

Lecture 8 - Circulating Beams and Imperfections

ACCELERATOR PHYSICS

MT 2014

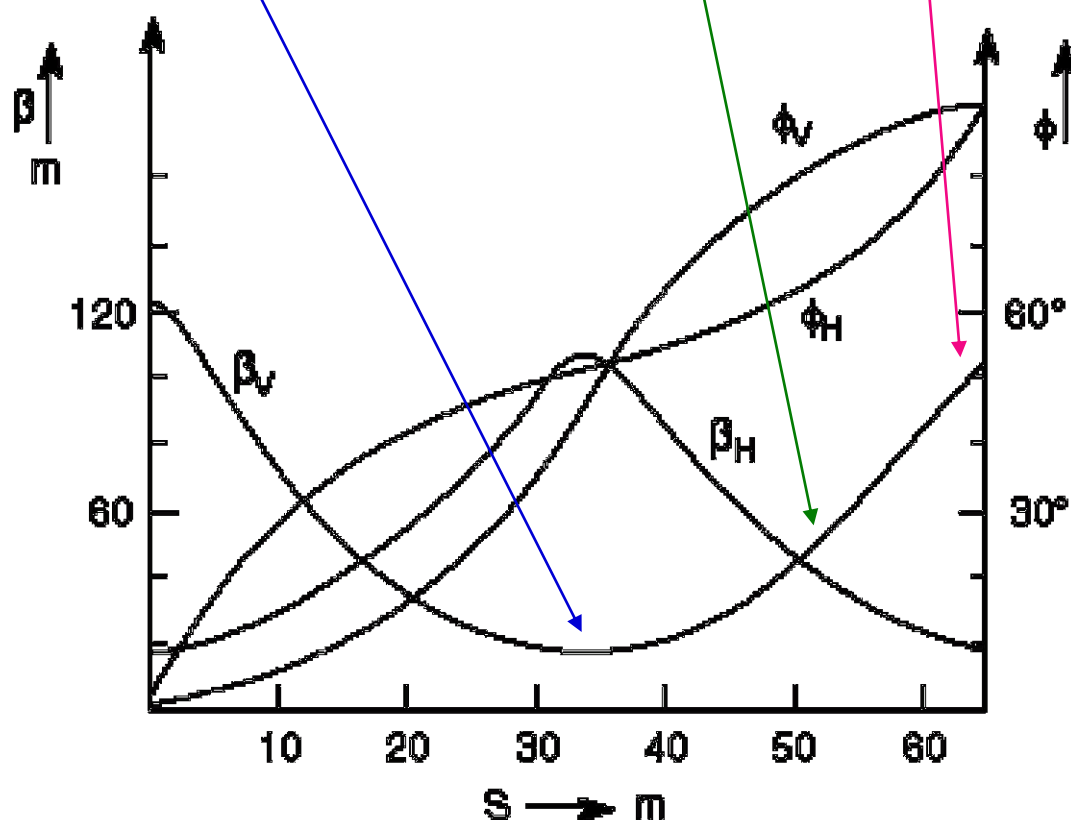
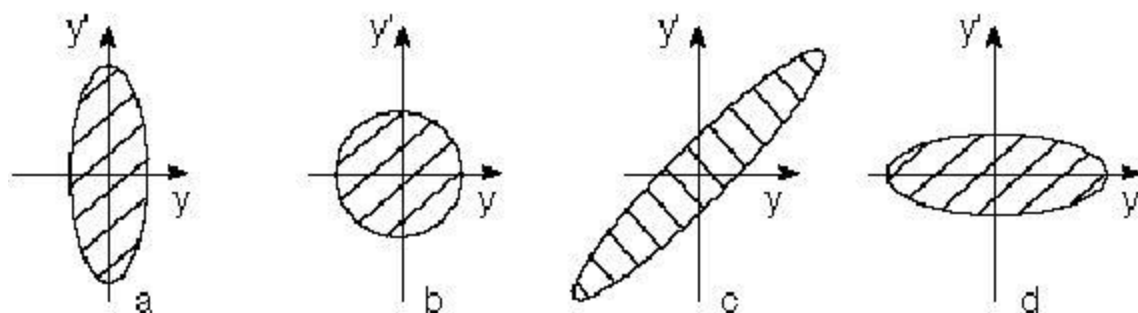
E. J. N. Wilson

- ◆ **Transition - does an accelerated particle catch up - it has further to go**
- ◆ **Phase jump at transition**
- ◆ **Synchrotron motion**
- ◆ **Synchrotron motion (continued)**
- ◆ **Large amplitudes**
- ◆ **Buckets**
- ◆ **Buckets**
- ◆ **Adiabatic capture**
- ◆ **A chain of buckets**

Lecture 7 - Circulating Beams and Errors - Contents

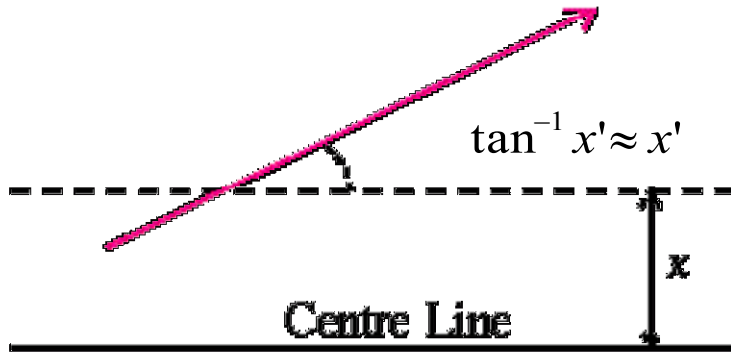
- ◆ **Phase space at various points in a lattice**
- ◆ **Adiabatic damping**
- ◆ **Adiabatic damping of proton beam**
- ◆ **Liouville's theorem - Strict Version**
- ◆ **Emittance definitions**
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- ◆ **Making an orbit bump grow**
- ◆ **Circle diagram**
- ◆ **Closed orbit in the circle diagram**
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- ◆ **Working daigram**

