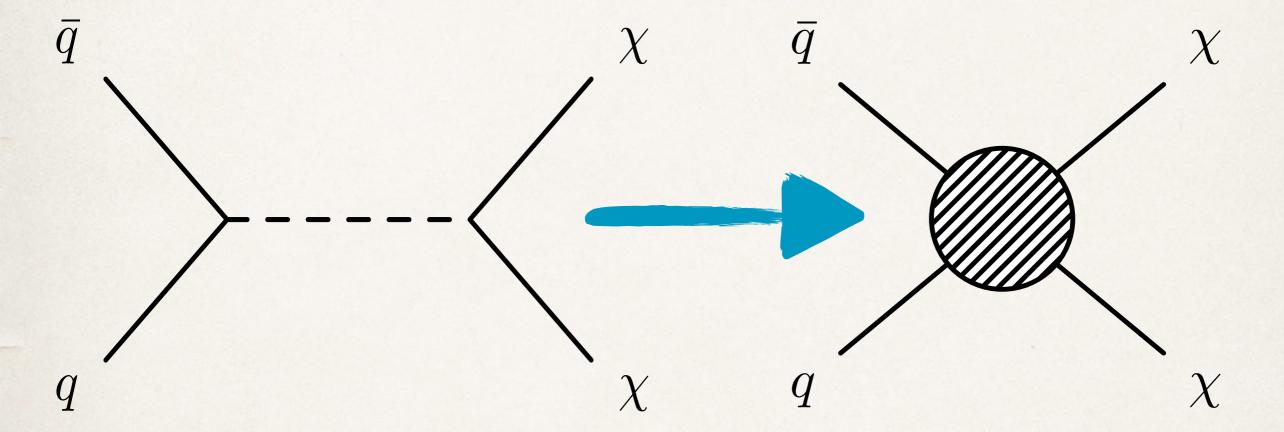


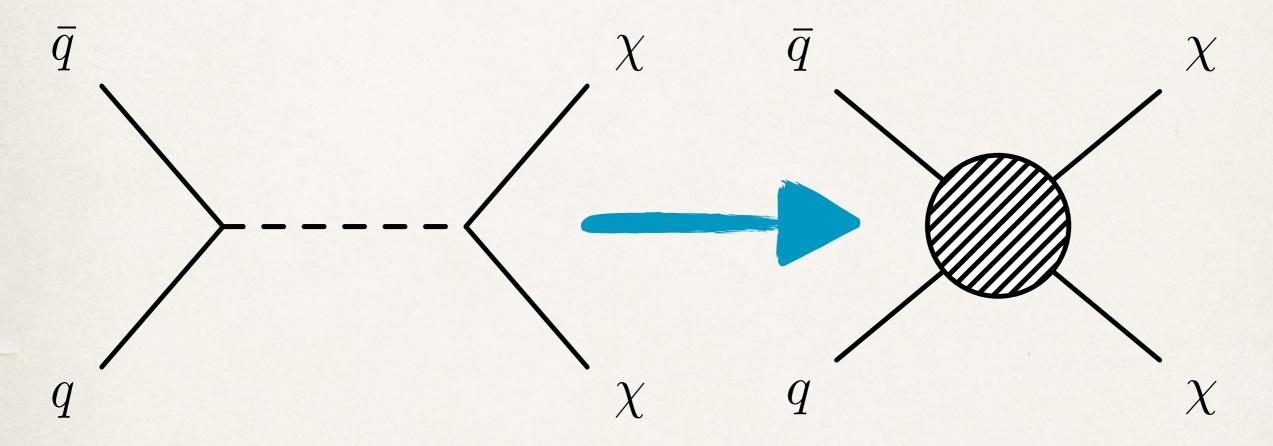
Effective Field Theories

- Integrate out the mediator, reducing parameters to $m_{\rm DM},\,M^*$ for each operator
- Easy to scan wide range of models, easy comparison between searches

