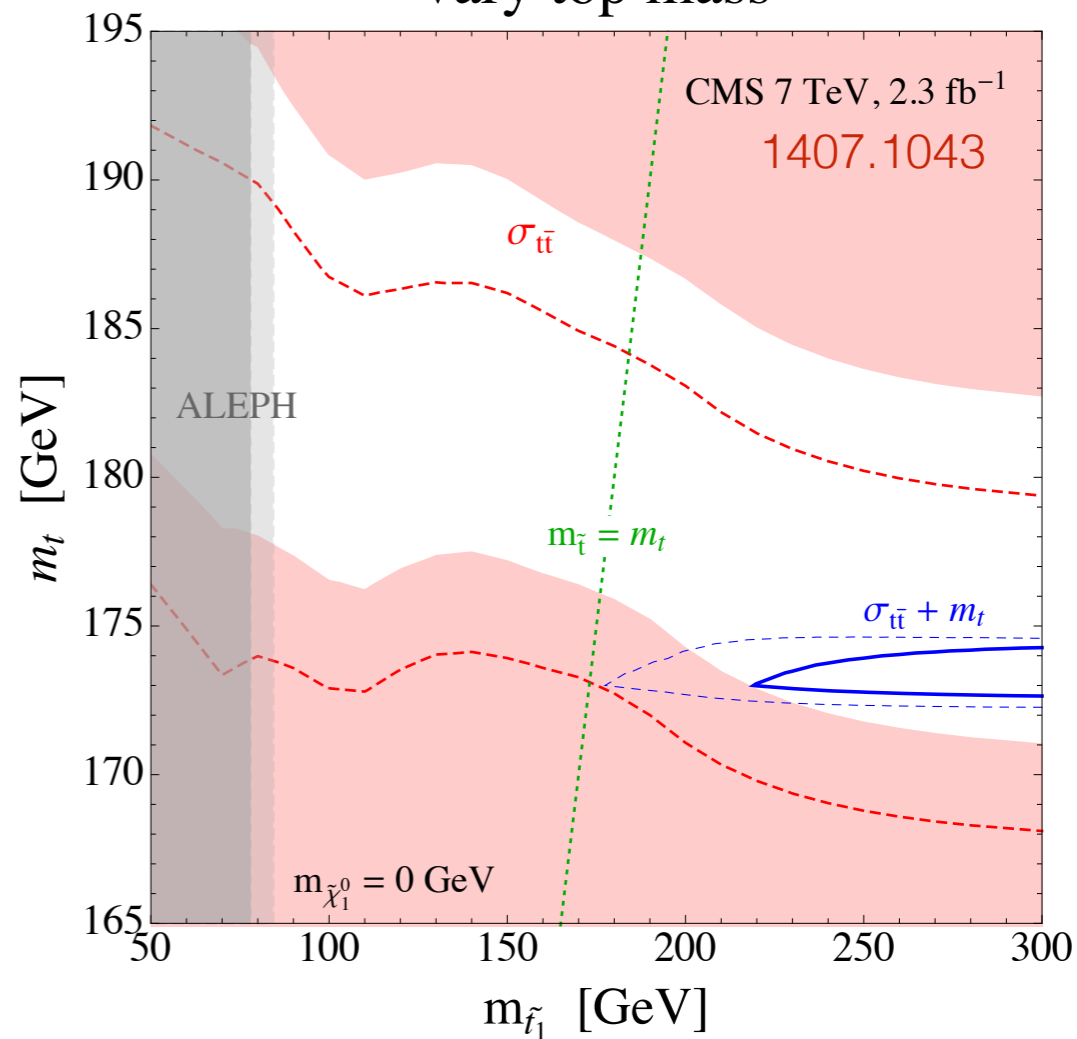
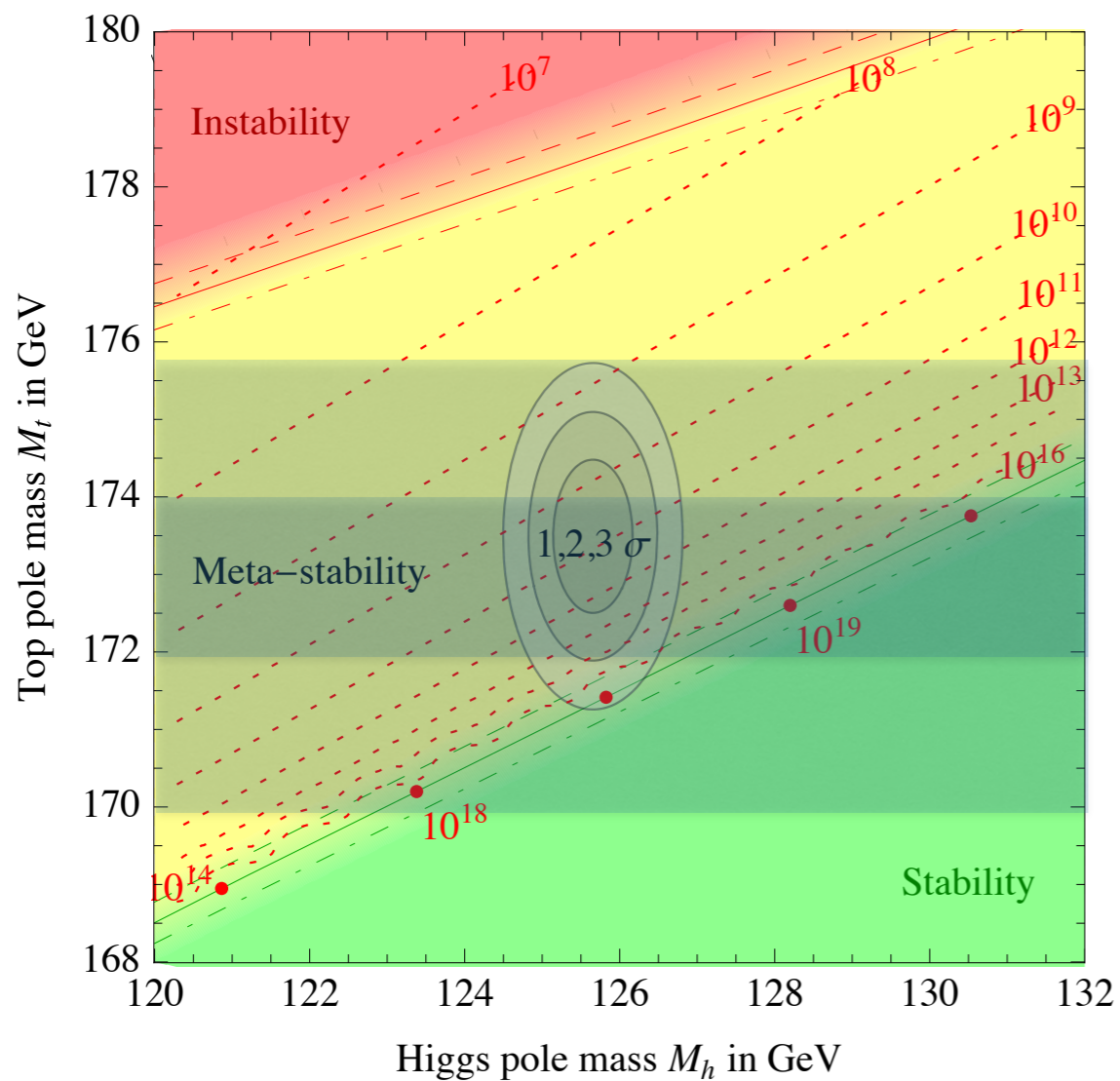
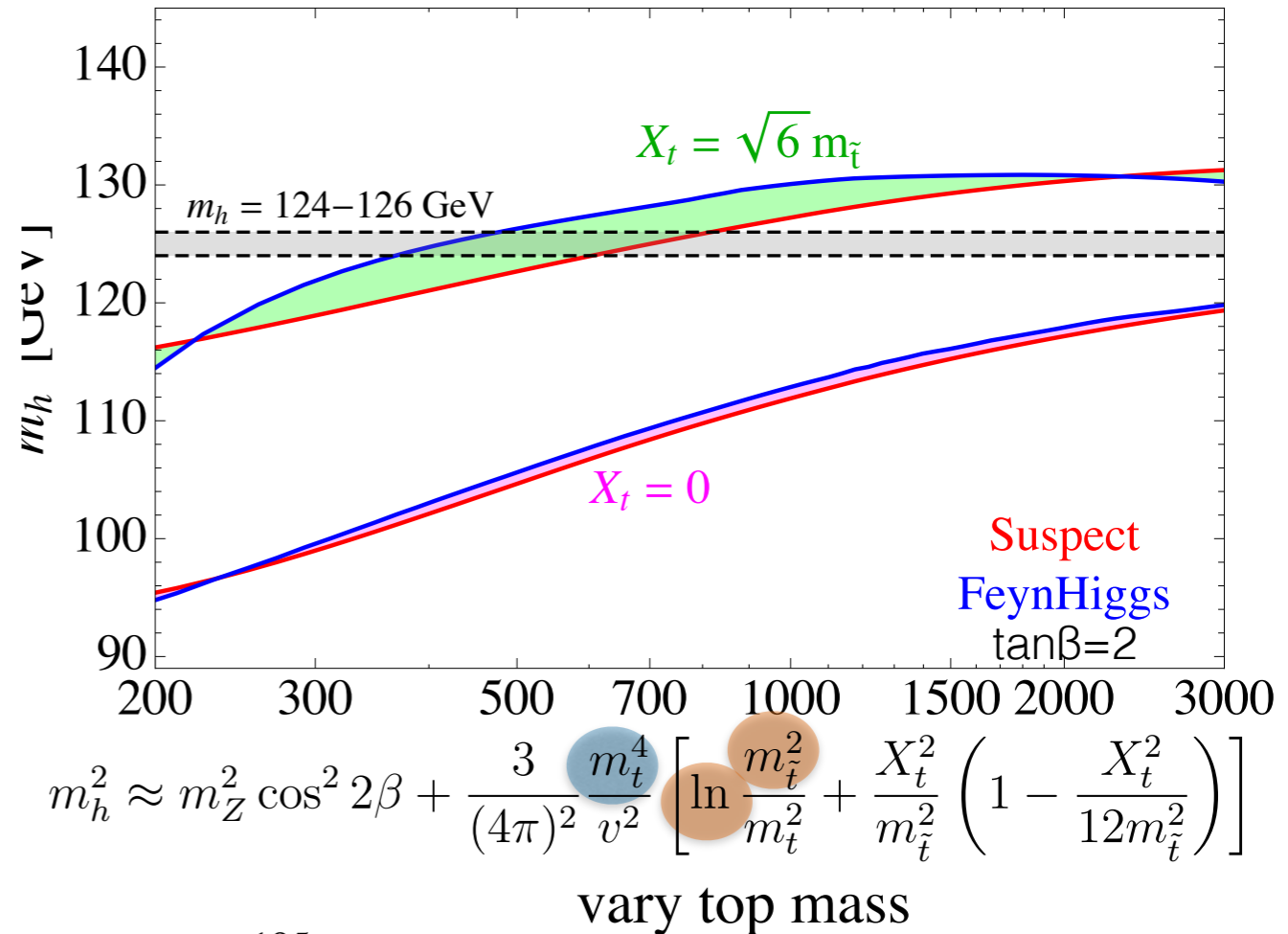
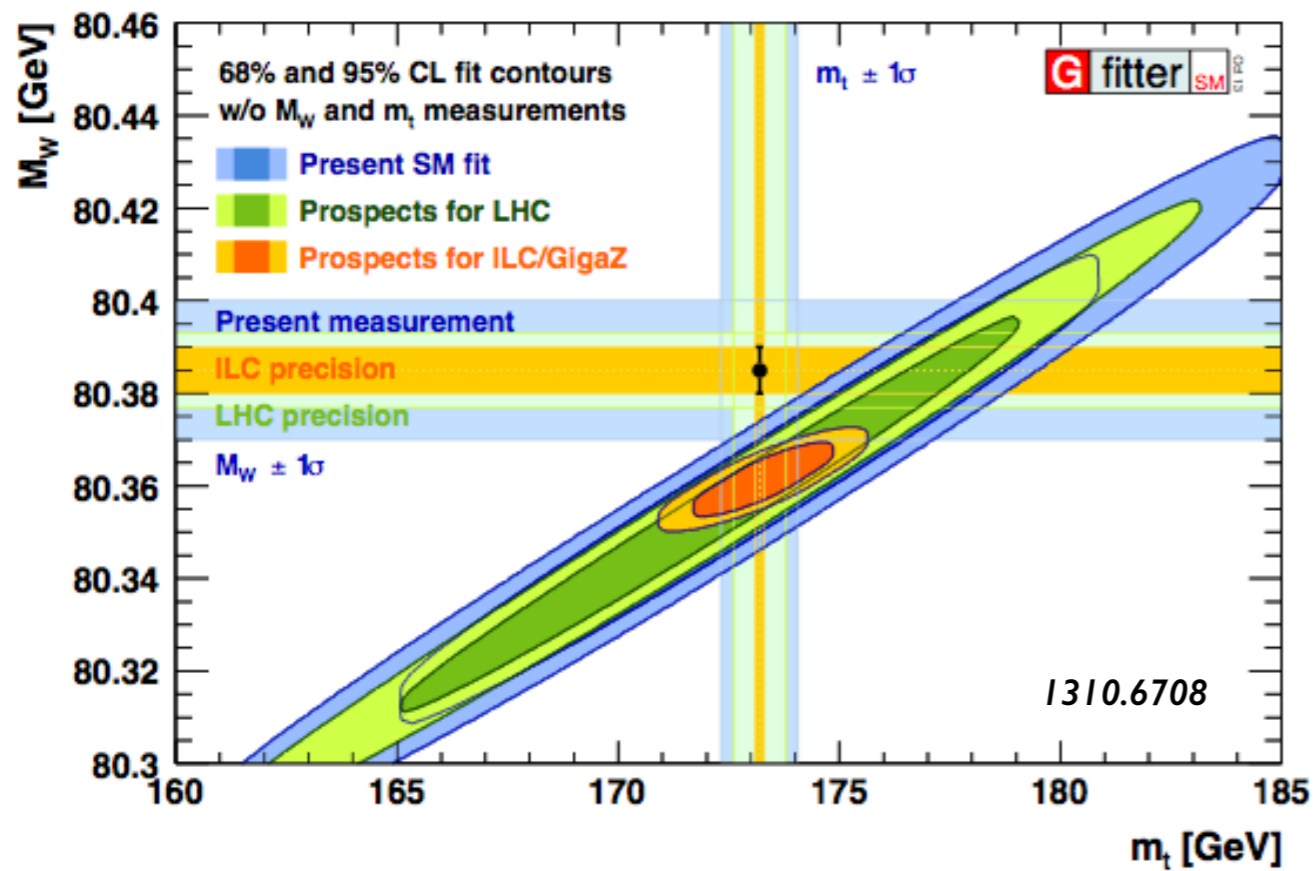


