

Tracker for Anti-Hydrogen Free-Fall measurement

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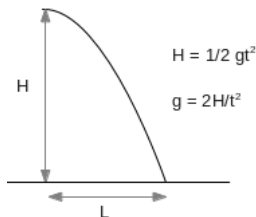
ETH Zurich

September 11, 2014

- 1 *Motivation*
- 2 *The Experiment*
 - Steps of experiment
- 3 *Shaping Velocity*
- 4 *Final Experiment measurement steps*
- 5 *My Work*
 - Micromegas Principle
- 6 *Simulations*
- 7 *Resolution for Vertex reconstruction Estimation*
 - Results from Simulation
- 8 *Micromegas Test Bench*
- 9 *Detector Design*
- 10 *Conclusion from Simulation*

- A direct test of the Equivalence Principle with antimatter.
- The acceleration imparted to a body by a gravitational field is independent of the nature of the body :
Inertial mass = Gravitational mass
- Indirect tests obtained by comparing the properties of particle and antiparticle. [1] [2] [3]
- No stringent constraints on direct tests with antimatter. (ALPHA ruled out ratios between the gravitational mass and the inertial mass of anti-hydrogen less than -65 and greater than +110) [4]

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Time of fall = T Table: Relative precision on \bar{g}

Temperature	Velocity Fluctuation	$\frac{\Delta \bar{g}}{\bar{g}}$	Detected \bar{H}
1K	100 m/sec	0.01	1.7×10^8
1mK	4 m/sec	0.01	1.7×10^5
1 μ K	0.1 m/sec	0.01	170
20 μ K	0.6 m/sec	0.01	3548
20 μ K	0.6 m/sec	0.01	9800

GBar values

$L=0.3\text{m}$ and $v_{\text{hor}}=1.5\text{m/s} \rightarrow h=10\text{ cm}$
 $(T(H) \sim 20\mu\text{K} \sim 1\text{ neV})$

GBar: cooled $\bar{H}^+ \rightarrow$ slow H

N.B: For GBar additional velocity fluctuation due to recoil of atom by laser (mentioned later)

Calculation

Classical Free-Fall: $Z=Z_0 + v_0 t + \frac{1}{2}gt^2$

Main Perturbation: $\rightarrow v_0$ (unknown)

Velocity fluctuation is Temperature equivalent:

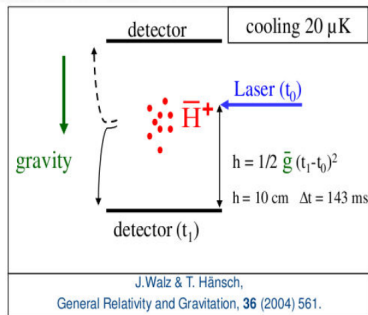
$$v_0 = \sqrt{\frac{2KT}{m}}$$

K- Boltzmann constant

T- Temperature

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- **Produce ion \bar{H}^+ .**
 $\bar{p} + oPs \rightarrow \bar{H} + e^-$
 $\bar{H} + oPs \rightarrow \bar{H}^+ + e^-$
- **Capture ion \bar{H}^+ .**
Segmented RF Paul Trap for capture.
- **Sympathetic cooling to $20\mu\text{K}$.**
Sympathetic cooling using Be^+ ions.
- **Photodetachment of e^+ .**
Using laser. Laser shot gives the start time.
- **Time of flight.**
Detection of \bar{H} annihilation gives the stop time and thus the time of flight



Improve velocity dispersion

Photodetachment of $e^+ \rightarrow$ Additional recoil velocity \rightarrow Additional velocity fluctuation

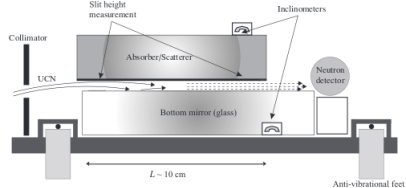
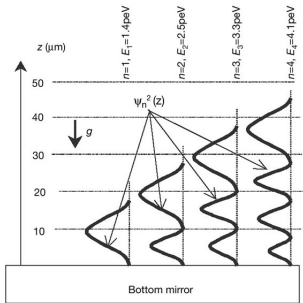
Threshold energy = \bar{H}^+ binding energy = 0.76 eV

