

# Low masses in the ALICE muon spectrometer

- ❑ Low masses and chiral symmetry
- ❑ Low masses in ALICE

# What are the low masses ?

Resonance	$\rho$	$\omega$	$\phi$	$J/\psi$	$\Upsilon$
Mass (MeV/c <sup>2</sup> )	770	782	1020	3097	9460
Width (MeV/c <sup>2</sup> )	150	8.4	4.4	0.087	0.052
$c\tau$ (fm)	1.3	23.4	45	2268	3752
BR $\mu^+ \mu^-$ (%)	$4.6 \cdot 10^{-3}$	$9.0 \cdot 10^{-3}$	$2.9 \cdot 10^{-2}$	5.9	2.5

- Very short mean lifetime

$$\tau_{\rho} \ll \tau_{QGP} \sim 10 \text{ fm} < \tau_{\omega, \phi} \ll \tau_{J/\psi, \Upsilon}$$

- Small B.R. in dimuon, but important production :

$$N_{\mu\mu} (\text{low masses}) \sim N_{\mu\mu} (J/\psi)$$

# Physics motivations

- Formation of a QGP leads to a general strangeness enhancement .
  - Study of the  $\phi$  ( $s\bar{s}$ ) production rate.
- Chiral symmetry restoration
  - Study of the  $\rho$  spectral function.

# The chiral symmetry

$$\mathcal{L}_{QCD} = -\frac{1}{4} G_a^{\mu\nu} G_{\mu\nu}^a + i \bar{u} \gamma^\mu D_\mu u + i \bar{d} \gamma^\mu D_\mu d - m_u \bar{u} u - m_d \bar{d} d + \dots$$

$$\psi = \begin{pmatrix} u \\ d \end{pmatrix}$$

If we do not take into account the gluons and the heavier quarks :

$$\mathcal{L} = i \bar{\psi} \gamma^\mu \partial_\mu \psi - \frac{m_u + m_d}{2} \bar{\psi} \psi - \frac{m_u - m_d}{2} \bar{\psi} \tau_3 \psi$$

$$m_u, m_d \approx 1,5 - 7 \text{ MeV} \ll \Lambda_{QCD} \approx 1 \text{ GeV}$$

$$\mathcal{L} = i \bar{\psi} \gamma^\mu \partial_\mu \psi$$

# The chiral symmetry (II)

$$\begin{aligned}\mathcal{L} &= i\bar{\psi}\gamma^\mu\partial_\mu\psi \\ &= i\bar{\psi}_L\gamma^\mu\partial_\mu\psi_L + i\bar{\psi}_R\gamma^\mu\partial_\mu\psi_R\end{aligned}$$

Left and right chiralities transform separately ( $SU(2)_L \otimes SU(2)_R$ )

↔ The Lagrangian is invariant under chiral transformations :

$$\psi_L \rightarrow e^{i\vec{\sigma} \cdot \vec{\epsilon}_L / 2} \psi_L$$

$$\psi_R \rightarrow e^{i\vec{\sigma} \cdot \vec{\epsilon}_R / 2} \psi_R$$

⇒ The QCD presents the chiral symmetry (in the light quarks limit).

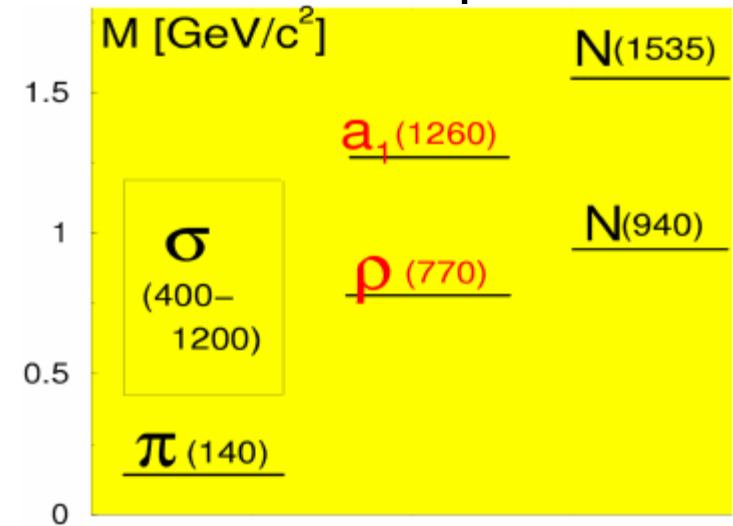
# Chiral partners

- Consequence of the chiral symmetry :
  - Each particle has a chiral partner with same spin and opposite parity.

$$\Pi(0^-) \leftrightarrow \sigma(0^+)$$

$$\rho(1^-) \leftrightarrow a_1(1^+)$$

$$N(940)(1/2^+) \leftrightarrow N(1535)(1/2^-)$$



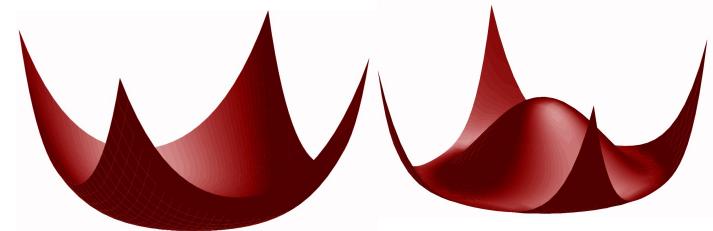
If chiral symmetry was realized then chiral partner would have the same mass. This is not the case.

⇒ Chiral symmetry is broken

# Chiral symmetry breaking

- Chiral symmetry is spontaneously broken :

- $\mathcal{L}_{\text{QCD}}$  is chiral symmetric, but
  - the vacuum is not ( $\sim$  Higgs mechanism).



