Naturalness at the 100 TeV scale

Nathaniel Craig UC Santa Barbara

FCC Week 2015



We should venture to higher energies/shorter distances purely for the sake of exploration. We should venture to higher energies/shorter distances purely for the sake of exploration.

However...

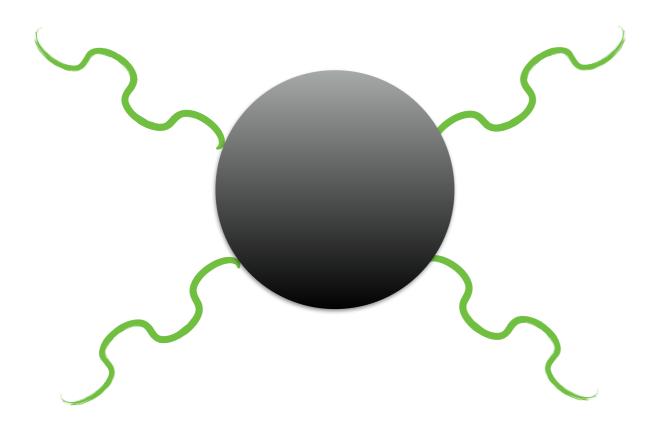
We should venture to higher energies/shorter distances purely for the sake of exploration.

However...

It is extremely useful to have specific motivation to shape the goals of future colliders.

Pre-Higgs Guarantee

Planning for colliders was simple following the discovery of W/Z bosons:



Consistency gave us a no-lose theorem.

Violation of unitarity in VV scattering

Standard Model breaks down parametrically near weak scale

Consistency demands new physics (e.g. Higgs mechanism)

Higgs discovery @ LHC

Post-Higgs Ambiguity

10¹⁸ GeV Quantum gravity?

10¹⁶ GeV Grand Unification?

The Standard Model with an elementary Higgs boson is consistent to energies far beyond the weak scale.

10¹² GeV Strong CP?

There are suggestive indications of interesting physics at various scales, but these need not be accessible at colliders.

We face an absence of no-lose theorems — but in their place we have *compelling strategies*.

10² GeV Weak scale

The naturalness strategy

In the SM, m_h is a parameter: not predicted, and worse, incalculable (elementary scalars are special).

In a theory where m_h is calculable, new physics beyond the SM enters at a scale Λ .

We see a *hierarchy problem*: quantum contributions to m_h are *at least* around this scale Λ .

$$\delta m_h \propto \Lambda$$

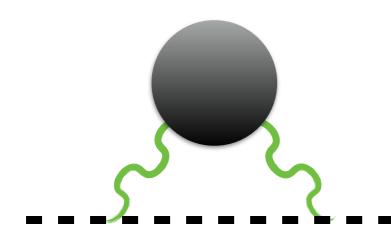
Natural if $\delta m_h \sim m_h$. ($\delta m_h \gg m_h$ unnatural or UV miracle)

The naturalness strategy

This is a *strategy* for new physics near m_h, not a *no-lose theorem*, because the theory does not break down if it is unnatural.

But naturalness has often been a very *successful* strategy.

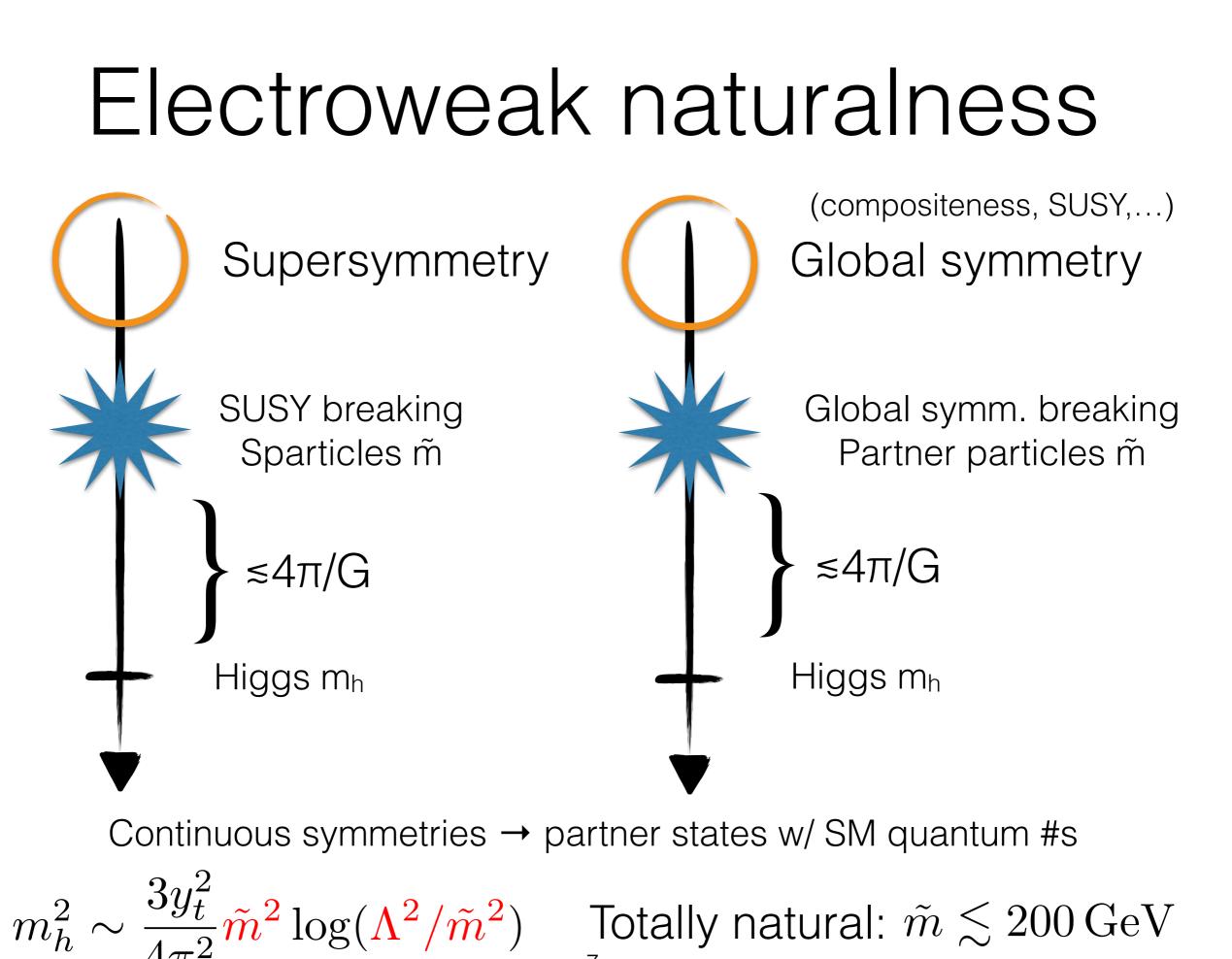
E.g. charged pions



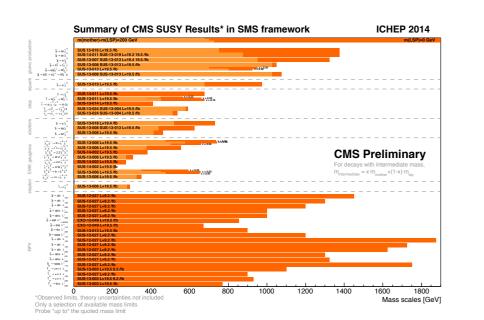
Electromagnetic contribution to the charged pion mass sensitive to the cutoff of the pion EFT.

