

TMDs at small-x: What we learnt and What to expect

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Nice things about transverse momentum distributions (TMDs)

- Universality and a universal language
 - DIS/Drell-Yan
 - eA/pA/AA(?), small-x wave functions of nucleus
- QCD dynamics
 - TMD evolution
 - Small-x evolution

QCD evolution at high energy

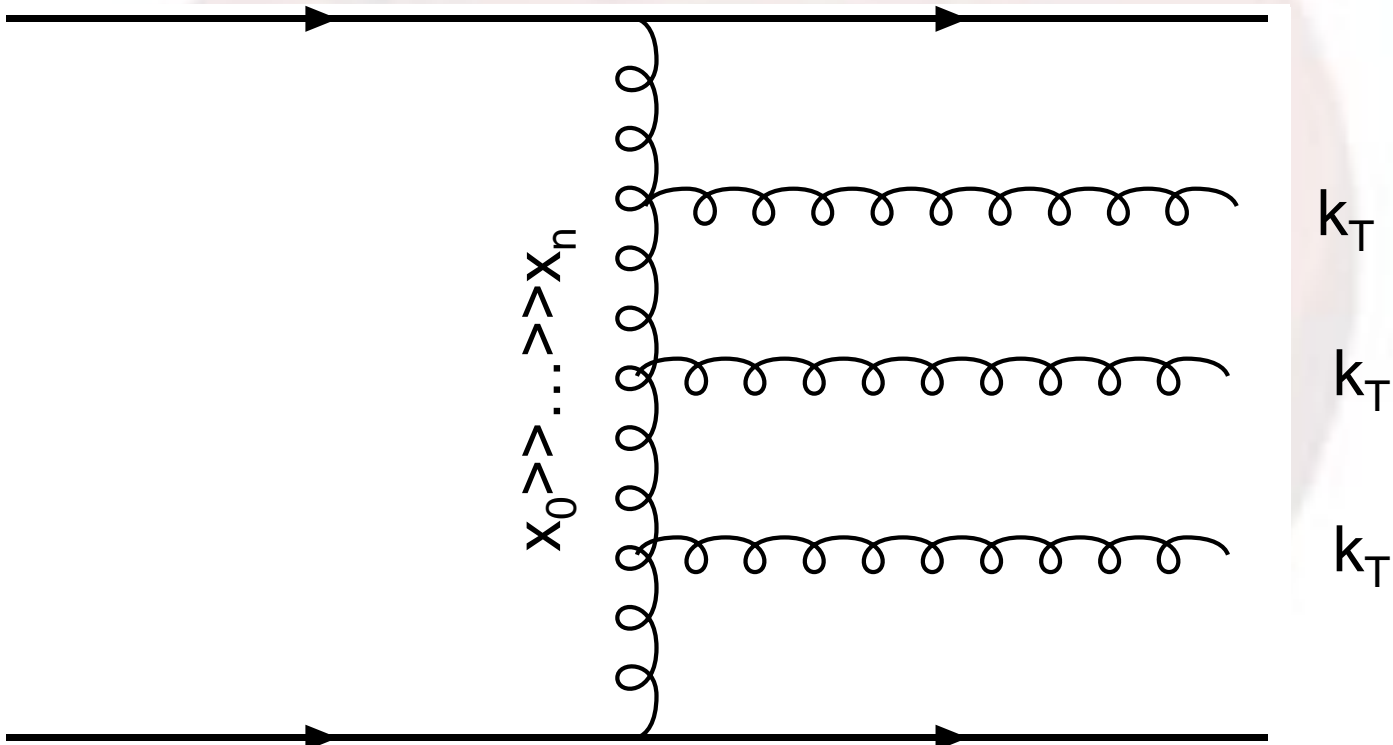
- BFKL/BK-JIMWLK (small-x)
- Sudakov (TMD)

Mueller-Xiao-Yuan 2013

Balitsky-Tarasov 2014

Tarasov's talk on Wednesday

High energy scattering



BFKL: $\frac{\partial \mathcal{F}}{\partial \ln(1/x)} = \mathcal{K} \otimes \mathcal{F}$ Un-integrated gluon distribution

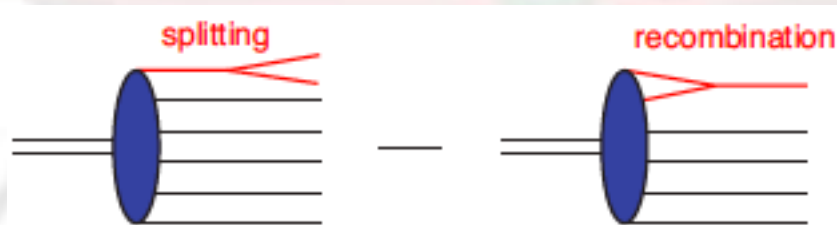
Non-linear term at high density

- Balitsky-Fadin-Lipatov-Kuraev, 1977-78

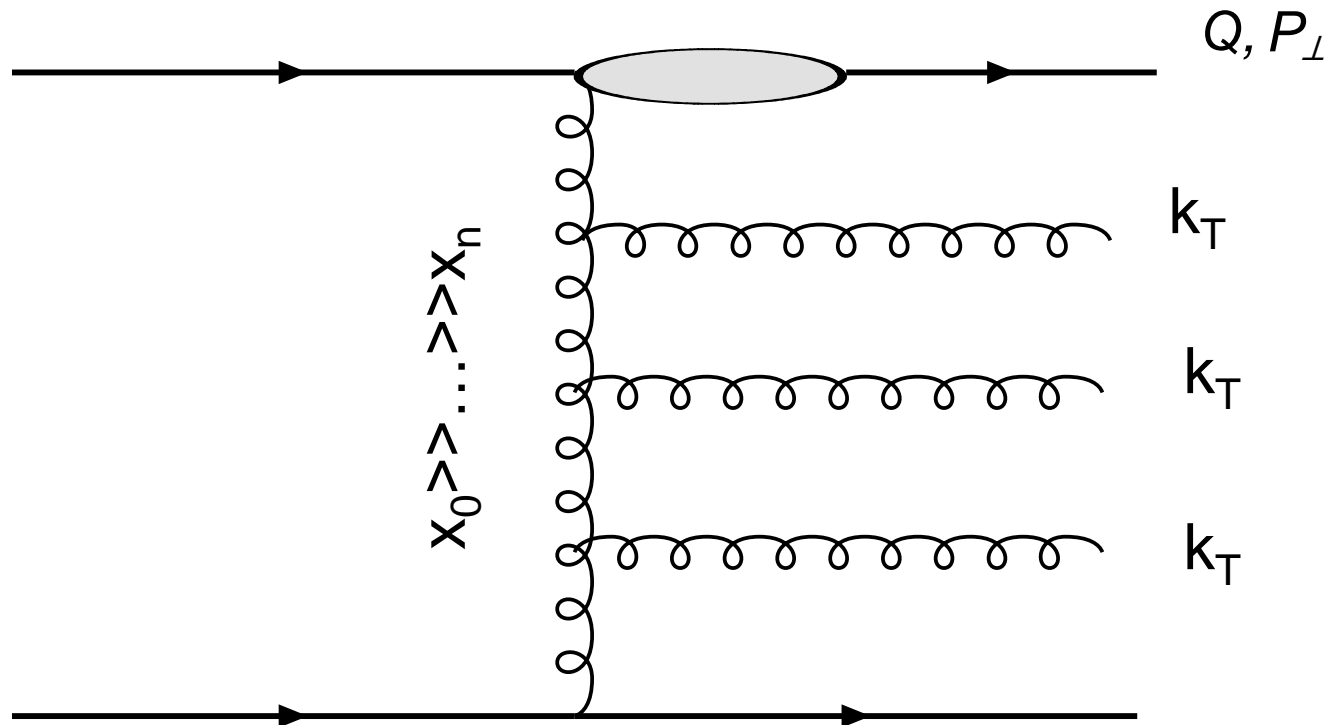
$$\frac{\partial N(x, r_T)}{\partial \ln(1/x)} = \alpha_s K_{\text{BFKL}} \otimes N(x, r_T)$$

