

Medium-Modified Fragmentation Functions

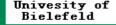
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Outline

- Medium induced gluon radiation
- Medium modified splitting functions
- Medium modified Sudakov form factor
 - Medium modified fragmentation functions
- Nuclear modification factor

Conclusions





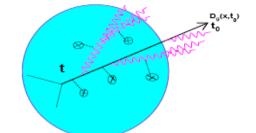
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Medium induced gluon radiation

- A hard parton produced in the early stage of Heavy Ion Collisions travels through the hot and dense QCD matter. The scattering centers induce successive gluon radiations.
- The hard parton looses virtuality from the initial scale t to the final hadronization one $t_0 \simeq \lambda_{QCD}^2$ Hadronization happens outside the medium.



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Medium induced gluon radiation is the standard explanation for Jet Quenching observed at RHIC.









The single inclusive distribution of medium induce gluons with energy ω and transverse momentum k_t from a parent parton of energy E:

$$\begin{split} &\omega \frac{dI}{d\omega dk_T} = \frac{\alpha_S C_R}{(2\pi)^2 \omega^2} 2Re \int_0^\infty dy_l \int_{y_l}^\infty d\overline{y_l} du e^{-i_T u} e^{-\frac{1}{2} \int_{\overline{y_l}}^\infty d\zeta \mathbf{n}(\boldsymbol{\zeta}) \sigma(\boldsymbol{u})} \\ &\frac{\partial}{\partial y} \frac{\partial}{\partial u} \int_{y=0=r(y_l)}^{u=r(\overline{y_l})} \mathcal{D} r e^{i \int_{y_l}^{\overline{y_l}} d\zeta \frac{\omega}{2} (r^2 - \frac{\mathbf{n}(\boldsymbol{\zeta}) \sigma(r)}{i\omega})} \end{split}$$

Two approximations

- Opacity expansion, in powers of $n(\zeta)\sigma(\zeta)$.
- ✓ multiple soft scattering: $n(\zeta)\sigma(r) \simeq \frac{1}{2}\hat{q}(\zeta)r^2$, the path integral is one of a harmonic oscilator.
 - $\hat{q}(\zeta)$ is the transport coefficient, $< q_T^2 > /\lambda$

The
$$\omega \frac{dI}{d\omega dk_T} \longrightarrow F(\frac{\omega}{\omega_c}, \kappa^2), \ \omega_C = \frac{1}{2}\hat{q}L^2, \ \kappa^2 = \frac{k_T^2}{\hat{q}L}$$



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Previous MMFF calculations

 A Poissonian distribution of independent radiations was assumed by BDMPS

 $P_E(\epsilon) = \sum_{n=0}^{\infty} \frac{1}{n!} \left[\prod_{i=1}^n \int d\omega_i \frac{d((\omega_i))}{d\omega} \right] \delta(\epsilon - \sum_{i=1}^n \frac{\omega_i}{E}) e^{-\int d\omega \frac{dI}{d\omega}}$

- The MMFF were calculated shifting the vacuum ones $D_{kh}^{(med)}(z,Q^2) = \int_0^1 d\epsilon P_E(\epsilon) \frac{1}{1-\epsilon} D_k(\frac{z}{1-\epsilon},Q^2)$
 - Limitations
 - The energy and momentum are not conserved
 - There is no evolution in virtuality
 The medium and veguum are treated of
 - The medium and vacuum are treated differently









Now Medium modified splitting functions

The total medium-induced gluon radiation spectrum:

$$\frac{dI}{d\omega \, d\mathbf{k}_{\perp}} = \omega \frac{dI^{\text{vac}}}{d\omega \, d\mathbf{k}_{\perp}} + \omega \frac{dI^{\text{med}}}{d\omega \, d\mathbf{k}_{\perp}}$$

• The vacuum case $\frac{dI^{vac}}{dzdk_T^2} = \frac{\alpha_s P(z)_{z \longrightarrow 1}^{vac}}{2\pi k_T^2}$, $P(z)_{z \longrightarrow 1}^{vac} \simeq \frac{2C_R}{1-z}$, z = 1 - x

 The ansatz is an extension of the former vacuum one to medium case: Salgado and Polosa (hep-ph/0607295)

$$\frac{dI}{dzdk_T^2} \overset{\dot{M}ED}{=} \frac{\alpha_s P(z)_{z \longrightarrow 1}^{MED}}{2\pi k_T^2} , \ P(z)_{z \longrightarrow 1}^{med} = \frac{2\pi zt}{\hat{q}L} F(\frac{\omega}{\omega_c}, \kappa^2)$$