- Order parameter of the chiral symmetry :

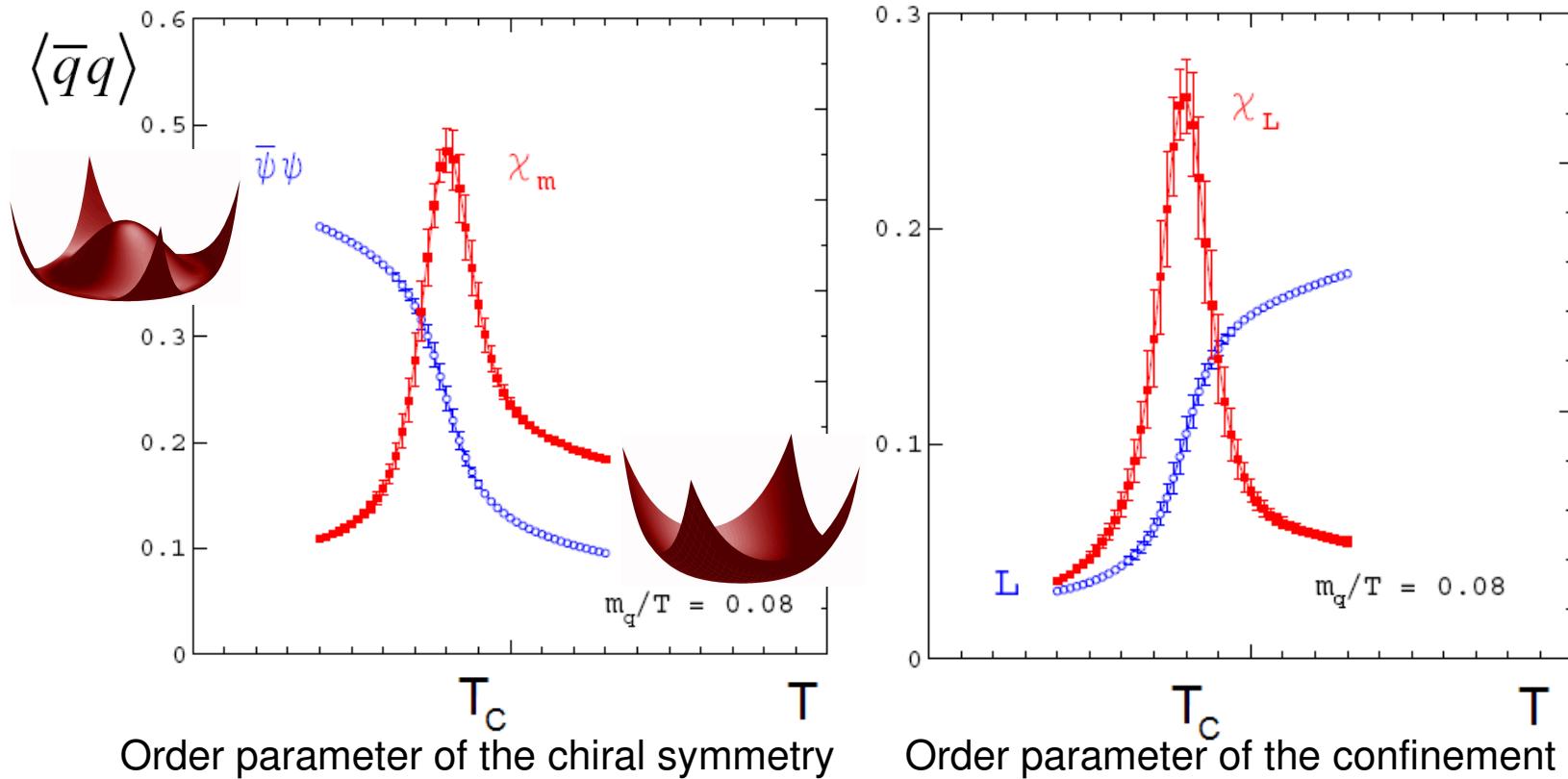
the quark condensate :  $\langle \bar{q}q \rangle = \frac{1}{2} \langle \bar{q}_R q_L + \bar{q}_L q_R \rangle$

- In the QCD vacuum :

$$\langle \bar{q}q \rangle_0 \approx -(240 \text{ MeV})^3$$

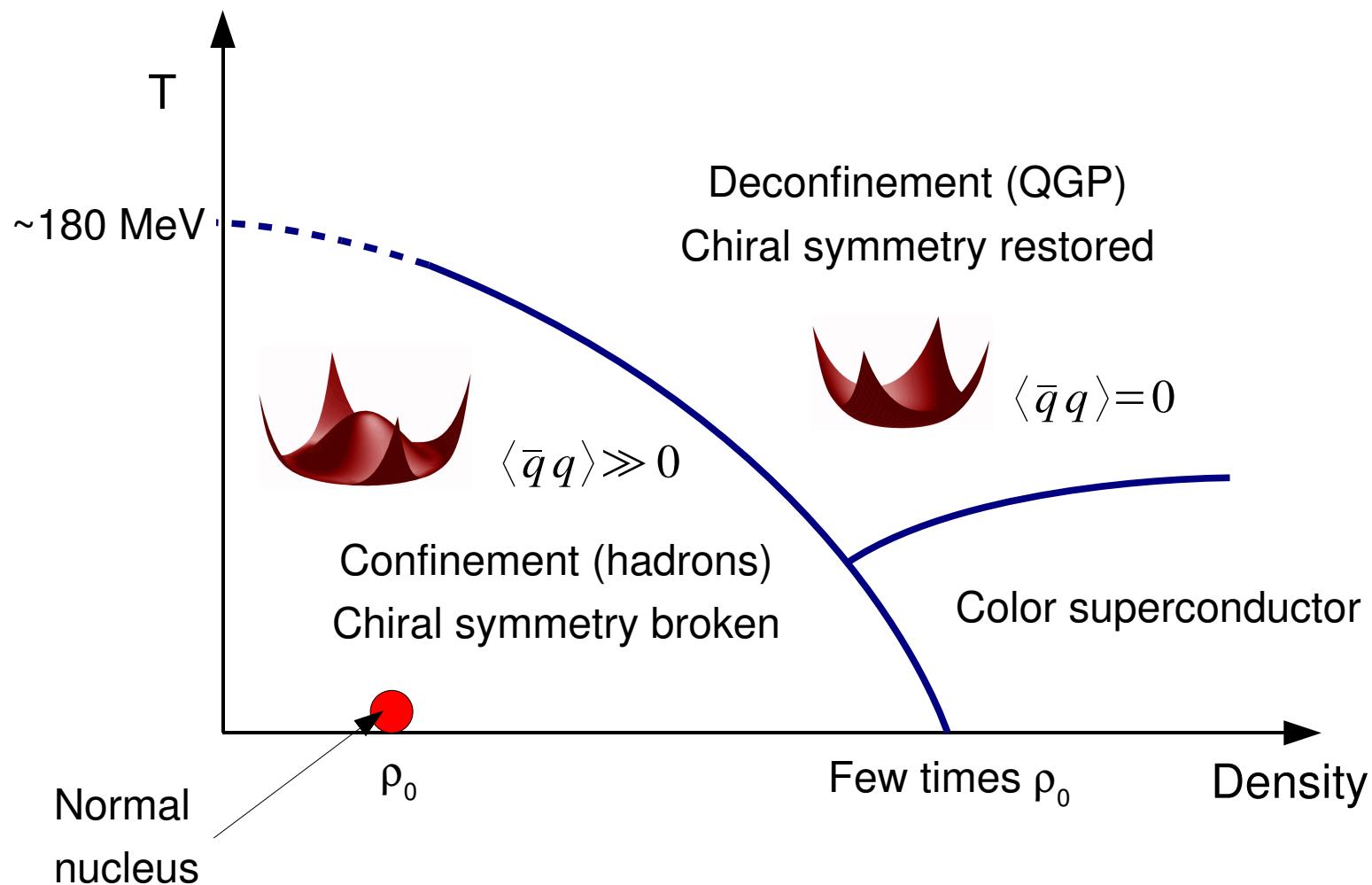
(measured from the decay of the pions)