Betatron phase space at various points in a lattice

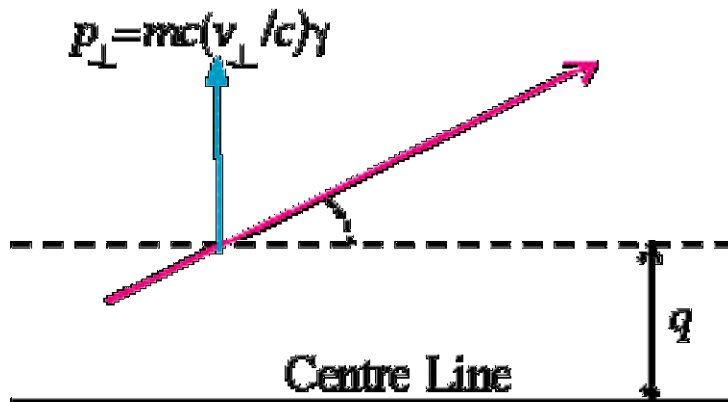


Adiabatic damping

◆ Accelerator coordinates



◆ Canonical coordinates of Hamilton

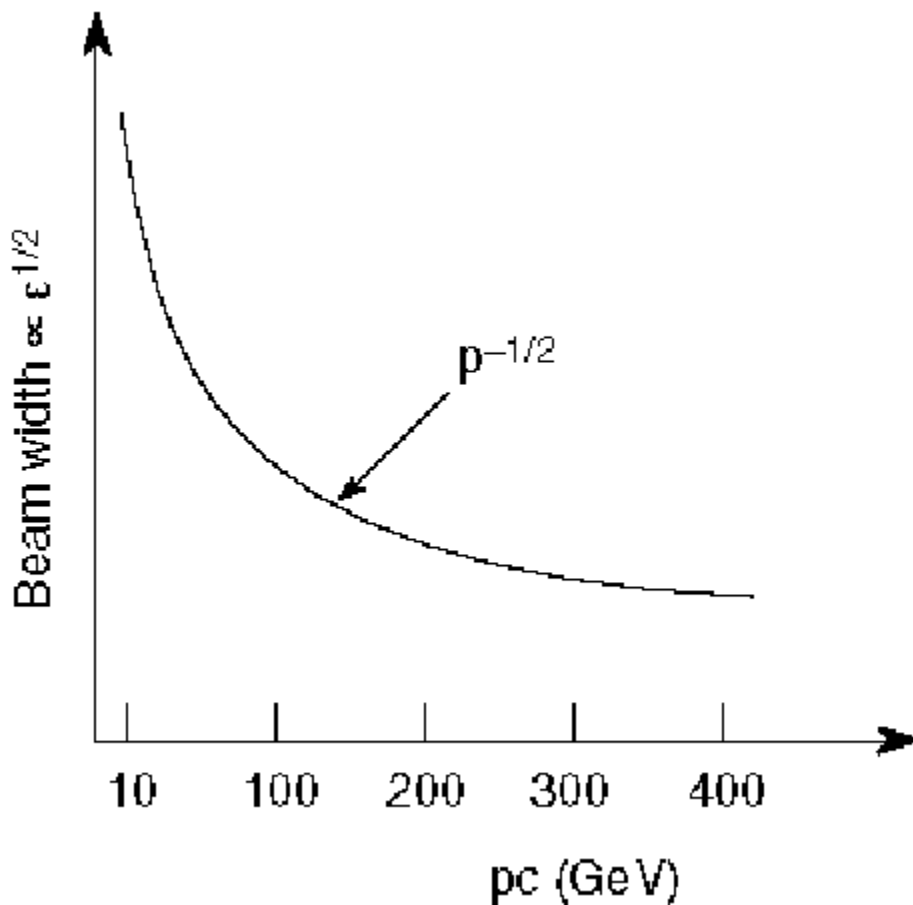


$$p = m \frac{dq}{dt} \gamma = m \frac{ds}{dt} \frac{dq}{ds} \gamma = m \alpha(\beta\gamma) x'$$

$$\int p dq = m \alpha(\beta\gamma) \int x' dx = p_0 \int x' dx$$

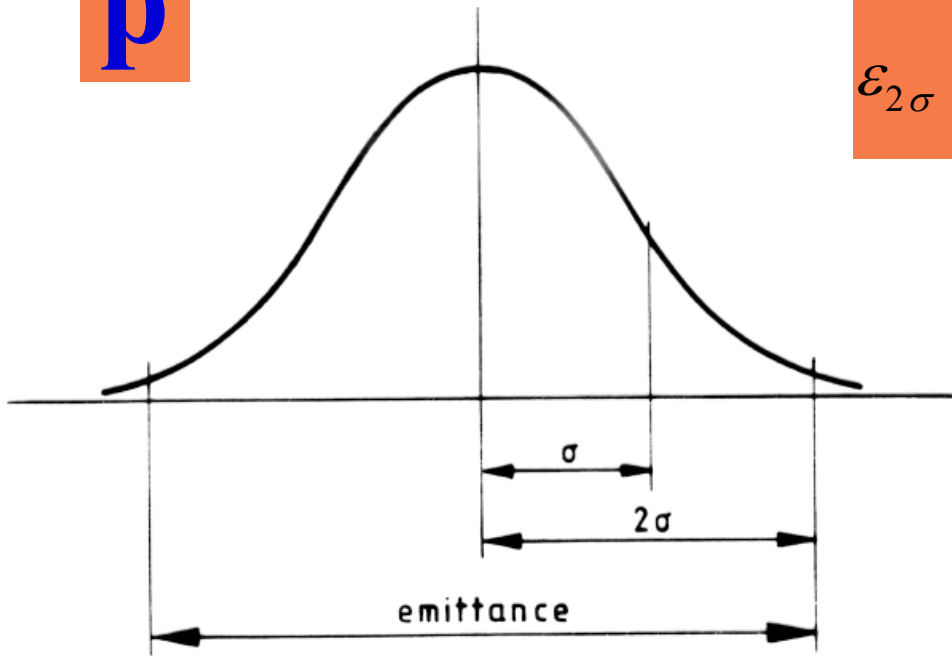
$$\text{Emittance} = \pi \epsilon = \int x' dx = \pi \epsilon_{\text{normalised}} / (\beta\gamma) \propto 1/p_0$$