The total splitting distribution is assumed to be the vacuum+medium ones: $P^{TOTAL}(z) = P^{VACUUM}(z) + P^{MEDIUM}(z)$



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6



Sudakov form factor

• The sudakov factor

$$\Delta_a(t, t_0^a) = e^{-\sum_{a-cc'} \int_{t_0^a}^t \frac{dt'}{t'} \int_{z_{min}(t')}^{1-z_{min}(t')} dz \frac{\alpha_S(t', z)}{2\pi} P_{ca}(z)}$$

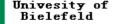
- $\Delta_a(t, t_0^a)$ means the probability for a parton not to branch while evolving from an initial virtuality t to a final scale t_0
 - We modify the sudakov factor:

$$\Delta_a(t, t_0^a) = e^{-\sum_{a-cc'} \int_{t_0^a}^t \frac{dt'}{t'} \int_{z_{min}(t')}^{1-z_{min}(t')} dz \frac{\alpha_S(t', z)}{2\pi} (P_{ca}(z)^{vac} + P_{ca}(z)^{med})}$$

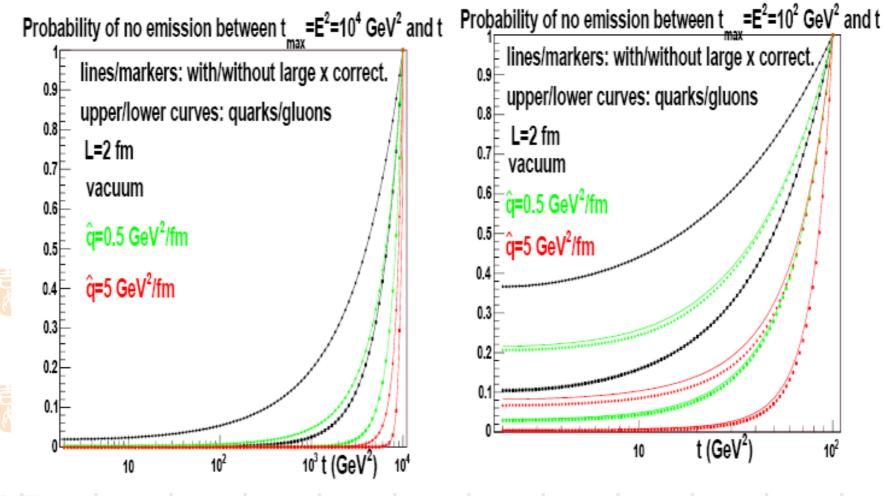




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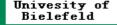


Medium modified sudakov factor





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DGLAP evolution equation can be written in terms of the sudakovs:

 $t\frac{\partial}{\partial t}\left(\frac{D_a^h(x,t)}{\Delta_a(t,t_0^a)}\right) = \int_x^{1-z_{min}(t)} \frac{dz}{z} \frac{\alpha_S(k_T^2,z)}{2\pi} P_{ba}(z) \frac{D_b^h(\frac{x}{z},t)}{\Delta_a(t,t_0^a)}$

- The renormalization scale is $t(1-z)z = k_T^2$
- The initial values for the FF we take the KKP ones at virtuality t



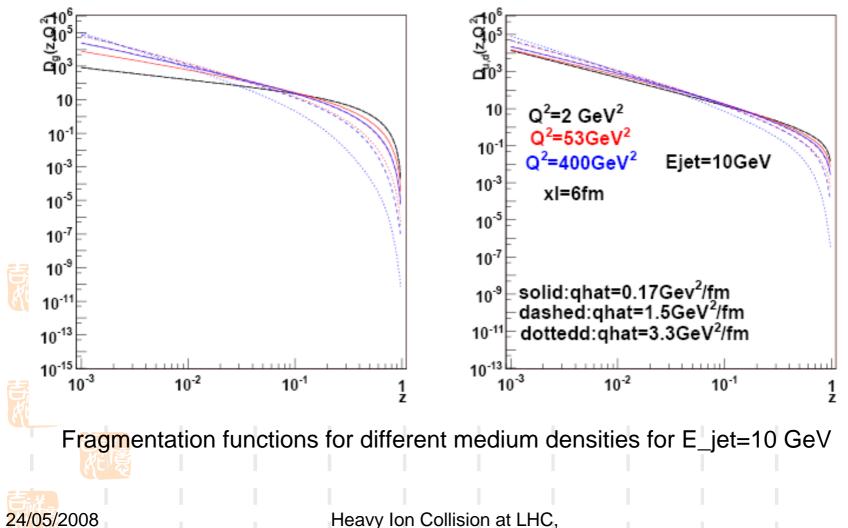
For each parton energy, $t_0 = 2GeV^2 t_0 < t < 4E^2$ and $t_0/t < z(t) < 1 - t_0/t$



Our evolution depends on the initial parton energy through the scale range in the Sudakov.

