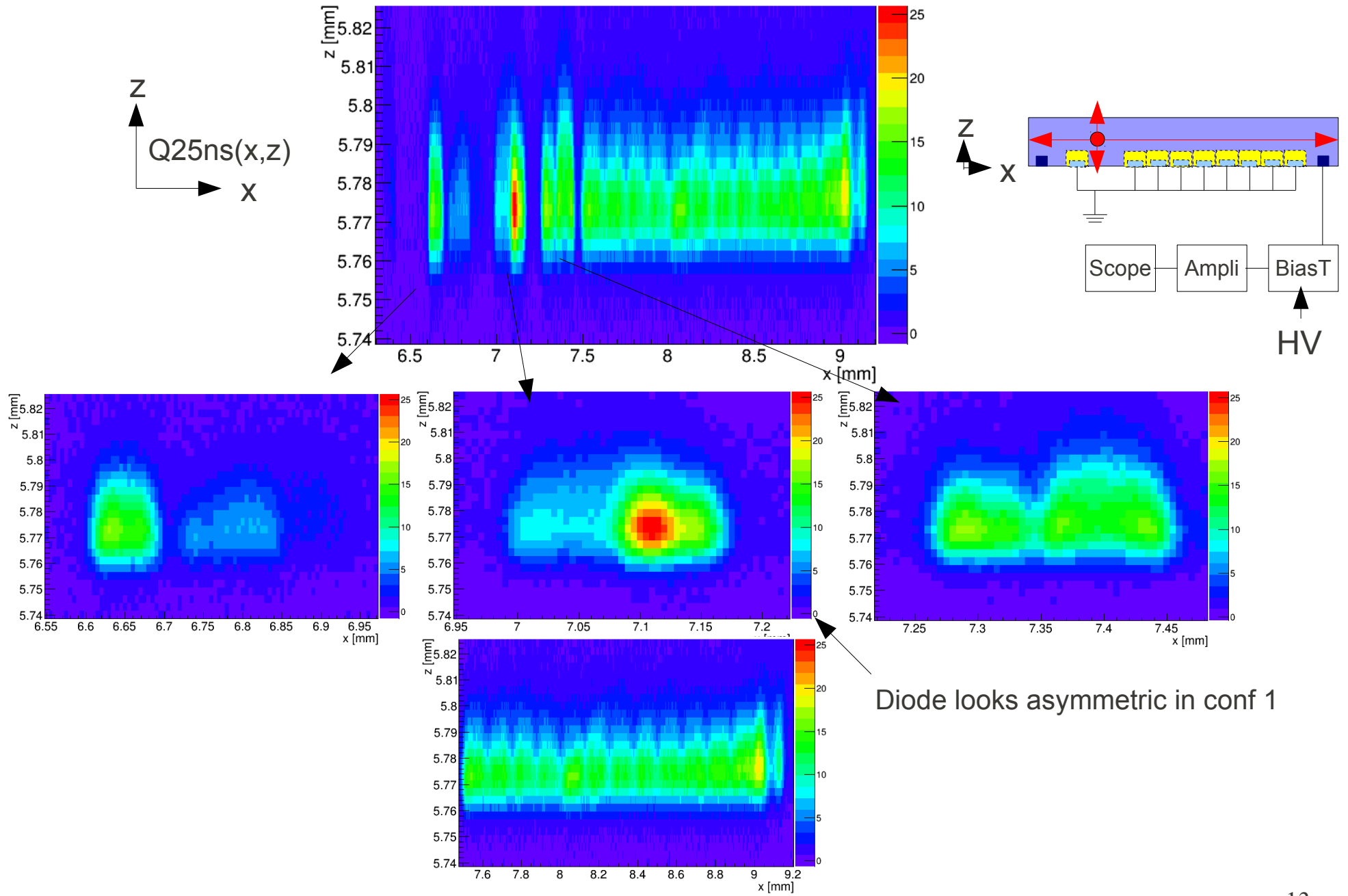
## Zoom on "fast" signal



# Full detector readout (configuration 1)



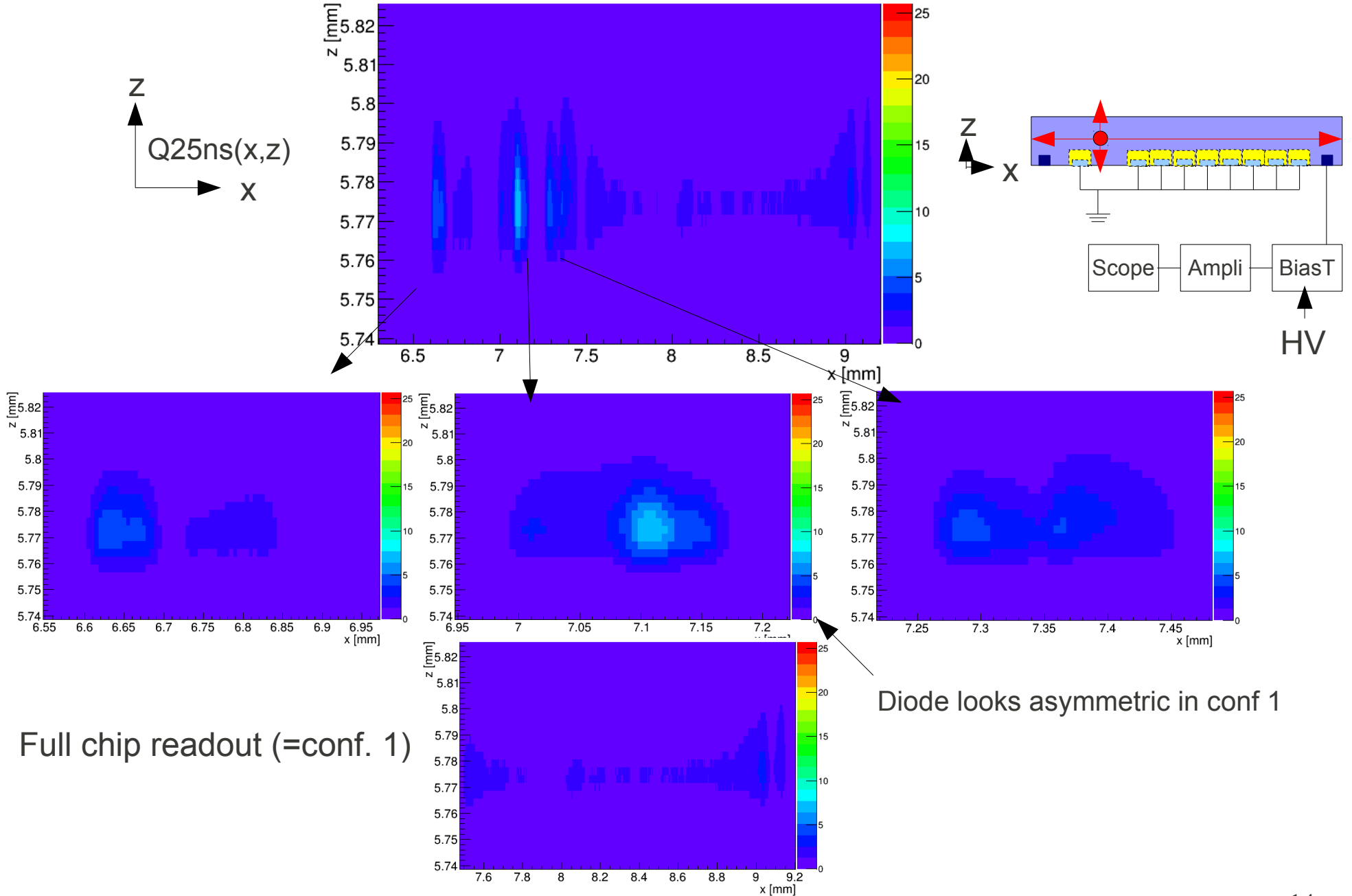
# 2e16 HVCMOS, full chip readout scan $Q_{25ns} = Q(x,z)$



Diode looks asymmetric in conf 1

$(\Delta x=10, \Delta z=5) \mu m$  resolution

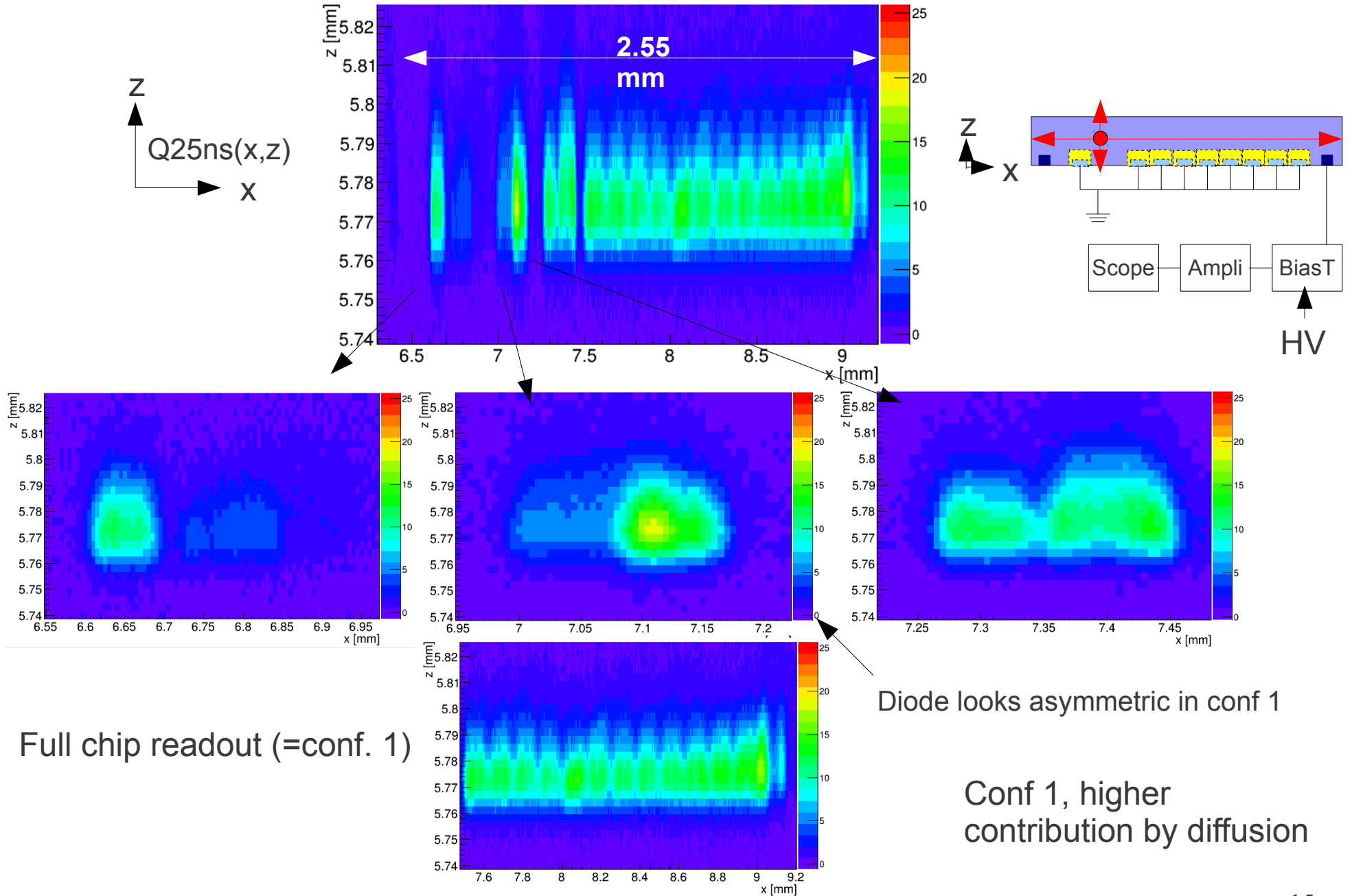
# 2e16 HVCMOS, full chip readout scan, only drift, Q(t<2.5 ns)



Full chip readout (=conf. 1)

$(\Delta x=10, \Delta z=5) \mu\text{m}$  resolution

# 2e16 HVCMOS, full chip readout scan, only diffusion, Q(t in [2.5,25] ns)



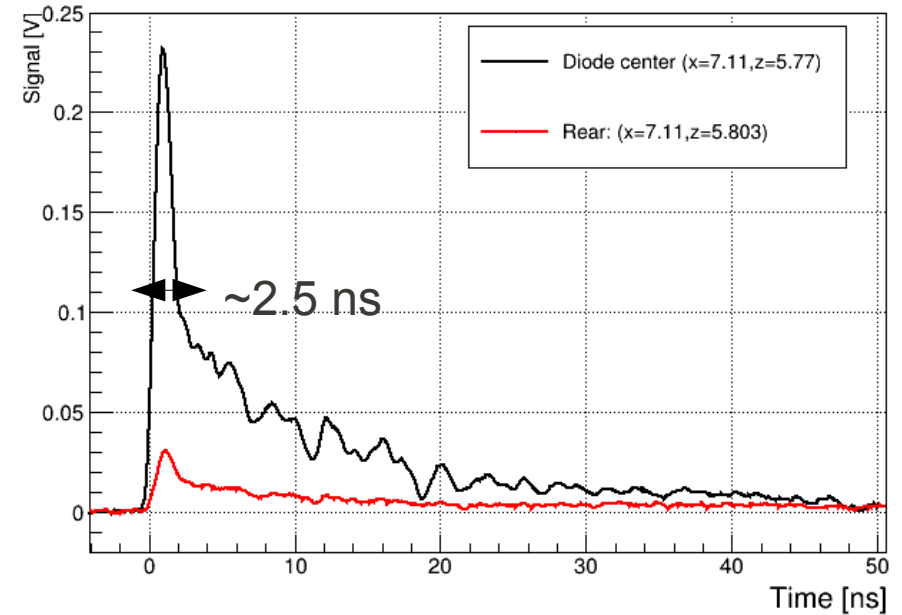
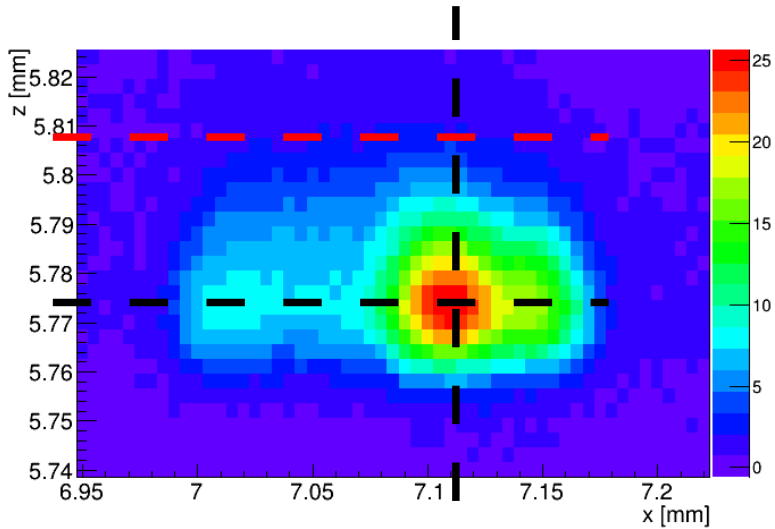
Full chip readout (=conf. 1)

Diode looks asymmetric in conf 1

Conf 1, higher contribution by diffusion

$(\Delta x=10, \Delta z=5) \mu\text{m}$  resolution

# 2e16 HVCMOS, zoom on diode in full detector readout

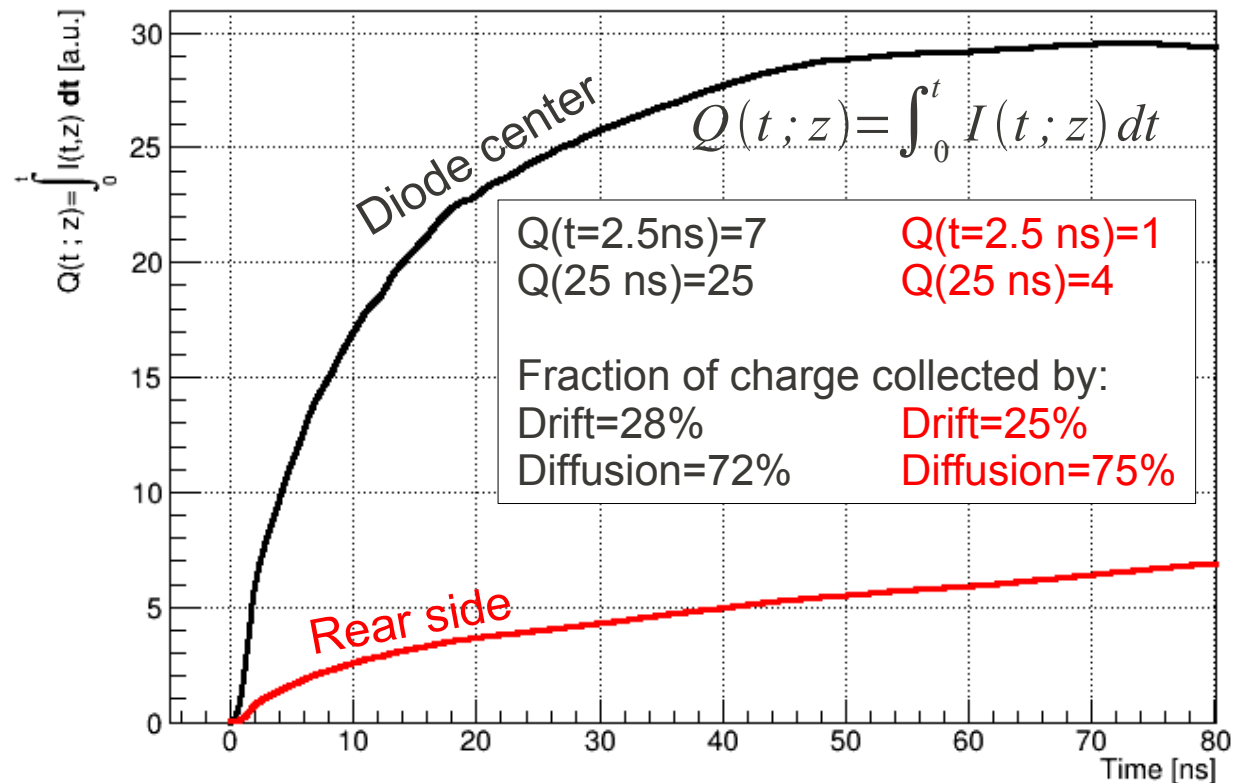


Distinctive long pulses due to diffusion have been washed out.

Very fast peak  $\sim 2.5$  ns wide associated to drift (as in the case of the unirradiated detector).

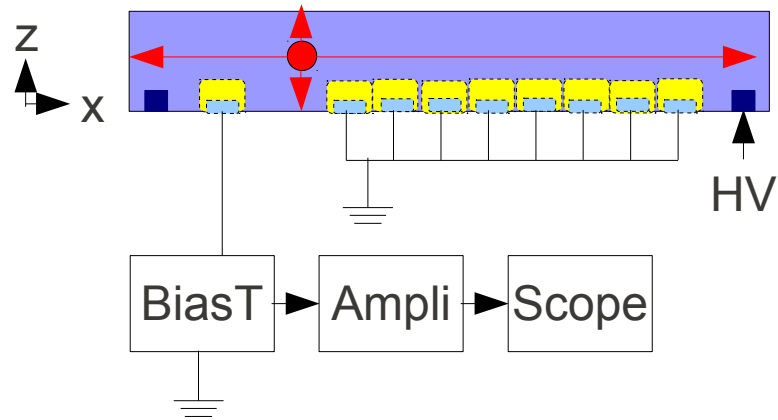
But diffusion still present even after  $2 \times 10^{16}$ .

At  $2e16$ , diffusion accounts for 75% of the collected signal (in 25 ns). It was 94% in the unirradiated detector.

