



# Pseudo-observables in Higgs decays

Admir Greljo

based on:

- (1) Gonzalez-Alonso, AG, Isidori, Marzocca, **Eur. Phys. J. C75 (2015) 3, 128**
- (2) Gonzalez-Alonso, AG, Isidori, Marzocca, **work in progress**

Portorož 2015 workshop, 08/04/2015

The idea is to provide **general** and yet **simple** parameterisation of New Physics (NP) effects in Higgs decays...

# Outline

- “kappa” formalism and beyond
- **P**seudo-**O**bservables (**POs**) in Higgs decays
- Parameter counting
- Symmetry limits
- Higgs **POs** from LEP data in the linear EFT

# “kappa” formalism

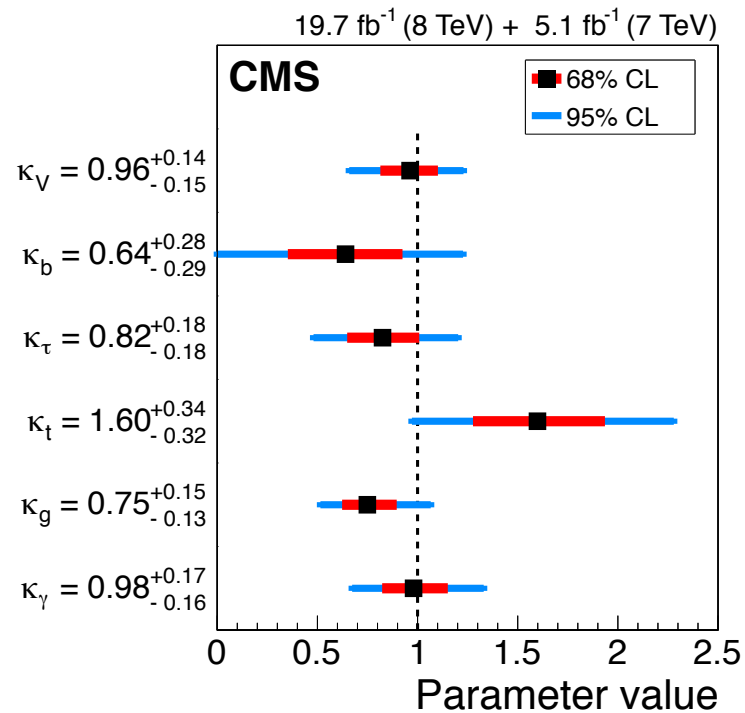
- Interim framework for the analysis of Higgs couplings  
(CERN Yellow Report 3, 1307.1347)

- Single narrow resonance with a mass  $\sim 125$  GeV
- The production and decay kinematics are SM-like
- Introduce “coupling strength” scaling factors  $\kappa_i$ , i.e.,

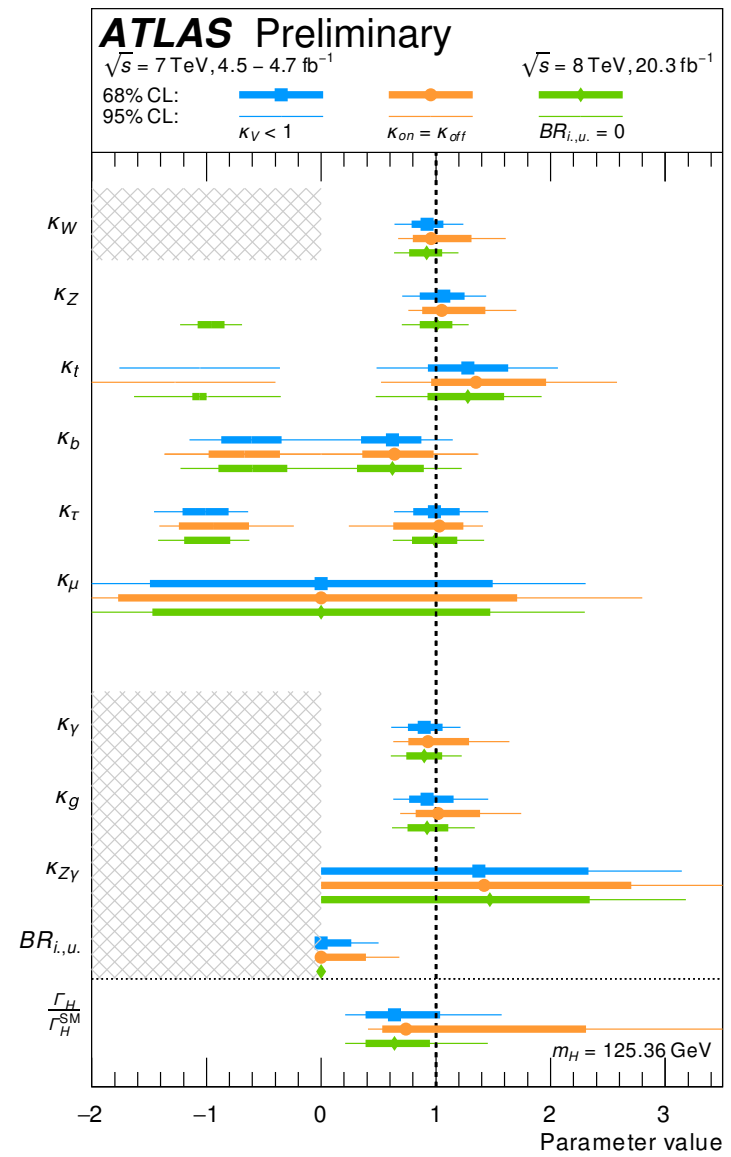
$$(\sigma \cdot \text{BR})(gg \rightarrow H \rightarrow \gamma\gamma) = \sigma_{\text{SM}}(gg \rightarrow H) \cdot \text{BR}_{\text{SM}}(H \rightarrow \gamma\gamma) \cdot \frac{\kappa_g^2 \cdot \kappa_\gamma^2}{\kappa_H^2}$$

