# Physics At VLHC Tao Han

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Exploring the Physics Frontier with Circular Colliders ACP Winter Conference, Jan. 27, 2015



Photo credit: Hitoshi Murayama FERMILAB-CONF-97/318-T

#### Summary of the Very Large Hadron Collider Physics and Detector Workshop

Physics at the high energy frontier beyond the LHC March 13-15, 1997

### (Bill Foster's initiative)

G. Anderson (Fermilab), U. Baur (SUNY at Buffalo), M. Berger (Indiana University), F. Borcherding (Fermilab), A. Brandt (Fermilab), D. Denisov (Fermilab, Co-Chair and Co-editor), S. Eno (University of Maryland), T. Han (University of California–Davis), S. Keller (Fermilab, Co-Chair and Co-editor), D. Khazins (Duke University), T. LeCompte (Argonne National Laboratory), J. Lykken (Fermilab), F. Olness (Southern Methodist University), F. Paige (Brookhaven National Laboratory), R. Scalise (Southern Methodist University), E. H. Simmons (Boston University), G. Snow (University of Nebraska–Lincoln), C. Taylor (Case Western Reserve University), J. Womersley (Fermilab).

#### PHYSICS AT 100-200 TeV

Tao Han, Univ. of Wisconsin-Madison

(1999 VLHC Annual Meeting, June. 28)

#### I. Brief Introduction:

Particle Physics and Colliders

#### II. Physics Expectations at the VLHC:

- Representative SM Physics
- Physics Beyond the SM

III. Physics at the High-Energy Frontier

Beyond the Naive Expectation

#### PHYSICS AT THE VLHC

Tao Han, Univ. of Wisconsin-Madison

#### (July 17, Snowmass 2001)

VLHC: The True Energy Frontier

 Invitation to Innovative Ideas for New Physics

Beyond the SM Physics

• New Threshold and Extended Reach

Theory Overview

in the light of future hadron colliders

Tao Han Univ. of Wisconsin - Madison

#### VLHC workshop, Fermilab, Oct. 16, 2003

The Standard Model as It Is The Need For Going Beyond SM The Role of Future Hadron Colliders

### Physics Issues in 1999 (TH's list)

- $M_W$ ,  $M_Z$ ? (Gauge symmetry breaking)
- $M_H \sim \mathcal{O}(M_Z)$ ? (natural EW scale)
- ? Supersymmetry?  $(M_Z M_{pl} \text{ hierarchy})$
- ?  $m_t$ ,  $m_f$ ,  $m_{\nu}$ ? (fermion masses and mixing)
- ?★ Techni-/top-color? (dynamical symm. brkng)
  - extra dimensions/low-scale gravity?
  - ? (gravity+hierarchy)
    - Superstring?
  - ? (quantum gravity/Theory of everything?)

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? • ... ...? (DM)

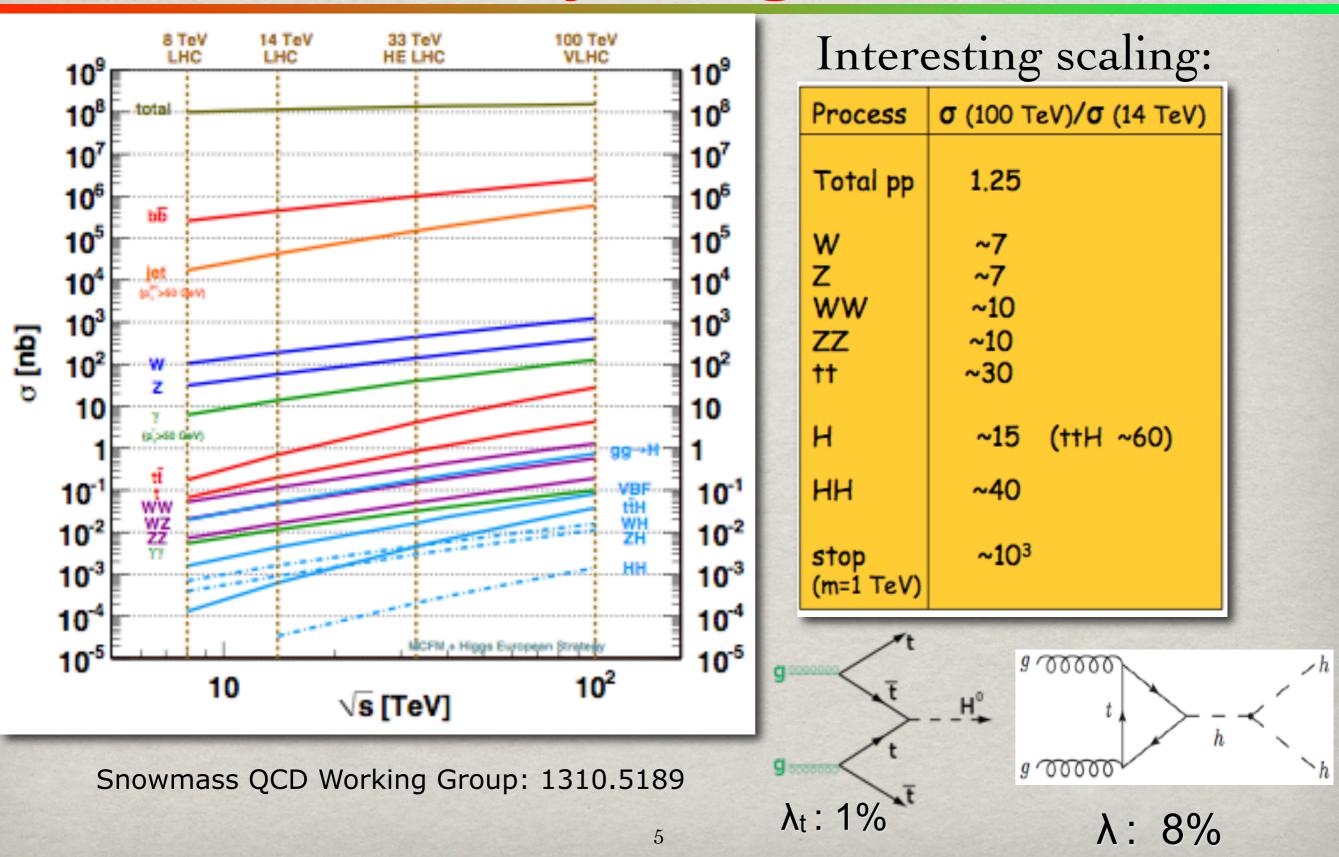
#### Issues for the Future (Starting now!)