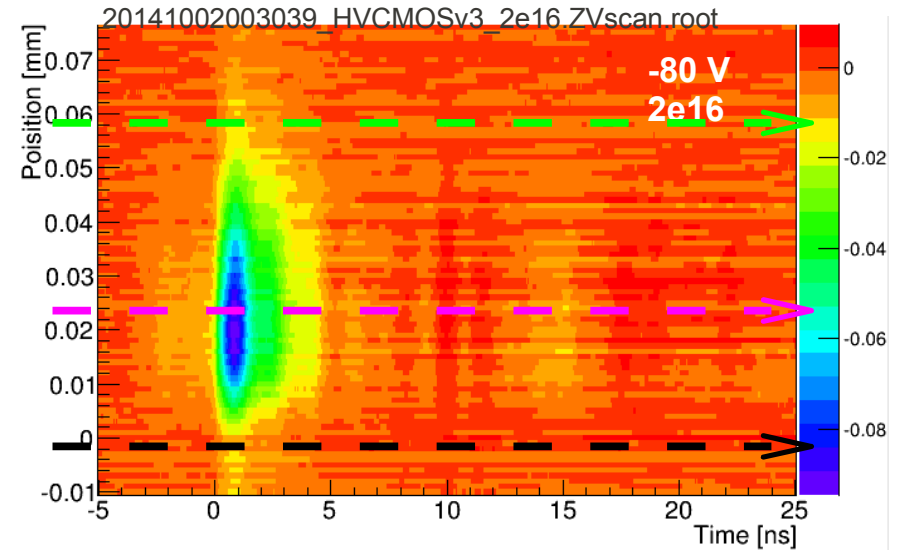
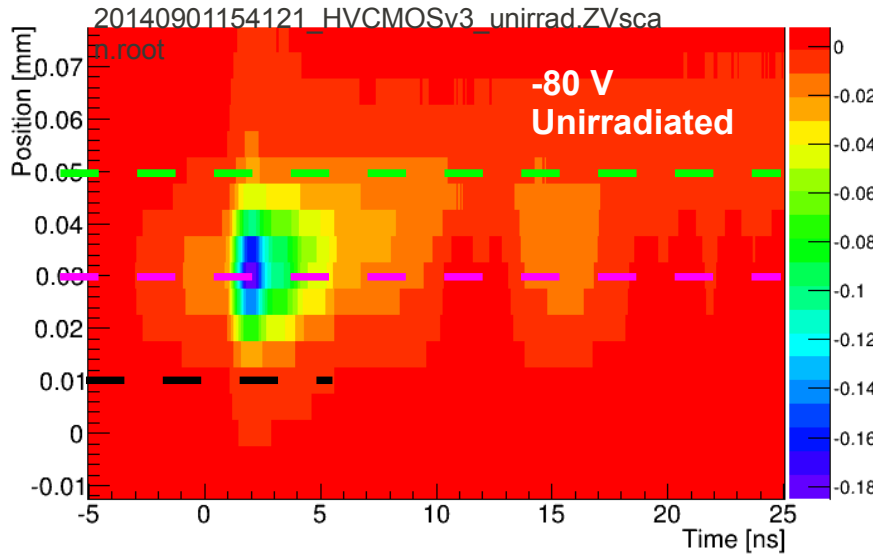




## Diode only readout (configuration 2)

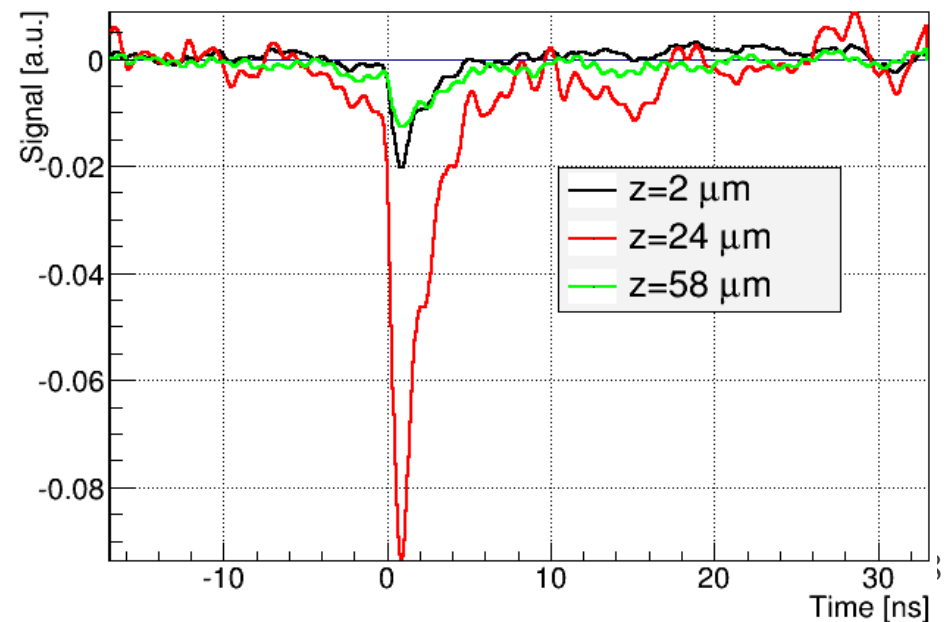
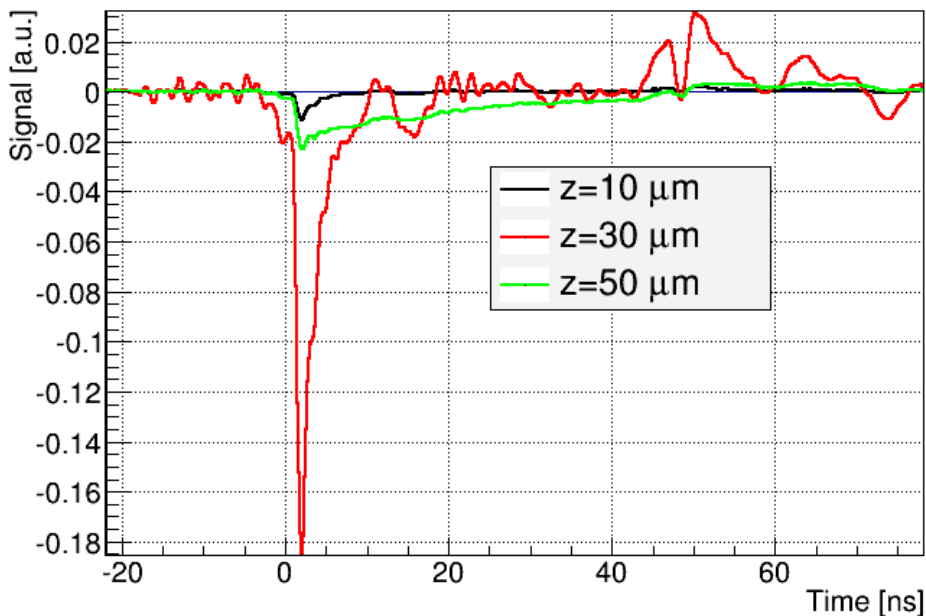


# Diode readout (conf. 2) signal maps: $V(z,t)$

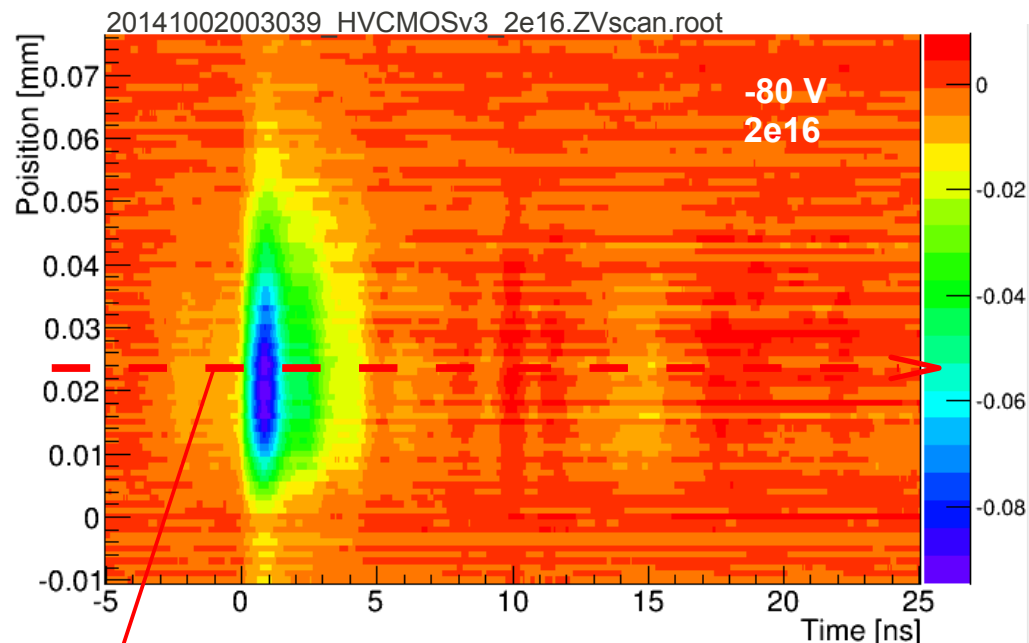
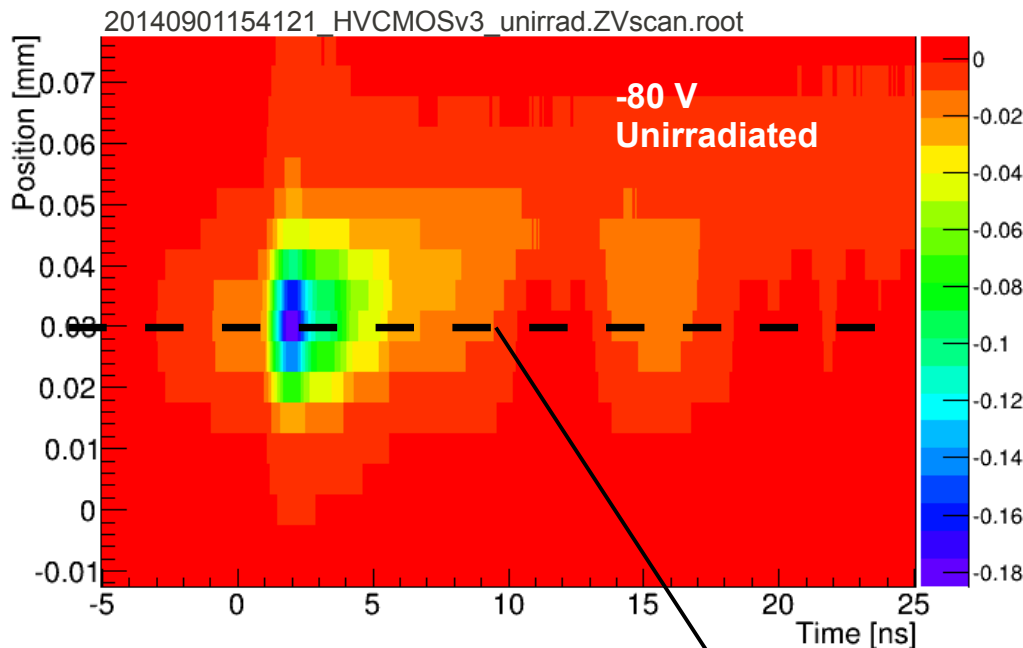


Difficult to assign drift and diffusion tags. Some diffusion still visible (see  $z=50 \mu\text{m}$ ), but mainly fast signal (from drift). Fast peak is wider. Arbitrarily taking 5 ns now as drift peak width,

Efficient region for charge creation at  $2e16$  seems wider.



# Diode readout (conf. 2), integrated charge

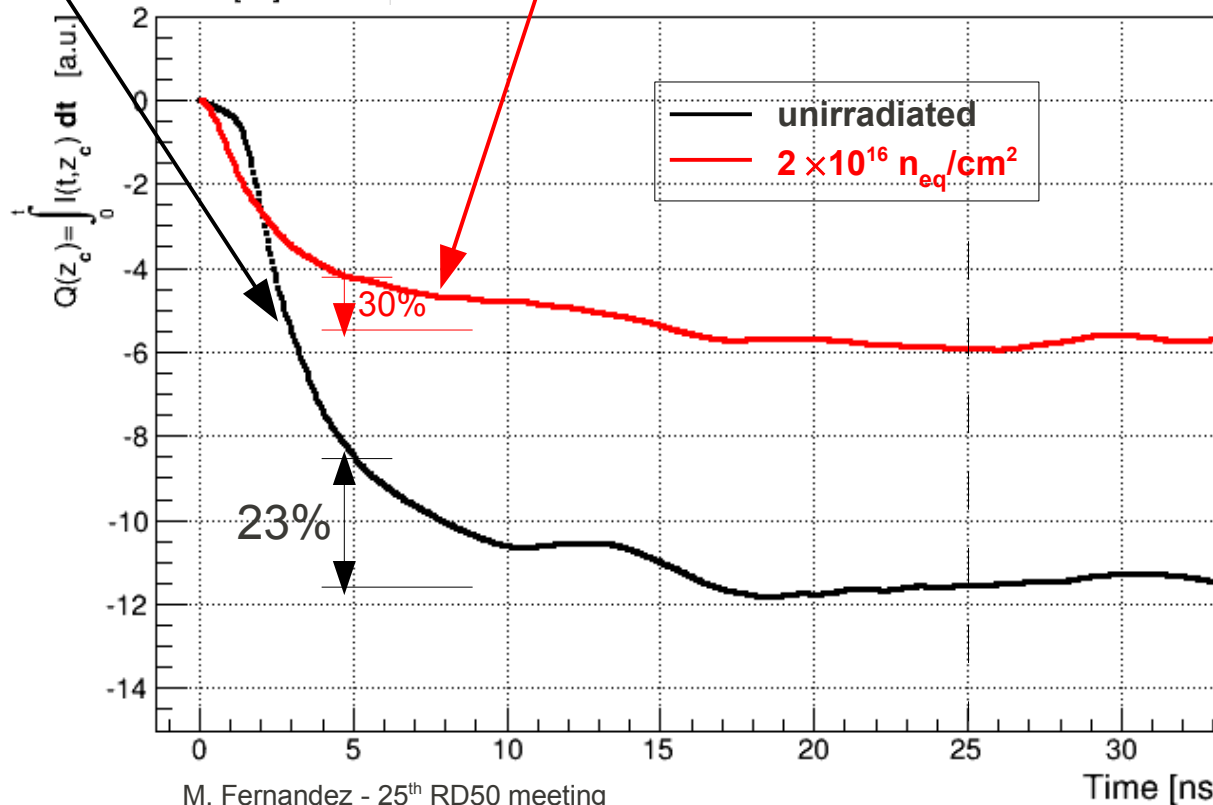


Shown to the right the charge as a function of integration time, at the center of the depleted region ( $z_c$ ):

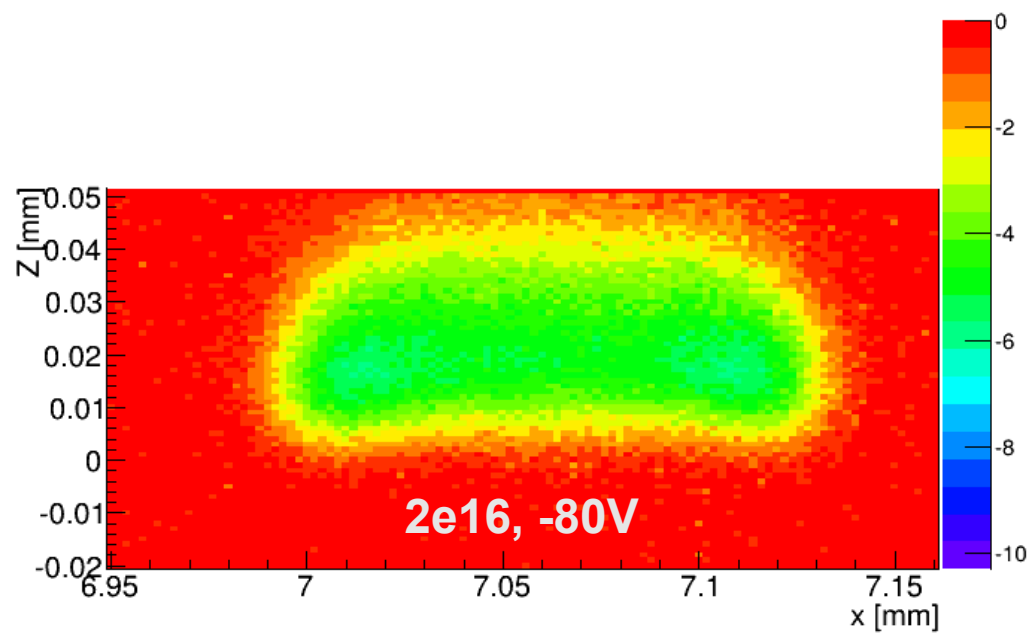
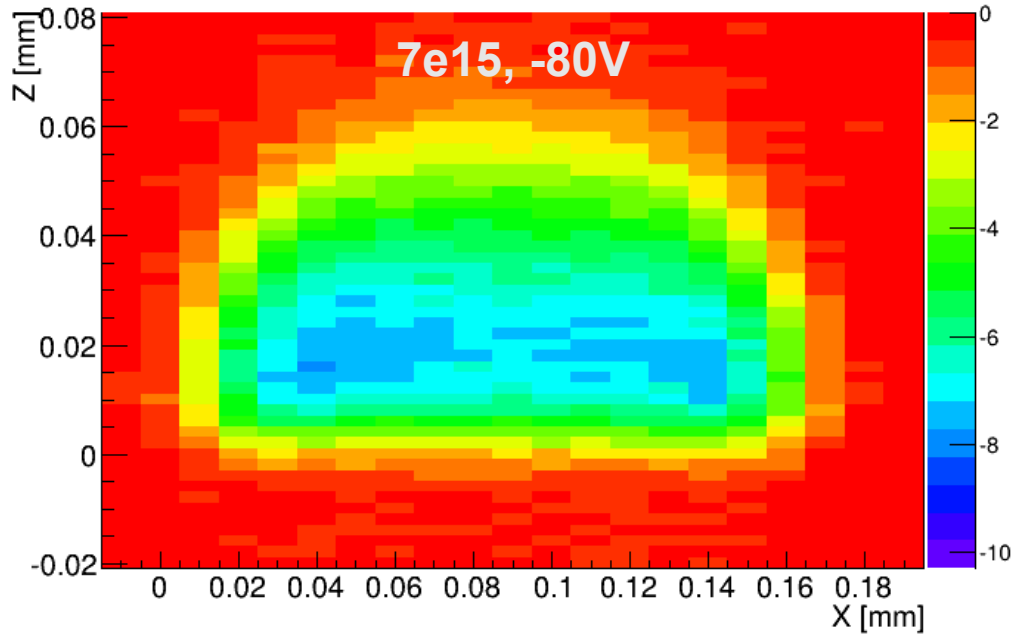
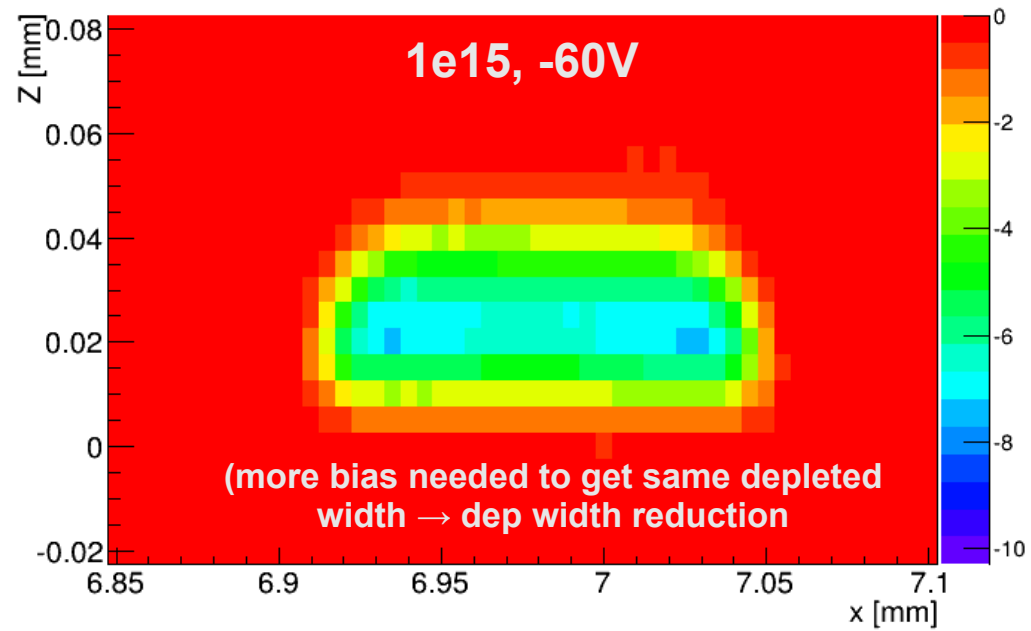
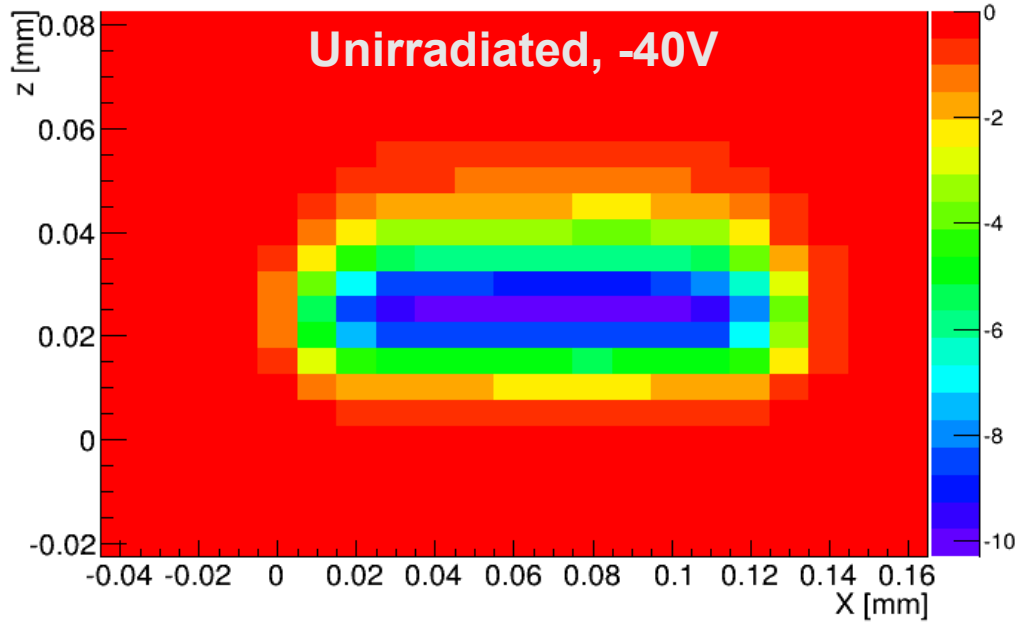
$$Q(t) = \int_0^t I(t) dt$$

