

Probing the Higgs-light-quarks couplings & uncharming the Higgs

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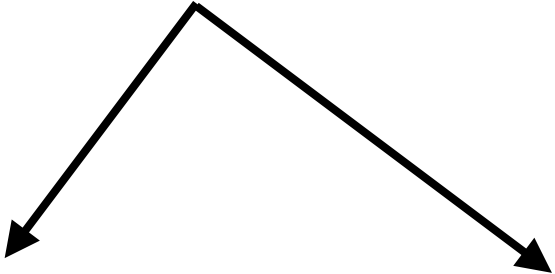
Delaunay, Golling, GP & Soreq (13)
Bodwin, Petriello, Stoynev & Velasco (13)
Kagan, GP, Petriello, Soreq, Stoynev & Zupan (14)
GP, Soreq, Stamou & Tobioka (Feb/15)



Aspen Winter 2015
Exploring the Physics Frontier with Circular Colliders

Status: two paths to measuring light-quark Yukawas

Recent rapid (th+exp) progress (things are still preliminary, tons of info missing)



Inclusive (c-tagging)

exclusive (formalism: Neubert's talk)

quarks	c	$udsc$
th.	Delaunay, Golling, GP & Soreq (13) GP, Soreq, Stamou & Tobioka (Feb/15)	Bodwin, Petriello, Stoynev & Velasco (13); Kagan, GP, Petriello, Soreq, Stoynev & Zupan; Bodwin, Chung, Ee, Lee & Petriello (14); Grossmann, Konig & Neubert (15); GP, Soreq, Stamou & Tobioka (Feb/15)
exp.	ATLAS-CONF-2013-068 ATLAS-CONF-2014-063	ATLAS: 1501.03276

Outline

- ◆ Intro': Higgs & flavor physics within the standard model (SM) & beyond.
- ◆ Charming the Higgs, an inclusive approach. (charm-tagging)
Recent developments, establishing Higgs-quark non-univ. & more.
- ◆ Brief: exclusive approach, unique window to Higgs-light quarks couplings.
- ◆ Some projections & summary.

Higgs & flavor physics within the SM

◆ Higgs in minimal SM, 2 roles:

- (i) induce electroweak (EW) gauge boson masses & unitarization (high-E consistency);
- (ii) induce fermion masses & unitarization (high-E consistency).

(i) was already tested in a quantitative way (ii) much less & mostly for 3rd gen'.
We focus on (ii), significant progress can be made.

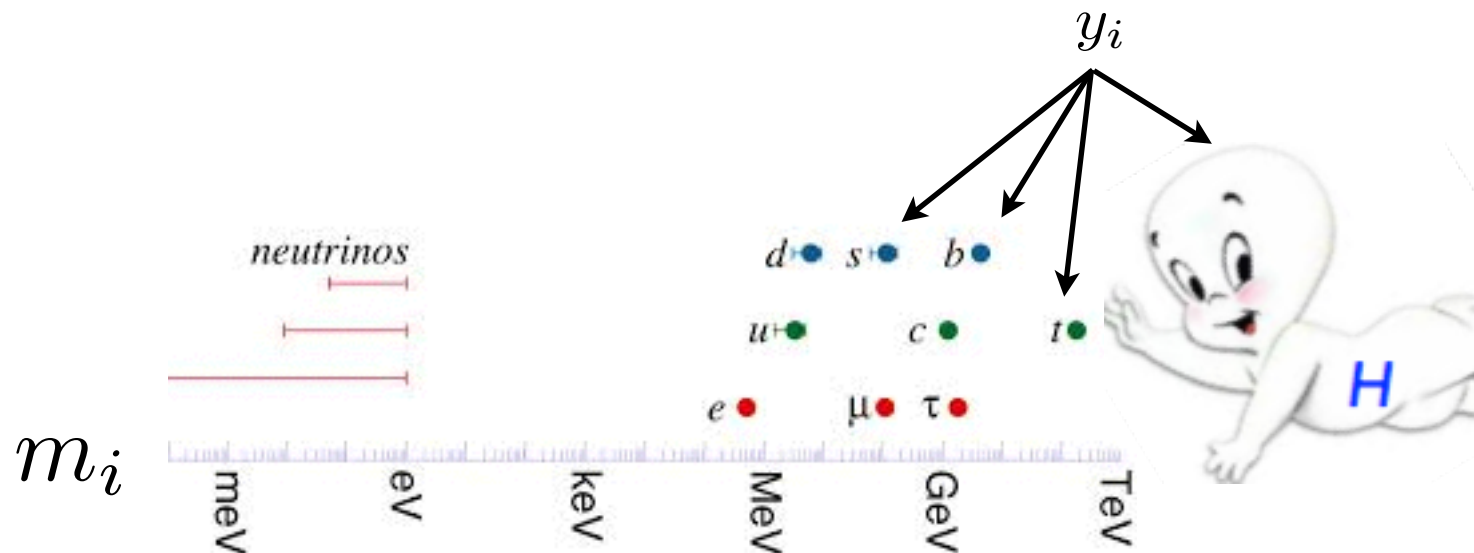
◆ Recall: in the SM we have 2 type of interactions:

- (i) gauge interactions: these are flavor blind/universal/same for all quarks;
- (ii) Yukawa interactions: generation-dependent, non-universal, but to a single scalar.

Higgs & flavor physics within the SM

- ◆ The above 2 facts + renormalizability leads to a simple relations, up two small corrections (1loop+GIM suppressed)

SM: Higgs couples like the mass - $y_i = m_i/v$

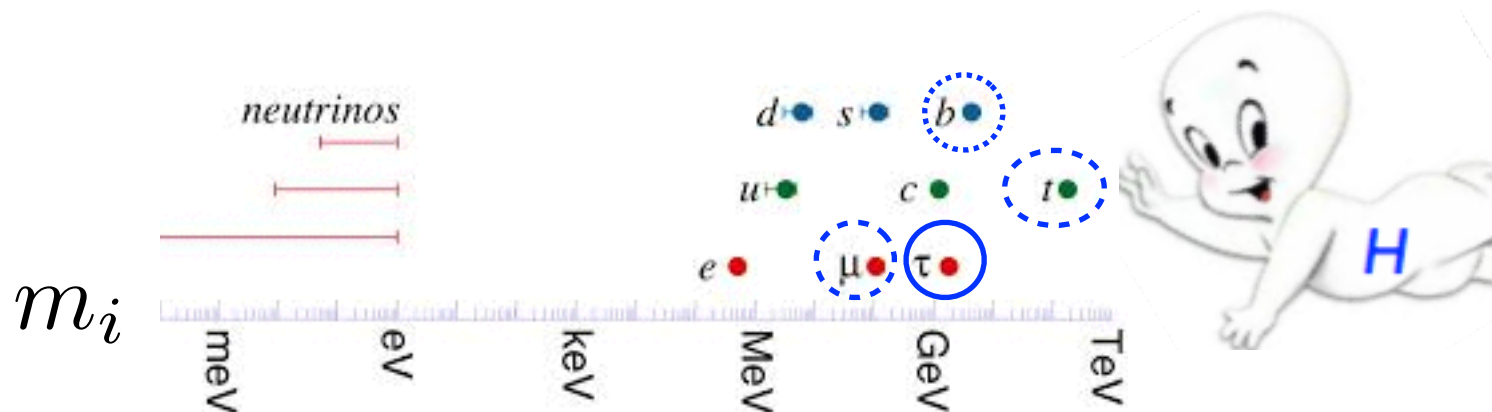


How much is known experimentally

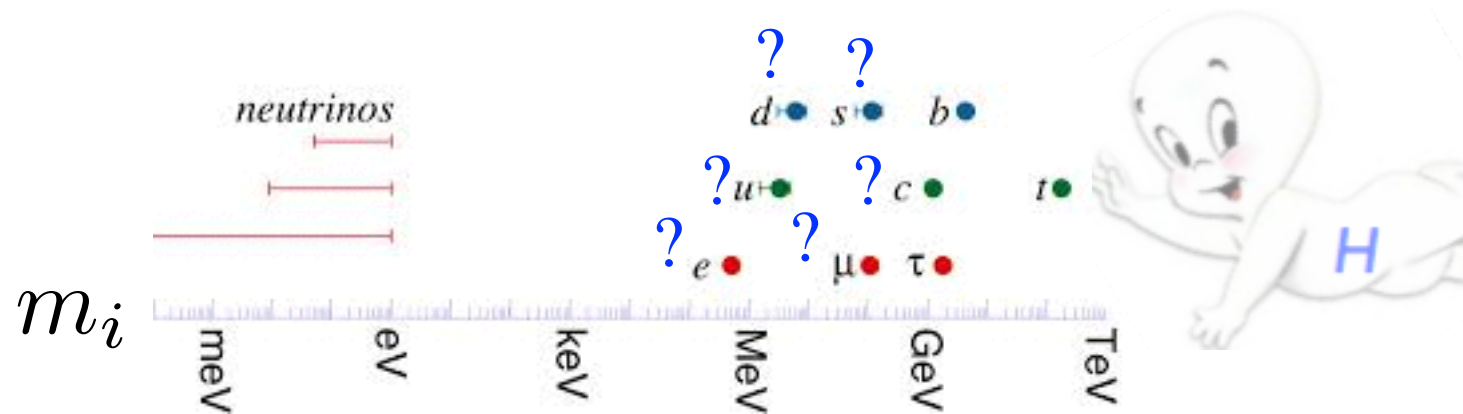
μ_x	ATLAS+CMS
τ	0.97 ± 0.23
b	0.71 ± 0.31
t	2.4 ± 0.81