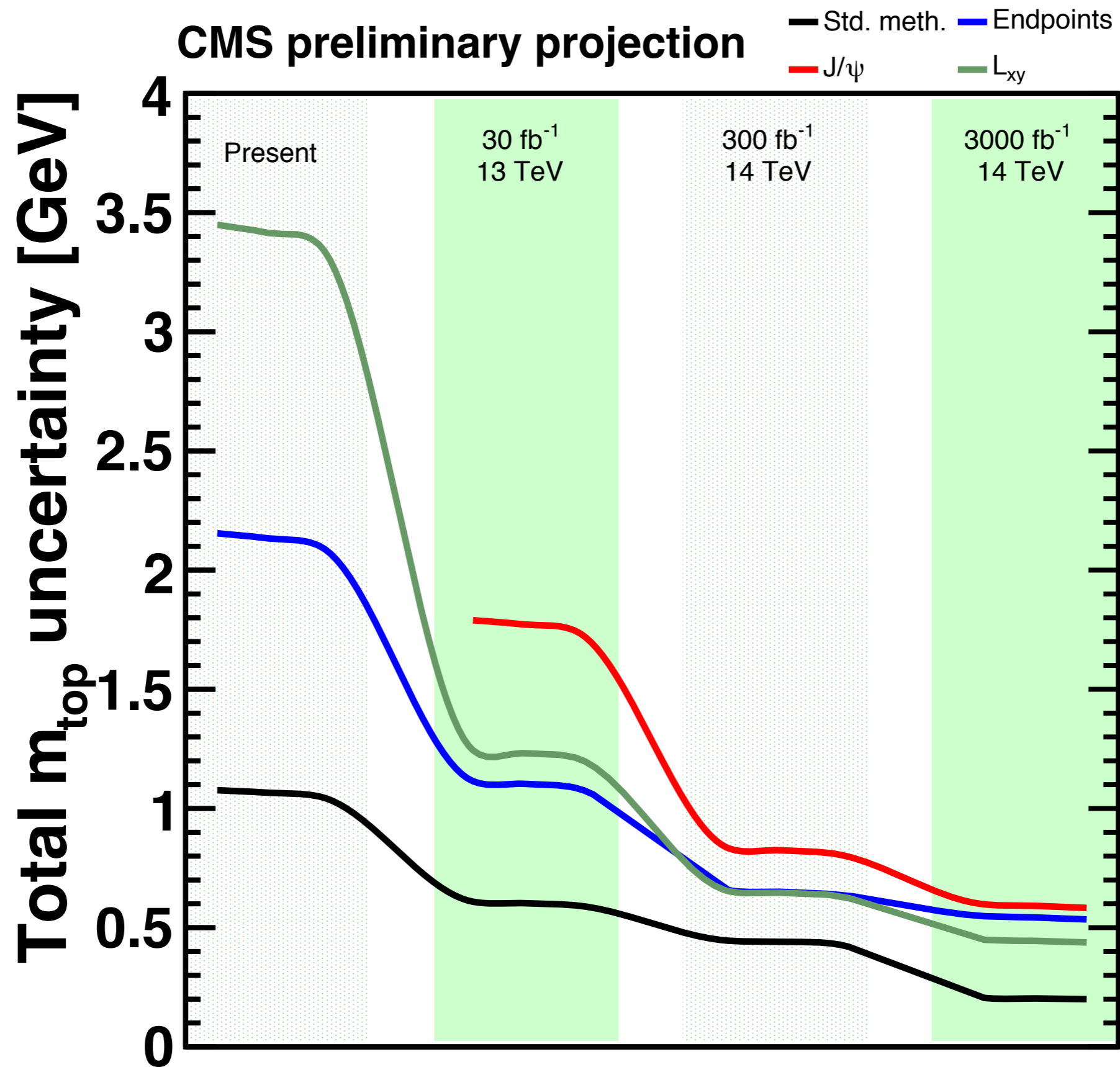
# Top quark mass from the bottom (at NLO)

Roberto Franceschini (CERN)  
April 9th 2015 (Portorož 2015)