At  $t=5$  ns, the  $2 \times 10^{16}$  detector collects half the charge of the unirradiated (drift)

Diffusion ( $t > 5$  ns) contributes:  
30 % for  $2 \times 10^{16}$   
23% for unirradiated

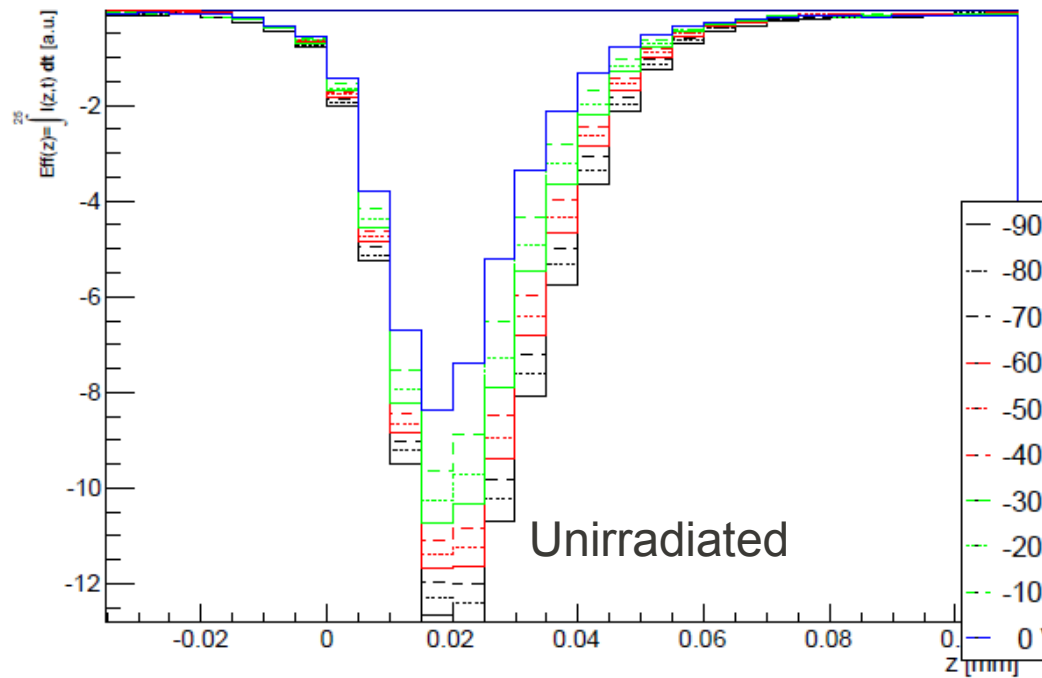


# Diode readout, charge maps (various voltages)

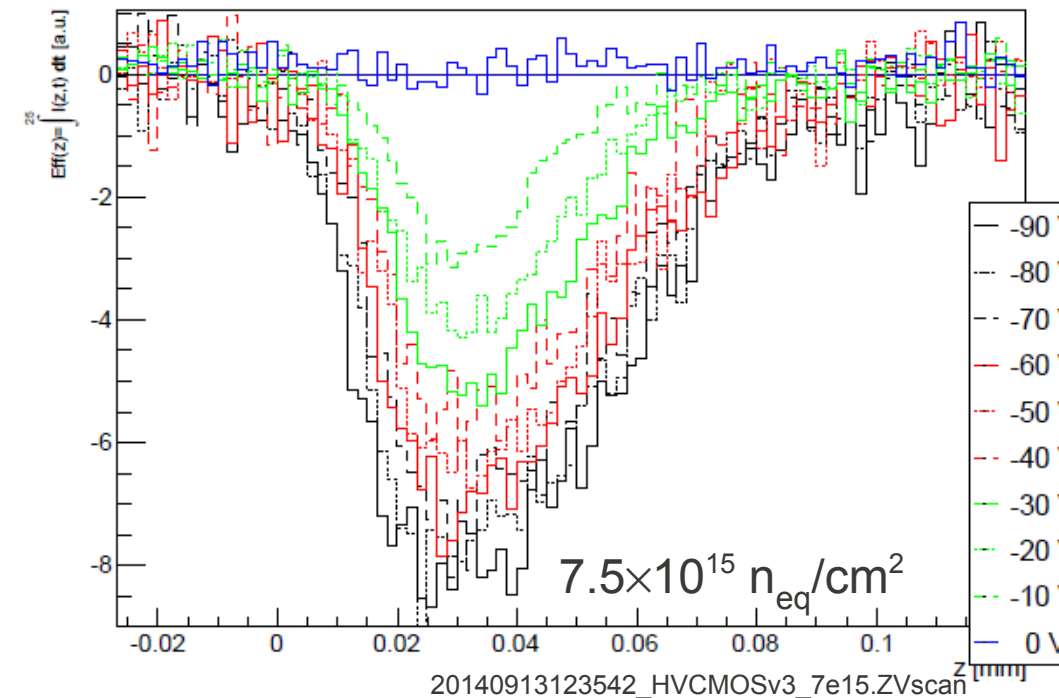
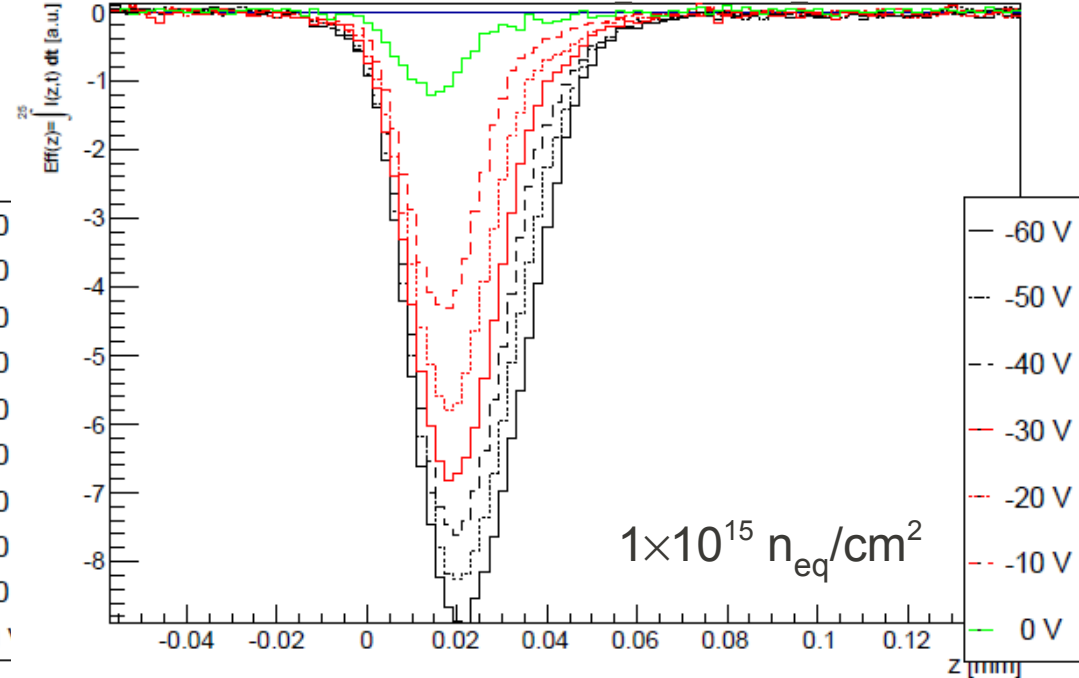


# Charge in 25 ns at the center of the diodes (mind different X-axis)

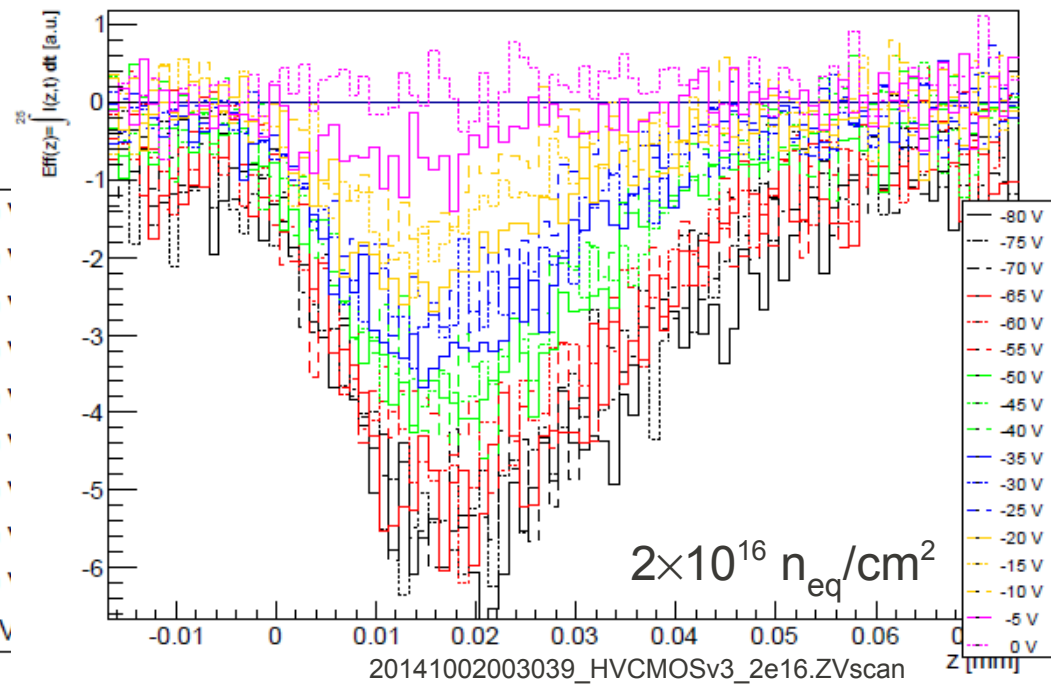
20140901154121\_HVCMOSv3\_unirrad.ZVscan



20140904145059\_HVCMOSv3\_irrad2.ZVscan

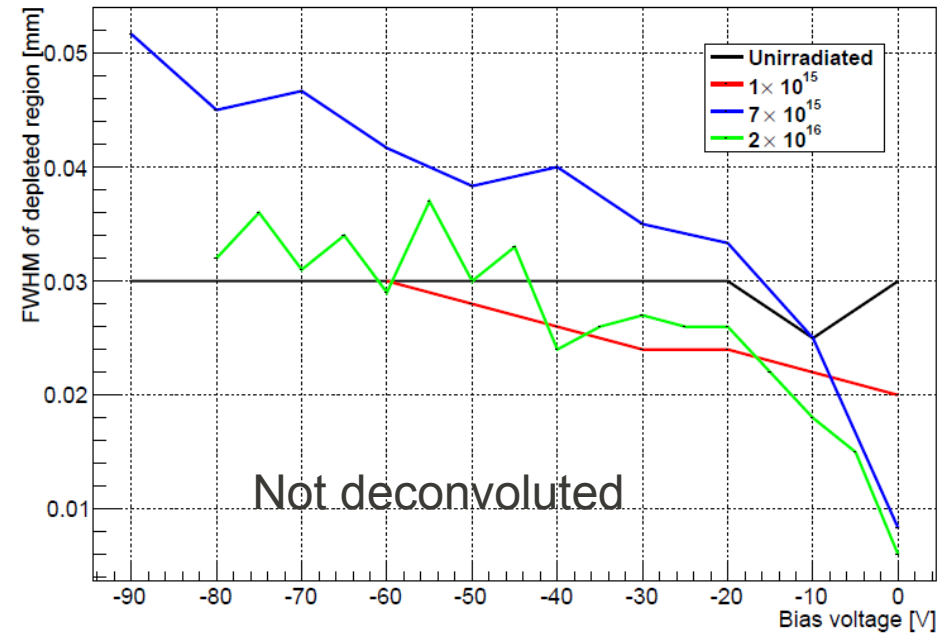
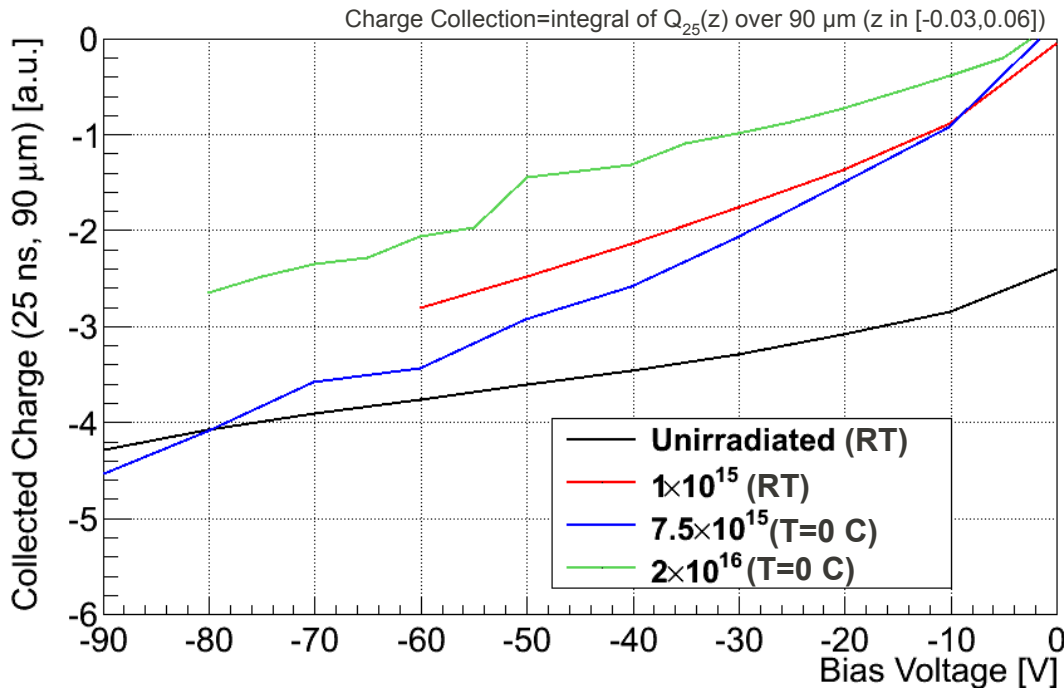


20140913123542\_HVCMOSv3\_7e15.ZVscan



20141002003039\_HVCMOSv3\_2e16.ZVscan

# HVCMOSv3, radiation hardness



Note measurements at different T! Absorption in Si decreases with decreasing T  $\rightarrow$  less signal produced at lower T  $\rightarrow$  CCV of  $7 \times 10^{15}$ ,  $2 \times 10^{16}$  are probably underestimated.

## Interpretation:

$$w = 0.3 \text{ } [\mu\text{m}] \times \sqrt{\rho(V_b + V_{bi})}$$

$$V_{dp} = 11 \left[ \frac{\Omega \cdot \text{cm}}{(\mu\text{m})^2} \right] \cdot \frac{d^2}{\rho_p} - V_{bi}$$

$1 \times 10^{15}$ : Depth of depleted region decreases  $\leftrightarrow$   $\rho$  decrease  $\leftrightarrow$   $V_{\text{dep}}$  increase. Less charge collected for the same voltage across the detector  $\rightarrow$  drop of collected charge.

$7 \times 10^{15}$ : Depleted width increases, we collect more charge for the same voltage. For bias  $> 80\text{V}$  it seems to overcome the unirradiated value

$2 \times 10^{16}$ : depletion width as unirradiated. However less charge collected  $\rightarrow$  trapping

Depletion width looks bigger than expected due to contribution of diffusion in 25 ns!!