- SM limit recovered for  $\kappa_i \rightarrow 1$

# “kappa” formalism



CMS-HIG-14-009



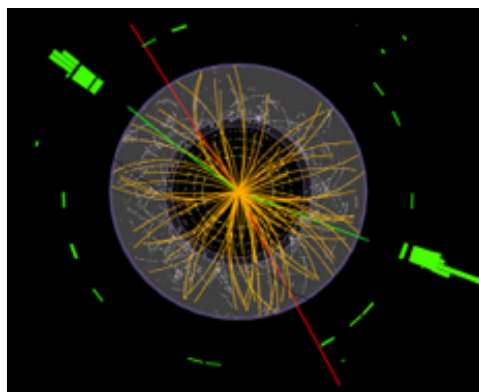
ATLAS-CONF-2015-007

# Beyond “kappa” formalism

- Higgs physics is entering a precision era
- LHC Run II: Precise measurements of the observed Higgs processes
- $\mathcal{O}(100)$  events expected soon in  $h \rightarrow 4\ell$
- Exploit full kinematics of the events  
(not only the total rate)
- ***Extended “kappa” formalism needed***

# Pseudo-Observables

Marzocca, HXSWG plenary meeting - Jan 2015



$$\begin{aligned}\mathcal{L} = & -\frac{1}{4} F_{\mu\nu} F^{\mu\nu} \\ & + i\bar{\Psi}\not{\partial}\Psi + h.c. \\ & + \bar{\Psi}_i y_{ij} \Psi_j \phi + h.c. \\ & + \frac{1}{2} \partial_\mu \phi^2 - V(\phi)\end{aligned}$$



Experimental  
data

*Fiducial cross sections,  
distributions,  
...*

Pseudo  
Observables

Lagrangian  
parameters

*Couplings,  
bare masses,  
Wilson coeff.,  
...*

- **POs** are limited set of “idealised” observables defined from “on-shell” amplitudes
- Ideally, **POs** should:
  1. provide general encoding of the experimental data
  2. be computable and encode all possible predictions of large set of theories

# POs in Higgs decays

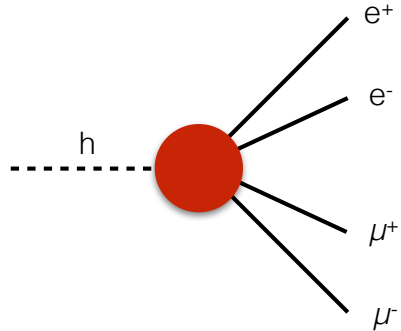
## Assumptions:

- $h(125)$  is a spin 0 particle
- Zero width approximation, “on-shell” single Higgs processes  
(factorisation of new physics effects in production and decay)
- No light New Physics  
Notion of an underlining Effective Field Theory (EFT), smooth kinematical distortions from the SM,  
Momentum expansion of the on-shell Higgs amplitudes.
- Power counting based on the canonical dimension:  
(1) Higgs, gauge bosons, derivatives (momenta)  $\sim 1$   
(2) Fermions  $\sim 3/2$
- Neglect contributions with dimension  $> 6$
- General enough to accommodate all the effect from next-to-leading order terms in the expansions of a generic linear **and** non-linear EFT
- No assumptions on custodial symmetry, CP, flavour universality or  $SU(2)_L$  properties of the Higgs, **we rather want to test it from data!**



# $h \rightarrow e^+ e^- \mu^+ \mu^-$

The most generic structure of the amplitude (helicity conserving):

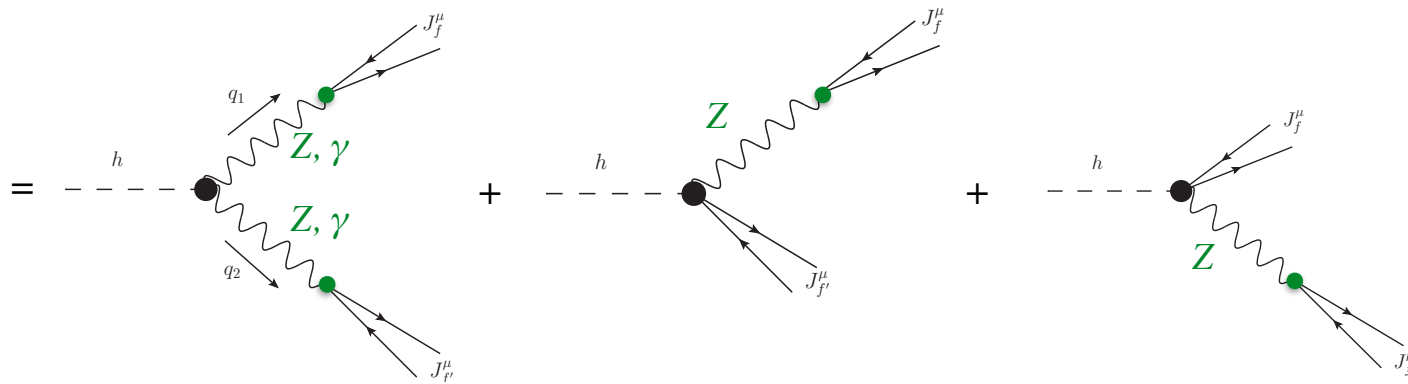


$$\mathcal{A} = i \frac{2m_Z^2}{v_F} \sum_{e=e_L, e_R} \sum_{\mu=\mu_L, \mu_R} (e \gamma_\alpha e) (\bar{\mu} \gamma_\beta \mu) \times \left[ F_1^{e\mu}(q_1^2, q_2^2) g^{\alpha\beta} + F_3^{e\mu}(q_1^2, q_2^2) \frac{q_1 \cdot q_2 g^{\alpha\beta} - q_2^\alpha q_1^\beta}{m_Z^2} + F_4^{e\mu}(q_1^2, q_2^2) \frac{\epsilon^{\alpha\beta\rho\sigma} q_{2\rho} q_{1\sigma}}{m_Z^2} \right]$$