- Balitsky-Kovchegov: Non-linear term, 98

$$\frac{\partial N(x, r_T)}{\partial \ln(1/x)} = \alpha_s K_{\text{BFKL}} \otimes N(x, r_T) - \alpha_s [N(x, r_T)]^2.$$



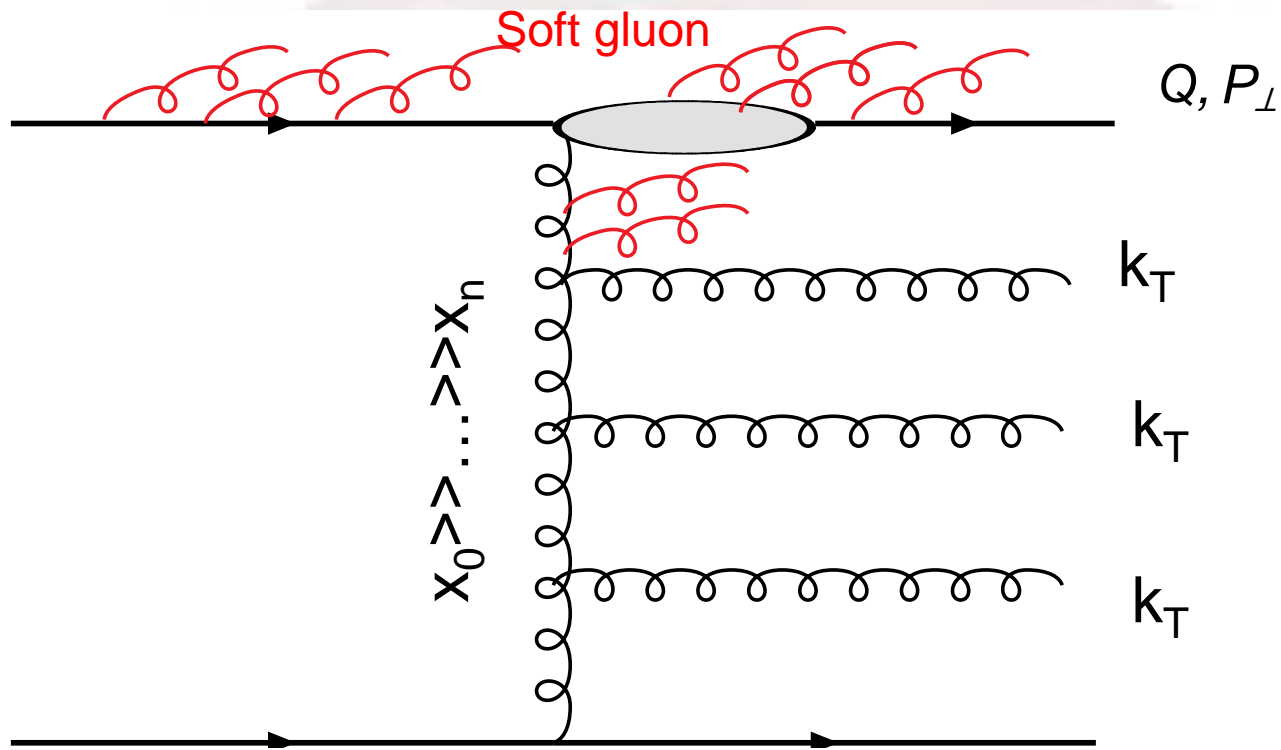
Hard processes at small-x



- Manifest dependence on un-integrated gluon distributions

□ Dominguez-Marquet-Xiao-Yuan, 2010

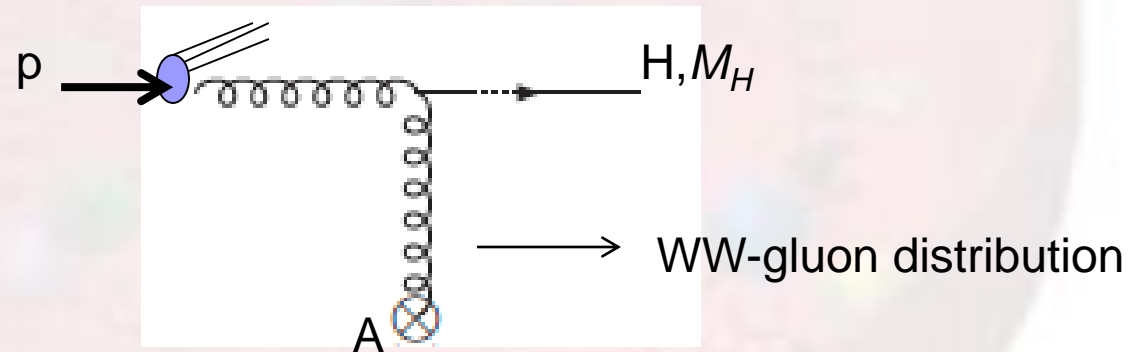
Additional dynamics comes in



- BFKL vs **Sudakov** resummations (LL)

