

Interacting pulsar winds in high-mass binaries

Guillaume Dubus (Grenoble)

with

Astrid Lamberts (Caltech)

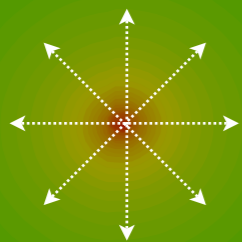
Sébastien Fromang (Saclay)

A&A, 581, A27 (2015)

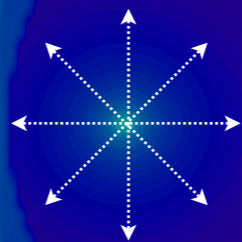
Emission powered by spindown of a pulsar

Ram pressure equilibrium between pulsar wind and stellar wind gives rise to small-scale pulsar wind nebula with a bow shock structure, set by $\eta = \frac{\dot{E}/c}{\dot{M}v_w}$

massive star wind



termination shock



pulsar wind powered by spindown

Termination shock is close to pulsar, about $10^4 R_{LC}$

High-mass gamma-ray binaries

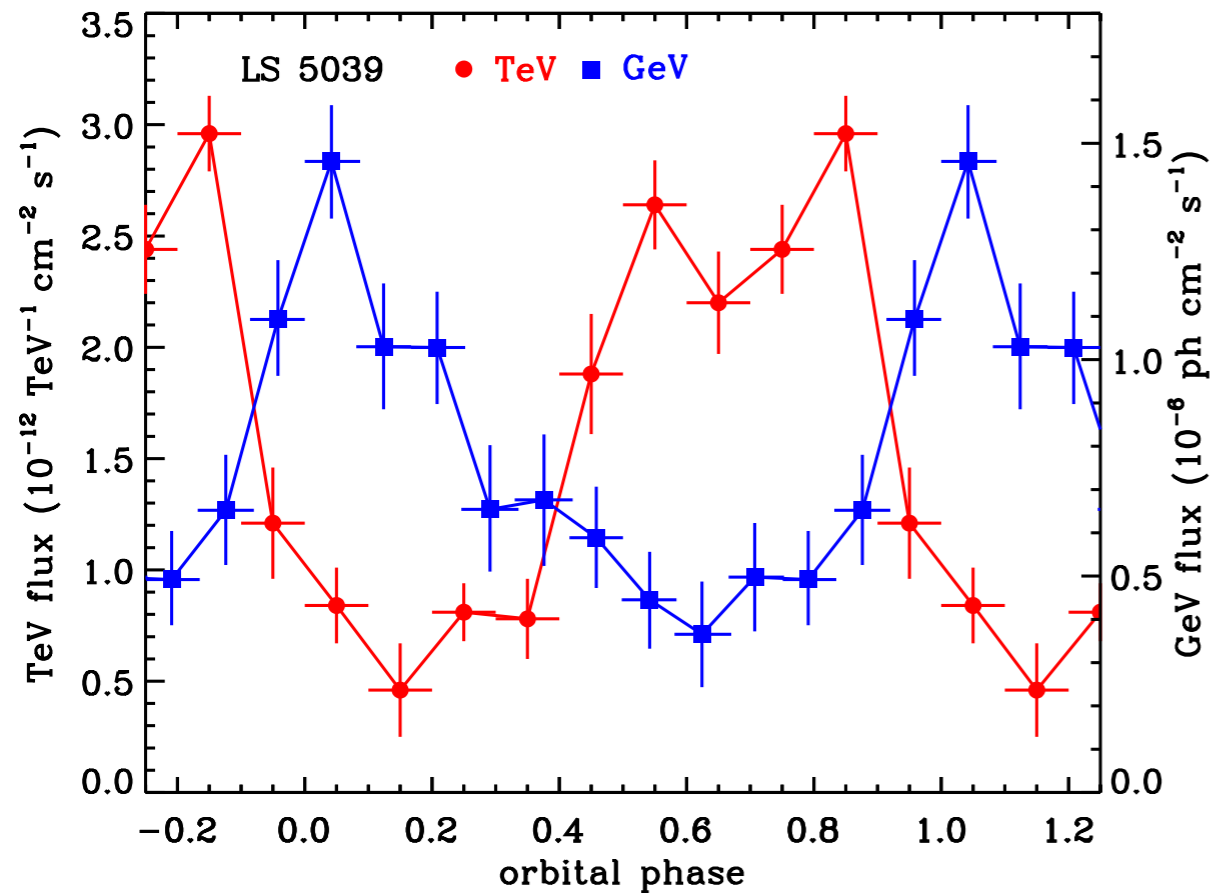
Sources of non-thermal emission up to GeV and TeV gamma-rays

system	star	radio pulsar	P_{orb} (days)	HE or VHE gamma-rays ?
PSR B1259-63	Be	yes	1237	yes
LS 5039	O	?	3.9	yes
LS I +61 303	Be	(?)	26.5	yes
HESS J0632+057	Be	?	320	yes
1FGL J1018.6-5856	O	?	16.6	yes