## Conclusions

- Tested 4 HVCMOSv3 to 4 different fluences: 0,  $1 \times 10^{15}$ ,  $7 \times 10^{15}$ ,  $2 \times 10^{16}$   $n_{eq}/cm^2$
- Only analog readout, no amplified signal from pixel matrix was used
- Full sensor readout (conf. 1) shows high contribution from diffusion, reducing after irradiation. For configuration 2 (diode only), drift dominates
- Different relative contributions to collected charge from drift and diffusion:

| Q contribution<br>conf.1 | Unirradiated      |                     | $2 \times 10^{16}$ |                    |
|--------------------------|-------------------|---------------------|--------------------|--------------------|
|                          | Drift<br>[2.5 ns] | diff<br>[2.5-25 ns] | Drift<br>[5 ns]    | diff.<br>[5-25 ns] |
| @drift max               | 6%                | 94%                 | 25%                | 75%                |
| Bias voltage [V]         | -40               |                     | -80                |                    |

| Q contribution<br>conf.2 | Unirradiated    |                   | $2 \times 10^{16}$ |                    |
|--------------------------|-----------------|-------------------|--------------------|--------------------|
|                          | Drift<br>[5 ns] | diff<br>[5-25 ns] | Drift<br>[5 ns]    | diff.<br>[5-25 ns] |
| @drift max               | 77%             | 23%               | 70%                | 30%                |
| Bias voltage [V]         | -80             |                   | -80                |                    |

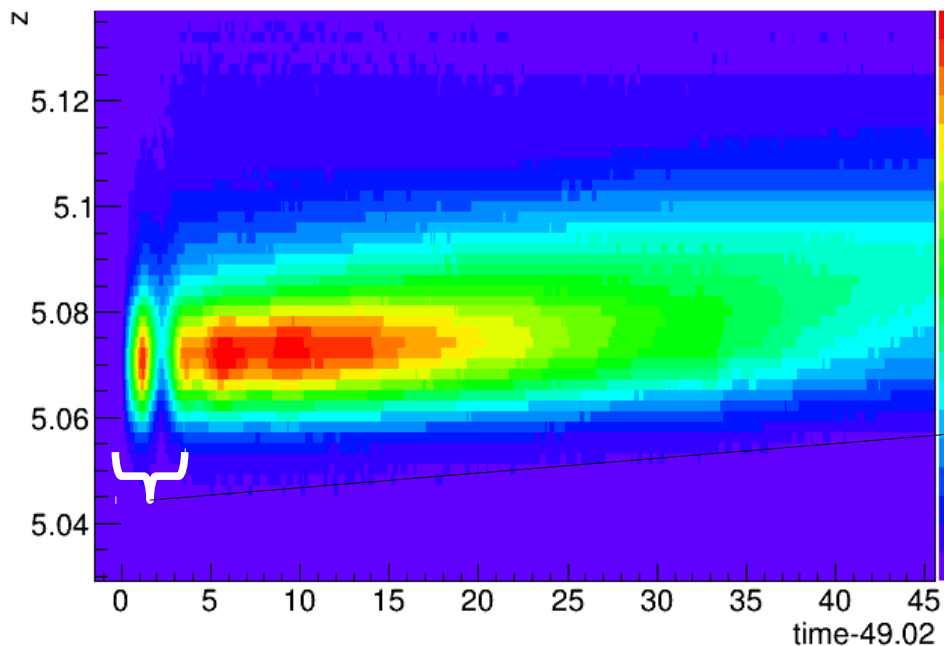
- Region efficient for charge creation reduces for  $1 \times 10^{15}$ , increases for  $7 \times 10^{15}$  and recovers the unirradiated level at  $2 \times 10^{16}$ .
- At  $1 \times 10^{15}$  the detector could be operated without cooling. Collected charge generally decreases down to 50% at -60V and  $2 \times 10^{16}$ .

① Many things can be improved in next measurements: PCB and mechanics. We need also more statistics as well.

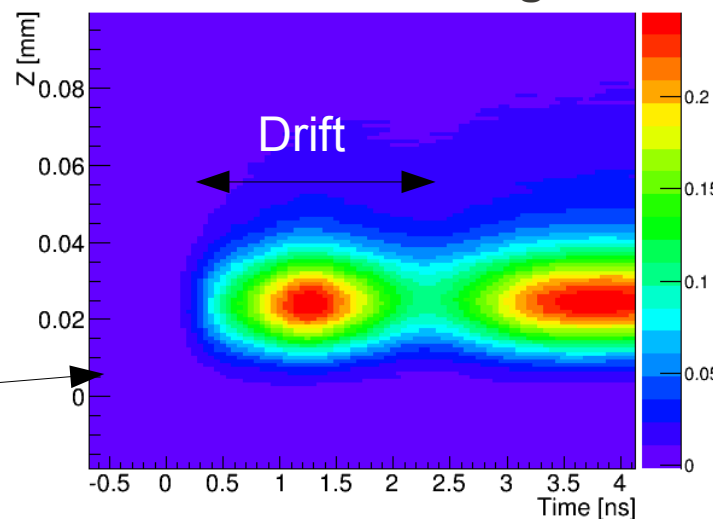
**BACKUP**



# Full chip readout: charge profiles $Q(z)$

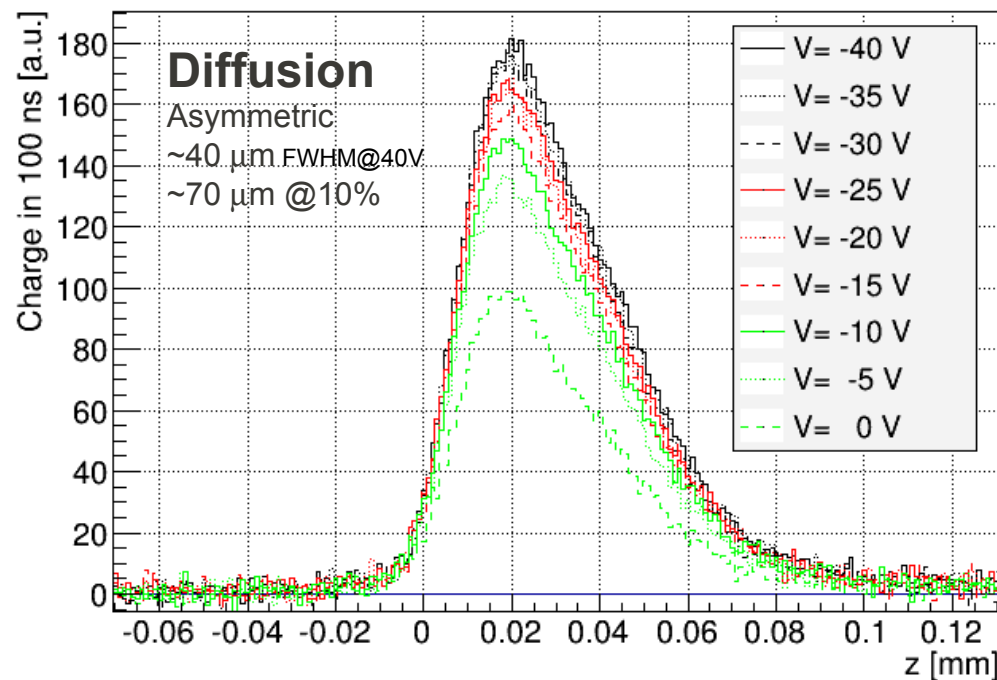
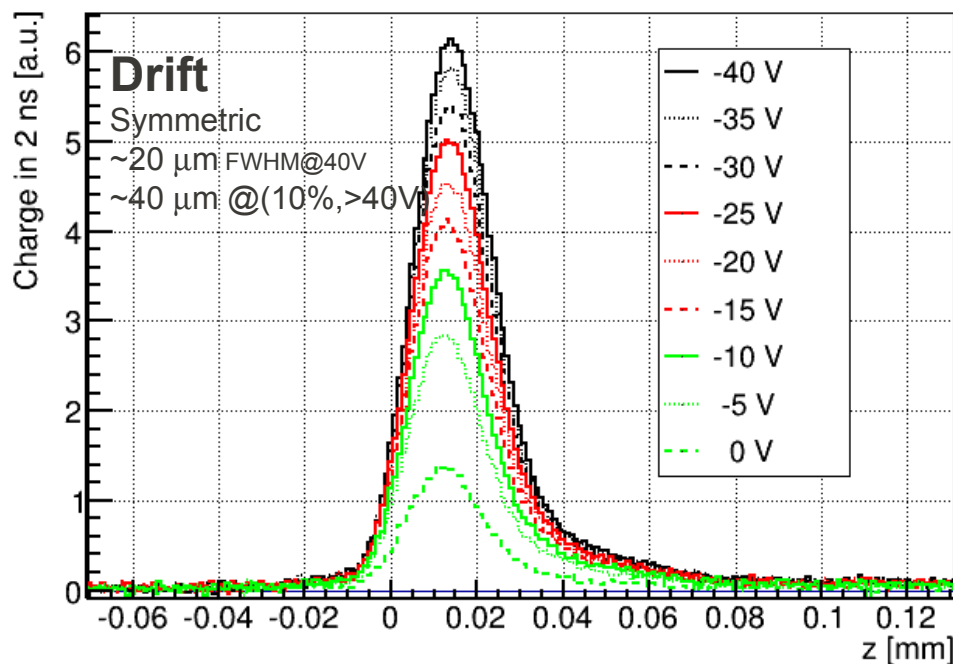


## Zoom on "fast" signal



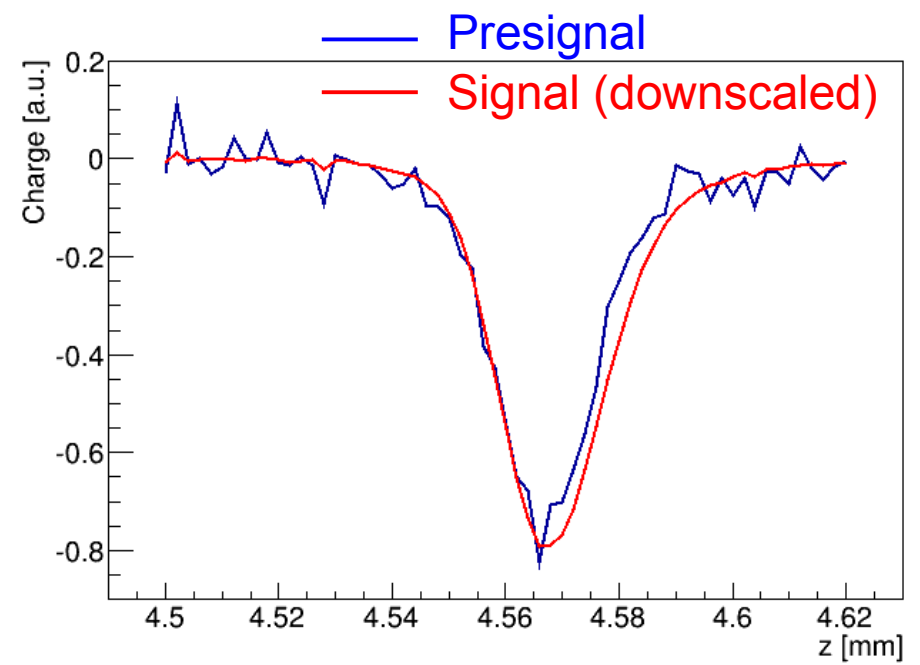
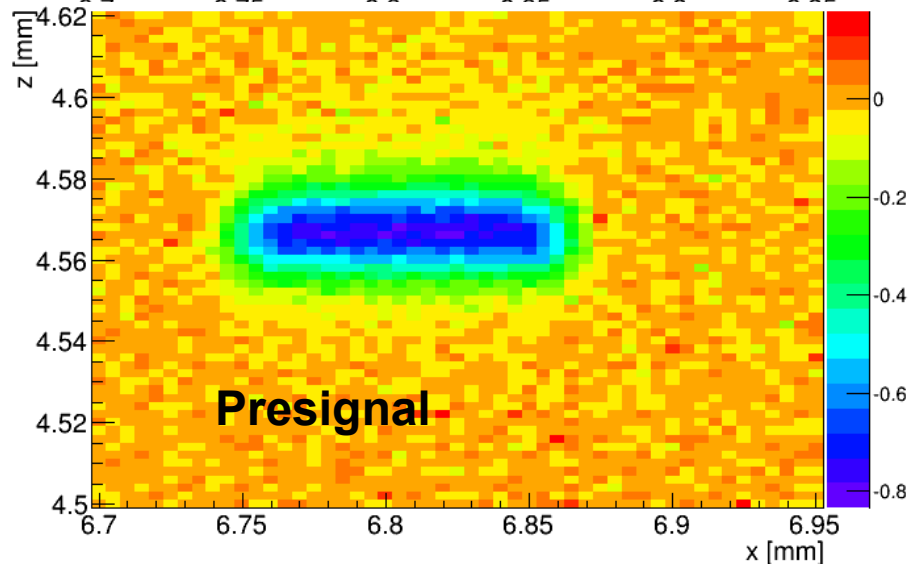
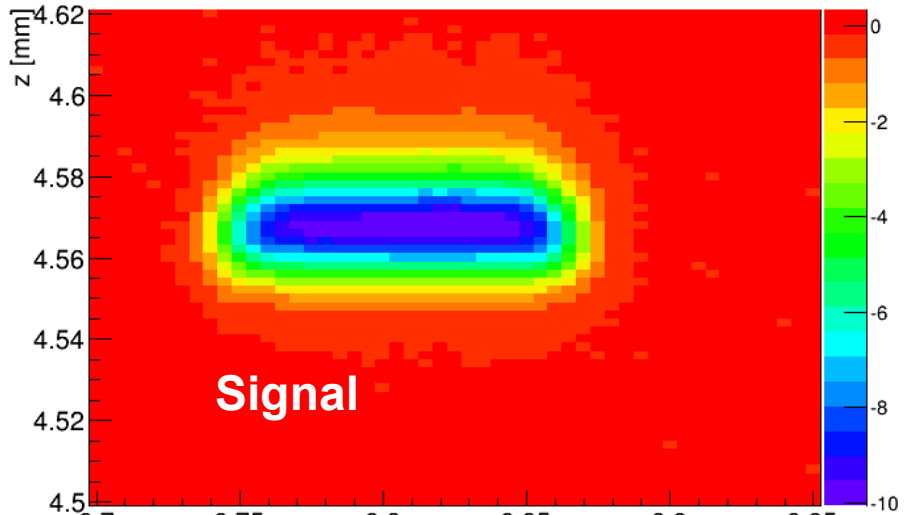
Drift peak does not shift with depth  
(see strip edge-TCT before for comparison)

## Diffusion $\Rightarrow$ charge profile c.o.g. is shifted

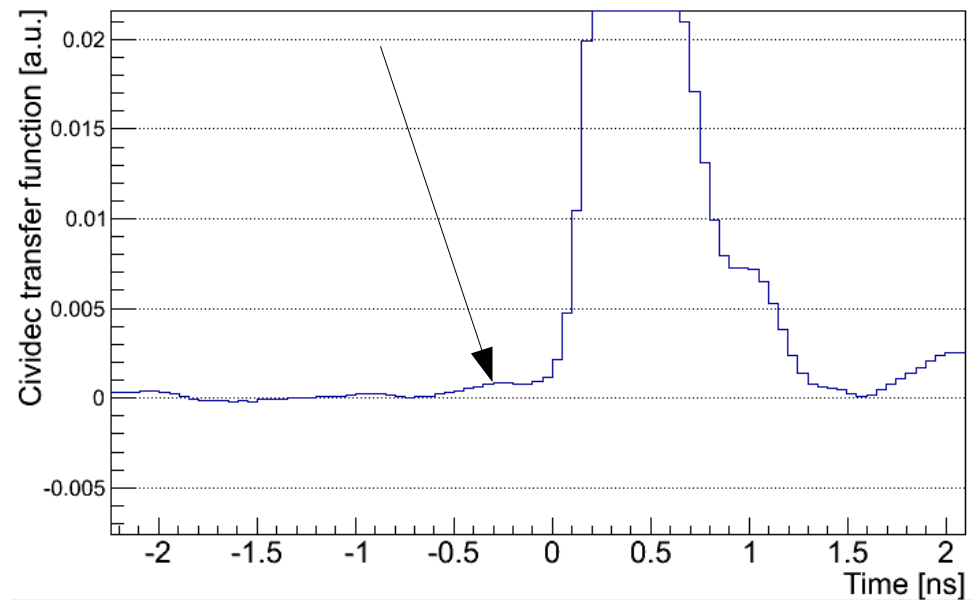


# Presignal:

Seen in conf 2 (diode), not in conf 1

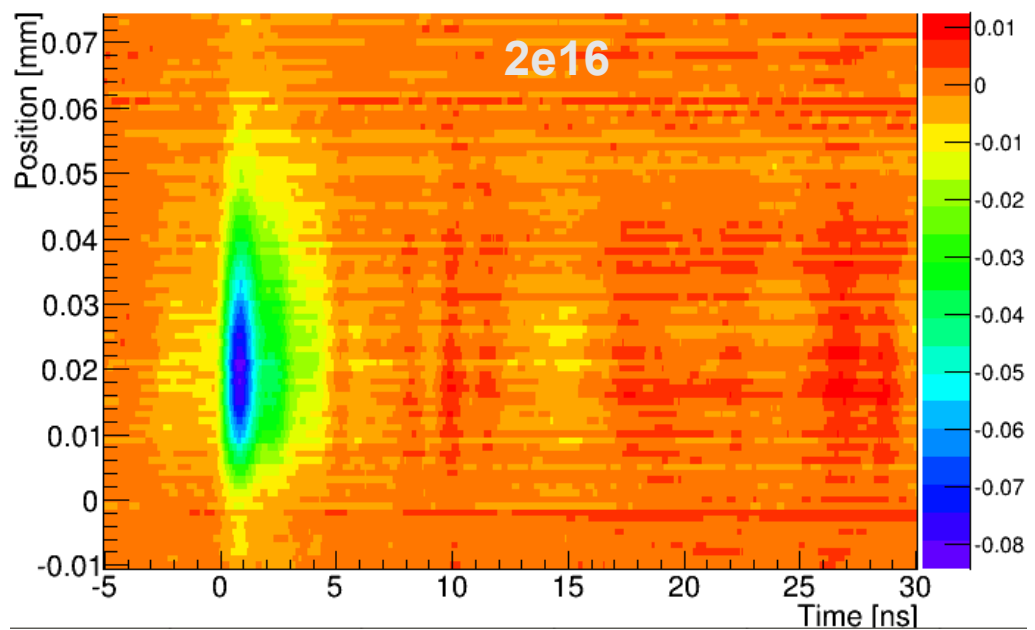
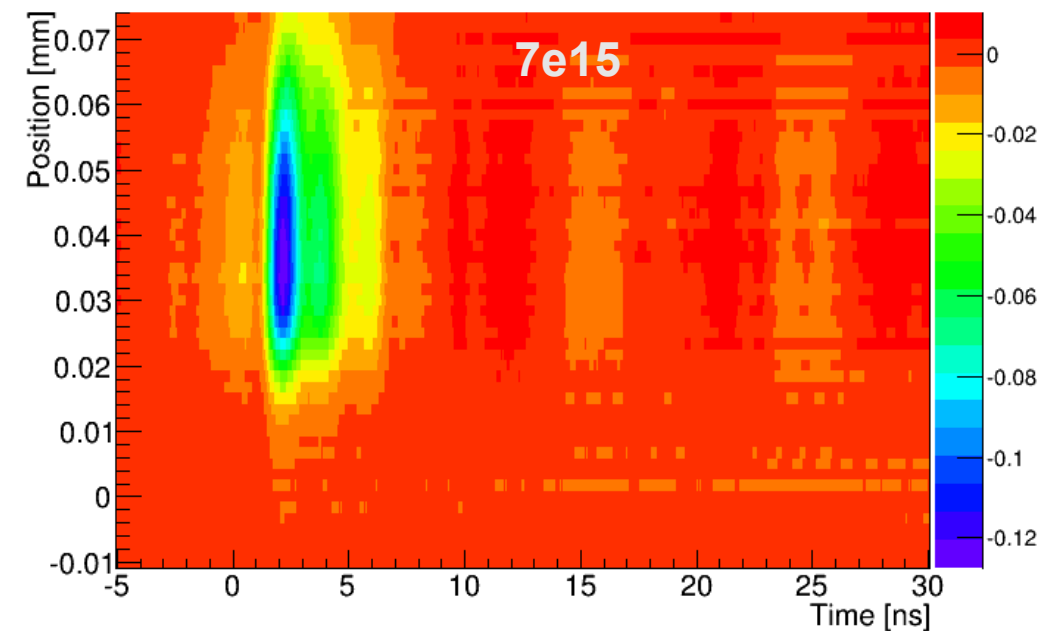
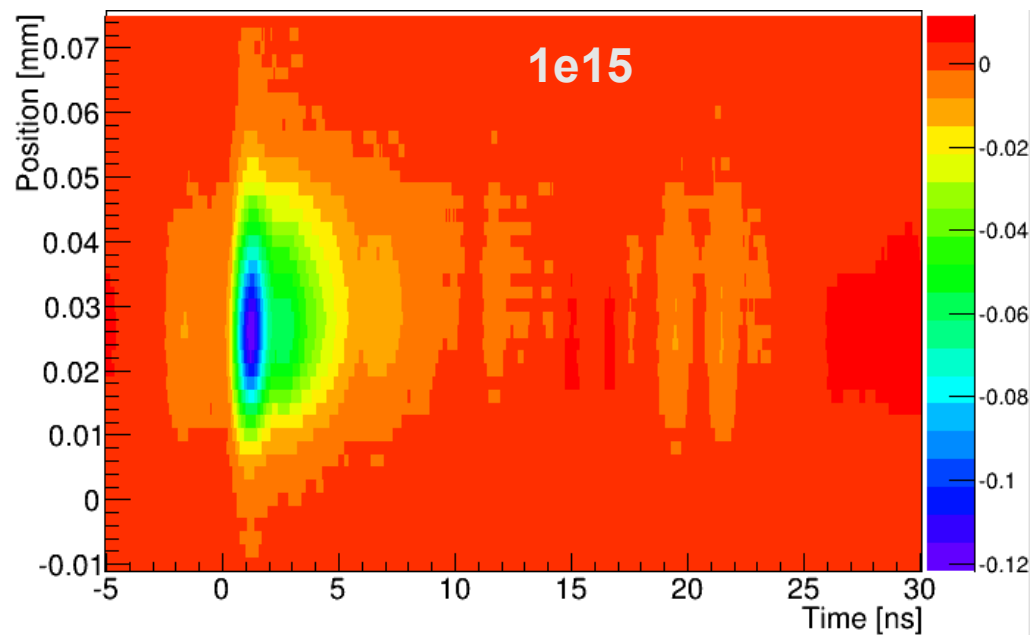
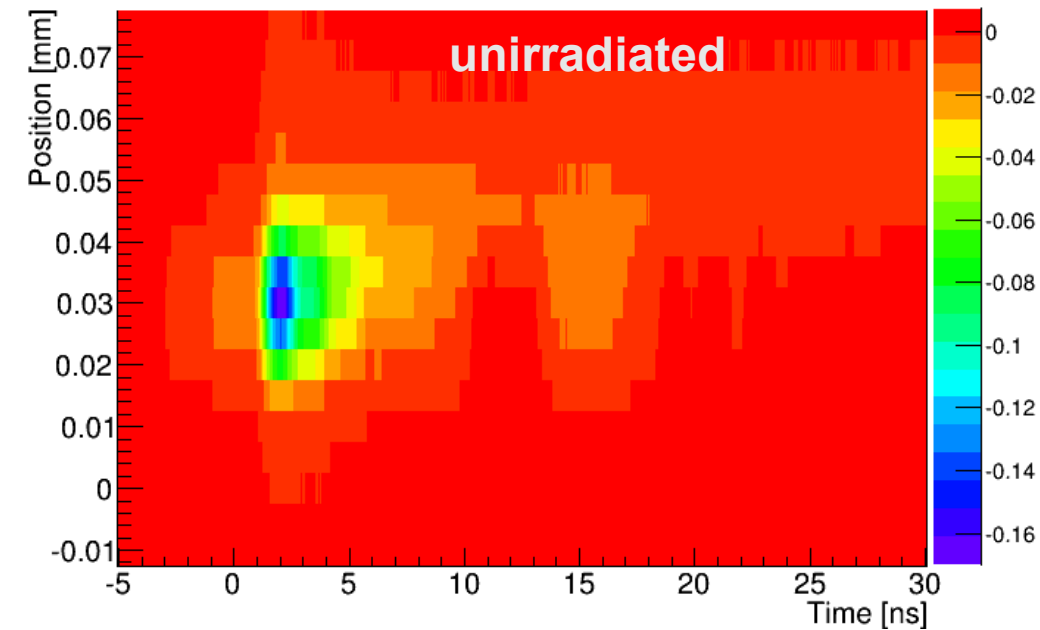