# Chiral symmetry restoration



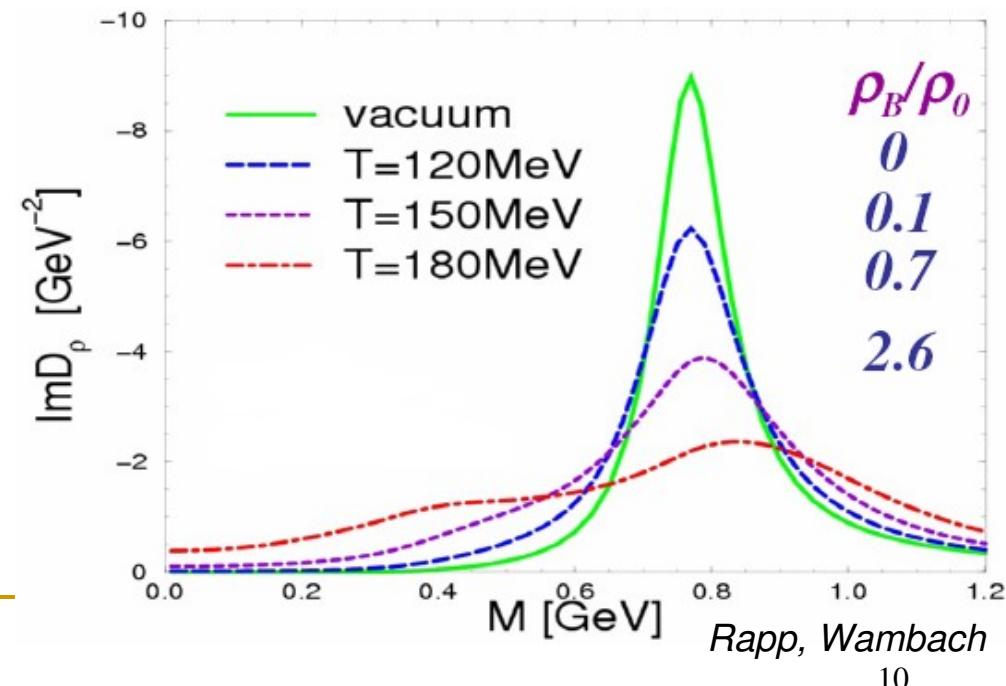
- Lattice QCD calculation show a drop of the quark condensate value when  $T=T_C \rightarrow$  phase transition: **chiral symmetry restoration**
- $T_C(\text{chiral})=T_C(\text{deconfinement})$  !

# Phase diagram



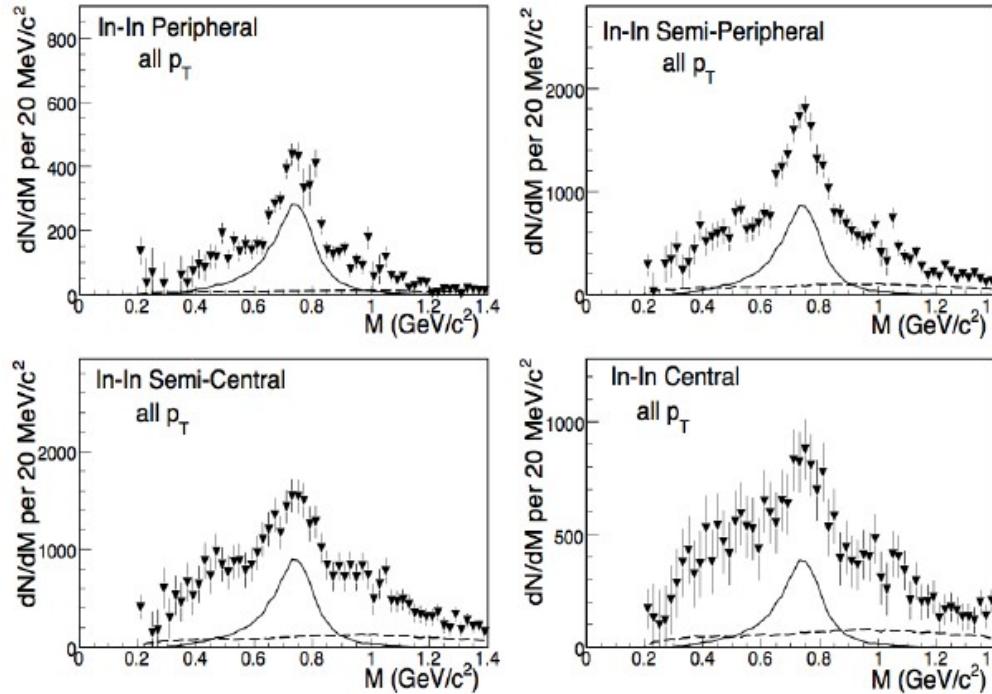
# Consequences on the $\rho$ spectrum

- When the temperature increases, the chiral symmetry restoration induces a mixing between chiral partners.
  - For example  $\rho$  and  $a_1$ .
- The  $\rho$  spectral function in a hot and dense medium has been calculated.
- Prediction :  
Broadening of the spectral function with the increase of the temperature and density.