μ_μ : $\sigma.\text{Br} < 7.0 (7.2)(\sigma.\text{Br})_{\text{SM}}$

Universal couplings
~260 times SM

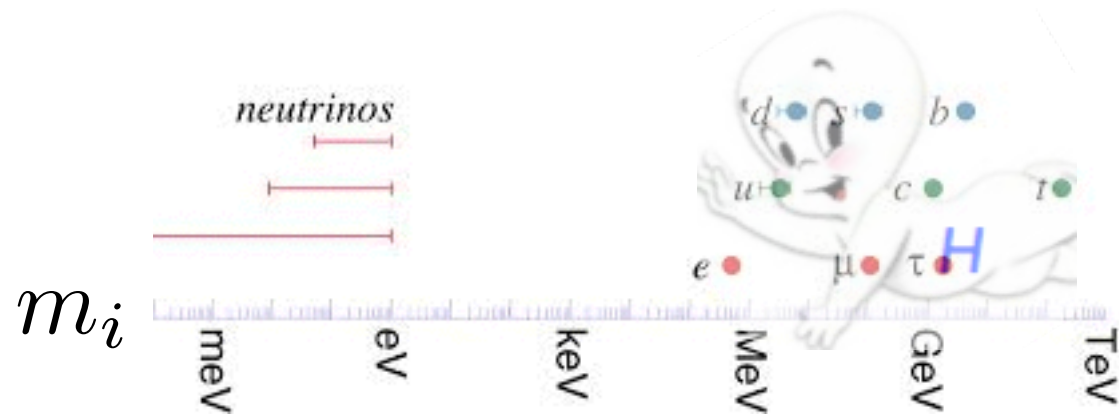


Higgs & flavor physics within the SM



Does the Higgs couples like the mass - $y_i = m_i/v$??

Higgs & flavor physics within the SM



Does the Higgs couples like the mass - $y_i = m_i/v$??

- ◆ This could dramatically change if non-SM exists, especially because the Higgs is light and its decay (& production) is controlled by small couplings. Let us see a trivial example.

Charming the Higgs

- Currently not much known directly on the charm Yukawa:

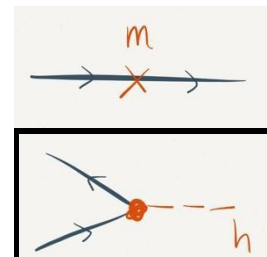
(i) SM - $y_c = m_c/v \sim 0.4\%$ $\Rightarrow BR(H \rightarrow c\bar{c}) \sim 4\%$, very non-trivial to observe...

- However, as $y_b \sim 2\%$ & $BR(H \rightarrow b\bar{b}) \sim 60\%$, Higgs collider pheno' is susceptible to small perturbation.

- Enlarging charm Yukawa by few leads to dramatic changes, for instance:

Delaunay, Golling, GP & Soreq (13)

$$\mathcal{L}_{\text{EFT}} \supset \lambda_{ij}^u \bar{Q}_i \tilde{H} U_j + \frac{g_{ij}^u}{\Lambda^2} \bar{Q}_i \tilde{H} U_j (H^\dagger H) + \text{h.c.}$$



$$= \frac{v}{\sqrt{2}} \left(\lambda_{ij}^u + g_{ij}^u \frac{v^2}{2\Lambda^2} \right),$$

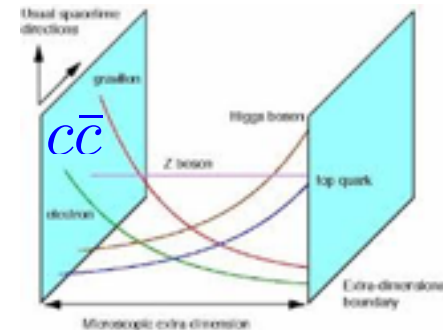
$$= \frac{1}{\sqrt{2}} \left(\lambda_{ij}^u + 3g_{ij}^u \frac{v^2}{2\Lambda^2} \right).$$

$$\Lambda \simeq \frac{44 \text{ TeV}}{\sqrt{c_e - 1}}$$

$$\mathcal{L}_0 = \frac{h}{v} \left[c_V (2m_W^2 W_\mu^+ W^{\mu-} + m_Z^2 Z_\mu Z^\mu) - \sum_q c_q m_q \bar{q} q - \sum_\ell c_\ell m_\ell \bar{\ell} \ell \right],$$

- Or is it simply technicolor for the light quarks ?

If you really care, more models: Delaunay, Grojean & GP (13); Kagan, GP, Volansky & Zupan (09); Dery, Efrati, Hiller, Hochberg & Nir (13); Giudice & Lebedev (08)



Charming the Higgs, current status & few projections

Delaunay, Golling, GP & Soreq (13)

- ◆ Ball park bounds are from Higgs “invisible” bound (assumes $c_v=1$):

if all other “visible” couplings set to SM values:

$$Br_{inv} \sim < 22\% \text{ @95\%CL}$$

adding a new physics source of ggh: $Br_{inv} \sim < 50\% \text{ @95\%CL}$

$BR(H \rightarrow b\bar{b})$ is significantly suppressed:

$$BR_{h \rightarrow b\bar{b}} = \frac{BR_{h \rightarrow b\bar{b}}^{\text{SM}}}{1 + (|c_c|^2 - 1)BR_{h \rightarrow c\bar{c}}^{\text{SM}}} \approx 40\% (20\%)$$

with $c_{gg} > 0$

$$\hat{c}_{gg} = c_{gg} + \left[1.3 \times 10^{-2} c_t - (4.0 - 4.3i) \times 10^{-4} c_b - (4.4 - 3.0i) \times 10^{-5} c_c \right],$$

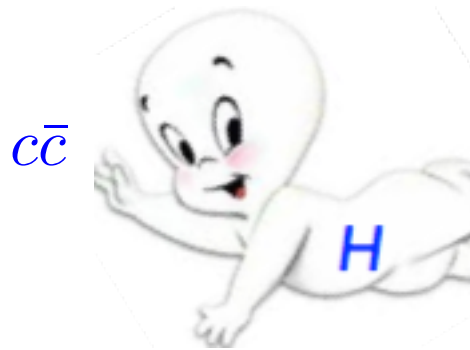
$$\sigma_{c\bar{c} \rightarrow h} \simeq 3.0 \times 10^{-3} |c_c|^2 \sigma_{gg \rightarrow h}^{\text{SM}},$$

assume instead a speculative $\epsilon_c = 40\%$ c-tagging efficiency:

$$\rightarrow \mu_{bb+cc} \approx 0.9 (0.6) \text{ @8TeV}$$

Uncharming the Higgs, establishing non-universality & more

GP, Soreq, Stamou & Tobioka (Feb/15)



Before talking about our work, 2 slides about an experimental break through

Charm tagging at the LHC

- ◆ In new ATLAS search for stop decay to charm + neutralino ($\tilde{t} \rightarrow c + \chi^0$) charm jet tagging has been employed for the first time at LHC

ATLAS-CONF-2013-068

- ◆ charm jets identified by combining “information from the impact parameters of displaced tracks and topological properties of secondary and tertiary decay vertices” using multivariate techniques
 - ‘medium’ operating point: c-tagging efficiency = 20%, rejection factor of 5 for b jets, 140 for light jets. #’s obtained for simulated $t\bar{t}$ events for jets with $30 < p_T < 200$, and calibrated with data

More recently, constraining (non-deg.) scharm

◆ An interesting viable possibility is anarchic squark spectrum. Nir & Seiberg (93)

◆ Scenario still viable and the bounds on scharm are very weak.