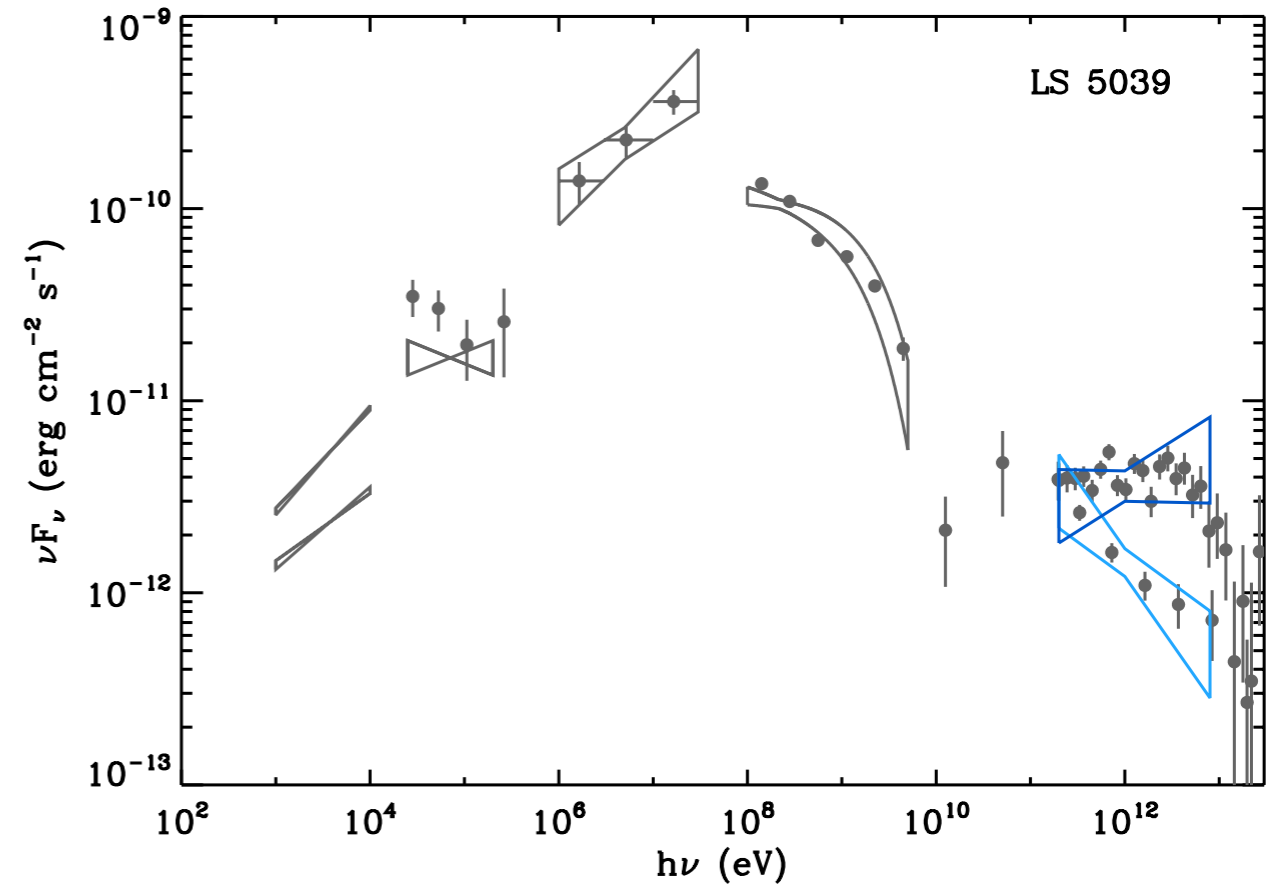
(*) there are also **low-mass** gamma-ray binaries (tomorrow)