# Experimental results : NA60

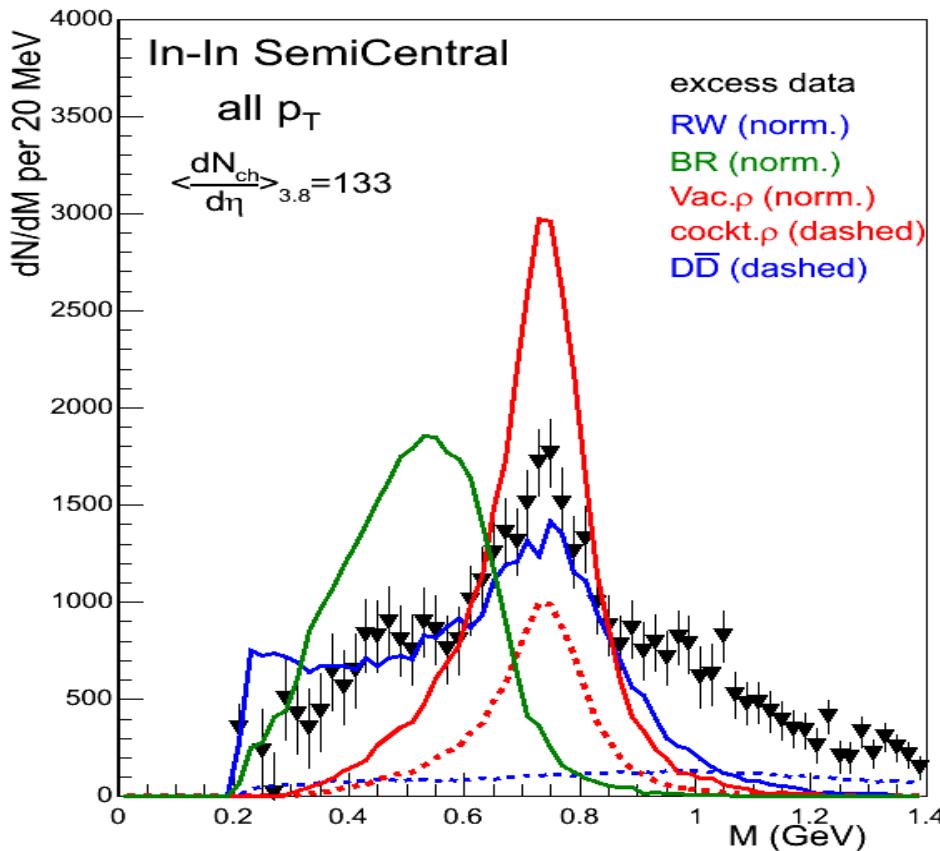
NA60 at SPS measured the rho spectral function in In-In collisions.



Eur.Phys.J.C 49 (2007) 235

- Net excess peaked at  $M_\rho$ , increasing with the centrality. Explained by  $\pi\pi \rightarrow \rho \rightarrow \mu^+\mu^-$
- Broadening with the centrality. Explained by in medium effects.

# Theory vs NA60

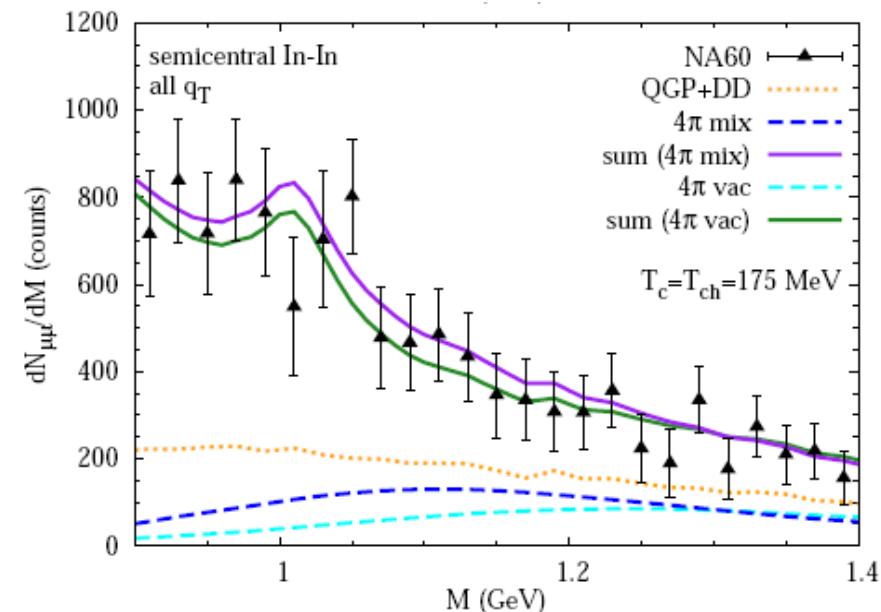
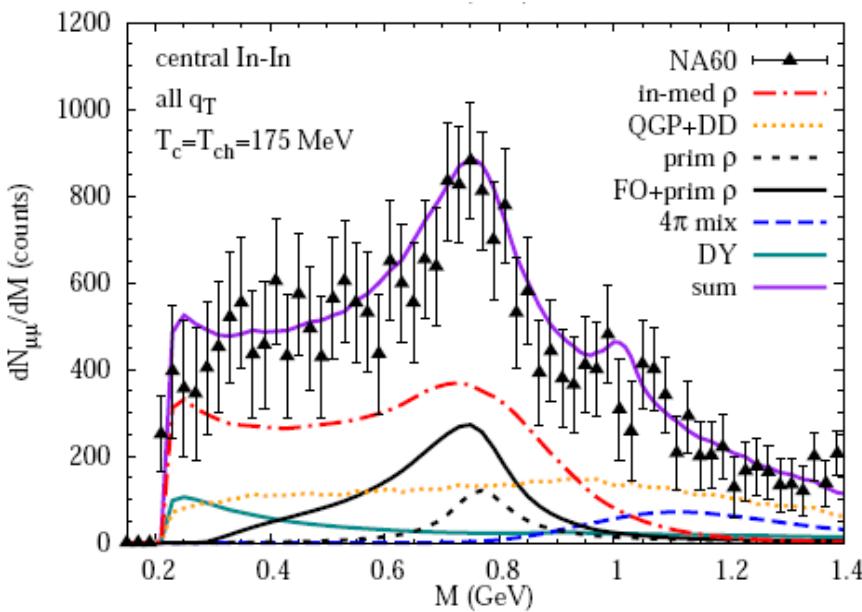


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The scenario of a rho spectrum broadening describes the data.

The scenarios of unmodified rho (red) and of mass variation (green) are not observed.