Adiabatic damping of proton beam



Emittance definitions

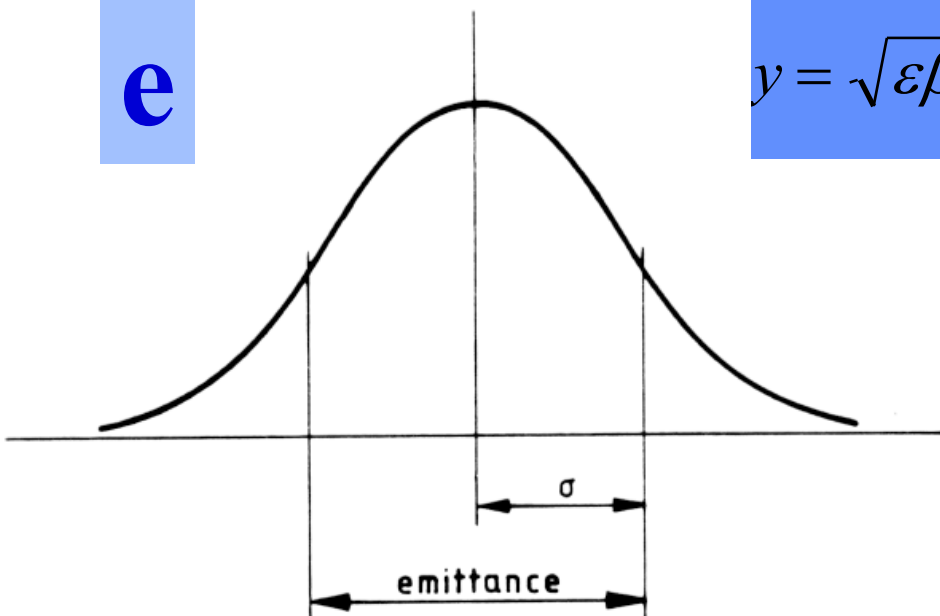
p



$$y = \sqrt{\varepsilon\beta} \quad \therefore \varepsilon = \frac{y^2}{\beta}$$

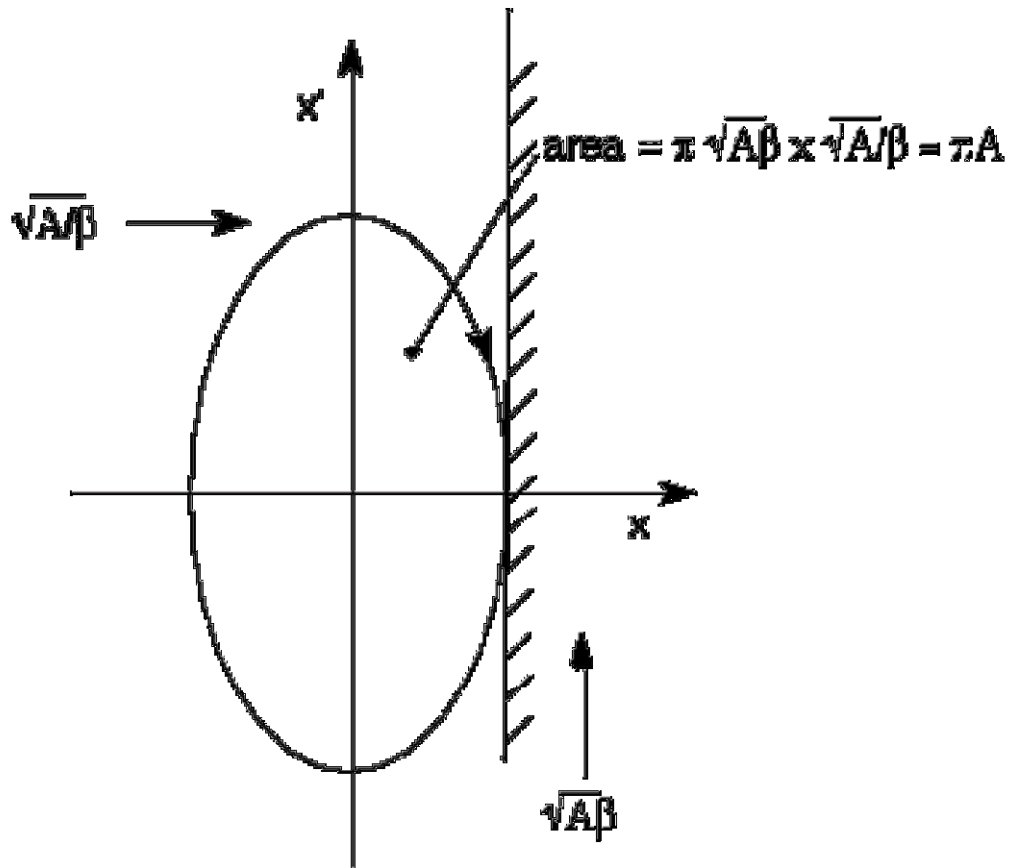
$$\varepsilon_{2\sigma} = \frac{(2\sigma)^2}{\beta} = \frac{4\sigma^2}{\beta}$$

e



$$y = \sqrt{\varepsilon\beta} \quad \varepsilon_{\sigma} = \frac{\sigma^2}{\beta}$$

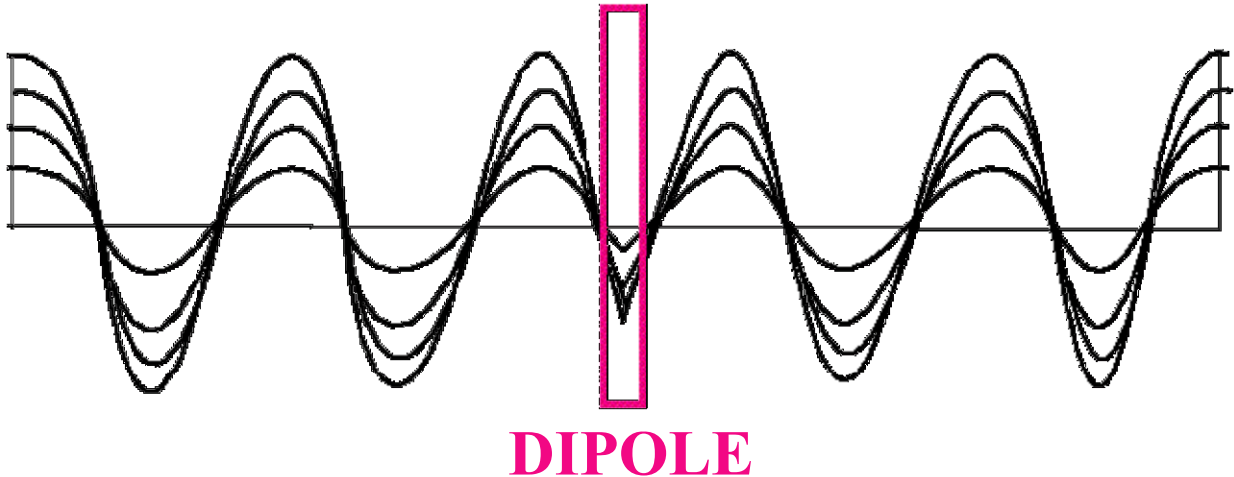
Acceptance



- ◆ Largest particle grazing an obstacle defines acceptance.
- ◆ Acceptance is equivalent to emittance

