

Twist-3 Spin Observables for Single-Hadron Production in DIS

(A. Metz, Temple University, Philadelphia)

- Introduction and Motivation
- A related observable: double-spin asymmetry A_{LT} for $\vec{\ell} N^\uparrow \rightarrow \ell X$
- Single-spin asymmetry A_{UT} (A_N) for $\ell N^\uparrow \rightarrow h X$
- Single-spin asymmetry A_{UT} (A_N) for $\ell N \rightarrow \Lambda^\uparrow X$
- Double-spin asymmetry A_{LT} for $\vec{\ell} N^\uparrow \rightarrow h X$
- Summary

talk mainly based on

arXiv:1407.5078, Gamberg, Kang, A.M., Pitonyak, Prokudin

arXiv:1411.6459, Kanazawa, A.M., Pitonyak, Schlegel

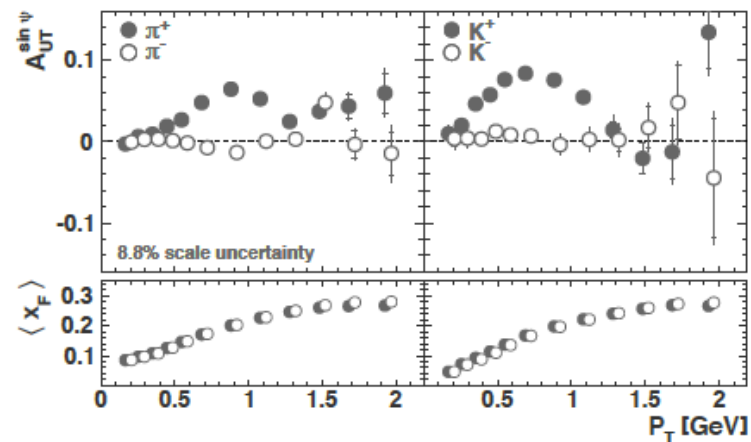
arXiv:1503.02003, Kanazawa, A.M., Pitonyak, Schlegel

Introduction and Motivation

1. Data exist for twist-3 spin asymmetries in $\ell N \rightarrow h X$

- example: A_N for $\ell N^\uparrow \rightarrow h X$ (HERMES, 2013 / JLab Hall A, 2013)

$$A_N = \frac{d\sigma^\uparrow - d\sigma^\downarrow}{d\sigma^\uparrow + d\sigma^\downarrow} \quad x_F = \frac{2P_{hL}}{\sqrt{s}}$$



(HERMES, 2013)

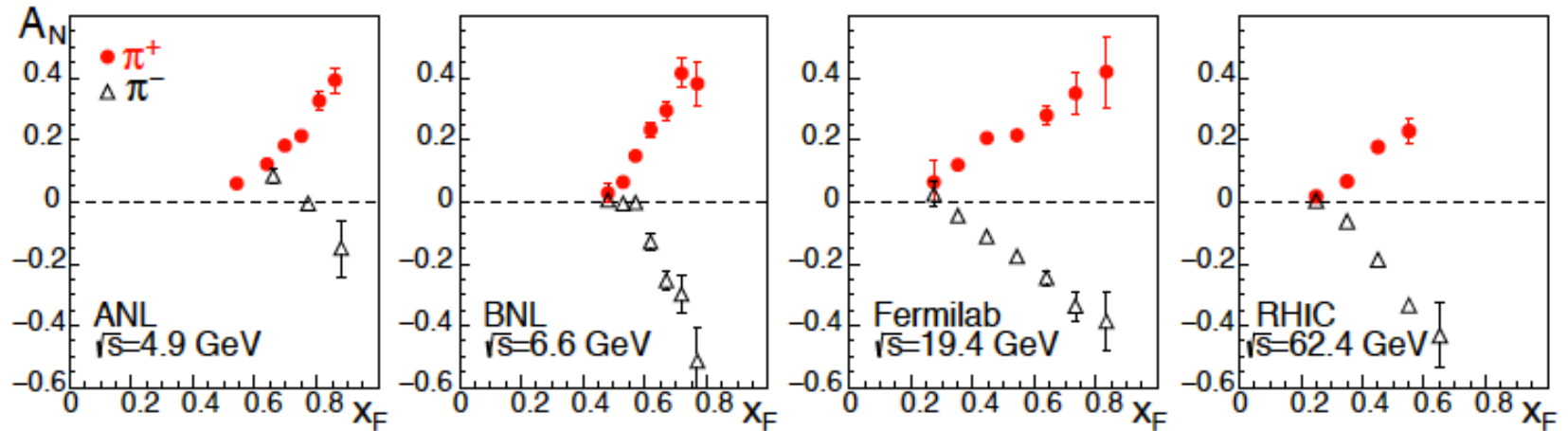
- more data: A_N for $\ell N \rightarrow \Lambda^\uparrow X$ (HERMES, 2014)

$$A_{LT} \text{ for } \vec{\ell} N^\uparrow \rightarrow h X \text{ (JLab Hall A, 2015)}$$

- can one understand those data, and what can one learn from them ?