Energy above threshold, ΔE , for laser $\sim 15\mu\text{eV}$

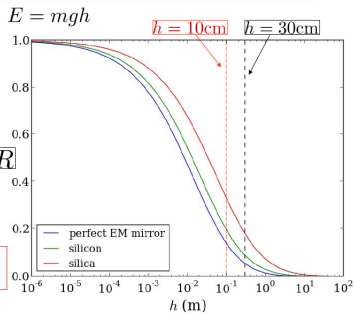
Velocity Shaping using Quantum Reflections \rightarrow Reduced velocity dispersion \rightarrow Get required precision for fewer statistics.

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 - Results from Simulation
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- 9 *Detector Design*
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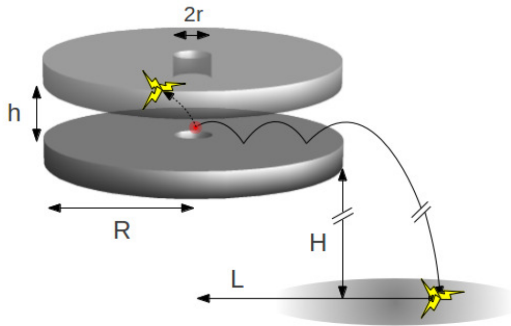
Quantum Reflection



The probability of neutron tunneling through the gravitational barrier at height z corresponding to the n^{th} quantum state is proportional to the square of the wave-function $\Psi_n^2(z)$. E_n is the energy of the n^{th} quantum state [5]



Due to the high efficiency of Quantum Reflections at small energies because of the Casimir interaction, atoms with sufficiently low vertical velocities may bounce off the bottom mirror. This phenomena has already been demonstrated for neutrons in a study of the neutron quantum states in the potential well formed by the earth's gravitational field and a horizontal mirror [6].



To shape the vertical velocity to reduce the dispersion, a similar shaping scheme with a bottom mirror is envisaged. For this $\Delta v = \sqrt{2\bar{g}h}$
 Smaller the height h , lesser will be the velocity dispersion and lesser the required statistics for the aimed precision on \bar{g} .

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The calculation of \bar{g} thus comprises calculating the following:

- Total time of flight(T) = Time of detection of annihilation-Time of laser shot.
- Time of free-fall = Total time of flight-Time in shaper.
- Time in shaper = $\frac{R}{v_{horizontal}}$
- $v_{horizontal} = \frac{L}{T}$

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Thus L needs to be measured accurately to improve resolution on \bar{g} .

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My work involves building the detector for the experiment to reconstruct the vertex of annihilation. For the first phase of the experiment to reach 1% precision on \bar{g} a vertex resolution $\sim 4\text{mm}$ is necessary.

- Initially a TPC was thought to be used for the pion track reconstruction. But my simulations showed a XY Microstrip Micromegas achieves much better spatial resolution.
- Advantages of the Microstrip Micromegas includes:
 - ▶ No need for drift time calculation for track points since the X-Y strips give directly the position.
 - ▶ Much reduced HV(few hundred volts instead of kV)
 - ▶ Much lighter detector.
 - ▶ Spatial resolution is better.
 - ▶ Lesser cost.

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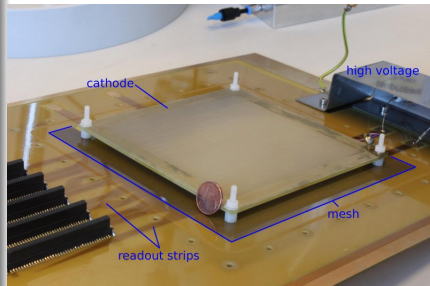
- Two-stage parallel plate avalanche chamber.

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- Ni electroformed micromesh cathodes.

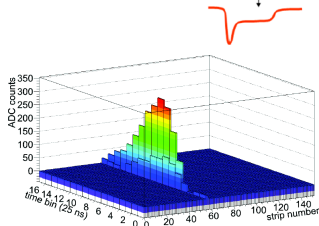
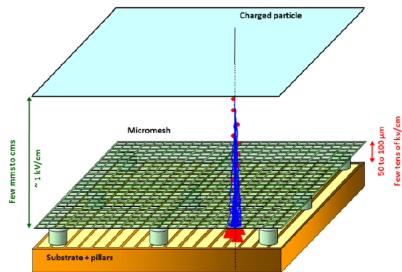
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- Amplification gap of 128 μm between the mesh and the anode strips of pitch $\sim 320\mu\text{m}$ with a high E-Field of 40-50kV/cm.