Gedalia, Kamenik, Ligeti & GP (12)
Mahbubani, Papucci, GP, Ruderman & Weiler (13)

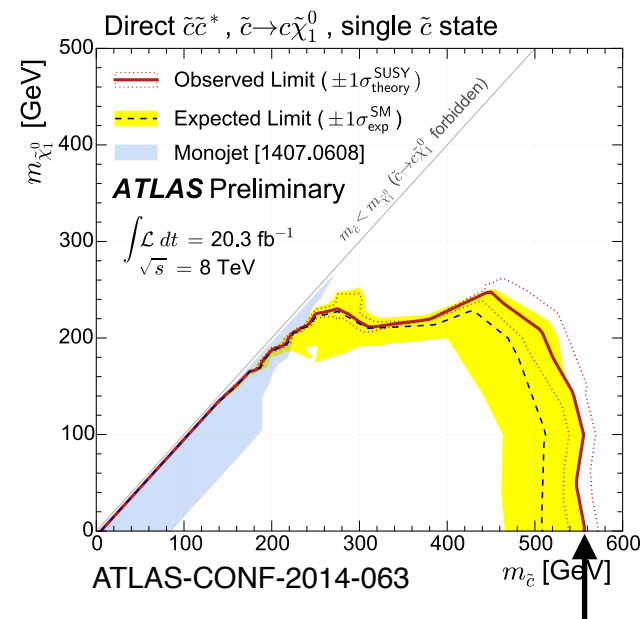
◆ Has potential consequences for naturalness (“flavorful naturalness”).

Blanke, Giudice, Paradisi, GP & Zupan (13)

◆ ATLAS: light scharm search \w new working point for charm tagging:

Search for Scalar-Charm Pair-Production with the ATLAS Detector in pp
Collisions at $\sqrt{s} = 8$ TeV

$$\epsilon_c = 19\% \quad \epsilon_b = 12\%$$



Executive sum.: Constraining Higgs-charm univ.

GP, Soreq, Stamou & Tobioka (Feb/15)

- ◆ Bottom line: can use existing data to constrain Higgs-quarks univ..
 - (i) Direct constraint: recast $VH(bb)$, taking advantage of 2 working point $c_c < 250$.
 - (ii) the recent ATLAS search to $h \rightarrow J/\psi\gamma$ (see later) yield $c_c < 210$;
(assumes Higgs coupling to two photons and/or four leptons is not significantly modified by new physics);
 - (iii) the direct measurement of the total width yield $c_c < 150$ (ATLAS), 120 (CMS).;
 - (iv) Global fit to the Higgs signal strength, $c_c < 6$.
 - (v) tth data $\Rightarrow c_t > 0.9$ (equivalence to $c_c > 280$).

#1 Direct constraint: recast $VH(bb)$ (preliminary)

GP, Soreq, Stamou & Tobioka (Feb/15)

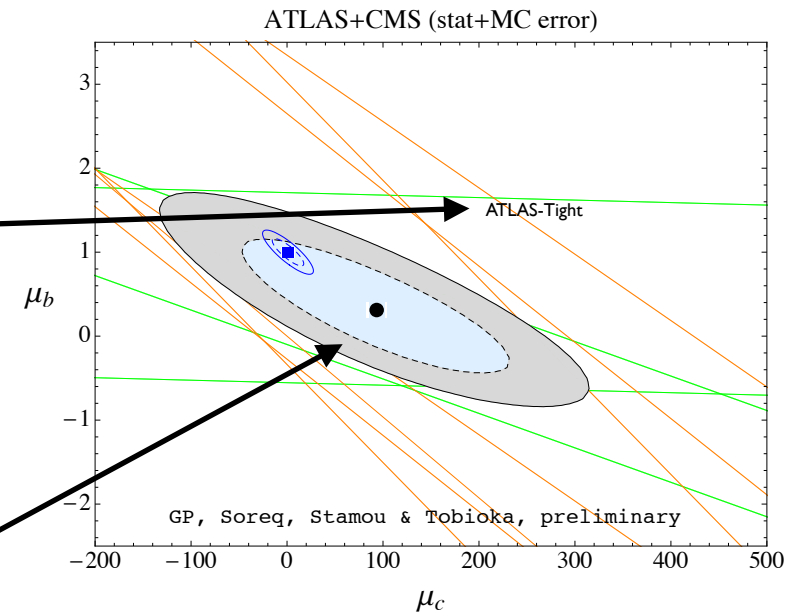
- ◆ Idea: use several charm-tagging working points of ATLAS & CMS in their $VH(bb)$ analysis.

$$\mu_b = \frac{\sigma}{\sigma_{SM}} \frac{BR_{b\bar{b}}}{BR_{b\bar{b}}^{SM}} \rightarrow \mu_b + \frac{Br_c^{SM}}{Br_b^{SM}} \frac{\epsilon_{c1} \epsilon_{c2}}{\epsilon_{b1} \epsilon_{b2}} \mu_c$$

where $\epsilon_{b_{1,2}}$ and $\epsilon_{c_{1,2}}$ are efficiencies to tag jets originating from bottom and charm quark, respectively. μ_c is normalized to be 1 in a case of the SM.

- ◆ Each working point yields flat direction:

ATLAS	Med	Tight	CMS	Loose	Med1	Med2	Med3
ϵ_b	70%	50%	ϵ_b	88%	82%	78%	71%
ϵ_c	20%	3.8%	ϵ_c	47%	34%	27%	21%

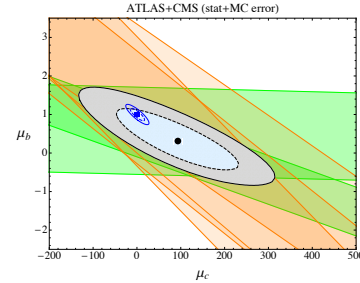


- ◆ However, combining points => bound.

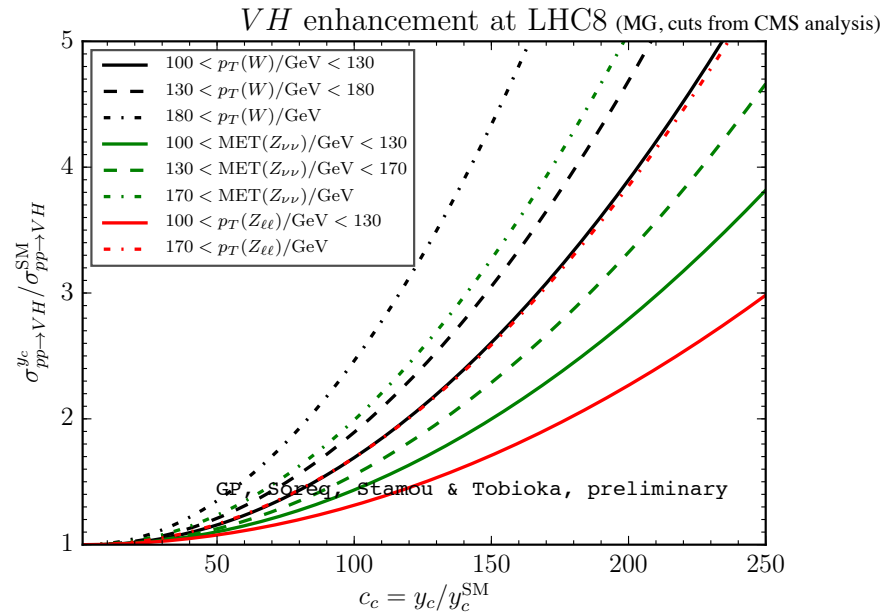
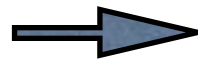
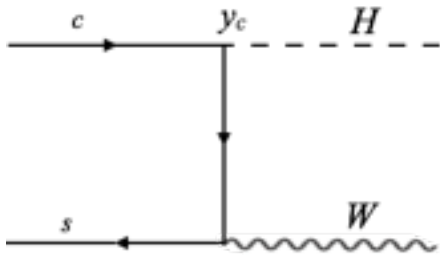
New production mechanism $VH(bb)$ (*preliminary*)

GP, Soreq, Stamou & Tobioka (Feb/15)

◆ $\mu_c = \frac{\sigma}{\sigma^{\text{SM}}} \frac{\text{Br}}{\text{Br}_c^{\text{SM}}} \Rightarrow \text{w SM } VH\text{-production } \mu_c < 30 \Rightarrow \text{no constraint on } y_c.$



◆ However $\mu_c < 30$ for large $c_c > 50$ new production mechanism:



No runaway for c_c



$c_c < 250.$

Constraining Higgs-quark universality #1 (model indep')

◆ New production eliminates the y_c runaway $\Rightarrow C_c < 250$ what about y_t ?

◆ ATLAS+CMS tth : $\mu_{tth}^{\text{avg}} = 2.41 \pm 0.81 \Rightarrow c_t > 0.9 \sqrt{\frac{\text{Br}_{h \rightarrow \text{relevant modes}}^{\text{SM}}}{\text{Br}_{h \rightarrow \text{relevant modes}}}} > 0.9$

$$\frac{C_c}{C_t} = \frac{y_t^{\text{SM}}}{y_c^{\text{SM}}} \frac{y_c}{y_t} = 280 \frac{y_c}{y_t} < 250 \Rightarrow y_c < y_t!$$

GP, Soreq, Stamou & Tobioka (Feb/15)

As shown below: the method works much better via real c-tagging working point.