Well-established phenomenology in LS 5039

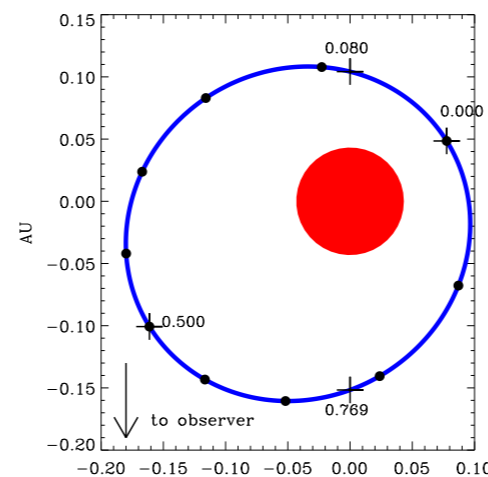
orbital modulation



spectrum



Good constrains on 3.9 day orbit



Possible gamma-ray emission sites

shocked
stellar wind

bow-shocked
pulsar wind

back-shocked
pulsar wind

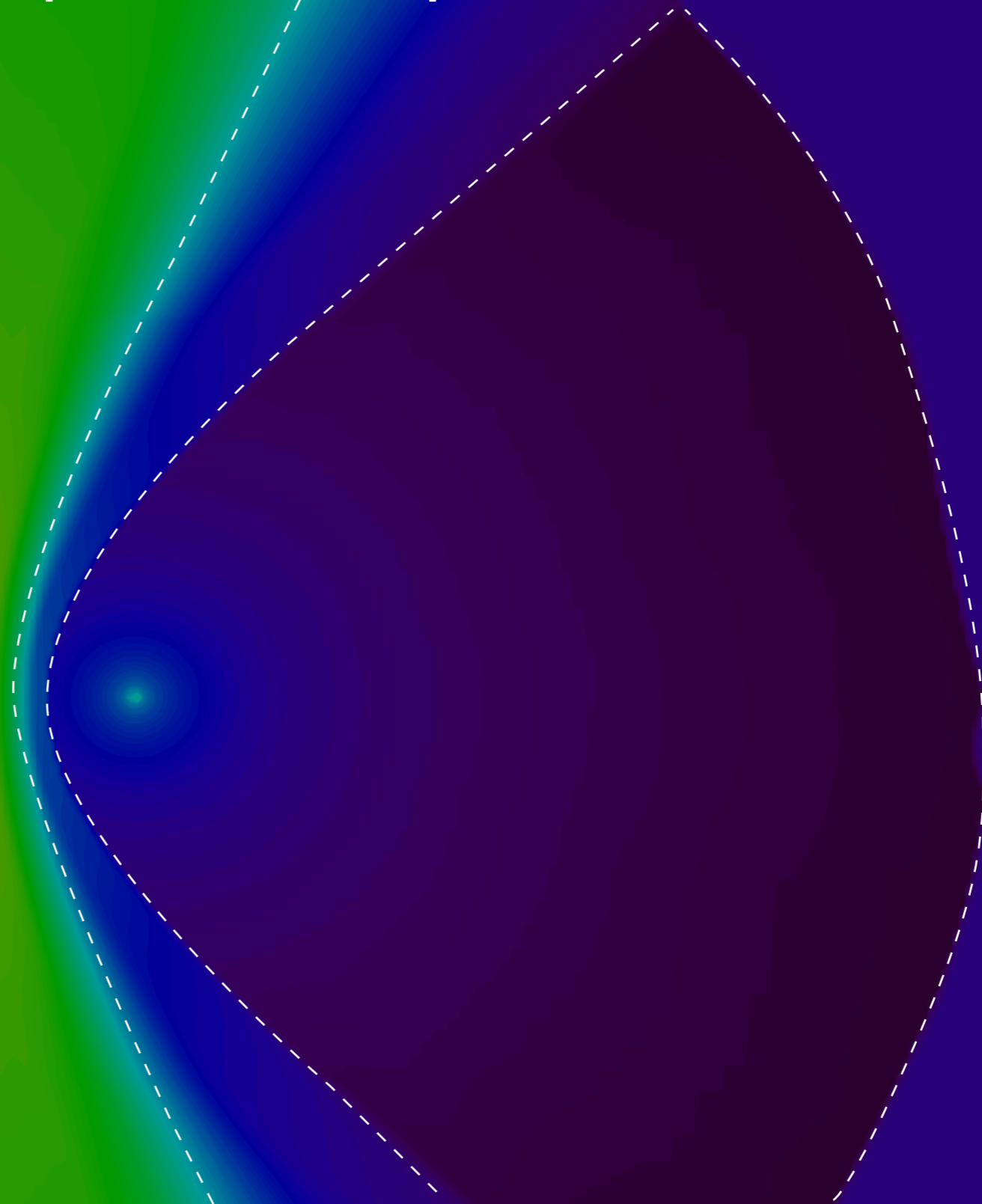
pulsar wind

pulsar
magnetosphere

Use relativistic hydro simulation as input

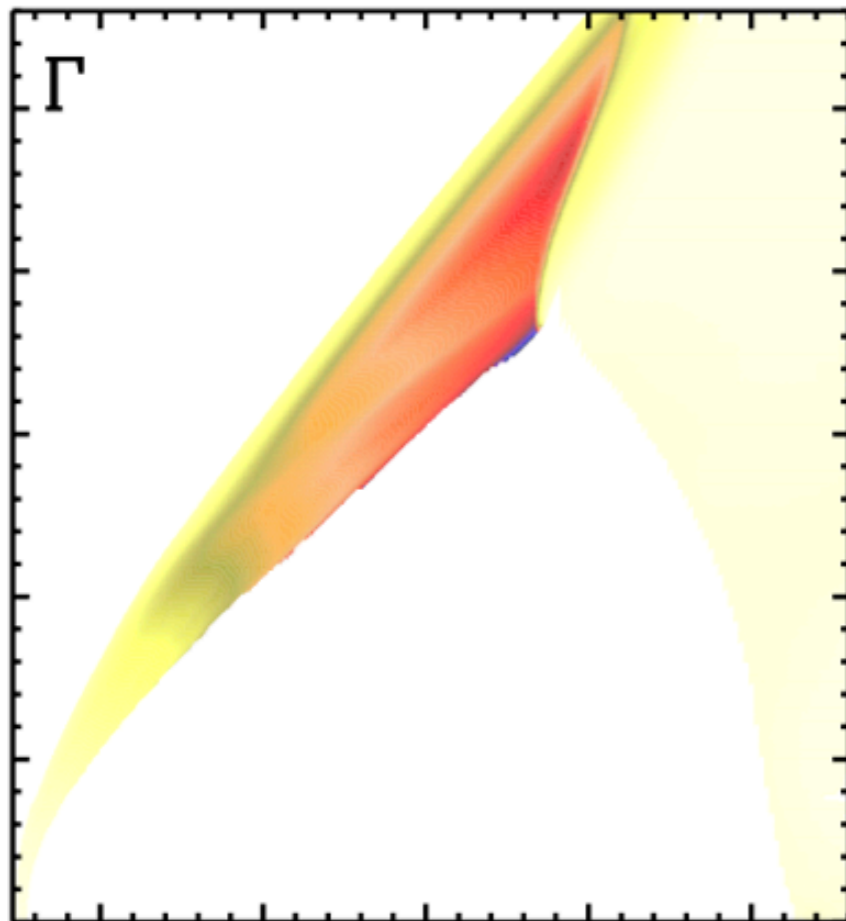
provides consistent description in shocked pulsar wind emission model of