2. Related asymmetries in processes like $pp \rightarrow h X$

- example: A_N for $pp^\uparrow \rightarrow \pi X$



(plot from Aidala, Bass, Hasch, Mallot, 2012)

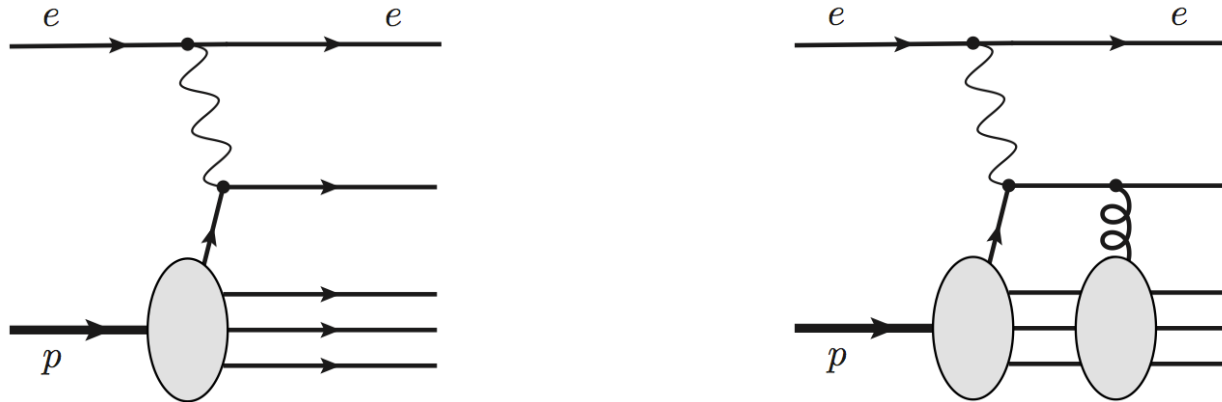
- data exist for A_N in $pp \rightarrow \Lambda^\uparrow X$ (Bunce et al, 1976 / ...)
- calculation available for A_{LT} for $\vec{p}p^\uparrow \rightarrow (h, \gamma, \text{jet}) X$
Liang, A.M., Pitonyak, Schäfer, Song, Zhou, 2012
- it is challenging to reveal the underlying physics of the available data
- maybe the asymmetries for ℓN help to understand the asymmetries for pp

3. Playground to solidify and streamline theory tools
(due to small number of Feynman graphs)
 - gauge-invariance of calculation
 - frame-independence of results
 - higher order corrections

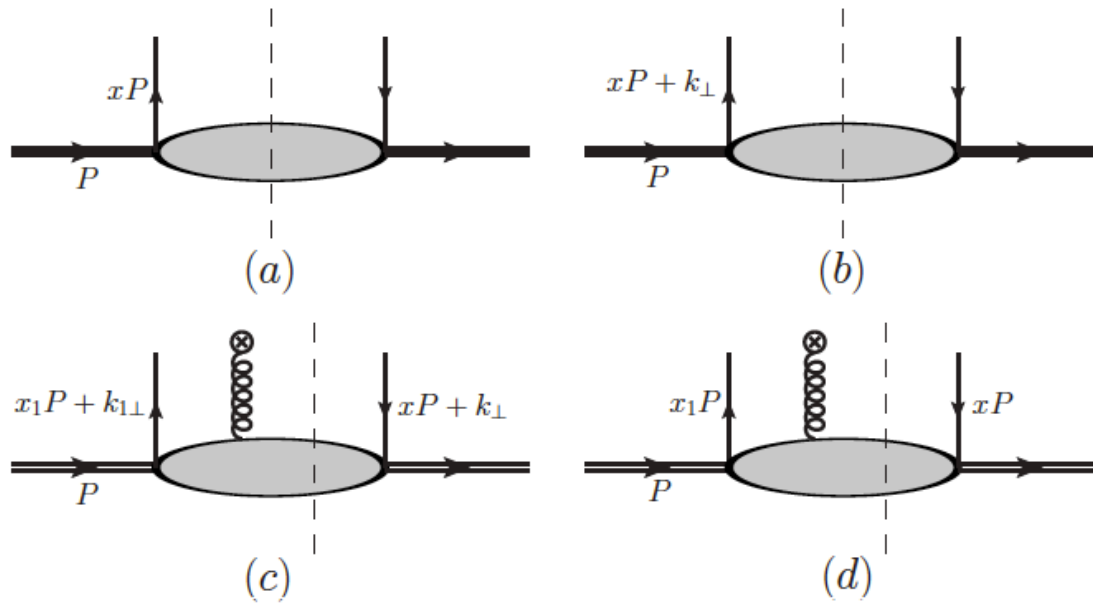
4. Explore potential of those observables for future measurements
 - **EIC** should be ideal for future studies
 - measurements possible for large transverse momenta $P_{h\perp}$

Reminder: double-spin asymmetry A_{LT} for $\vec{\ell} N^\uparrow \rightarrow \ell X$

- Re-scattering of struck quark matters at twist-3 (gluon with physical polarization)



- Contributing correlators after factorization



- collinear quark-quark correlator at twist-3 $\rightarrow g_T(x)$
- k_\perp -dependent quark-quark correlator $\rightarrow \tilde{g}(x) = \int d^2\vec{k}_\perp \frac{\vec{k}_\perp^2}{2M^2} g_{1T}(x, \vec{k}_\perp^2)$
- (collinear) quark-gluon-quark correlator $\rightarrow F_{FT}(x, x_1) \quad G_{FT}(x, x_1)$