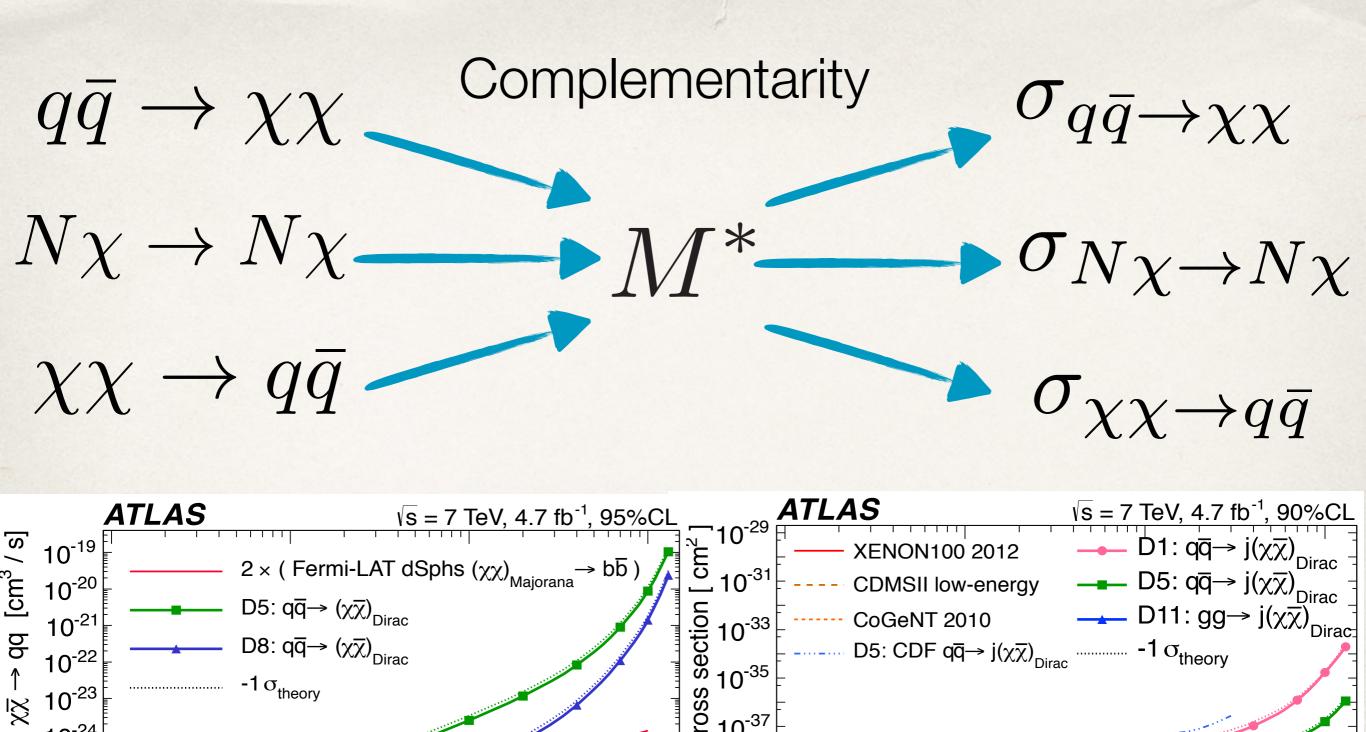


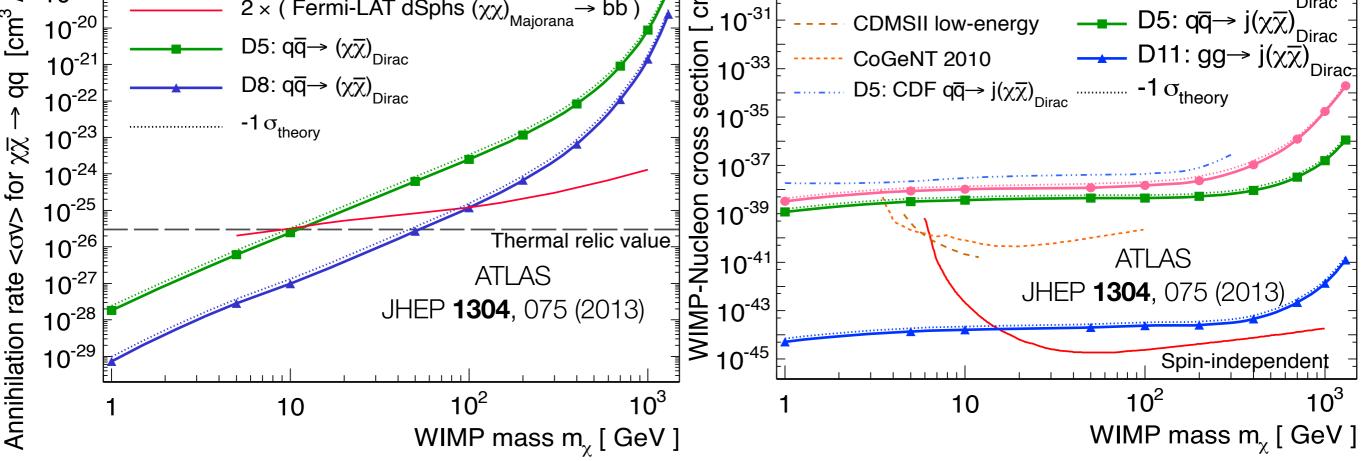
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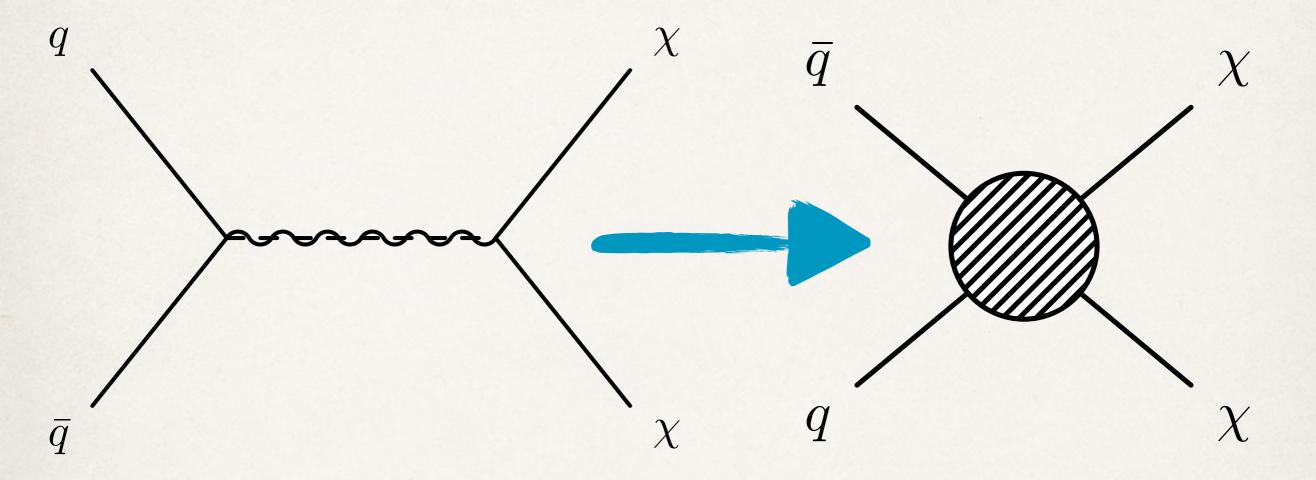


D5 =
$$(\bar{\chi}\gamma^{\mu}\chi)(\bar{q}\gamma_{\mu}q)$$
 M1 = $(\chi\chi)(\bar{q}q)$
D8 = $(\bar{\chi}\gamma^{\mu}\gamma^{5}\chi)(\bar{q}\gamma_{\mu}\gamma^{5}q)$ C3 = $(\chi^{\dagger}\partial_{\mu}\chi)(\bar{q}\gamma^{\mu}q)$



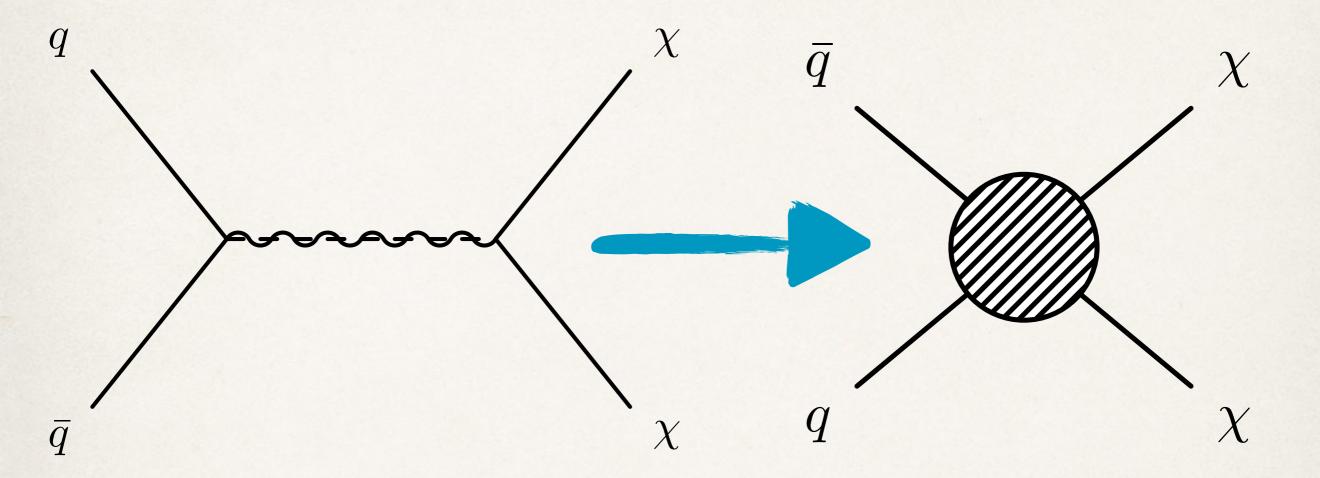


Effective field theories



$$\frac{g_q g_{\rm DM}}{M^2 - Q_{\rm tr}^2} \simeq \frac{g_q g_{\rm DM}}{M^2} \equiv \frac{1}{M^{*2}}$$