Constraining Higgs-quark universality #2+3

- ◆ Width bound: $\Gamma_h < 2.6 \text{ GeV (ATLAS)}, \Gamma_h < 1.7 \text{ GeV (CMS)} \Rightarrow \boxed{C_c < 150, 120 .}$

GP, Soreq, Stamou & Tobioka (Feb/15)

- ◆ Interpretation of ATLAS recent $h \rightarrow J/\psi\gamma$ (1501.03276): $\sigma(pp \rightarrow h) \times \text{BR}_{h \rightarrow J\psi\gamma} < 33 \text{ fb}$,

- ◆ As discussed below, this implies: $\Gamma_{h \rightarrow J/\psi\gamma} = 1.42 (\kappa_\gamma - 0.087\kappa_c)^2 \times 10^{-8} \text{ GeV}$

Bodwin, Petriello, Stoynev & Velasco (13); Bodwin, Chung, Ee, Lee & Petriello (14)

- ◆ Getting rid of production: $\mathcal{R}_{J/\psi,Z} = \frac{\sigma(pp \rightarrow h) \times \text{BR}_{h \rightarrow J/\psi\gamma}}{\sigma(pp \rightarrow h) \times \text{BR}_{h \rightarrow ZZ^* \rightarrow 4\ell}} = \frac{\Gamma_{h \rightarrow J/\psi\gamma}}{\Gamma_{h \rightarrow ZZ^* \rightarrow 4\ell}} = 2.79 \frac{(\kappa_\gamma - 0.087\kappa_c)^2}{\kappa_V^2} \times 10^{-2}$,

$$\mathcal{R}_{J/\psi,Z} = \frac{33 \text{ fb}}{\mu_{ZZ^*} \sigma^{\text{SM}} \text{BR}_{h \rightarrow ZZ^* \rightarrow 4\ell}^{\text{SM}}} < 9.32$$



$$\boxed{c_c < 210c_V + 11c_\gamma .}$$

(LEP: $c_V = 1.08 \pm 0.07$)

GP, Soreq, Stamou & Tobioka (Feb/15)

Finally global analysis

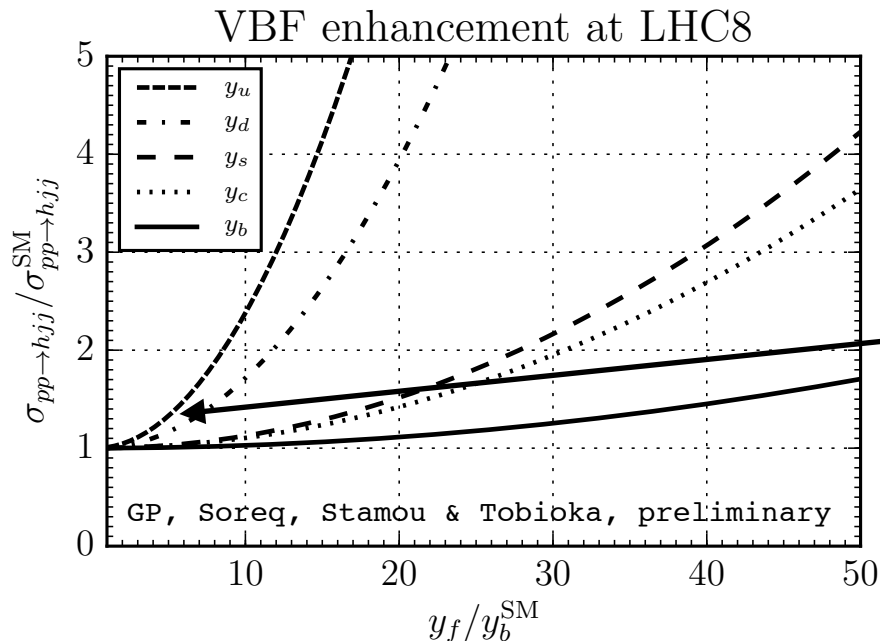
- ◆ The conventional way of doing the fit leads to: $\mathcal{C}_C < 6$.
- ◆ It is equivalence to the invisible (untagged) Higgs decay bound, driven by VBF:

$$\mu_{\text{VBF} \rightarrow h \rightarrow WW^*} = 1.27^{+0.44+0.30}_{-0.40-0.21} = 1.27^{+0.53}_{-0.45},$$

\Leftrightarrow

$$\mu_{\text{VBF} \rightarrow h \rightarrow WW^*} = \left(\kappa_V^2 + \bar{\sigma}_{\text{VBF}}^{\text{non-SM}} \right) \frac{\kappa_V^2}{R_\Gamma}$$

where $R_\Gamma = \Gamma_h / \Gamma_h^{\text{SM}}$

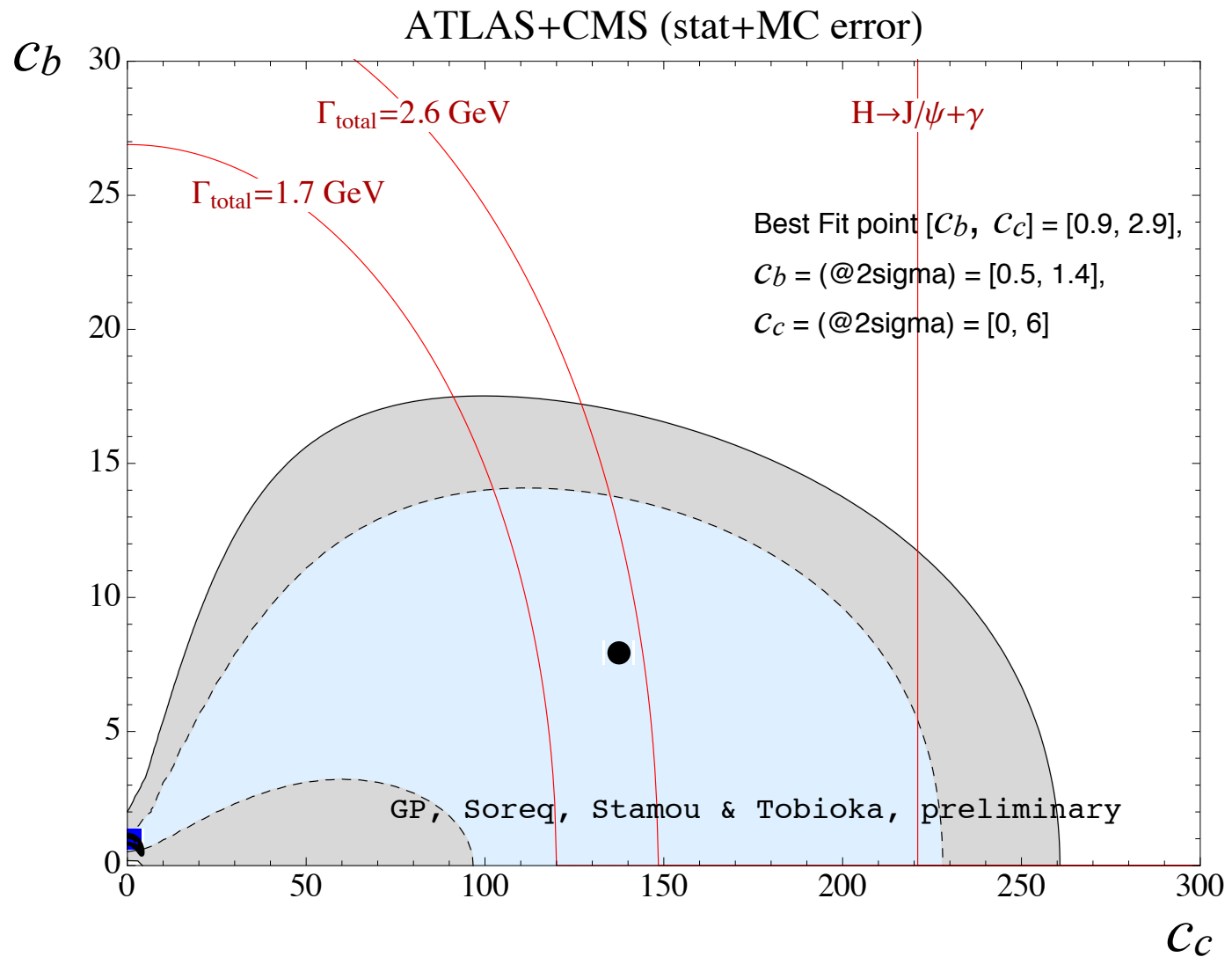


always set to zero, however not necessarily negligible.