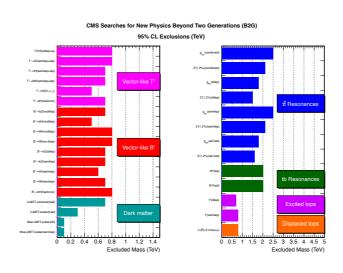
$$\delta m^2 \sim \frac{3e^2}{16\pi^2} \Lambda^2$$

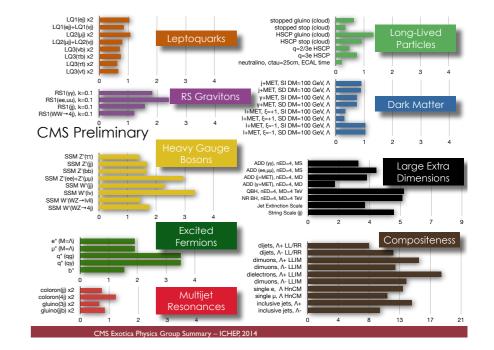
Naturalness suggests Λ ~850 MeV. Rho meson (new physics!) enters at 770 MeV.



A physics driver @ LHC

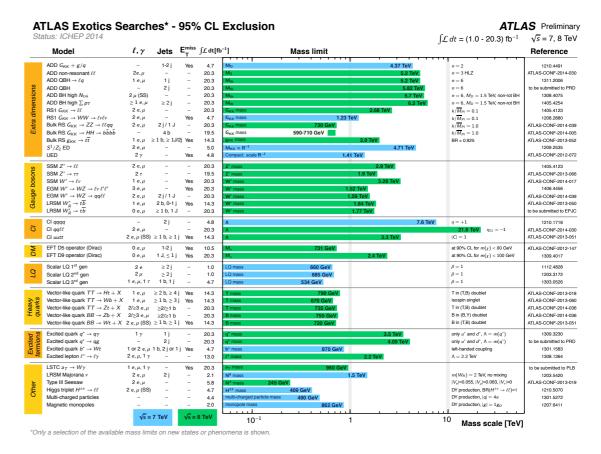




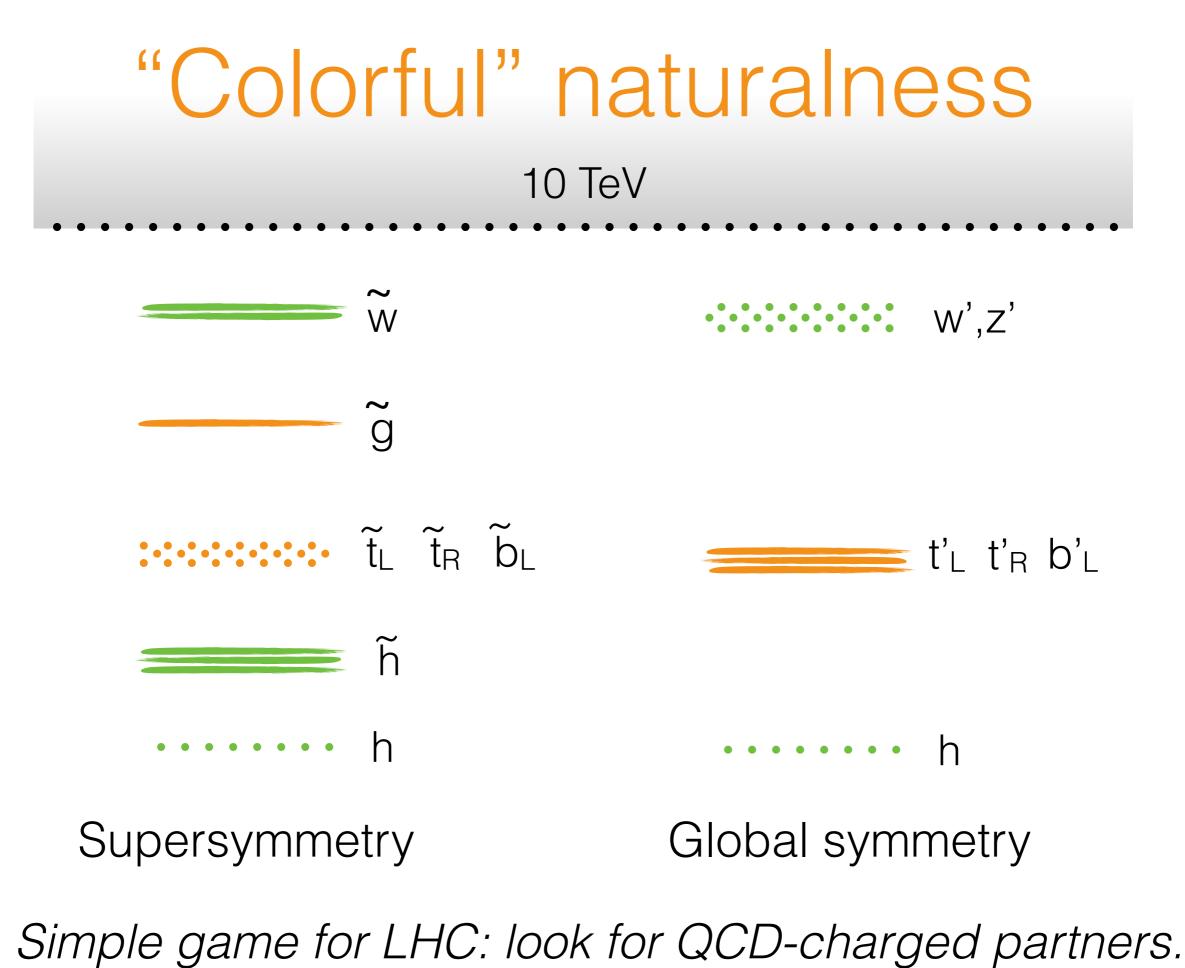


170 of these 226 channels tied to naturalness

Model	e, μ, τ, γ	Jets	$E_{\rm T}^{\rm miss}$	$\int \mathcal{L} dt [\mathbf{fb}]$	¹] Mass limit		Reference
$ \begin{array}{c} \text{MSUGRA/CMSSM} \\ \bar{q}\bar{q}, \bar{q} \rightarrow q \tilde{\ell}_{1}^{0} \\ \bar{q}, \bar{q} \rightarrow q \tilde{\ell}_{1}^{0} \\ \bar{z}, \bar{z}, \bar{q} \rightarrow q \tilde{\ell}_{1}^{0} \\ \bar{z}, $	$\begin{matrix} 0 \\ 0 \\ 1 \\ \gamma \\ 0 \\ 1 \\ e, \mu \\ 2 \\ e, \mu \\ 1.2 \\ \tau + 0.1 \\ \ell \\ 2 \\ \gamma \\ 1 \\ e, \mu + \gamma \\ \gamma \\ 2 \\ e, \mu (Z) \end{matrix}$	2-6 jets 2-6 jets 0-1 jet 2-6 jets 3-6 jets 0-3 jets 0-2 jets - 1 b 0-3 jets	Yes Yes Yes Yes Yes Yes Yes Yes Yes	20.3 20.3 20.3 20.3 20 20 20.3 20.3 4.8 4.8 5.8	4.2 1.7 TeV m(4 850 GeV m(7 250 GeV m(8 1.33 TeV m(2 1.32 TeV m(2 1.32 TeV m(8 1.52 TeV m(8 1.5 TeV m(8 619 GeV m(8 619 GeV m(90 GeV m(m($ \bar{q} - m(\bar{z})$ $ \bar{t}_1^{(2)} - 0 \cdot G\bar{c}^2 V, m([1^{40} \text{ gen.} \bar{q}] - m([2^{40} \text{ gen.} \bar{q}])$ $\bar{q} - m(\bar{c})^2, \bar{t}_1^{(2)} - 0 \cdot G\bar{c} V, m(\bar{c}^2) - 0.5(m(\bar{c}_1^{2}) + m(\bar{z}));$ $ \bar{t}_1^{(2)} - 0 \cdot G\bar{c} V, \bar{t}_1^{(2)} - 0.5(m(\bar{c}_1^{2}) + m(\bar{z}));$ $\bar{t}_1^{(2)} - 50 \cdot G\bar{c} V, \bar{t}_1^{(2)} - 50 \cdot G\bar{c} V,$ $ \bar{t}_1^{(2)} - 50 \cdot G\bar{c} V$ $ \bar{t}_1^{(2)} - 50 \cdot G\bar{c} V$	1405.7875 1405.7875 1411.1559 1405.7875 1501.03555 1501.03555 1407.0603 ATLAS-CONF-2014-001 ATLAS-CONF-2012-142 211.1167 ATLAS-CONF-2012-152
Gravitino LSP $\vec{g} \rightarrow b \vec{b} \vec{k}_{1}^{0}$ $\vec{g} \rightarrow t \vec{k}_{1}^{0}$ $\vec{g} \rightarrow t \vec{k}_{1}^{0}$ $\vec{g} \rightarrow t \vec{k}_{1}^{0}$ $\vec{g} \rightarrow t \vec{k}_{1}^{0}$	0 0 0-1 e, µ 0-1 e, µ	3 b 7-10 jets 3 b 3 b 3 b	Yes Yes Yes Yes Yes	20.3 20.1 20.3 20.1 20.1	ž 1.25 TeV π.(ž 1.1 TeV π.(ž 1.34 TeV π.((G)>1.8 × 10 ⁻⁴ eV, m(g)=m(q)=1.5 TeV [ξ ⁰] <400 GeV [ξ ⁰] <500 GeV [ξ ⁰] <500 GeV [ξ ⁰] <400 GeV [ξ ⁰] <500 GeV	1502.01518 1407.0600 1308.1841 1407.0600 1407.0600
$\begin{array}{c} \bar{b}_1 \bar{b}_1, \bar{b}_1 \rightarrow b \bar{k}_1^0 \\ \bar{b}_1 \bar{b}_1, \bar{b}_1 \rightarrow b \bar{k}_1^0 \\ \bar{b}_1 \bar{b}_1, \bar{b}_1 \rightarrow b \bar{k}_1^0 \\ \bar{f}_1 \bar{f}_1, \bar{f}_1 \rightarrow b \bar{k}_1^0 \\ \bar{f}_1 \bar{f}_1 \bar{f}_1 \bar{f}_1 - b \bar{k}_1^0 \bar{f}_1 \\ \bar{f}_1 \bar{f}_1 \bar{f}_1 - b \bar{k}_1^0 \bar{f}_1 \\ \bar{f}_1 \bar{f}_1 \bar{f}_1 \bar{f}_1 - b \bar{k}_1^0 \bar{f}_1 \\ \bar{f}_1 \bar{f}_1 \bar{f}_1 \bar{f}_1 - b \bar{k}_1^0 \bar{f}_1 \\ \bar{f}_1 \bar{f}_1 \bar{f}_1 \bar{f}_1 \bar{f}_1 - b \bar{k}_1^0 \bar{f}_1 \\ \bar{f}_1 \bar{f}_1 \bar{f}_1 \bar{f}_1 - b \bar{k}_1^0 \bar{f}_1 \\ \bar{f}_1 \bar{f}_1 \bar{f}_1 \bar{f}_1 \bar{f}_1 - b \bar{k}_1 \bar{f}_1 \\ \bar{f}_1 \bar{f}_1 \bar{f}_1 \bar{f}_1 \bar{f}_1 - b \bar{k}_1 \bar{f}_1 \\ \bar{f}_1 \bar{f}_1 \bar{f}_1 \bar{f}_1 \bar{f}_1 \bar{f}_1 \bar{f}_1 \bar{f}_1 \bar{f}_1 - b \bar{k}_1 \bar{f}_1 \\ f$	0 2 e, μ (SS) 1-2 e, μ 2 e, μ 0-1 e, μ 0 m 2 e, μ (Z) 3 e, μ (Z)	2 b 0-3 b 1-2 b 0-2 jets 1-2 b ono-jet/c-t 1 b 1 b	Yes Yes Yes Yes Yes ag Yes Yes Yes	20.1 20.3 4.7 20.3 20.3 20.3 20.3 20.3	δ1 275-440 GeV m() f1 110-167 GeV 230-460 GeV m() f2 90-191 GeV 215-530 GeV m() f1 90-240 GeV m() m() f1 90-240 GeV m() m()	[t])<90 GeV [t])=2 m(t]) [t])=2 m(t]) [t])=55 GeV [t])=1 GeV [t])=1 GeV [t])=156 GeV [t])>150 GeV [t])<250 GeV	1308.2631 1404.2500 1209.2102, 1407.0583 1403.4853, 1412.4742 1407.0583,1406.1122 1407.0608 1403.5222 1403.5222
$\begin{array}{c} \tilde{\ell}_{1,\mathbf{x}}\tilde{\ell}_{1,\mathbf{x}}, \tilde{\ell} \rightarrow \tilde{\ell}\tilde{\chi}_{1}^{0} \\ \tilde{\chi}_{1}^{\dagger}\tilde{\chi}_{1}^{-}, \tilde{\chi}_{1}^{\dagger} \rightarrow \tilde{\ell}_{1} (\tilde{r}) \\ \tilde{\chi}_{1}^{\dagger}\tilde{\chi}_{1}^{-}, \tilde{\chi}_{1}^{\dagger} \rightarrow \tilde{r} \tau (\tilde{r}) \\ \tilde{\chi}_{1}^{\dagger}\tilde{\chi}_{1}^{-}, \tilde{\chi}_{1}^{\dagger} \rightarrow \tilde{r} \tau (\tilde{r}) \\ \tilde{\chi}_{1}^{\dagger}\tilde{\chi}_{2}^{-} \rightarrow W_{1}^{0} 2\tilde{\chi}_{1}^{0} \\ \tilde{\chi}_{1}^{\dagger}\tilde{\chi}_{2}^{0} \rightarrow W_{1}^{0} 2\tilde{\chi}_{1}^{0} \\ \tilde{\chi}_{1}^{\dagger}\tilde{\chi}_{2}^{0} \rightarrow W_{1}^{0} \tilde{\chi}_{1}^{0} \\ \tilde{\chi}_{2}^{\dagger}\tilde{\chi}_{1}^{-}, \tilde{\chi}_{2}^{0} \tilde{\chi}_{1}^{-} \tilde{\chi}_{2}^{0} \\ \tilde{\chi}_{2}^{\dagger}\tilde{\chi}_{1}^{-} \tilde{\chi}_{2}^{0} \tilde{\chi}_{1}^{-} \tilde{\chi}_{2}^{0} \\ \tilde{\chi}_{2}^{\dagger}\tilde{\chi}_{1}^{-} \tilde{\chi}_{2}^{0} \\ \tilde{\chi}_{1}^{\dagger}\tilde{\chi}_{2}^{0} \rightarrow W_{1}^{0} \tilde{\chi}_{1}^{0} \\ \tilde{\chi}_{2}^{0} \\ \tilde{\chi}_{1}^{-} \tilde{\chi}_{2}^{0} \\ \tilde{\chi}_{1}^{-} \tilde{\chi}_{2}^{0} \\ \tilde{\chi}_{1}^{0} \\ \tilde{\chi}_{1}$	2 e,μ 2 e,μ 2 τ 3 e,μ 2-3 e,μ -/γγ e,μ,γ 4 e,μ	0 0 0-2 jets 0-2 b 0	Yes Yes Yes Yes Yes Yes	20.3 20.3 20.3 20.3 20.3 20.3 20.3 20.3	រី <mark>វ 140-455 GeV</mark> ៣៣ នី <mark>វ 100-350 GeV ៣៥</mark> នីវិ ₁ .នី ¹ 700 GeV ៣៥ វីវិ ₁ .ដី នីវិ ₁ .