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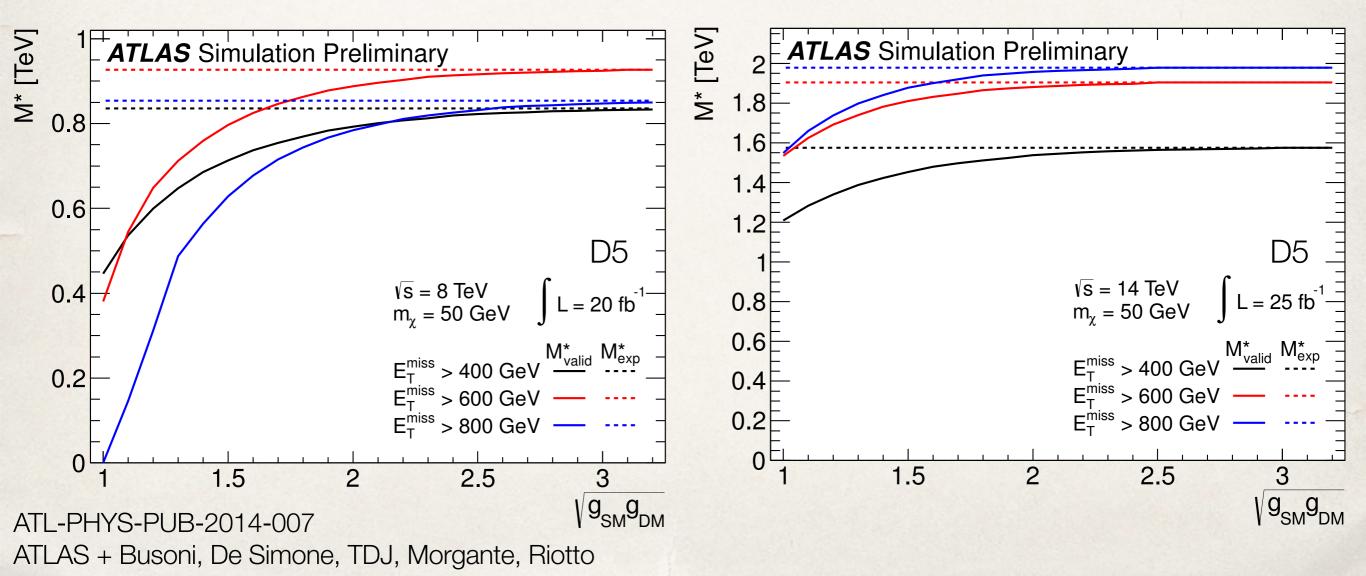
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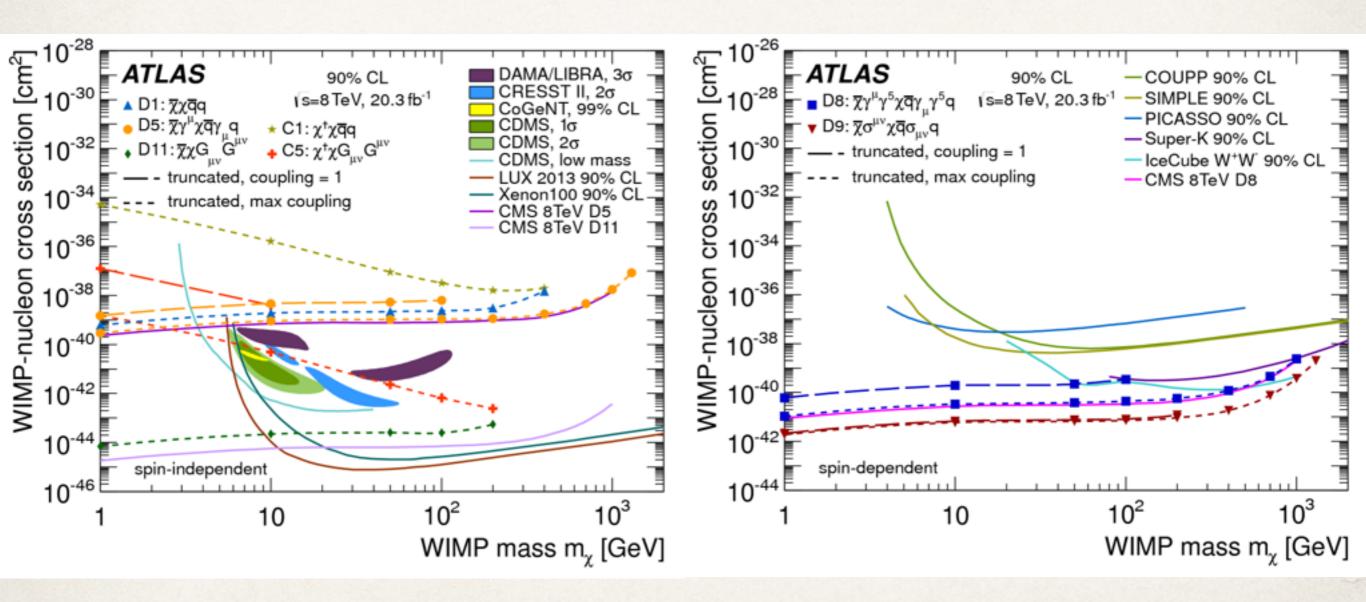
Rescaling EFT constraints

• For a given $\sqrt{g_q g_\chi}$, cut all events that don't pass

$$M \equiv \sqrt{g_q g_\chi} M^* \ge Q_{\rm tr}$$



Rescaling EFT constraints



Eur. Phys. J. C (2015) 75:299 ATLAS + Busoni, De Simone, TDJ, Morgante, Riotto

- Minimum particle content necessary to describe new physics;
 Usually one DM candidate and one mediator
- Enlarged parameter space: Minimum of 4 parameters
- Even the simplest models have more unless you start making choices and assumptions