Sudakov resummation at small- x

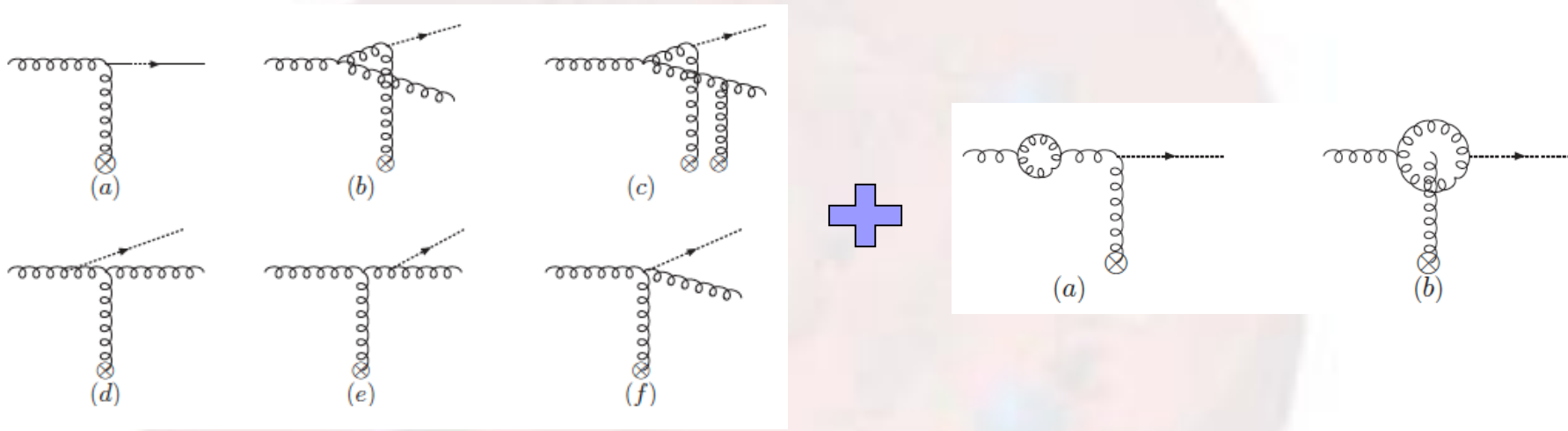
- Take massive scalar particle production $p+A \rightarrow H+X$ as an example to demonstrate the double logarithms, and resummation



$$\frac{d\sigma^{(\text{LO})}}{dy d^2k_{\perp}} = \sigma_0 \int \frac{d^2x_{\perp} d^2x'_{\perp}}{(2\pi)^2} e^{ik_{\perp} \cdot r_{\perp}} x_0 g_p(x_0) S^{(WW)}(x_{\perp}, x'_{\perp})$$

$$S_Y^{WW}(x_{\perp}, y_{\perp}) = - \left\langle \text{Tr} \left[\partial_{\perp}^{\beta} U(x_{\perp}) U^{\dagger}(y_{\perp}) \partial_{\perp}^{\beta} U(y_{\perp}) U^{\dagger}(x_{\perp}) \right] \right\rangle_Y$$

Explicit one-loop calculations



$$x_0 g_p(x_0) \int \frac{d\xi}{\xi} \mathbf{K}_{DMMX} \otimes S^{WW}(x_\perp, y_\perp) + \left(-\frac{1}{\epsilon}\right) S^{WW}(x_\perp, y_\perp) \mathcal{P}_{g/g} \otimes x_0 g(x_0),$$

- Collinear divergence \rightarrow DGLAP evolution
- Small- x divergence \rightarrow BK-type evolution

Dominguiz-Mueller-Munier-Xiao, 2011

Soft vs Collinear gluons

- Radiated gluon momentum

$$k_g = \alpha_g p_1 + \beta_g p_2 + k_{g\perp} ,$$

- Soft gluon, $\alpha \sim \beta \ll 1$
- Collinear gluon, $\alpha \sim 1, \beta \ll 1$
- Small- x collinear gluon, $1 - \beta \ll 1, \alpha \rightarrow 0$
 - Rapidity divergence

Final result

- Double logs at one-loop order

$$\frac{d\sigma^{(\text{LO+NLO})}}{dyd^2k_{\perp}} \Big|_{k_{\perp} \ll Q} = \sigma_0 \int \frac{d^2x_{\perp} d^2x'_{\perp}}{(2\pi)^2} e^{ik_{\perp} \cdot r_{\perp}} S_{Y=\ln 1/x_a}^{WW}(x_{\perp}, x'_{\perp}) x g_p(x, \mu^2 = \frac{c_0^2}{r_{\perp}^2}) \left\{ 1 + \frac{\alpha_s}{\pi} C_A \left[\beta_0 \ln \frac{Q^2 r_{\perp}^2}{c_0^2} - \frac{1}{2} \left(\ln \frac{Q^2 r_{\perp}^2}{c_0^2} \right)^2 + \frac{\pi^2}{2} \right] \right\},$$

- Include both BFKL (BK) and Sudakov

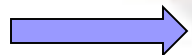
$$\frac{d\sigma^{(\text{resum})}}{dyd^2k_{\perp}} \Big|_{k_{\perp} \ll Q} = \sigma_0 \int \frac{d^2x_{\perp} d^2x'_{\perp}}{(2\pi)^2} e^{ik_{\perp} \cdot r_{\perp}} e^{-S_{\text{sud}}(Q^2, r_{\perp}^2)} S_{Y=\ln 1/x_a}^{WW}(x_{\perp}, x'_{\perp}) \times x g_p(x, \mu^2 = c_0^2/r_{\perp}^2) \left[1 + \frac{\alpha_s}{\pi} \frac{\pi^2}{2} N_c \right],$$

Sudakov leading double logs+small-x logs in hard processes

- Each incoming parton contributes to a half of the associated color factor in Sudakov
 - Initial gluon radiation, aka, TMDs

$$\frac{d\sigma}{dy_1 dy_2 dP_\perp^2 d^2 k_\perp} \propto H(P_\perp^2) \int d^2 x_\perp d^2 y_\perp e^{ik_\perp \cdot (x_\perp - y_\perp)} \widetilde{W}_{x_A}(x_\perp, y_\perp)$$