# Theory vs NA60, latest results



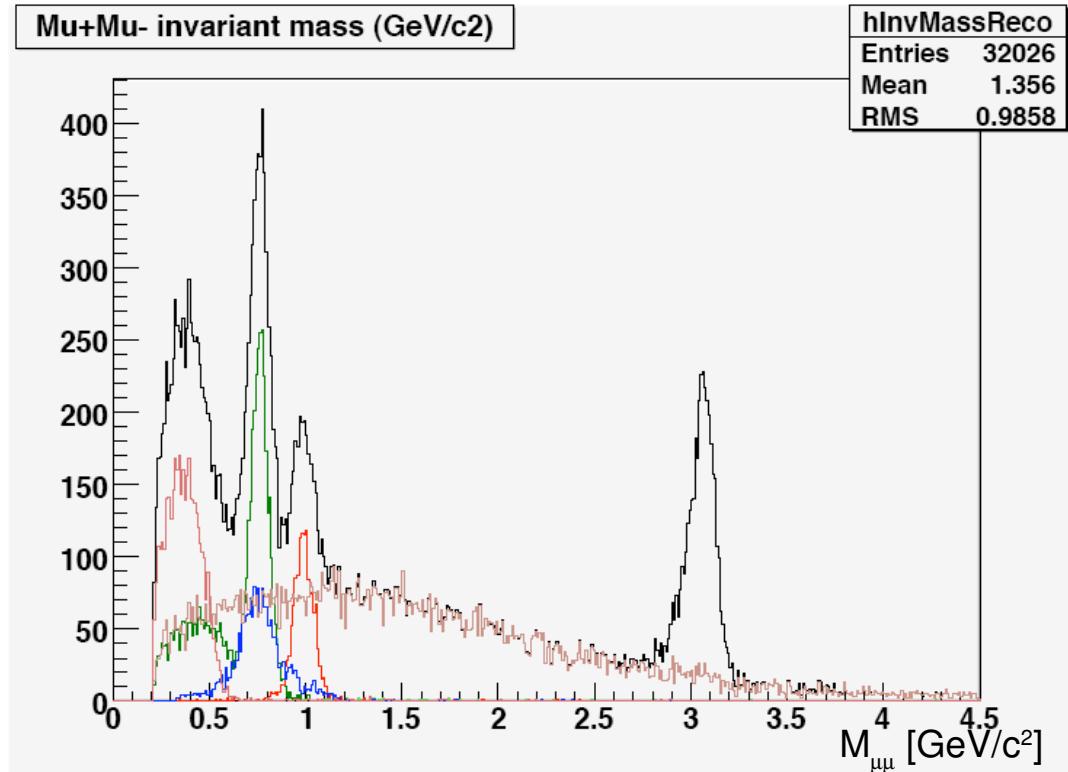
Van Hess, Rapp 04/08

Recent calculations reproduce well the data in all mass range (left panel)  
 → In medium effects are well understood.

Direct chiral mixing has a too small effect to be proven (right panel)  
 → perhaps in ALICE ?

# Prediction for ALICE

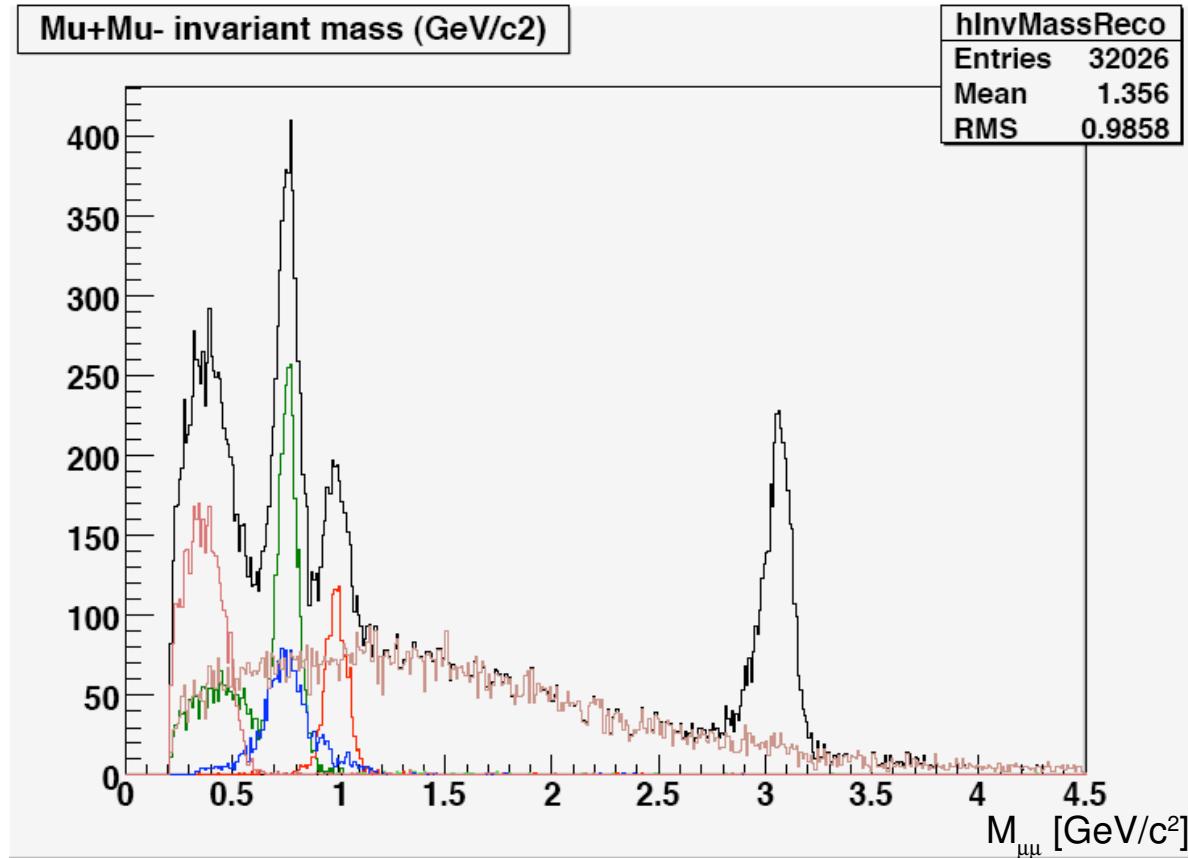
# ALICE: pp



Expected signal in the ALICE dimuon spectrometer, without the uncorrelated background (~50000 pp events of PDC06,  $p_T$  cut at 0.5  $\text{GeV}/c$ ).