#### Chris Quigg, IAS-HKUST Jan. 19, 2015

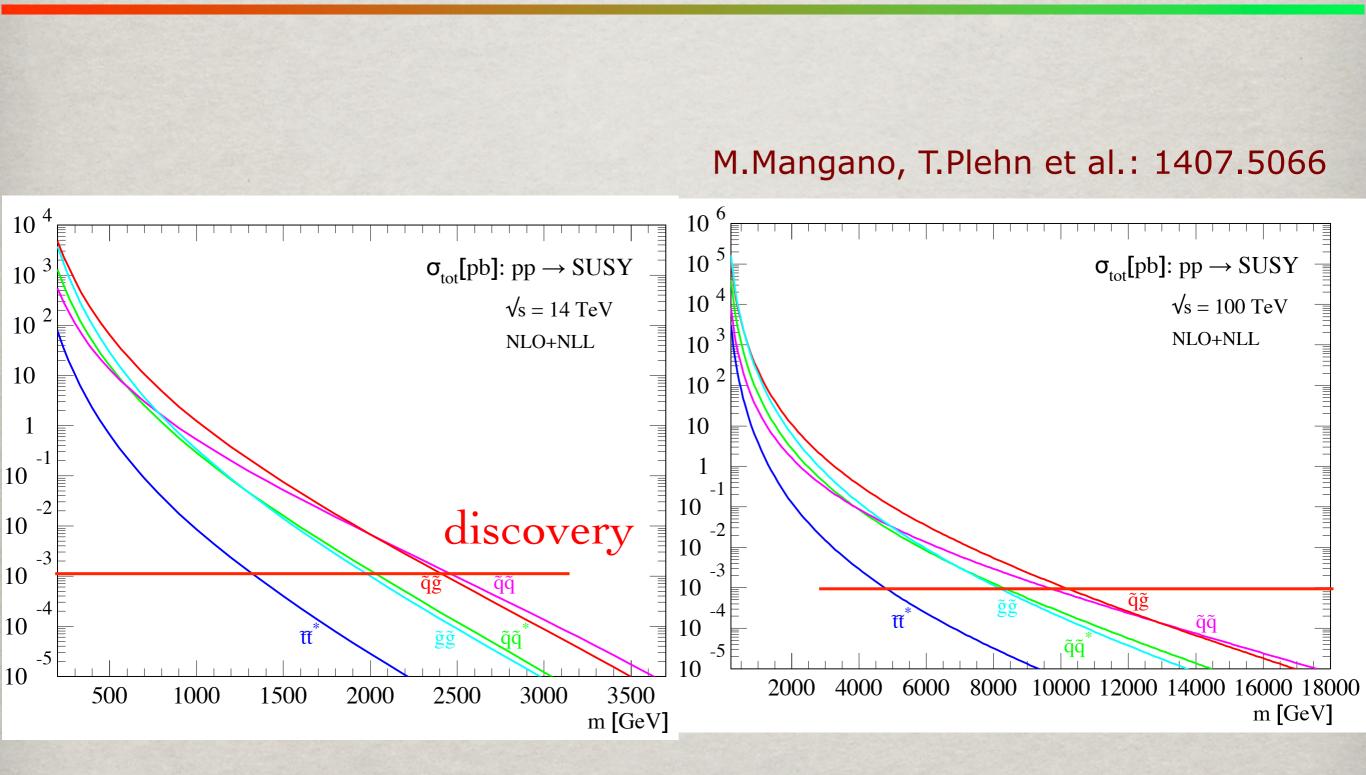
I. What is the agent of EWSB? There is a Higgs	6. Do the different CC behaviors of LH, RH fermions
Might there be several?	reflect a fundamental asymmetry in nature's laws?
2. Is the Higgs boson elementary or composite	7. What will be the next symmetry we recognize? Are
does it interact with itself? What triggers EWS	there additional heavy gauge bosons? Is nature
3. Does the Higgs boson give mass to fermions	supersymmetric? Is EW theory contained in a GUT?
only to the weak bosons? What sets the masses	8 Are all flavor-changing interactions governed by the
mixings of the quarks and leptons? (How) is ferr	standard-model Yukawa couplings? Does "minimal
mass related to the electroweak scale?	flavor violation" hold? If so why?
4. Are there new flavor symmetries that give in	9. Are there additional sequential quark & lepton
into remnion masses and mixings:	generations? Or new exotic (vector-like) fermions?
5. What stabilizes the Higgs-boson mass below	10. What resolves the strong CP problem?
$11 \sqrt{bet even the device metters? Any flavour etc.}$	16 What explains the baryon asymmetry of the

11. What are the dark matters? Any flavor str 16. What explains the baryon asymmetry of the
12. Is EWSB an emergent phenomenon connuniverse? Are there new (CC) CP-violating phases? with strong dynamics? How would that alter conception of unified theories of the strong, would observation, or more stringent limits, on and electromagnetic interactions?
13. Is EWSB related to gravity through extra spacetime dimensions?
14. What resolves the vacuum energy proble 19. At what scale are the neutrino masses set? Do
15. (When we understand the origin of EWS they speak to the TeV scale, unification scale, Planck lessons does EWSB hold for unified theories scale, ...?
20. How are we prisoners of conventional thinking?