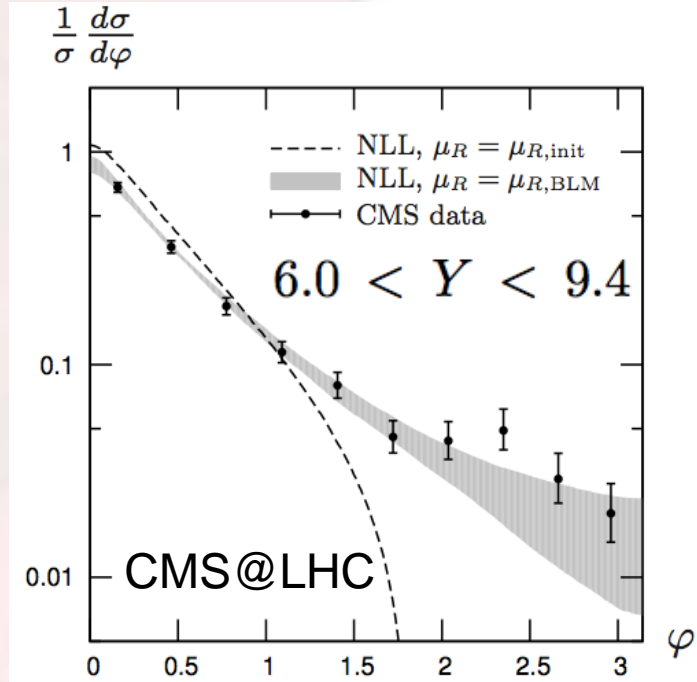
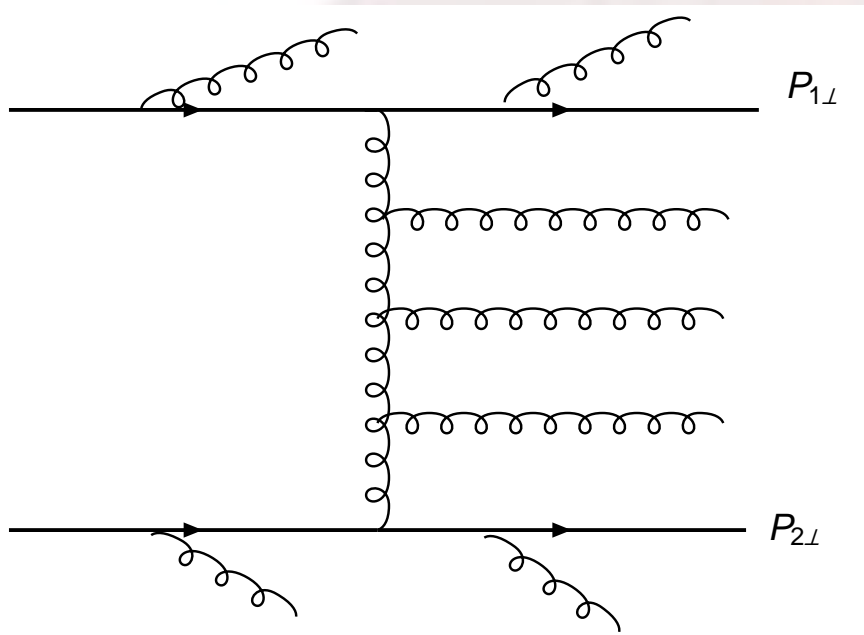
Sudakov



$$H(P_\perp^2) \int d^2 x_\perp d^2 y_\perp e^{ik_\perp \cdot R_\perp} e^{-S_{sud}(P_\perp, R_\perp)} \widetilde{W}_{x_A}(x_\perp, y_\perp)$$

Mueller-Xiao-Yuan 2013

Dijet with large rapidity gap



Sudakov resummation will dominate
Small angle distribution (Mueller-Xiao-Yuan)

Ducloue, Szymanowski, Wallon
1309.3229, only take into account
BFKL

Work in progress,

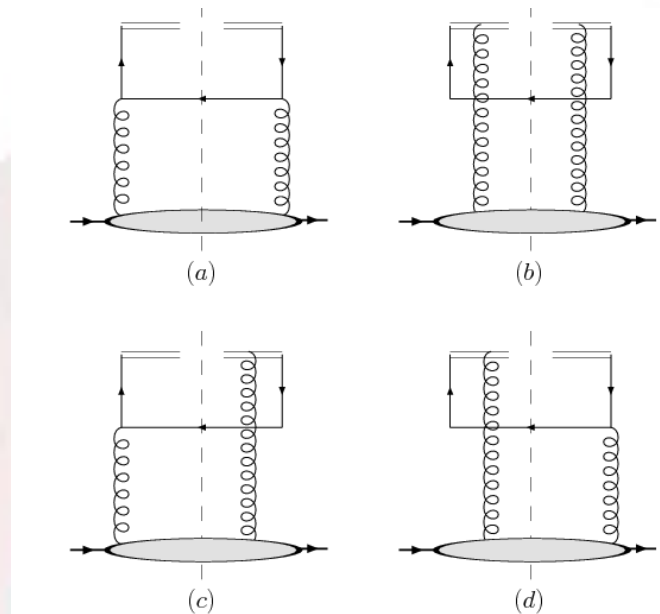
TMDs at small-x: what we do we know?

- We hope to extract these TMDs from experiments by applying the resummation formula

$$\frac{d\sigma}{dy_1 dy_2 dP_{\perp}^2 d^2k_{\perp}} \propto H(P_{\perp}^2) \int d^2x_{\perp} d^2y_{\perp} e^{ik_{\perp} \cdot R_{\perp}} e^{-S_{sud}(P_{\perp}, R_{\perp})} \widetilde{W}_{x_A}(x_{\perp}, y_{\perp})$$

- Small-x TMDs can be calculated from CGC, instead of parameterization
 - Prediction power