Seems as “feature” of the amplifier.



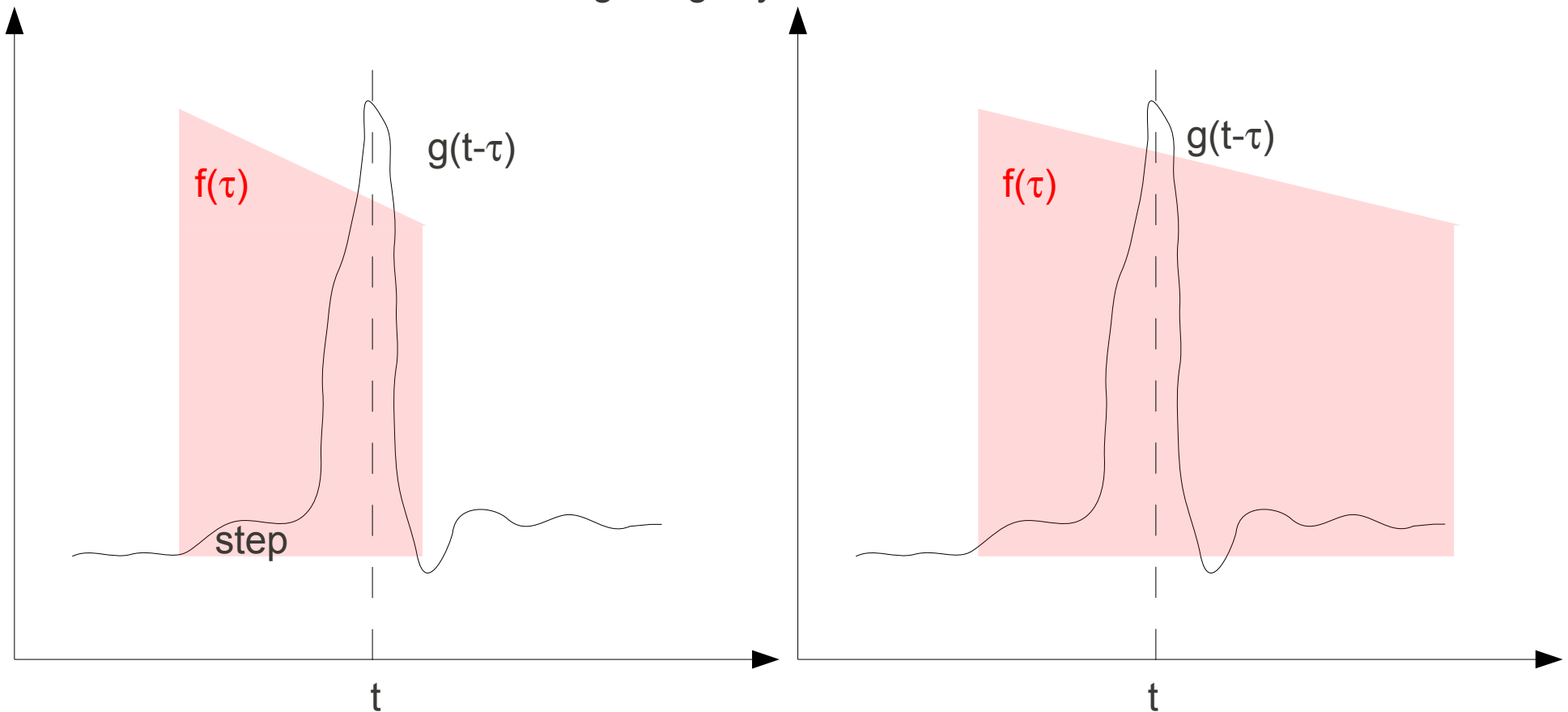
Bulk~10 Ohm.cm  
Implant~

# Diode readout, -60V, signal maps



$$(f * g)(t) \stackrel{\text{def}}{=} \int_{-\infty}^{\infty} f(\tau) g(t - \tau) d\tau$$

f ≡ Ideal signal, g ≡ System Transfer Function



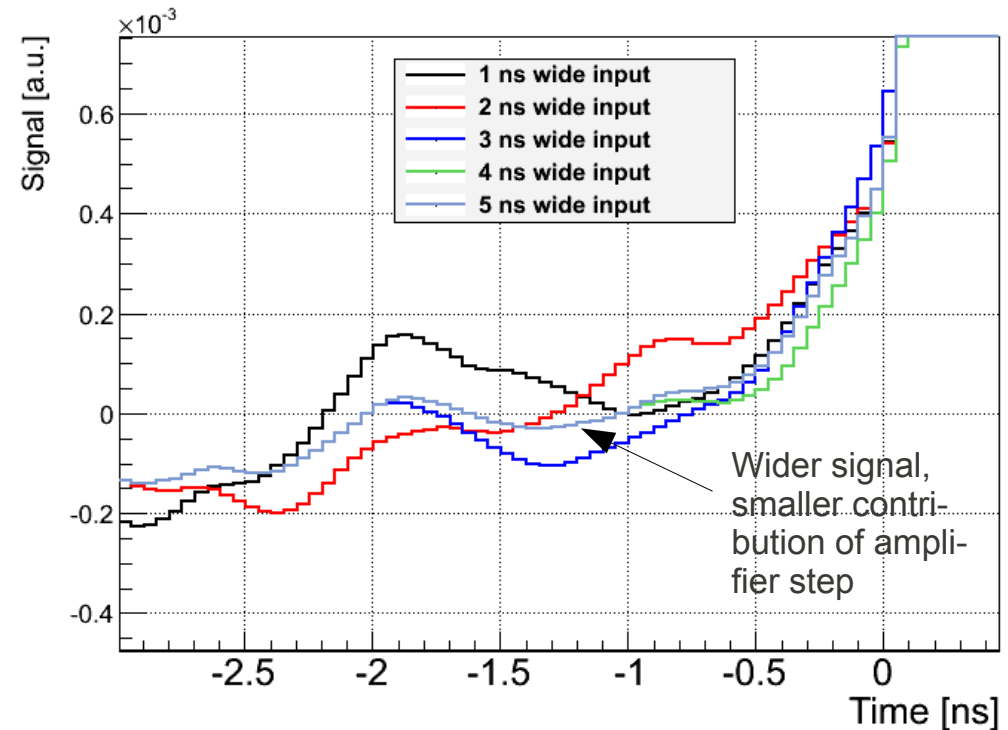
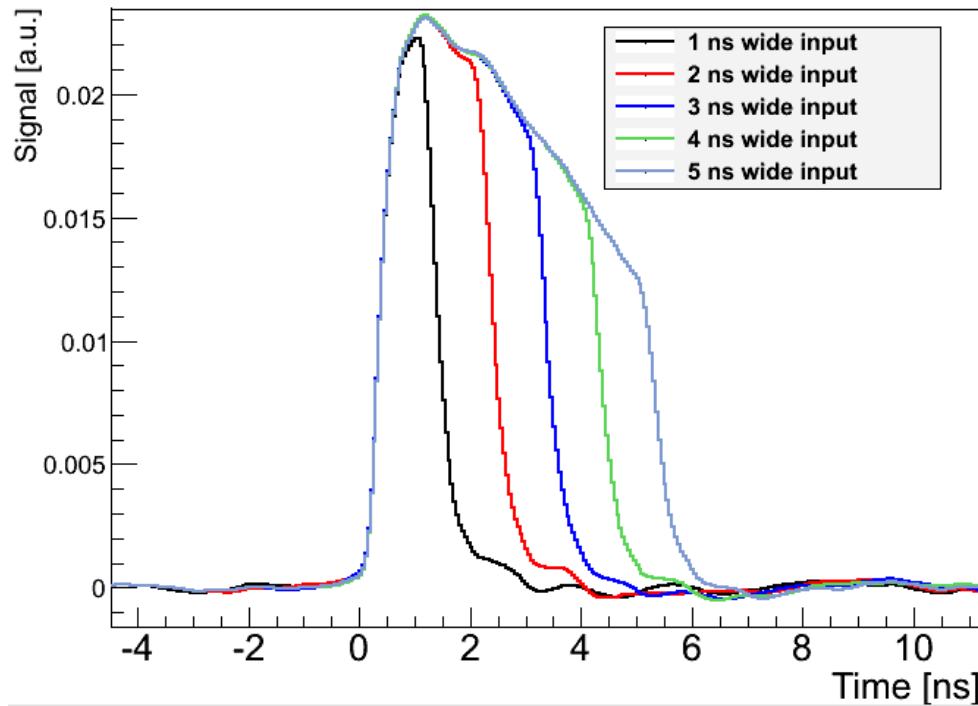
Convolution integral expands from  $(-\infty, \infty)$

Narrow signals: relative contribution of amplifier step, at point “t”, is higher for small signals

Wide signals: at the same position, the integral from  $(-\infty, \infty)$  adds up more signal, and the step contribution is smaller

$$(f * g)(t) \stackrel{\text{def}}{=} \int_{-\infty}^{\infty} f(\tau) g(t - \tau) d\tau$$

f ≡ Ideal signal, g ≡ System Transfer Function

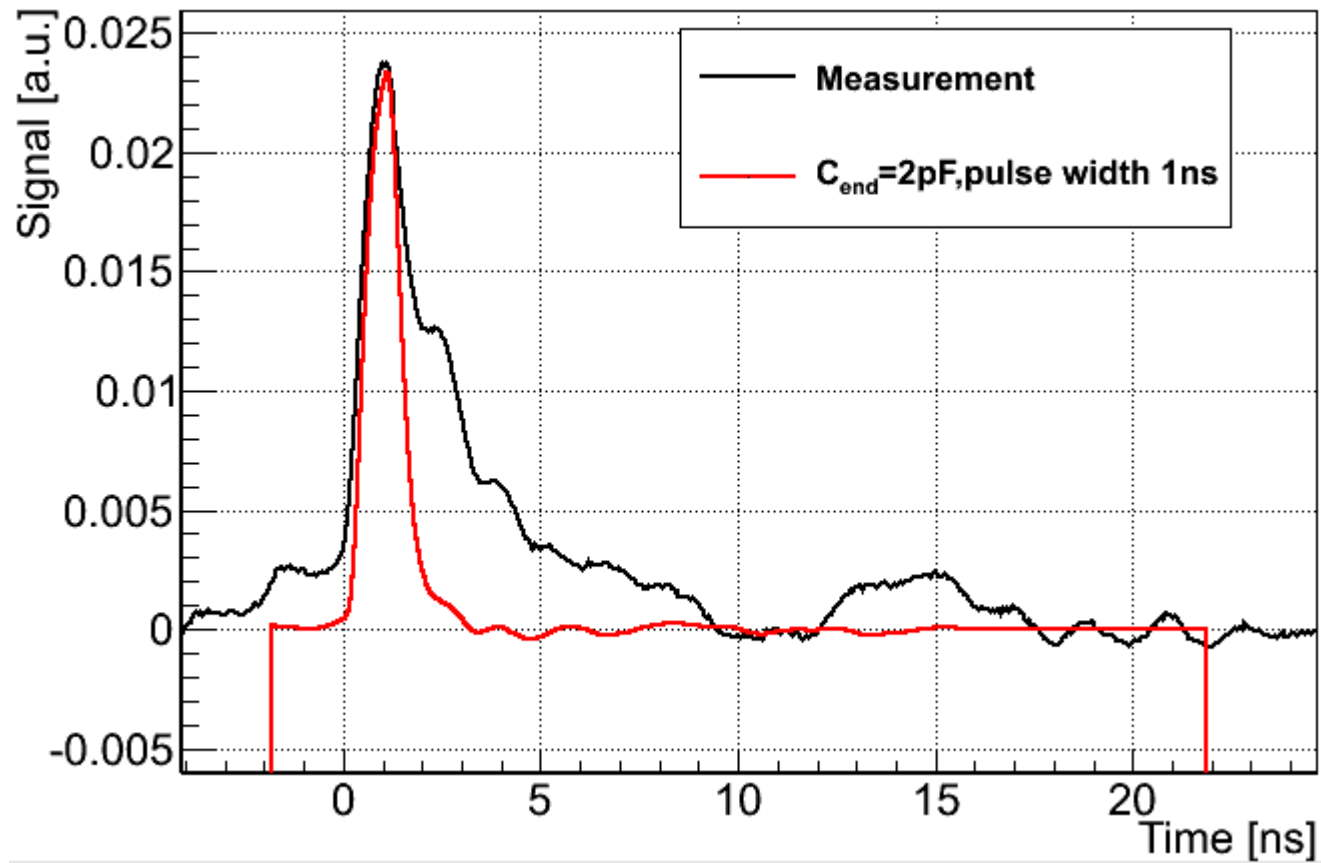


Convolution integral expands from  $(-\infty, \infty)$

Narrow signals: relative contribution of amplifier step, at point “t”, is higher for small signals

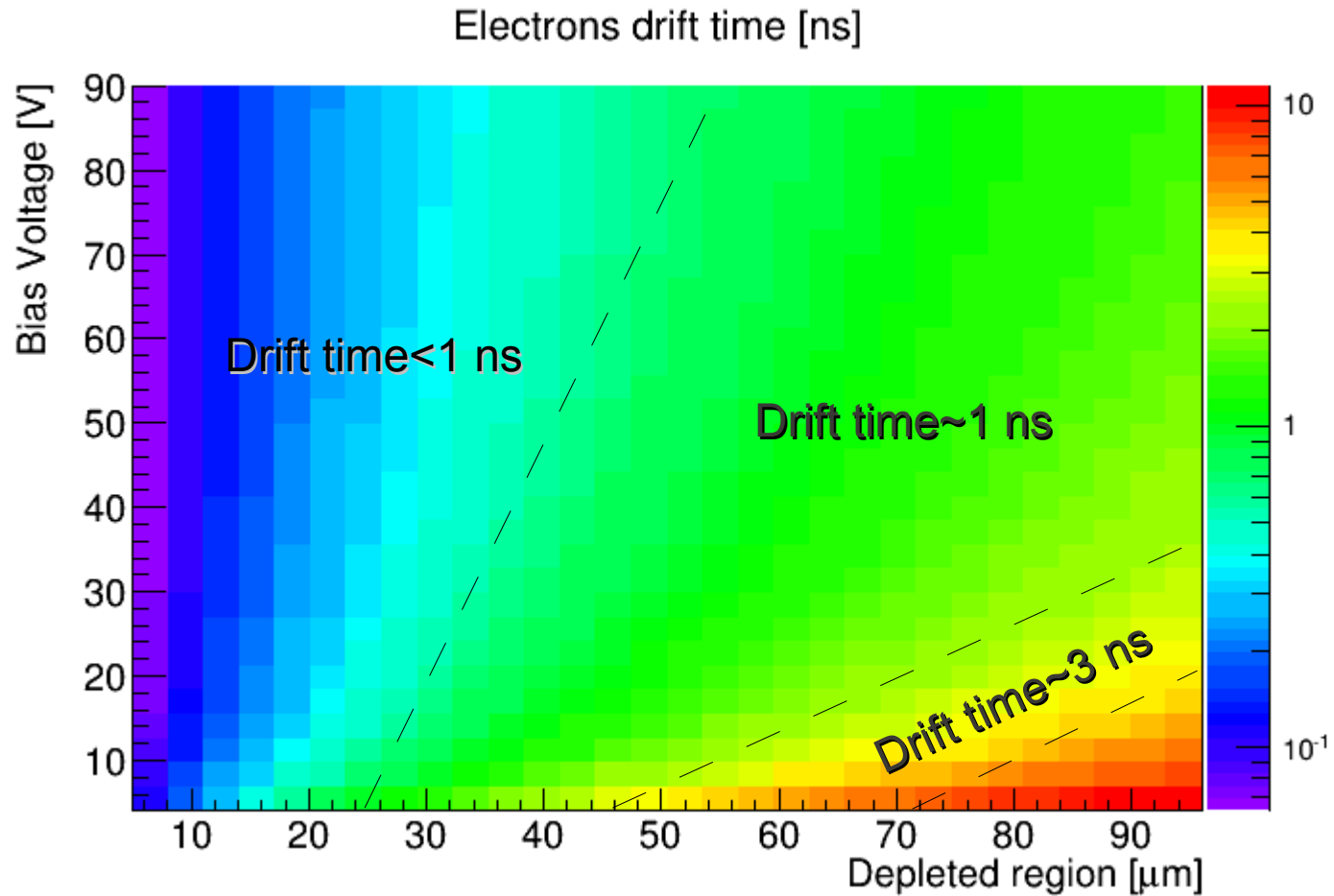
Wide signals: at the same position, the integral from  $(-\infty, \infty)$  adds up more signal, and the step contribution is smaller