Marzocca, HXSWG plenary meeting - Jan 2015

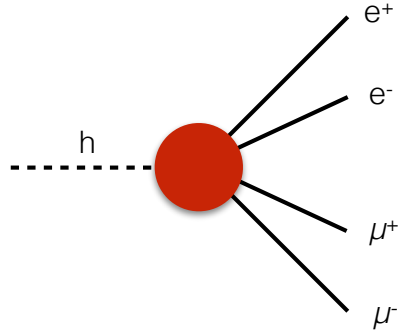
Gonzalez-Alonso, AG, Isidori, Marzocca

Momentum expansion of the form factors around the physical poles



# $h \rightarrow e^+ e^- \mu^+ \mu^-$

The most generic structure of the amplitude:



$$A = i \frac{2m_Z^2}{v_F} \sum_{e=e_L, e_R} \sum_{\mu=\mu_L, \mu_R} (\bar{e} \gamma_\alpha e) (\bar{\mu} \gamma_\beta \mu) \times \left[ F_1^{e\mu}(q_1^2, q_2^2) g^{\alpha\beta} + F_3^{e\mu}(q_1^2, q_2^2) \frac{q_1 \cdot q_2 g^{\alpha\beta} - q_2^\alpha q_1^\beta}{m_Z^2} + F_4^{e\mu}(q_1^2, q_2^2) \frac{\varepsilon^{\alpha\beta\rho\sigma} q_{2\rho} q_{1\sigma}}{m_Z^2} \right]$$

Momentum expansion of the form factors around the physical poles

$$A = i \frac{2m_Z^2}{v_F} \sum_{e=e_L, e_R} \sum_{\mu=\mu_L, \mu_R} (\bar{e} \gamma_\alpha e) (\bar{\mu} \gamma_\beta \mu) \times \left[ \left( \kappa_{ZZ} \frac{g_Z^e g_Z^\mu}{P_Z(q_1^2) P_Z(q_2^2)} + \frac{\epsilon_{Ze}}{m_Z^2} \frac{g_Z^\mu}{P_Z(q_2^2)} + \frac{\epsilon_{Z\mu}}{m_Z^2} \frac{g_Z^e}{P_Z(q_1^2)} \right) g^{\alpha\beta} + \left( \epsilon_{ZZ} \frac{g_Z^e g_Z^\mu}{P_Z(q_1^2) P_Z(q_2^2)} + \kappa_{Z\gamma} \epsilon_{Z\gamma}^{\text{SM-1L}} \left( \frac{e Q_\mu g_Z^e}{q_2^2 P_Z(q_1^2)} + \frac{e Q_e g_Z^\mu}{q_1^2 P_Z(q_2^2)} \right) + \kappa_{\gamma\gamma} \epsilon_{\gamma\gamma}^{\text{SM-1L}} \frac{e^2 Q_e Q_\mu}{q_1^2 q_2^2} \right) \frac{q_1 \cdot q_2 g^{\alpha\beta} - q_2^\alpha q_1^\beta}{m_Z^2} + \left( \epsilon_{ZZ}^{\text{CP}} \frac{g_Z^e g_Z^\mu}{P_Z(q_1^2) P_Z(q_2^2)} + \epsilon_{Z\gamma}^{\text{CP}} \left( \frac{e Q_\mu g_Z^e}{q_2^2 P_Z(q_1^2)} + \frac{e Q_e g_Z^\mu}{q_1^2 P_Z(q_2^2)} \right) + \epsilon_{\gamma\gamma}^{\text{CP}} \frac{e^2 Q_e Q_\mu}{q_1^2 q_2^2} \right) \frac{\varepsilon^{\alpha\beta\rho\sigma} q_{2\rho} q_{1\sigma}}{m_Z^2} \right]$$

In the SM at tree level:  $\kappa_X \rightarrow 1$ ,  $\epsilon_X \rightarrow 0$

NLO EW - Prophecy4F code, hep-ph/0604011

Systematic inclusion of radiative QED corrections possible

Isidori et al, work in progress

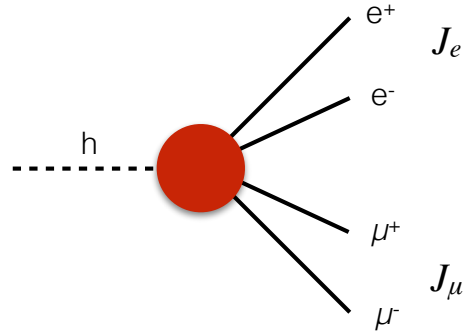
$$P_Z(q^2) = q^2 - m_Z^2 + im_Z \Gamma_Z$$

$$\epsilon_{\gamma\gamma}^{\text{SM-1L}} \simeq 3.8 \times 10^{-3},$$

$$\epsilon_{Z\gamma}^{\text{SM-1L}} \simeq 6.7 \times 10^{-3}$$

# $h \rightarrow e^+ e^- \mu^+ \mu^-$

The most generic structure of the amplitude:



$$A = i \frac{2m_Z^2}{v_F} \sum_{e=e_L, e_R} \sum_{\mu=\mu_L, \mu_R} (\bar{e} \gamma_\alpha e) (\bar{\mu} \gamma_\beta \mu) \times \left[ F_1^{e\mu}(q_1^2, q_2^2) g^{\alpha\beta} + F_3^{e\mu}(q_1^2, q_2^2) \frac{q_1 \cdot q_2 g^{\alpha\beta} - q_2^\alpha q_1^\beta}{m_Z^2} + F_4^{e\mu}(q_1^2, q_2^2) \frac{\epsilon^{\alpha\beta\rho\sigma} q_{2\rho} q_{1\sigma}}{m_Z^2} \right]$$

