

UV (in)sensitivity of Higgs Inflation

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with Marieke
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(in preparation).



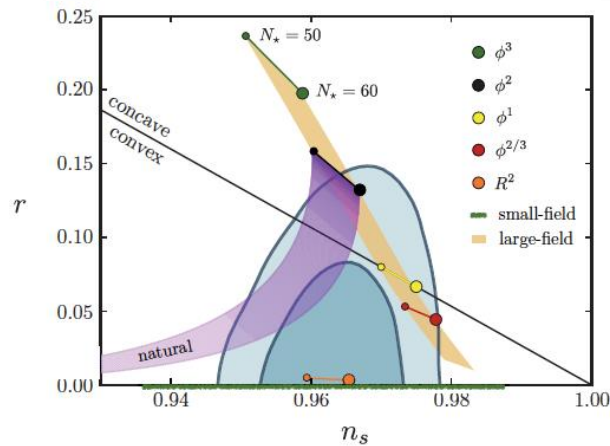
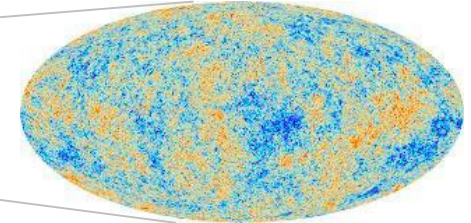
TEXAS Symposium 2015,
Geneva,
14/12/2015

Motivation

- Inflation $\ddot{a}(t) > 0$

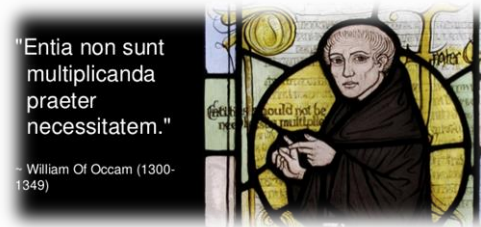
A.Guth, A.Linde, P. Steinhard'80

- Who is the Inflaton?



- Request for simplicity, \implies
SM
LHC , PLANCK
- Consistent simplicity

Only one scalar field in the



Overview

- Higgs Inflation
 - Tree level
 - Quantum aspects (Unitarity, renormalizability ..)
- UV sensitivity
 - UV completion
 - RG flow and renormalization scale
 - Results for the CMB parameters

Higgs Inflation

'08

$$S = \mathcal{S}_{SM} + \frac{M_{pl}^2}{2} \int \sqrt{-g} \left(1 + \frac{2}{M_{pl}^2} \xi \mathcal{H}^t \mathcal{H} \right) R$$

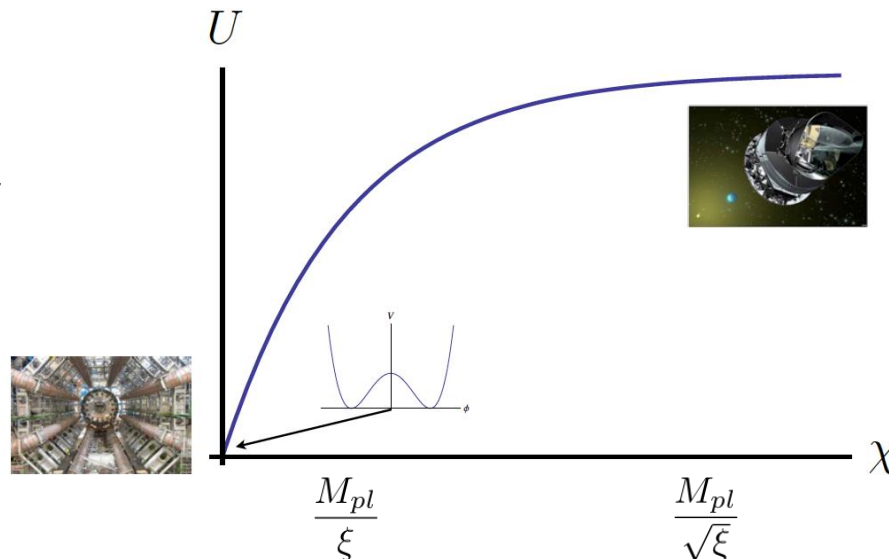
• Einstein frame

$$g_{E\mu\nu} = \Omega^2 g_{\mu\nu}$$

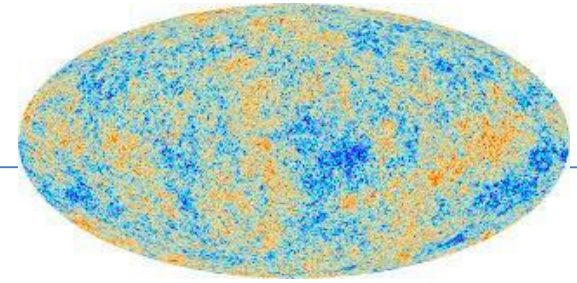
$$\gamma(\phi) (\partial_\mu \phi)^2 = (\partial_\mu \chi)^2$$