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m DM}$$
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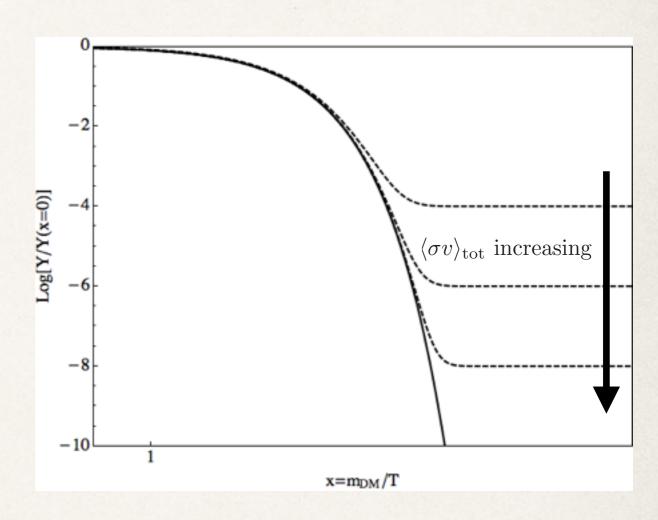
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Thermal relic dark matter

- Dark matter in thermal equilibrium at large T
- When m_{DM} > T, comoving abundance drops exponentially
- As universe expands, abundance freezes out
- Annihilation rate controls abundance at freezeout



$$\langle \sigma v \rangle_{\rm tot} \simeq \frac{4.8 \times 10^{-10} \, {\rm GeV}^{-2}}{\Omega_{\rm DM} h^2}$$

- Relic density is sensitive to the annihilation channels of the model
- Collider searches assume minimal channels:
 Mediator couples to quarks only, with equal coupling to each flavour
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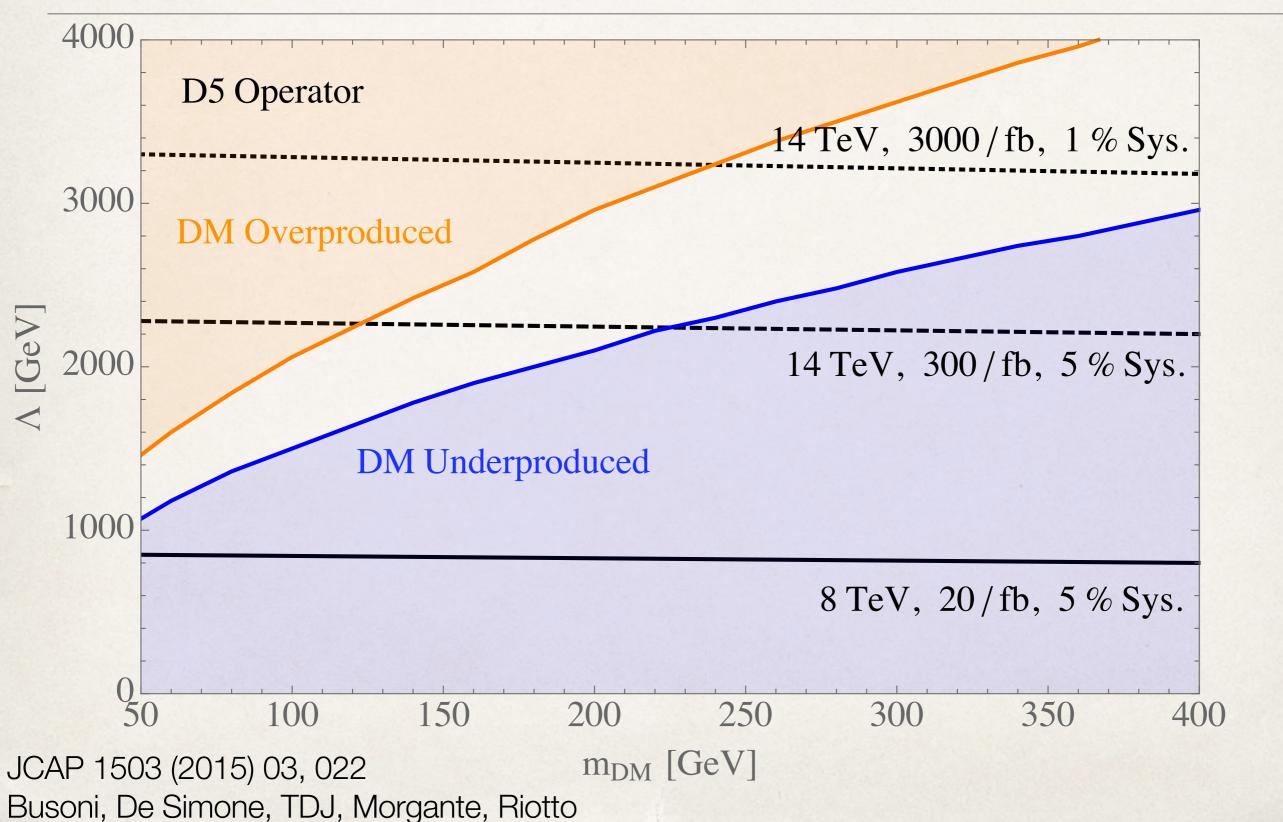
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$$1.0 \times 10^{-9} \,\mathrm{GeV}^{-2} \lesssim \langle \sigma v \rangle_{\mathrm{min}} \lesssim 4.0 \times 10^{-9} \,\mathrm{GeV}^{-2}$$