- shock extension
- doppler boosting
- adiabatic cooling

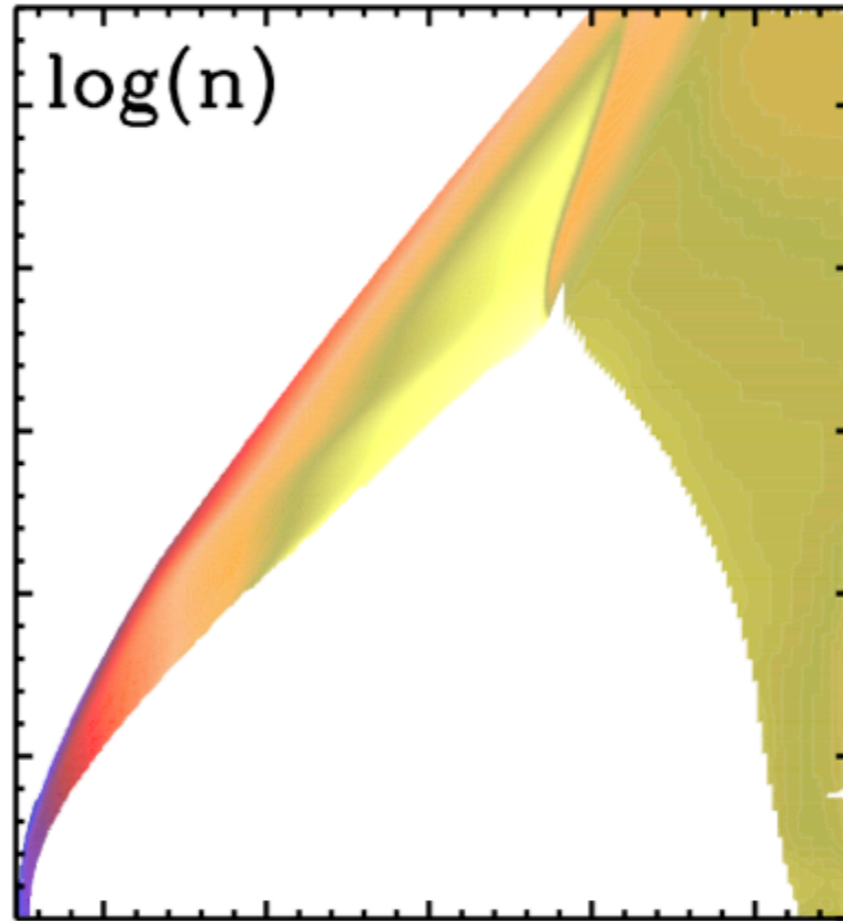


Conditions in the shocked pulsar wind

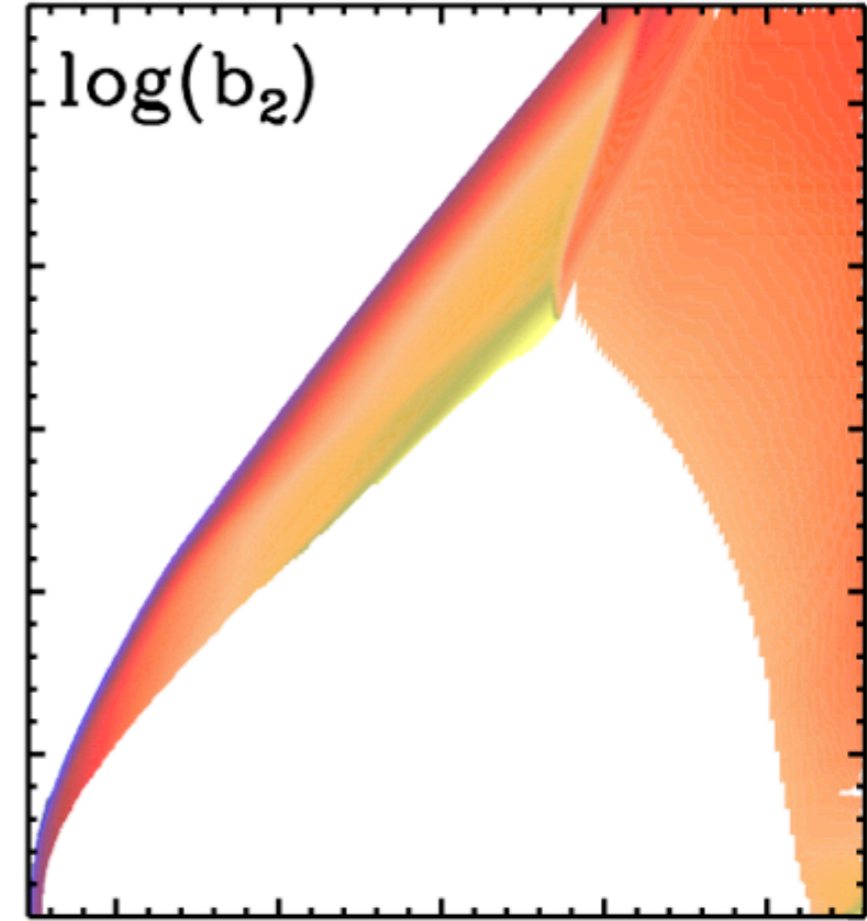
Lorentz factor



density



magnetic field

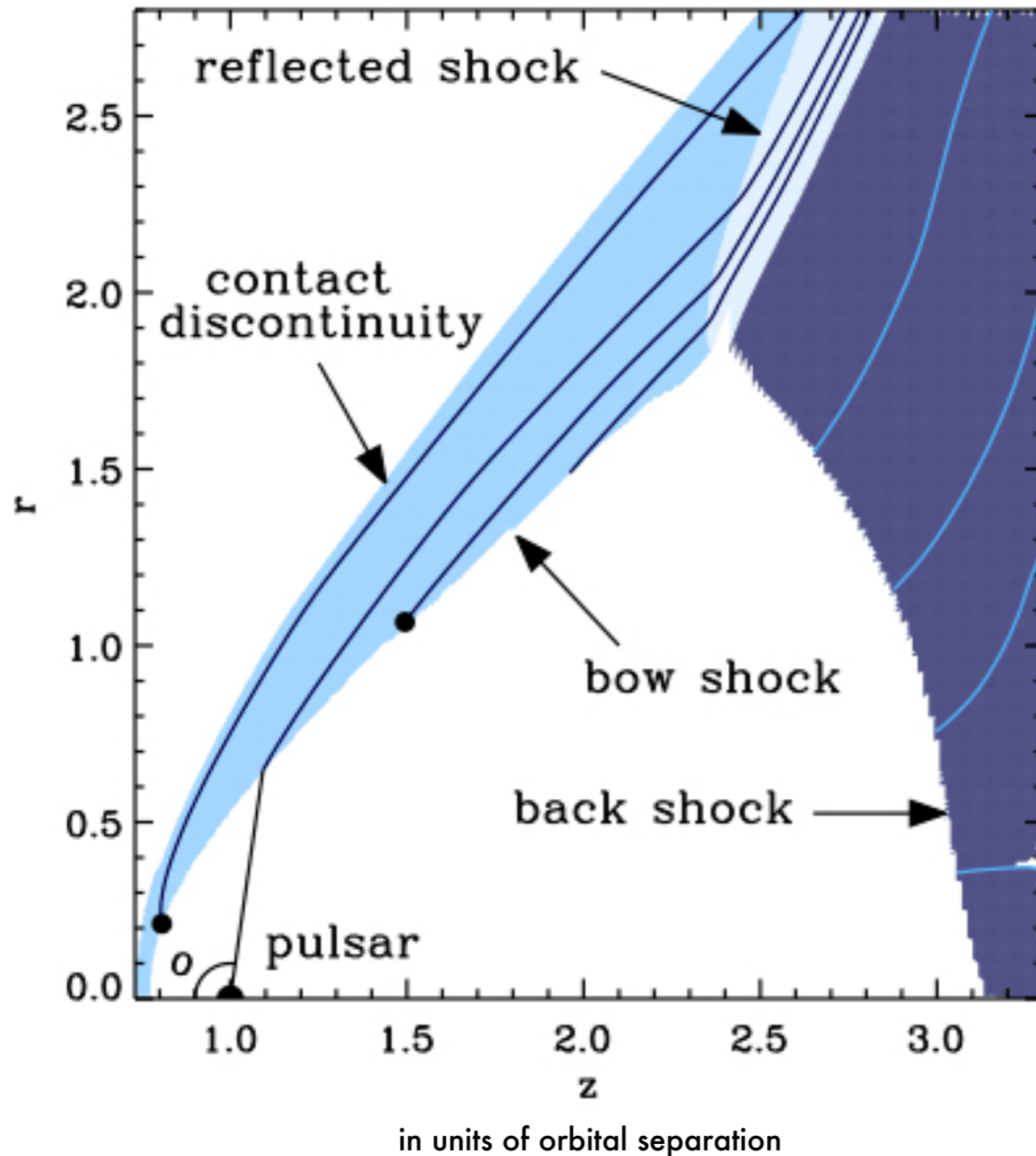


$B \sim 1/d_p$ at shock (assumed perpendicular)