$$S = \frac{1}{2} \int \sqrt{-g_E} R(g_E) - \int \sqrt{-g_E} \left[\frac{1}{2} \partial_\mu \chi \partial^\mu \chi + U(\chi) \right]$$

$$U(\chi) = \frac{\lambda(\phi^2(\chi) - v_{ew}^2)^2}{4\Omega^4(\chi)}$$



Density perturbations

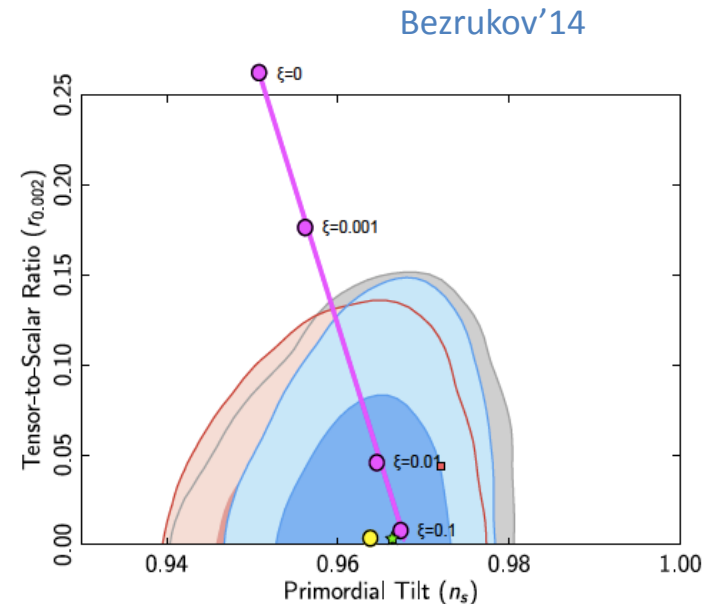


$$N_{\star} = \int_{\chi_E}^{\chi_{\star}} \frac{U}{dU/d\chi} \frac{d\chi}{M_{pl}} \implies \phi_{\star} \simeq 9.13 M_{pl} / \sqrt{\xi}$$

$$\Delta_{\mathcal{R}} \left(\propto \frac{\lambda}{\xi^2} \right) \simeq 2.2 \cdot 10^{-9} \implies \xi \simeq 47000 \sqrt{\lambda}$$

$$n_s \simeq 1 - \frac{2}{N_{\star}} - \frac{3}{N_{\star}^2} \simeq 0.967$$

$$r = 16\epsilon_{\star} \simeq \left(1 + \frac{1}{6\xi_{\star}}\right) \frac{12}{N_{\star}^2} \simeq 0.0031$$



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Unitarity bound

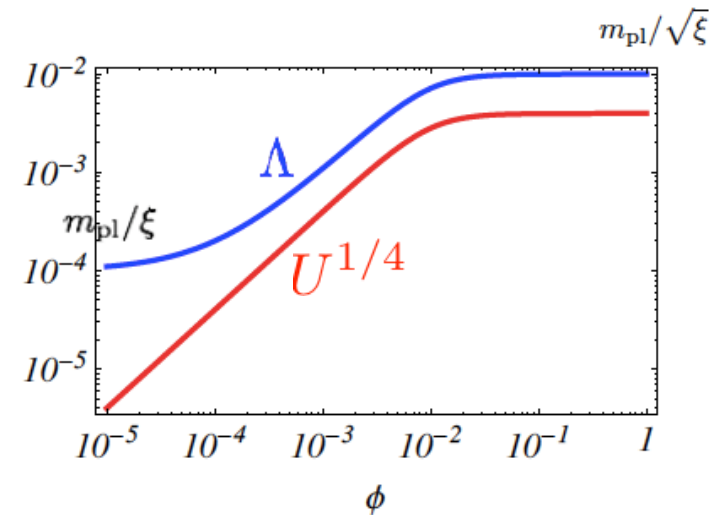
$$S \supset \sqrt{-g} \frac{\xi}{M_{pl}^2} \phi^2 R \xrightarrow{g_{\mu\nu} = \eta_{\mu\nu} + M_{pl}^{-1} h_{\mu\nu}} \frac{1}{M_{pl}/\xi} \phi \square h_{\mu\nu}$$