Momentum expansion of the form factors around the physical poles

$$A = i \frac{2m_Z^2}{v_F} \sum_{e=e_L, e_R} \sum_{\mu=\mu_L, \mu_R} (\bar{e} \gamma_\alpha e) (\bar{\mu} \gamma_\beta \mu) \times \left[ \begin{aligned} & \left( \kappa_{ZZ} \frac{g_Z^e g_Z^\mu}{P_Z(q_1^2) P_Z(q_2^2)} + \frac{\epsilon_{Ze}}{m_Z^2} \frac{g_Z^\mu}{P_Z(q_2^2)} + \frac{\epsilon_{Z\mu}}{m_Z^2} \frac{g_Z^e}{P_Z(q_1^2)} \right) g^{\alpha\beta} + \\ & + \left( \epsilon_{ZZ} \frac{g_Z^e g_Z^\mu}{P_Z(q_1^2) P_Z(q_2^2)} + \kappa_{Z\gamma} \epsilon_{Z\gamma}^{\text{SM-1L}} \left( \frac{e Q_\mu g_Z^e}{q_2^2 P_Z(q_1^2)} + \frac{e Q_e g_Z^\mu}{q_1^2 P_Z(q_2^2)} \right) + \kappa_{\gamma\gamma} \epsilon_{\gamma\gamma}^{\text{SM-1L}} \frac{e^2 Q_e Q_\mu}{q_1^2 q_2^2} \right) \frac{q_1 \cdot q_2 g^{\alpha\beta} - q_2^\alpha q_1^\beta}{m_Z^2} + \\ & + \left( \epsilon_{ZZ}^{\text{CP}} \frac{g_Z^e g_Z^\mu}{P_Z(q_1^2) P_Z(q_2^2)} + \epsilon_{Z\gamma}^{\text{CP}} \left( \frac{e Q_\mu g_Z^e}{q_2^2 P_Z(q_1^2)} + \frac{e Q_e g_Z^\mu}{q_1^2 P_Z(q_2^2)} \right) + \epsilon_{\gamma\gamma}^{\text{CP}} \frac{e^2 Q_e Q_\mu}{q_1^2 q_2^2} \right) \frac{\epsilon^{\alpha\beta\rho\sigma} q_{2\rho} q_{1\sigma}}{m_Z^2} \end{aligned} \right]$$

*Annotations:*  $hZ_\mu Z^\mu, hZ^\mu \partial^\nu V_{\mu\nu}$  (pointing to  $\kappa_{ZZ}$  term);  $hZ_\mu \bar{f} \gamma^\mu f, hZ^\mu \partial^\nu V_{\mu\nu}$  (pointing to  $\epsilon_{Ze}, \epsilon_{Z\mu}$  terms);  $hV_{\mu\nu} V^{\mu\nu}$  (pointing to  $\epsilon_{ZZ}$  term);  $h\epsilon^{\mu\nu\rho\sigma} V_{\mu\nu} V_{\rho\sigma}$  (pointing to  $\epsilon_{ZZ}^{\text{CP}}$  term);  $V = Z, \gamma$  (referring to the SM-1L and CP terms).

# Parameter counting

Consider decays:  $h \rightarrow WW^*$ ,  $h \rightarrow ZZ^*$ ,  $h \rightarrow Z\gamma$ ,  $h \rightarrow \gamma\gamma$

## Neutral currents

$h \rightarrow e^+e^-\mu^+\mu^-$

$h \rightarrow \mu^+\mu^-\mu^+\mu^-$

$h \rightarrow e^+e^-e^+e^-$

$h \rightarrow \gamma e^+e^-$

$h \rightarrow \gamma\mu^+\mu^-$

$h \rightarrow \gamma\gamma$

$\kappa_{ZZ}, \kappa_{Z\gamma}, \kappa_{\gamma\gamma}, \epsilon_{ZZ}$  ,

$\epsilon_{Z\gamma}^{CP}, \epsilon_{\gamma\gamma}^{CP}, \epsilon_{ZZ}^{CP}$  ,

$\epsilon_{Ze_L}, \epsilon_{Ze_R}, \epsilon_{Z\mu_L}, \epsilon_{Z\mu_R}$

11

## Charged currents

$h \rightarrow e^+\mu^- \nu\nu$

$h \rightarrow e^-\mu^+ \nu\nu$

$\kappa_{WW}, \epsilon_{WW}, \epsilon_{WW}^{CP}$  ,

$\epsilon_{W_e}, \epsilon_{W_\mu}$  (complex)

7

## N. & C. interference

$h \rightarrow e^+e^- \nu\nu$

$h \rightarrow \mu^-\mu^+ \nu\nu$

$\epsilon_{Z\nu_e}, \epsilon_{Z\nu_\mu}$

2

- Compared to only 4 parameter in the present “kappa” formalism
- Easy to implement in MC event generator, AG, Marzocca, private code