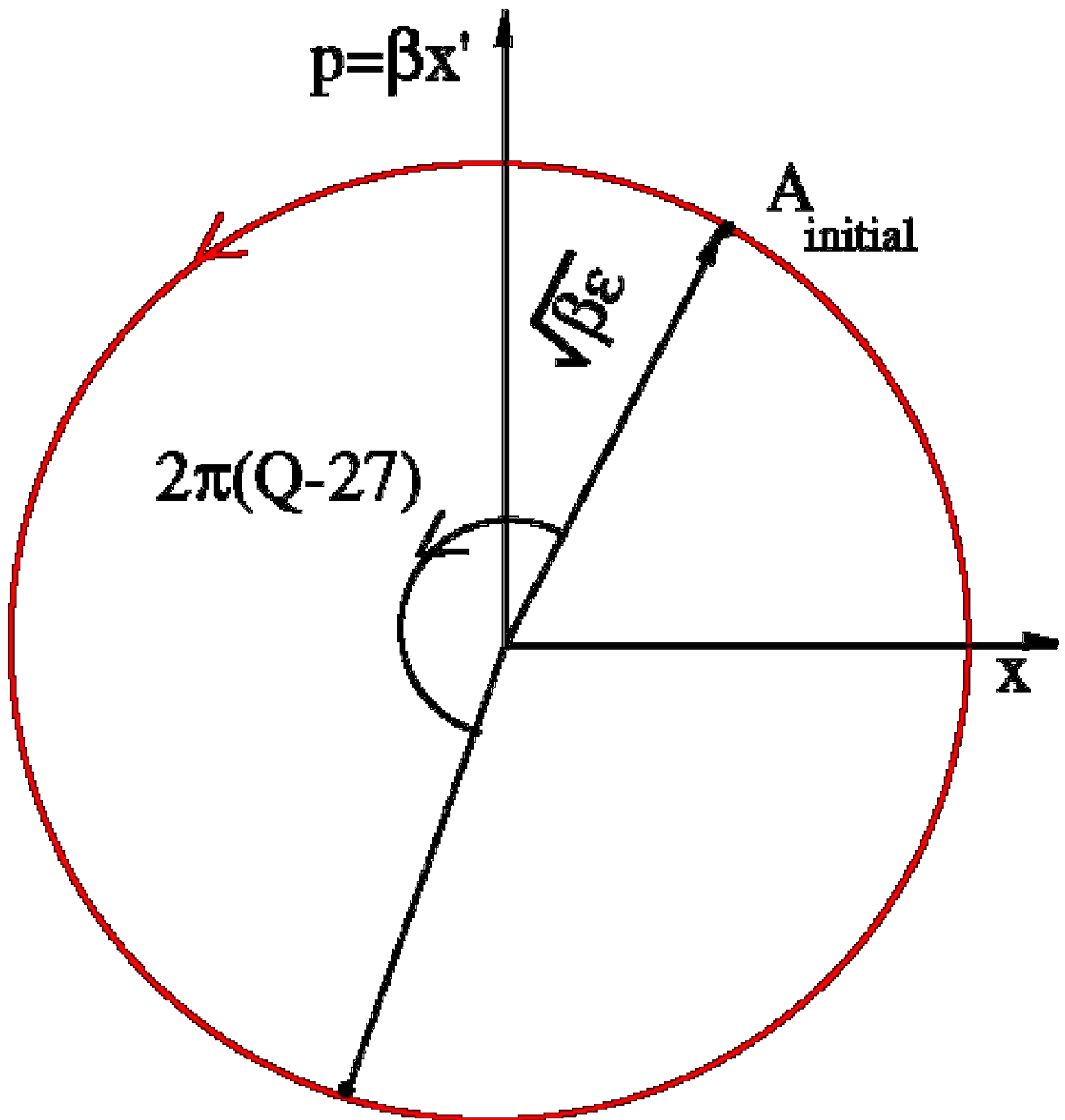
$$A = \frac{\hat{x}^2}{\beta}$$

Making an orbit bump grow



- ◆ As we slowly raise the current in a dipole:
- ◆ The zero-amplitude betatron particle follows a distorted orbit
- ◆ The distorted orbit is **CLOSED**
- ◆ It is still obeying Hill's Equation
- ◆ Except at the kink (dipole) it follows a betatron oscillation.
- ◆ Other particles with finite amplitudes oscillate about this new closed orbit

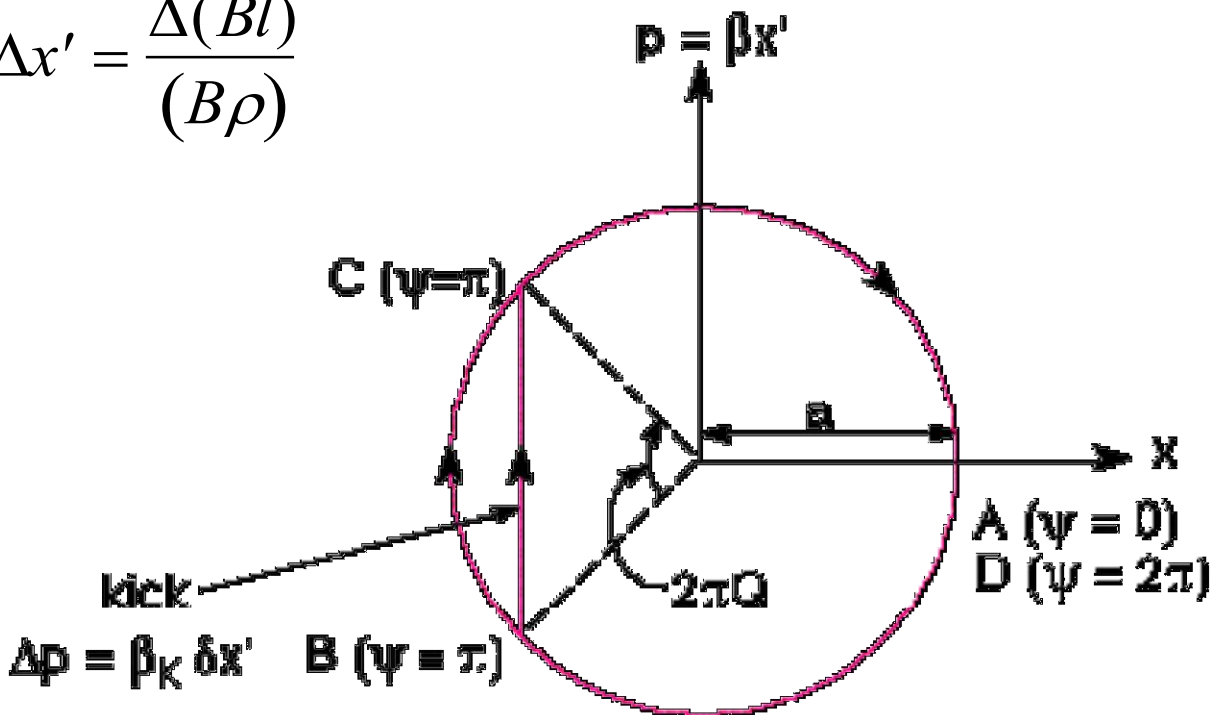
Circle diagram



B (after one turn and
27+ 0.6 betatron oscillations)