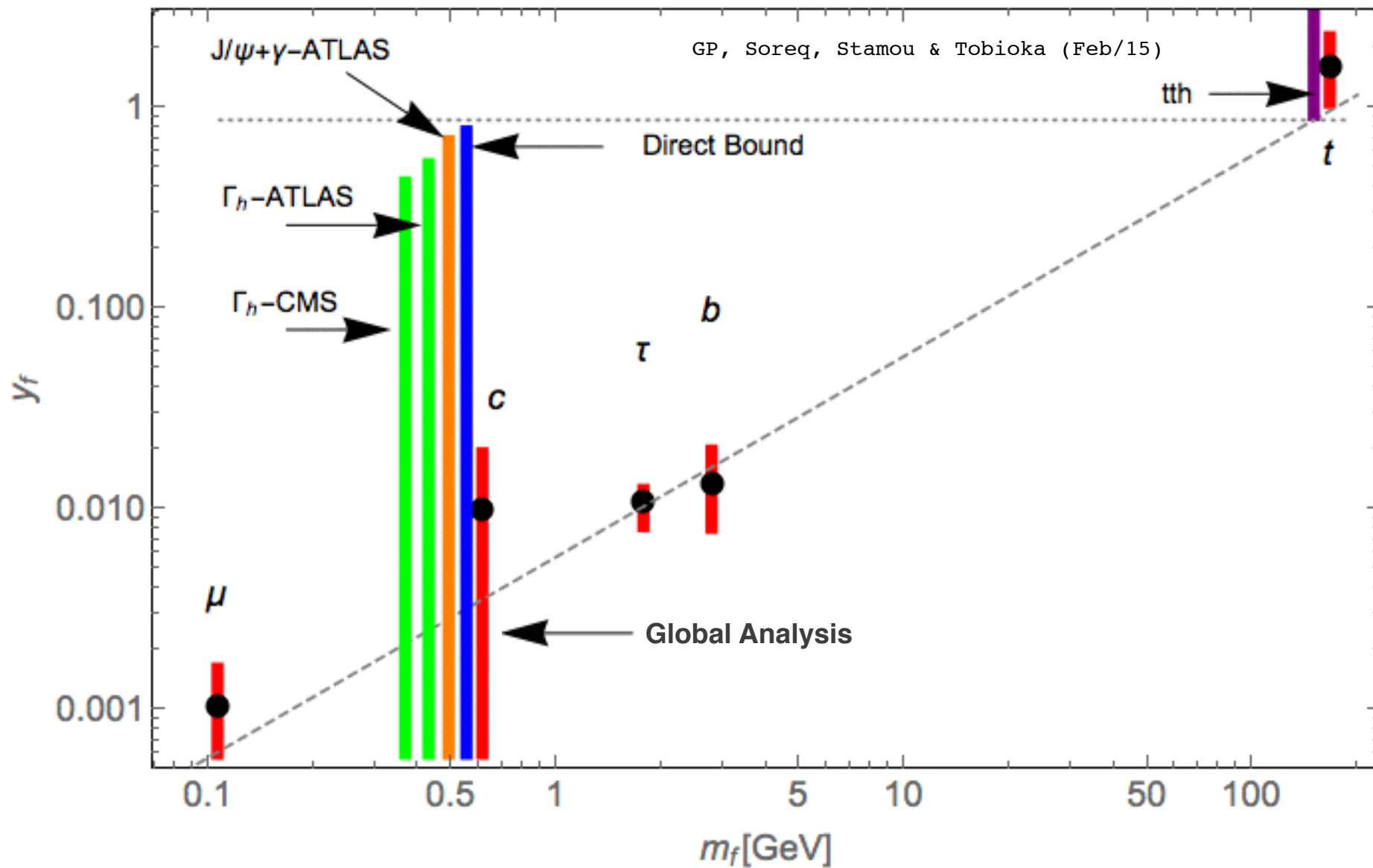
Currently small effect but might not be in the future.

VBF & Vh can be compared to other machines, leptons? hadrons?

Showing all constraints together



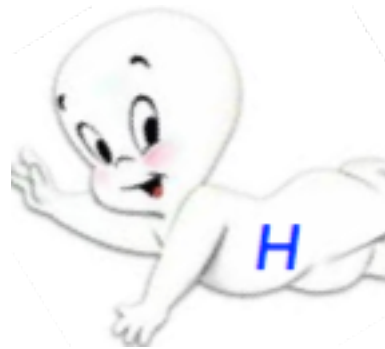
Preliminary summary/money plot ...



An Exclusive Window onto Higgs Yukawa Couplings to light quarks

Bodwin, Petriello, Stoynev & Velasco (13)
Kagan, GP, Petriello, Soreq, Stoynev & Zupan (14)

$s\bar{b}$ $s\bar{s}$
 $d\bar{b}$ $d\bar{d}$
 $d\bar{s}$ $u\bar{u}$



Exclusive path towards Higgs-light quark couplings

◆ Use the eff. Lagrangian:
$$\mathcal{L}_{\text{eff}} = - \sum_{q=u,d,s} \bar{\kappa}_q \frac{m_b}{v} h \bar{q}_L q_R - \sum_{q \neq q'} \bar{\kappa}_{qq'} \frac{m_b}{v} h \bar{q}_L q'_R + h.c.$$
$$+ \kappa_Z m_Z^2 \frac{h}{v} Z_\mu Z^\mu + 2\kappa_W m_W^2 \frac{h}{v} W_\mu W^\mu + \kappa_\gamma A_\gamma \frac{\alpha h}{\pi v} F^{\mu\nu} F_{\mu\nu},$$

Notice that:
$$\bar{\kappa}_q = y_q / y_b^{\text{SM}}, \quad (\text{sorry different notation})$$

in the SM:

$$\bar{\kappa}_s = m_s / m_b \simeq 0.020$$

$$\bar{\kappa}_d = m_d / m_b \simeq 1.0 \cdot 10^{-3}$$

$$\bar{\kappa}_u = m_u / m_b \simeq 4.7 \cdot 10^{-4}$$

$$\kappa_\gamma = \kappa_V = 1$$

Exclusive path towards Higgs-light quark couplings

Kagan, GP, Petriello, Soreq, Stoynev & Zupan (14)

◆ Use the eff. Lagrangian:

$$\mathcal{L}_{\text{eff}} = - \sum_{q=u,d,s} \bar{\kappa}_q \frac{m_b}{v} h \bar{q}_L q_R - \sum_{q \neq q'} \bar{\kappa}_{qq'} \frac{m_b}{v} h \bar{q}_L q'_R + h.c.$$

$$+ \kappa_Z m_Z^2 \frac{h}{v} Z_\mu Z^\mu + 2\kappa_W m_W^2 \frac{h}{v} W_\mu W^\mu + \kappa_\gamma A_\gamma \frac{\alpha h}{\pi v} F^{\mu\nu} F_{\mu\nu},$$

Notice that: $\bar{\kappa}_q = y_q / y_b^{\text{SM}},$

where generically: $|\bar{\kappa}_u| < 0.98, \quad |\bar{\kappa}_d| < 0.93, \quad |\bar{\kappa}_s| < 0.70$

varying only one at the time (95%CL)

$$|\bar{\kappa}_u| < 1.3, \quad |\bar{\kappa}_d| < 1.4, \quad |\bar{\kappa}_s| < 1.4$$

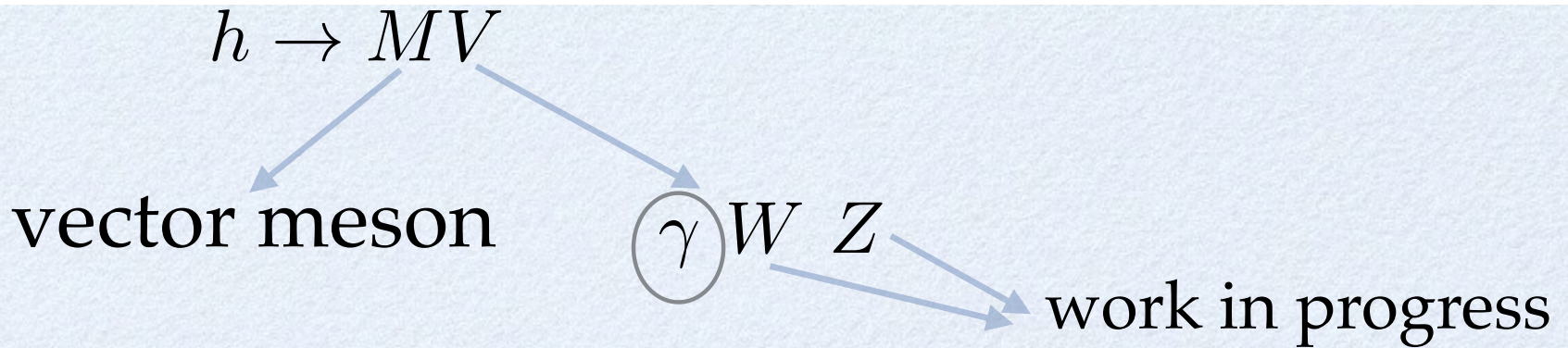
varying all couplings (95%CL)