# Symmetry limits

Consider decays:  $h \rightarrow WW^*$ ,  $h \rightarrow ZZ^*$ ,  $h \rightarrow Z\gamma$ ,  $h \rightarrow \gamma\gamma$

## Neutral currents

$$h \rightarrow e^+e^-\mu^+\mu^-$$

$$h \rightarrow \mu^+\mu^-\mu^+\mu^-$$

$$h \rightarrow e^+e^-e^+e^-$$

$$h \rightarrow \gamma e^+e^-$$

$$h \rightarrow \gamma\mu^+\mu^-$$

$$h \rightarrow \gamma\gamma$$

$$\kappa_{ZZ}, \kappa_{Z\gamma}, \kappa_{\gamma\gamma}, \epsilon_{ZZ},$$

$$\epsilon_{Z\gamma}^{CP}, \epsilon_{\gamma\gamma}^{CP}, \epsilon_{ZZ}^{CP},$$

$$\epsilon_{Ze_L}, \epsilon_{Ze_R}$$

9

FU

## Charged currents

$$h \rightarrow e^+\mu^- \nu\nu$$

$$h \rightarrow e^-\mu^+ \nu\nu$$

$$\kappa_{WW}, \epsilon_{WW}, \epsilon_{WW}^{CP},$$

$$\epsilon_{We_L} \quad (\text{complex})$$

5

## N. & C. interference

$$h \rightarrow e^+e^- \nu\nu$$

$$h \rightarrow \mu^-\mu^+ \nu\nu$$

$$\epsilon_{Z\nu_e}$$

1

## Flavour universality

$$\epsilon_{Ze_L} = \epsilon_{Z\mu_L},$$

$$\epsilon_{Ze_R} = \epsilon_{Z\mu_R},$$

$$\epsilon_{Z\nu_e} = \epsilon_{Z\nu_\mu},$$

$$\epsilon_{We_L} = \epsilon_{W\mu_L}.$$

# Symmetry limits

Consider decays:  $h \rightarrow WW^*$ ,  $h \rightarrow ZZ^*$ ,  $h \rightarrow Z\gamma$ ,  $h \rightarrow \gamma\gamma$

## Neutral currents

$$h \rightarrow e^+e^-\mu^+\mu^-$$

$$h \rightarrow \mu^+\mu^-\mu^+\mu^-$$

$$h \rightarrow e^+e^-e^+e^-$$

$$h \rightarrow \gamma e^+e^-$$

$$h \rightarrow \gamma\mu^+\mu^-$$

$$h \rightarrow \gamma\gamma$$

$\kappa_{ZZ}, \kappa_{Z\gamma}, \kappa_{\gamma\gamma}, \epsilon_{ZZ}$ ,

$\epsilon_{Ze_L}, \epsilon_{Ze_R}$

6

FU + CP

## Charged currents

$$h \rightarrow e^+\mu^- \nu\nu$$

$$h \rightarrow e^-\mu^+ \nu\nu$$

$\kappa_{WW}, \epsilon_{WW}$ ,

$\text{Re}(\epsilon_{We_L})$  (complex)

3

## N. & C. interference

$$h \rightarrow e^+e^- \nu\nu$$

$$h \rightarrow \mu^-\mu^+ \nu\nu$$

$\epsilon_{Z\nu_e}$

1

## CP conservation

$$\epsilon_{ZZ}^{CP} = \epsilon_{Z\gamma}^{CP} = \epsilon_{\gamma\gamma}^{CP} = \epsilon_{WW}^{CP} = \text{Im}\epsilon_{We_L} = \text{Im}\epsilon_{W\mu_L} = 0$$

# Symmetry limits

Consider decays:  $h \rightarrow WW^*$ ,  $h \rightarrow ZZ^*$ ,  $h \rightarrow Z\gamma$ ,  $h \rightarrow \gamma\gamma$

## Neutral currents

$$h \rightarrow e^+e^-\mu^+\mu^-$$

$$h \rightarrow \mu^+\mu^-\mu^+\mu^-$$

$$h \rightarrow e^+e^-e^+e^-$$

$$h \rightarrow \gamma e^+e^-$$

$$h \rightarrow \gamma\mu^+\mu^-$$

$$h \rightarrow \gamma\gamma$$

$$\kappa_{ZZ}, \kappa_{Z\gamma}, \kappa_{\gamma\gamma}, \epsilon_{ZZ},$$

$$\epsilon_{Ze_L}, \epsilon_{Ze_R}$$

6

## Charged currents

$$h \rightarrow e^+\mu^- \nu\nu$$

$$h \rightarrow e^-\mu^+ \nu\nu$$

$$\text{Re}(\epsilon_{We_L}) \quad (\text{complex})$$

1

## N. & C. interference

$$h \rightarrow e^+e^- \nu\nu$$

$$h \rightarrow \mu^-\mu^+ \nu\nu$$

0

FU + CP + cust. symm.

## Custodial symmetry

$$\epsilon_{WW} = c_w^2 \epsilon_{ZZ} + 2c_w s_w \epsilon_{Z\gamma} + s_w^2 \epsilon_{\gamma\gamma},$$

$$\epsilon_{WW}^{CP} = c_w^2 \epsilon_{ZZ}^{CP} + 2c_w s_w \epsilon_{Z\gamma}^{CP} + s_w^2 \epsilon_{\gamma\gamma}^{CP},$$

$$\kappa_{WW} - \kappa_{ZZ} = -\frac{2}{g} \left( \sqrt{2} \epsilon_{We_L^i} + 2c_w \epsilon_{Ze_L^i} \right),$$

$$\epsilon_{We_L^i} = \frac{c_w}{\sqrt{2}} (\epsilon_{Z\nu_L^i} - \epsilon_{Ze_L^i}),$$

\*The BSM sector is invariant under the custodial symmetry group

$G = \text{SU}(2)_L \times \text{SU}(2)_R \times \text{U}(1)_X$ , spontaneously broken to

$H = \text{SU}(2)_{L+R} \times \text{U}(1)_X$ .