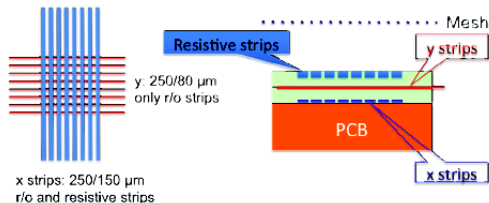
- Two-stage parallel plate avalanche chamber.
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- 3mm conversion gap between drift cathode and the mesh with E-Field of 100-1000V/cm.
- Amplification gap of $128\ \mu\text{m}$ between the mesh and the anode strips of pitch $\sim 320\ \mu\text{m}$ with a high E-Field of 40-50kV/cm.
- $150\ \mu\text{m}$ diameter epoxy spacers placed every 2mm between anode and mesh to maintain parallelism.



- Incoming particle ionizes gas (eg: Ar- C_4H_{10} , Ar CF_4 etc. mixture few mbars above atmospheric pressure).
- Electrons drift towards the mesh from the conversion gap under applied electric field.
- Electrons enter the amplification region and produces an avalanche like secondary ionization due to strong electric field.
- Signal induced on the strips is a sum of the electron and ion signal (e^- signal rise time ~ 1 ns followed by the ion tail ~ 100 ns depending on gas mixture and amplification gap) .



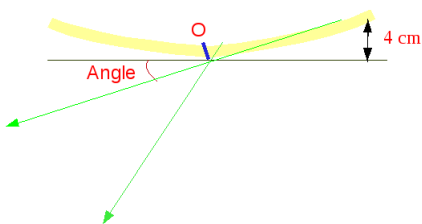
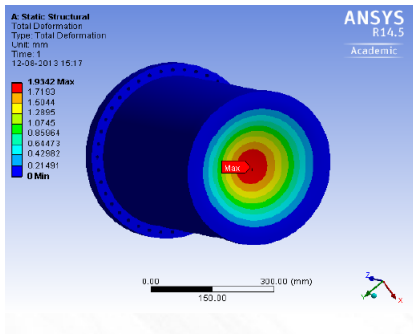
XY Micromegas for two co-ordinate readout from the same gas gap:



Crossed strips (R16)

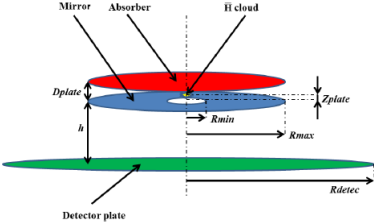
Each layer separated by 70 μm thick layer of insulator. Varying strip width to compensate for the weak capacitive coupling to the X-strips. Achieved charge ratio X:Y=1:2.4. Spatial resolution achieved for each strip layer ~ 80 microns.

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- 10 *Conclusion from Simulation*

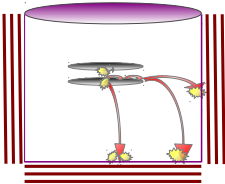


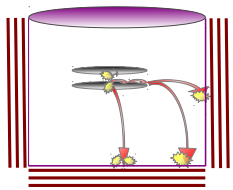
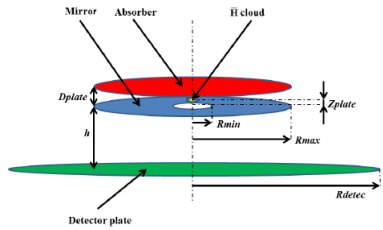
- The resolution is limited by the intrinsic detector resolution and multiple scattering in the annihilation plane and the detector material.
- It would be ideal to make the annihilation plane as thin as possible to limit multiple scattering but due to vacuum deformation we can't make it too thin.
- Vacuum deformation was simulated using ANSYS for different thicknesses and materials to estimate the minimum plane thickness required.
- The plane could have been thinner for a round bottom cylinder but that worsens the resolution due to the longer lever arm for angular tracks as shown in the figure.
- For a flat bottom cylinder simulations showed a plane thickness ~ 4 mm of Stainless Steel is the accepted option (staying under the elastic deformation limit) with a 2 mm deformation.

