On homogeneous and isotropic Universe

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Notation

$$\begin{split} (\mathbb{M},g) &- \text{geodesically complete (pseudo)Riemannian manifold (space-time)} \\ g_{\alpha\beta}(t,x), \quad \text{sign } g_{\alpha\beta} = (+---) &- \text{Lorentzian signature metric} \\ \left\{ x^{\alpha} \right\} &= \left\{ t, x^{\mu} \right\}, \quad \alpha = 0, 1, 2, 3; \quad \mu = 1, 2, 3 \quad \text{- coordinates} \end{split}$$

<u>Theorem</u>. Let a four dimensional space-time be the topological product $\mathbb{M} = \mathbb{R} \times \mathbb{S}$ where $t \in \mathbb{R}$ is a time coordinate $x \in \mathbb{S}$ is a three dimensional space of constant curvature. We suppose that sufficiently smooth homogeneous and isotropic metric of Lorentzian signature is given on it. Then there is such coordinate system t, x^{μ} in some neighbourhood of each point where metric has the form

(1)
$$ds^{2} = dt^{2} + a^{2}\hat{g}_{\mu\nu}dx^{\mu}dx^{\nu},$$

the Friedmann metric

where a(t) > 0 is an arbitrary function (scale factor) and $\hat{g}_{\mu\nu}(x)$ is a negative definite three dimensional constant curvature metric on S depending only on space coordinates.

Brief History

A. Friedmann. Zs. Phys. 10(1922)377; ibid. 21(1924)326.

The first papers were metric (1) was used in cosmology. No theorem.

G. Lemaitre. Ann. Soc. Sci. (Bruxelles) 47A(1927)49; ibid. A53(1933)51. Metric (1) was used in several cosmological models. No theorem.

H. P. Robertson. *Proc. Nat. Acad. Sci.* 15(1929)822; *Rev. Mod. Phys.* 5(1933)62; *Ap. J.* 82(1935)284.

> In the first two papers, the theorem was formulated but not proved. Instead, he referred to: D. Hilbert. *Mathematische Annalen,* 15(1924)1. The first part of the proof is given in a general case.

> G. Fubini. Annali di Matematica, Pura Appl.[3], 9 (1904) 33. The second part of the proof is given in one way. In the third paper, metric (1) was obtained in a different way by considering a set of observers with given properties.

R. C. Tolman. *Proc. Nat. Acad. Sci.* 16(1930) 320, ibid. 409, ibid. 511; Metric (1) was obtained from different assumptions. In particular he assumed spherical symmetry and used Einstein's equations.

A. G. Walker. Proc. London Math. Soc. Ser.2 42(1936) 90.

The theorem was proved in one way.

For proof, see, for example, S. Weinberg. Gravitation and cosmology.1972

Constant curvature space S

 $\hat{K} = -1, 0, 1$ - Gaussian curvature

Spherical coordinates on $\ensuremath{\mathbb{S}}$:

$$\hat{g}_{\mu\nu}dx^{\mu}dx^{\nu} = \begin{cases} d\chi^2 + \sin^2\chi \left(d\vartheta^2 + \sin^2\vartheta d\varphi^2 \right), & \hat{K} = 1, \\ d\chi^2 + \chi^2 \left(d\vartheta^2 + \sin^2\vartheta d\varphi^2 \right), & \hat{K} = 0, \\ d\chi^2 + \sinh^2\chi \left(d\vartheta^2 + \sin^2\vartheta d\varphi^2 \right), & \hat{K} = -1 \end{cases}$$

$$ds^2 = dt^2 + a^2 \hat{g}_{\mu\nu} dx^\mu dx^\nu$$

Stereographic coordinates on S:

$$\hat{g}_{\mu\nu}dx^{\mu}dx^{\nu} = \frac{\eta_{\mu\nu}dx^{\mu}dx^{\nu}}{\left(1+b_{0}x^{2}\right)^{2}}$$

where:
$$\eta_{\mu\nu} \coloneqq \operatorname{diag}(---)$$

 $x^2 \coloneqq \eta_{\mu\nu} x^{\mu} x^{\nu} \le 0$
 $\hat{K} \coloneqq -12b_0$

$$x \in \mathbb{R}^{3}, \quad \hat{K} = 1,$$

$$x \in \mathbb{R}^{3}, \quad \hat{K} = 0,$$

$$x \in \mathbb{B}_{r}^{3}, \quad \hat{K} = -1,$$

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Coordinate transformation

$$g_{\alpha\beta} = \begin{pmatrix} 1 & 0 \\ 0 & \frac{a^2 \eta_{\mu\nu}}{\left(1 + b_0 x^2\right)^2} \end{pmatrix}$$

- Friedmann metric in stereographic coordinates

Coordinate transformation $x^{\mu} \mapsto \frac{x^{\mu}}{-}$

$$g_{\alpha\beta} = \begin{pmatrix} 1 + \frac{\dot{b}^2 x^2}{4b^2 (1 + bx^2)^2} & \frac{\dot{b}x_{\nu}}{2b (1 + bx^2)^2} \\ \frac{\dot{b}x_{\mu}}{2b (1 + bx^2)^2} & \frac{\eta_{\mu\nu}}{(1 + bx^2)^2} \end{pmatrix}$$