## VLHC LEADS ENERGY FRONTIER Rich Physics @ VLHC

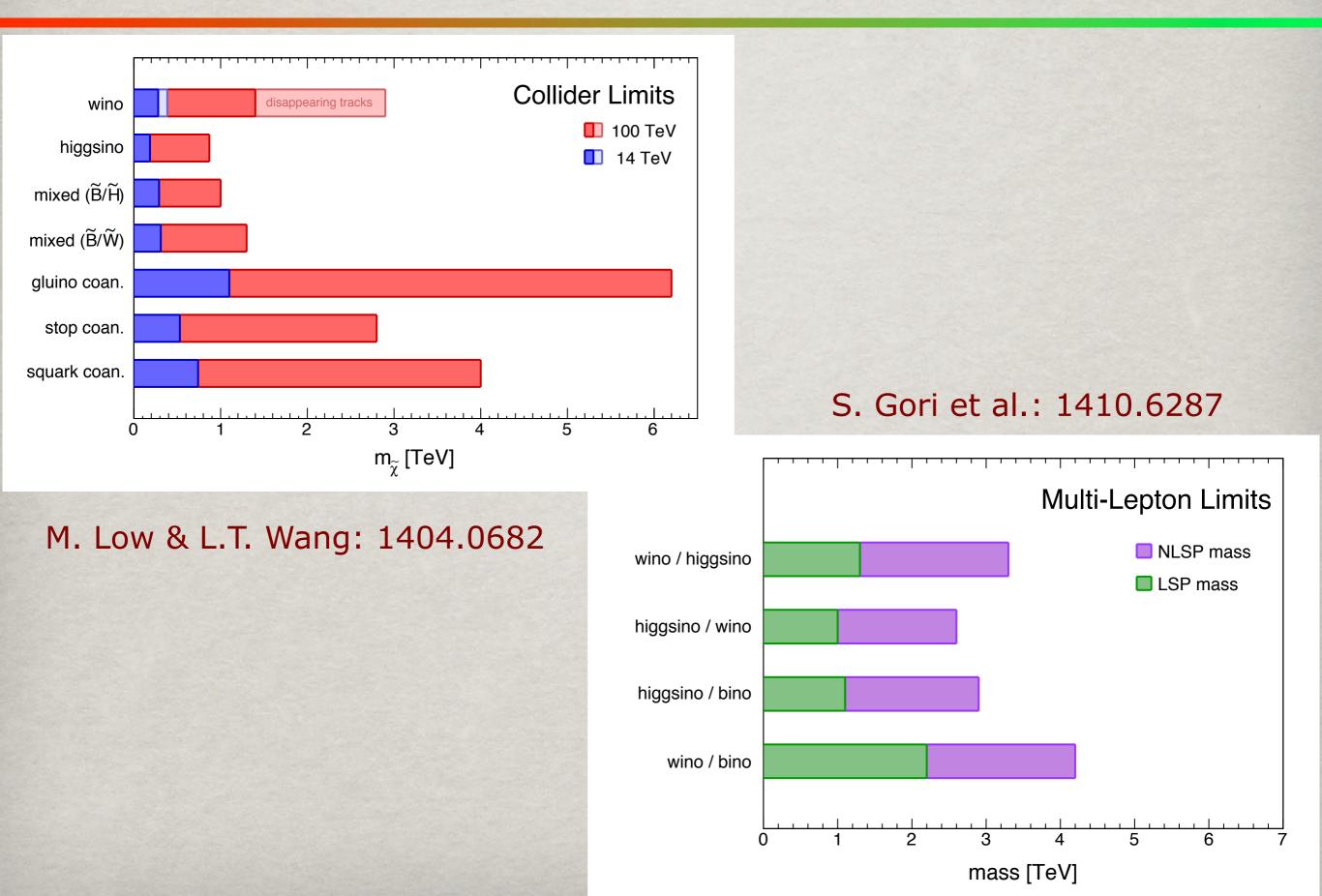


## SUSY @ VLHC



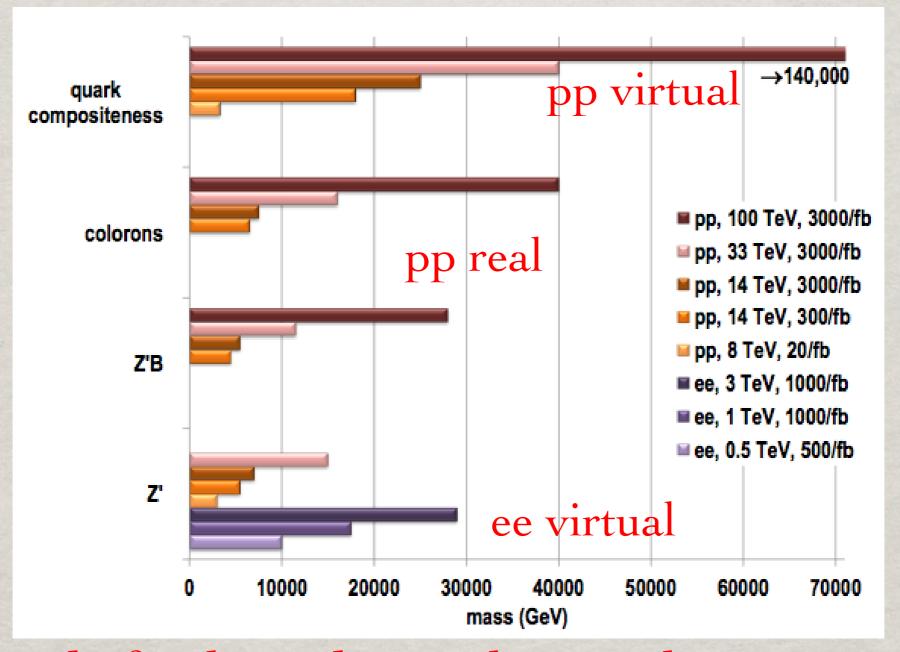
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## SUSY DM @ VLHC



## **New Particle Searches at VLHC**

#### Snowmass NP report, 1311.0299



VLHC leads for broad searches at the energy frontier. *Look forward to the many inspiring talks this week!* <sup>8</sup>

## **NEW BEHAVIOR OF "OLD PHYSICS"** At the new energy frontier VLHC: $v/\sqrt{s} \sim 2.5 \times 10^{-3}$

The EW gauge bosons & the top quark pretty much "massless": the EW symmetry is "restored" In the collision processes

- Final state particles > narrow jets, radiations
- New physics @ heavier scales  $\rightarrow W^{\pm}/Z/H/top$
- → Studying  $W^{\pm}/Z/H/top$  at higher energies:
- bread & butter (new) phenomena within SM
- first step toward understanding O(10 TeV) scale physics



**TOP QUARK INITIATED PROCESSES** TH, J. Sayre, S. Westhoff: 1411.2588 With  $m_t << E_{cm}$ , The top quark *IS* as massless at the VLHC as *b*-quark at the Tevatron:  $m_h / E_{TeV} \sim 3.5 / 1 \times 10^3 \sim 3.5 \times 10^{-3}$  $m_{t}/E_{VLHC} \sim 160/50 \mathrm{x} 10^{3} \sim 3.2 \mathrm{x} 10^{-3}$ When a heavy scale M is involved, so that  $\alpha_{s} \ln(M^{2}/m_{t}^{2}) \sim O(1) \rightarrow M \sim (50-100) m_{t}$ then the collinear large logs need to be resummed  $\rightarrow$  top quarks as partons 8 6000000 - H · g receeded 10

## TOP QUARK INITIATED PROCESSES With $m_t \sim v_t$ ,

The top quark may hold the key to new physics:

- Most sensitive to the "naturalness" issue.
- Vacuum stability

• .... TH, J. Sayre, S. Westhoff: 1411.2588 Examples for new physics:  $t\bar{t} \rightarrow X$ 

spin 0: neutral scalar  $H^0$ :  $i\frac{y}{\sqrt{2}}$ ; pseudo scalar  $A^0$ :  $i\frac{y}{\sqrt{2}}\gamma_5$ ; charged scalar  $H^+$ :  $i\frac{y}{\sqrt{2}}(g_L P_L + g_R P_R)$ ;

spin 1: color – singlet vector/axial vector  $Z^{\prime 0}$ ,  $W^{\prime +}$ :  $ig\gamma^{\mu}(g_V - g_A\gamma_5)$ ; color – octet vector/axial vector  $g_{KK}$ :  $ig_s\gamma^{\mu}(g_V - g_A\gamma_5) t^a$ ; spin 2: tensor G:  $-i\frac{\kappa}{8}[\gamma^{\mu}(p_t - p_{\bar{t}})^{\nu} + \gamma^{\nu}(p_t - p_{\bar{t}})^{\mu} - 2g^{\mu\nu}(\not{p}_t - \not{p}_{\bar{t}} - 2m_t)].$ 

### New Physics Examples: $tt \to X$

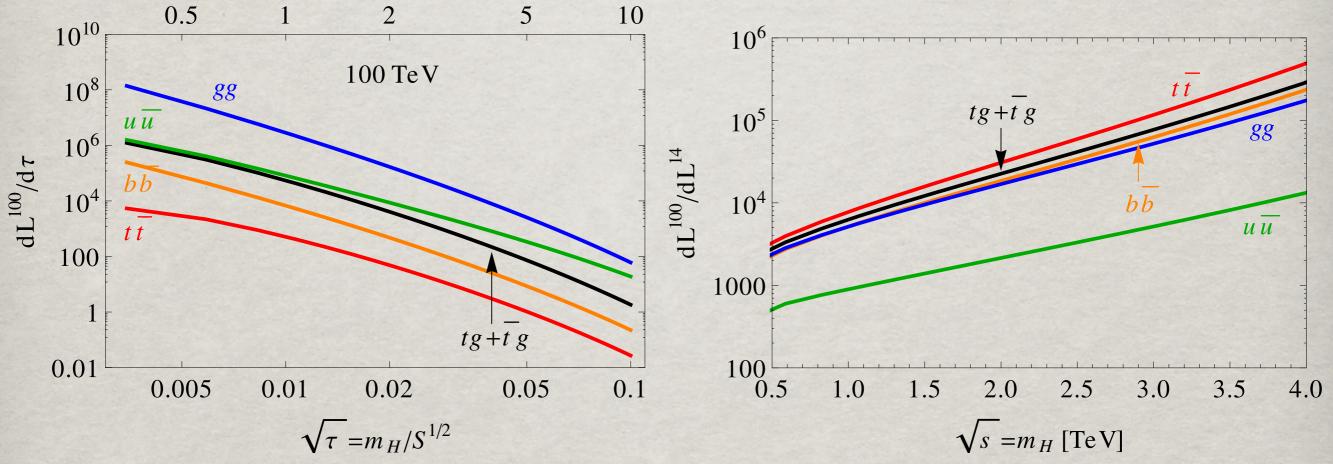
#### TH, J. Sayre, S. Westhoff: 1411.2588

Table 1: Spin- and color-averaged squared matrix elements for the production of an on-shell heavy particle of mass  $m_H = \sqrt{s}$  from heavy-quark fusion and the corresponding threshold behavior. The number of colors is denoted by  $N_c$  and the SU(3) invariant as  $C_F = 4/3$ . Subscripts T and L indicate transverse and longitudinal polarization, respectively. The kinematic factors are  $\beta_{ij}^2 = \ell_{ij}(1 - (m_i + m_j)^2/s)$  and  $\ell_{ij} = 1 - (m_i - m_j)^2/s$ , as well as the couplings  $g_{S,P} = (g_L \pm g_R)/2$  in terms of chiral couplings  $g_L$  and  $g_R$ .

process	$\overline{\sum} \mathcal{M} ^2$	threshold behavior	
$t\bar{t} \to H^0$	$\frac{y^2s}{4N_c} \ \beta_{t\bar{t}}^2$	P-wave	
$t\bar{t} \to A^0$	$\frac{y^2s}{4N_c}$	S-wave	
$t\bar{b} \to H^+$	$\frac{y^2 s}{4N_c} \left( g_S^2 \beta_{t\bar{b}}^2 / \ell_{t\bar{b}} + g_P^2 \ell_{t\bar{b}} \right)$	same as $H^0$ , $A^0$ , with an extra $\ell$	
$t\bar{t} \to Z_T^{\prime 0}$	$\frac{g^2s}{N_{\rm g}} \left(g_V^2 + g_A^2\beta_{t\bar{t}}^2\right)$	vector: S-wave; axial-vector: P-wave	
$t\bar{t} \to Z_L^{\prime 0}$	$\frac{\frac{g^2s}{N_c}}{\frac{g^2s}{N_c}} g_V^2(2m_t^2/s)$	fermion mass suppression	
$t\bar{b} \to W_T^{\prime +}$	$\frac{g^2 s}{N_c} \left( g_V^2 \ell_{t\bar{b}} + g_A^2 \beta_{t\bar{b}}^2 / \ell_{t\bar{b}} \right)$	same as $Z_T^{\prime 0}$ , with an extra $\ell$	
$t\bar{b} \to W_L^{\prime +}$	$\frac{g^2 s}{N_c} \left( g_V^2 \ell_{t\bar{b}} \frac{(m_t + m_b)^2}{2s} + g_A^2 \beta_{t\bar{b}}^2 \frac{(m_t - m_b)^2}{2s\ell_{t\bar{b}}} \right)$	fermion mass suppression	
$t\bar{t} \to g_{KK}$	$C_F \frac{g_s^2 s}{N_c} \left( g_V^2 (1 + 2m_t^2/s) + g_A^2 \beta_{t\bar{t}}^2 \right)$	same as $Z'^0$	
$t\bar{t} \to G$	$\frac{\kappa^2 s^2}{32N_c} \ (1 + 8m_t^2/3s)\beta_{t\bar{t}}^2$	P-wave	