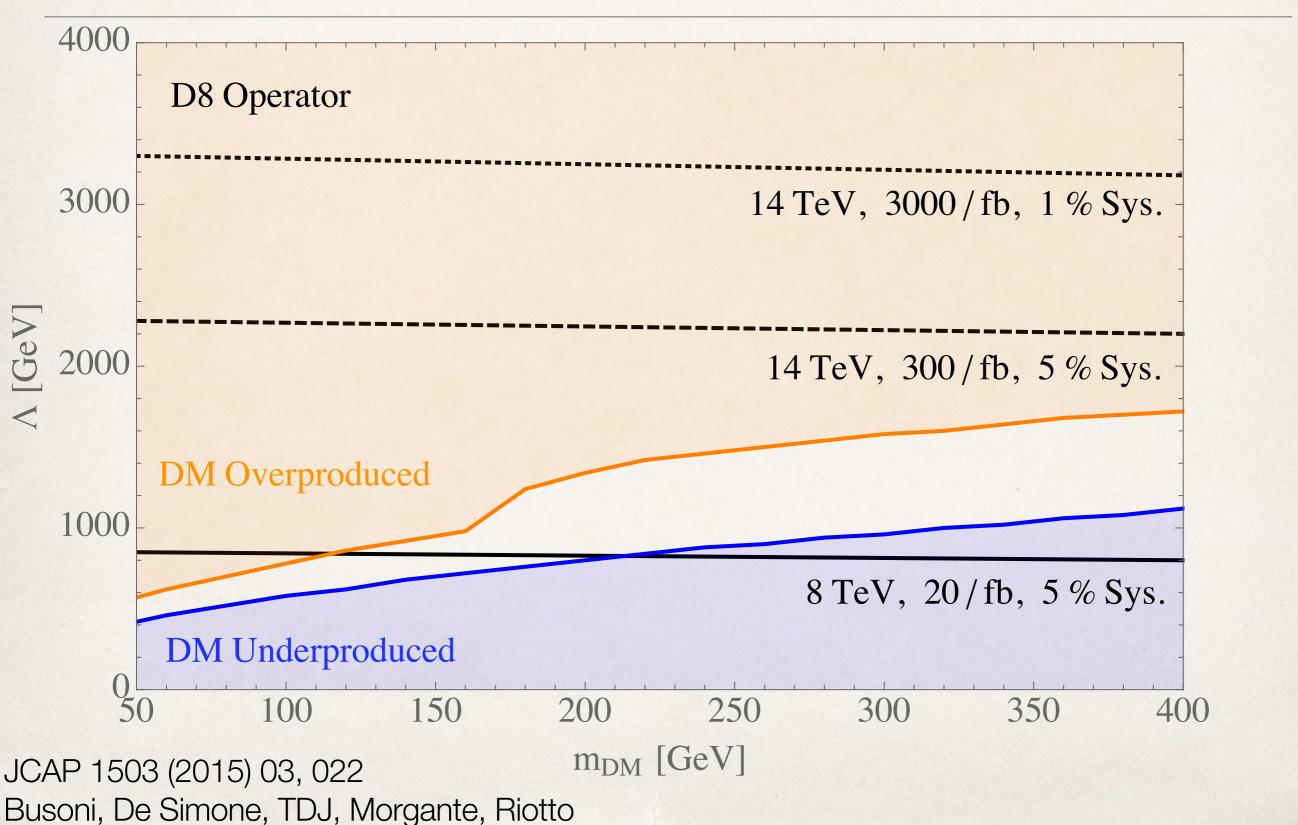
Effective operator results

$$D5 = \frac{1}{\Lambda^4} (\bar{\chi}\gamma^{\mu}\chi)(\bar{q}\gamma_{\mu}q)$$



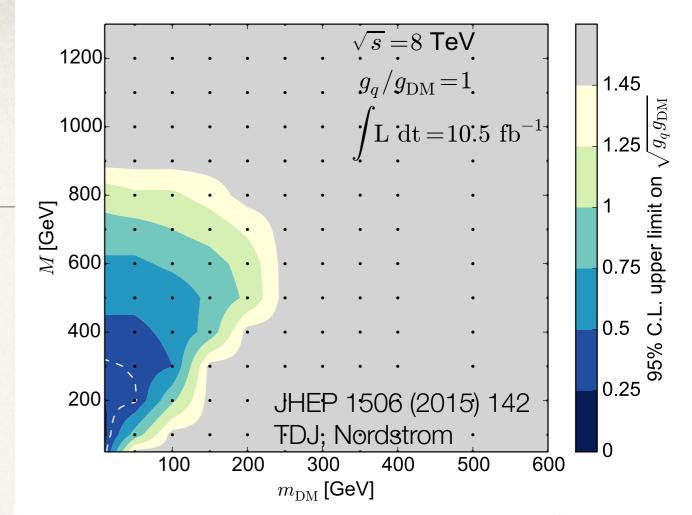
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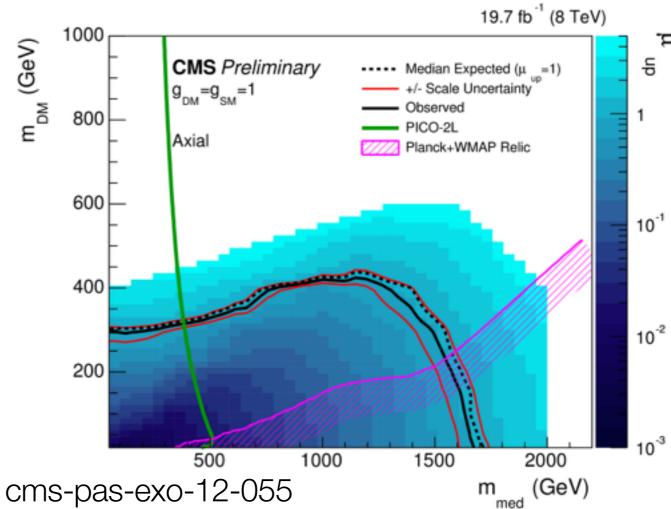
$$D8 = \frac{1}{\Lambda^4} (\bar{\chi}\gamma^{\mu}\gamma^5\chi)(\bar{q}\gamma_{\mu}\gamma^5q)$$



Simplified models

- We can prioritize the model space - e.g. direct detection is better for vector models, LHC for axial-vector
- The parameter space is still large!
 2 choices:
 - Scan over coupling, or
 - fixed benchmark
- Scan is more comprehensive and intuitive but technically difficult: Width changes at each point in 4D parameter space
- Benchmarks allow comparison between experiments