## Comparison **calculated signal** vs measured (conf 2)

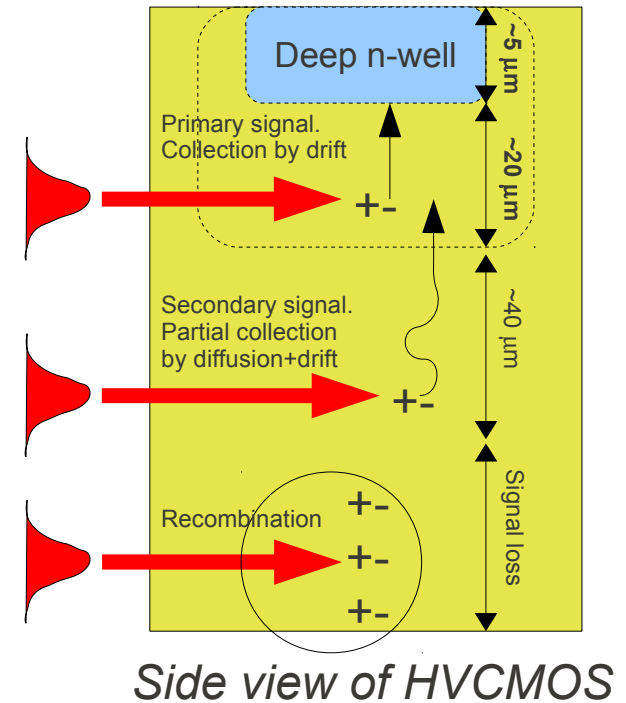
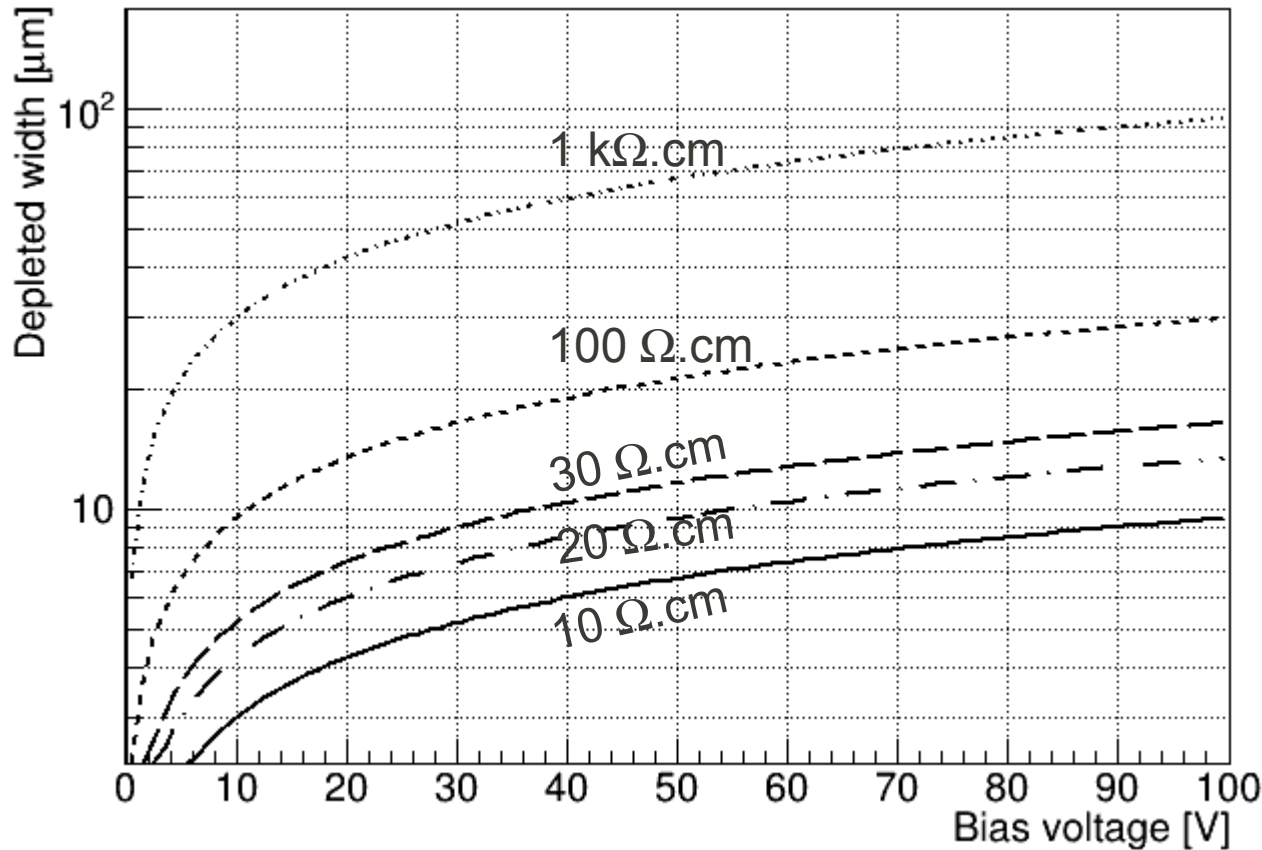


Capacitance is a second order effect in the width of the pulse.  
Oscillations coming from the PCB

Collection time needed for electrons to traverse  $X \mu\text{m}$  of depleted region at given bias voltage



# Depletion width in p-type material (formula from Spieler, page 66)



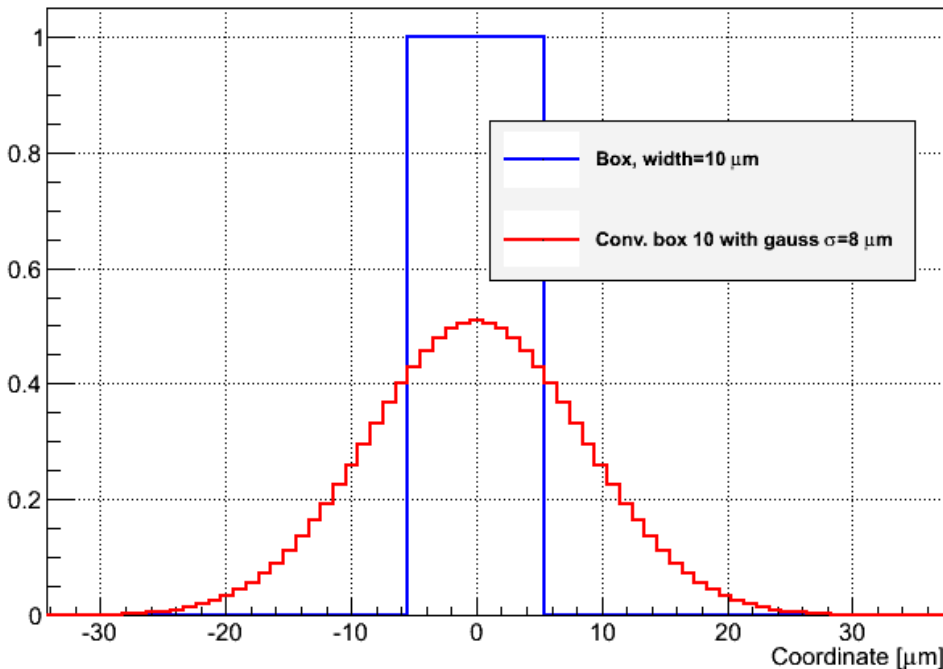
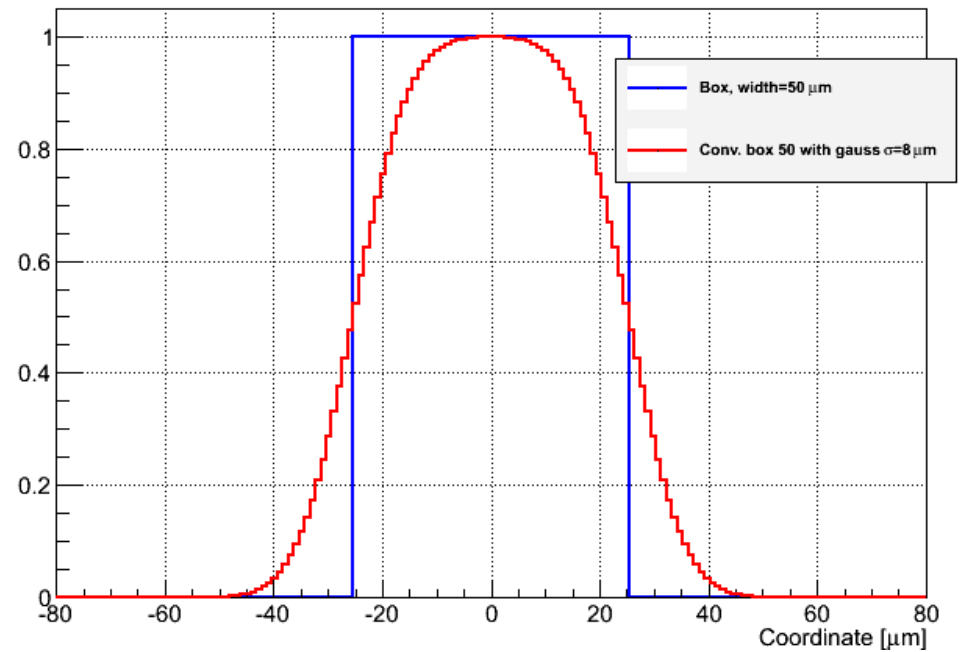
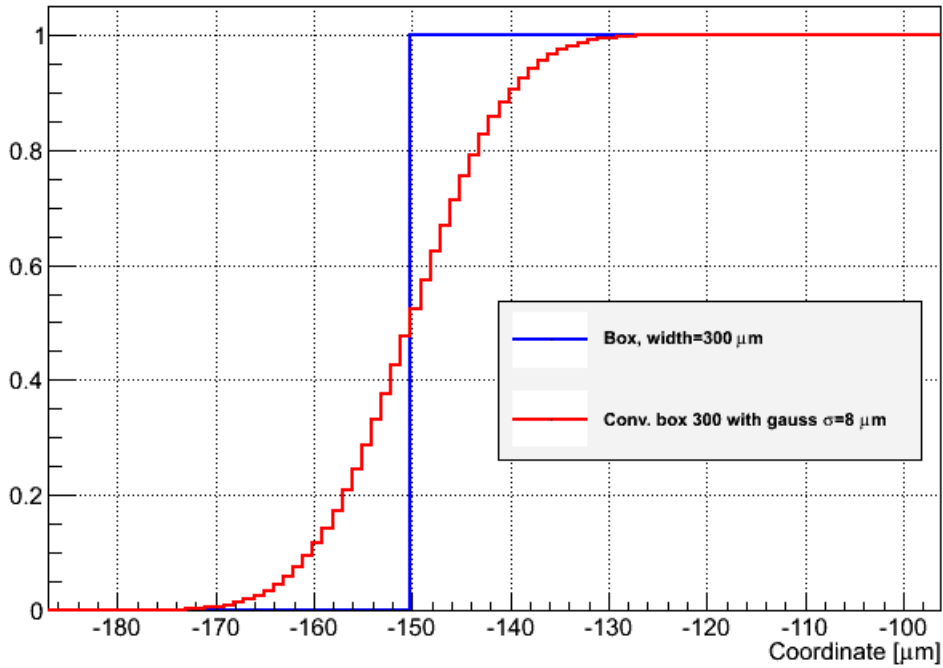
For 10 Ω.cm, the expected depletion voltage is  $V_{dep} \sim 1.1 \cdot (\text{thickness}[\mu\text{m}])^2$   
 For 10 Ω.cm we expect  $\sim 10 \mu\text{m}$  depleted bulk at 90 V

10 Ω.cm =  $1.4 \times 10^{15}$  (p-type)

| Nbulk |        | P bulk |         |
|-------|--------|--------|---------|
| cm-3  | Ω.cm   | cm-3   | Ω.cm    |
| 1e12  | 4415.3 | 1E12   | 13267.4 |
| 1e13  | 442.0  | 1E13   | 1327.5  |
| 1e14  | 44.5   | 1E14   | 133.1   |
| 1e15  | 4.6    | 1E15   | 13.5    |
| 1e16  | 0.5    | 1E16   | 1.45    |
| 1e17  | 0.09   | 1E17   | 0.197   |
| 1e18  | 0.023  | 1E18   | 0.041   |
| 1e19  | 0.005  | 1E19   | 0.009   |



# Effect of convoluting gaussian signal with depleted box-like region



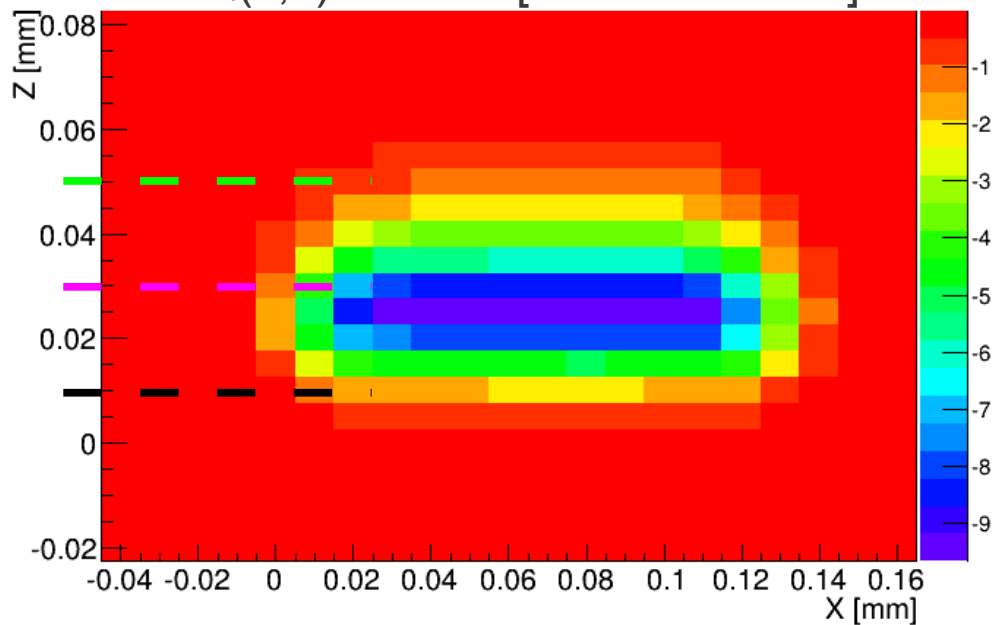
For “thick” detectors ( $\sim 300 \mu\text{m}$ ), signal is collected already 30 μm before the beam center enters the detector

If thickness  $< 50 \mu\text{m}$ , signal collection starts  $\sim 15 \mu\text{m}$  before

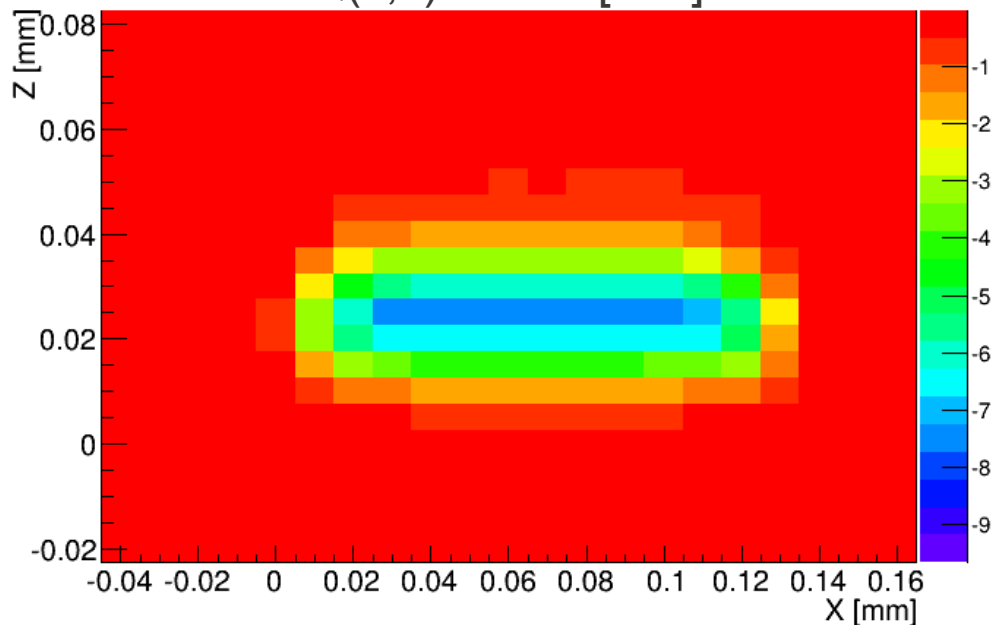
To have an active region  $\sim 80$  (40) μm as measured with the laser, the box thickness is 50(10) μm

# Unirradiated HVCMOS, diode readout

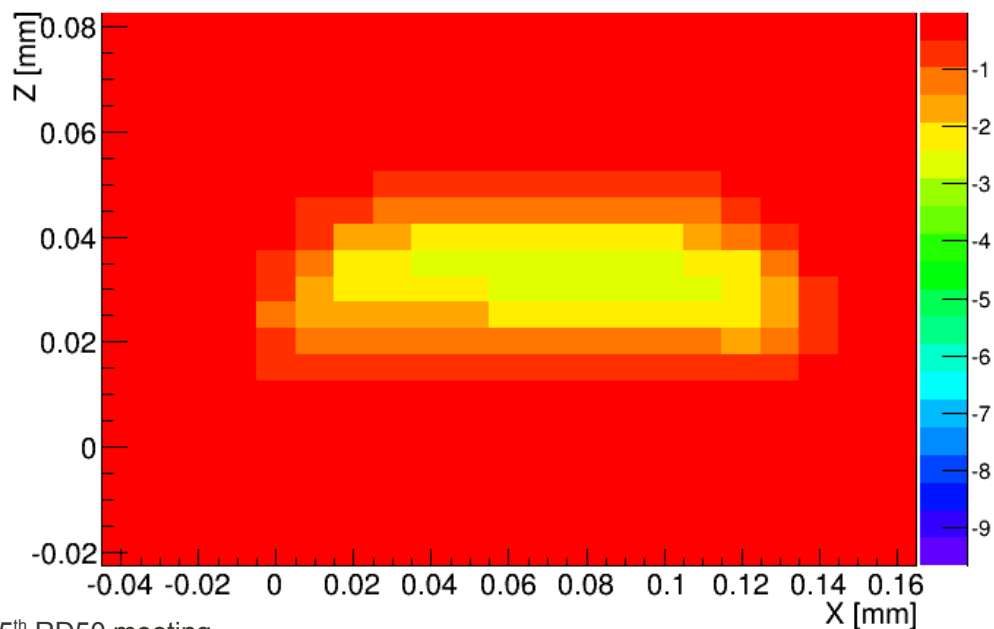
Q(x,z) t < 10 ns [drift+ "diffusion"]



Q(x,z) t < 5 ns [drift]

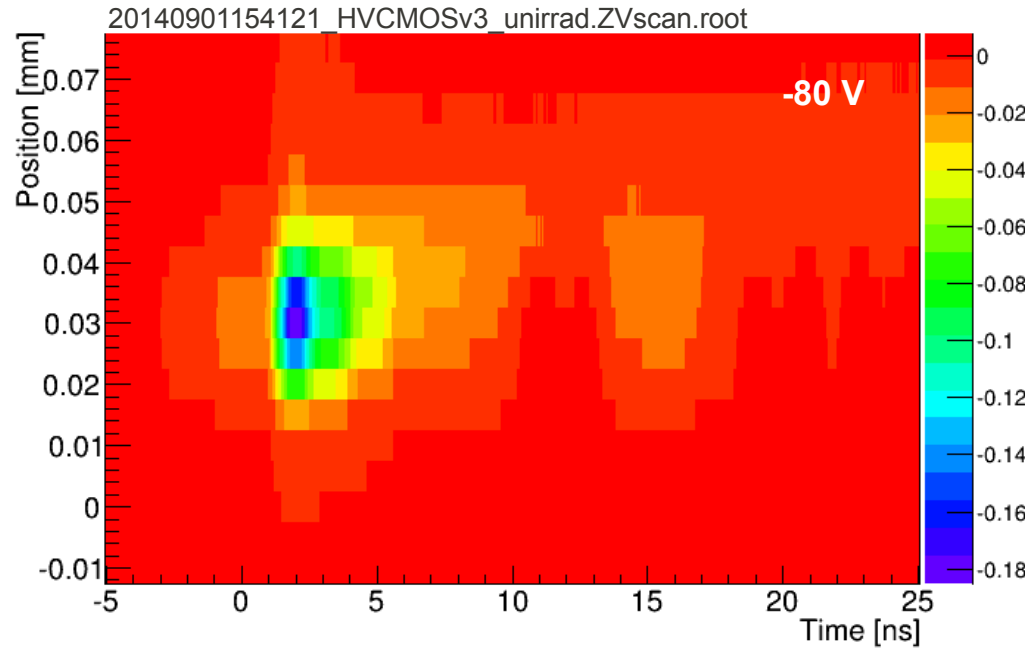


Q(x,z) t in [5-10] ns [ "diffusion" ]

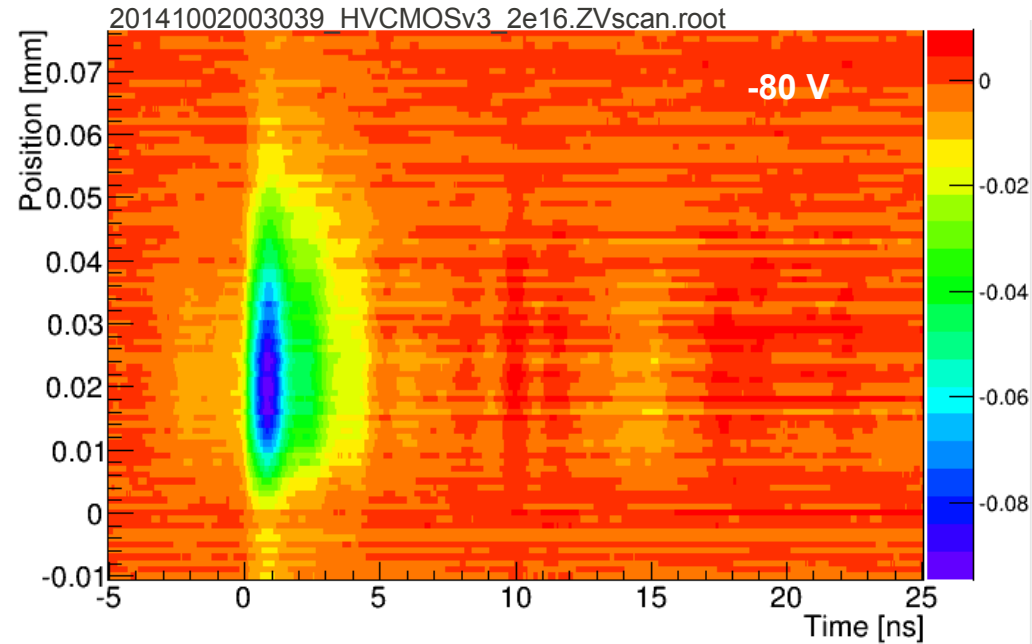


Relative contributions of drift and diffusion depend on the threshold by eye that we set.