# Linear vs Non-linear EFT

- **Non-linear EFT:**

An effective decoupling of the Higgs boson from the Goldstone-boson components of the  $SU(2)_L \times U(1)_Y / U(1)_{em}$  symmetry breaking.

- EW symmetry is non-linearly realised, derivative expansion over the cutoff
- All Higgs **POs** independent

- **Linear EFT:**

Higgs boson is part of an  $SU(2)_L$  doublet field **H**. Higher-dimensional operators are constructed in terms of the **H** field. The physical Higgs boson appears in operators contributing also to non-Higgs processes.

- **(1)** Some Higgs **POs** constrained from LEP data
- **(2)** Relations among Higgs **POs** due to the accidental custodial symmetry present in some of the  $D = 6$  operators

$$\begin{aligned}\epsilon_{WW} &= c_w^2 \epsilon_{ZZ} + 2c_w s_w \epsilon_{Z\gamma} + s_w^2 \epsilon_{\gamma\gamma}, \\ \epsilon_{WW}^{CP} &= c_w^2 \epsilon_{ZZ}^{CP} + 2c_w s_w \epsilon_{Z\gamma}^{CP} + s_w^2 \epsilon_{\gamma\gamma}^{CP}, \\ \epsilon_{We_L^i} &= \frac{c_w}{\sqrt{2}} (\epsilon_{Z\nu_L^i} - \epsilon_{Ze_L^i})\end{aligned}$$

**Violation of (1) and (2) would point towards non-linear realisation of EW symmetry!**

**Test from Higgs data!!!**



# Higgs POs from *LEP* data

Gonzalez-Alonso, AG,  
Isidori, Marzocca,  
work in progress

Assuming linear EFT with dimension 6 operators:

$$\epsilon_{Zf} = \frac{2m_Z}{v} (\delta g^{Zf} - (c_\theta^2 T_f^3 + s_\theta^2 Y_f) \mathbf{1}_3 \delta g_{1,z} + t_\theta^2 Y_f \mathbf{1}_3 \delta \kappa_\gamma) ,$$

$$\epsilon_{Wf} = \frac{\sqrt{2}m_W}{v} (\delta g^{Wf} - c_\theta^2 \mathbf{1}_3 \delta g_{1,z})$$

LEP I

LEP II

$$\delta \epsilon_{ZZ} = \delta \epsilon_{\gamma\gamma} + \frac{c_{2\theta}}{s_\theta c_\theta} \delta \epsilon_{Z\gamma} - \frac{1}{c_\theta^2} \delta \kappa_\gamma ,$$