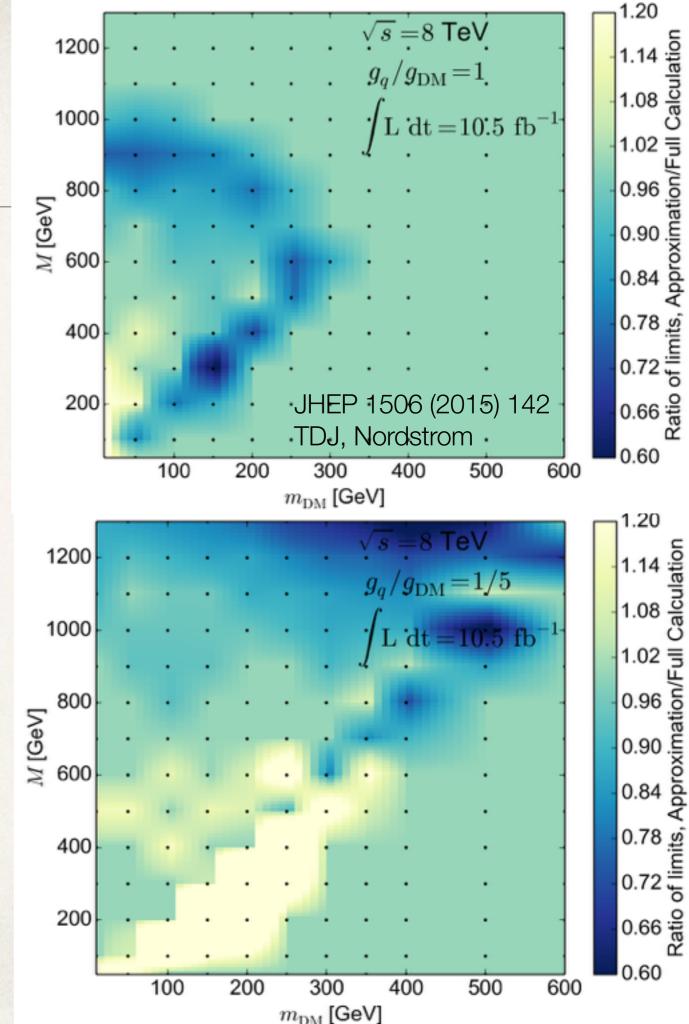


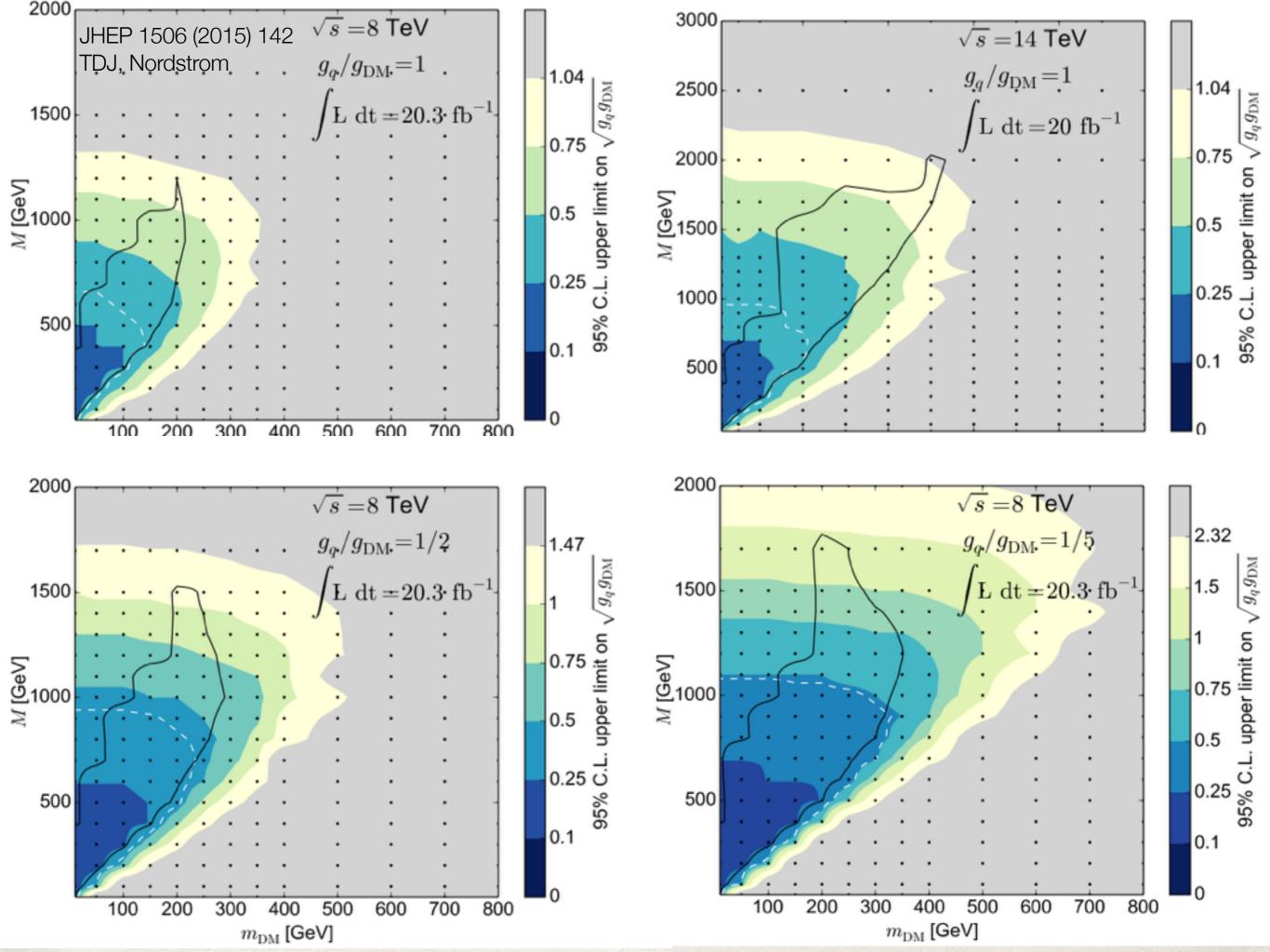
Simplified models

 Some approximations can help

$$\sigma \propto \begin{cases} g_q^2 g_{\rm DM}^2 / \Gamma_{\rm OS} & \text{if } M > 2m_{\rm DM} \\ g_q^2 g_{\rm DM}^2 & \text{if } M < 2m_{\rm DM} \end{cases}$$

Avoids scan over gq, gDM





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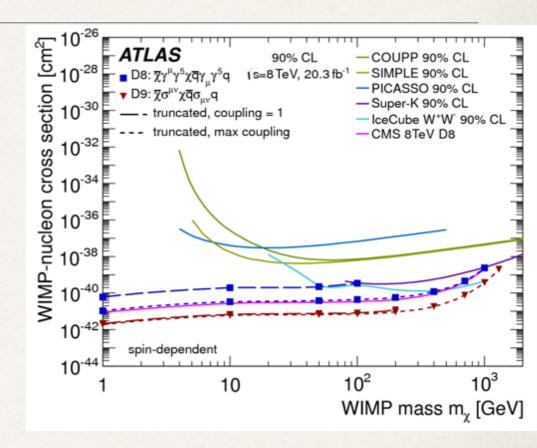
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$$\bar{\chi}\gamma_{\mu}\gamma_{5}\chi\,\bar{f}\gamma^{\mu}f,\ \bar{\chi}\gamma_{\mu}\gamma_{5}\chi\,\bar{f}\gamma^{\mu}\gamma_{5}f$$
 $\longrightarrow \vec{s}_{\chi}\cdot\vec{s}_{N},\ i\vec{s}_{\chi}\cdot(\vec{s}_{N}\times\vec{q}),\ \vec{s}_{\chi}\cdot\vec{v}^{\perp}$

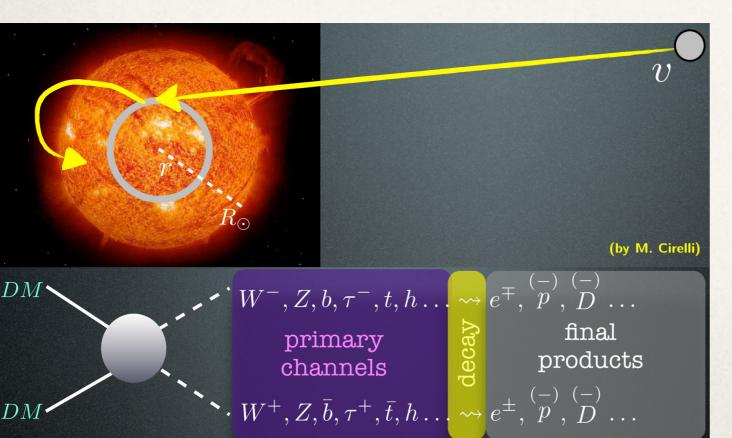
OSD Suppressed, but neither $\sigma_{\rm SI}$ nor $\sigma_{\rm SD}$