Work in Progress with K. Agashe, D. Kim and M. Schulze



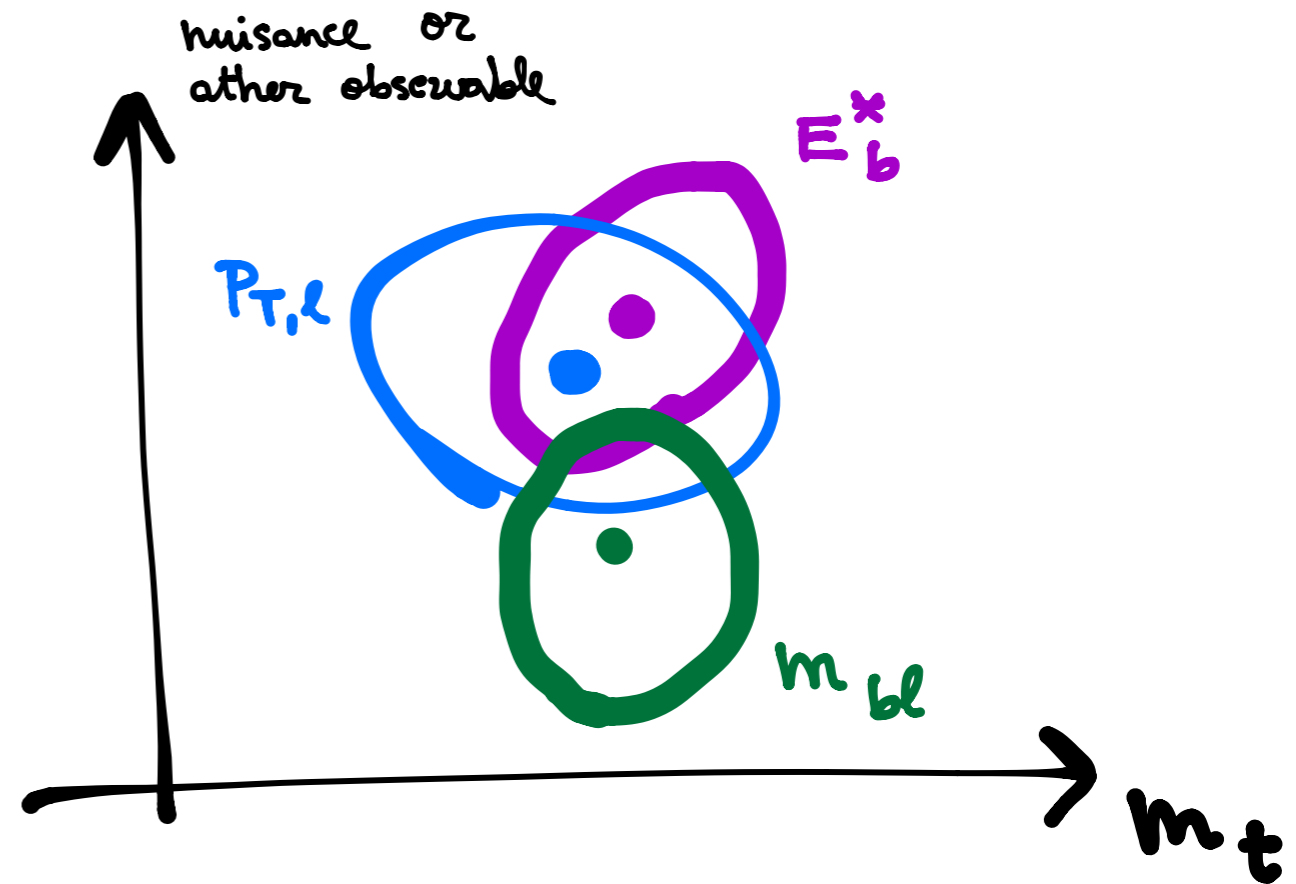
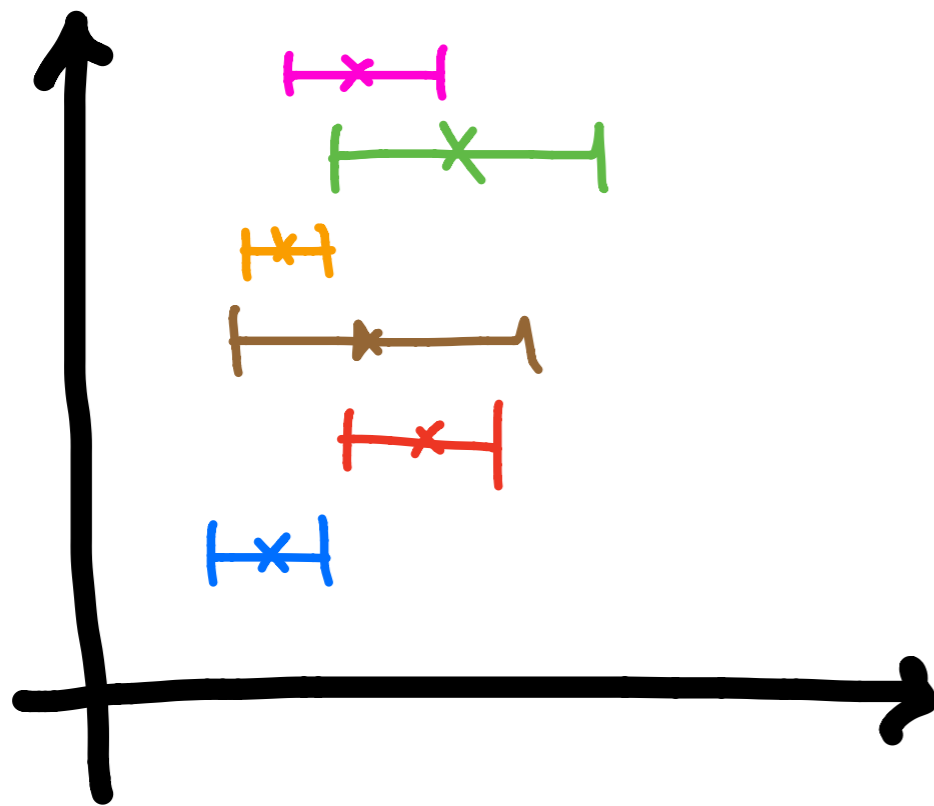
# Global strategy



CMS-PAS-FTR-13-017

1310.0799 - Juste,  
Mantry, Mitov, Penin,  
Skands, Varnes, Vos,  
Wimpenny -  
Determination of the  
top quark mass circa  
2013: methods,  
subtleties, perspective

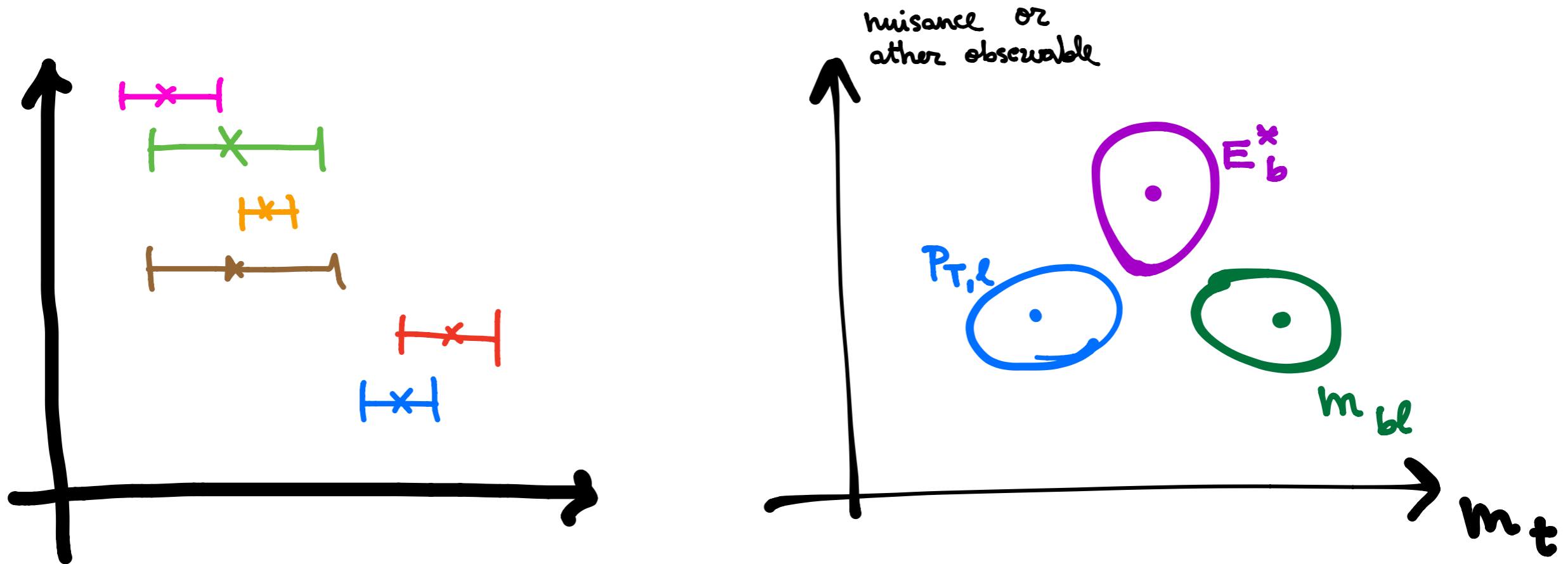
# Many measurements



The strength of the future LHC top mass measurement will build on the **diversity of methods**  
⇒ not very useful to talk about “*single best measurement*”

# Many measurements

due to different hypothesis, different mass measurement methods can result in significantly disagreeing measurements: **QCD or new physics effect?**



The strength of the future LHC top mass measurement will build on the **diversity of methods**  
⇒ not very useful to talk about “*single best measurement*”

# (Alternative) Methods

- Energy Peaks [1209.0772 + WIP](#)
- Generalized Medians [1405.2395](#)
- Leptonic Mellin moments [1407.2763](#)
- B-hadron life-time -  $L_{xy}$  [hep-ex/0501043](#)
- J/ $\psi$  exclusive [hep-ph/9912320](#)
- $d\sigma(ttj)$  [1303.6415](#)
- Inclusive  $\sigma(tt)$  [1307.1907](#)



# Energy Peaks

# Lorentz *variant* quantities

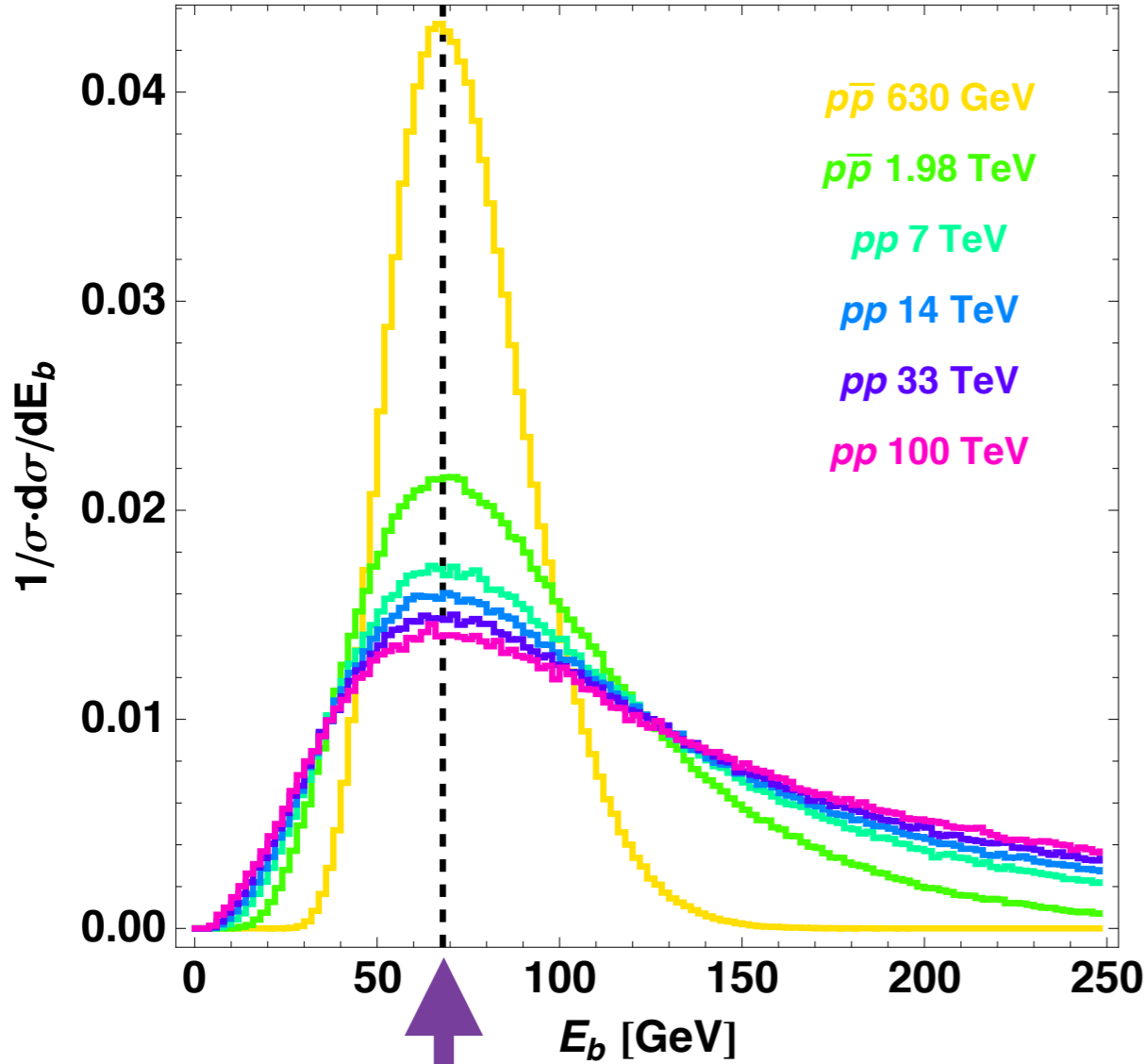
Given suitable conditions, Lorentz  
variant quantities can tell us a lot about  
the invariants