Closed orbit in the circle diagram

$$\Delta x' = \frac{\Delta(Bl)}{(B\rho)}$$



**Tracing a closed orbit for one turn
in the circle diagram with a single kick.
The path is ABCD.**

$$\frac{\Delta p}{2} = \frac{\beta_k \delta x'}{2} = a \sin\left(\frac{2\pi Q}{2}\right)$$

$$a = \frac{\beta_k \delta x'}{2 \sin \pi Q} \quad \text{elsewhere} \quad \hat{x} = a \sqrt{\frac{\beta(s)}{\beta_k}} = \frac{\sqrt{\beta_k \beta(s)}}{2 \sin \pi Q} \delta x'$$

Uncorrelated errors

- ◆ A random distribution of dipole errors

- ◆ Take the r.m.s. average of $\delta y_i' = \Delta(B\ell)/(B\rho)$
- ◆ Weighted according to the β_k values
- ◆ The expectation value of the amplitude is:

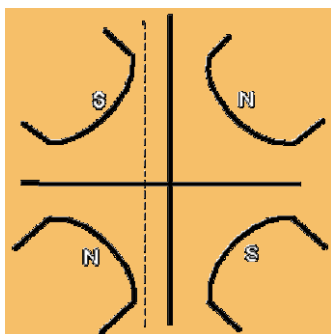
$$\langle y(s) \rangle = \frac{\sqrt{\beta(s)}}{2\sqrt{2} \sin \pi Q} \sqrt{\sum_i \beta_i \delta y_i'^2}$$

- ◆ Kicks from the N magnets in the ring.

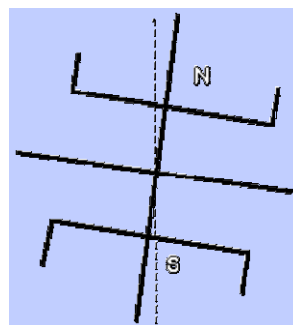
$$\approx \frac{\sqrt{\beta(s)\bar{\beta}}}{2\sqrt{2} \sin \pi Q} \sqrt{N} \frac{(\Delta B\ell)_{rms}}{B\rho}$$

- ◆ The factor $\sqrt{2}$ takes into account the averaging over both sine and cosine phases
- ◆ A further factor 2 safety is applied to include 98% of all sample distributions.

Sources of distortion



Δy



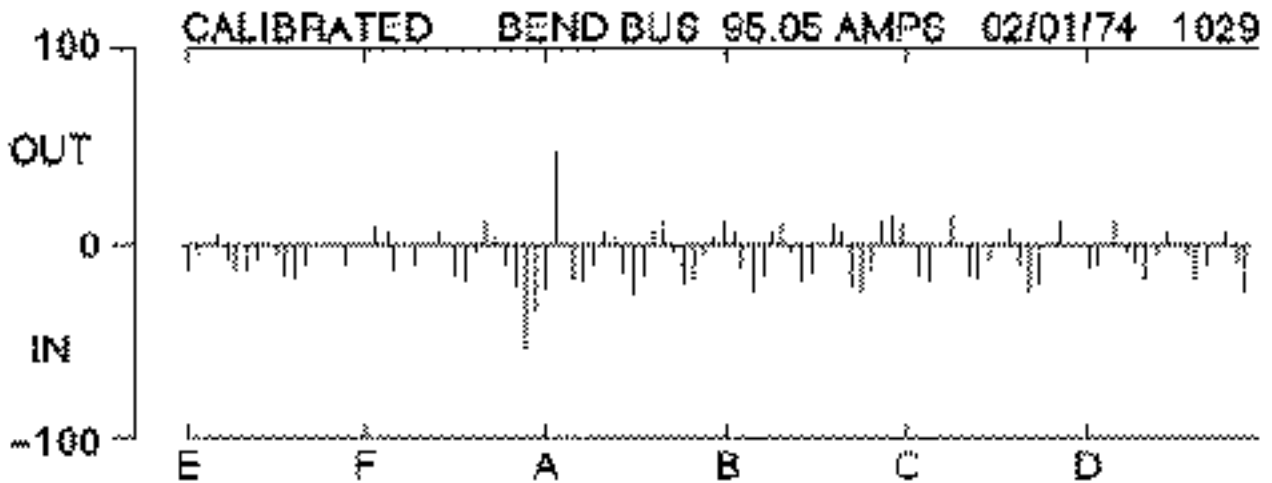
Δ

Table 1

Sources of Closed Orbit Distortion

| Type of element | Source of kick | r.m.s. value | $\langle \Delta B / (B\rho) \rangle_{rms}$ | plane |
|---|----------------|------------------------------|--|--------|
| Gradient magnet | Displacement | $\langle \Delta y \rangle$ | $k_i l_i \langle \Delta y \rangle$ | x, z |
| Bending magnet (bending angle = θ_i) | Tilt | $\langle \Delta \rangle$ | $\theta_i \langle \Delta \rangle$ | z |
| Bending magnet | Field error | $\langle \Delta B/B \rangle$ | $\theta_i \langle \Delta B/B \rangle$ | x |
| Straight sections (length = d_i) | Stray field | $\langle \Delta B_s \rangle$ | $d_i \langle \Delta B_s \rangle / (B\rho)_{inj}$ | x, z |

FNAL MEASUREMENT



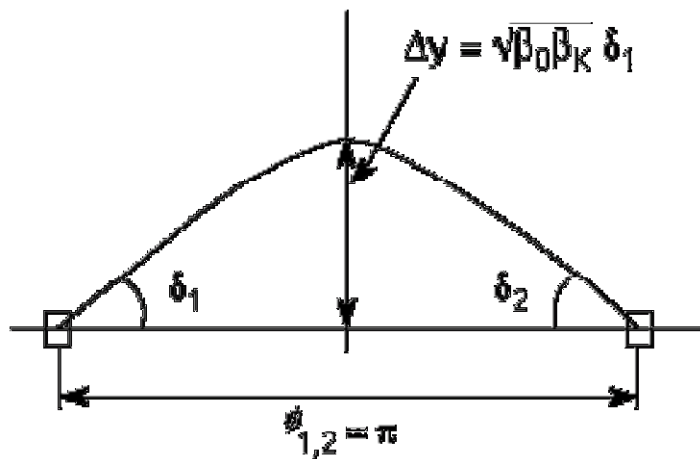
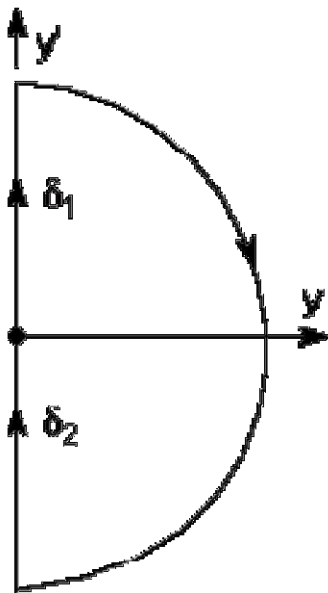
- ◆ Historic measurement from FNAL main ring
- ◆ Each bar is the position at a quadrupole
- ◆ +/- 100 is width of vacuum chamber
- ◆ Note mixture of 19th and 20th harmonic
- ◆ The Q value was 19.25

