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# Planck 2015 results on Dark Energy and Modified Gravity

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### **Outline**

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2 Beyond the Λ[CDM model](#page-8-0)

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#### 2 Beyond the Λ[CDM model](#page-8-0) **[Background Constraints](#page-16-0) [Perturbations constraints](#page-18-0)**



# Minimal ΛCDM model

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Since the '90s and the discovery of the late accelerated expansion of the Universe, the standard cosmological model has been the most efficient in describing our observations.

This model relies on a Cold Dark Matter component to describe the evolution of cosmic structures and on a cosmological constant  $\Lambda$  to account for the accelerated expansion phase.

The ΛCDM gives predictions for cosmological observables in terms of 6 standard parameters

 $\{\Omega_b, \Omega_{cdm}, n_s, A_s, H_0, \tau\}$ 

#### The Cosmic Microwave Background

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The minimal ΛCDM model explains quite efficiently cosmological observables, including the most primordial available: the Cosmic Microwave Background



This relic radiation from the Big Bang carries information on the density distribution and the primordial phases of the Universe

#### CMB and the late time Universe

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CMB photons travel through the more recent Universe and are therefore affected also by more recent physical mechanisms, e.g. CMB lensing and ISW effect.



CMB photons contain also informations on the more recent phases of the Universe and can be used to test cosmological models on a wide time range.

### The Planck Satellite

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Early in 2015, the ESA Planck satellite released the most up to date CMB data.

Accurate foreground control and extreme sensitivity.



### Planck 2015

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Temperature and polarization spectra extracted from CMB maps are in very good agreement with the predictions of the minimal 6-parameters ΛCDM model



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# Testing  $\Lambda$

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ΛCDM is a good fit to CMB data and within this framework accurate constraints on cosmological parameters can be obtained.

CMB allows to test the assumption of a Cosmological Constant  $(\Lambda)$  as the responsible for the late time acceleration.

Abandoning this paradigm impacts the evolution of the Universe through changes of

- **background expansion**
- $\blacksquare$  evolution of cosmological perturbations

CMB power spectra are affected by these modifications

### What to test for?

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If we want to test  $\Lambda$  we have to describe departures from it: parametrized deviations

find peculiar properties of your model and parametrise them (e.g.  $w(z) = -1$ )

specific alternative models assume a specific model and test whether or not it better fits the data

### What to test for?

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specific alternative models assume a specific model and test whether or not it better fits the data

Both approaches can be applied to the two broad classes of alternatives to Λ

$$
G_{\mu\nu}=8\pi G T_{\mu\nu}
$$

### What to test for?

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 $\Gamma_{\mu\nu} = 8\pi G T_{\mu\nu}$ 

Modified Gravity modifications to GR lagrangian

Dark Energy additional energy components

### Planck models and parametrizations

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Planck collaboration tried to investigate deviations from ΛCDM moving between these approaches

### Planck models and parametrizations

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#### Planck collaboration tried to investigate deviations from ΛCDM moving between these approaches



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### Data combinations

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These parametrizations are investigated combining Planck with additional observations

- **Planck baseline**: Planck  $TT + low-\ell$  polarization
- **background probes:** 
	- BAO: SDSS (Ross et al 2014), BOSS (Anderson et al. 2014), 6dFGS (Beutler et al. 2011)
	- SN: JLA (Betoule et al. 2013)
	- $H_0$  conservative prior (Efstathiou 2014)
- perturbation probes:
	- RSD (BOSS DR11, Samushia et al. 2014)
	- WL (CFHTLens, Kilbinger et al. 2013, Heymans et al. 2013), with an ultra conservative cut of non linear scales

### <span id="page-16-0"></span>Equation of State

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Assuming the background expansion deviates from the standard ΛCDM, the EoS parameter can depart from the constant -1

### Equation of State

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Assuming the background expansion deviates from the standard ΛCDM, the EoS parameter can depart from the constant  $-1 \Rightarrow w(z) = w_0 + w_a(\frac{z}{1+z})$  $\frac{z}{1+z}$ 



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<span id="page-18-0"></span>Focusing on perturbations evolution, Planck analysis exploited two different parametrizations of departure from standard ΛCDM driven behaviour:



Top-down approach Bottom-up approach

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constraints

■ Top-down approach→parametrize your theory

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■ Top-down approach $\rightarrow$ parametrize your theory Action for scalar tensor theories with only one extra dynamical field, preserving background isotropy and homogeneity

$$
S = \int d^4x \sqrt{-g} \{ \frac{m_0^2}{2} [1 + \Omega(\tau)] R + \Lambda(\tau) +
$$
  

$$
f(c, \hat{M}^2, \bar{M}_1^3, \bar{M}_2^4, \bar{M}_3^2, m_2^2) \}
$$

Planck analysed the case where  $\Omega(\tau)$  is the only free function (non minimally coupled K-essence)