- Friedmann metric after coordinate transformation

where $b(t) := \frac{b_0}{a^2(t)}$ - function on time

$$\dot{b} := \frac{db}{dt}$$

- time derivative

Homogeneous and isotropic metric?

$$g_{\alpha\beta} = \begin{pmatrix} 1 & 0 \\ 0 & \frac{a^2 \eta_{\mu\nu}}{\left(1 + bx^2\right)^2} \end{pmatrix}$$

Changing topology of space in time: b(t) = 0

$$R_{\alpha\beta} - \frac{1}{2}g_{\alpha\beta}R = -\frac{1}{2}T_{\alpha\beta}$$
 - Einstein's equations

 $T_{\alpha\beta} = \begin{pmatrix} E & 0 \\ 0 & -Ph_{m} \end{pmatrix}$ - homogeneous and isotropic energy-momentum tensor

$$h_{\mu\nu}(t,x) \coloneqq \frac{a^2 \eta_{\mu\nu}}{\left(1 + bx^2\right)^2}$$

$$R_{0\mu} = -\frac{4bx_{\mu}}{\left(1+bx^2\right)^2} = 0 \implies b = \text{const}$$

Definition. Diffeomorphism $i: x \mapsto x'$ is called *isometry* if it preserves metric $g(x) = i^* g(x')$ where i^* is the pullback of i.

n coordinates:
$$g_{\alpha\beta}(x) = \frac{\partial x'^{\gamma}}{\partial x^{\alpha}} \frac{\partial x'^{\delta}}{\partial x^{\beta}} g_{\gamma\delta}(x')$$

Infinitesimal isometry is generated by Killing vector field $K = K^{\alpha} \partial_{\alpha}$ $x^{\alpha} \mapsto x'^{\alpha} = x^{\alpha} + \varepsilon K^{\alpha}, \quad \varepsilon \ll 1$

$$L_{K}g = 0 \quad \Leftrightarrow \quad \nabla_{\alpha}K_{\beta} + \nabla_{\beta}K_{\alpha} = 0 \quad \text{where} \quad K_{\alpha} \coloneqq g_{\alpha\beta}K^{\beta}$$
$$\nabla_{\alpha}K_{\beta} \coloneqq \partial_{\alpha}K_{\beta} - \Gamma_{\alpha\beta}{}^{\gamma}K_{\gamma}$$

<u>Definition</u>. A manifold \mathbb{M} is called *homogeneous at a point* $p \in \mathbb{M}$ if there are infinitesimal isometries which map this point into any other point from some neighborhood \mathbb{U}_p . A manifold \mathbb{M} is called *homogeneous* if it is homogeneous at each point $p \in \mathbb{M}$.

Definition. A manifold \mathbb{M} is called *isotropic at a point* $p \in \mathbb{M}$ if there are such infinitesimal isometries with the Killing 1-forms $K = dx^{\alpha}K_{\alpha} \in \Lambda_1(p)$, preserving this point, that the external derivative $dK \in \Lambda_2(p)$ take any value in the space of two forms $\Lambda_2(p)$. A manifold \mathbb{M} is called *isotropic* if it is isotropic at each point $p \in \mathbb{M}$.

<u>Theorem</u>. Any isotropic (pseudo-)Riemannian manifold (\mathbb{M}, g) is homogeneous.

<u>Theorem</u>. Maximal dimension of the Lie algebra of infinitesimal isometries is equal to n(n+1)/2. If dimension of Lie algebra is maximal, then this manifold is isotropic and homogeneous, and it is a space of constant curvature.

See, for example, S. Weinberg. Gravitation and cosmology.1972

$$R_{\alpha\beta\gamma\delta} = \frac{R}{n(n-1)} \Big(g_{\alpha\gamma} g_{\beta\delta} - g_{\alpha\delta} g_{\beta\gamma} \Big), \quad R = \text{const} - \text{constant curvature space}$$

Manifolds with maximally symmetric submanifolds

 $\mathbb{M} = \mathbb{R} \times \mathbb{S}$ - space-time $t \in \mathbb{R}$ - time $x \in \mathbb{S}$ - constant curvature space slices

n = 4 $K_i = K_i^{\mu}(x)\partial_{\mu}, \quad i = 1,...,6$ - Killing vectors on space slices §

$$g_{\alpha\beta} = \begin{pmatrix} g_{00}(t,x) & g_{0\nu}(t,x) \\ g_{\mu0}(t,x) & h_{\mu\nu}(t,x) \end{pmatrix}$$

 $h_{\mu\nu}(t,x)$ - constant curvature metric on \mathbb{S} depending on t as a parameter Infinitesimal isometry: $t \mapsto t' = t$ $x^{\mu} \mapsto x'^{\mu} = x^{\mu} + \varepsilon K^{\mu}, \quad \varepsilon \ll 1$ (*)