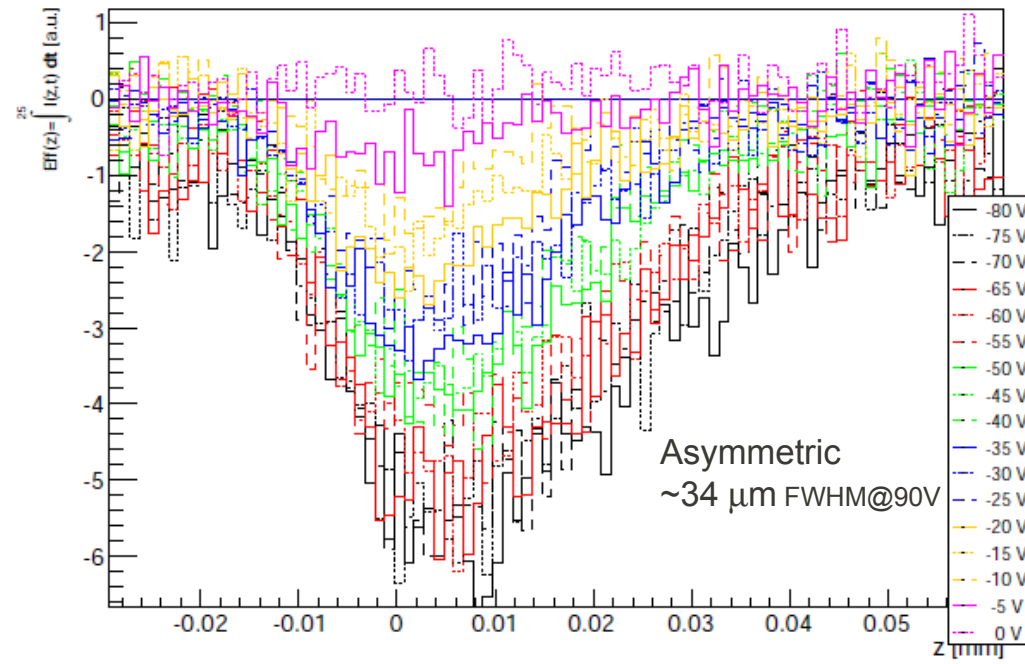
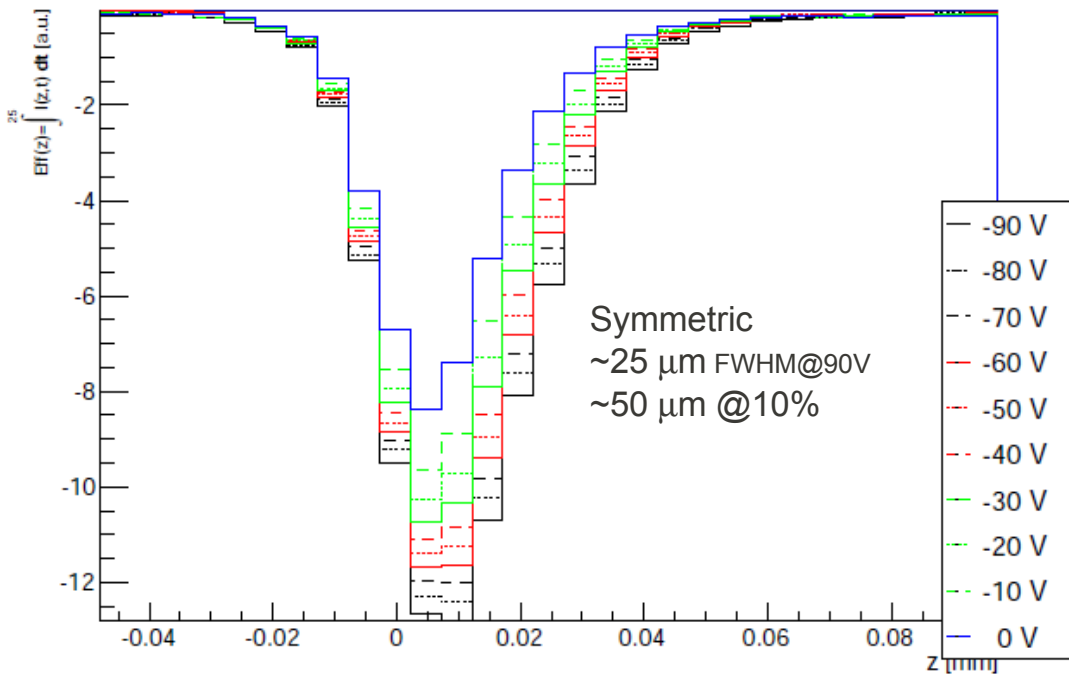
# Diode readout (conf. 2): depleted region



Some diffusion still visible, but mainly signal coming from drift

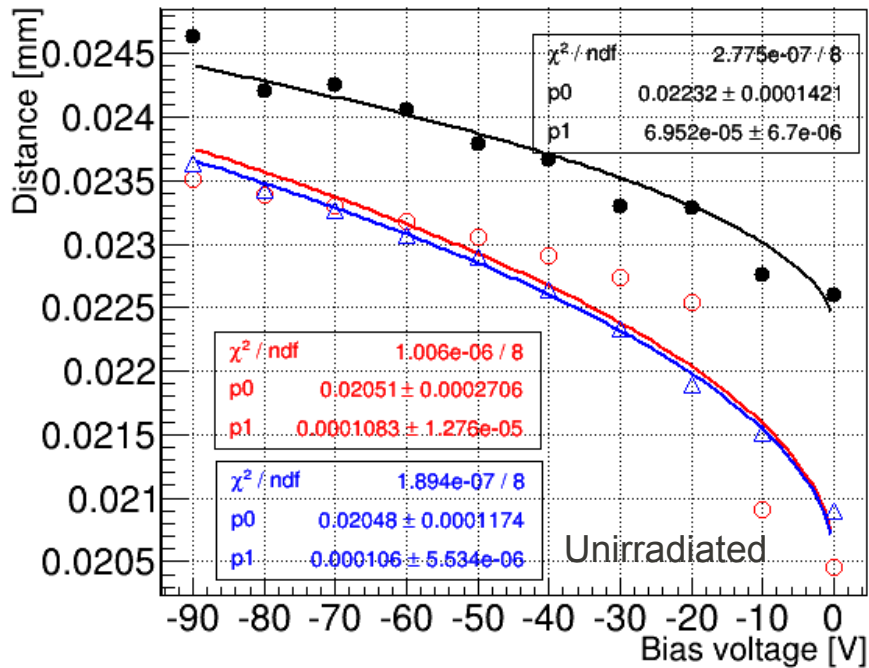


Diffusion contribution is gone. Efficient region for charge production seems bigger

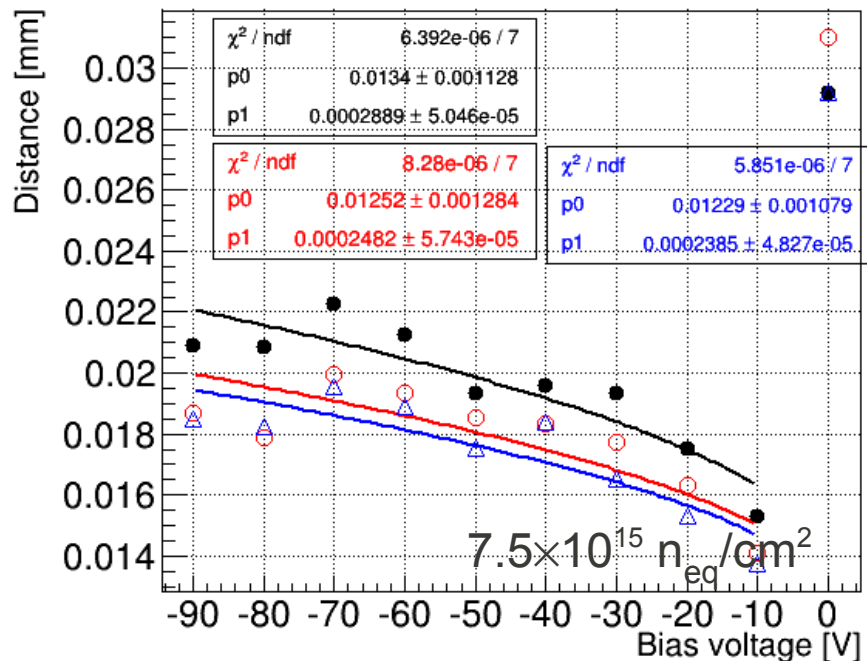
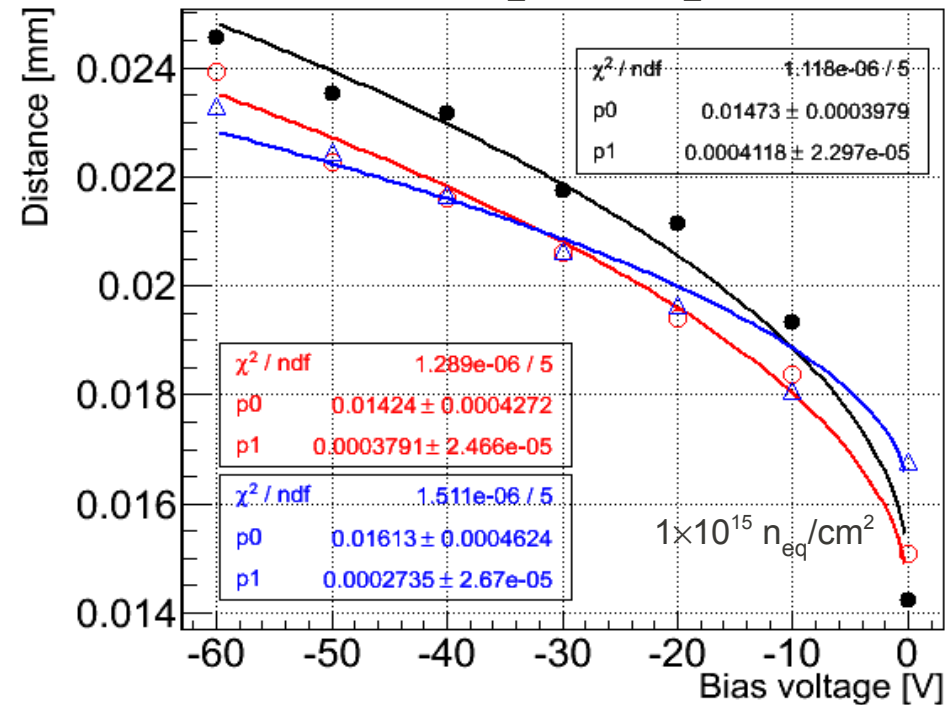


# Center of depleted region

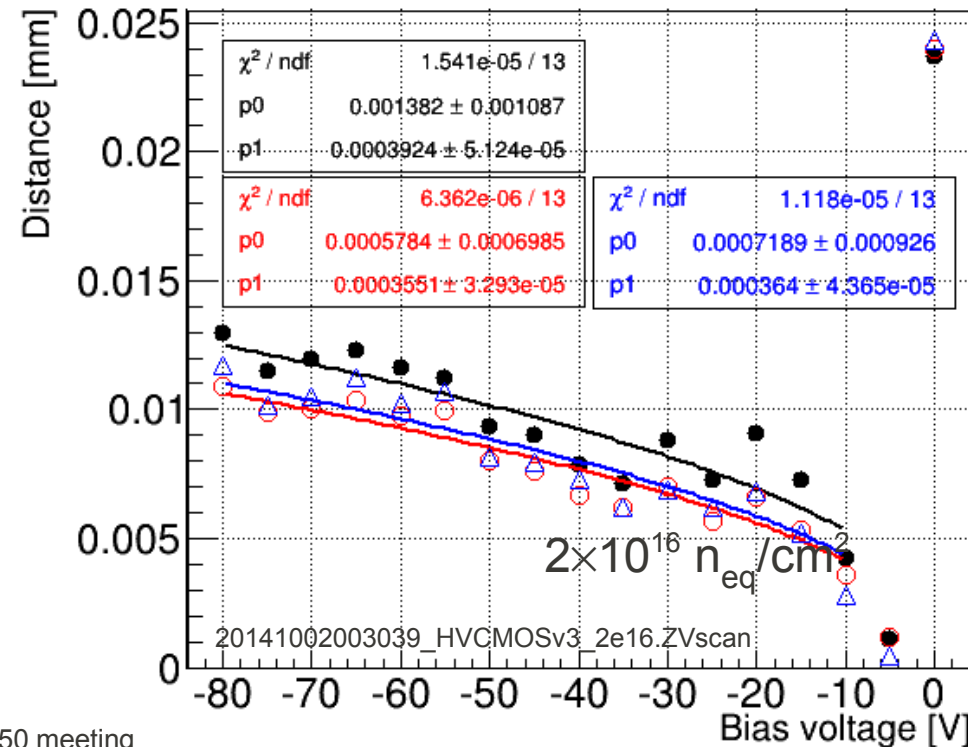
20140901154121\_HVCMOSv3\_unirrad.ZVscan