EFTCAMB (Hu, Raveri, Silvestri, Frusciante 2014)

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■ Top-down approach→parametrize your theory

$$
\Omega(a) = e^{\frac{\alpha_{M0}}{\beta}a^{\beta}} - 1 \quad \Omega(a) = \alpha_{M0}a
$$

Bellini, Sawicki 2014

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\Omega(a) = e^{\frac{\alpha_{M0}}{\beta}a^{\beta}} - 1 \quad \Omega(a) = \alpha_{M0}a
$$

Bellini, Sawicki 2014

Bottom-up approach→parametrize your "observables"

$$
k^{2}\Psi = 4\pi Ga^{2}\mu(a,k)\rho\Delta
$$

$$
\frac{\Phi}{\Psi} = \eta(a,k)
$$

MGCAMB (Hojjati, Pogosian, Zhao 2011)

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■ Top-down approach $\rightarrow$ parametrize your theory

$$
\Omega(a) = e^{\frac{\alpha_{M0}}{\beta}a^{\beta}} - 1 \quad \Omega(a) = \alpha_{M0}a
$$

Bellini, Sawicki 2014

Bottom-up approach→parametrize your "observables"  $\mathcal{L}_{\mathcal{A}}$ 

$$
\mu(z,k) = 1 + E_{11} \Omega_{DE}(z) \frac{1 + c_1 (\lambda H/k)^2}{1 + (\lambda H/k)^2}
$$

$$
\eta(z,k) = 1 + E_{22} \Omega_{DE}(z) \frac{1 + c_2 (\lambda H/k)^2}{1 + (\lambda H/k)^2}
$$

# Effects on CMB

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While testing for deviations from the standard perturbations evolution, background expansion of ΛCDM is assumed.

CMB is affected by the late modified evolution of cosmological perturbations: ISW effect and CMB lensing



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# (some) Planck results

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#### Top-down approach

- ΛCDM limit in agreement with all data combinations
- possible hint of tensions between the datasets





#### Bottom-up approach

- **Mild tension with**  $\Lambda$ CDM  $(\approx 2\sigma)$
- Tension enhanced including perturbation probes (RSD and WL)

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#### Lensing reconstruction

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From Planck CMB maps it's possible to reconstruct the CMB lensing potential power spectrum through quadratic estimators



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### The amplitude of lensing potential

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Planck collaboration measured how much the amplitude of the CMB lensing potential power spectrum deviates from the  $\Lambda$ CDM prediction  $(C_{\ell} = A_L C_{\ell}^{\Lambda CDM})$ , both through temperature and polarization spectra and through the extraction of the lensing potential from CMB maps.



#### The importance of being lensed

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Results obtained from lensing extraction seem "more in agreement" with ΛCDM.

This feature affects significantly the slight tension with the standard model found with the MGCAMB approach



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### Conclusions

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- Overall agreement betwen Planck and the ΛCDM model
- **E** "Agnostic" analysis, based on several parametrizations and approaches
- Improvement of previous bounds on both background and  $\mathbf{r}$ perturbation parametrizations
- Planck+external datasets results show some marginal tensions with ΛCDM

### Planck Collaboration

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# EoS PCA

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The PCA approach allows to avoid any assumptions on the behaviour of the EoS. The value of  $w(z)$  is free to vary in different redshift bins. Model independent, but larger errors.



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# Early Dark Energy

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This kind of models account for a time evolving DE with non vanishing density at early times

$$
\Omega_{de}(a) = \frac{\Omega_{de}^{0} - \Omega_{e}(1 - a^{-3w_0})}{\Omega_{de}^{0} + \Omega_{m}^{0}a^{3w_0}} + \Omega_{e}(1 - a^{-3w_0})
$$



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#### <span id="page-34-0"></span>Effective Field Theory for DE/MG

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In general there are 9 functions of time that include majority of Modified Gravity models

$$
S = \int d^4x \sqrt{-g} \{ \frac{m_0^2}{2} [1 + \Omega(\tau)] R + \Lambda(\tau) - a^2 c(\tau) \delta g^{00} + \frac{M_2^4(\tau)}{2} (a^2 \delta g^{00})^2 - \bar{M}_1^3(\tau) 2a^2 \delta g^{00} \delta K_\mu^\mu - \frac{\bar{M}_2^2(\tau)}{2} (\delta K_\mu^\mu)^2 - \frac{\bar{M}_3^2(\tau)}{2} \delta K_\nu^\mu \delta K_\mu^\nu + \frac{a^2 \hat{M}^2(\tau)}{2} \delta g^{00} \delta R^{(3)} + m_2^2(\tau) (g^{\mu\nu} + n^\mu n^\nu) \partial_\mu (a^2 g^{00}) \partial_\nu (a^2 g^{00}) \} + S_m[\chi_i, g_{\mu\nu}]
$$

Describes scalar tensor theories with one extra dynamical d.o.f.

Gubitosi et al. 2012