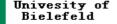


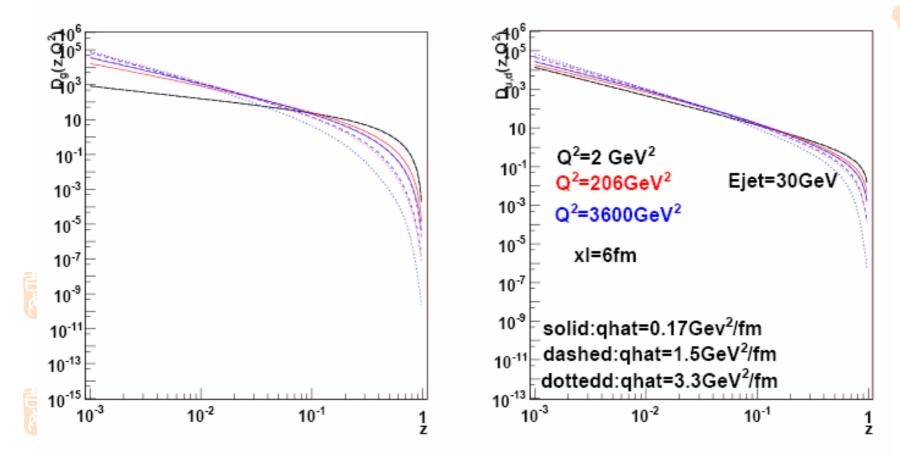
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10



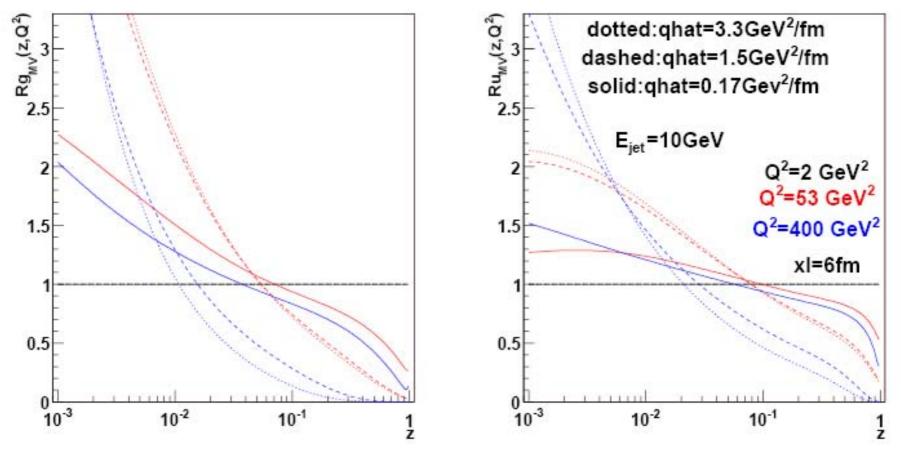


Fragmentation functions for different medium densities for E_jet=30 GeV



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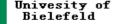