20140904145059\_HVCMOSv3\_irrad2.ZVscan

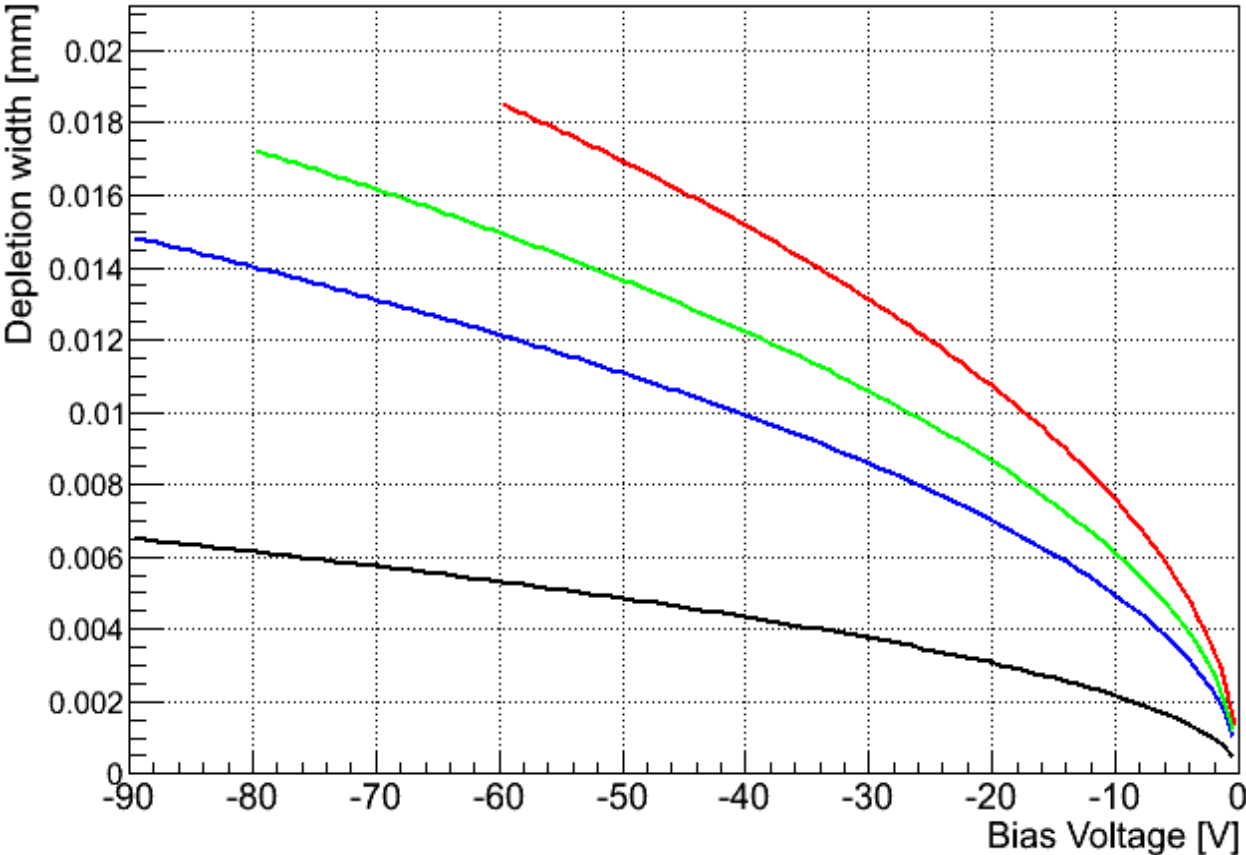


20140913123542\_HVCMOSv3\_7e15.ZVscan



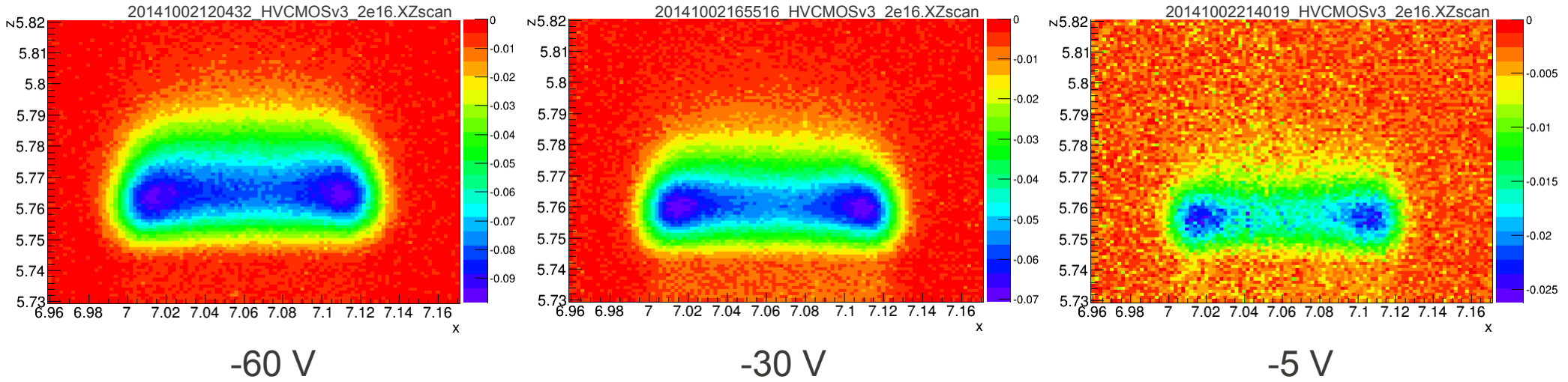
20141002003039\_HVCMOSv3\_2e16.ZVscan

# Calculated depletion width from fit to center of depleted region

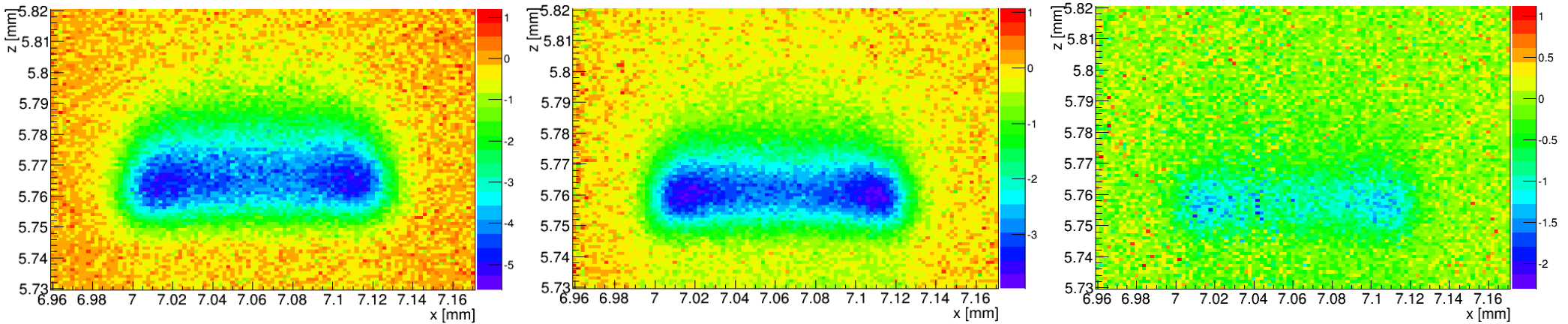


# 2D scan (XZ) at different bias voltages, $2 \times 10^{16}$ neq/cm<sup>2</sup>, configuration 2

## Signal amplitude

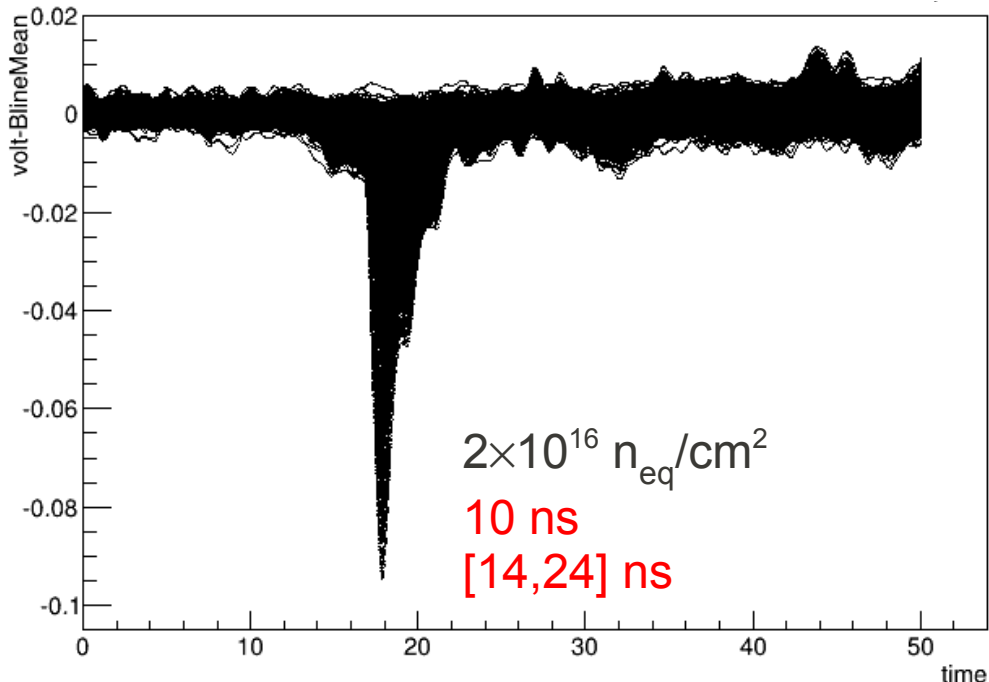
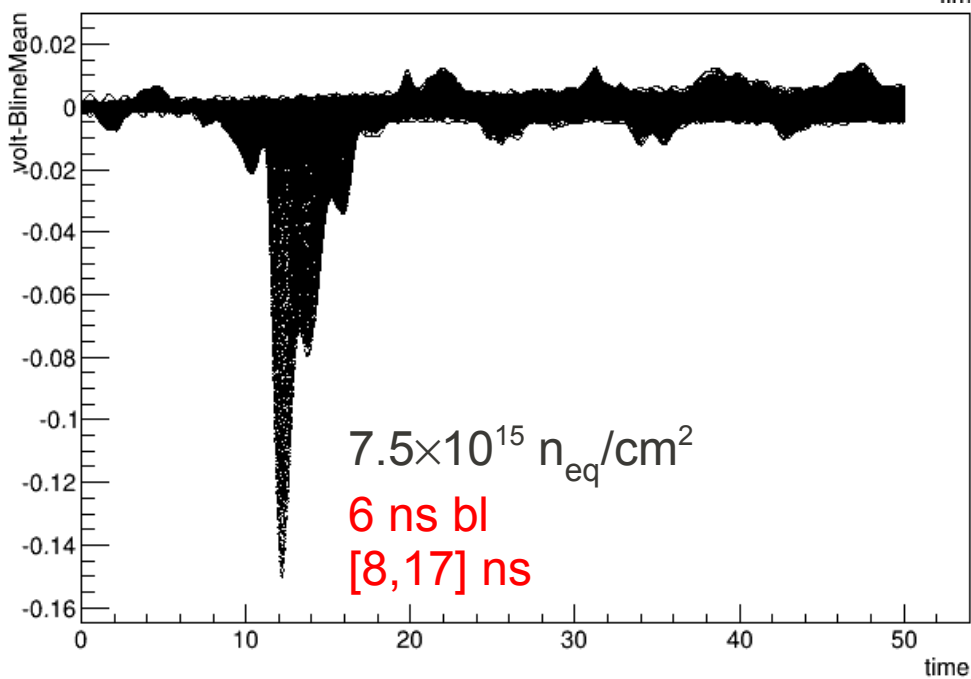
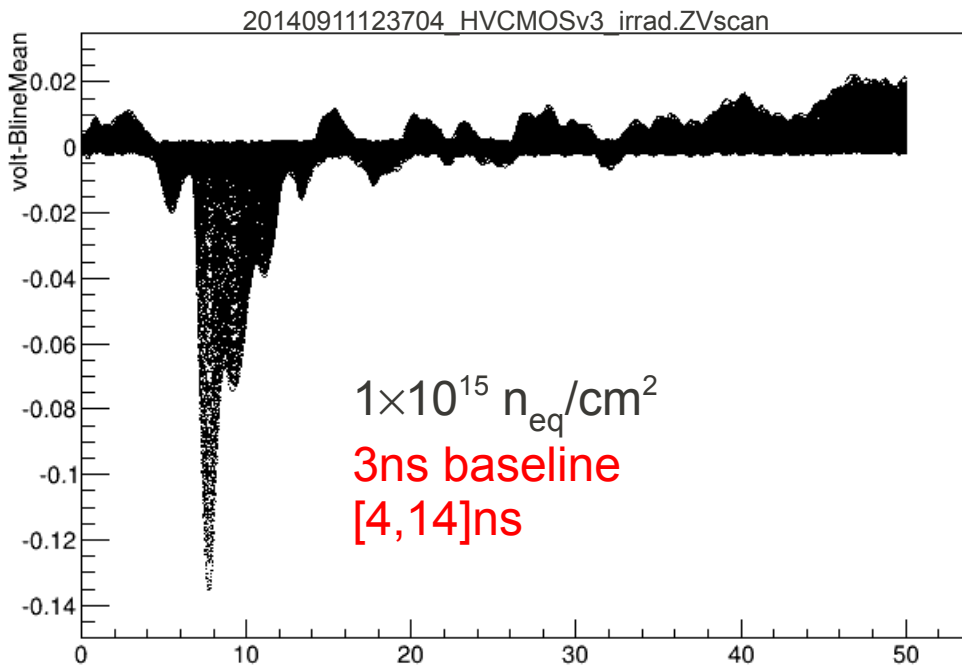
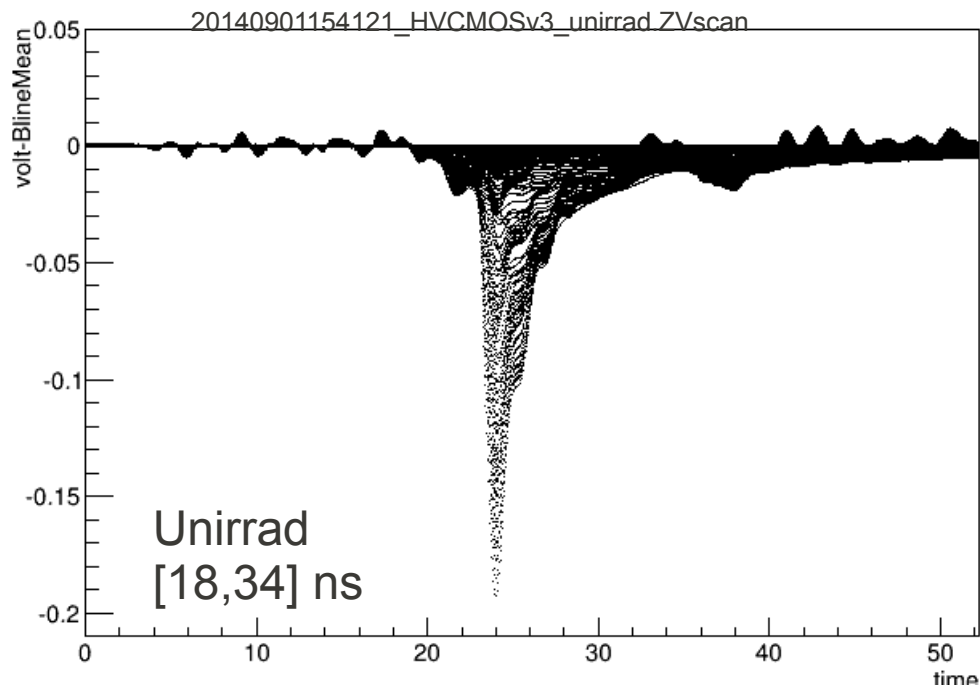


## Charge in 10 ns



Diode looks symmetric in conf 1

# All signals, all voltages [-90,0], conf. 2 (diode)

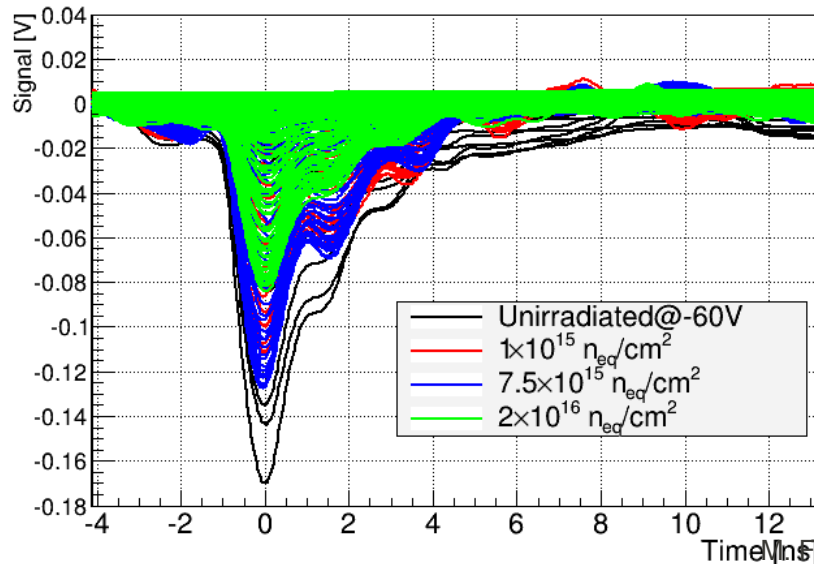
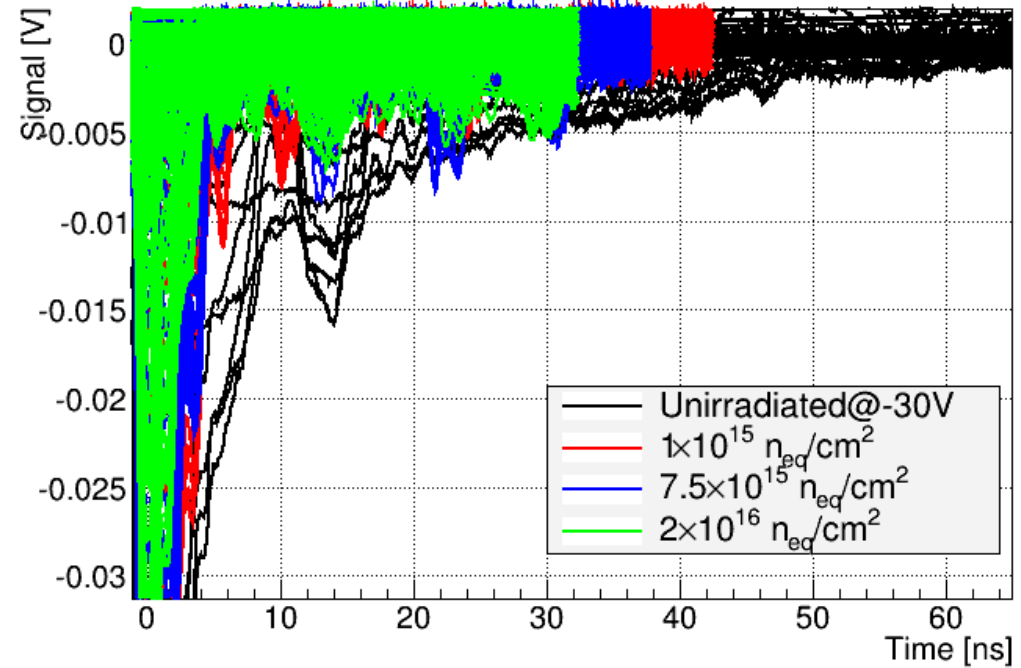
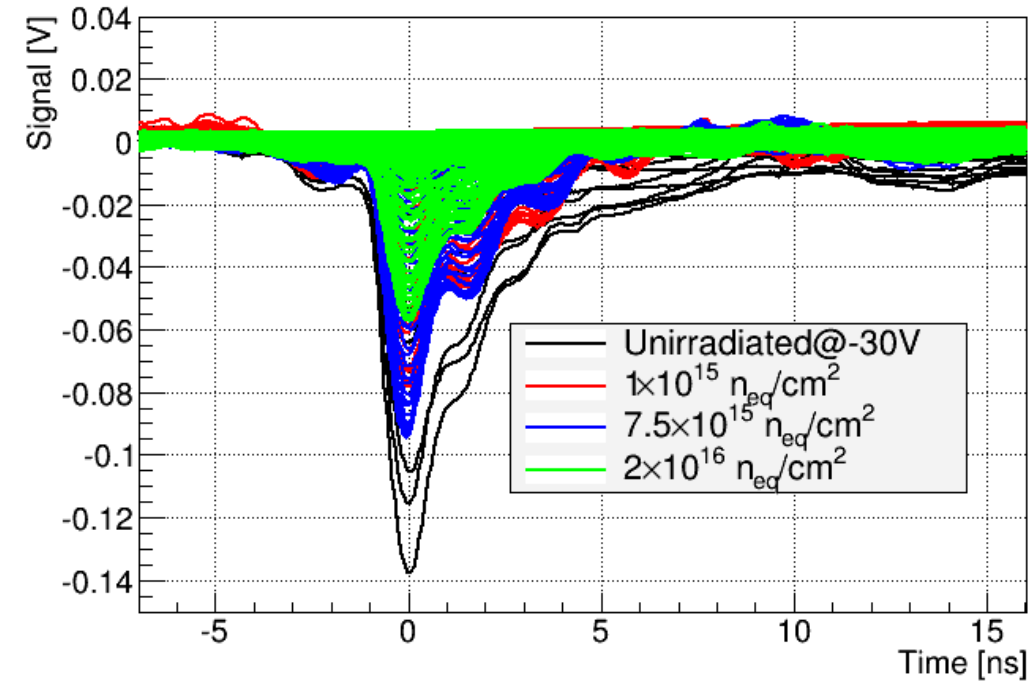


20140913123542\_HVCMOSv3\_7e15.ZVscan

20141002003039\_HVCMOSv3\_2e16.ZVscan

## All in one plot (diode configuration)

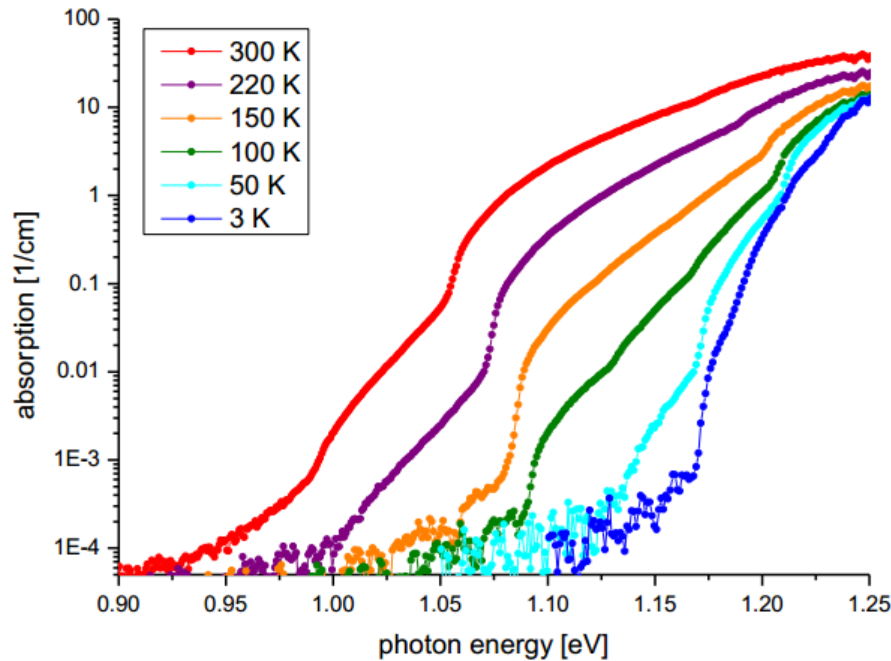
Diode configuration catches mainly drift. For unirradiated detector diffusion still visible (1 order of magnitude smaller than drift effect)







# Temperature dependent absorption



- for lower temperatures phonons freezing out (beginning with the higher energy phonons) → absorption below the gap is getting smaller

$\alpha = \frac{4\pi k}{\lambda}$  → K decreases → optical absorption decreases → less signal collected