- Exploit relations between functions

- relation due to QCD equation of motion

$$x g_T(x) = \int dx_1 \left[G_{DT}(x, x_1) - F_{DT}(x, x_1) \right]$$

- Final result

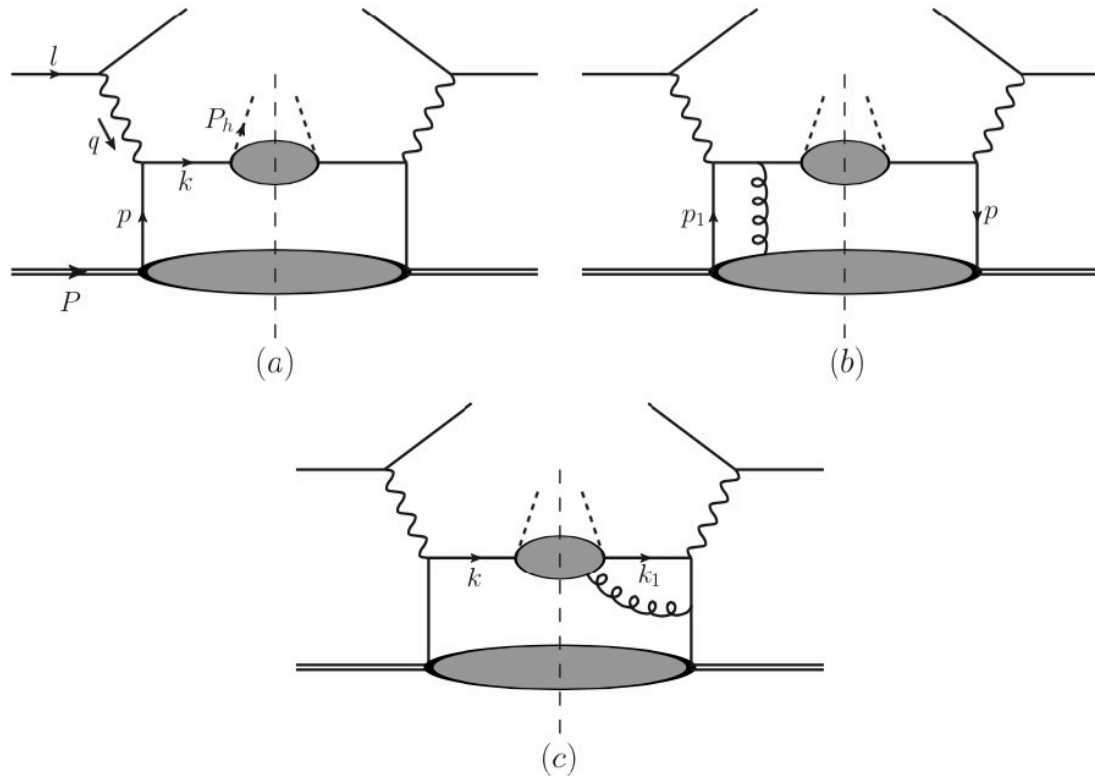
$$\frac{l'^0 d\sigma_{LT}}{d^3\vec{l}'} = -\frac{8 \alpha_{em}^2 x_B^2 \sqrt{1-y} M}{Q^5} \lambda_\ell |\vec{S}_\perp| \cos \phi_S \sum_q e_q^2 g_T^q(x_B)$$

- twist-3 effect
- final result looks rather simple
- comparable twist-3 observables may have more complicated structure

Single-spin asymmetry A_N for $\ell N^\uparrow \rightarrow h X$

(Gamberg, Kang, A.M., Pitonyak, Prokudin, 2014)

- Feynman diagrams for LO calculation



- twist-3 effects associated with nucleon, and with fragmentation process
- large scale for pQCD calculations provided by $P_{h\perp}$
- in LO formalism, hadron recoils against (undetected) final state lepton $\rightarrow Q^2$ large

- Analytical result

$$\begin{aligned}
P_h^0 \frac{d\sigma_{UT}}{d^3\vec{P}_h} = & -\frac{8\alpha_{\text{em}}^2}{S} \varepsilon_{\perp\mu\nu} S_{P\perp}^\mu P_{h\perp}^\nu \sum_q e_q^2 \int_{z_{\text{min}}}^1 \frac{dz}{z^3} \frac{1}{S+T/z} \frac{1}{x} \\
& \times \left\{ -\frac{\pi M}{\hat{u}} D_1^{h/q}(z) \left(F_{FT}^q(x, x) - x \frac{dF_{FT}^q(x, x)}{dx} \right) \left[\frac{\hat{s}(\hat{s}^2 + \hat{u}^2)}{2\hat{t}^3} \right] \right. \\
& \quad + \frac{M_h}{-x\hat{u} - \hat{t}} h_1^q(x) \left\{ \left(\hat{H}^{h/q}(z) - z \frac{d\hat{H}^{h/q}(z)}{dz} \right) \left[\frac{(1-x)\hat{s}\hat{u}}{\hat{t}^2} \right] \right. \\
& \quad \left. \left. + \frac{1}{z} H^{h/q}(z) \left[\frac{\hat{s}(\hat{s}^2 + (x-1)\hat{u}^2)}{\hat{t}^3} \right] + 2z^2 \int_z^\infty \frac{dz_1}{z_1^2} \frac{1}{\frac{1}{z} - \frac{1}{z_1}} \hat{H}_{FU}^{h/q, \mathfrak{S}}(z, z_1) \left[\frac{x\hat{s}^2\hat{u}}{\xi_z \hat{t}^3} \right] \right\} \right\}
\end{aligned}$$