TMD quark at small- x

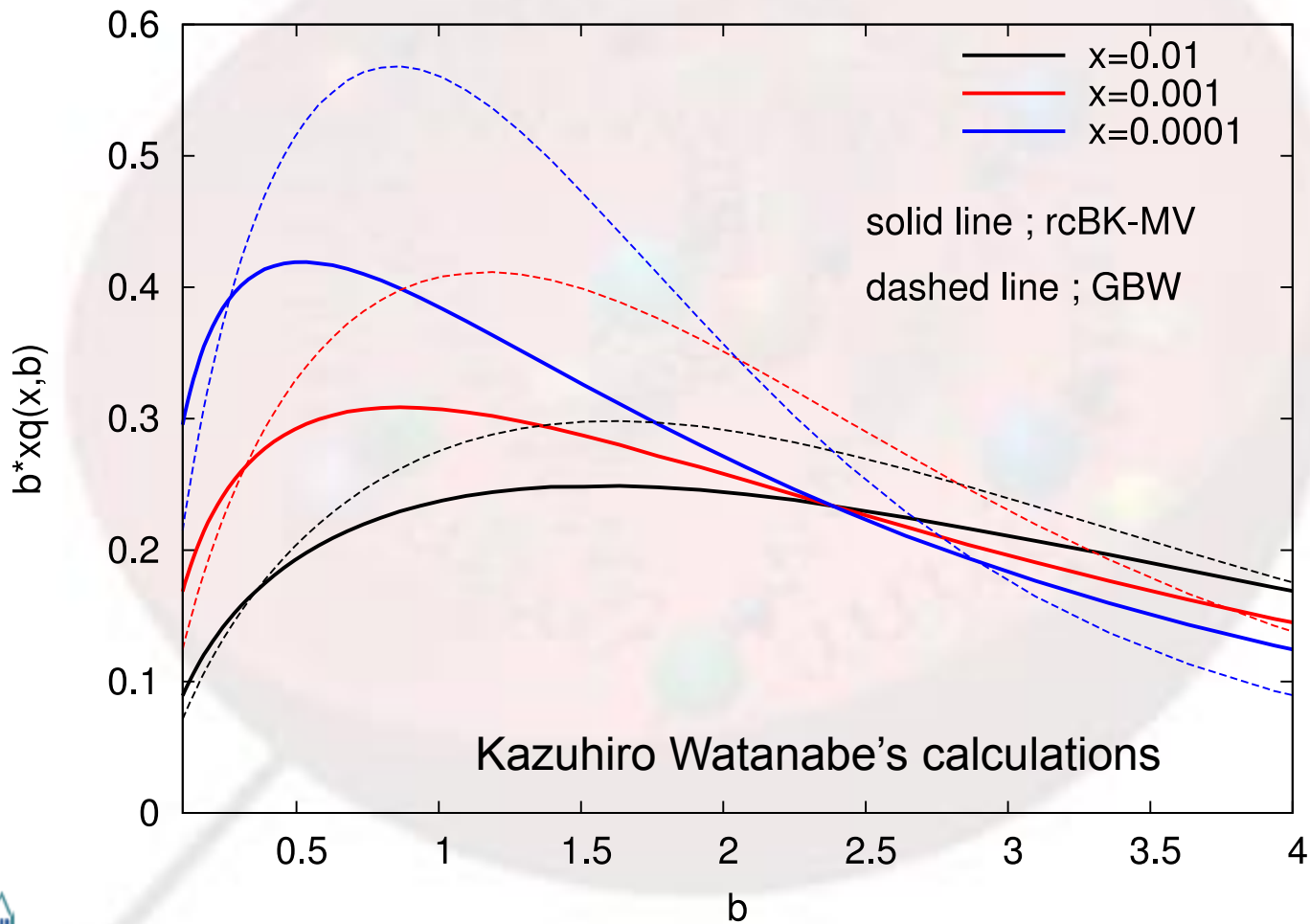


McLerran-Venugopalan 98

$$q(x, k_{\perp}) = \frac{N_c}{8\pi^4} \int \frac{dx'}{x'^2} \int d^2b d^2q_{\perp} F(q_{\perp}, x') A(q_{\perp}, k_{\perp})$$

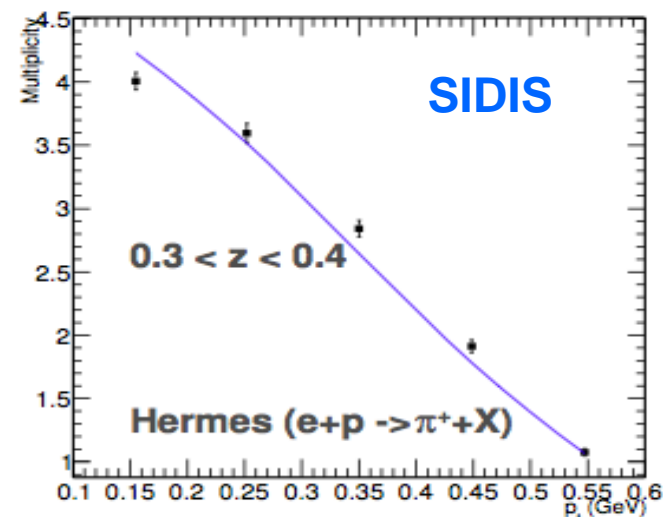
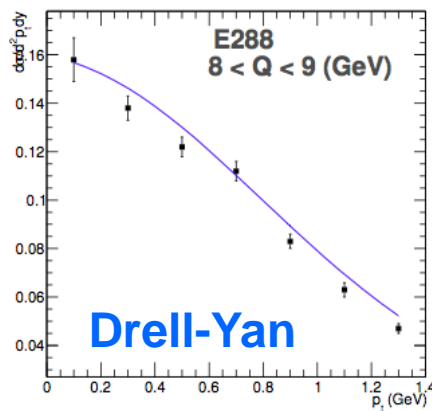
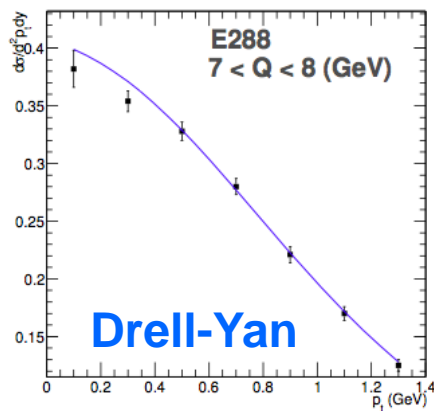
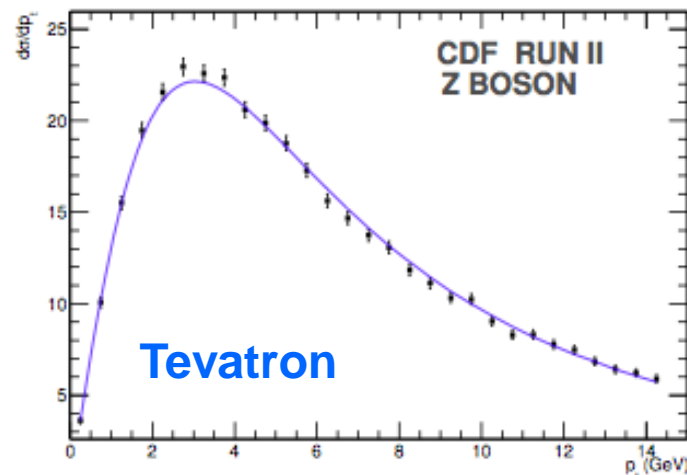
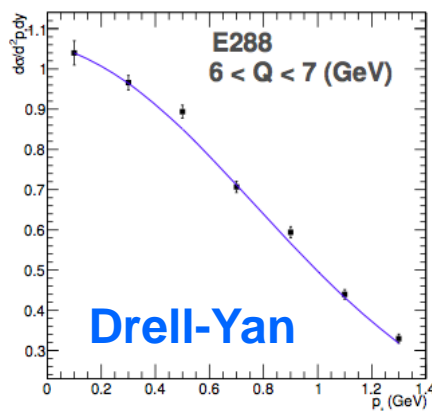
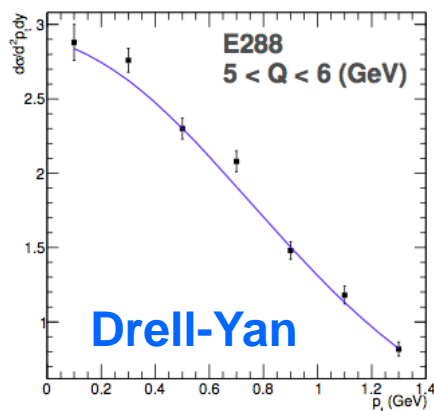
- Can be calculated from the dipole amplitude, and can be applied to DIS and Drell-Yan processes