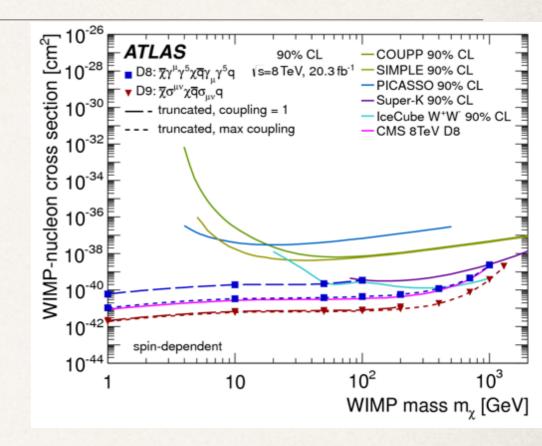
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- Solar neutrinos provide a unique window on DM scattering
- DM accumulates in the sun

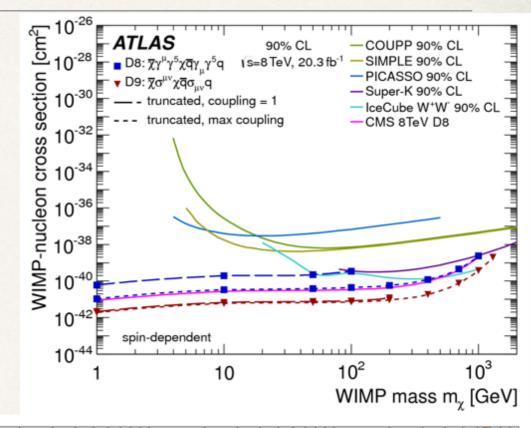
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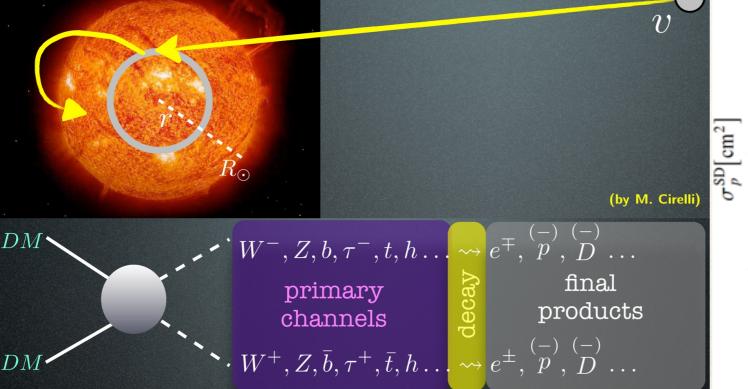


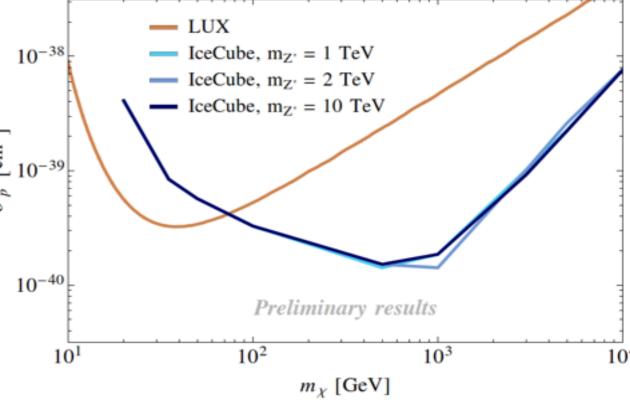


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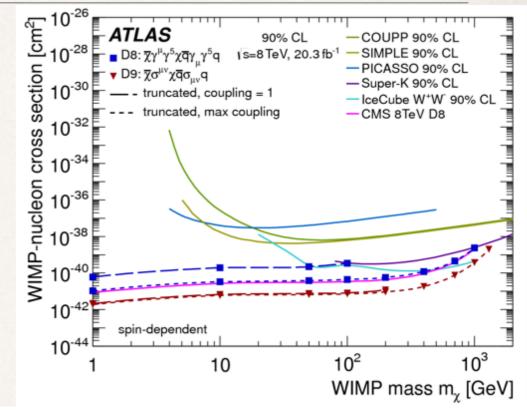