ដី 420 GeV ៣៥	$\begin{split} &\tilde{t}_{1}^{2}){=}OGeV \\ &\tilde{t}_{1}^{2}){=}OGeV, m(\tilde{c}, \tilde{v}){=}0.5(m(\tilde{c}_{1}^{+}),m(\tilde{c}_{1}^{0}))) \\ &\tilde{t}_{1}^{2}){=}OGeV, m(\tilde{c}, \tilde{v}){=}0.5(m(\tilde{c}_{1}^{+}),m(\tilde{c}_{1}^{0}))) \\ &\tilde{t}_{1}^{2}){=}m(\tilde{c}_{1}^{2}), m(\tilde{c}_{1}^{-}){=}0.5(m(\tilde{c}_{1}^{+}),m(\tilde{c}_{1}^{0}))) \\ &\tilde{t}_{1}^{2}){=}m(\tilde{c}_{1}^{2}), m(\tilde{c}_{1}^{0}){=}0, sleptons decoupled \\ &\tilde{t}_{1}^{2}){=}m(\tilde{c}_{1}^{2}), m(\tilde{c}_{1}^{0}){=}0.5(m(\tilde{c}_{1}^{0}),m(\tilde{c}_{1}^{0})) \\ &\tilde{t}_{1}^{2}){=}m(\tilde{c}_{1}^{2}), m(\tilde{c}_{1}^{0}){=}0.5(m(\tilde{c}_{1}^{0}),m(\tilde{c}_{1}^{0})) \end{split}$	1403.5294 1403.5294 1407.0350 1402.7029 1403.5294, 1402.7029 1501.07110 1405.5086
Direct $\tilde{\chi}_{1}^{+}\tilde{\chi}_{1}^{-}$ prod., long-lived \tilde{x} Stable, stopped \bar{g} R-hadron Stable \bar{g} R-hadron GMSB, stable $\bar{\tau}, \tilde{\chi}_{1}^{0} \rightarrow \bar{\tau}(\bar{e}, \bar{\mu}) + \tau(\bar{e}, \bar{\mu}) + \tau(\bar{e}, \bar{\mu}) - \tau(\bar{e}, \bar{e}, \bar{\mu}) - \tau(\bar{e}, \bar{\mu}) - \tau(e$	0 trk	1 jet 1-5 jets - - -	Yes Yes - Yes -	20.3 27.9 19.1 19.1 20.3 20.3	832 GeV m() ã 1.27 TeV x₁² 537 GeV 10 x₁² 435 GeV 2	$[\tilde{\tau}_1^{+}] -m[\tilde{\tau}_1^{0}] = 160 \text{ MeV}, \tau(\tilde{\tau}_1^{+}) = 0.2 \text{ ns}$ $[\tilde{\tau}_1^{0}] = 100 \text{ GeV}, 10 \mu \text{s} < \tau(\tilde{g}) < 1000 \text{ s}$ $ <\text{tan}\beta < 50$ $ \text{cr}(\tilde{\tau}_1^{0}) < 3 \text{ ns}, \text{SPS8 model}$ $5 < cr<156 \text{ mm}, \text{BR}(\mu) = 1, m[\tilde{\tau}_1^{0}] = 108 \text{ GeV}$	1310.3675 1310.6584 1411.6795 1411.6795 1409.5542 ATLAS-CONF-2013-092
$ \begin{array}{l} LFV \ pp \rightarrow \overline{v}_\tau + X, \ \overline{v}_\tau \rightarrow e + \mu \\ LFV \ pp \rightarrow \overline{v}_\tau + X, \ \overline{v}_\tau \rightarrow e(\mu) + \tau \\ Bilinear \ RPV \ CMSSM \\ \vec{x}, \ \vec{x}, \ \vec{x}_1 \rightarrow WX_1^0, \ \vec{x}_1^0 \rightarrow e\overline{v}_\mu, e\mu \overline{v}, \\ \vec{x}, \ \vec{x}_1^\tau, \ \vec{x}_1^\tau \rightarrow WX_1^0, \ \vec{x}_1^0 \rightarrow \tau \tau \overline{v}_e, e\tau \overline{v}, \\ \vec{x}, \ \vec{x}, \ \vec{x}, \ \vec{x} \rightarrow WX_1^0, \ \vec{x}_1^0 \rightarrow \tau \tau \overline{v}_e, e\tau \overline{v}, \\ \vec{g} \rightarrow qq \\ \vec{g} \rightarrow \vec{t}, \ \vec{t}, \ \vec{h} \rightarrow bs \end{array} $		- 0-3 b - - 6-7 jets 0-3 b	- Yes Yes Yes - Yes	4.6 4.6 20.3 20.3 20.3 20.3 20.3 20.3	τ. 1.1 TeV X ['] ₁₁ φ. έ 1.35 TeV m(δ ¹ 750 GeV m(χ ¹ 450 GeV m($\begin{split} &_{11} = 0.10, \mathcal{A}_{112} = 0.05 \\ &_{11} = 0.10, \mathcal{A}_{122} = 0.05 \\ &_{22}^{2} = m(\tilde{g}), c_{74,2} = C mm \\ &_{71}^{2} > 0.2 c_{74} = C mm \\ &_{71}^{2} > 0.2 c_{71} = C (\tilde{g}), \mathcal{A}_{112} = 0 \\ &_{71}^{2} > 0.2 c_{71} = C (\tilde{g}), \mathcal{A}_{112} = 0 \\ &_{71}^{2} = B(k) = BR(k) = BR(k) = BR(k) = 0\% \end{split}$	1212.1272 1212.1272 1404.2500 1405.5086 1405.5086 ATLAS-CONF-2013-091 1404.250
Scalar charm, $\bar{c} \rightarrow c \bar{\chi}_1^0$ $\sqrt{s} = 7 \text{ TeV}$	0 = 8 TeV	2 c	Yes 8 TeV	20.3	2 490 GeV m((₹ ⁰)<200 GeV	1501.01325