$$|\bar{\kappa}_{qq'}| < 0.6 (1) \quad \text{for } q, q' \in u, d, s, c, b \text{ and } q \neq q'$$

same for the flavor violating case

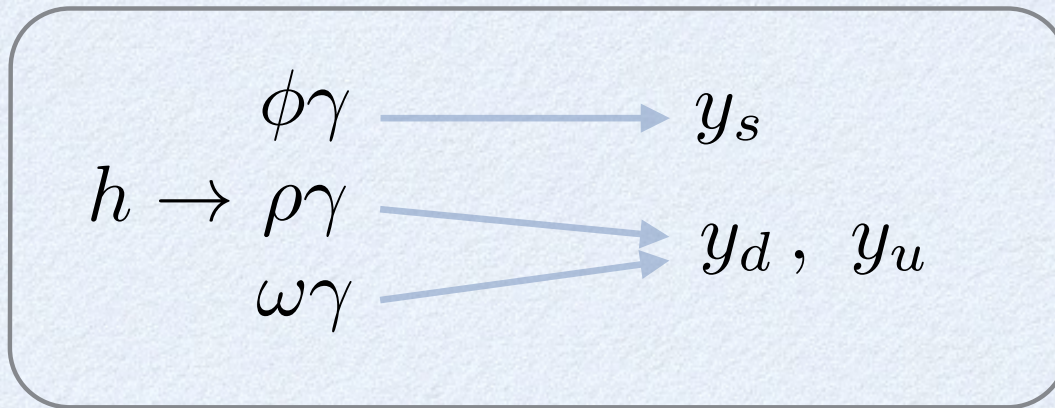
(FCNC non-robust bound: $|\bar{\kappa}_{bs}| < 8 \cdot 10^{-2}$ Harnik, Kopp & Zupan; Blankenburg, Ellis, Isidori, (12))

The main idea



Bodwin, Petriello,
 Stoynev, Velasco
 1306.5770

$$h \rightarrow J/\psi \gamma \longrightarrow \gamma_c$$



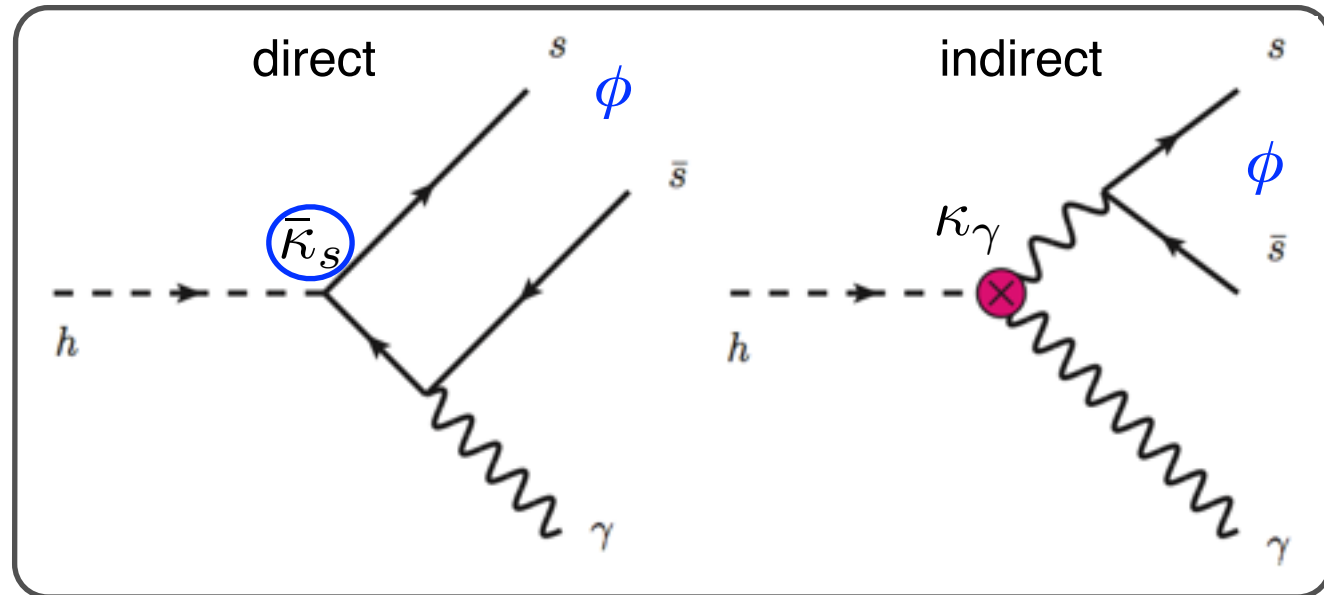
Kagan, GP, Petriello, Soreq, Stoynev & Zupan (14)

Adding off-diagonal: $h \rightarrow \bar{B}^{0*} \gamma, h \rightarrow \bar{B}^{0*} \gamma, h \rightarrow K^{0*} \gamma, h \rightarrow D^{0*} \gamma$

Ex.: $h \rightarrow \phi\gamma$

$$\Gamma_{h \rightarrow \phi\gamma} = \frac{1}{8\pi} \frac{1}{m_h} |M_{ss}^\phi|^2,$$

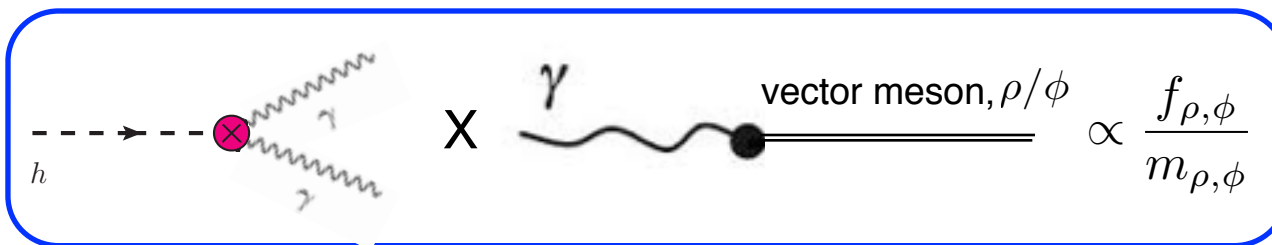
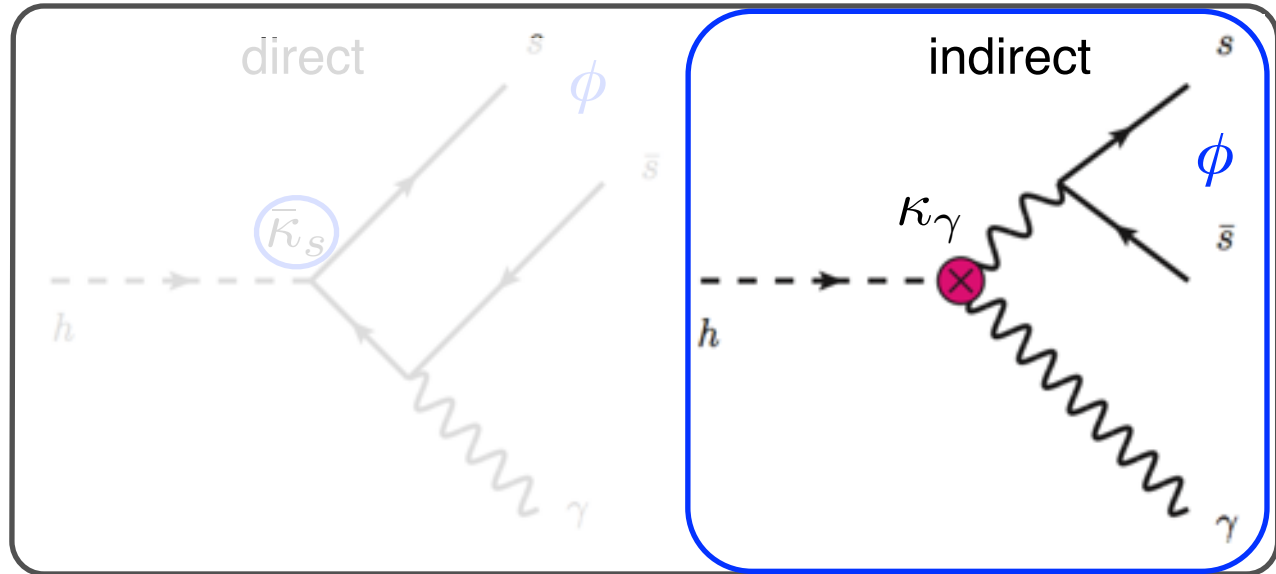
- ◆ Two paths to get $h \rightarrow \phi\gamma$:



- ◆ Let us understand them one by one.

