Studying Small-x Physics at an ElC Sartre, a Generator for Diffractive Physics

Thomas Ullrich (BNL/Yale) Dallas, April 29, 2015

Electron-Ion Collider

Investigate with precision universal dynamics of gluons

Ultimate QCD machine:

- The world's first polarized electron-polarized proton collider
- The world's first electron-heavy ion collider
- Luminosities: a hundred to up to a thousand times HERA
- Fine resolution inside proton down to 10⁻¹⁸ meters

EIC:

- NP Long Range Plan Resolution Meeting (April '15): Green Light
- Possible sites: BNL (eRHIC), JLab (MEIC)
 - Add ERL+FFAG Recirculating e Rings to RHIC facility
 - Figure-8 Ring-Ring Collider, use of CEBAF

What Drives Interest In e+A Collisions?

Gluon Saturation

- Can we find experimentally evidence of non-linear QCD dynamics?
- What is the dynamics of gluon saturation? How does it evolve?
- Is the Color Glass Condensate the correct theory in this realm?





Enhancement of Q_S with A: saturation regime reached at significantly lower energy in nuclei (and lower cost)

 $\ln Q^2$

Key Measurements - Diffraction

Diffractive physics will be *the* major component of the e+A program at an EIC



Close relative of DIS

t : momentum transfer squared
 M_X : mass of diffractive final-state
 characterized by large rapidity gap
 mediated by color neutral exchange (e.g. 2
 or more gluons)

- Came into limelight with discovery at HERA: diffractive events (~15% of total DIS rate)
- High sensitivity to gluon density:
 σ ~ [g(x,Q²)]² due to color-neutral exchange
- Only known process where spatial gluon distributions can be extracted



Issue: Missing Event Generators

- Requirements:
 - DIS Events
 - Diffractive Events

- saturation and non-saturation picture
- e+p (cross checks with existing data)
- e+A incl. nuclear effects
- Plethora of e+p generators from HERA times
 - PYTHIA6, LEPTO, CASCADE, HERACLES, DJANGOH, ...
- Many p+p/p+A/A+A generators (RHIC, LHC)
 - ► HIJING, DPMJET, EPOS, ...
- e+A

Issue: Missing Event Generators

- Requirements:
 - DIS Events
 - Diffractive Events

- saturation and non-saturation picture
- e+p (cross checks with existing data)
- e+A incl. nuclear effects
- Plethora of e+p generators from HERA times
 - PYTHIA6, LEPTO, CASCADE, HERACLES, DJANGOH, ...
- Many p+p/p+A/A+A generators (RHIC, LHC)
 - ► HIJING, DPMJET, EPOS, ...
- e+A

▶ ...

- EIC community \Rightarrow do-it-yourself
 - Sartre 1 for exclusive diffractive VM production (T. Toll and TU)
 - Sartre 2 for inclusive diffractive events (T. Toll)
 - DIS event generator pending
 - o currently: patch of PYTHIA6+DPMJet+Fluka

Incoherent: nucleus dissociation ($f \neq i$)

$$egin{aligned} \sigma_{ ext{incoh}} &\propto & \sum_{f
eq i} \langle i|\mathcal{A}|f
angle^{\dagger} \langle f|\mathcal{A}|i
angle \ &= & \sum_{f} \langle i|\mathcal{A}|f
angle^{\dagger} \langle f|\mathcal{A}|i
angle - \langle i|A|i
angle^{\dagger} \langle i|A|i
angle \ &= & \langle i|\left|\mathcal{A}
ight|^{2}|i
angle - \left|\langle i|\mathcal{A}|i
angle
ight|^{2} &= \langle |\mathcal{A}|^{2}
angle - \left|\langle \mathcal{A}
angle
ight|^{2} \end{aligned}$$

Incoherent CS is variance of amplitude ⇒ measure of fluctuating source density Coherent CS reflects the average source density

$$\frac{\mathrm{d}\sigma_{\mathrm{total}}}{\mathrm{d}t} = \frac{1}{16\pi} \langle |\mathcal{A}|^2 \rangle$$





Use Dipole Model: bSat and bNonSat

$$\Delta = \sqrt{-t}$$





Sartre: Wood-Saxon

Nucleus:

- small x \Rightarrow large $\lambda \Rightarrow$ coherently probes whole nucleus for x $\ll A^{-1/3} m_N/R_N \sim 10^{-2}$
- Position of nucleon in nucleus is not an observable
- To calculate CS need to average over all possible states of nucleon configurations Ω

$$\frac{\mathrm{d}\sigma^{\gamma^* \mathrm{A}}}{\mathrm{d}t}(x, Q^2, t) = \frac{1}{16\pi} \frac{1}{C_{\max}} \sum_{i=1}^{C_{\max}} \left| \mathcal{A}(x, Q^2, t, \Omega_i) \right|^2$$

Use Dipole Model: bSat and bNonSat

$$\Delta = \sqrt{-t}$$



Amplitude in *ep*:

$$\stackrel{\rightarrow Vp}{\rightarrow} (x, Q, \Delta) = i \int dr \int \frac{dz}{4\pi} \int d^2 \mathbf{b} (\Psi_V^* \Psi) (r, z)$$

$$\times 2\pi r J_0 ([1-z]r\Delta) e^{-i\mathbf{b}\cdot\Delta} \frac{d\sigma_{q\bar{q}}^{(p)}}{d^2\mathbf{b}} (x, r, \mathbf{b})$$



Sartre: Wood-Saxon

- Analytic average not exact
- Need C_{max} ~ 800, more at large t

$$\frac{\mathrm{d}\sigma^{\gamma^* \mathrm{A}}}{\mathrm{d}t}(x, Q^2, t) = \frac{1}{16\pi} \frac{1}{C_{\max}} \sum_{i=1}^{C_{\max}} \left| \mathcal{A}(x, Q^2, t, \Omega_i) \right|^2$$





Dependence on nucleon configurations in the amplitude is entirely contained in this dipole cross-section.

bSat:
$$\frac{1}{2} \frac{\mathrm{d}\sigma_{q\bar{q}}^{(A)}}{\mathrm{d}^{2}\mathbf{b}}(x, r, \mathbf{b}, \Omega) = 1 - \exp\left(-\frac{\pi^{2}}{2N_{C}}r^{2}\alpha_{S}(\mu^{2})xg(x, \mu^{2})\sum_{i=1}^{A}T(|\mathbf{b} - \mathbf{b}_{i}|)\right)$$

bNonSat:
$$\frac{\mathrm{d}\sigma_{q\bar{q}}^{(A)}}{\mathrm{d}^{2}b} = \frac{\pi^{2}}{N_{C}}r^{2}\alpha_{s}(\mu^{2})xg(x, \mu^{2})\sum_{i=1}^{A}T(|\mathbf{b} - \mathbf{b}_{i}|)$$



Wave overlap function $\Psi^*\Psi$ falls steeply for large dipole radii

- J/ψ not sensitive to saturation.
- Need to look at φ, or ρ that "see" more of the dipole amplitude

$$\mathcal{A}_{T,L}^{\gamma^* p \to V p}(x, Q, \Delta) = i \int \mathrm{d}r \int \frac{\mathrm{d}z}{4\pi} \int \mathrm{d}^2 \mathbf{b} (\Psi_V^* \Psi) (r, z)$$
$$\times 2\pi r J_0([1-z]r\Delta) e^{-i\mathbf{b}\cdot\Delta} \frac{\mathrm{d}\sigma_{q\bar{q}}^{(p)}}{\mathrm{d}^2 \mathbf{b}} (x, r, \mathbf{b})$$





Wave overlap function $\Psi^*\Psi$ falls steeply for large dipole radii

- J/ψ not sensitive to saturation.
- Need to look at φ, or ρ that "see" more of the dipole amplitude

$$\mathcal{A}_{T,L}^{\gamma^* p \to V p}(x, Q, \Delta) = i \int \mathrm{d}r \int \frac{\mathrm{d}z}{4\pi} \int \mathrm{d}^2 \mathbf{b} (\Psi_V^* \Psi) (r, z)$$
$$\times 2\pi r J_0([1-z]r\Delta) e^{-i\mathbf{b}\cdot\Delta} \frac{\mathrm{d}\sigma_{q\bar{q}}^{(p)}}{\mathrm{d}^2 \mathbf{b}} (x, r, \mathbf{b})$$