Field is passive so induction equation $\Rightarrow B \sim \Gamma n r$

see Lamberts et al. 2013 and Dubus et al. 2015 for code & setup

Particle evolution in a post-processing step



1- inject particles at shock

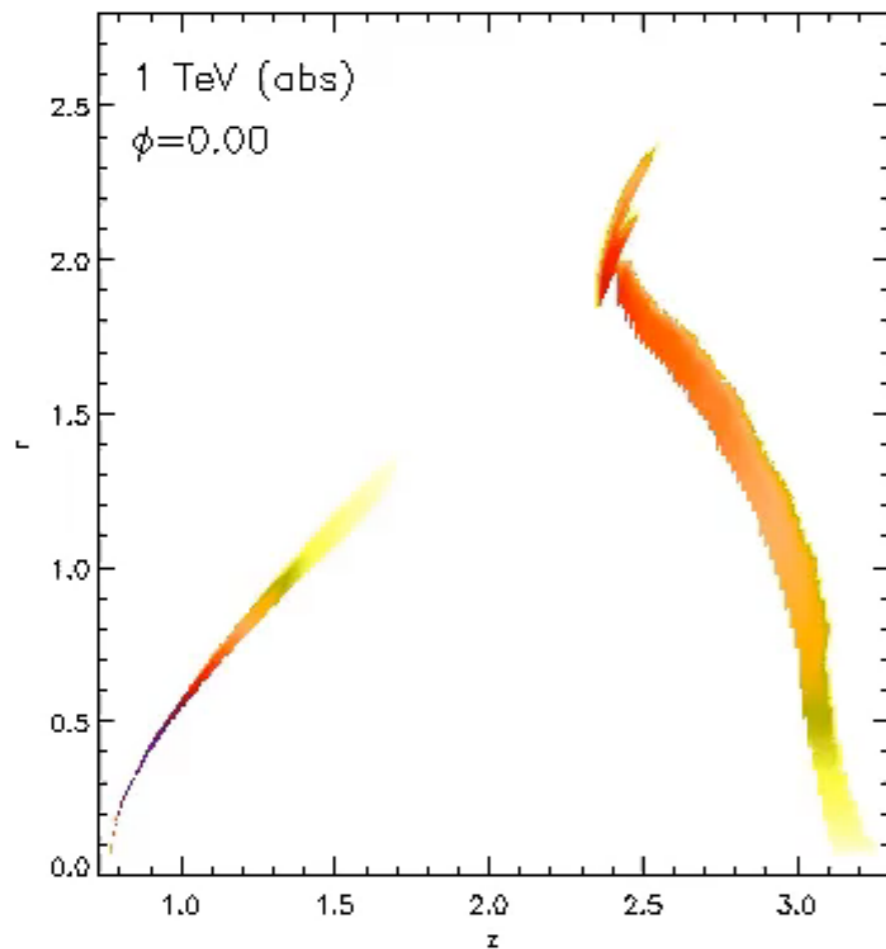
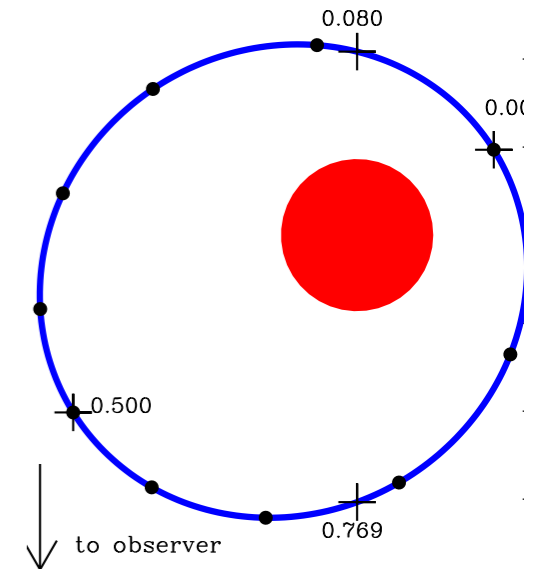
four global parameters:

- power-law slope
- acceleration timescale (E_{\max})
- pulsar power (E_{\min})
- magnetisation (B)

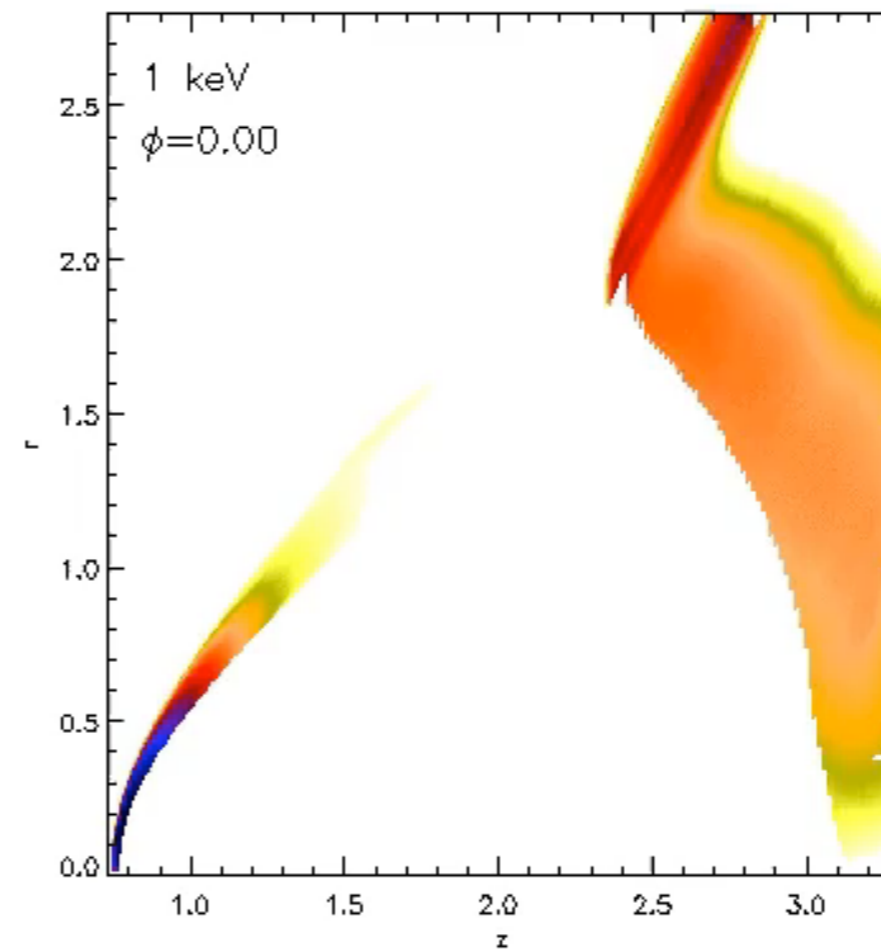
2- evolve along streamlines with radiative & adiabatic cooling.