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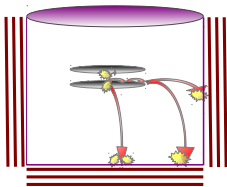
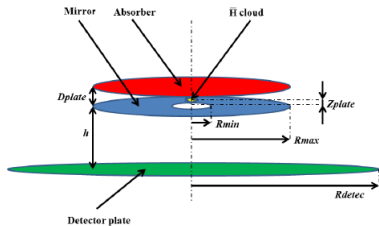


- A Geant4 code was written to simulate anti-hydrogen annihilation.

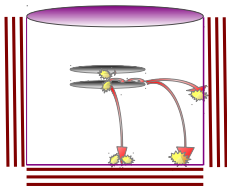
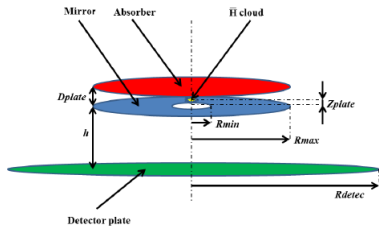




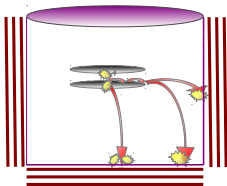
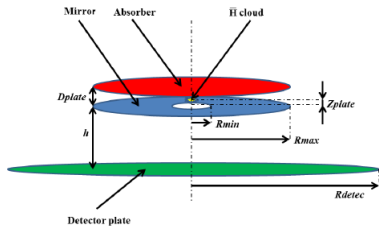
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- The motion of the anti-hydrogen through the shaper was added to the code and its subsequent position on the detector was obtained. The parameters considered for the simulation:



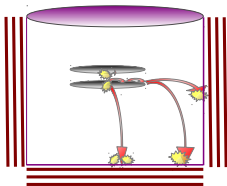
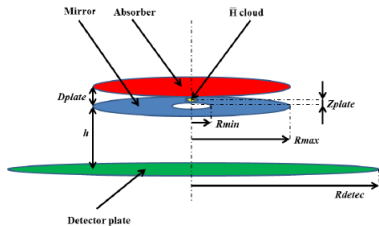
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 - ▶ Free-fall height = 10cm
 - ▶ $\Delta E = 15 \mu\text{eV}$ (laser energy above threshold)
 - ▶ $T = 10 \mu\text{K}$, $\epsilon = 3$ (ratio of horizontal to vertical temperature)
 - ▶ $Z_{\text{Plate}} = 1.5\text{mm}$ $D_{\text{Plate}} = 3\text{mm}$
 - ▶ $R_{\text{min}} = 5\text{mm}$ $R_{\text{max}} = 30\text{mm}$



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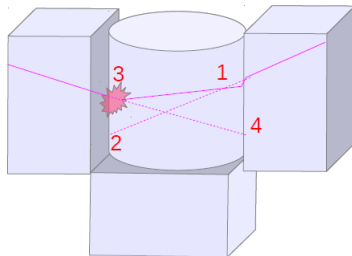
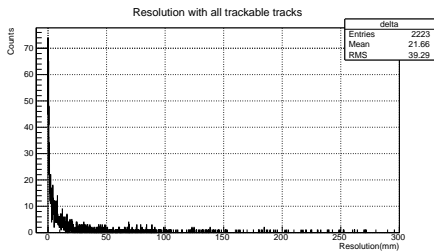
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- The annihilation chamber is a 50cm diameter 4mm thick cylinder with 50cm \times 50cm Micromegas modules surrounding the cylinder on the four sides and bottom. In this configuration the corners at the sides are not covered.