Burgess '09, Barbon '09, Hertzberg '10

- Cutoff field dependent
- Considering Gauge interaction (Goldstone)

$$\mathcal{M}(\theta\theta \rightarrow \theta\theta) > 1$$

$$\Lambda_{gauge}(\phi_0) \sim \left(\frac{M_{pl}}{\xi}, \phi_0, \frac{M_{pl}}{\sqrt{\xi}} \right)$$

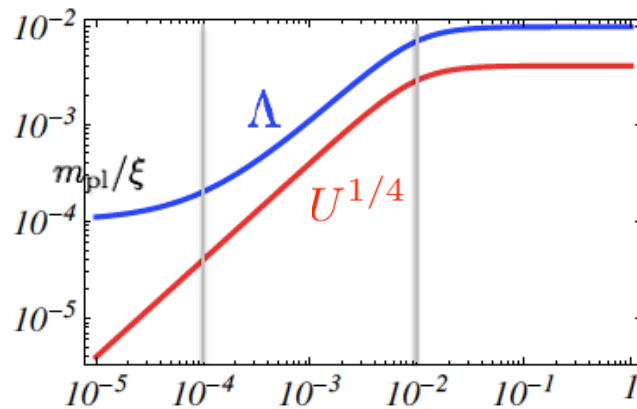


Bezrukov et al. '11, Burgess '14, Prokopec & Weenink '14

Renormalizability in EFT sense

D.George, S.Mooij, M.Postma '14
'15

- Small regime $\phi \ll \frac{M_{pl}}{\xi}$ $\delta_s = \xi\phi$
- Mid regime $\frac{M_{pl}}{\xi} \ll \phi \ll \frac{M_{pl}}{\sqrt{\xi}}$ $\xi \rightarrow \delta_m^{-2}\xi, \phi \rightarrow \delta_m^{\frac{3}{2}}\phi$
- Large regime $\phi \gg \frac{M_{pl}}{\sqrt{\xi}}$ $\delta = 1/\xi\phi$



Demand: at every order a finite number of counter terms

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UV completion



$$\frac{\mathcal{L}_{HI}}{\sqrt{-g}} + \sum_i \frac{c_i}{\Lambda^n} \mathcal{O}_i^{n+4}$$

Bezrukov, Rubio and Shap.'14,
Burgess, Patil, Trott '14

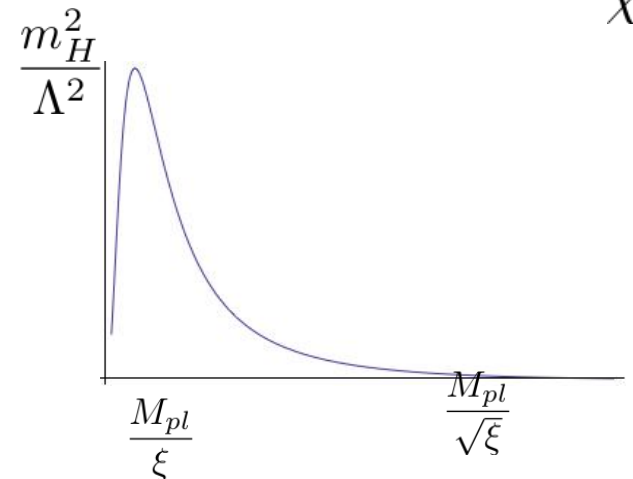
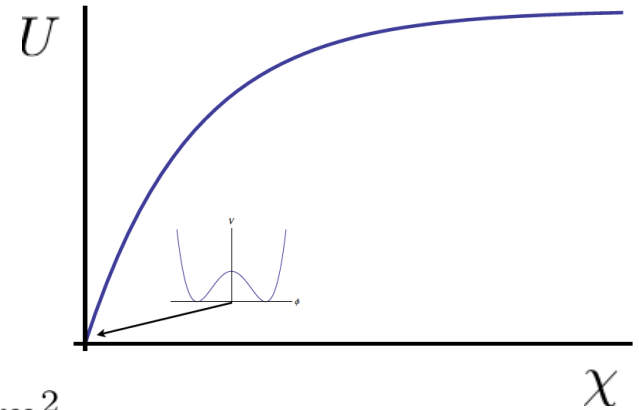
- Which shape for the suppression scale?