# How special is this invariance?

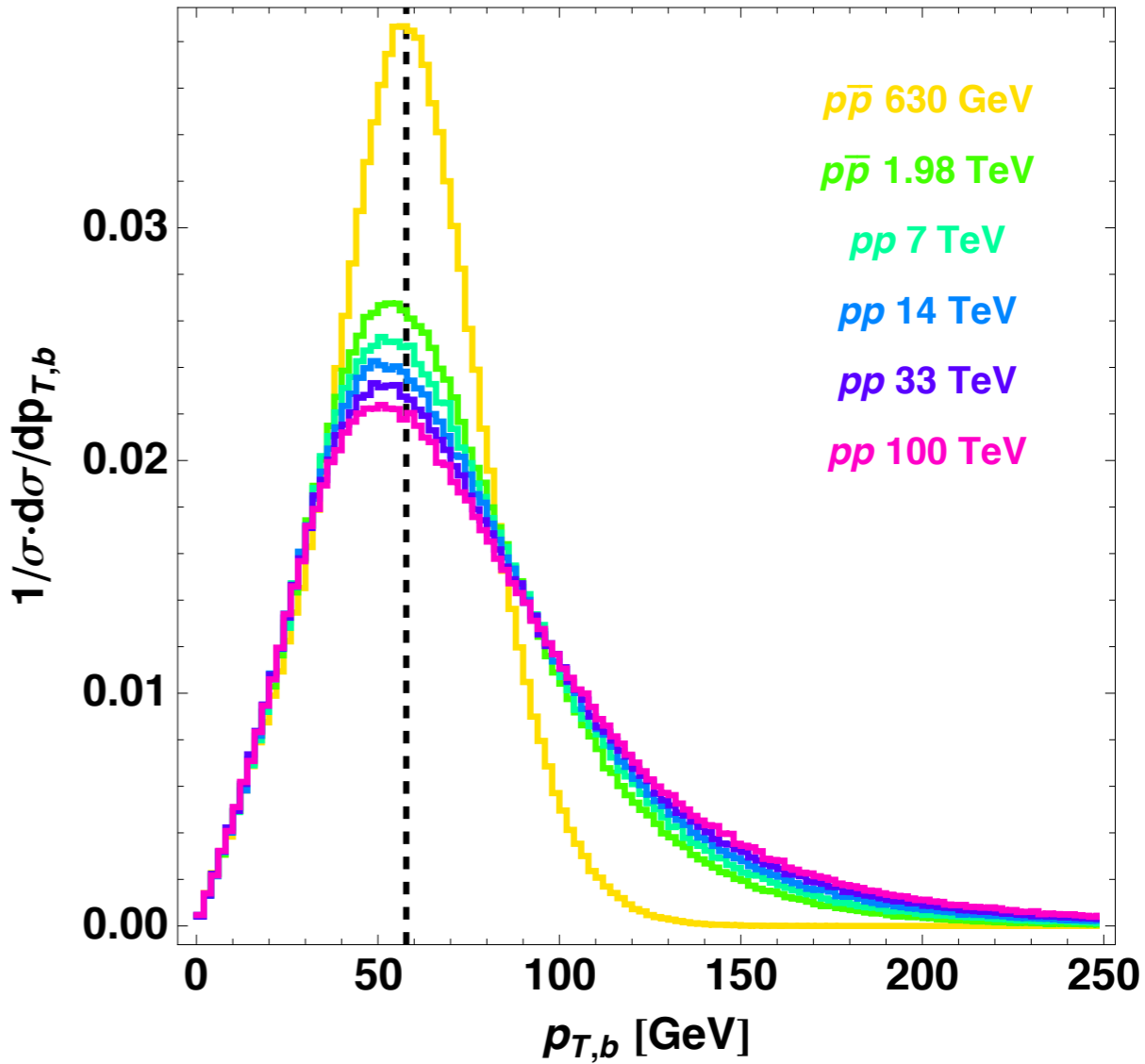
1209.0772 - Agashe, Franceschini and Kim

Shape changes, peak doesn't!



$$E_b^* = \frac{m_t^2 - m_W^2 + m_b^2}{2m_t}$$

Shape changes, peak does too



The sensitivity to the **boost distribution** is the key

# The Breit-Wigner peak substitute?



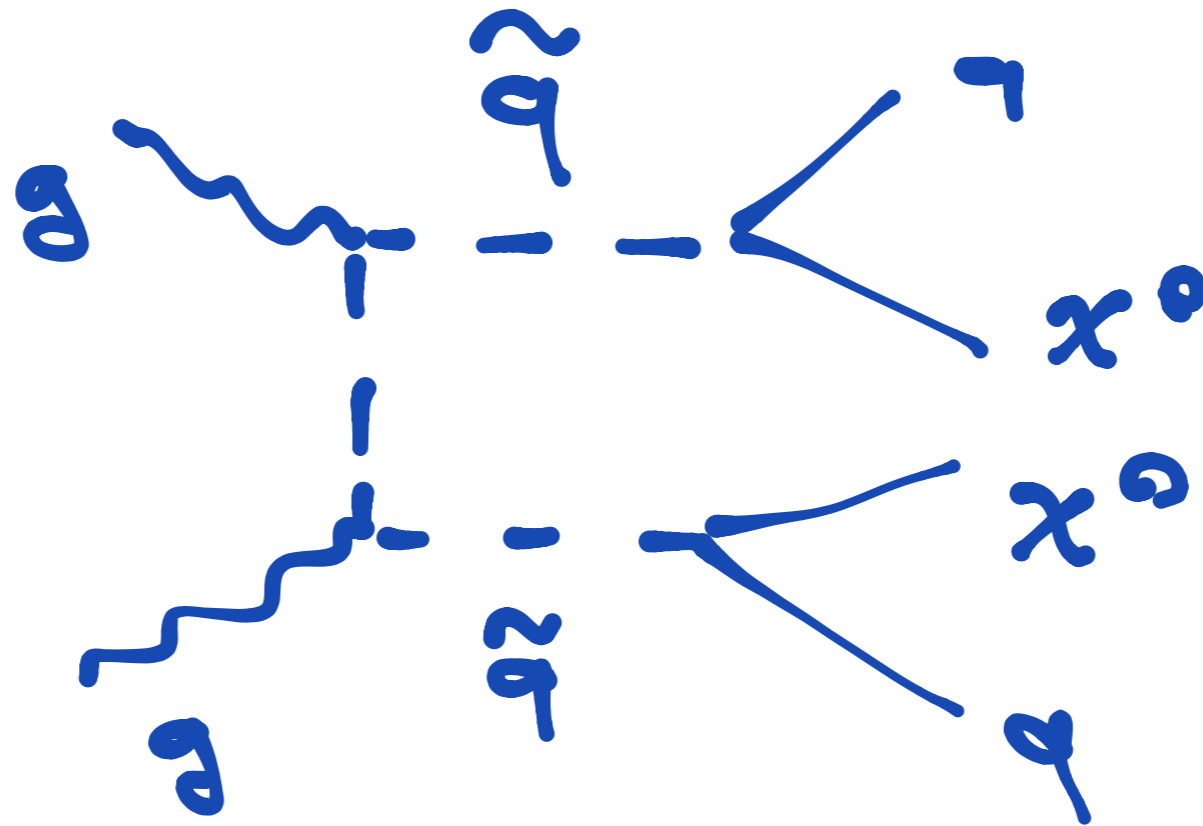
$$(P_{\mu^+} P_{\mu^-})^2 \rightarrow m_Z^2$$