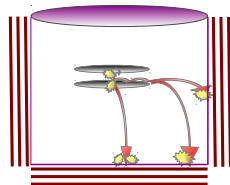
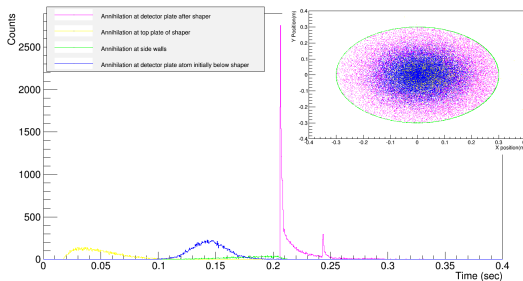
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- The resolution for vertex reconstruction is estimated and compared tracking all tracks and only the ones through the bottom detector modules with a time-cut.

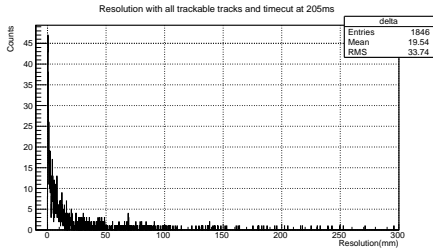
- 1 *Motivation*
- 2 *The Experiment*
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The plot shows a long tail which is because of the annihilation on the side walls. With the side annihilation there is always an ambiguity in the position of annihilation with there being two possibilities on two wall faces.



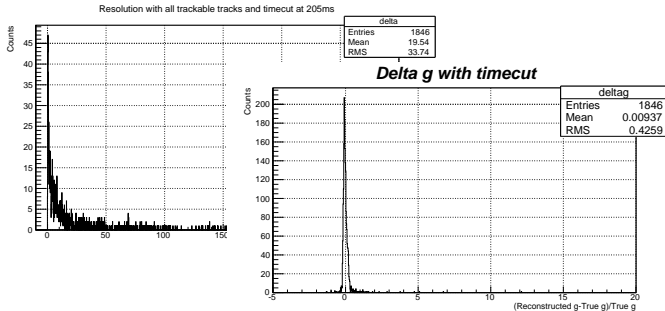
The time distribution shows a time cut at $\sim 205\text{ms}$ will include only the events with annihilation near or at the bottom plate.





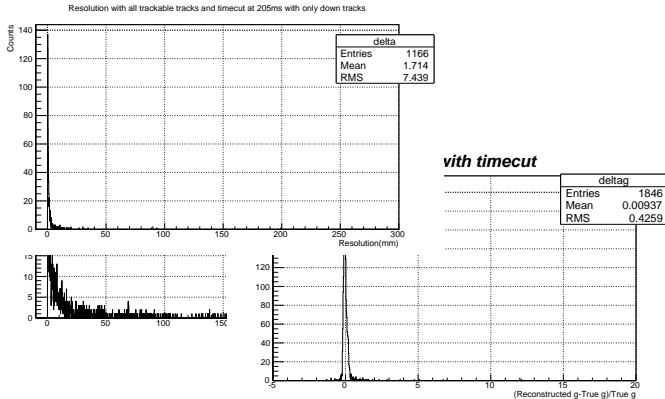
T

With the time cut we have 83% of the events through the shaper with $\sigma_g = 0.43$. We still have the tail as we include the side tracks for reconstruction. Using only the down tracks for vertex reconstruction we eliminate the tail with an improved $\sigma_g = 0.13$. Thus to achieve 1% resolution on g we will need around 170 detected events with this scheme.



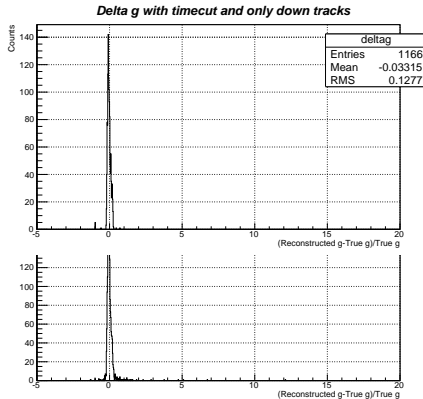
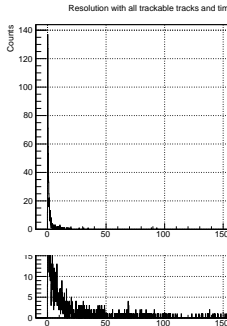
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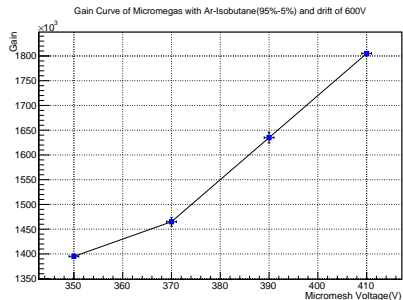
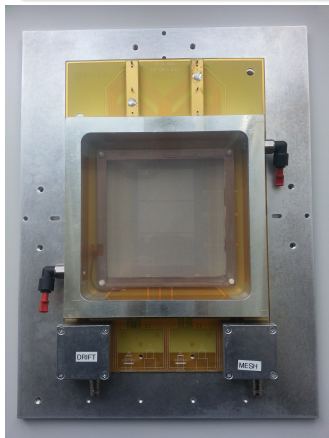
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Table: Efficiency of tracking