Wave overlap function $\Psi^*\Psi$ falls steeply for large dipole radii

- J/ψ not sensitive to saturation.
- Need to look at φ, or ρ that "see" more of the dipole amplitude

$$\mathcal{A}_{T,L}^{\gamma^* p \to V p}(x, Q, \Delta) = i \int \mathrm{d}r \int \frac{\mathrm{d}z}{4\pi} \int \mathrm{d}^2 \mathbf{b} (\Psi_V^* \Psi) (r, z)$$
$$\times 2\pi r J_0([1-z]r\Delta) e^{-i\mathbf{b}\cdot\Delta} \frac{\mathrm{d}\sigma_{q\bar{q}}^{(p)}}{\mathrm{d}^2 \mathbf{b}} (x, r, \mathbf{b})$$



EIC: Exclusive Vector Meson Production



- Sartre event generator (bSat & bNonSat = linearized bSat)
- As expected: big difference for ϕ less so for J/ψ
- Note: A^{4/3} scaling strictly only valid at large Q²

EIC: Exclusive Vector Meson Production



- Sartre event generator (bSat & bNonSat = linearized bSat)
- As expected: big difference for ϕ less so for J/ψ
- Note: A^{4/3} scaling strictly only valid at large Q²

Exclusive Diffractive VM Production: $d\sigma/dt$

- In general in e+A: cannot detect the outgoing nucleus and measure its momentum
- Exception: exclusive VM production
 - Need only to measure e' and the VM



$$\begin{split} t &= (p_A - p_{A'})^2 = (p_{\rm VM} + p_{e'} - p_e)^2 \\ &\approx (p_T^{e'} + p_T^{\rm VM})^2 \end{split}$$

Vector Meson Production: do/dt



- Find: Typical diffractive pattern for coherent (non-breakup) part
- As expected: J/ψ less sensitive to saturation than ϕ
- Need this sliced in x bins \Rightarrow luminosity hungry
- Crucial: t resolution and reach

Spatial Gluon Distribution from do/dt

Diffractive vector meson production: $e + Au \rightarrow e' + Au' + J/\psi$

• Momentum transfer $t = |\mathbf{p}_{Au} - \mathbf{p}_{Au'}|^2$ conjugate to b_T



- Converges to input F(b) rapidly: |t| < 0.1 almost enough
- Fourier transformation requires $\int \mathcal{L} dt > 1 \text{ fb}^{-1}/A$

Spatial Gluon Distribution from do/dt



- J/ψ perfect for obtaining F(b) in both cases sat and non-sat
- less so since coherence distorts F(b)
- Note: difference in F(b) of ϕ and J/ψ reveals saturation

Spatial Gluon Distribution from do/dt



- J/ψ perfect for obtaining F(b) in both cases sat and non-sat
- less so since coherence distorts F(b)
- Note: difference in F(b) of ϕ and J/ ψ reveals saturation

e+A: Diffractive over Total Cross-Section

- Predicted to be enhanced in e+A compared to e+p at large β, i.e. small M_X² (β = xQ²(Q²-M_X²))
 - β = momentum fraction of the struck parton with respect to the Pomeron



Kowalski et al. Phys.Rev. C78 (2008) 045201





Inclusive Diffractive Events: Sartre 2

Sartre 2

- Allows for simulation of inclusive diffractive events
- Dipole model: $qar q, qar q g, \ldots$
- States handed over to Pythia 8 for parton showering





Simulations show little dependence on W (as observed at HERA)

Sartre 2: Diffractive over Total Cross-Section



Saturation model calculations (Sartre event generator & analytic calculations by T. Lappi) now include $\overline{q}qg$ that affect the ratio at low β . Sat-simulations describe HERA results in ep.

Sartre 2: Diffractive over Total Cross-Section



Saturation model calculations (Sartre event generator & analytic calculations by T. Lappi) now include $\overline{q}qg$ that affect the ratio at low β . Sat-simulations describe HERA results in ep.

Take Away Message

The diffractive physics e+A program at an EIC is unprecedented, allowing the study of matter in a new emergent regime.

- We developed a generator for diffractive events
 - in e+A and e+p
 - for saturation and non-saturation
 - exclusive vector mesons in Sartre 1
 - inclusive diffractive events in Sartre 2 (in progress)
- It helped us to simulate key measurements at an EIC
 - Indispensable for optimization of detector and IR design
- More work to do on DIS generators for e+A