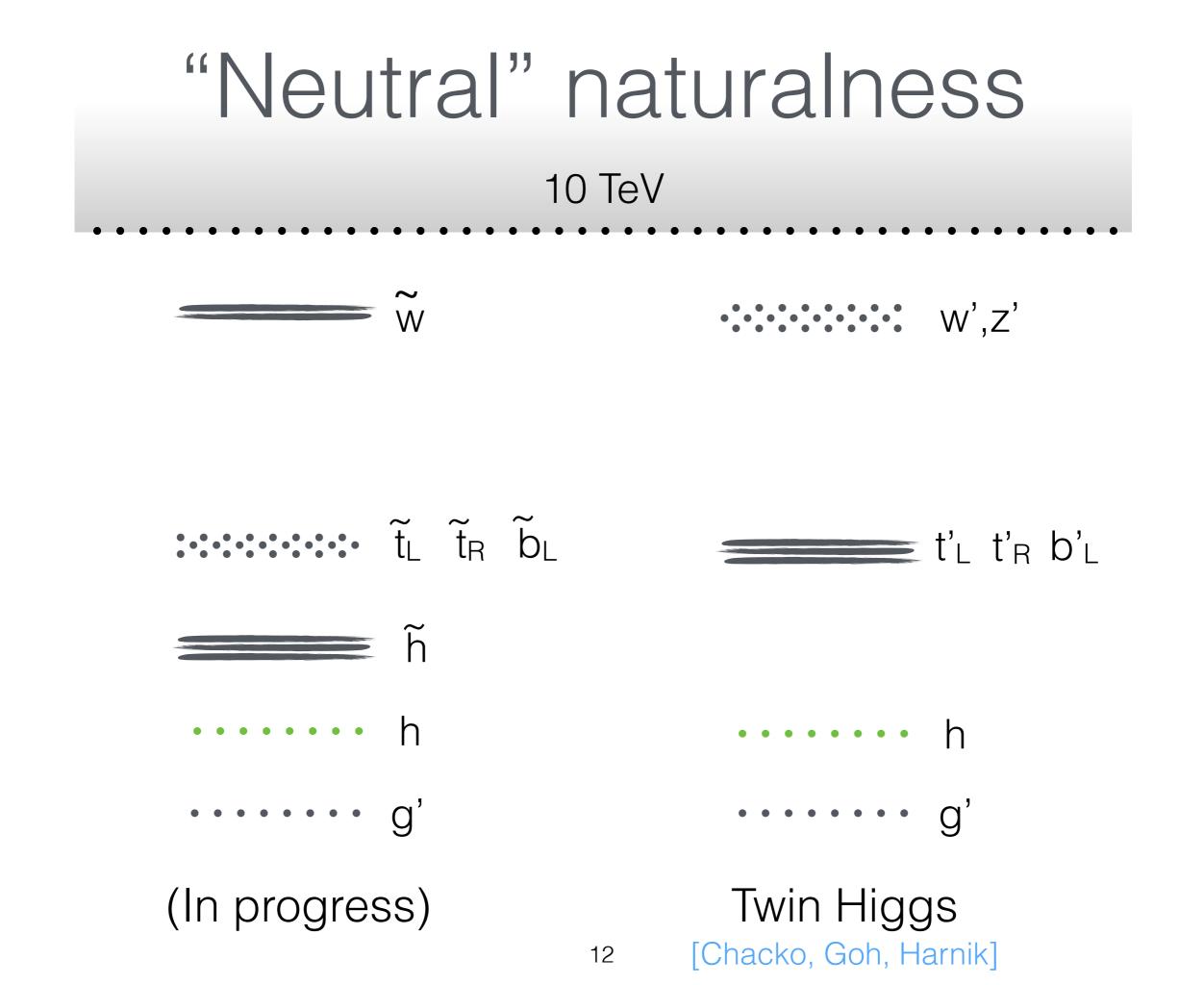
Signs of naturalness



Colorful naturalness Experimental handles

- SUSY: Direct searches (and indirect searches).
 - Look for colored partner states (stops, gluinos)
 - Look for O(loop*v/m) Higgs coupling deviations.
- Global: Direct and indirect searches.
 - Look for colored partner states (vector-like t')
 - Look for O(v/f) Higgs coupling deviations.

This is our current search program for naturalness. But...



Neutral naturalness Experimental handles

- SUSY: Direct searches (and indirect searches).
 - Look for off-shell Higgs portal.
 - Look for O(loop*v/m) Higgs coupling deviations.
 - Look for the UV completion.
- Global: Direct and indirect searches.
 - Look for O(v/f) Higgs coupling deviations.
 - Look for displaced decays [NC, Katz, Strassler, Sundrum]
 - Look for the UV completion.

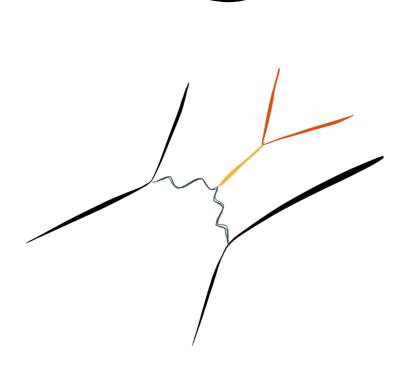
Neutral naturalness

Higgs couplings: accustomed to looking for corrections to loop-level couplings ($h \rightarrow \gamma \gamma$, gg), but even loops of neutral states can be seen.

[NC, Englert, McCullough; Henning, Lu, Murayama; NC, Farina, McCullough, Perelstein]

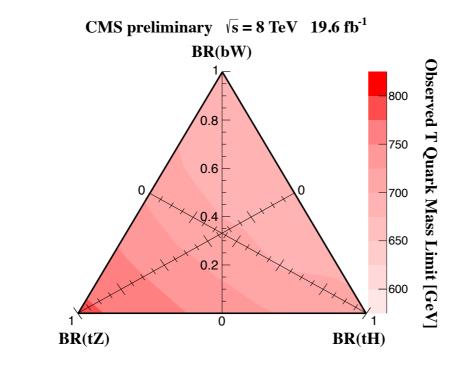
$$\frac{c_H}{m_\phi^2} \left(\partial_\mu |H|^2 \right)^2 \to \delta \sigma_{Zh} = -2c_H \frac{v^2}{m_\phi^2}$$

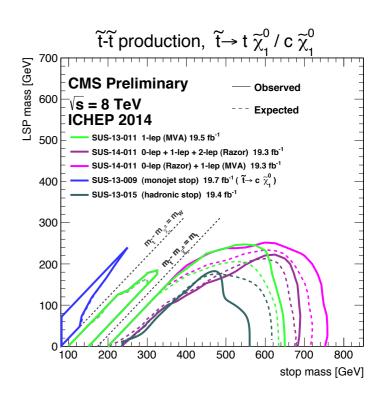
Direct searches: states lighter than m_h/2 easily constrained by Higgs width; if heavier than m_h/2, can still produce via an off-shell Higgs. Look for associated production + invisible. [Curtin, Meade, Yu; NC, Lou, McCullough, Thalapilli]

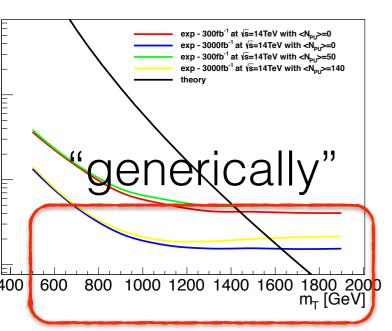


Colorful naturalness

@LHC



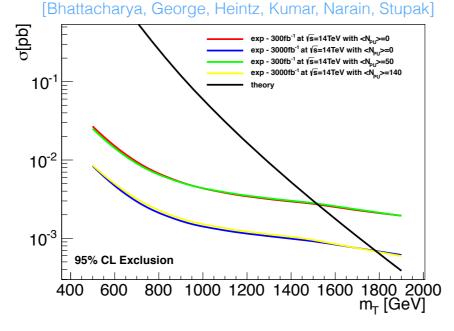


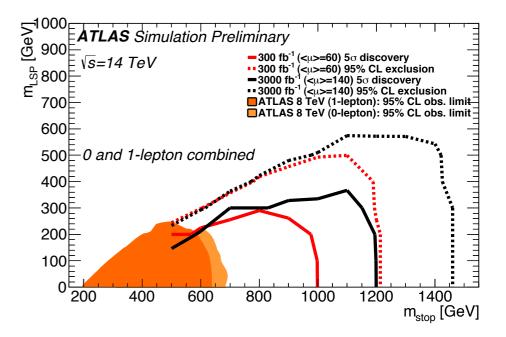


Where we are:

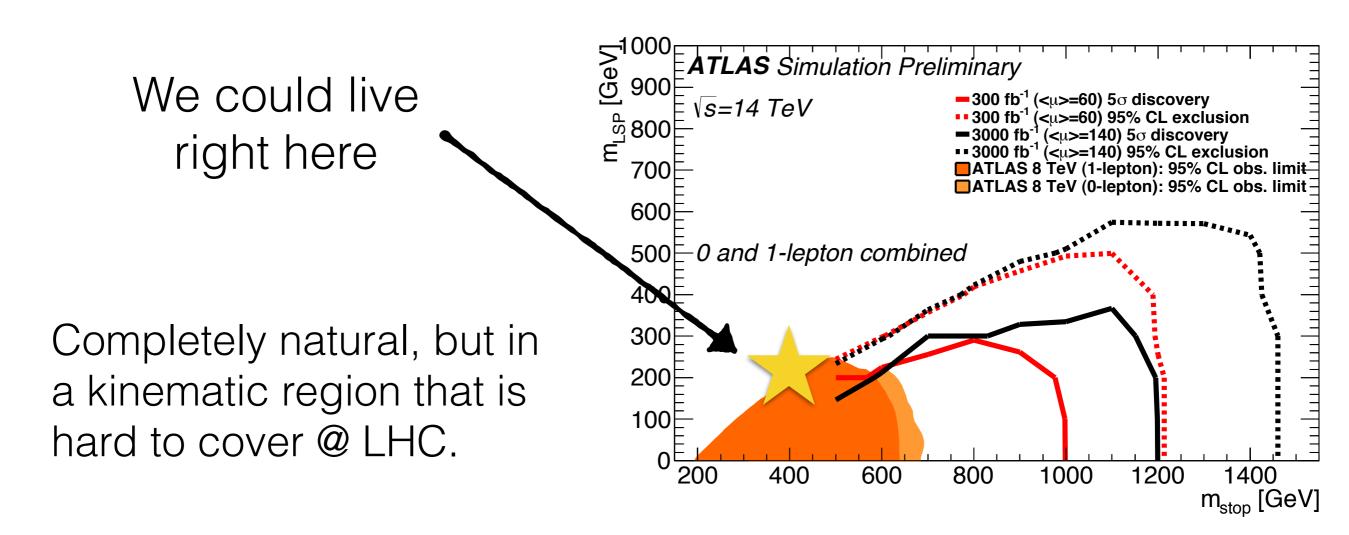
"generically"

~7% tuning level





And yet...



...or the conventional collider signals could be eroded or reduced by modest complications of the theory (RPV, stealth, etc.)

Neutral naturalness

@LHC

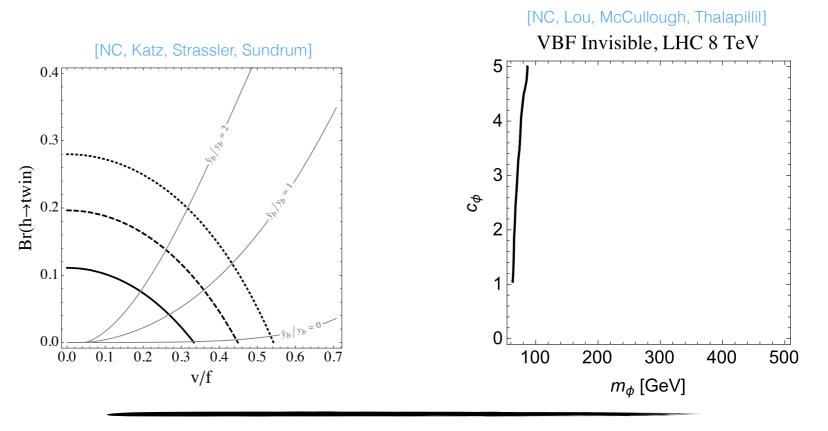


Where we are:

natural (at worst 30% for global)

Where we'll be @ end of LHC:

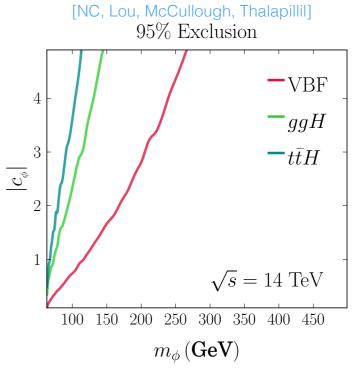
natural (at worst 20% for global)