Ex.: $h \rightarrow \phi\gamma$, indirect contribution

- Two paths to get $h \rightarrow \phi\gamma$:

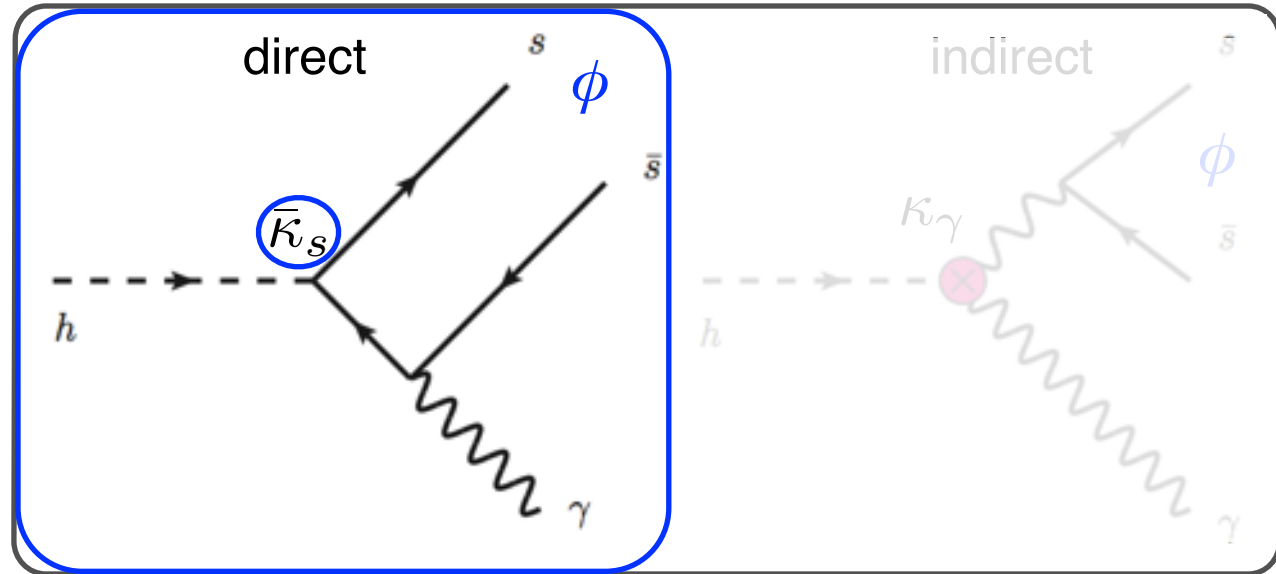


from experiment, $\phi \rightarrow e^+e^-$
 $(f_\phi = 0.235(5) \text{ GeV})$

$$(M_{ss}^\phi)_{\text{indir}} \approx \frac{f_{\rho,\phi}}{m_{\rho,\phi}} \kappa_\gamma A_\gamma \frac{4\alpha m_h^2}{\pi v}$$

Ex.: $h \rightarrow \phi\gamma$, direct contribution

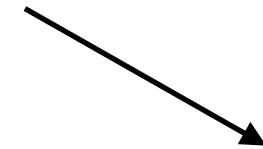
- Two paths to get $h \rightarrow \phi\gamma$:



$$\propto \frac{\bar{\kappa}_s m_b}{v} \frac{f_{\rho,\phi}^\perp}{m_h}$$

("local" structure : $\bar{s}\sigma_{\mu\nu}s \times F^{\mu\nu}$)

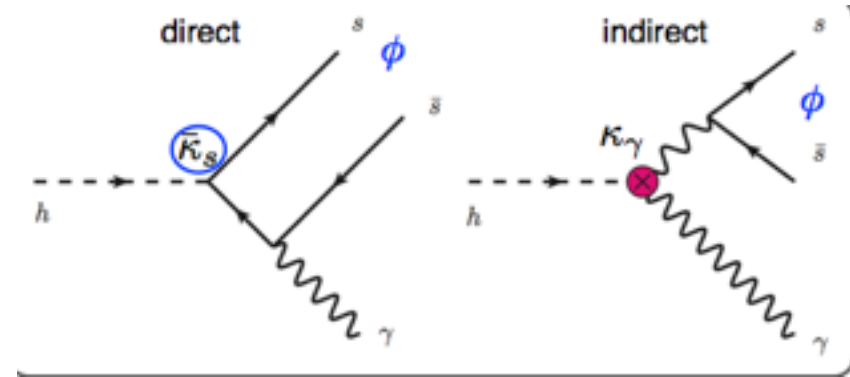
from experiment, $\phi \rightarrow e^+e^-$
 $(f_\perp^\phi = 0.191(28))$



$$(M_{ss}^\phi)_{\text{dir}} \approx \frac{\bar{\kappa}_s m_b}{v} f_\phi^\perp$$

Final result for the $\text{BR}(h \rightarrow \phi\gamma)$

$$\Gamma_{h \rightarrow \phi\gamma} = \frac{1}{8\pi} \frac{1}{m_h} |M_{ss}^\phi|^2,$$



◆ The resulting sensitivity:

$$\frac{\text{BR}_{h \rightarrow \phi\gamma}}{\text{BR}_{h \rightarrow b\bar{b}}} = \frac{\kappa_\gamma \left[(3.0 \pm 0.13)\kappa_\gamma - 0.78\bar{\kappa}_s \right] \cdot 10^{-6}}{0.57\bar{\kappa}_b^2},$$

Similar holds
for 1st generation:

$$\frac{\text{BR}_{h \rightarrow \rho\gamma}}{\text{BR}_{h \rightarrow b\bar{b}}} = \frac{\kappa_\gamma \left[(1.9 \pm 0.15)\kappa_\gamma - 0.24\bar{\kappa}_u - 0.12\bar{\kappa}_d \right] \cdot 10^{-5}}{0.57\bar{\kappa}_b^2},$$

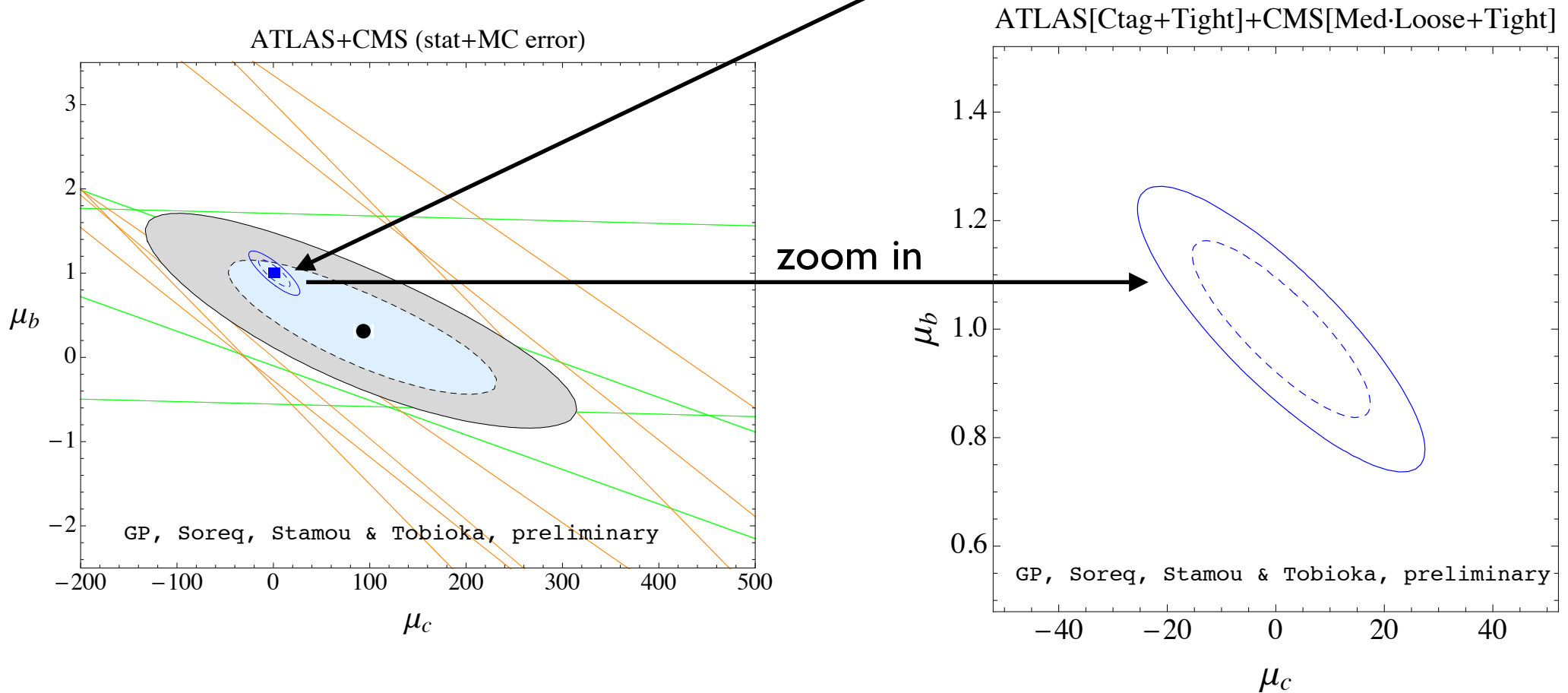
$$\frac{\text{BR}_{h \rightarrow \omega\gamma}}{\text{BR}_{h \rightarrow b\bar{b}}} = \frac{\kappa_\gamma \left[(1.6 \pm 0.17)\kappa_\gamma - 0.59\bar{\kappa}_u - 0.29\bar{\kappa}_d \right] \cdot 10^{-6}}{0.57\bar{\kappa}_b^2},$$