# The Breit-Wigner peak substitute?



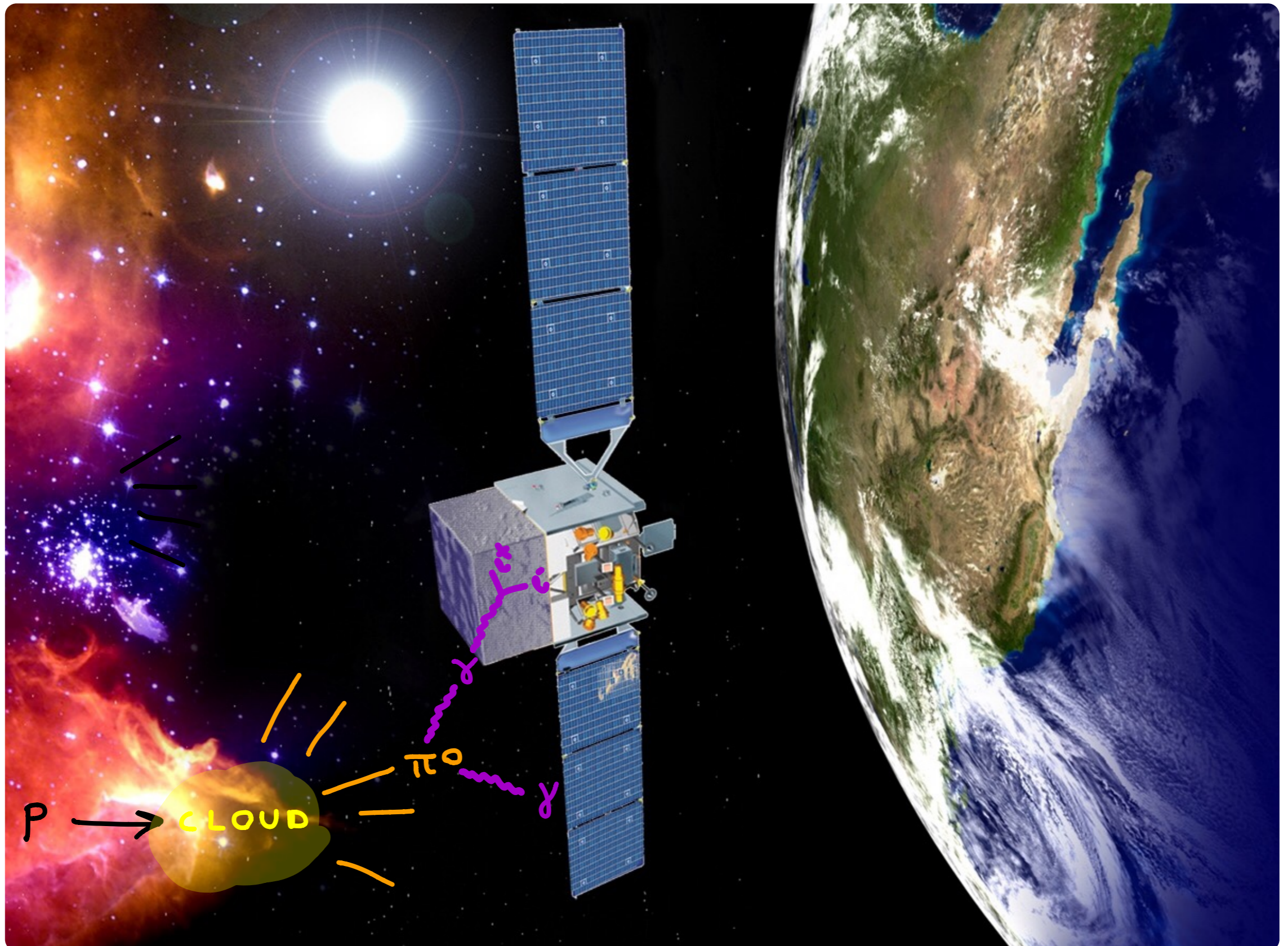
?

# The Breit-Wigner peak substitute?



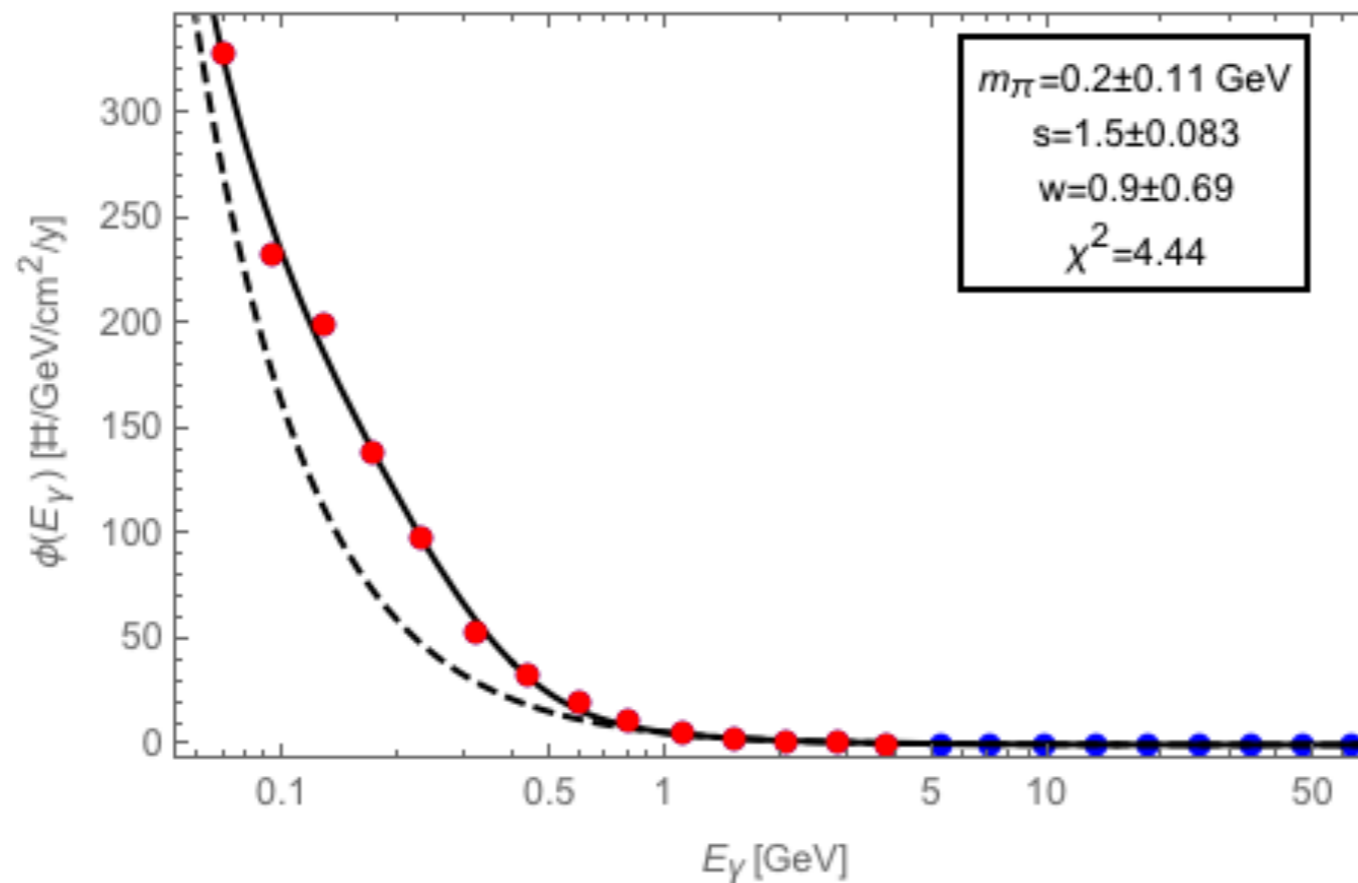
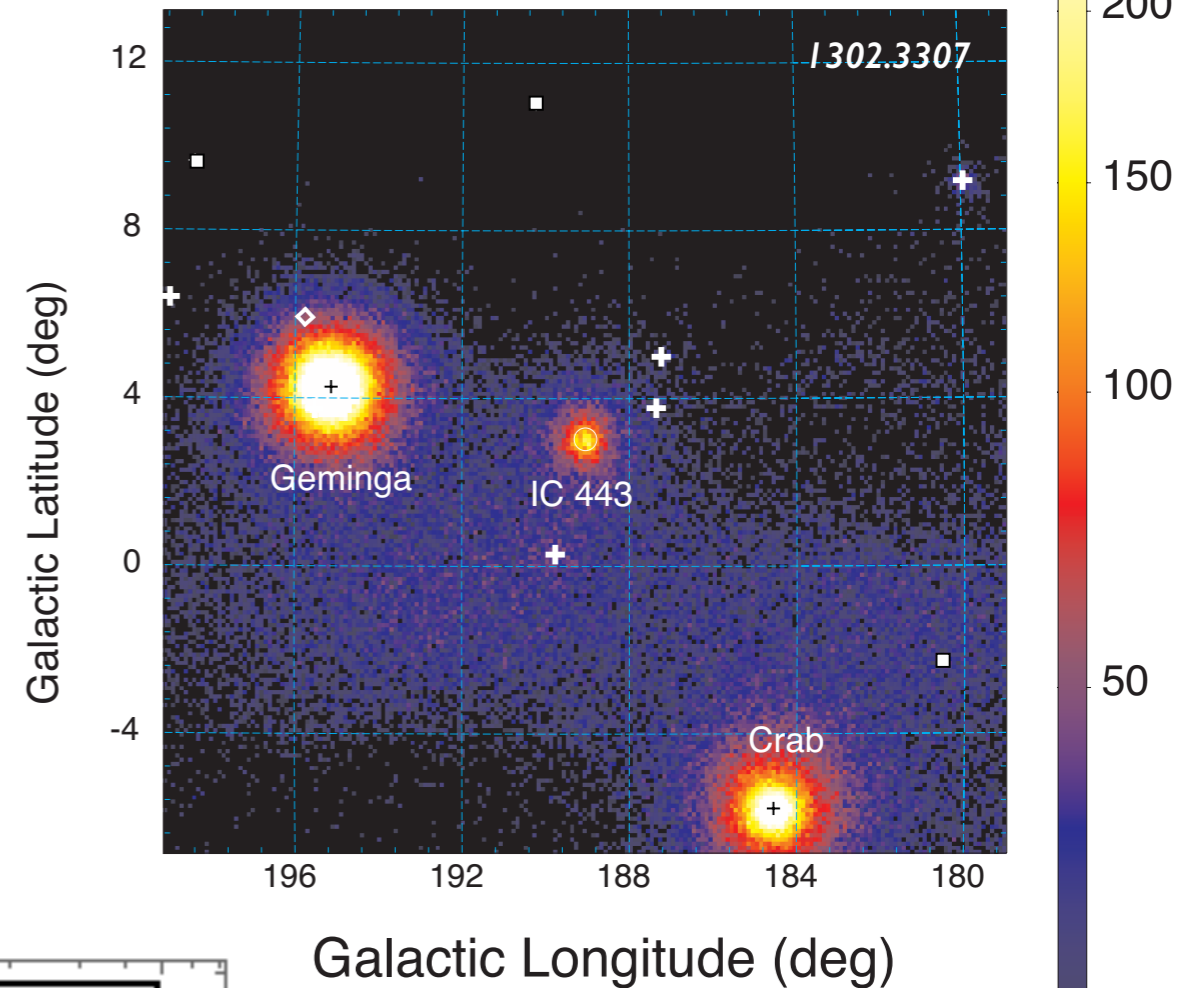
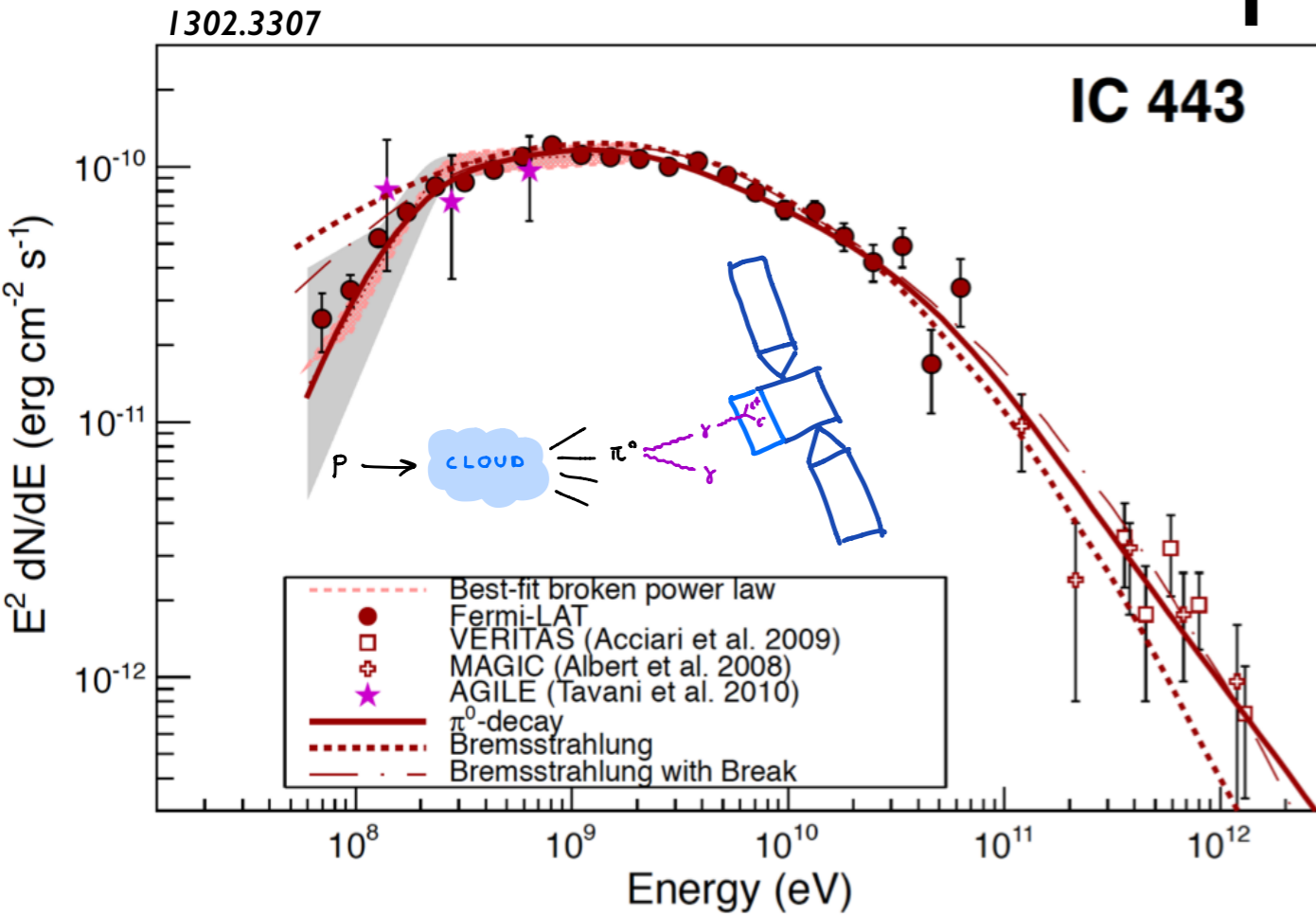
???