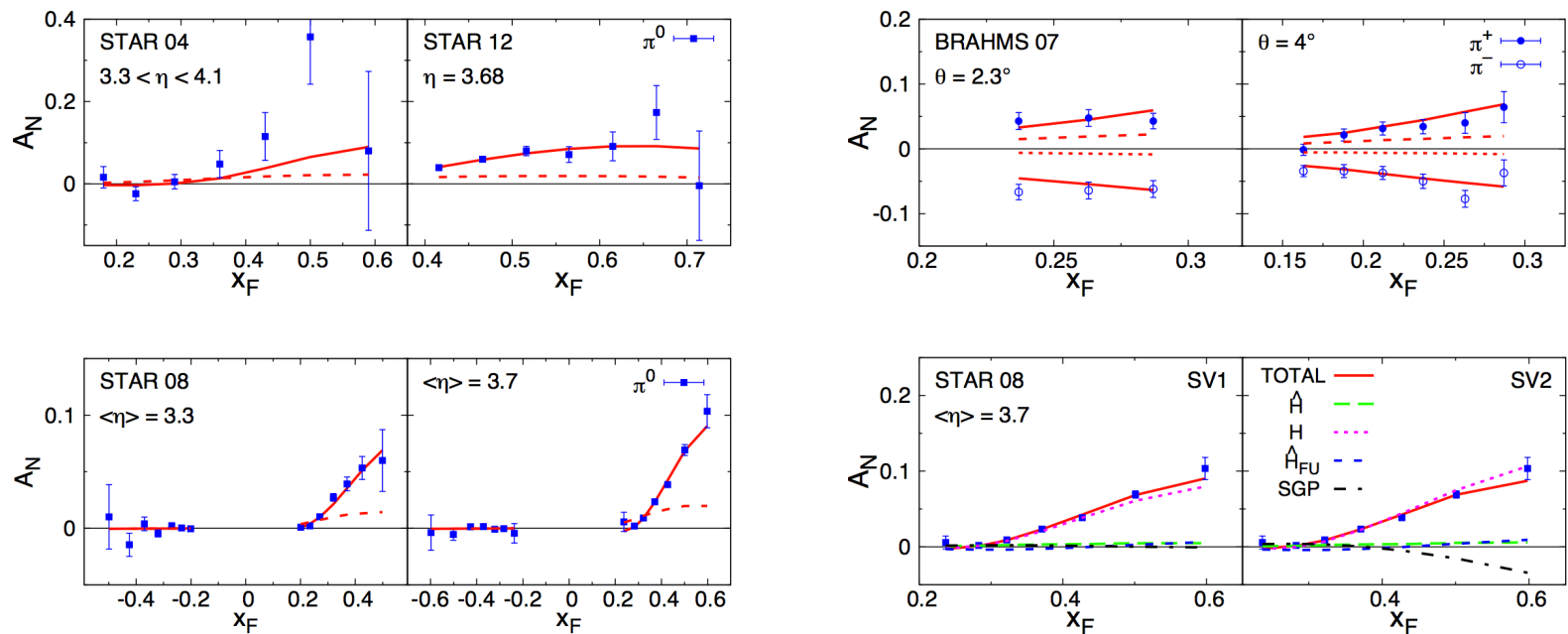
- Ingredients for numerics

- twist-2 FF D_1 (de Florian, Sassot, Stratmann, 2007)
- Qiu-Sterman function F_{FT} , from Sivers function f_{1T}^\perp (Anselmino et al, 2008)
- transversity h_1 (Anselmino et al, 2013)
- twist-3 FF \hat{H} , from Collins function H_1^\perp (Anselmino et al, 2013)

- twist-3 FF H and $\hat{H}_{FU}^{\mathfrak{S}}$ enter A_N for $pp^\uparrow \rightarrow h X$ (A.M., Pitonyak, 2012)
use model-independent relation

$$H^{h/q}(z) = -2z\hat{H}^{h/q}(z) + 2z^3 \int_z^\infty \frac{dz_1}{z_1^2} \frac{1}{\frac{1}{z} - \frac{1}{z_1}} \hat{H}_{FU}^{h/q, \mathfrak{S}}(z, z_1)$$