TMD quark at small-x



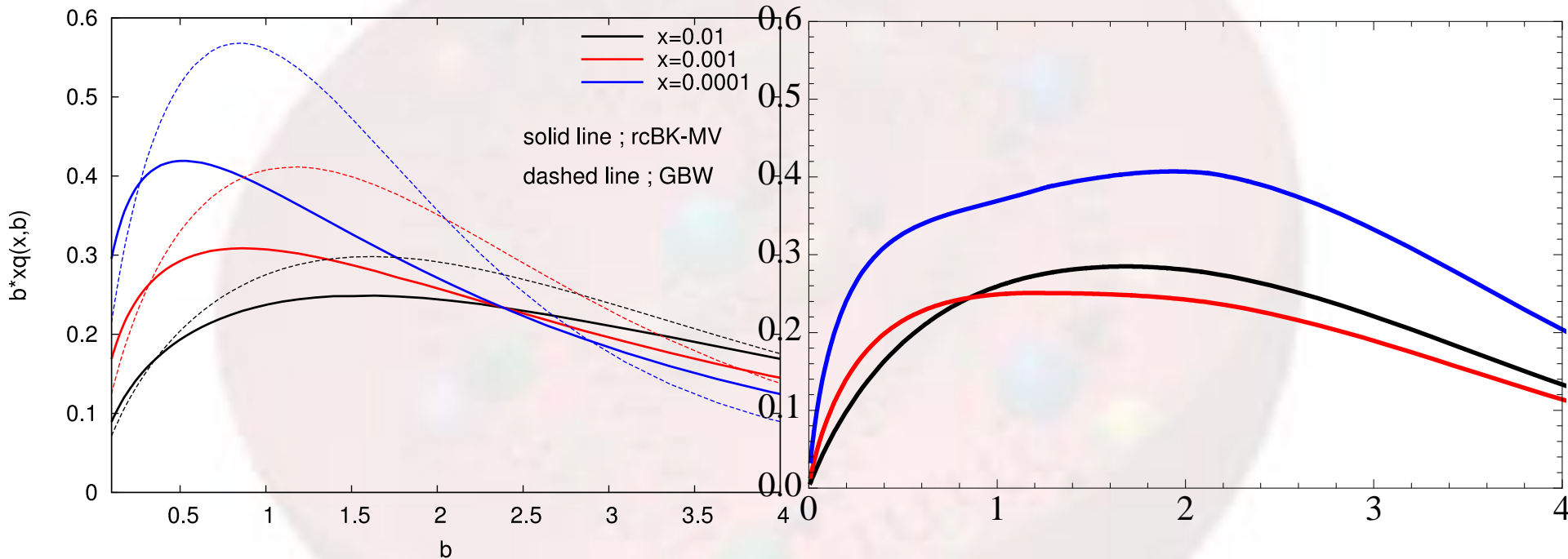
What we know the TMD quarks (not small-x)