# Cosmic peaks (Stecker 1971)



# Cosmic peaks

(Stecker 1971)



- Data (FERMI IC433)
- - -  $E^{-s}$  (background)
- $E^{-s} + \text{Exp}(-w(\frac{2E}{m_\pi} + \frac{m_\pi}{2E}))$
- Fitted data

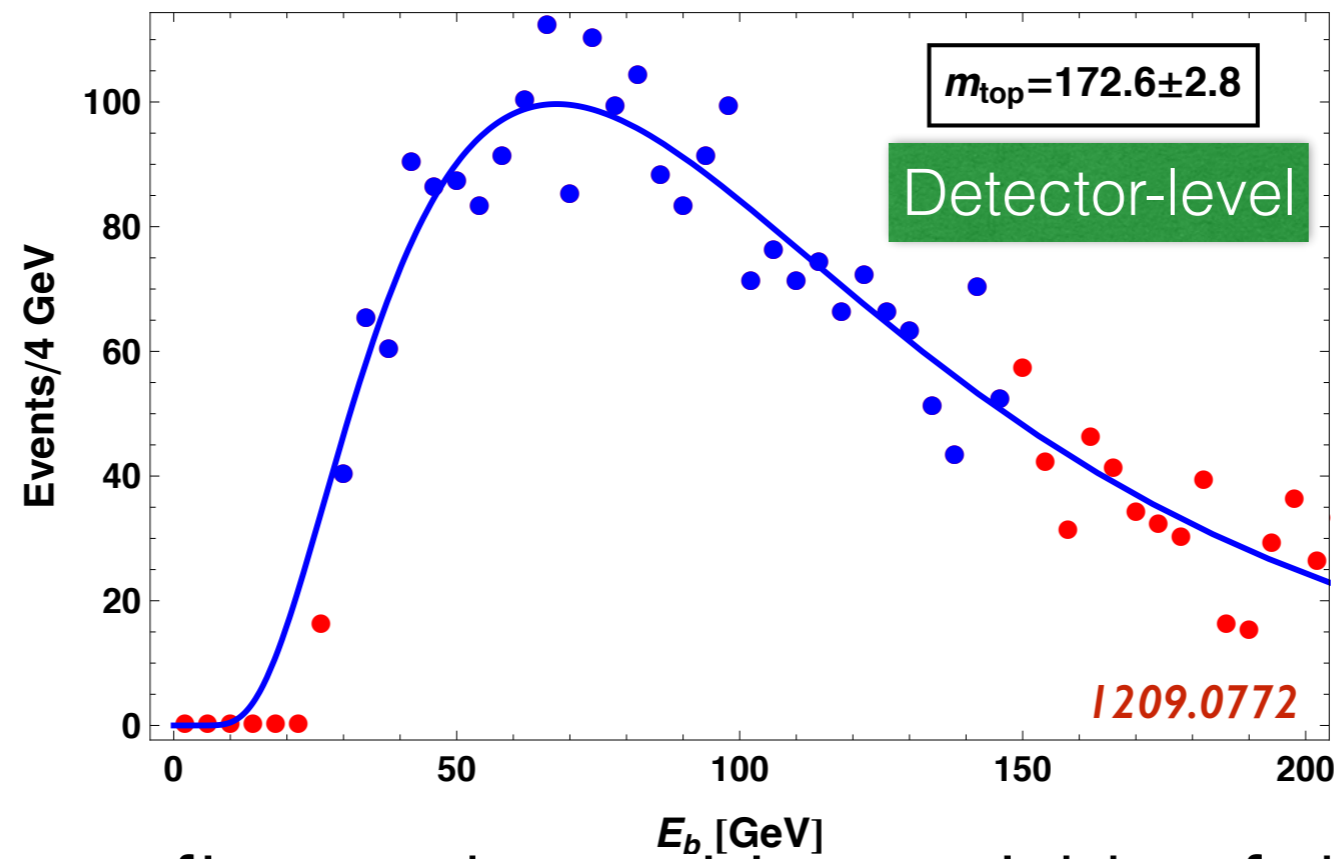
- properties similar to Lorentz invariants

Useful in practice?

$$E_b^x = \frac{m_t^2 - m_w^2 + m_b^2}{2m_t}$$

# b-jet energy (LO+PS)

100 pseudo-experiments from [MadGraph5+Pythia6.4+Delphes](#) (**ATLAS-2012-097**)



2-parameters fit: peak position, width of the distribution

*Proof of the concept: 5/fb LHC 7 TeV*

**$m_{\text{top}} = 173.1 \pm 2.5 \text{ GeV (stat)}$**

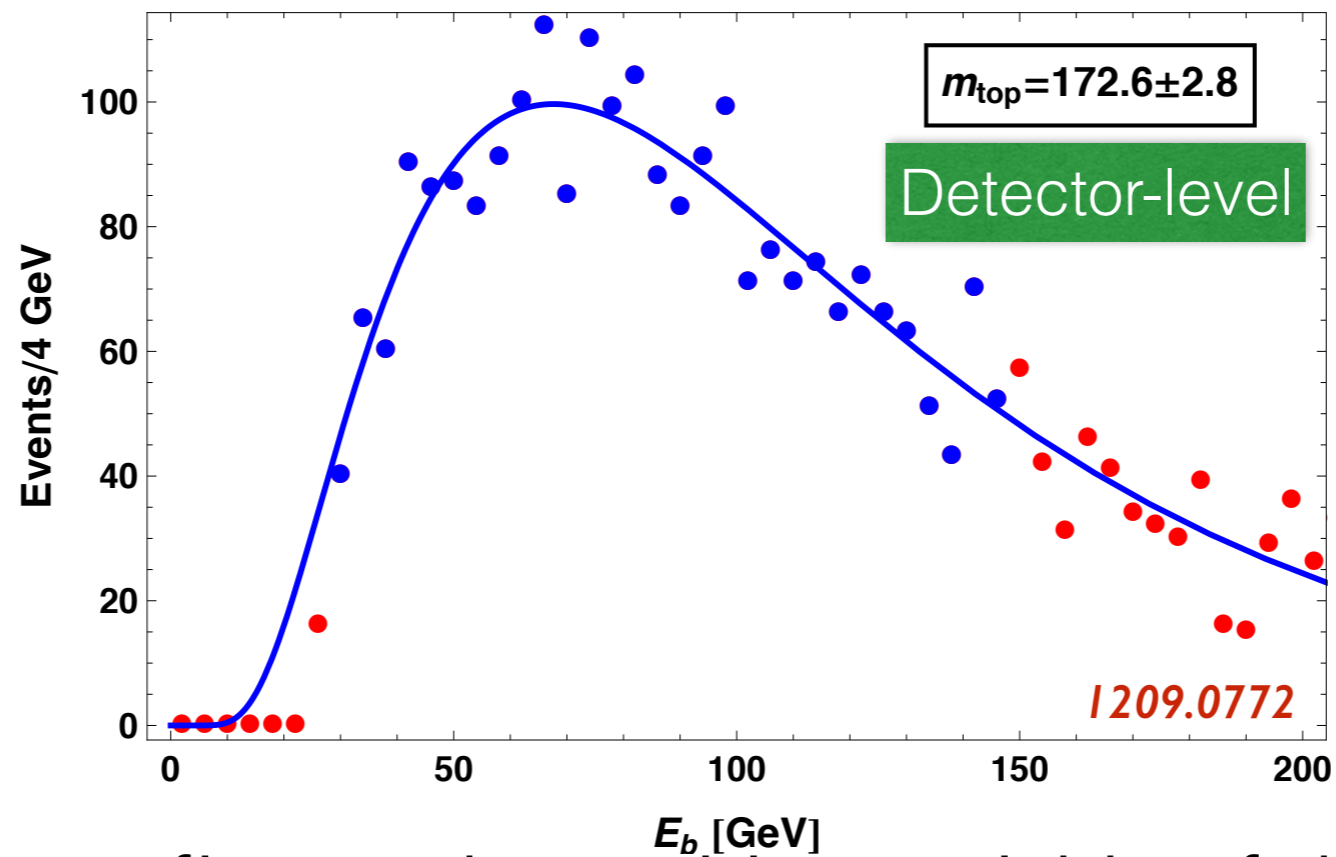
1209.0772 - Agashe Franceschini and Kim

message: LO effects are well under control → CMS at work!