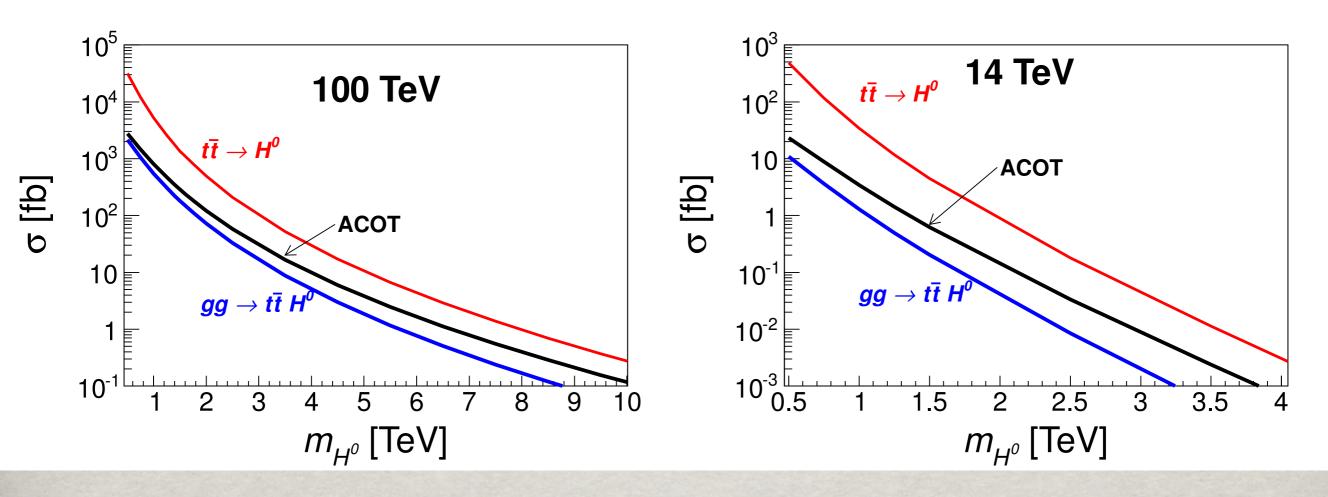
### Partonic luminosities

$$\sigma_{pp \to H+X}(S) = \sum_{i,j} \int_{m_H^2/S}^1 dx_1 \int_{m_H^2/(x_1S)}^1 dx_2 f_i(x_1,\mu) f_j(x_2,\mu) \hat{\sigma}_{ij \to H}(s)$$
  
$$\equiv \sum_{i,j} \int_{m_H^2/S}^1 d\tau \ \frac{dL_{ij}}{d\tau} \ \hat{\sigma}_{ij}(s), \qquad \frac{dL_{ij}}{d\tau}(\tau,\mu) = \int_{\tau}^1 \frac{dx}{x} f_i(x,\mu) f_j(\tau/x,\mu)$$
  
$$\sqrt{s} = m_H [\text{TeV}]$$

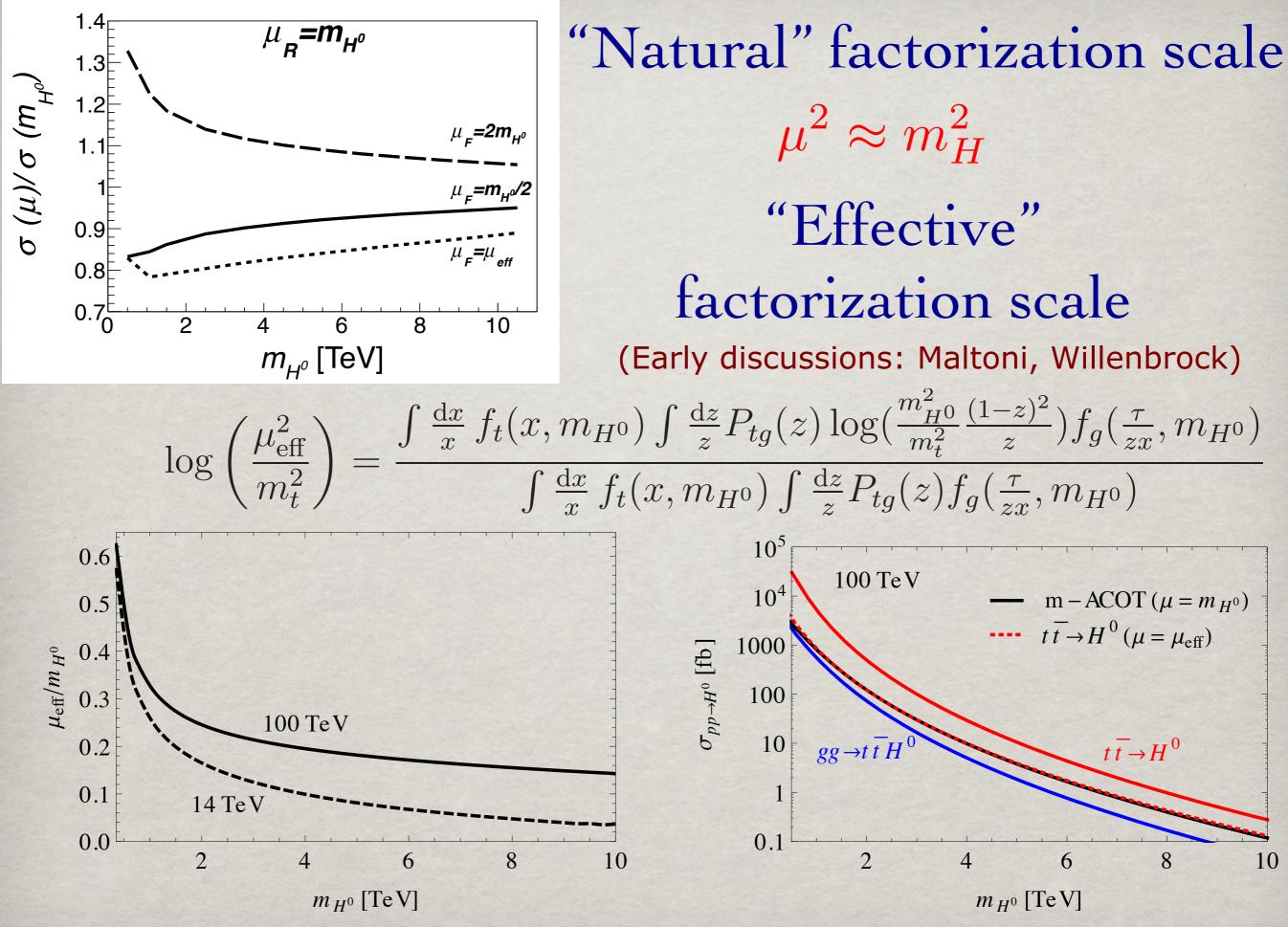


Top lumi tracking gg, reaching few% of bb! Relevant range:  $0.002 \leq \bar{x} \leq 0.1$ , for  $200 \text{ GeV} \leq \sqrt{s} \leq 10 \text{ TeV}$ . Lumi(gg, bb, tt @100/14) increased by  $1000 - 10^5$  for 500 GeV - 4 TeV! 13 TH, J. Sayre, S. Westhoff: 1411.2588