Contributions of the rho (blue), omega (green), phi (red), eta (dark red) and open c/b (brown).

# ALICE: pp (II)



With uncorrelated background (fast simulations for 1 month of data taking).

The signal over noise ratio is still good (S/B  $\sim 0.3 - 1$ ).

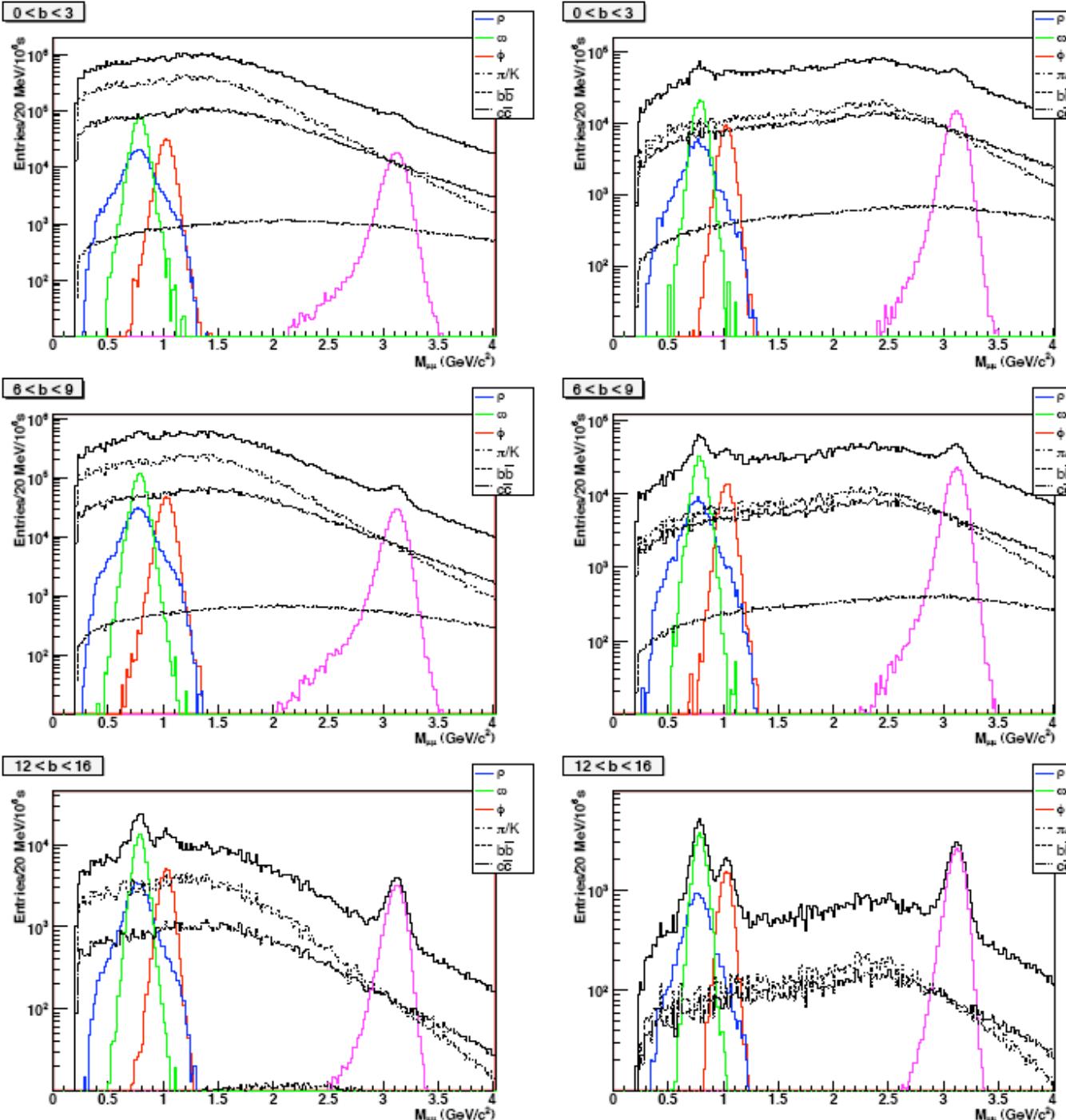
# Pb-Pb

Fast simulations, for  
 $p_T > 0.5 \text{ GeV}$  (left) and  
 $p_T > 1 \text{ GeV}$  (right),

for 3 different centralities

Very important  
 background (especially at  
 low  $p_T$ ).

Problem : in medium  
 effects more important  
 at low  $p_T$ .



# Signal over background ratio in Pb-Pb

For one month of data taking :

Meson	b range (fm)	S[ $\times 10^3$ ]	B[ $\times 10^3$ ]	S/B	Significance
$\rho$	0 - 3	255	16703	0.02	62
	6 - 9	381	9672	0.04	120
	12 - 16	41	167	0.25	92
$\omega$	0 - 3	503	8485	0.06	168
	6 - 9	755	5178	0.15	310
	12 - 16	82	89	0.92	199
$\phi$	0 - 3	210	10668	0.02	64
	6 - 9	323	6125	0.05	127
	12 - 16	34	112	0.31	91

S/B are small, but significance is high thanks to high statistics.

# Conclusions

- Low masses dimuons are important for the study of chiral symmetry.
- This study is possible in ALICE, thanks to very high statistics.
- But : an important work on the signal extraction will be necessary.

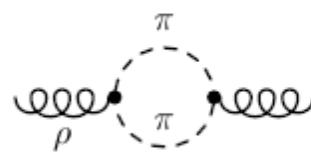
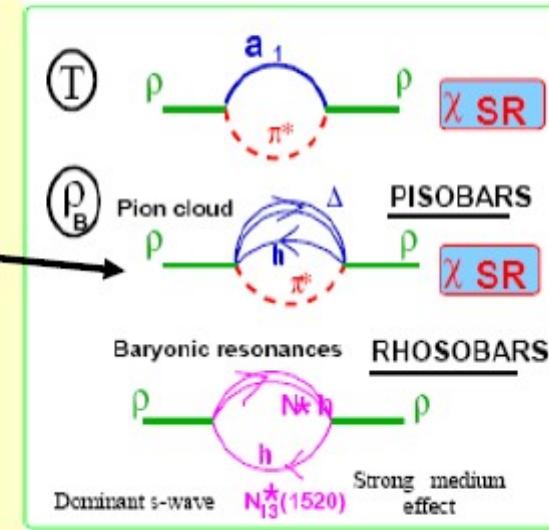
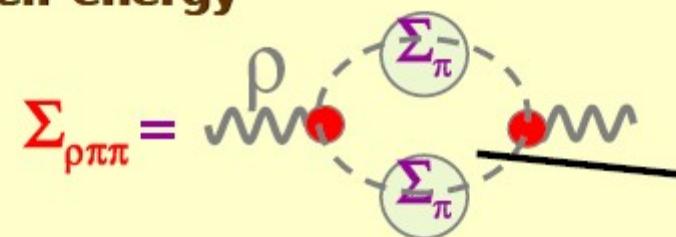
# Backup