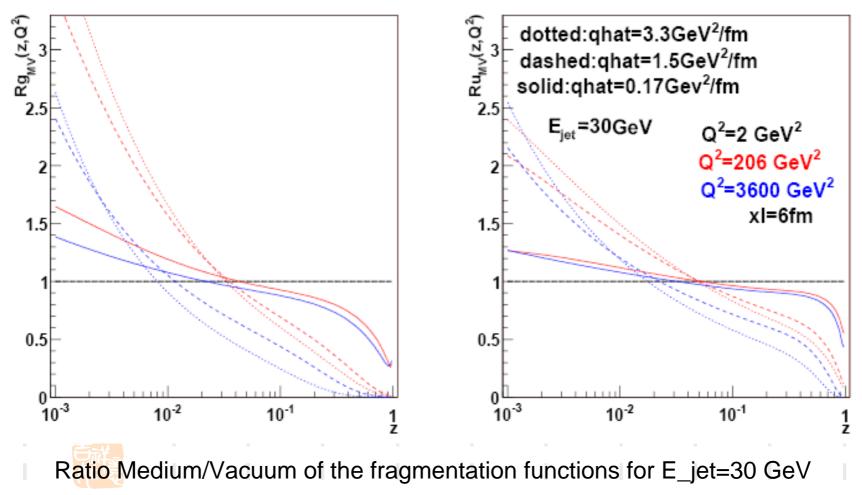


Ratio Medium/Vacuum of the fragmentation functions for E_jet=10 GeV



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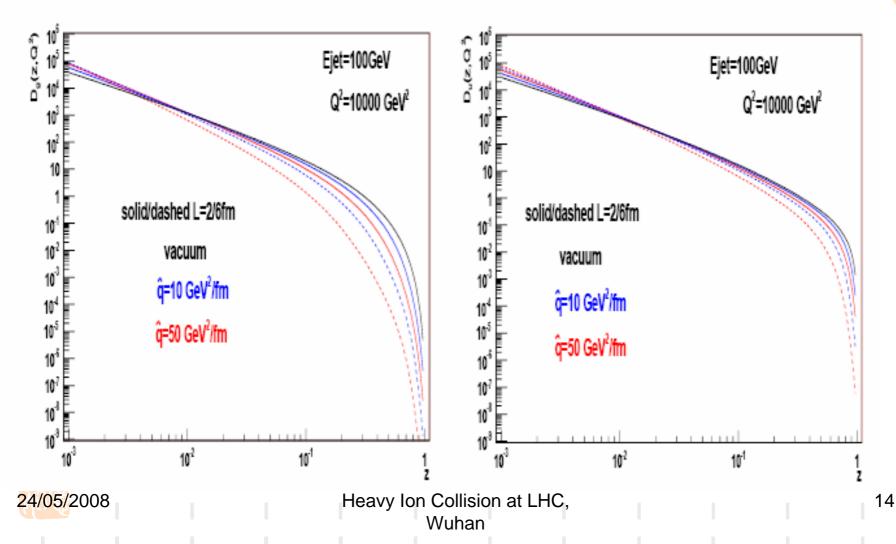




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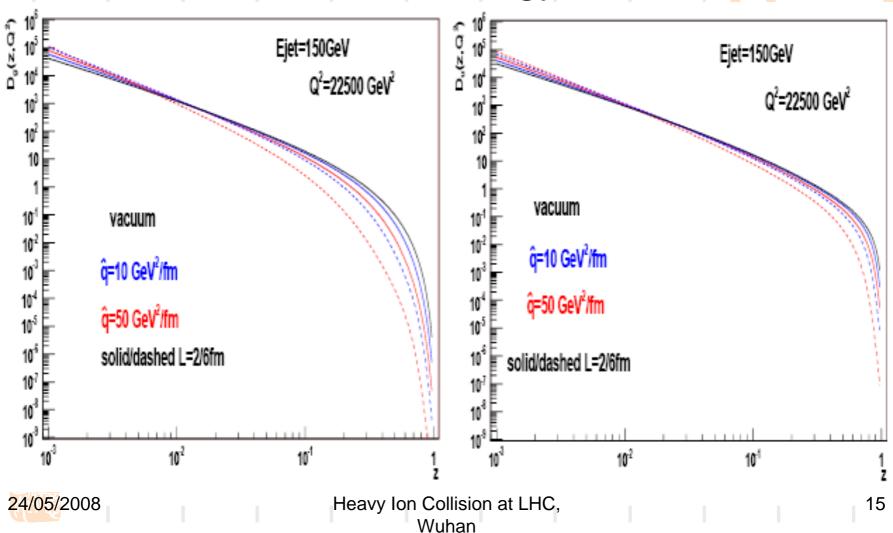


Medium modified fragmentation functions: LHC energy





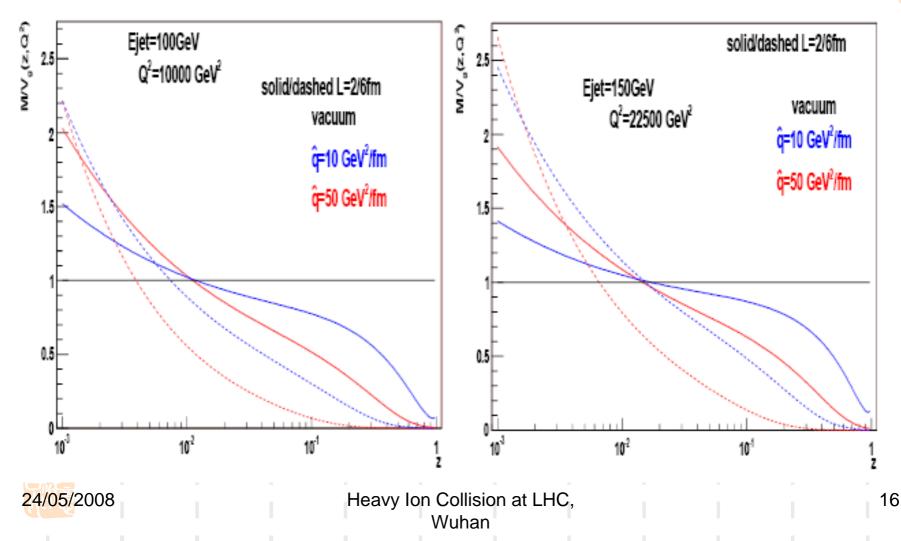
Medium modified fragmentation functions: LHC energy