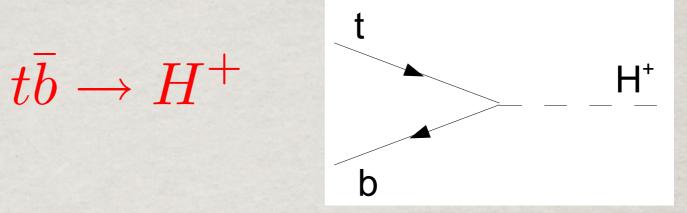
## 5-flavor vs. 6-flavor: (ACOT: massive top with careful subtraction)



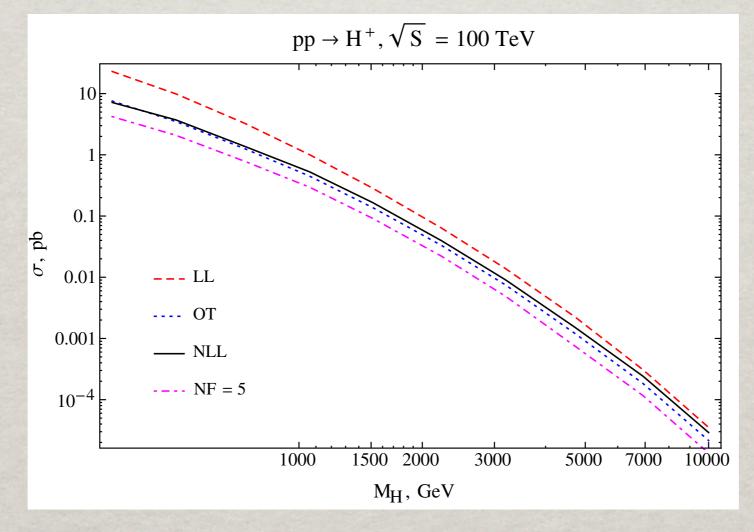
- 5-flavors usually underestimate the rate (better at low M)
- 6-flavors usually overestimate the rate (better at high M) (too much resummation) → proper treatment needed
- Higher CM Energies better approximation



### Another recent work \*

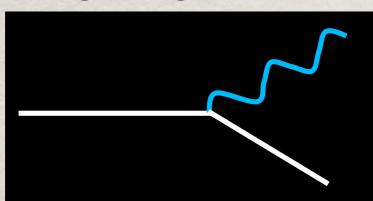


### Full NLO $O(\alpha_s)$ calculations, including NLO PDF's.



\* Dawson, Ismail, Low, 2014

## **GAUGE-BOSON INITIATED PROCESSES** At colliding energies **E** >> M<sub>W</sub>, EW gauge bosons are new "gluons"!

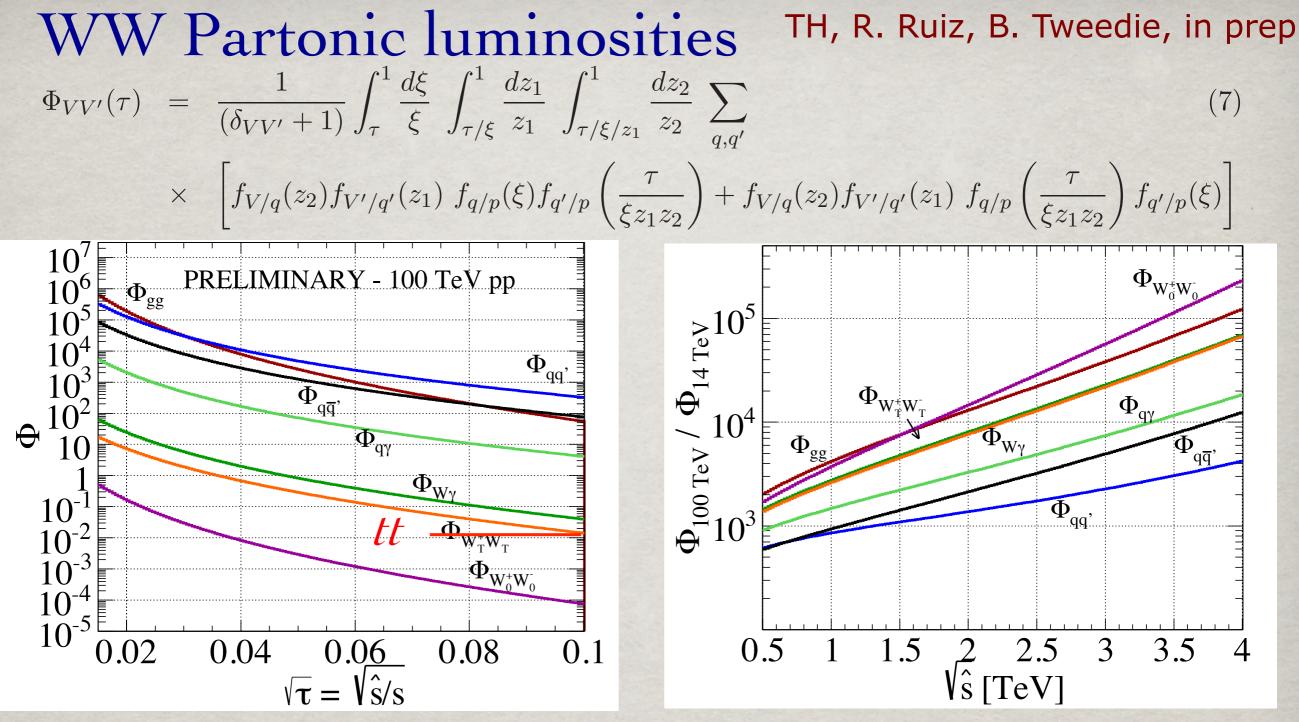


In the EW theory:

"Effective W-Approximation" (VBF → h is seen by ATLAS/CMS) S. Dawson, 1985; G. Kane et al., 1984; Chanowitz & Gailard, 1984