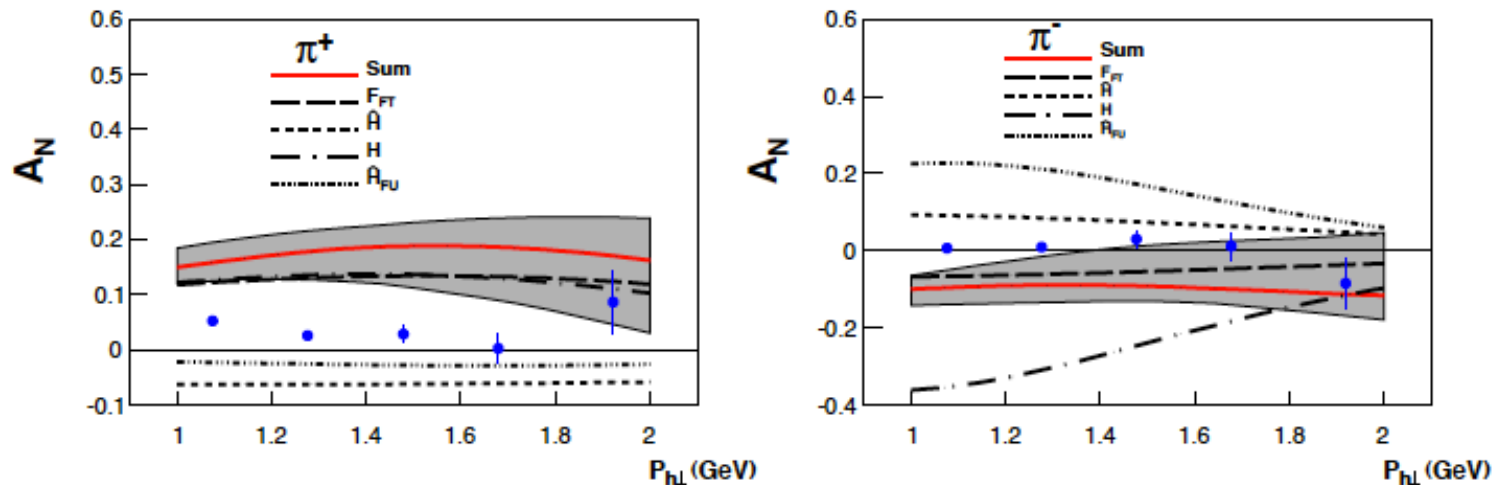
and fit to RHIC data for A_N (Kanazawa, Koike, A.M., Pitonyak, 2014)



- * good fit can be obtained ($\chi^2/\text{d.o.f} = 1.03$)
- * A_N dominated by twist-3 FFs (beyond moment of Collins function)
- * data on A_N for $\ell N^\uparrow \rightarrow h X$ may allow cross check

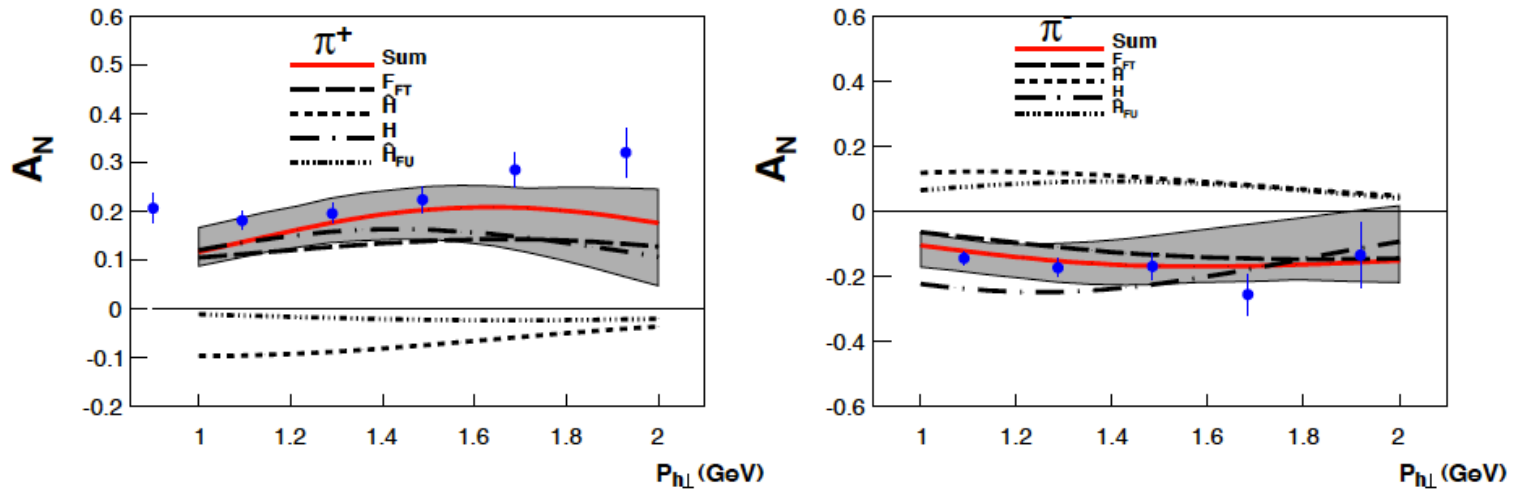
- Numerical results and discussion

- comparison with HERMES data (JLab data are at very low $P_{h\perp}$)