 $\mathbb{T}(\mathbb{S}) \ni \left\{ K^{\mu}(x) \right\} \mapsto \left\{ 0, K^{\mu}(t, x) \right\} \in \mathbb{T}(\mathbb{M}) \quad \text{- extension from } \mathbb{S} \text{ to } \mathbb{M}$

<u>Theorem</u>. Let metric on $\mathbb{M} = \mathbb{R} \times \mathbb{S}$ be invariant with respect to transformations (*), then there is a coordinate system in which the metric is block diagonal

$$ds^{2} = dt^{2} + h_{\mu\nu}dx^{\mu}dx^{\nu}, \qquad (**)$$

where $h_{\mu\nu}(t,x)$ is the metric of constant curvature on \mathbb{S} for all $t \in \mathbb{R}$ depending on time t as a parameter. In this coordinate system the Killing vector fields do not depend on time K = K(x).

Proof.

D. Hilbert. Mathematische Annalen, 15(1924)1 Eisenhart, Continuous Groups of Transformations, (1933) Fix one space slice t = const with coordinates x^{μ} , $\mu = 1, 2, 3$. Take orthogonal (timelike) vector field n on this slice. Construct geodesics going through each point of the space slice along n. Let s be the canonical parameter along geodesics. Choose the coordinate system (s, x^{μ}) on MI.

Then $g_{00} = 1$ by construction.

Geodesic equations $\implies \partial_0 g_{0\mu} = 0$ with initial condition $g_{0\mu}(0, x) = 0$

 \implies unique solution $g_{0\mu}(s,x) = 0 \implies$ Metric has form (**)

For metric (**) the (00) and $(\mu\nu)$ components of the Killing equations are identically satisfied.

The (0μ) component $\implies \partial_s K^{\mu} = 0 \implies K^{\mu} = K^{\mu}(x)$

Theorem. Under the conditions of the previous theorem metric on space slices has the form

$$h_{\mu\nu}(t,x) = a^2(t)\hat{g}_{\mu\nu}(x)$$

where a(t) > 0 is the scale factor and $\hat{g}_{\mu\nu}(x)$ is a constant curvature metric on space slices \mathbb{S} , depending only on x.

Proof.

Killing equations:

$$\nabla_{\mu}K_{\nu} + \nabla_{\nu}K_{\mu} = 0 \iff h_{\mu\rho}\partial_{\nu}K^{\rho} + h_{\nu\rho}\partial_{\mu}K^{\rho} + K^{\rho}\partial_{\rho}h_{\mu\nu} = 0$$

Previous theorem: K = K(x)

$$\dot{h}_{\mu\rho}\partial_{\nu}K^{\rho} + \dot{h}_{\nu\rho}\partial_{\mu}K^{\rho} + K^{\rho}\partial_{\rho}\dot{h}_{\mu\nu} = 0 \quad \Leftrightarrow \quad L_{K}\dot{h}_{\mu\nu} = 0$$

 $h_{\mu\nu} = f(t)h_{\mu\nu}$ - the most general second rank tensor

$$t \mapsto t' \qquad dt' = f(t)dt$$

 $\frac{dh_{\mu\nu}}{dt'} = h_{\mu\nu} \quad \Leftrightarrow \quad h_{\mu\nu}(t', x) = Ce^{t'} \hat{g}_{\mu\nu}(x), \quad C = \text{const}$

Two functions ?

$$g_{\alpha\beta} = \begin{pmatrix} 1 & 0 \\ 0 & \frac{a^2 \eta_{\mu\nu}}{\left(1 + bx^2\right)^2} \end{pmatrix}$$

 $\nabla_{\mu}K_{\nu} + \nabla_{\nu}K_{\mu} = 0$ - Killing equation is satisfied

Killing vector in stereographic coordinates does depend on time

The metric $g_{\alpha\beta}$ is not homogeneous and isotropic in four dimensions.

 $R^{(4)} = R^{(4)}(t, x^2)$ is not homogeneous.

For homogeneous and isotropic matter b = const due to Einstein's equations

Conclusion

<u>Theorem</u>. Let space-time be the topological product $\mathbb{M} = \mathbb{R} \times \mathbb{S}$. Let it be homogeneous and isotropic. Then, up to a coordinate transformation, the most general metric is

$$ds^2 = dt^2 + a^2 \hat{g}_{\mu\nu} dx^{\mu} dx^{\nu},$$

where $\hat{g}_{\mu\nu}(x)$ is the metric of constant curvature on \mathbb{S} and a(t) > 0 is a scale factor. In this coordinate system the Killing vector fields do not depend on time K = K(x)

Definition. A space-time $\mathbb{M} = \mathbb{R} \times \mathbb{S}$ is called homogeneous and isotropic if 1) every section of constant time is a three dimensional space of constant curvature,

2) extrinsic curvature of a submanifold $\mathbb{S} \to \mathbb{M}$ is homogeneous and isotropic.

Extrinsic curvature

$$K_{\mu\nu} = -\frac{1}{2}\frac{d}{dt}h_{\mu\nu}(t,x)$$