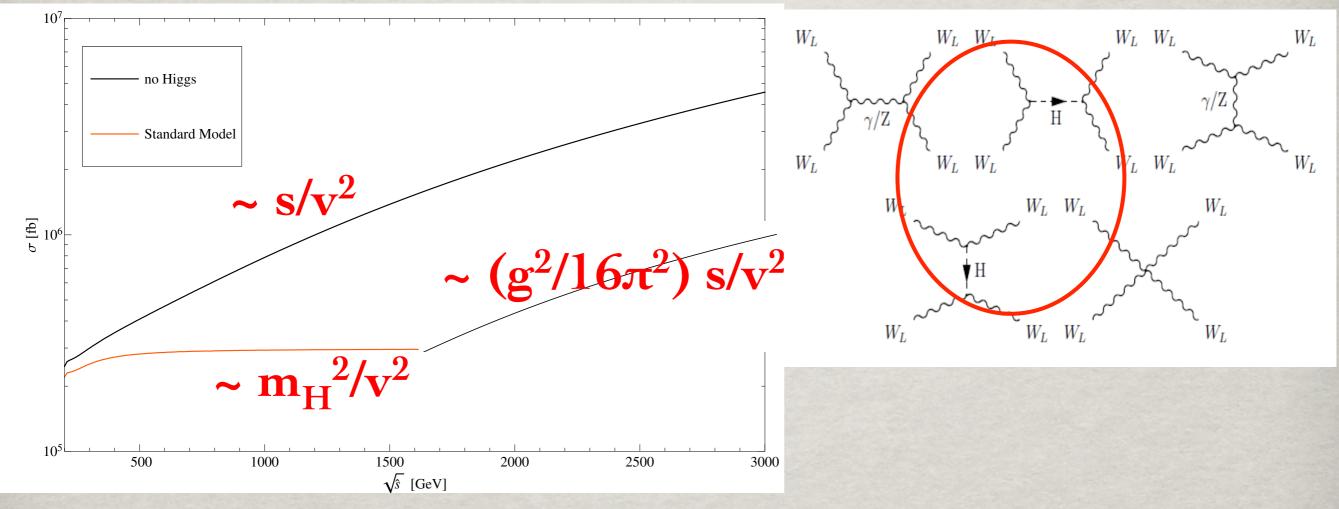
 $P_{q \to qV_T} = (g_V^2 + g_A^2) \frac{\alpha_2}{2\pi} \frac{1 + (1 - x)^2}{x} \ln \frac{Q^2}{\Lambda^2}$  $P_{q \to qV_L} = (g_V^2 + g_A^2) \frac{\alpha_2}{\pi} \frac{1 - x}{x}$ 

- $V_T$  radiation the same as  $g, \gamma : |\mathcal{M}|^2 \sim p_T^2$ :
  - "dead cone" at  $p_T \rightarrow 0$
  - log-enhancement at high  $p_T$  & soft x
- $V_L$  radiation no collinear enhancement/suppression, not the same as a scalar radiation.



Lumi(W<sup>+</sup><sub>T</sub>W<sup>-</sup><sub>T</sub>) similar size to lumi(tt); Lumi(W<sup>+</sup><sub>T</sub>W<sup>-</sup><sub>T</sub>) ~ Lumi(W<sup>±</sup>  $\gamma$ ), Electro=weak Lumi(W<sup>+</sup><sub>L</sub>W<sup>-</sup><sub>L</sub>) 100 times smaller: Goldstones Lumi(100/14) increased by 1000 – 10<sup>5</sup> for 500 GeV - 4 TeV!

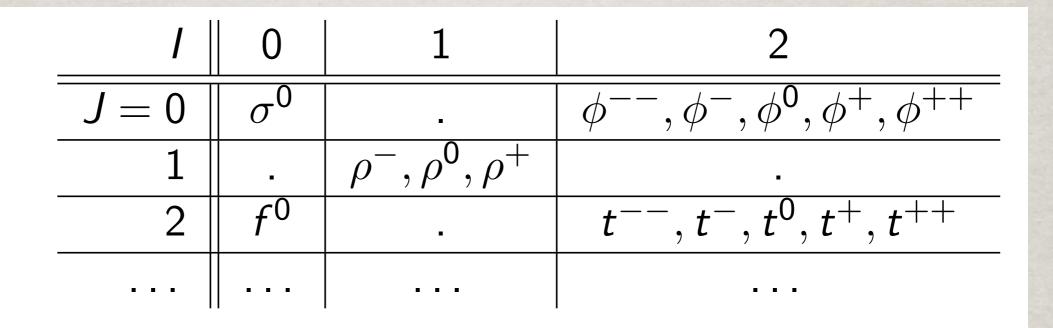
# *W<sub>L</sub>W<sub>L</sub>* Scattering: The existence of a light, weakly coupled Higgs boson unitarize the WW amplitude:



- Consistent perturbative theory up to  $\Lambda$  (?)
- New strong dynamics effects may still exist, but "delayed" to  $\rho^2/\Lambda^2$ .

## $W_L W_L$ Scattering:

### Different channels are sensitive to different physics:



- ► I = 0: resonant in  $W^+W^-$  and ZZ scattering
- ► I = 1: resonant in  $W^+Z$  and  $W^-Z$  scattering
- ▶ I = 2: resonant in  $W^+W^+$  and  $W^-W^-$  scattering

Equally important:  $WW \rightarrow HH$ , *tt* for H<sup>3</sup> & top couplings.

## MULTI GAUGE-BOSON PRODUCTION FROM PROMPT PRODUCTION

### At 100 TeV:

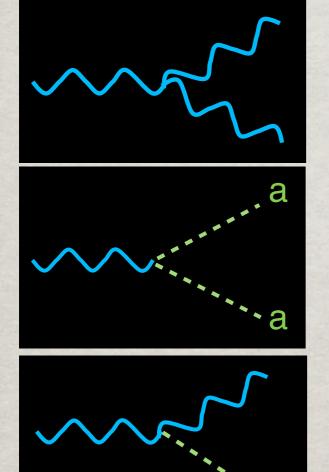
. . . .