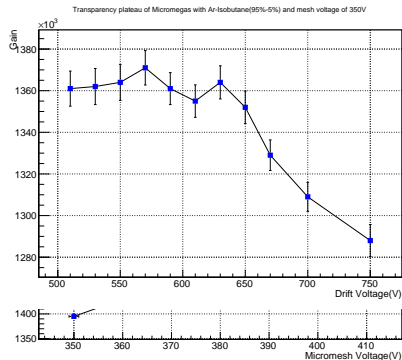
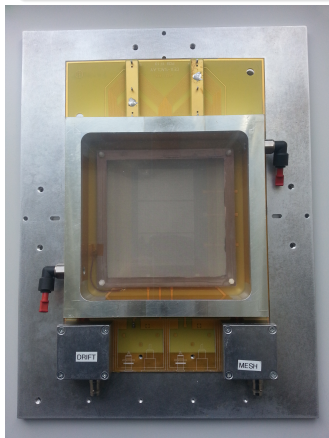
	Full Coverage at the sides	Corners not covered
Total number of events simulated	10000	10000
Number of events through shaper	2223	2223
Events with time cut at 205ms	1846	1846
Events with the time cut and at least 1 trackable track through the bottom modules	1166	1166
Events with the time cut and at least 2 trackable tracks with 1 going through the bottom modules	956	934
Events with the time cut and at least 3 trackable tracks with 1 going through the bottom modules	425	403

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A test bench is set up at ETH, Zurich with a 1-D 10cm×6cm Micromegas prototype received from Saclay. Preliminary tests to characterize the detector have already been performed reading signals off the mesh. Electronics to read the strips will be available soon.

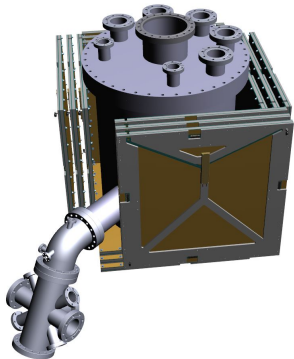


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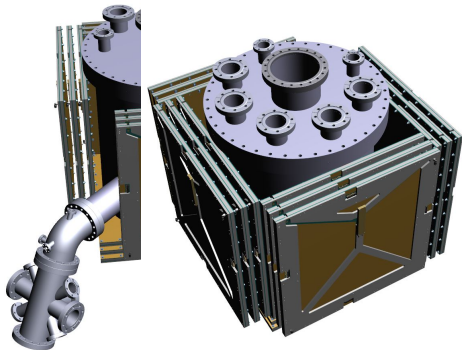


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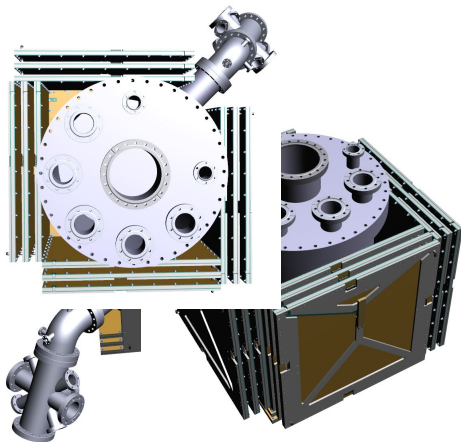
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- The production of a prototype (a triplet) with mechanics and electronics is foreseen following which it will be tested with cosmics.


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Thank You!!!