# b-jet energy (LO+PS)

100 pseudo-experiments from [MadGraph5+Pythia6.4+Delphes](#) (**ATLAS-2012-097**)



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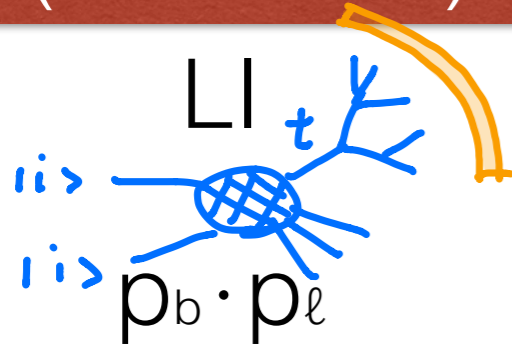
$$m_{\text{top}} = 173.1 (1 \pm \alpha/\pi) \pm 2.5 \text{ GeV (stat)}$$

1209.0772 - Agashe Franceschini and Kim

message: LO effects are well under control → CMS at work!

# variations around Lorentz Invariance

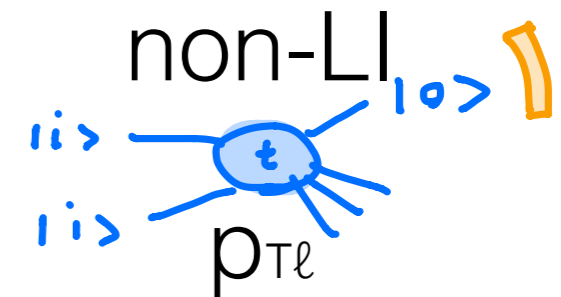
needs two  
particles  
(combinations)



needs just one particle

“pheno”-LI

$$\hat{E}_b$$



radiation in decays  
breaks true-LI due to  
reconstruction

radiation in decays  
breaks pheno-LI  
due to 3-body

end-point is safe w.r.t  
radiation in decay

exclusiveness  
breaks pheno-LI

in practice we need the  
tail, which is sensitive to  
radiation

what is the “small parameter”  $\Delta_{TH}$   
that “breaks” (true or effective) LI?

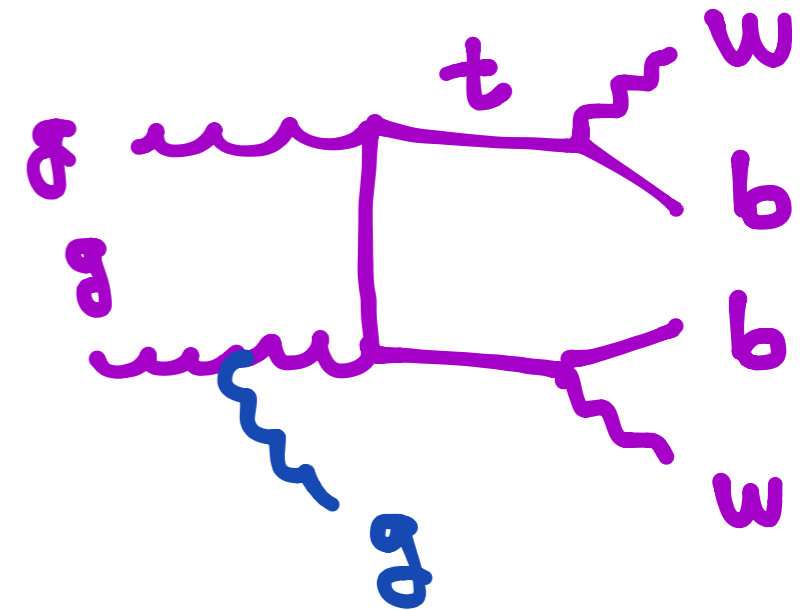
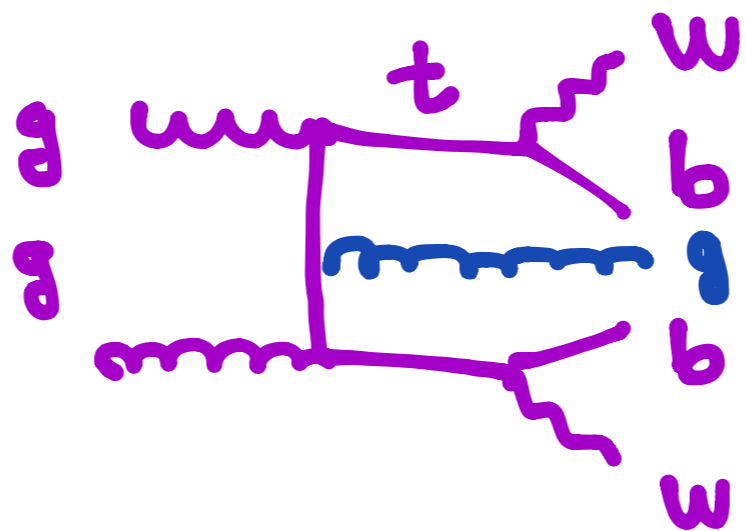
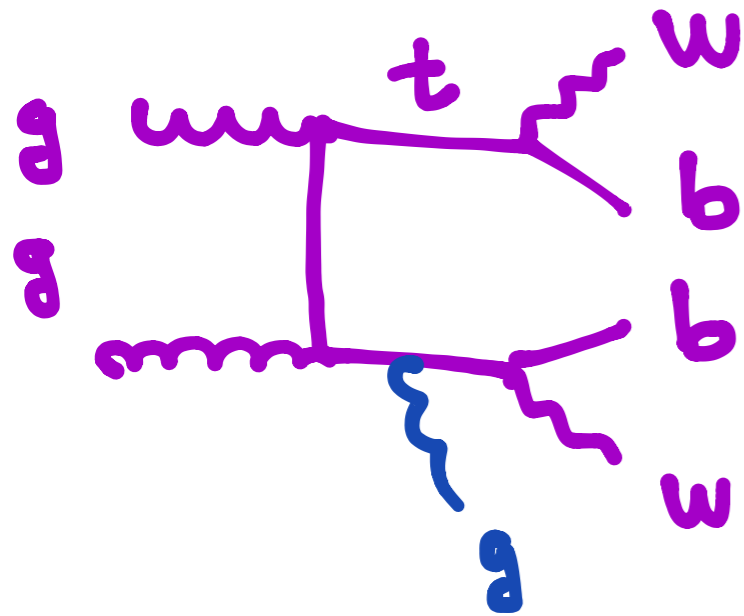
very encouraging LO  
result with b-jet energy

after having explored a number of **new physics applications** of this idea

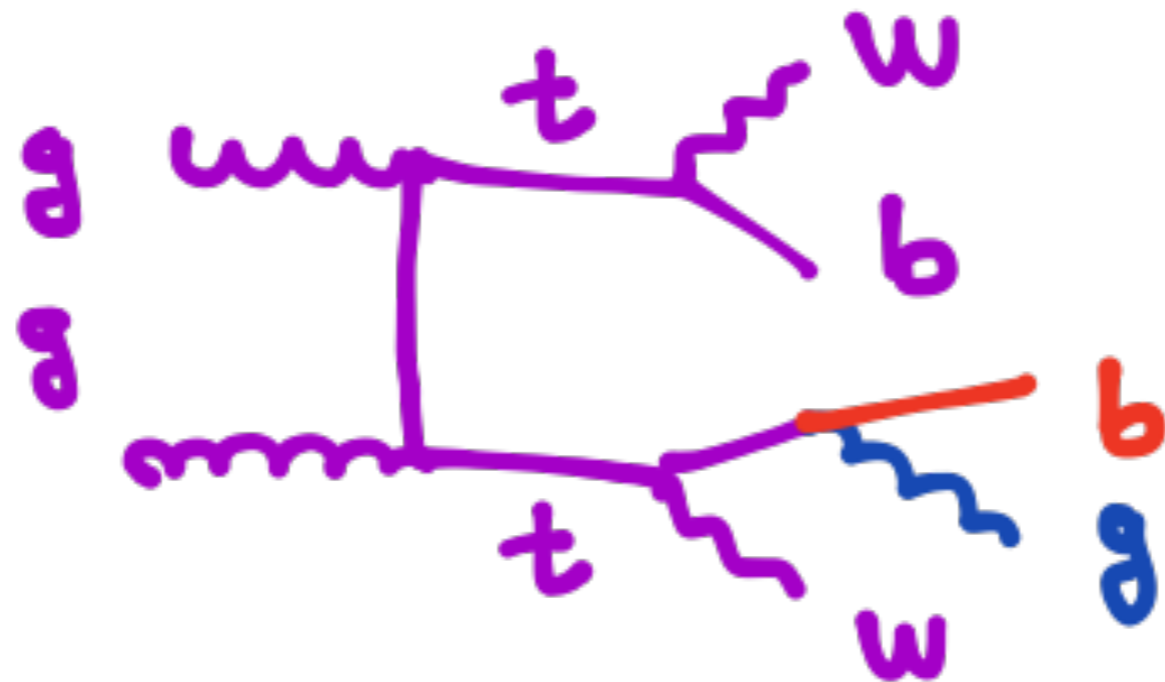
- 1212.5230 - Agashe, RF, Kim, Wardlow
- 1309.4776 - Agashe, RF, Kim
- 1403.3399 - Chen, Davoudiasl, Kim
- 1503.03836 - Agashe, RF, Kim, Wardlow
- Agashe, RF, Kim, Hong - WIP