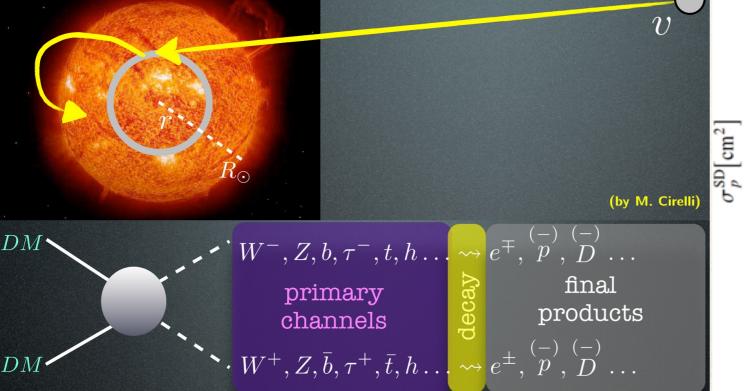


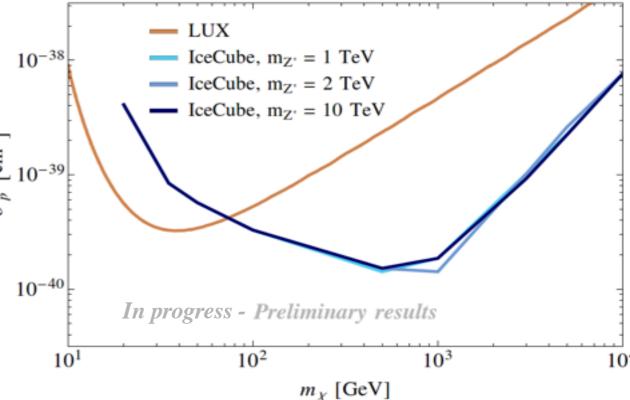


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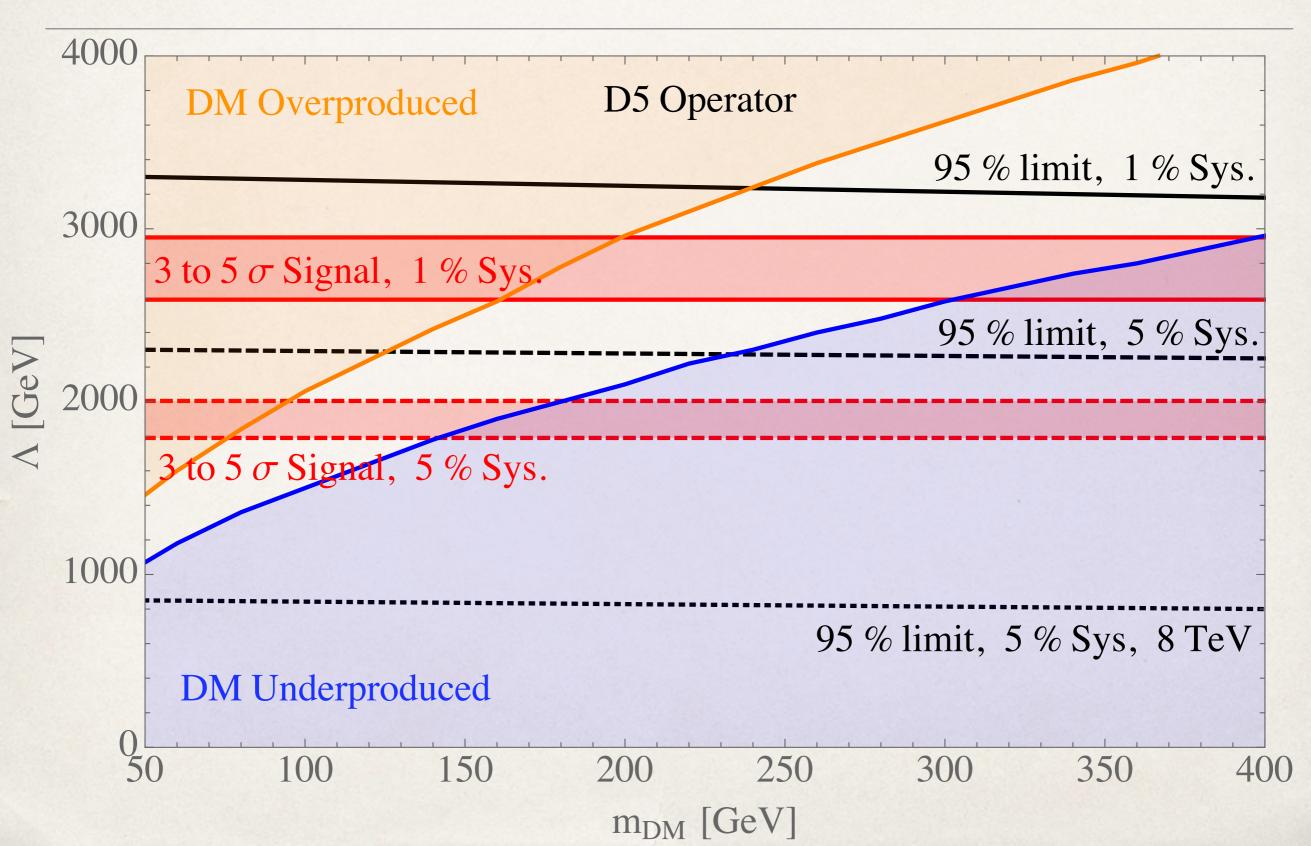


Conclusion

- EFTs remain a useful tool for comparison between experiments, but have limited validity for LHC searches.
 Move to simplified models is necessary
- The difficulties associated with simplified models are eased by considering all available information and taking a different approach to relic density constraints
- LHC is strongest in searches for models which have suppressed scattering and annihilation rates, but building consistent models is important

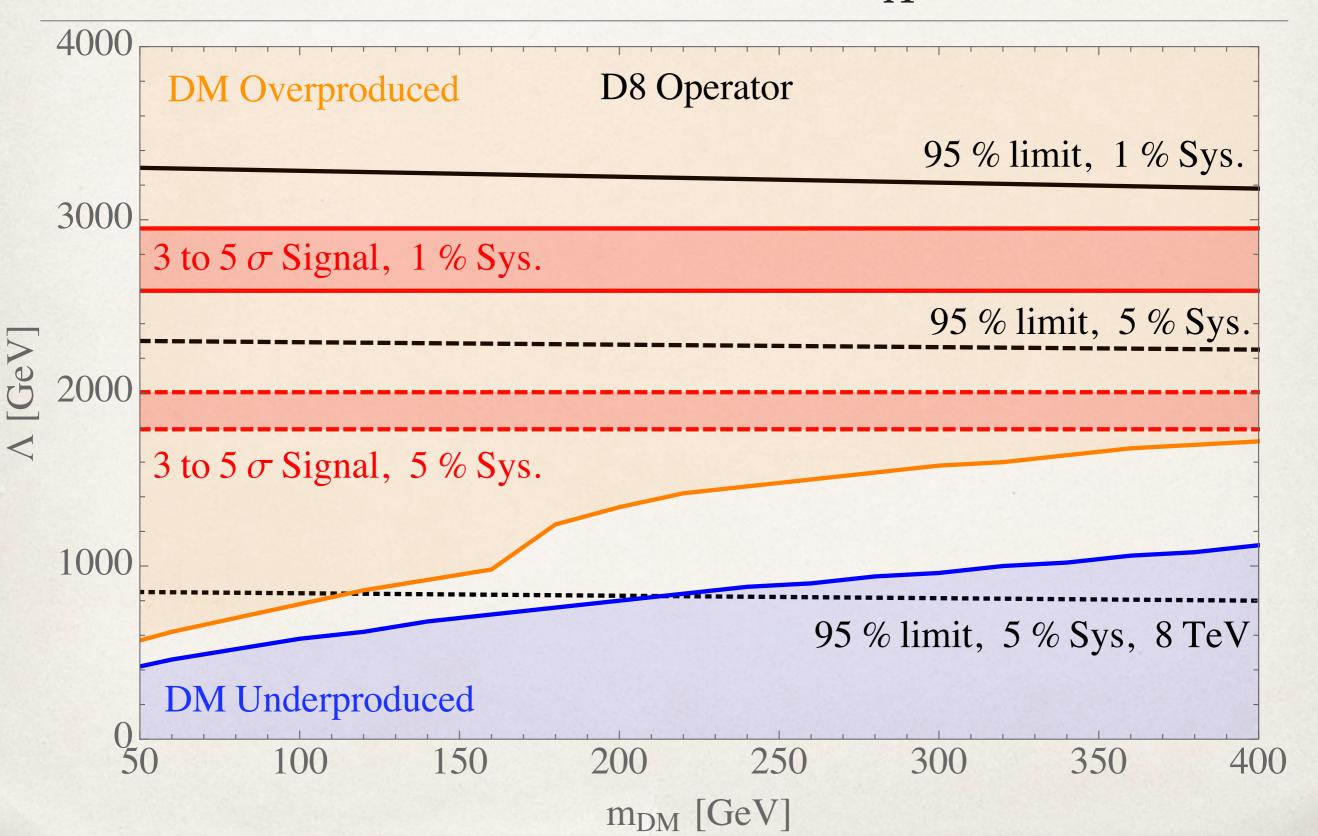
Backup

Effective operator results



Effective operator results

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Simplified model results