Diad bump

- ◆ Simplest bump is from two equal dipoles 180 degrees apart in betatron phase. Each gives:

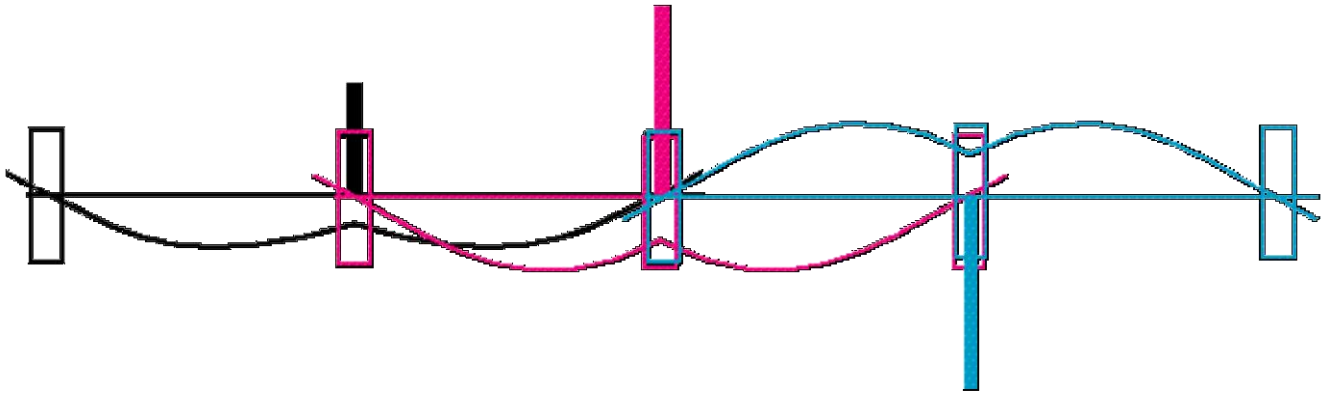
$$\delta = \frac{\Delta(B1)}{B\rho}$$

- ◆ The trajectory is : $y(s) = \delta \sqrt{\beta(s) \beta_k} \sin(\phi - \phi_0)$



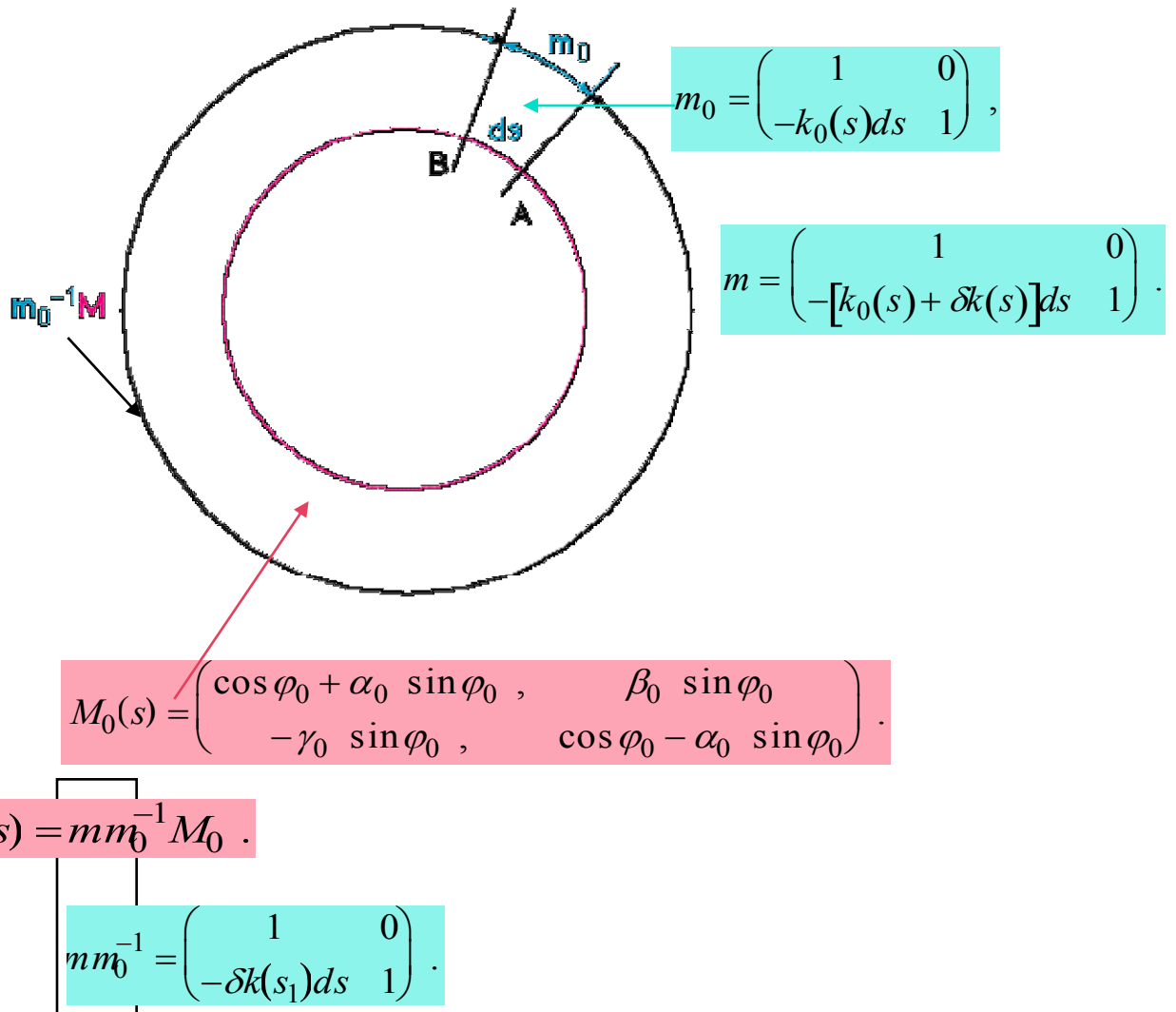
- ◆ The matrix is
$$\begin{pmatrix} y \\ y' \end{pmatrix} = \begin{pmatrix} (\sqrt{\beta} / \sqrt{\beta_0}) (\cos \Delta\phi + \alpha_0 \sin \Delta\phi) & \sqrt{\beta_0 \beta} \sin \Delta\phi \\ (-1 / \sqrt{\beta_0 \beta}) \{ (\alpha - \alpha_0) \cos \Delta\phi + (1 + \alpha \alpha_0) \sin \Delta\phi \} & (\sqrt{\beta} / \sqrt{\beta_0}) (\cos \Delta\phi - \alpha \sin \Delta\phi) \end{pmatrix} \begin{pmatrix} y_0 \\ y'_0 \end{pmatrix}$$