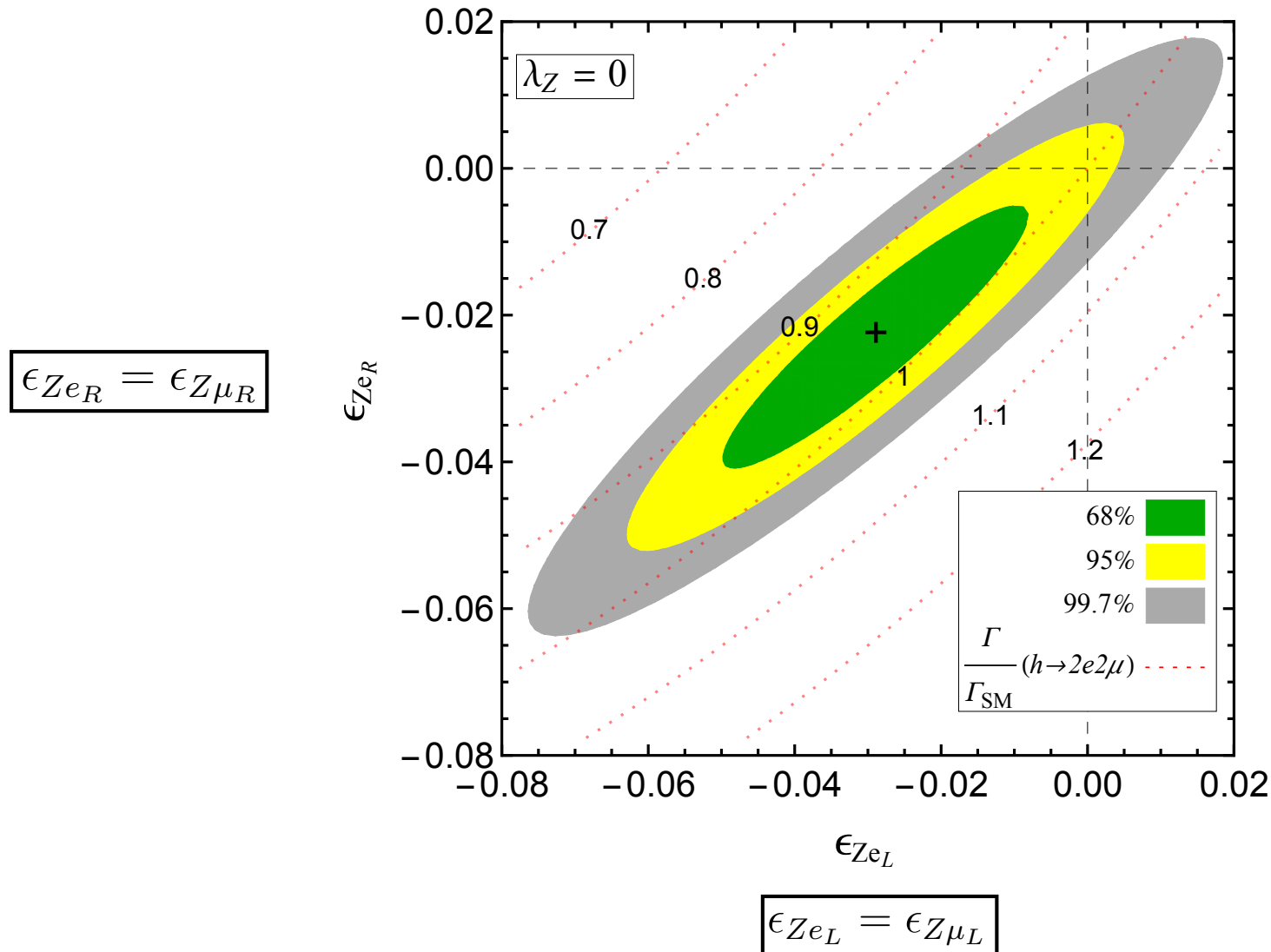
LEP II

# Higgs POs from *LEP* data

Gonzalez-Alonso, AG,  
Isidori, Marzocca,  
work in progress

Assuming linear EFT with dimension 6 operators:

Using the EW fits from Falkowski, Riva and Efrati, Falkowski, Soreq we find:

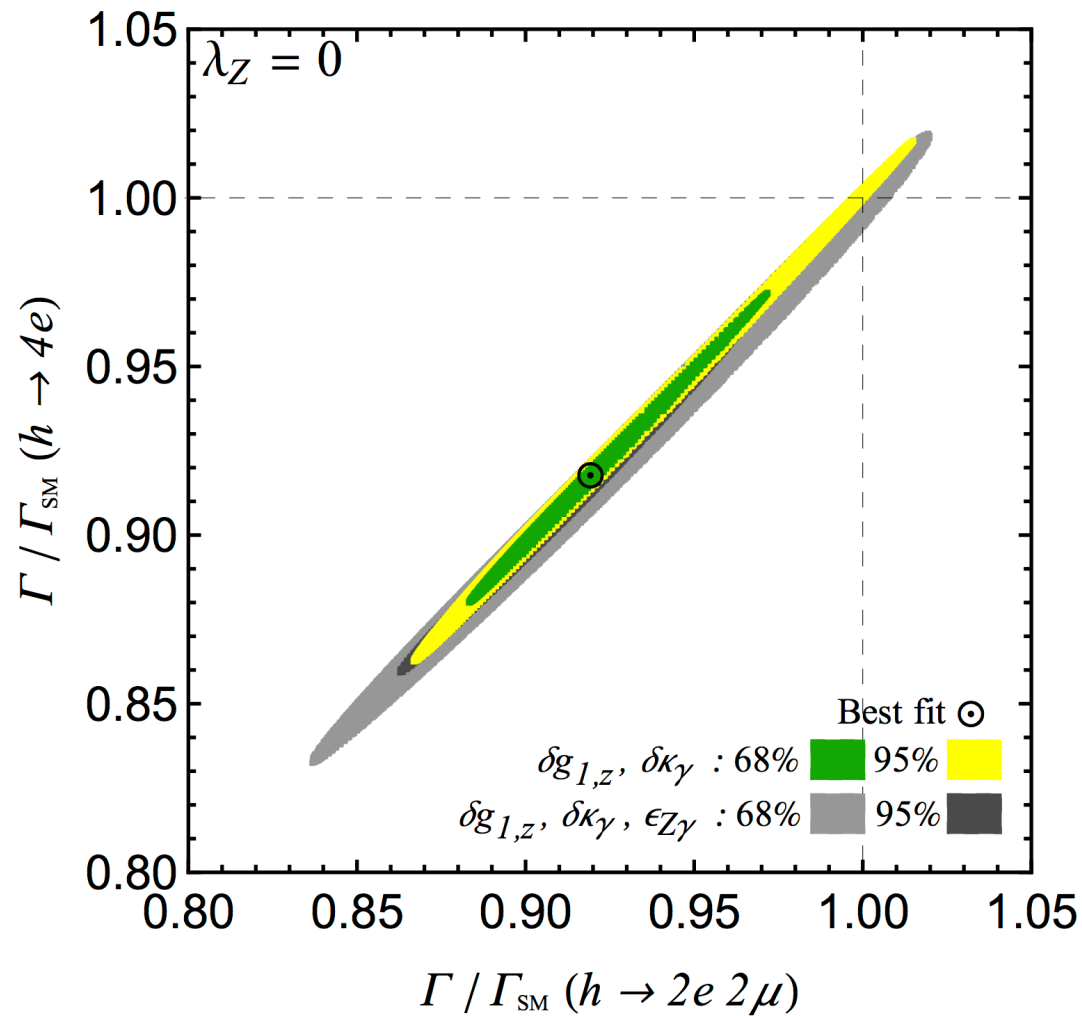


Flavour  
universality

# Higgs POs from *LEP* data

Gonzalez-Alonso, AG,  
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Assuming linear EFT with dimension 6 operators:

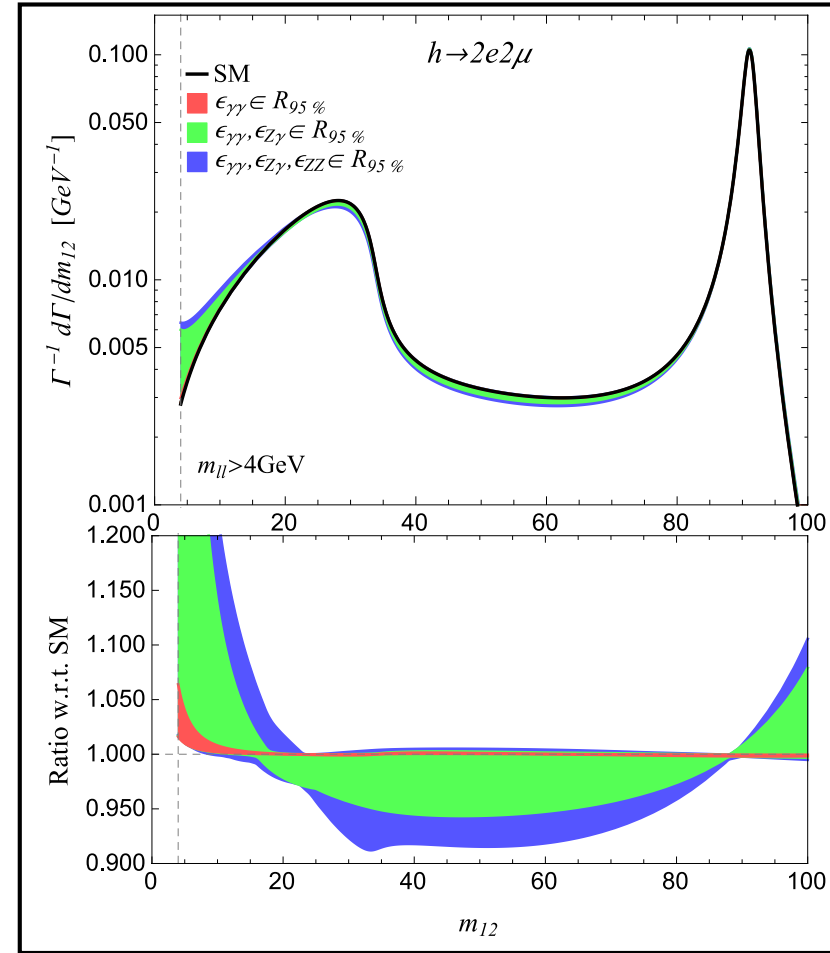
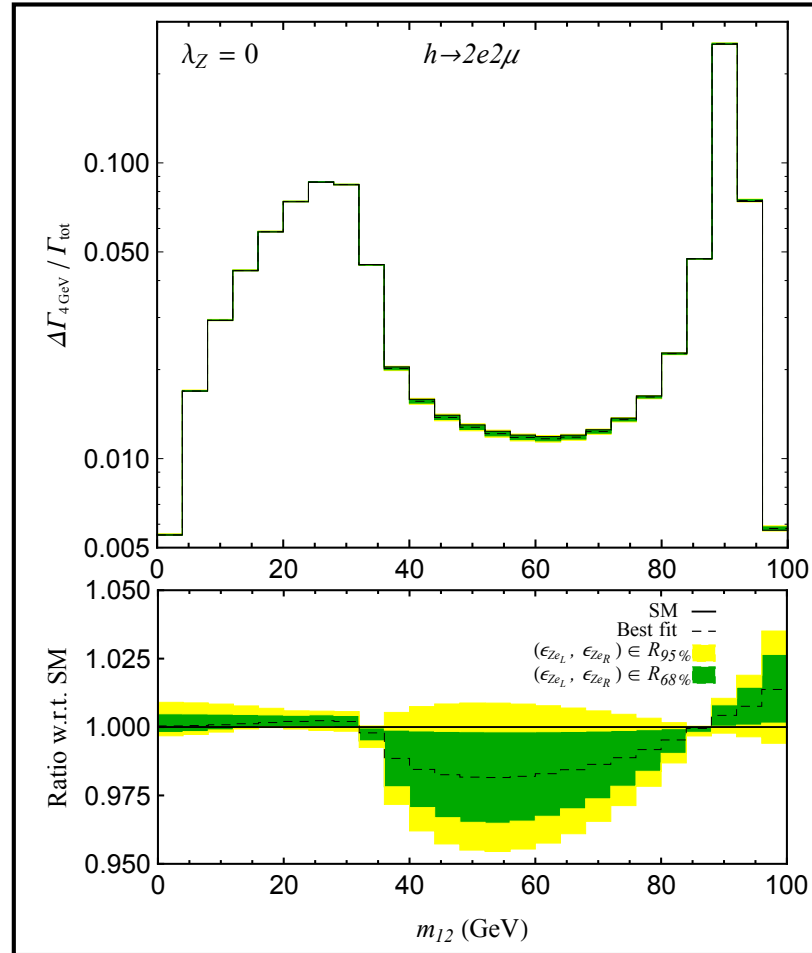


Flavour  
universality

# Higgs POs from *LEP* data

Gonzalez-Alonso, AG,  
Isidori, Marzocca,  
work in progress

Assuming linear EFT with dimension 6 operators:



Small deviations in the shape expected.

# Conclusions

- We propose a set of **POs** to characterise NP in Higgs decays
  - A. General encoding of the experimental results
  - B. Easily computable in large set of theories
  - C. Not a substitute (or competition) to **EFT** approach, rather an intermediate step
  - D. Dressing with QED radiation possible
- FU, CP, custodial symmetry, linear or non-linear EFT not assumed!  
**We should keep our eyes open.**
- Linear EFT > firm predictions on  $h \rightarrow 4\ell$  shape & LFU  
(possibility to falsify them with Higgs data would be a "double discovery": NP +  $h(125)$  non-pure  $SU(2)_L$  doublet)
- Work to do: compute the projections for LHC run 2  
(in collaboration with experimental groups)