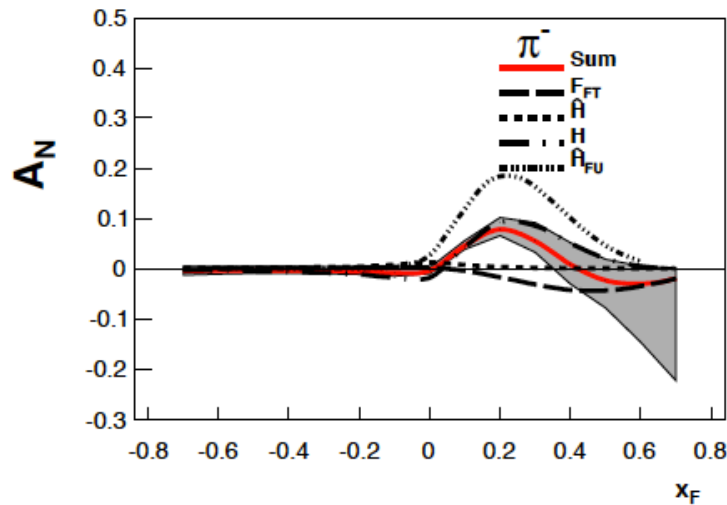
- * error band based on uncertainties of f_{1T}^\perp , h_1 , H_1^\perp only
- * relatively poor comparison with data, especially for π^+ production
- * potential reasons for discrepancy:
 - (1) no error band for twist-3 FF \hat{H}_{FU}^S and hence for FF H
 - (2) (significant) other source(s) for A_N in $pp^\uparrow \rightarrow hX$
 - (3) leading order formalism not appropriate for rather low $P_{h\perp}$ of available data; HERMES: even data at highest $P_{h\perp}$ dominated by quasi-real photo-production \rightarrow calculation of NLO correction needed

- better comparison with "DIS" sub-set of HERMES data ($Q^2 \geq 1 \text{ GeV}^2$)

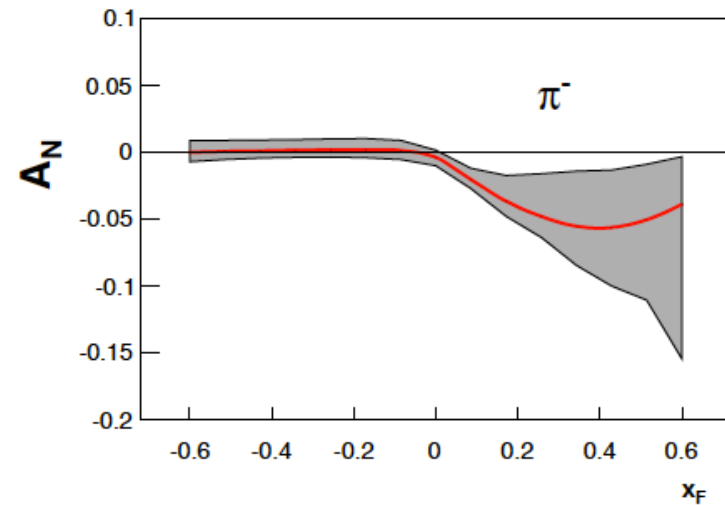


- prospects for measurement at an EIC ($\sqrt{S} = 63 \text{ GeV}$, $P_{h\perp} = 3 \text{ GeV}$)

with $\hat{H}_{FU}^{\mathcal{S}}$



without $\hat{H}_{FU}^{\mathcal{S}}$



* description of A_N in $pp^\uparrow \rightarrow hX$ through twist-3 fragmentation may be checked

Single-spin asymmetry A_N for $\ell N \rightarrow \Lambda^\uparrow X$

(Kanazawa, A.M., Pitonyak, Schlegel, 2015)

- Same Feynman graphs as for A_N in $\ell N^\uparrow \rightarrow h X$
- Analytical results
 - twist-3 distribution contribution

$$\frac{P_h^0 d\sigma_{LC}^{Dist}}{d^3\vec{P}_h} = \frac{8\pi M\alpha_{em}^2}{S} \epsilon_{\perp}^{P_{h\perp} S_{hT}} \sum_q e_q^2 \int_{z_{min}}^1 \frac{dz}{xz^3} \frac{1}{S+T/z} \frac{1}{\hat{u}} H_1^q(z) \left(x \frac{dH_{FU}^q(x, x)}{dx} \right) \left[-\frac{\hat{s}^2 \hat{u}}{\hat{t}^3} \right]$$

* no non-derivative term (first observed by Zhou, Yuan, Liang, 2008)

- twist-3 fragmentation contribution

$$\begin{aligned} \frac{P_h^0 d\sigma_{LC}^{Frag}}{d^3\vec{P}_h} &= \frac{2M_h\alpha_{em}^2}{S} \epsilon_{\perp}^{P_{h\perp} S_{hT}} \sum_q e_q^2 \int_{z_{min}}^1 \frac{dz}{xz^3} \frac{1}{S+T/z} \frac{1}{-\hat{t}-x\hat{u}} f_1^q(x) \\ &\times \left[z \frac{d\hat{D}_T^q(z)}{dz} \hat{\sigma}_D + \hat{D}_T^q(z) \hat{\sigma}_N + \frac{1}{z} D_T^q(z) \hat{\sigma}_2 + \int_z^\infty \frac{dz_1}{z_1^2} \frac{1}{1/z - 1/z_1} \frac{1}{\xi} \hat{D}_{FT}^{q, \mathfrak{S}}(z, z_1) \hat{\sigma}_3 \right] \end{aligned}$$