Overlapping beam bumps



- ◆ Each colour shows a triad bump centred on a beam position measurement.
- ◆ A computer calculates the superposition of the currents in the dipoles and corrects the whole orbit simultaneously

Gradient errors



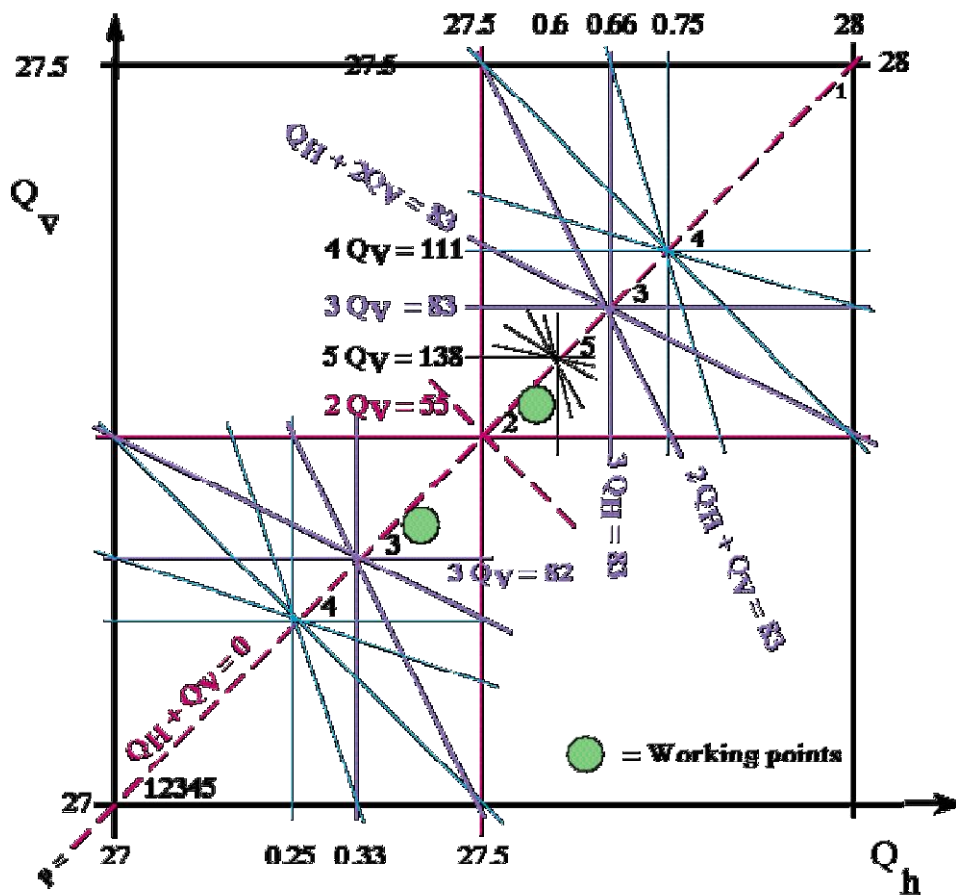
$$M = \begin{pmatrix} \cos \phi_0 + \alpha_0 \sin \phi_0 & \beta_0 \sin \phi_0 \\ -\delta k(s)ds(\cos \phi_0 + \alpha_0 \sin \phi_0) - \gamma \sin \phi_0 & -\delta k(s)ds\beta_0 \sin \phi_0 + \cos \phi_0 - \alpha_0 \sin \phi_0 \end{pmatrix}.$$

$$\Delta(\text{Tr } M) / 2 = \Delta(\cos \phi) = -\Delta\phi \sin \phi_0 = \frac{\sin \phi_0}{2} \beta_0(s) \delta k(s) ds$$