$$\Lambda = \Lambda_{gauge}(\phi)$$

- Which form for the higher d operators?

$$\mathcal{L}_{new} \supset \frac{(\mathcal{H}^\dagger \mathcal{H})^3}{\Lambda^2} + \dots \quad \mathcal{L}_{new} \supset \frac{m_h^2 \mathcal{O}^{(4)}}{\Lambda^2}$$

- Preserving the quasi-shift symmetry
- Effect only where really needed!



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RG flow

$$\beta_i = \mu \frac{\partial \lambda_i}{\partial \ln \mu} = \beta_i^{HI} + \delta \beta_i$$

Bezrukov, Grubinov, Shaposhnikov
 Barvinsky, Kamenshchik, Kiefer,
 Starobinsky,
 Simone, Hertzberg, Wilzcek
 George, Mooij, Postma
 Burgess, Patil, Trott

- Corrected Potential

$$U = \frac{\lambda \phi^4}{4 \left(1 + \frac{\xi \phi^2}{M_{pl}^2}\right)} + U^{(1)} + 2\text{-loop} + \dots \supset (-1)^{f_i} c_i \frac{m^4(\phi)}{64\pi^2} \ln \left(\frac{m^2(\phi)}{\mu^2} \right)$$

$$m(\phi) = \frac{m_{SM}(\phi)}{\Omega(\phi)}$$

- RG improvement

$$U(\phi, g_i, \mu) = U(\phi, g_i(t), \mu(t))$$

$$\mu(t) = \mu e^t, \quad \frac{dg_i(t)}{dt} = \beta_i(g_j(t))$$

$$\mu(t) \sim \frac{\phi}{\Omega} \equiv \frac{\phi}{\sqrt{1 + \xi(t)\phi^2}}$$

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Effect on the CMB observables

$$U_{eff} = \frac{\lambda(t(\phi))\phi^4}{4\left(1 + \frac{\xi(t(\phi))\phi^2}{M_{pl}^2}\right)}$$

$$t = \ln\left(\frac{\mu}{m_t}\right) = \ln\left(\frac{\phi}{m_t\sqrt{1 + \xi(t)\phi^2}}\right)$$

Inflationary regime $\phi^2 \gg M_{pl}^2/\xi$

$$\longrightarrow \boxed{\delta = \frac{1}{\xi\phi^2} \ll 1}$$

$$\eta \approx -\frac{4}{3} \frac{1}{\left(1 + \frac{1}{6\xi}\right)} \mathcal{Z} \delta + O(\delta^2)$$

$$\epsilon \approx \frac{4}{3} \frac{1}{\left(1 + \frac{1}{6\xi}\right)} \mathcal{Z}^2 \delta^2 + O(\delta^3)$$

$$\downarrow$$

$$\frac{1 + \frac{\beta_\lambda}{4\lambda}}{1 + \frac{\beta_\xi}{2\xi}}$$

$$N_\star \approx \frac{3}{4} \mathcal{Z}_\star^{-1} \left(1 + \frac{1}{6\xi_\star}\right) \frac{1}{\delta_\star}$$

$$n_s \simeq 1 - \frac{2}{N_\star} + O(N_\star^{-2})$$

$$r \simeq \frac{12}{N_\star^2} \left(1 + \frac{1}{6\xi_\star}\right) + O(N_\star^{-3})$$

Same as tree level results

!