Sun-Issacson-Yuan-Yuan, 2014

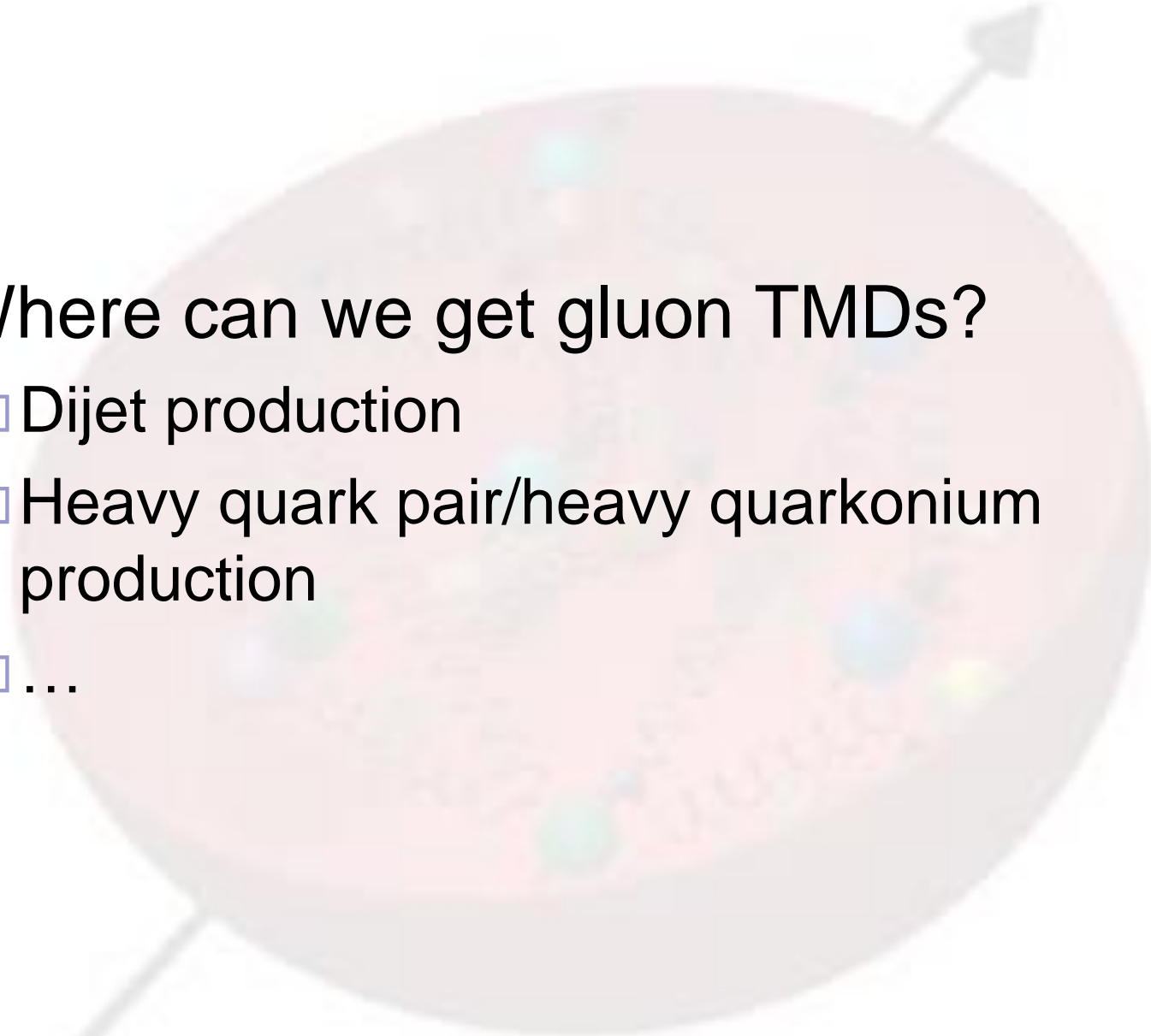


See also, BLNY 2002

TMD quark at small-x: CGC vs Collinear



- Realistic comparison will shed light on the TMD quarks at small-x (work in progress)

- 
- Where can we get gluon TMDs?
 - Dijet production
 - Heavy quark pair/heavy quarkonium production
 - ...

TMD gluon from quarkonium productions

Sun, C.-P. Yuan, F. Yuan, PRD 2013
earlier work: Bergers-Qiu-Wang, 2004

■ Low pt distribution

$$\frac{d\sigma}{d^2P_\perp dy} \Big|_{P_\perp \ll M} = \frac{1}{(2\pi)^2} \int d^2b e^{i\vec{P}_\perp \cdot \vec{b}} W(b, M, x_1, x_2)$$

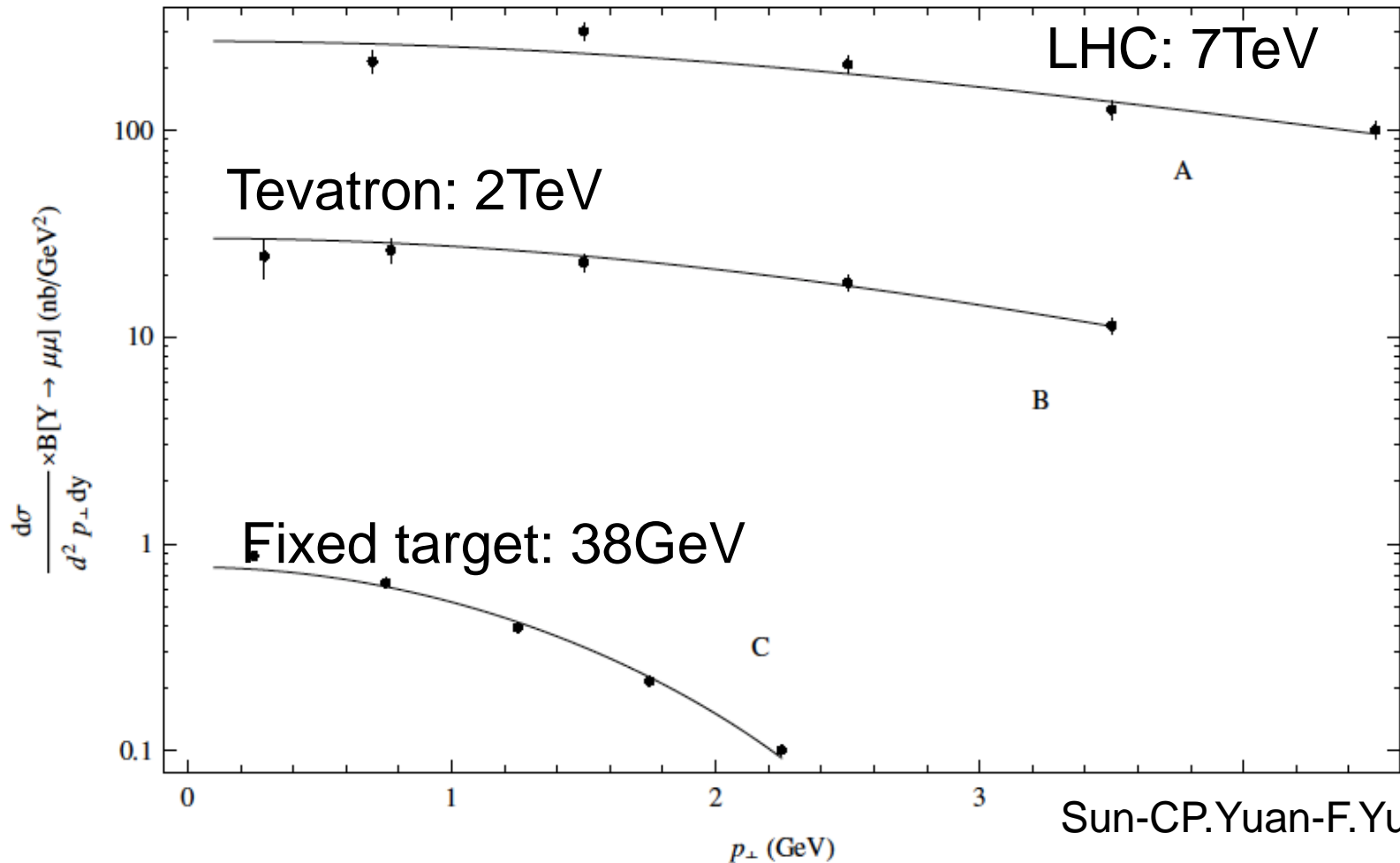
$$W(b, M^2) = e^{-S_{Sud}(M^2, b, C_1, C_2)} W(b, C_1, C_2)$$