## e. Many body approach (R. Rapp, J. Wambach, G. Colangelo)

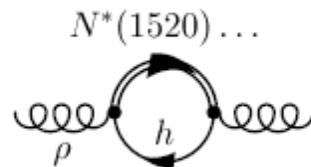
**$\rho$  propagator**

$$D_\rho(M, \vec{q}, \mu_{B,T}) = (M^2 - m_\rho^2 - \Sigma_{\rho\pi\pi} - \Sigma_{\rho B} - \Sigma_{\rho M})^{-1}$$

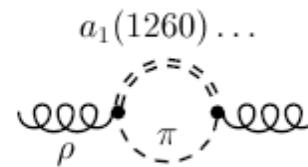
**$\rho$  self-energy**



(a)



(b)



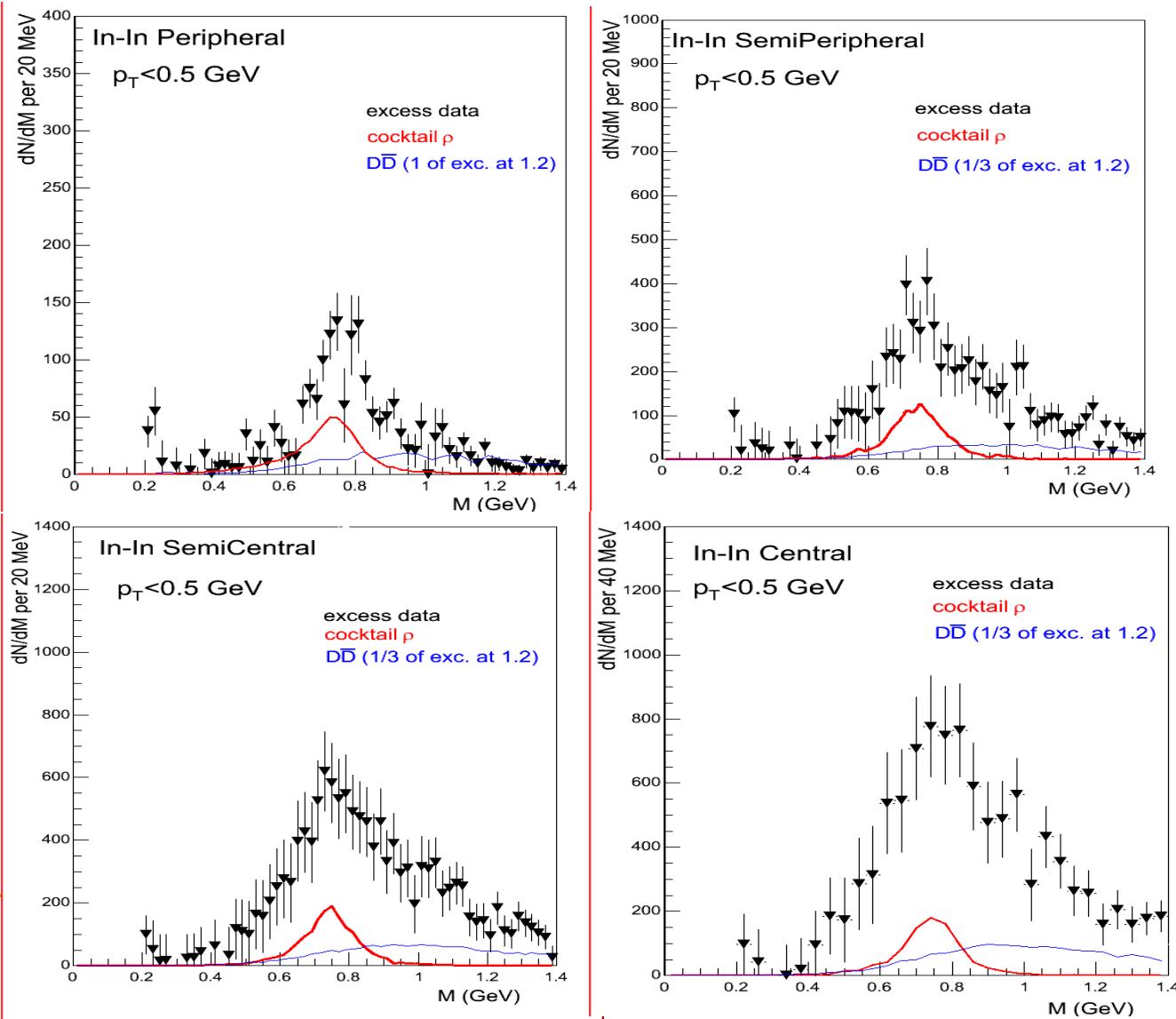
(c)