Running insensitivity

- $Z_* \approx 0$

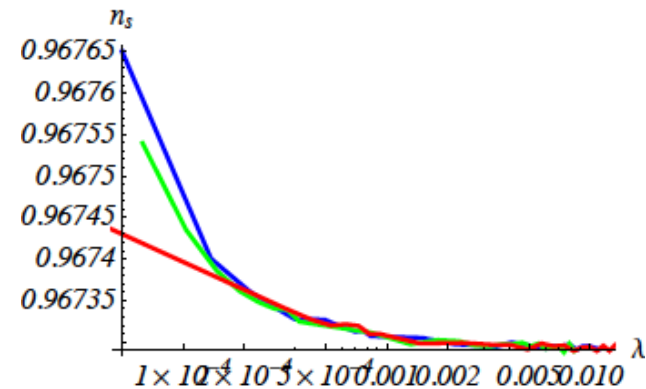
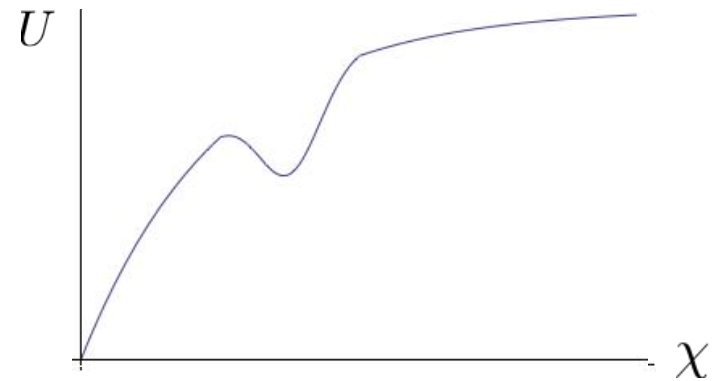
$$U_{\chi}/U \propto Z\delta$$

$$U_{\chi\chi}/U \propto -Z\delta$$

- Numerical results, example

For Standard Model
same as [A.Kyle'14](#)

- Second order, h.o. corrections $O(10^{-4})$ and suppressed



Conclusions

- Higgs inflation at tree level gives great prediction (for α_s, r)
- The theory is not renormalizable, we need at least a particular UV completion
- Threshold corrections to the slow roll parameter cancel in the expression for α_s and r independently of the form of this UV completion

RG flow

$$\Gamma[\phi_{cl}] = S[\phi_{cl}] + \Gamma^{1\text{-loop}} + \dots$$

$$V_{CW} = V_t(\phi_{cl}) + \frac{1}{64\pi^2} \sum_i (-1)^{f_i} s_i m_i^4(\phi_{cl}) \ln \left(\frac{m_i^2(\phi_{cl})}{\mu^2} - c_i \right)$$

Coleman-Weinberg '73

$$g_E = \Omega^2 g_J$$

Quantum corrections two routes:

Jordan

$$V_J(\phi) + V_{JCW}^{(1)} \xrightarrow{E} \frac{V_J(\phi)}{\Omega^4} + \frac{V_J^{(1)}}{\Omega^4}$$

Einstein

$$V_E(\phi) + V_{CW}^{(1)} = \frac{V_J}{\Omega^4} + V_E^{(1)}$$

- ⇒
- Different results in the GB sector
 - Back reaction negligible only in the Einstein frame
 - Approx. FLRW
- $\sim O(\epsilon), \epsilon \ll \eta$

M.Postma '14

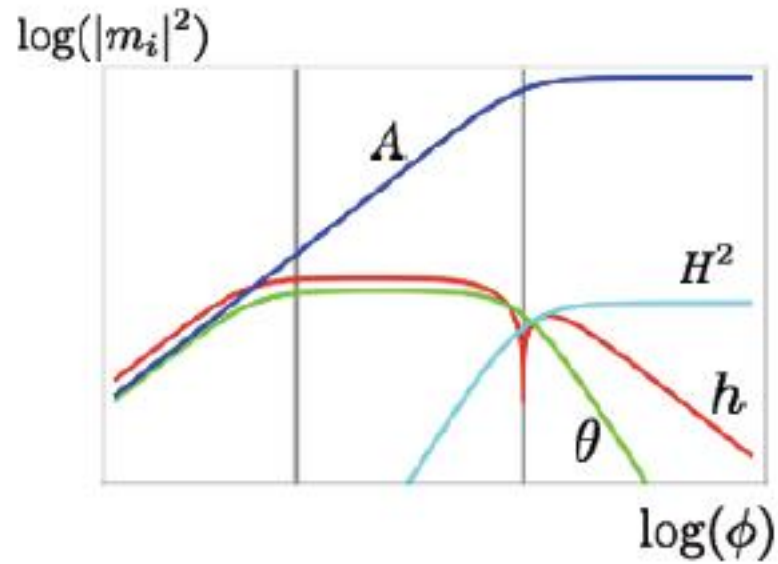
S.Mooij, M.Postma '11

D. George, S.Mooij,

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Jacopo Fumagalli

Masses



Higgs instability