$$2\pi\Delta Q = \Delta\phi = \frac{\beta(s)\delta k(s)ds}{2}.$$

$$\Delta Q = \frac{1}{4\pi} \int \beta(s) \delta k(s) ds.$$

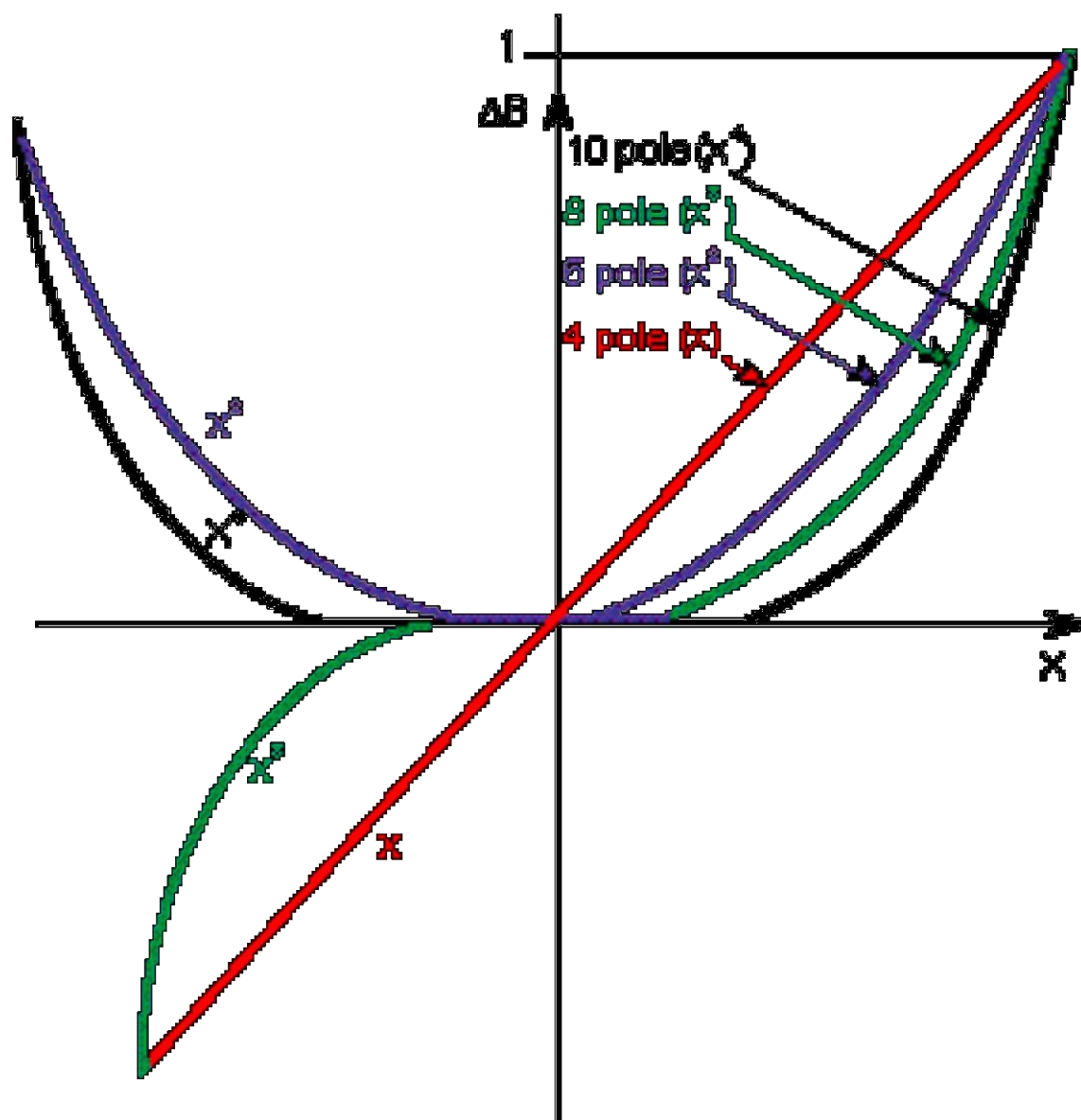
Q diagram



$$nQ = p ,$$

$$1Q_H + mQ_V = p ,$$

Multipole field shapes



Physics of Chromaticity

- ◆ The Q is determined by the lattice quadrupoles whose strength is:

$$k = \frac{1}{(B\rho)} \frac{dB_z}{dx} \propto \frac{1}{p}$$

- ◆ Differentiating:

$$\frac{\Delta k}{k} = -\frac{\Delta p}{p} .$$

- ◆ Remember from gradient error analysis

$$\Delta Q = \frac{1}{4\pi} \int \beta(s) \delta k(s) ds .$$

- ◆ Giving by substitution

$$\Delta Q = \frac{1}{4\pi} \int \beta(s) \Delta k(s) ds = \left[\frac{-1}{4\pi} \int \beta(s) k(s) ds \right] \frac{\Delta p}{p} .$$

$$\Delta Q = Q' \frac{\Delta p}{p}$$

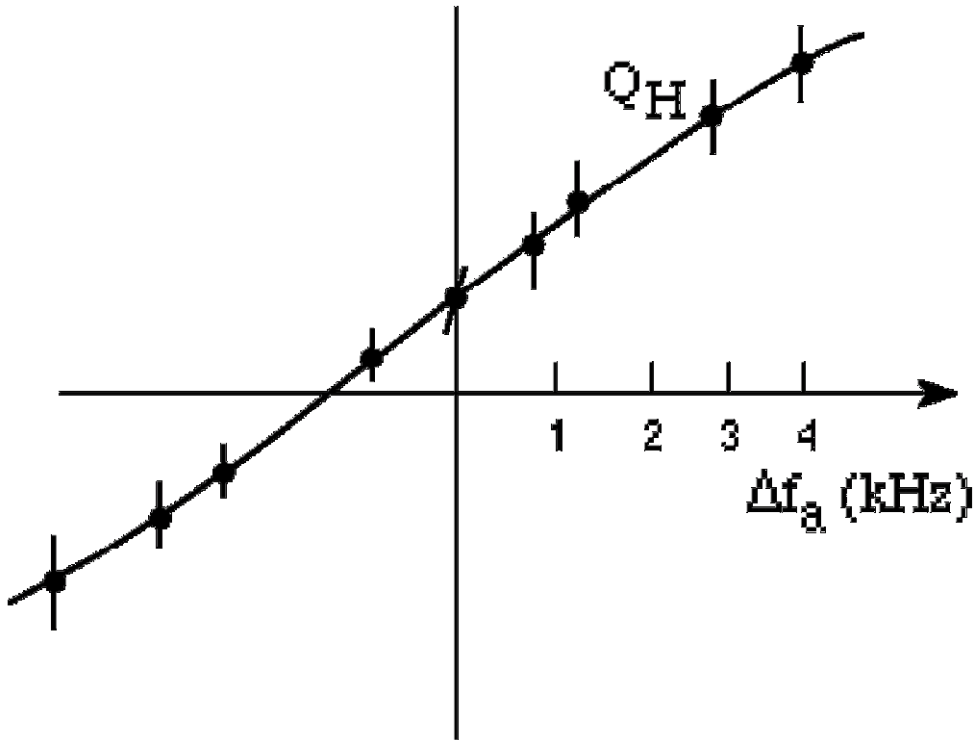
Q' is the chromaticity

- ◆ “Natural” chromaticity

$$Q' = -\frac{1}{4\pi} \oint \beta(s) k(s) ds \approx -1.3Q$$

N.B. Old books say $\xi = \frac{p}{Q} \frac{dQ}{dp} = \frac{Q'}{Q}$

Measurement of Chromaticity



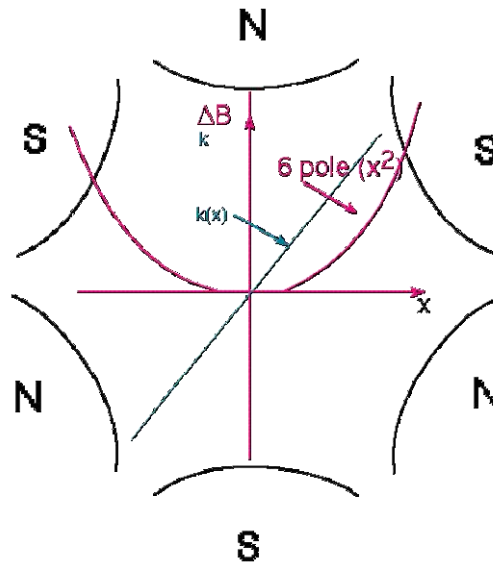
- ◆ We can steer the beam to a different mean radius and a different momentum by changing the rf frequency and measure Q

$$\Delta f_a = f_a \eta \frac{\Delta p}{p} \quad \Delta r = D_{av} \frac{\Delta p}{p}$$

- ◆ Since $\Delta Q = Q' \frac{\Delta p}{p}$

- ◆ Hence $\therefore Q' = f_a \eta \frac{dQ}{df_a}$

Correction of Chromaticity



- ◆ Parabolic field of a 6 pole is really a gradient which rises linearly with x
- ◆ If x is the product of momentum error and dispersion

$$\Delta k = \frac{B'' D}{(B\rho)} \frac{\Delta p}{p}.$$

- ◆ The effect of all this extra focusing cancels **chromaticity**

$$\Delta Q = \left[\frac{1}{4\pi} \int \frac{B''(s)\beta(s)D(s)ds}{(B\rho)} \right] \frac{dp}{p}.$$

- ◆ Because gradient is opposite in v plane we must have two sets of opposite polarity at F and D quads where betas are different

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