95% CL bounds

 $\frac{v/f}{(\text{ATLAS})} \lesssim 0.31 \, (0.25)$

 $Br(inv.) \lesssim 10\%$

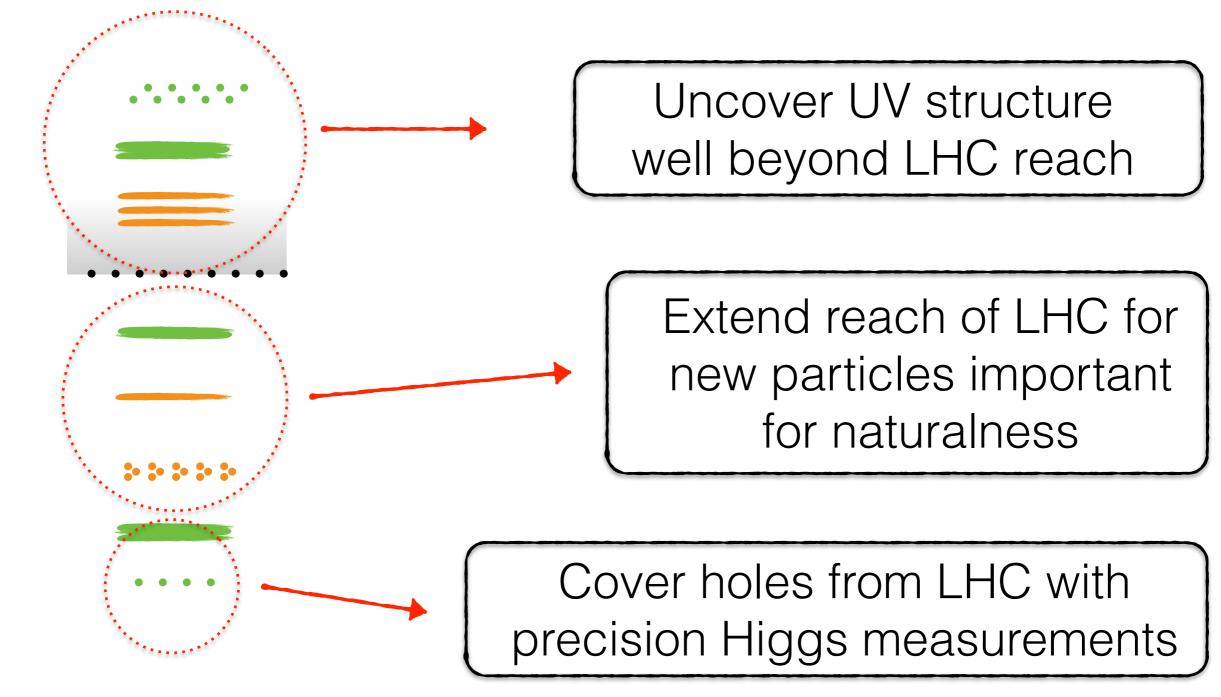


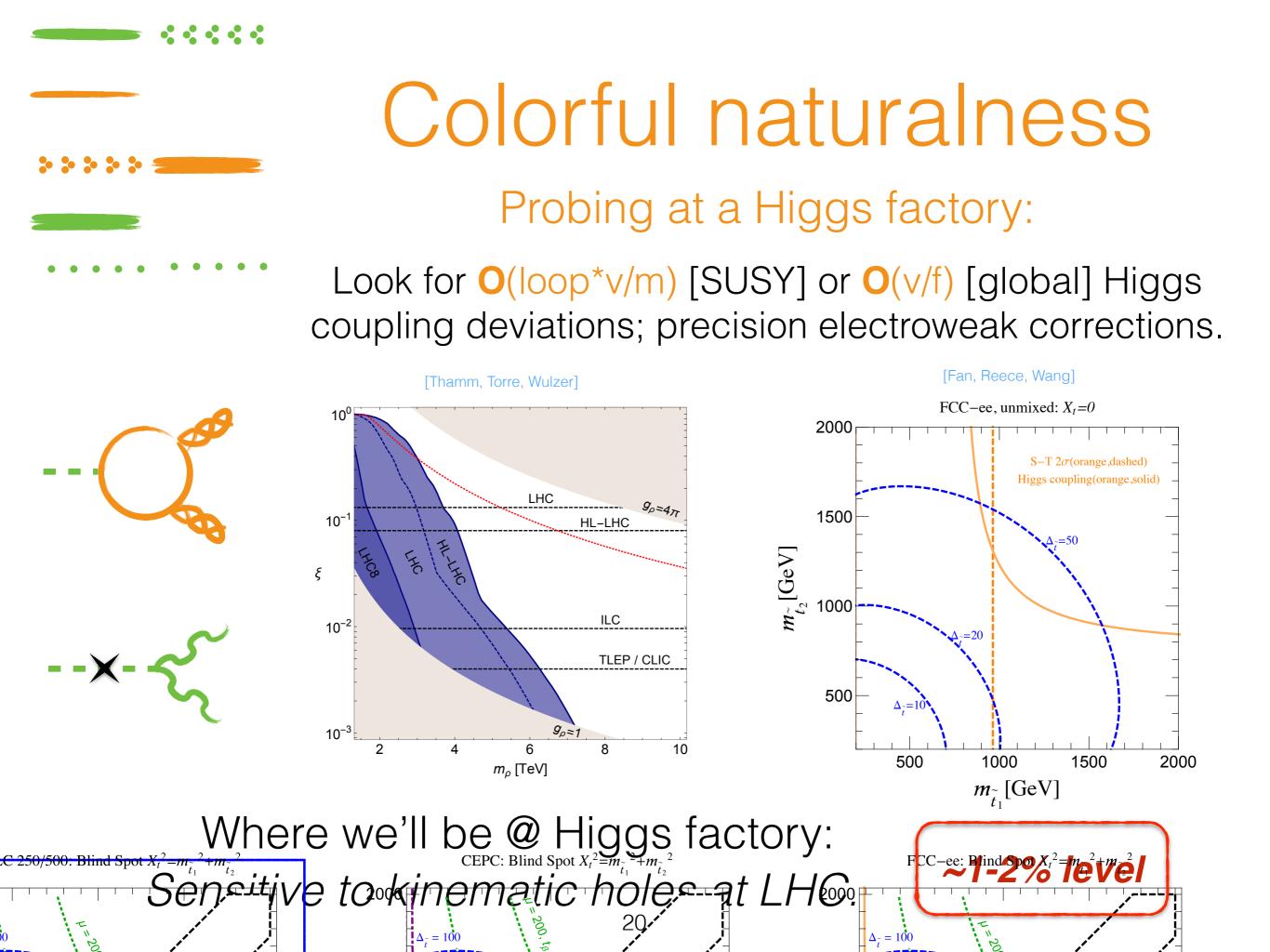
The question of electroweak naturalness cannot be settled at the LHC.

Settling the question of naturalness is a compelling strategy for future colliders.