* (again) contributions from three sources

- Discussion
 - identified all relevant parton correlation functions and their relations
 - first complete result in twist-3 collinear factorization
 - results obtained in both light-cone gauge and Feynman gauge
(for the first time for twist-3 fragmentation contribution to transverse SSA)
 - checked fragmentation contribution to A_N for $\ell N^\uparrow \rightarrow h X$ in Feynman gauge
 - numerical estimate needed
 - ingredients available for calculation of A_N for process like $pp \rightarrow \Lambda^\uparrow X$

Double-spin asymmetry A_{LT} for $\vec{\ell} N^\uparrow \rightarrow h X$

(Kanazawa, A.M., Pitonyak, Schlegel, 2014)

- Same Feynman graphs as before, but somewhat different treatment of kinematics
- Analytical results (calculation in lepton-nucleon cm frame)
 - results obtained in both light-cone gauge and Feynman gauge
 - twist-3 distribution contribution

$$\frac{P_h^0 d\sigma_{LT}^{Dist}(\lambda_\ell, \vec{S}_\perp)}{d^3\vec{P}_h} = -\frac{8\alpha_{em}^2}{S} M \vec{P}_{h\perp} \cdot \vec{S}_\perp \lambda_\ell \sum_q e_q^2 \int_{z_{min}}^1 \frac{dz}{z^3} \frac{1}{S+T/z} \frac{1}{x\hat{u}} D_1^{h/q}(z) \\ \times \left\{ \left(\tilde{g}^q(x) - x \frac{d\tilde{g}^q(x)}{dx} \right) \left[\frac{\hat{s}(\hat{s} - \hat{u})}{2\hat{t}^2} \right] + x g_T^q(x) \left[\frac{\hat{u}}{2\hat{t}} \right] + \int dx_1 G_{DT}^q(x, x_1) \left[\frac{\hat{u}(\hat{s} - \hat{u})}{\xi\hat{t}^2} \right] \right\}$$

* (again) contributions from three sources

- twist-3 fragmentation contribution

$$\frac{P_h^0 d\sigma_{LT}^{Frag}(\lambda_\ell, \vec{S}_\perp)}{d^3\vec{P}_h} = -\frac{8\alpha_{em}^2}{S} M_h \vec{P}_{h\perp} \cdot \vec{S}_\perp \lambda_\ell \sum_q e_q^2 \int_{z_{min}}^1 \frac{dz}{z^3} \frac{1}{S+T/z} \frac{1}{zx\hat{t}} h_1^q(x) E^{h/q}(z) \left[-\frac{\hat{s}}{\hat{t}} \right]$$

* final result becomes rather simple

- Comparison between lepton-nucleon cm-frame and nucleon-hadron cm-frame
 - agreement after taking into account the following relation (LIR):
(Bukhvostov, Kuraev, Lipatov, 1984 / ...)

$$g_T(x) = g_1(x) - \frac{2}{x} \int dx_1 \frac{1}{\xi} G_{DT}(x, x_1)$$

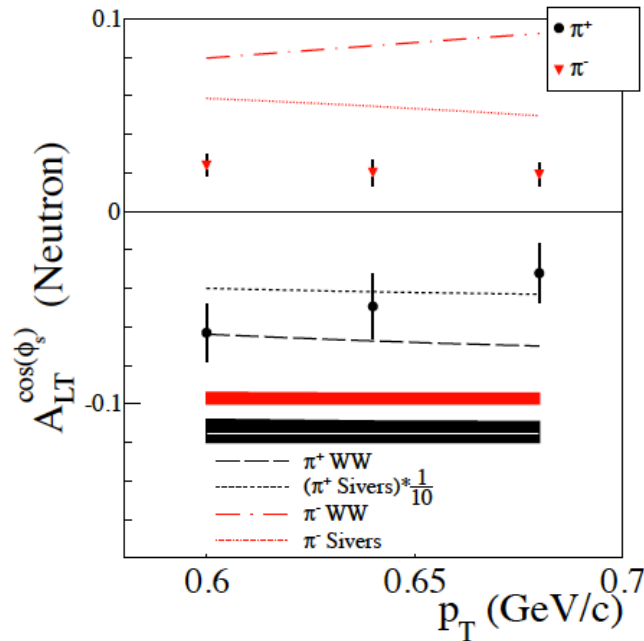
- matters here only for twist-3 distribution contribution
- LIR also allows one to simplify the final result

$$\frac{P_h^0 d\sigma_{LT}^{Dist}(\lambda_\ell, \vec{S}_\perp)}{d^3 \vec{P}_h} = -\frac{8\alpha_{em}^2}{S} M \vec{P}_{h\perp} \cdot \vec{S}_\perp \lambda_\ell \sum_q e_q^2 \int_{z_{min}}^1 \frac{dz}{z^3} \frac{1}{S + T/z} \frac{1}{x\hat{u}} D_1^{h/q}(z)$$

$$\times \left\{ \left(\tilde{g}^q(x) - x \frac{d\tilde{g}^q(x)}{dx} \right) \left[\frac{\hat{s}(\hat{s} - \hat{u})}{2\hat{t}^2} \right] + x g_T^q(x) \left[\frac{-\hat{s}\hat{u}}{\hat{t}^2} \right] + x g_1^q(x) \left[\frac{\hat{u}(\hat{s} - \hat{u})}{2\hat{t}^2} \right] \right\}$$