Few projections

HL projection, inclusive c-tagging (sorry haven't finish the FCC ones)

Combining medium & charm-tagging points: $\epsilon_{b,c} = (0.7, 0.2), (0.12, 0.19)$

Not including new IBL @ ATLAS => poten' improve things!
(yesterday's dinner-discussion with Anadi Canepa, correspondence w Tim Scanlon)



Of course $O(5\%)$ for y_c for lepton colliders (ex. of complementarity)

Exclusive modes, projections

Kagan, GP, Petriello, Soreq, Stoynev & Zupan (14)

- focus on $h \rightarrow \phi\gamma$, use **Pythia 8.1**
 - main decay modes: $\phi \rightarrow K^+K^-$ (49%), K_LK_S (34%), $\pi^+\pi^-\pi^0$ (15%)
 - for $pp \rightarrow h \rightarrow \phi\gamma$ at 14TeV LHC in 70 to 75% cases the kaons/pions and the prompt photon have $|\eta| < 2.4$
 - within the minimal fiducial volume of the ATLAS and CMS experiments
 - adopt the geometrical acceptance factor $A_g = 0.75$
 - do not include other efficiency or trigger factors
- assume $\kappa_\gamma = 1$, negligible background, 3σ reach

two detectors

one detector

\sqrt{s} [TeV]	$\int \mathcal{L} dt$ [fb $^{-1}$]	# of events (SM)	$\bar{\kappa}_s > (<)$	$\bar{\kappa}_s^{\text{stat.}} > (<)$
14	3000	770	0.56 (−1.2)	0.27 (−0.81)
33	3000	1380	0.54 (−1.2)	0.22 (−0.75)
100	3000	5920	0.54 (−1.2)	0.13 (−0.63)

no theory error

5x SM strange Yukawa

Future experiments

- only a few events expected at e^+e^- colliders
 - ILC, ILC with luminosity upgrade, CLIC
 - probably too small for observation of $h \rightarrow \phi\gamma$
- ≈ 30 events expected at FCC-ee (TLEP)
 - too small to probe a deviation from the SM prediction
- $h \rightarrow \phi\gamma$ measurements unique to future hadron machines

Conclusions

- ◆ Is the Higgs-mechanism behind the light quark masses?.
- ◆ Order one modifications to Higgs light quark (charm) coupling lead to dramatic change in Higgs pheno'.
- ◆ Charm coupling is constrained via charm-tagging, or exclusively.
- ◆ The light quarks can be potentially probed via exclusive decays.
- ◆ Looked at $h \rightarrow M\gamma$, with $h \rightarrow \phi\gamma$ most promising.
- ◆ Established higgs-quarks non-universality.

Backups

Experimental sensitivity

- focus on $h \rightarrow \phi\gamma$, use $\mathcal{P} = 1 - \mathcal{P}_{\text{SM}}$
 - main channels: $h \rightarrow \pi^+\pi^-\pi^0$ (15%), $h \rightarrow \text{kaons/pions}$
 - $\mathcal{S} = |\text{BR}_{h \rightarrow \phi\gamma} - \text{BR}_{h \rightarrow \phi\gamma}^{\text{SM}}| / (\delta\text{BR}_{h \rightarrow \phi\gamma})$, where $(\delta\text{BR}_{h \rightarrow \phi\gamma})^2 = \text{BR}_{h \rightarrow \phi\gamma} / (\sigma_h \mathcal{L} A_g) + (\delta\text{BR}_{h \rightarrow \phi\gamma}^{\text{th}})^2$ (ATLAS and CMS)
 - adopt the geometric acceptance $A_g = 0.75$
 - do not include other efficiency or trigger factors
 - assume $\kappa_\gamma = 1$, negligible background, 3σ reach

two detectors
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Thoughts about experimental strategy

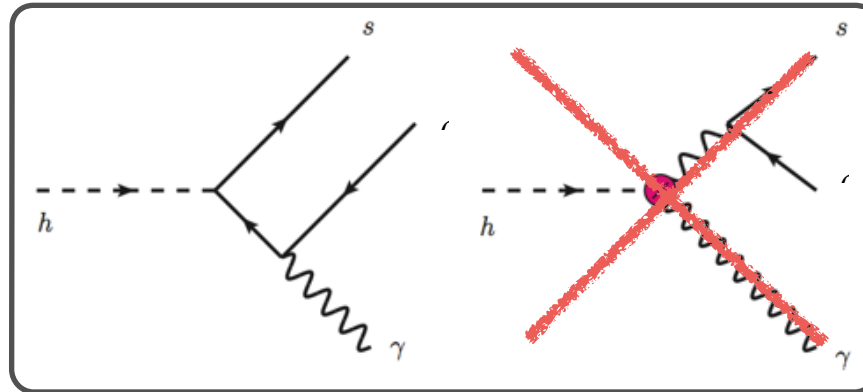
- for $h \rightarrow \phi\gamma$ decay most promising $\phi \rightarrow K^+K^-$
 - near collinearity of the photon and the ϕ -jet in the transverse plane
 - jet sub-structure information
 - two close high- p_T tracks in a narrow cone
 - di-track invariant mass distribution assuming kaons
 - 1.5% (better than 15 MeV) resolution (CMS)
- can probably be used to significantly cut on the background
 - on jet+ γ QCD backgrounds
 - on $h \rightarrow \phi\gamma+n\pi^0$, $\eta^{(\prime)}$ (\rightarrow neutr.) γ
- dedicated trigger probably required to enhance the reach

Thoughts about experimental strategy

- $h \rightarrow \rho^0 \gamma$ mode
 - $Br(\rho^0 \rightarrow \pi^+ \pi^-) \sim 100\%$
 - relatively clean mode, similar to $\phi \rightarrow K^+ K^-$ decay
- $h \rightarrow \omega \gamma$ mode
 - $Br(\omega \rightarrow \pi^+ \pi^- \pi^0) \sim 89\%$
 - harder to trigger on
 - hard-to-identify π^0 smears the observable quantities
 - a detailed experimental study required

Flavor violating couplings

Kagan, GP, Petriello, Soreq, Stoynev & Zupan (14)



- FV modes $h \rightarrow \bar{B}_s^{0*} \gamma, h \rightarrow \bar{B}^{0*} \gamma, h \rightarrow \bar{K}^{0*} \gamma, h \rightarrow D^{0*} \gamma$
 - can probe $\bar{\kappa}_{bs, sb}, \bar{\kappa}_{bd, db}, \bar{\kappa}_{sd, ds}$ and $\bar{\kappa}_{cu, uc}$
- $h \rightarrow \bar{K}^{0*} \gamma$ similar expr. as $h \rightarrow \phi \gamma$
 - but only direct amplitude
- for $\bar{\kappa}_{ds} \sim O(1) \Rightarrow Br(h \rightarrow \bar{K}^{0*} \gamma) \sim O(10^{-8})$
 - not observable at planned future colliders

$$\frac{BR_{h \rightarrow \bar{B}_s^{0*} \gamma}}{BR_{h \rightarrow b\bar{b}}} = \frac{(2.1 \pm 1.0) \cdot 10^{-7} |\bar{\kappa}_{bs}|^2 + |\bar{\kappa}_{sb}|^2}{0.57 \bar{\kappa}_b^2} \cdot \frac{1}{2},$$

The production via $q\bar{q} \rightarrow h$ in pp with 8TeV are (the SM is 19 pb)

$$\sigma_{u\bar{u} \rightarrow h} = \left(\frac{y_u}{y_b^{\text{SM}}} \right)^2 9.16 \text{ pb},$$

$$\sigma_{d\bar{d} \rightarrow h} = \left(\frac{y_d}{y_b^{\text{SM}}} \right)^2 6.29 \text{ pb},$$

$$\sigma_{s\bar{s} \rightarrow h} = \left(\frac{y_s}{y_b^{\text{SM}}} \right)^2 1.67 \text{ pb},$$

$$\sigma_{c\bar{c} \rightarrow h} = \left(\frac{y_c}{y_b^{\text{SM}}} \right)^2 0.83 \text{ pb}.$$