Naturalness at future colliders

Three major opportunities



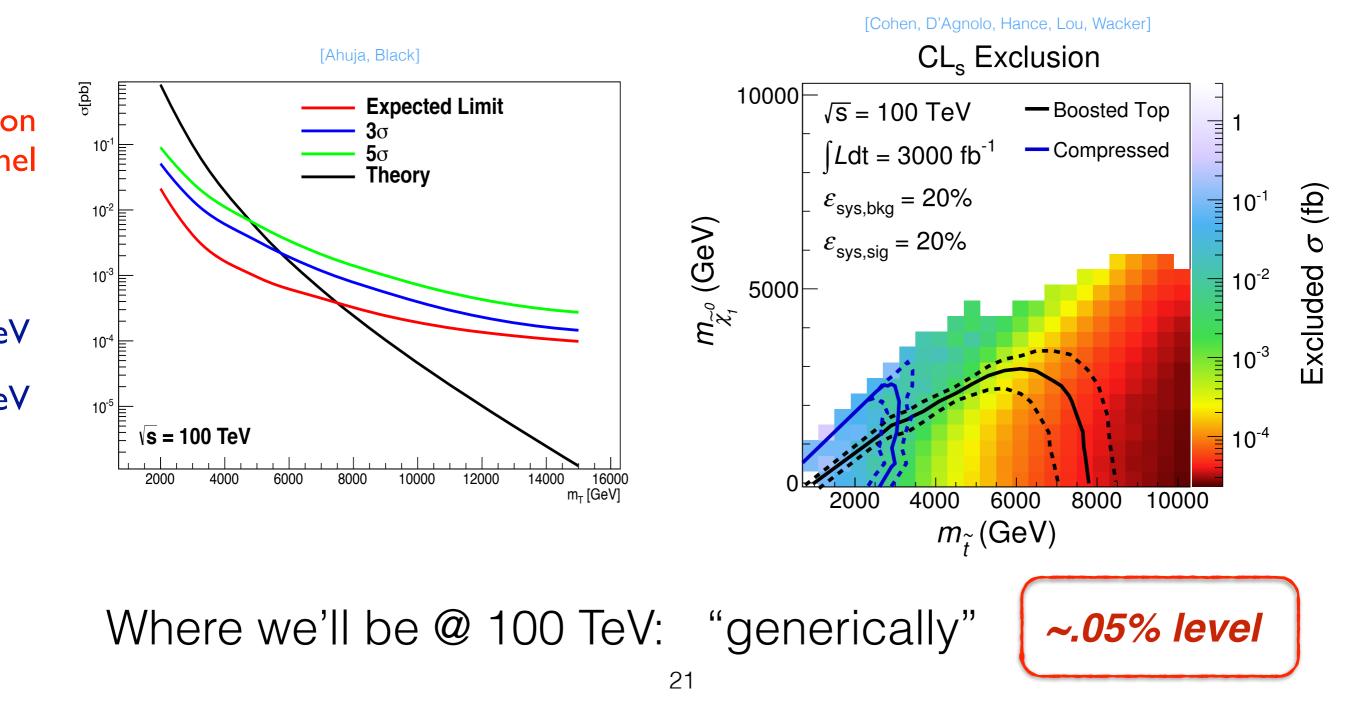




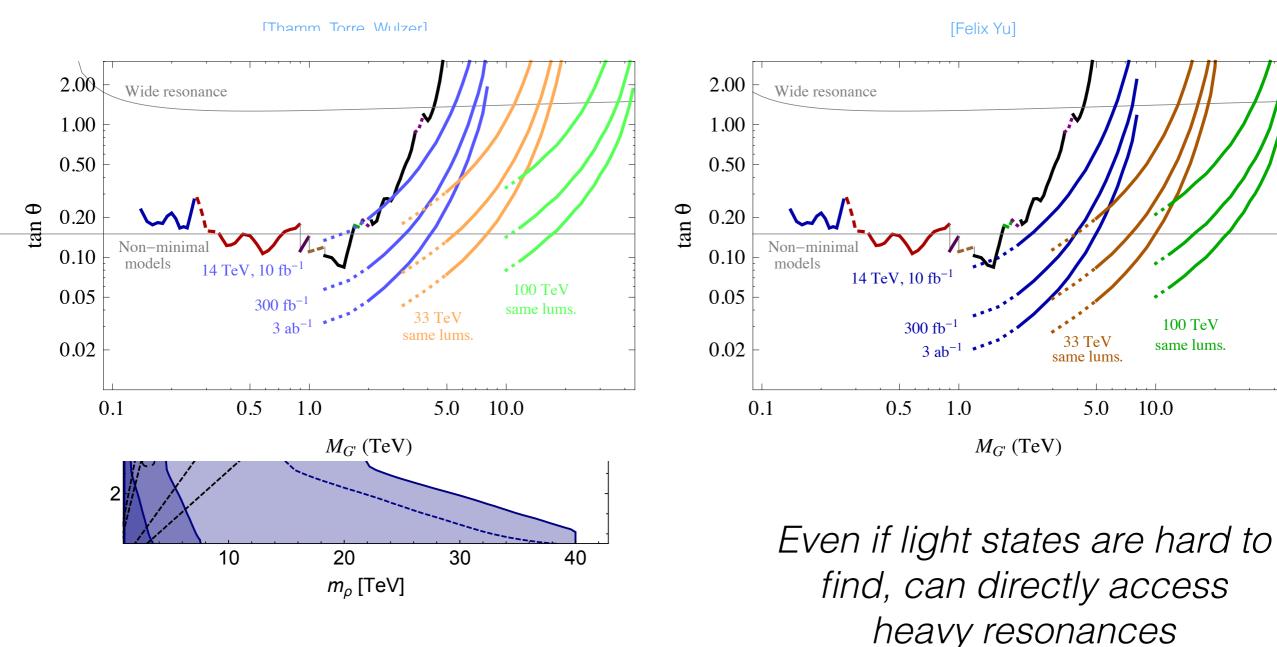
Resu

Colorful naturalness Probing at 100 TeV:

Look for the light partner states

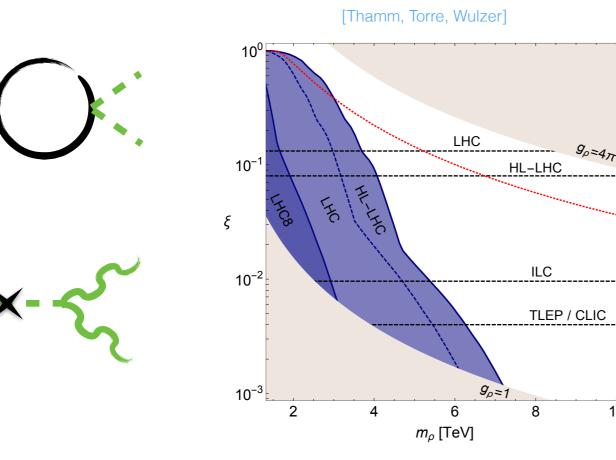


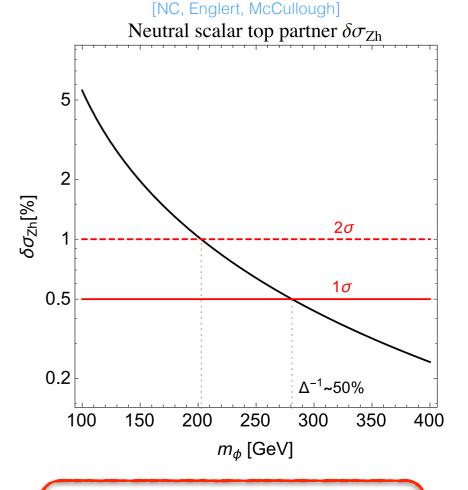




Neutral naturalness Probing at a Higgs factory:

Look for $O(loop^*v/m)$ oblique [SUSY] or O(v/f) [global] Higgs coupling deviations.





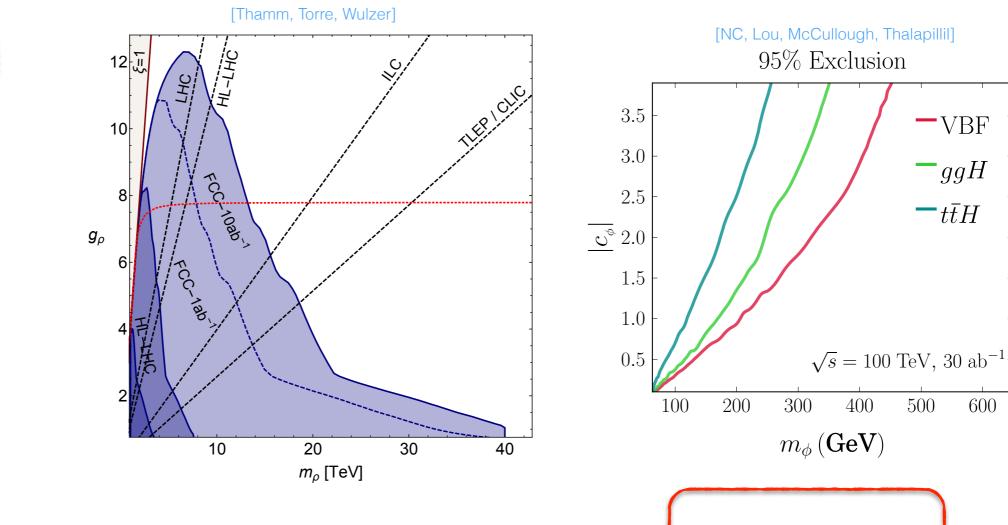
Where we'll be @ Higgs factory:

~1% level (global) ~50% level (SUSY)

10

Even if the light natural states are neutral, there are heavier states with SM charges

Neutral naturalness Probing at 100 TeV Look for the UV completion, or probe light states via the Higgs portal.



Where we'll be @ 100 TeV:



25

 Naturalness is one of the most compelling motivations for new physics near the weak scale.

 Naturalness is one of the most compelling motivations for new physics near the weak scale.

 The LHC will eventually probe conventional "colorful" theories to (at best) the ~1% level.

 Naturalness is one of the most compelling motivations for new physics near the weak scale.

 The LHC will eventually probe conventional "colorful" theories to (at best) the ~1% level.

But it will leave kinematic regions in conventional theories

 and all regions of more novel theories — essentially untested, and the status of naturalness truly unresolved.

 Naturalness is one of the most compelling motivations for new physics near the weak scale.

 The LHC will eventually probe conventional "colorful" theories to (at best) the ~1% level.

 But it will leave kinematic regions in conventional theories — and all regions of more novel theories — essentially untested, and the status of naturalness truly unresolved.

 A Higgs factory & 100 TeV collider can uniformly probe natural symmetry-based theories to the ~1% level with powerful complementarity.