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Medium modified fragmentation functions: LHC energy



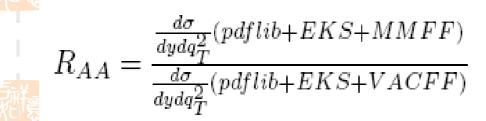


Particle production

A typical hard cross section can be written in the form:

 $\sigma^{ABh} = f_A(x1, Q^2) f_B(x2, Q^2) \otimes \sigma(x1, x2, Q^2) \otimes D_{i \longrightarrow h}(z, Q^2)$

• We define the nuclear modification factor as:



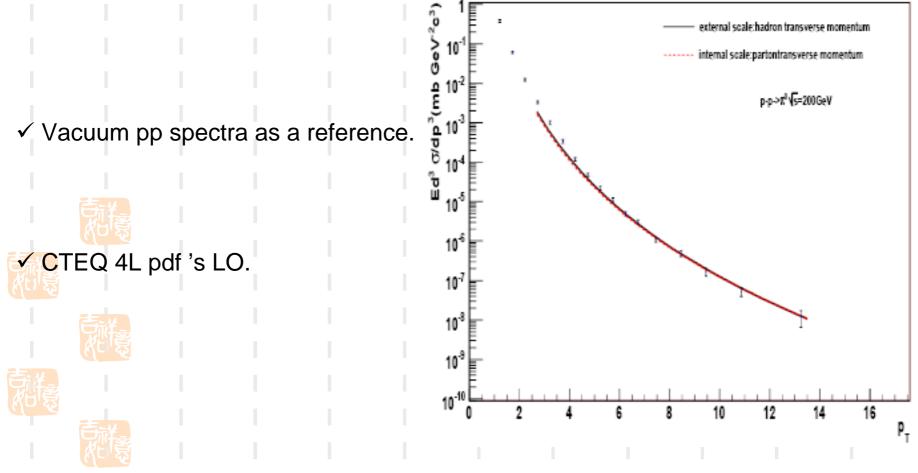


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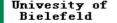


PP reference

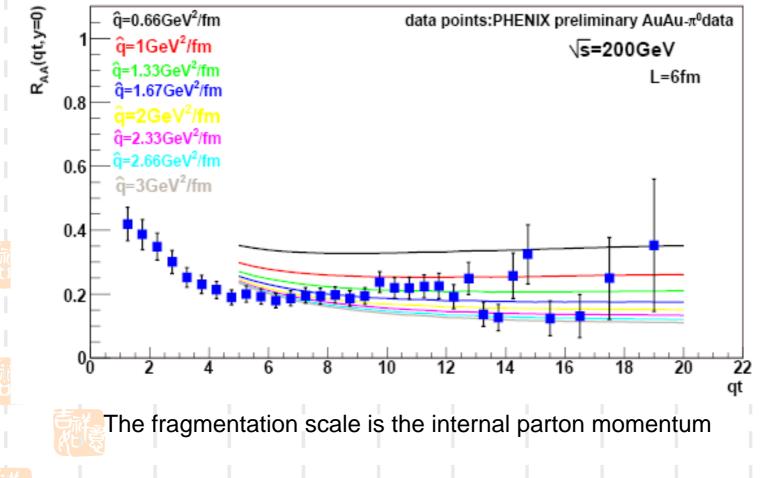




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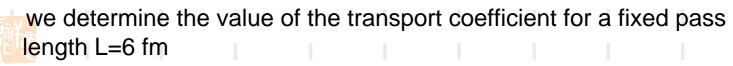
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- Sudakov form factors in medium
- Medium modified DGLAP evolution via Sudakov factor
- Some phenomenological applications:

 - Comparison to experimental data: Nuclear modification factor



$$\implies \hat{q} \simeq 1 GeV^2/fm$$



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