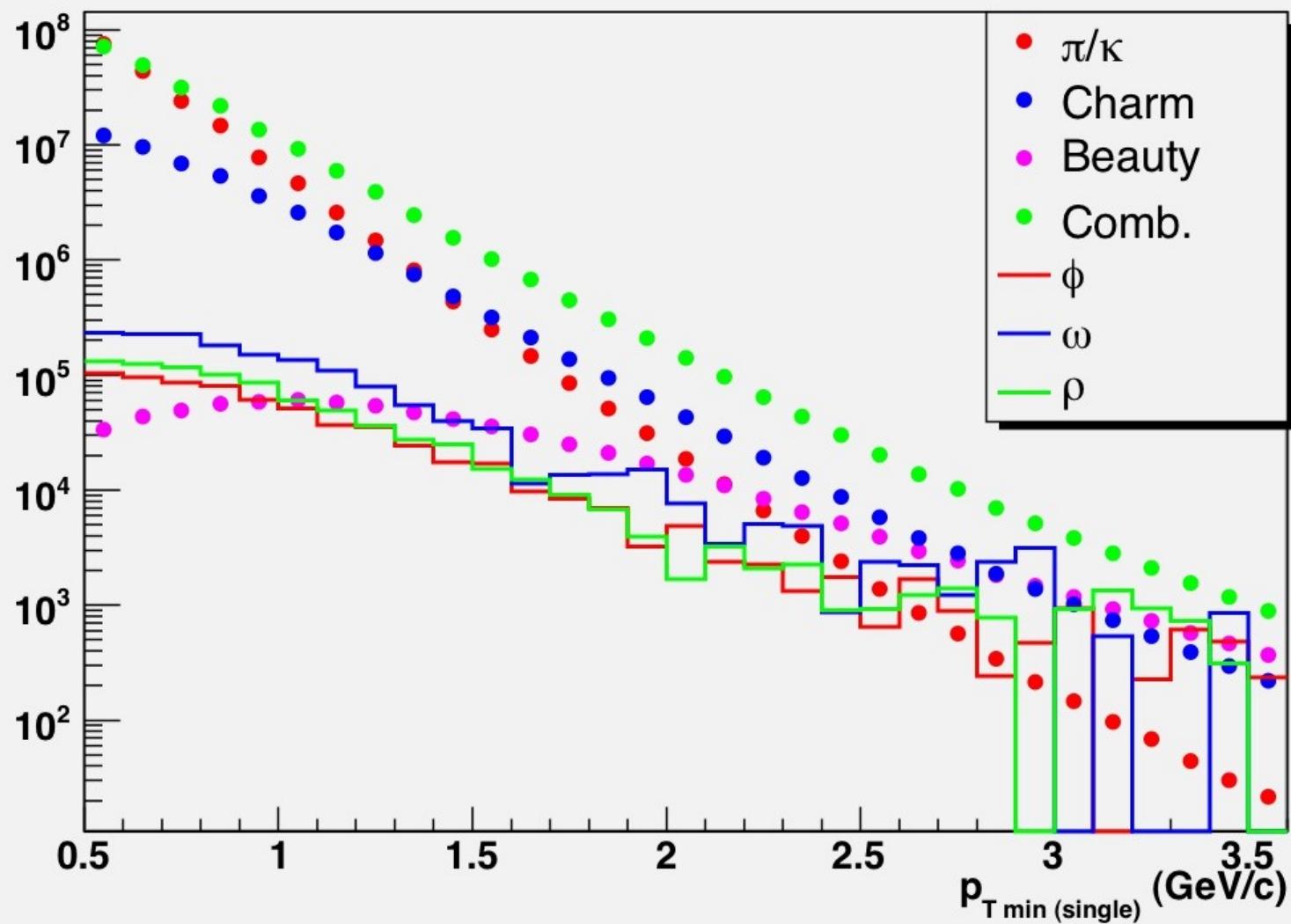
Figure 2: Sources of medium effects induced by interactions of the  $\rho$  meson in hot and dense hadronic matter: (a) renormalization of its pion cloud due to modified pion propagators, and direct interactions of the  $\rho$  meson with (b) baryons and (c) mesons, typically approximated by baryon- and meson-resonance excitations [11, 12].

**Table 1.5:** Idem as Table 1.5 using the "Low"  $p_T$  trigger condition ( $p_T^{single} > 1 \text{ GeV/c}$ ).

$p_T^{single} > 1 \text{ GeV/c}$					
Meson	b range (fm)	S[ $\times 10^3$ ]	B[ $\times 10^3$ ]	S/B	Significance
$\rho$	0 - 3	67	861	0.08	70
	3 - 6	123	1058	0.12	114
	6 - 9	104	505	0.21	133
	9 - 12	50	106	0.47	128
	12 - 16	11	8	1.37	81
$\omega$	0 - 3	137	458	0.30	178
	3 - 6	236	564	0.42	265
	6 - 9	206	267	0.77	300
	9 - 12	100	57	1.74	253
	12 - 16	23	4	5.40	140
$\phi$	0 - 3	61	569	0.11	78
	3 - 6	109	710	0.15	121
	6 - 9	93	339	0.28	142
	9 - 12	45	72	0.63	133
	12 - 16	10	5	1.71	80

# NA60, low pT





**Table 1.6:** Expected resonance signal (S), background (B) yields (in unit of  $10^3$ ), signal-to-background ratio (S/B) and significance ( $S/\sqrt{S+B}$ ) for  $\rho$ ,  $\omega$  and  $\phi$  for five bins in  $p_T$  for the most central bin ( $0 < b < 3$ ). For each column, the first number gives the result using the "Natural"  $p_T$  trigger cut ( $p_T^{single} > 0.5$  GeV/c) and the number in parenthesis the result using the "Low"  $p_T$  trigger cut ( $p_T^{single} > 1$  GeV/c).

Meson	$p_T$ range (GeV/c)	S[ $\times 10^3$ ]	B[ $\times 10^3$ ]	S/B	$S/\sqrt{S+B}$
$\rho$	0 - 2	104 (0.4)	12142 (39)	0.01 (0.01)	30 (1.8)
	2 - 4	136 (49)	4243 (751)	0.03 (0.06)	65 (55)
	4 - 6	20 (15)	138 (58)	0.14 (0.27)	50 (57)
	6 - 8	3.7 (3.3)	13 (4.9)	0.29 (0.68)	29 (37)
	8 - 10	0.8 (0.7)	2.0 (0.7)	0.38 (1.02)	15 (19)
$\omega$	0 - 2	210 (0.8)	6176 (24)	0.03 (0.03)	83 (4.9)
	2 - 4	274 (100)	2234 (398)	0.12 (0.25)	173 (142)
	4 - 6	42 (33)	74 (32)	0.56 (1.05)	122 (131)
	6 - 8	6.9 (6.4)	7.0 (2.7)	0.99 (2.39)	59 (67)
	8 - 10	1.5 (1.4)	1.1 (0.4)	1.39 (3.65)	29 (33)
$\phi$	0 - 2	89 (0.9)	7487 (36)	0.01 (0.02)	32 (4.7)
	2 - 4	106 (43)	2213 (450)	0.05 (0.10)	69 (62)
	4 - 6	19 (15)	85 (36)	0.22 (0.42)	58 (67)
	6 - 8	2.9 (2.5)	8.1 (3.1)	0.36 (0.80)	28 (34)
	8 - 10	0.8 (0.7)	1.4 (0.5)	0.54 (1.51)	16 (21)