3D calculation for particle emission

- **Synchrotron and inverse Compton emission**
eccentric orbit, relativistic effects, anisotropic scattering & pair production on stellar photons
- **phase-dependent spectra & orbital modulation**

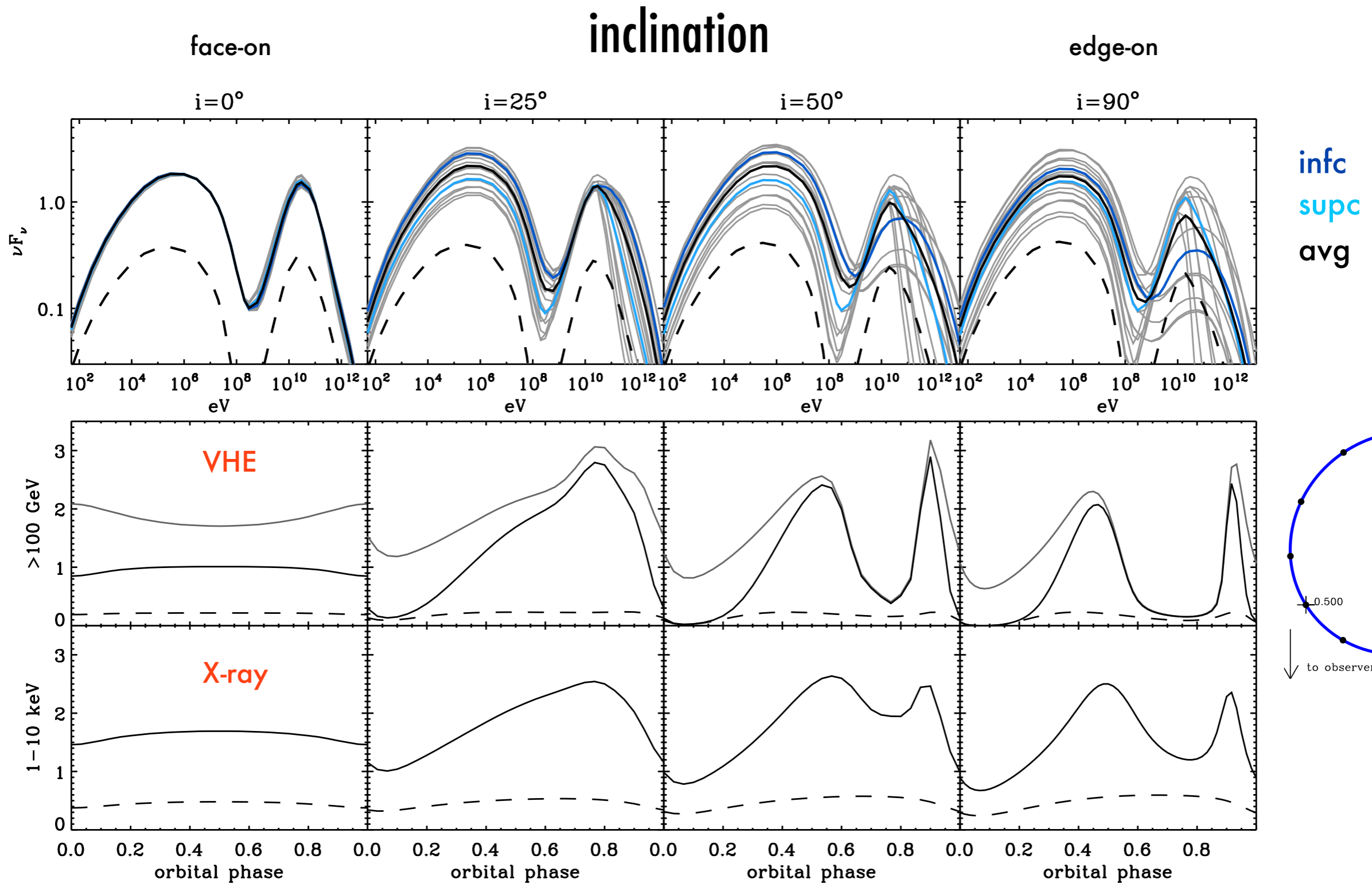


in units of orbital separation



in units of orbital separation

Doppler boosting shapes the modulation



Influence of particle injection parameters

infc
supc

magnetic field

σ

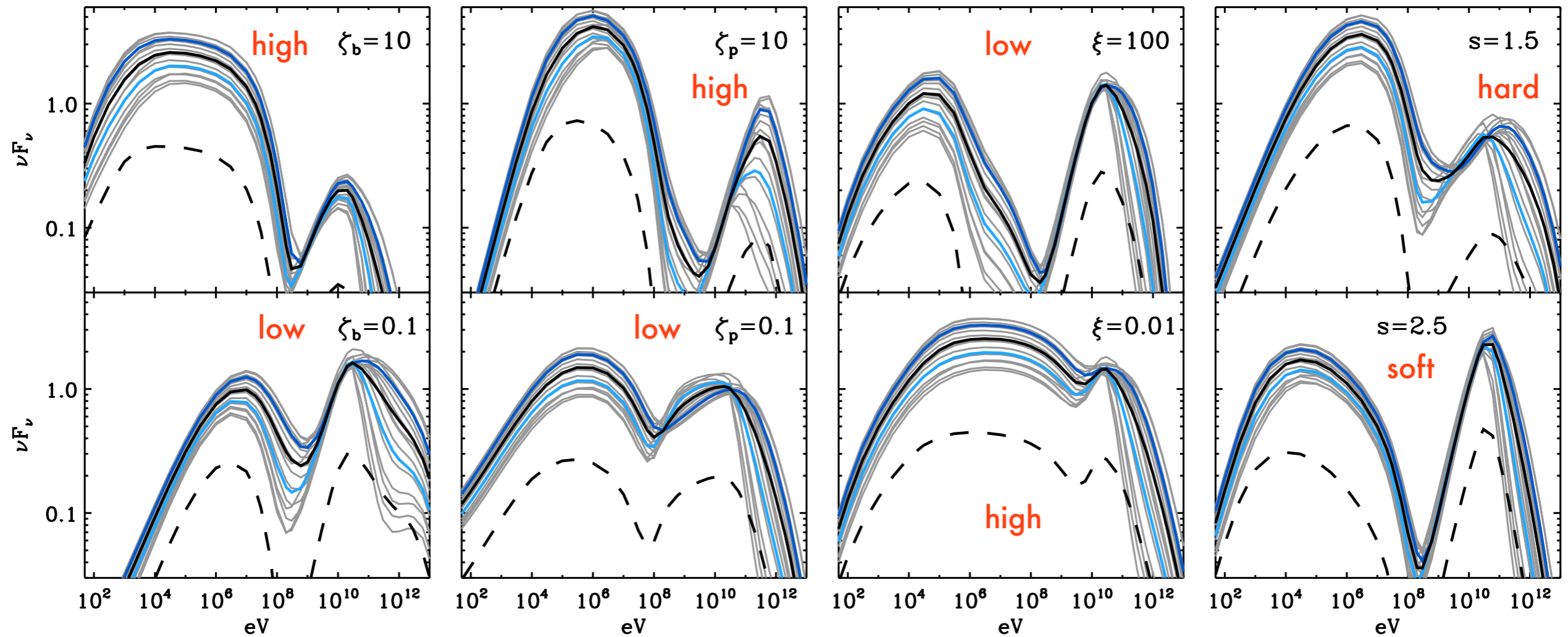
E_{\min}

pulsar power

E_{\max}