■ Non-perturbative input

$$W(b) = W(b_*) W^{NP}(b)$$

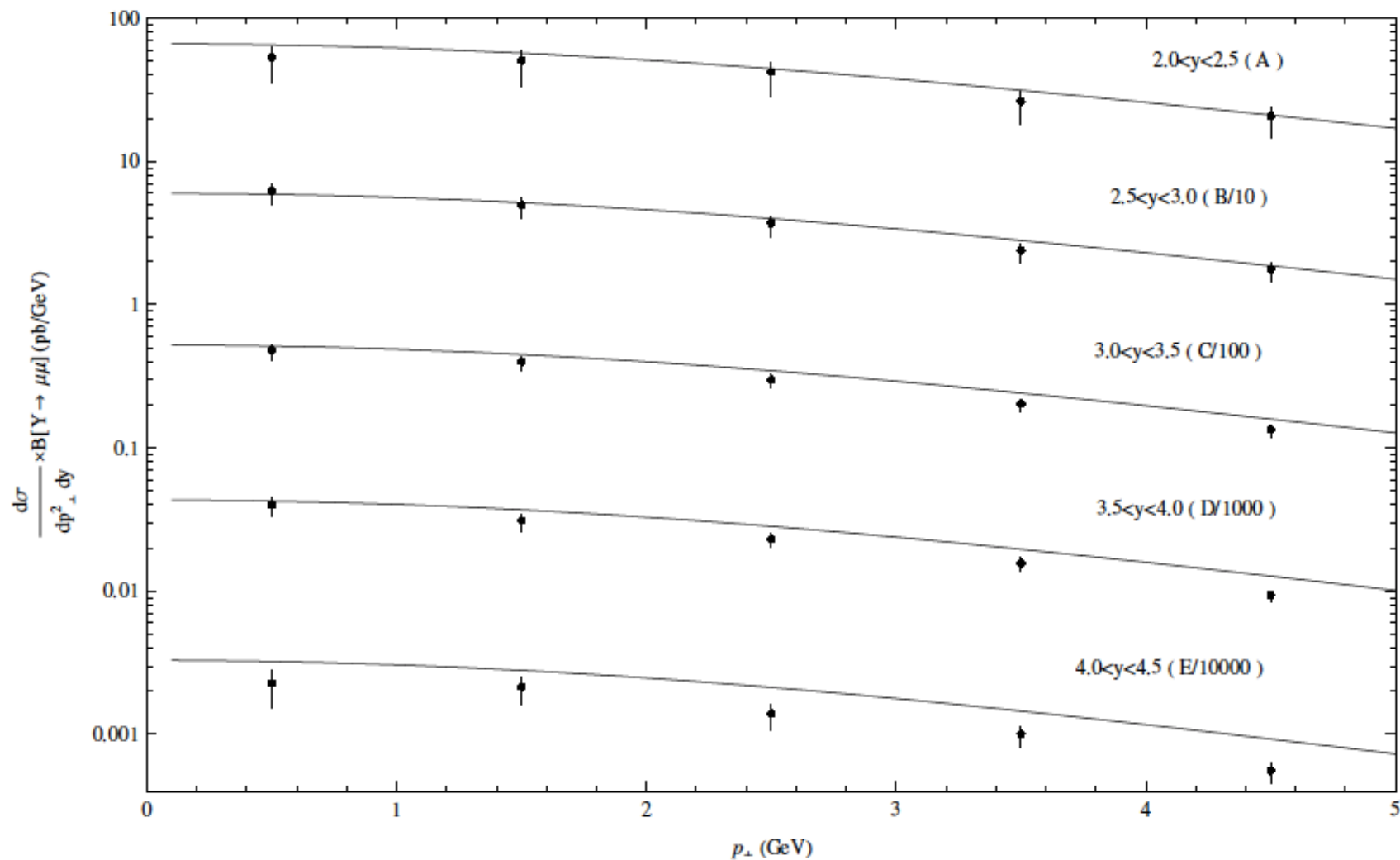
$$W^{NP}(b) = \exp \left[-g_1 - g_2 \ln \left(\frac{Q}{2Q_0} \right) - g_1 g_3 \ln(100x_1 x_2) \right] b^2$$

Quarkonium production as a probe to the gluon TMDs (Y)



Sun-CP.Yuan-F.Yuan, 2013

More data from LHC



Gluon TMD: CGC vs Collinear

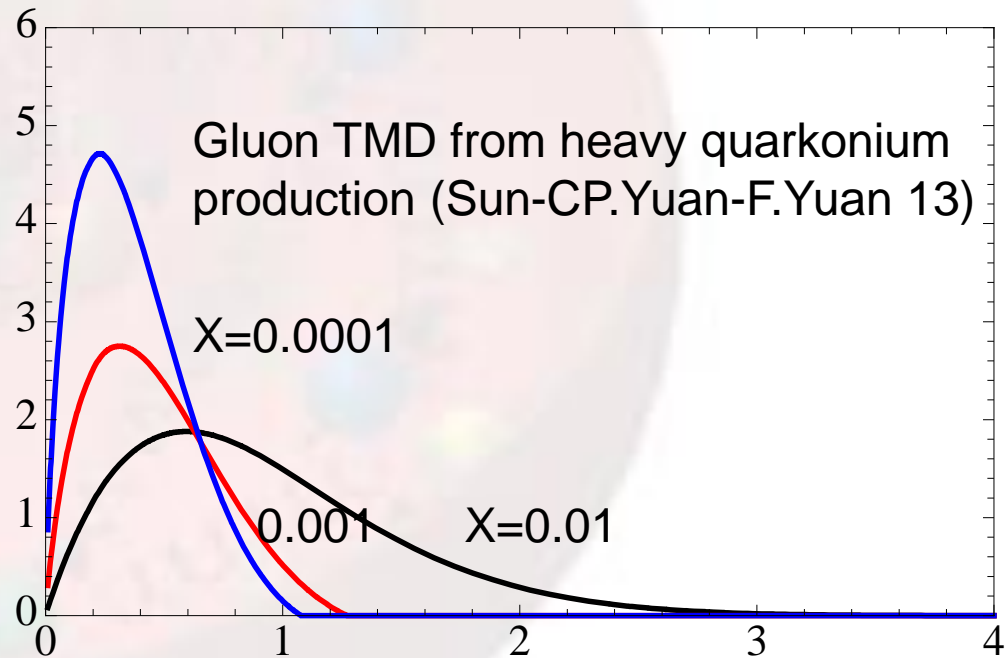
Model-dependent way
??

Qiu-Sun-Xiao-Yuan 14

Ma-Venugopalan 14/15

Ducloué-Lappi-Mäntysaari 15

(work in progress)



Conclusion

- We start to have a quantitative picture on TMDs at small- x
 - Which connect different physics
- More theoretical works are needed
- Experiments are crucial to answer some of the important questions
 - LHC, and EIC