WW	<b>σ=770 pb</b>
WWW	<b>σ=2 pb</b>
WWZ	<b>σ=1.6 pb</b>
WWWW	<b>σ=15 fb</b>
WWWZ	<b>σ=20 fb</b>

M. Mangano's talk

Each W costs you a factor of ~ 1/100 (EW coupling)

MULTI GAUGE-BOSON PRODUCTIONFROM SPLITTING/SHOWERING:At colliding energies  $E >> M_v$ ,In EW gauge boson splitting:J. Chen, TH, B. Tweedie, in prep



$$P_{V_T \to V_T V_T'} = \frac{\alpha_2}{2\pi} \left[\frac{1}{x(1-x)} + x(1-x)\right] \ln \frac{Q^2}{M_W^2}$$

$$P_{V_T \to V_L V'_L} = \frac{\alpha_2}{4\pi} x (1-x) \ln \frac{Q^2}{M_W^2}$$

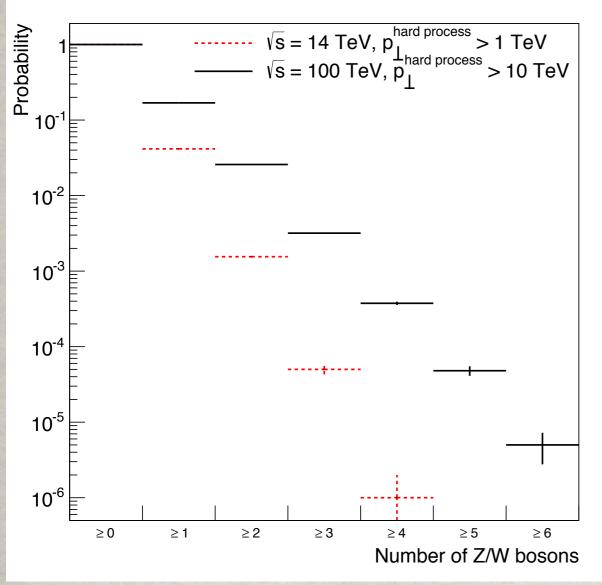
$$P_{V_T \to V_T H} = \frac{\alpha_2}{4\pi} \frac{1-x}{x}$$

- $V_T$  the "new gluons"!
- $V_L/H$  radiations the Goldstone Eq. Theo.

### **SPLITTING PROBABILITIES:**

		J. Chen, TH, B. Twee	edie, in prep
Split	Form	Rate: E=1TeV	10 TeV
$q \rightarrow qV_T$	$-2.8 \times 10^{-3} \ln^2(E/M_W)$	1.7%	(7%)
$q \rightarrow qV_L \frac{P}{P}$	$T_{1.4x}10^{-3} \ln (E/M_W)$ proportional to $gv$	0.5%	1%
$V_T \rightarrow V_T V_T$	$0.01 \mathrm{x} \ln^2(\mathrm{E/M_W})$	6%	(22%)
$V_T \rightarrow V_L V_L$	$4 x 10^{-4} \ln (E/M_W)$	) 0.15%	0.3%
$\rightarrow V_L h$	ET same pure gauge	e couplings	
$V_L \rightarrow V_T V_L$	$2x10^{-3}\ln^2(E/M_W)$	v) 1%	(4%)
$\rightarrow V_{T}h$	same ET		
$h \rightarrow V_T V_L$	same <sup>L1</sup>		
$V_T^* \rightarrow ff'$	$0.04 \mathrm{x} \ln(\mathrm{E/M}_{\mathrm{W}})$	5%	(10%)
$V_T \rightarrow V_T V_L$	$\frac{0.01 \text{x} \ln(\text{E/M}_{W})}{ET 3 \text{x} 10^{-4}}$	) 2%	(5%)
$\rightarrow V_Th$	ET 3x10-4 proportio	onal to $gv 0.03\%$	0.03%

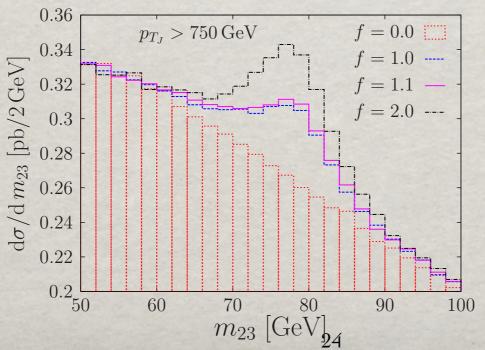
### MULTI GAUGE-BOSON FROM SHOWERING:

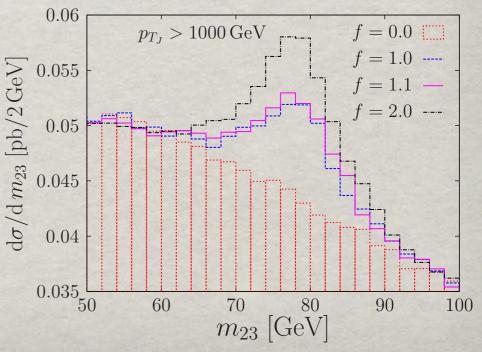


Christiansen, Sjostrand: 1401.5238 At higher energies, each *W* costs you a factor of ~ 1/10 ! We are in the process of developing a more complete EW showering code.

J. Chen, TH, B. Tweedie

Kruass, Petrov, Schonher, Spannowsky: 1403.4788





## NEW PHYSICS WITH ENERGETIC/ MULTI TOPS/GAUGE-BOSONS

SUSY examples:  $\tilde{b}\tilde{b}^* \to t\chi^- \bar{t}\chi^+$ ,  $\tilde{t}W^- \tilde{t}^*W^+ \to 4W^{\pm} b\bar{b}$ . Heavy quark examples: *TT*', *BB*', ...

Energetic  $W^{\pm}$ , Z, H, t as new radiation sources from heavy W', Z' decays &  $W_L W_L$  scattering

## OVERALL

- \* With the Higgs discovery, the SM is healthier than ever, valid to a scale up to  $\Lambda \sim ?$ tune) 10 But the Higgs sector fine-tuned  $\delta$ : 10-6 EW 10-9 \* VLHC will take the lead for searches: 10-12  $\tilde{g}, \tilde{t}, \tilde{b}, \chi^{\pm,0}, \dots H^{\pm}, A^{0}; W^{\pm'}, Z' \dots$ 10-15  $10^{-18}$ Vacc 10-21 Stability The *top*, *W*, *Z*, *H* may hold the key for 10-24 10-27 discovery!
  - Searching for new physics starts from  $10^{-30}$   $\frac{1}{10^3}$   $\frac{1}{10^6}$   $\frac{1}{10^9}$   $\frac{1}{10^{12}}$   $\frac{1}{10^{15}}$   $\frac{1}{10^{18}}$ understanding old physics in the new regime:  $\Lambda$  (GeV)

**GUTs** 

- *top*, *W*,*Z* may behave as partons to produce new heavy states;
- *top, W,Z,H* may serve as new radiation sources;
   and may help reveal new heavy states.
  - Thus, need precise understanding of the dynamics/kinematics WHILE NEW PHYSICS SEARCHES EXCITING, SM PHYSICS REMAINS RICH AT VLHC!