acceleration efficiency

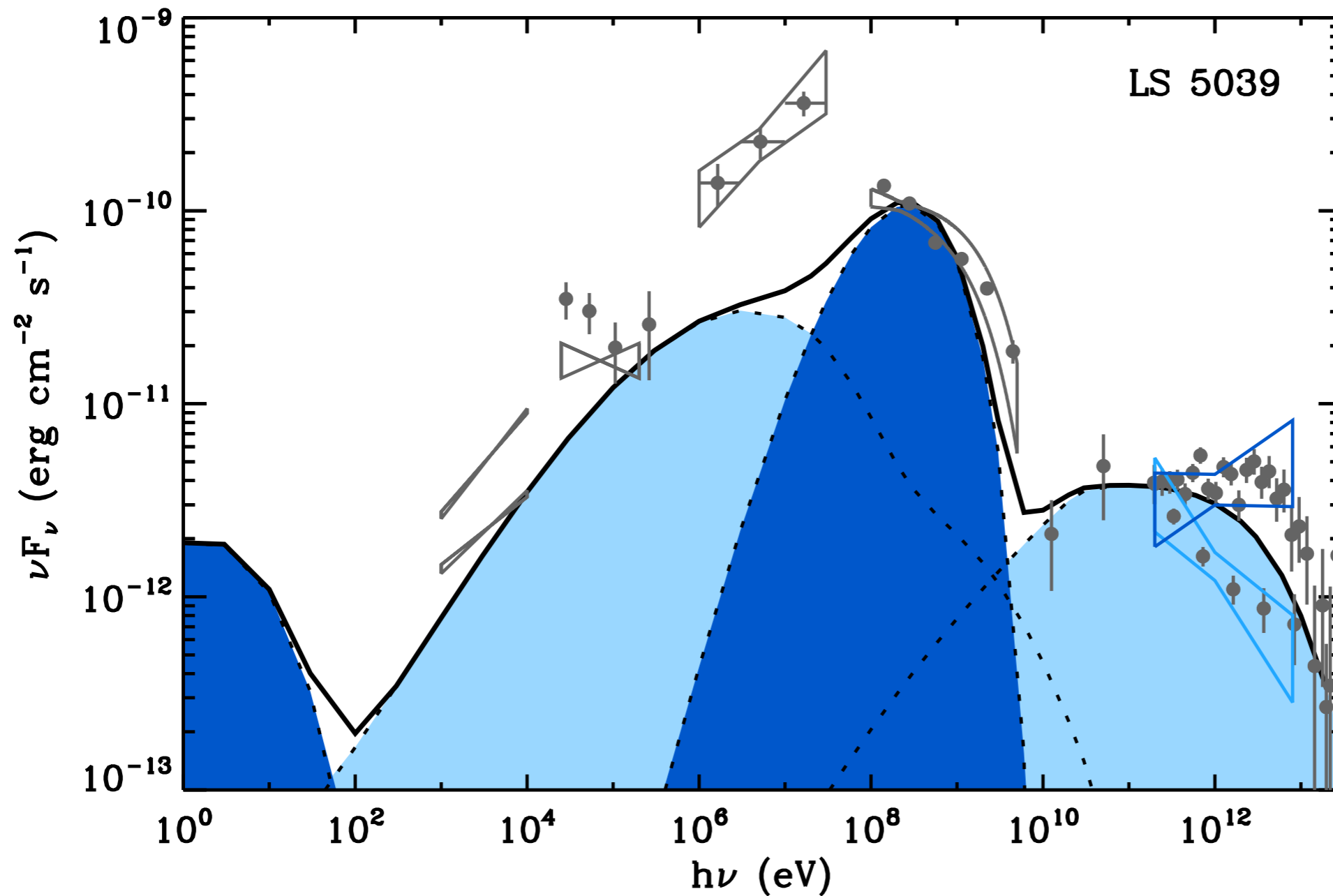
slope



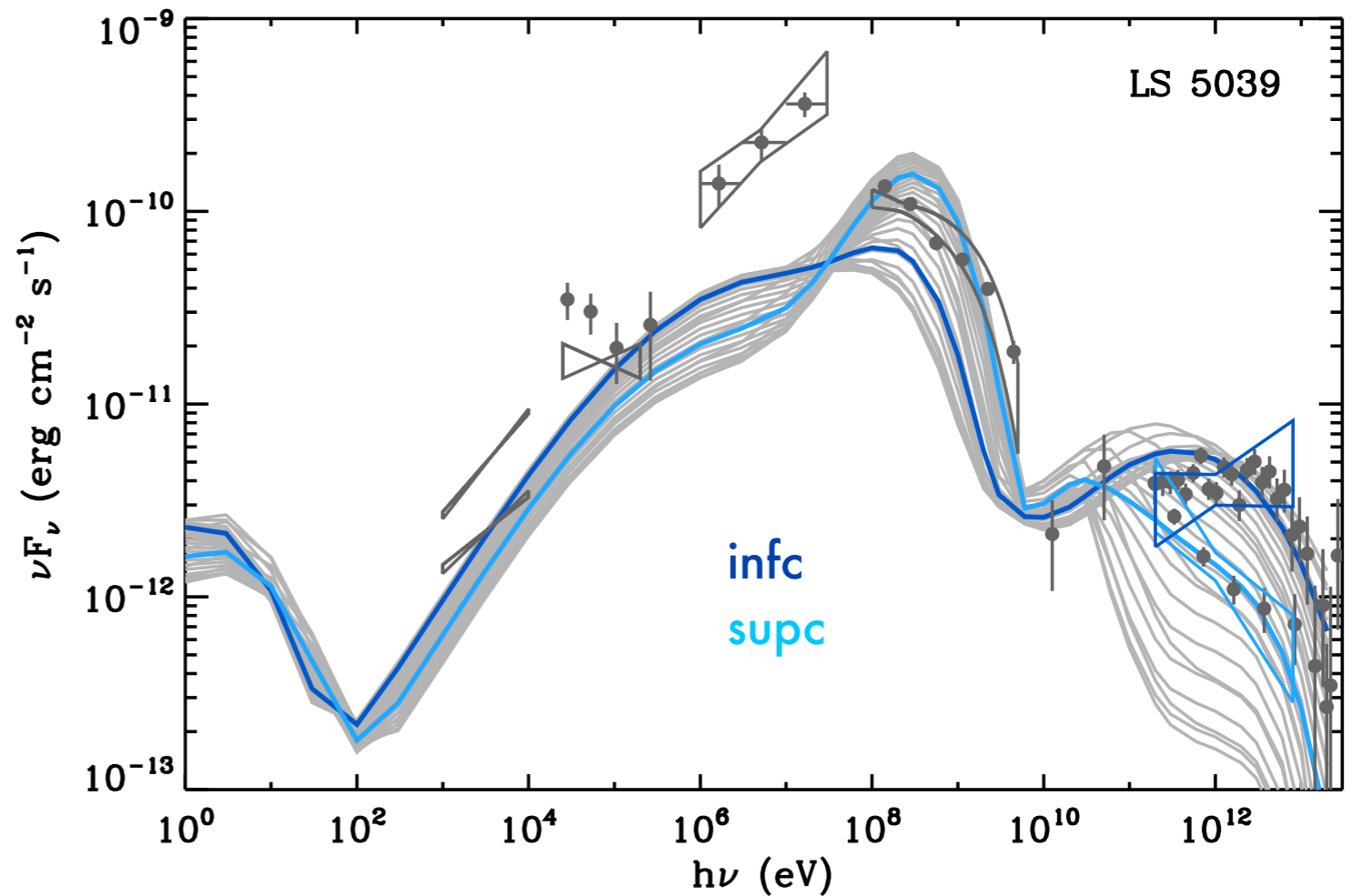
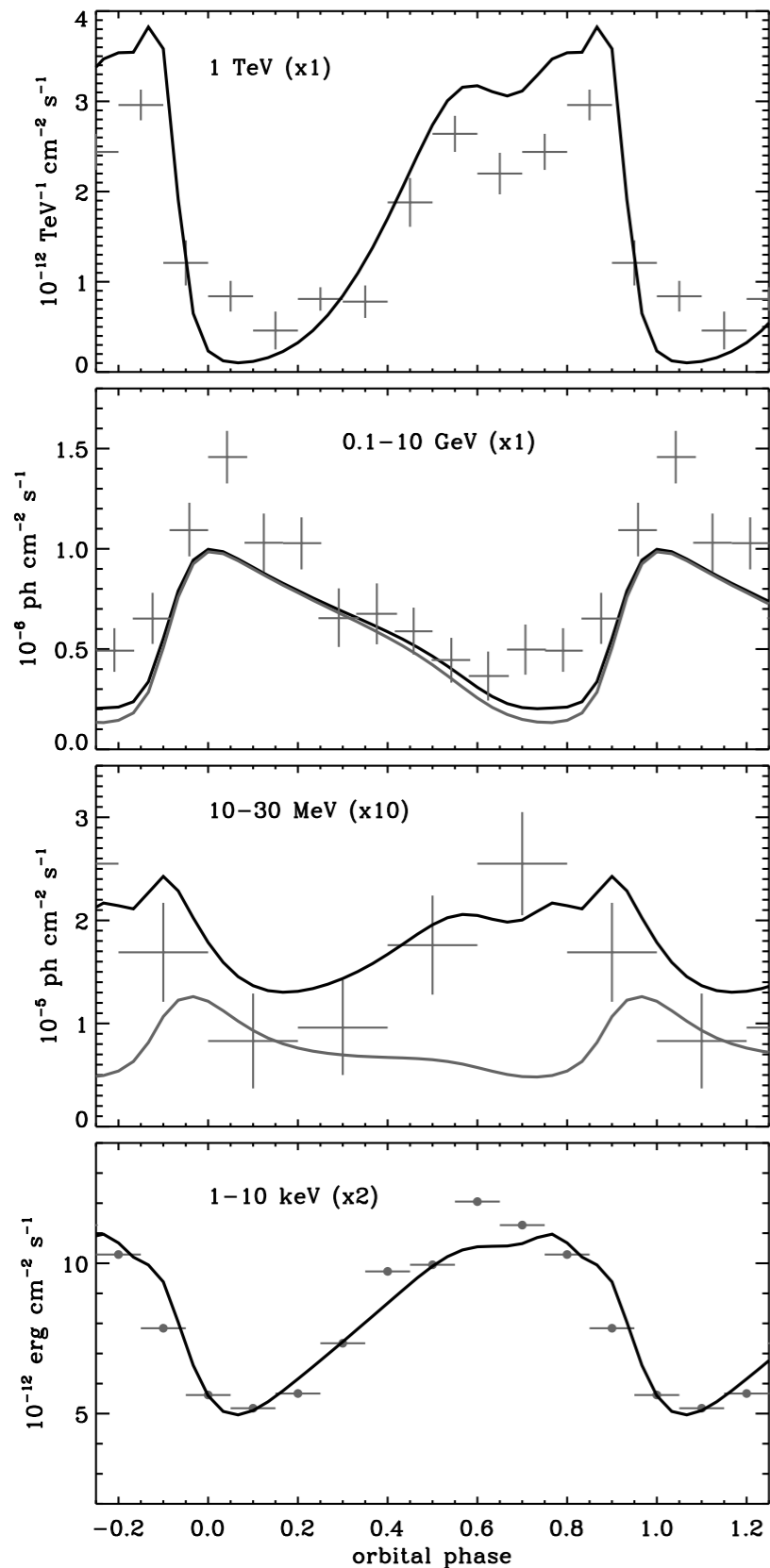
Tough to reproduce hard VHE spectrum + level of X-ray emission in LS 5039

LS 5039 requires two particle populations

Power-law with hard slope $E^{-1.5}$ + Maxwellian with $\Gamma \approx 5000$



Orbital variability in reasonable agreement



modulation set by Doppler boost $\Rightarrow i \approx 35^\circ$
(for assumed $\eta=0.1$)

First RHD-based model of LS 5039's emission

- emission concentrated at head of bow shock
- modulation mostly due to Doppler boosting
- power-law is radiatively-efficient, adiabatic cooling negligible
- maxwellian carries most of the energy but radiatively-inefficient

- particle acceleration extremely fast (Bohm)
- hard power-law, pulsar wind magnetisation $\langle\sigma\rangle\approx 1$
- shock-driven reconnection in gamma-ray binaries ?

requires (Sironi & Spitkovsky 2011)

$$\kappa \frac{R_{lc}}{R_{ts}} \geq 1 \quad \text{easier in binaries !}$$

Dubus, Lamberts, Fromang 2015, A&A, 581, A27