- * in particular, numerics becomes much easier (g_1 instead of G_{DT})

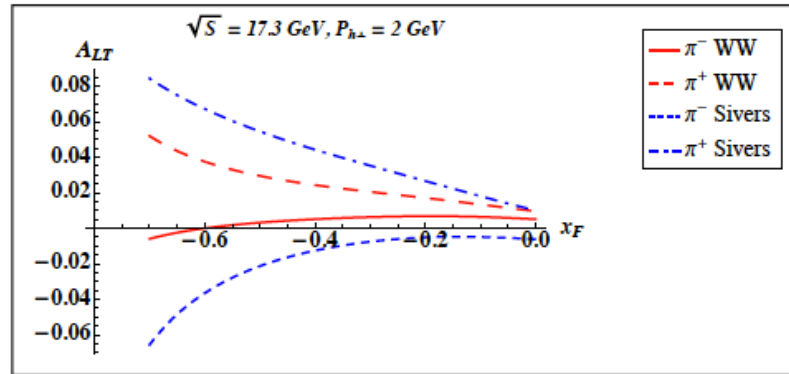
- Numerical results and discussion (twist-3 distribution term only)
 - use Wandzura-Wilceck approximation for g_T
 - two models for \tilde{g}
 - (1) WW-type approximation: $\tilde{g}(x) = g_{1T}^{(1)}(x) \approx x \int_x^1 \frac{dy}{y} g_1(y)$
 - (2) large N_c analysis suggests: $\tilde{g}(x) \approx -f_{1T}^{(1)}(x)$ (Pobylitsa, 2003)
 - comparison with data from JLab Hall A



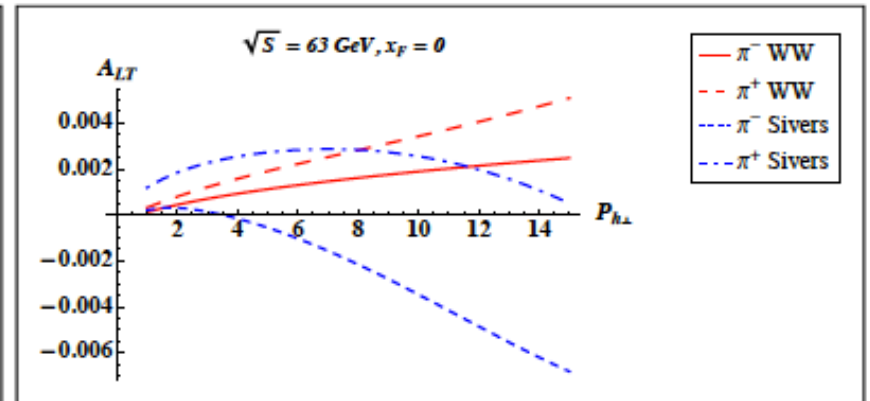
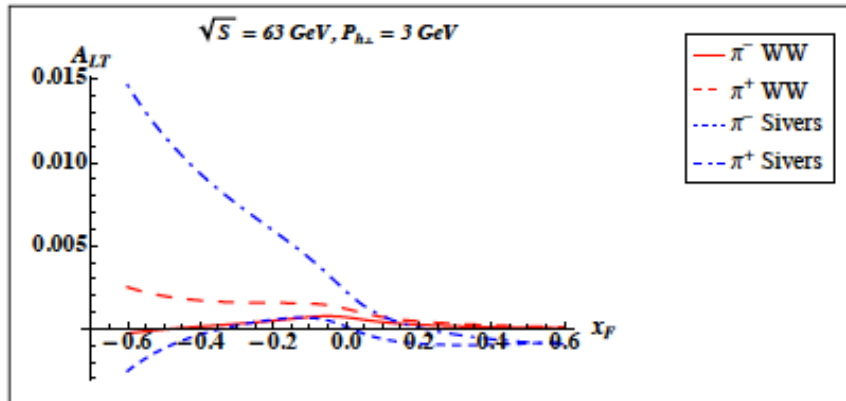
(JLab Hall A, 2015)

- * agreement in sign
- * quantitative description only for π^+ for WW-type approximation
- * but, keep in mind that $P_{h\perp}$ of data certainly too low
- * in general, A_{LT} may allow one to study the TMD g_{1T}

- prediction for COMPASS ($\sqrt{S} = 17.3 \text{ GeV}$, $P_{h\perp} = 2 \text{ GeV}$)



- prediction for an EIC ($\sqrt{S} = 63 \text{ GeV}$)



- A_{LT} becomes very small at higher energies

Summary

- Twist-3 spin asymmetries for $\ell N \rightarrow h X$ are interesting "new" observables
- They may give new insight into non-perturbative parton correlation functions
- They may give new insight into corresponding asymmetries for hadronic collisions
- First data available for three such asymmetries
- LO calculations available for those three asymmetries
 - helped to solidify theory tools (gauge invariance, frame independence)
 - phenomenology is at exploratory stage
- Calculation of higher order corrections needed