study of perturbative effects at  
fixed NLO nearing completion

# corrections to the production mechanism



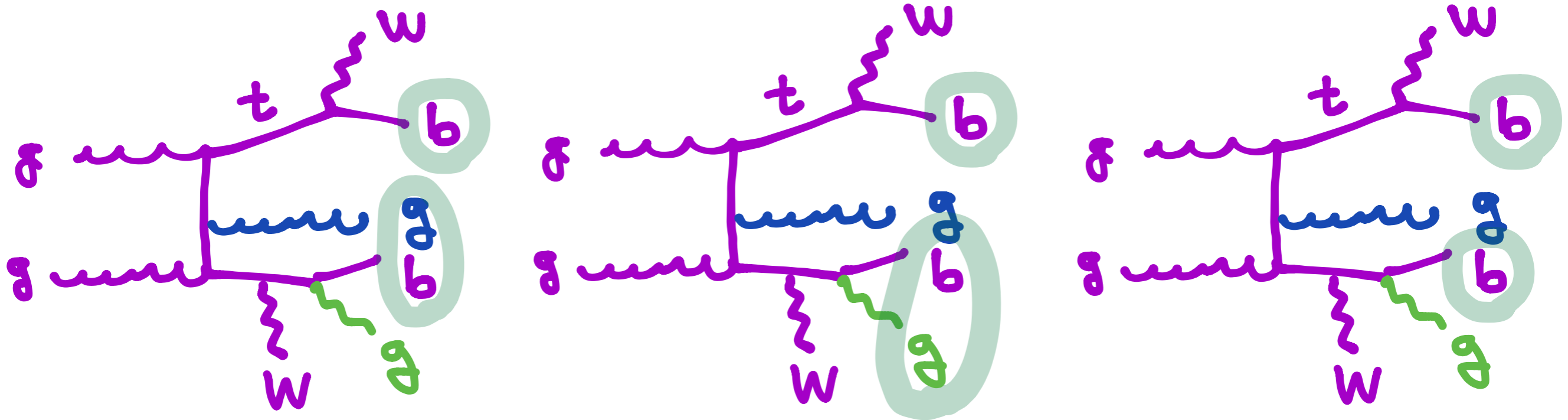
# corrections to the top decay



# NLO: production & decay

(MCFM)

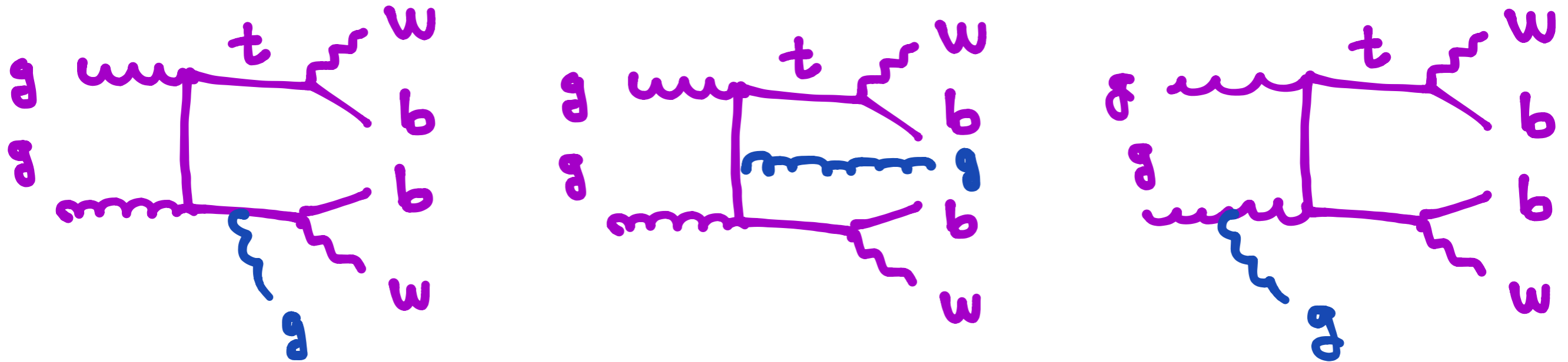
*Agashe, Franceschini, Kim, Schulze - in preparation*



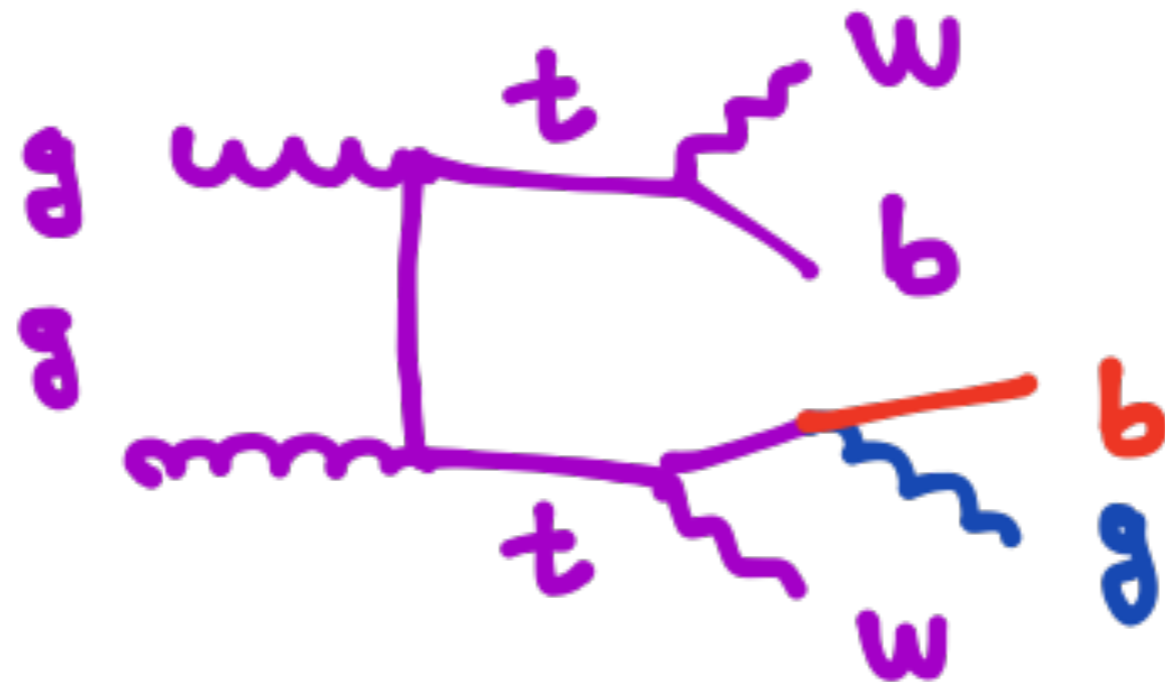
anti-kT jets

- resolved gluon from the top decay
- merged “extraneous” gluon
- reclustered bottom-gluon jets

# Energy peak protected from some NLO

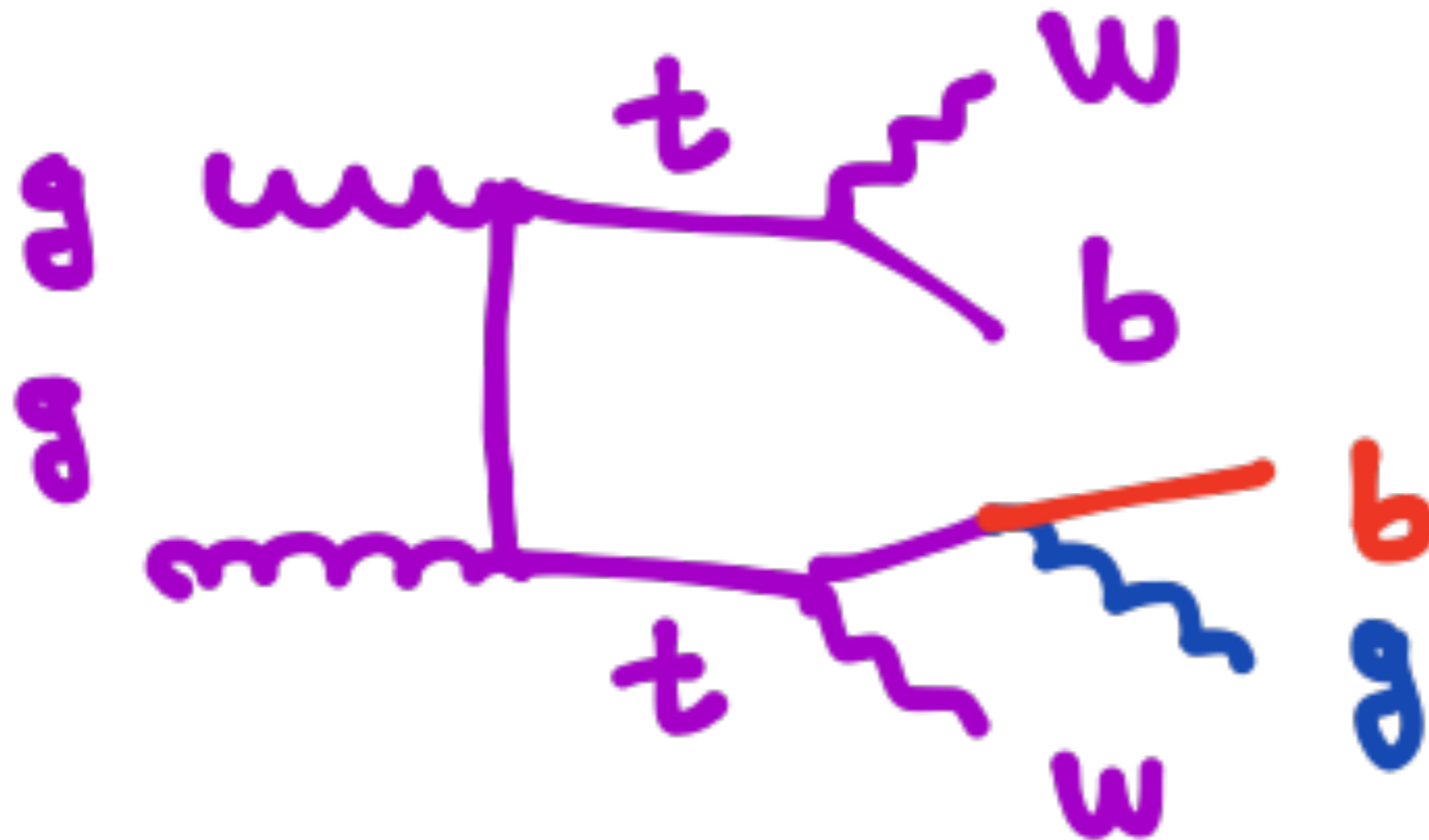


need to compute radiation in decay

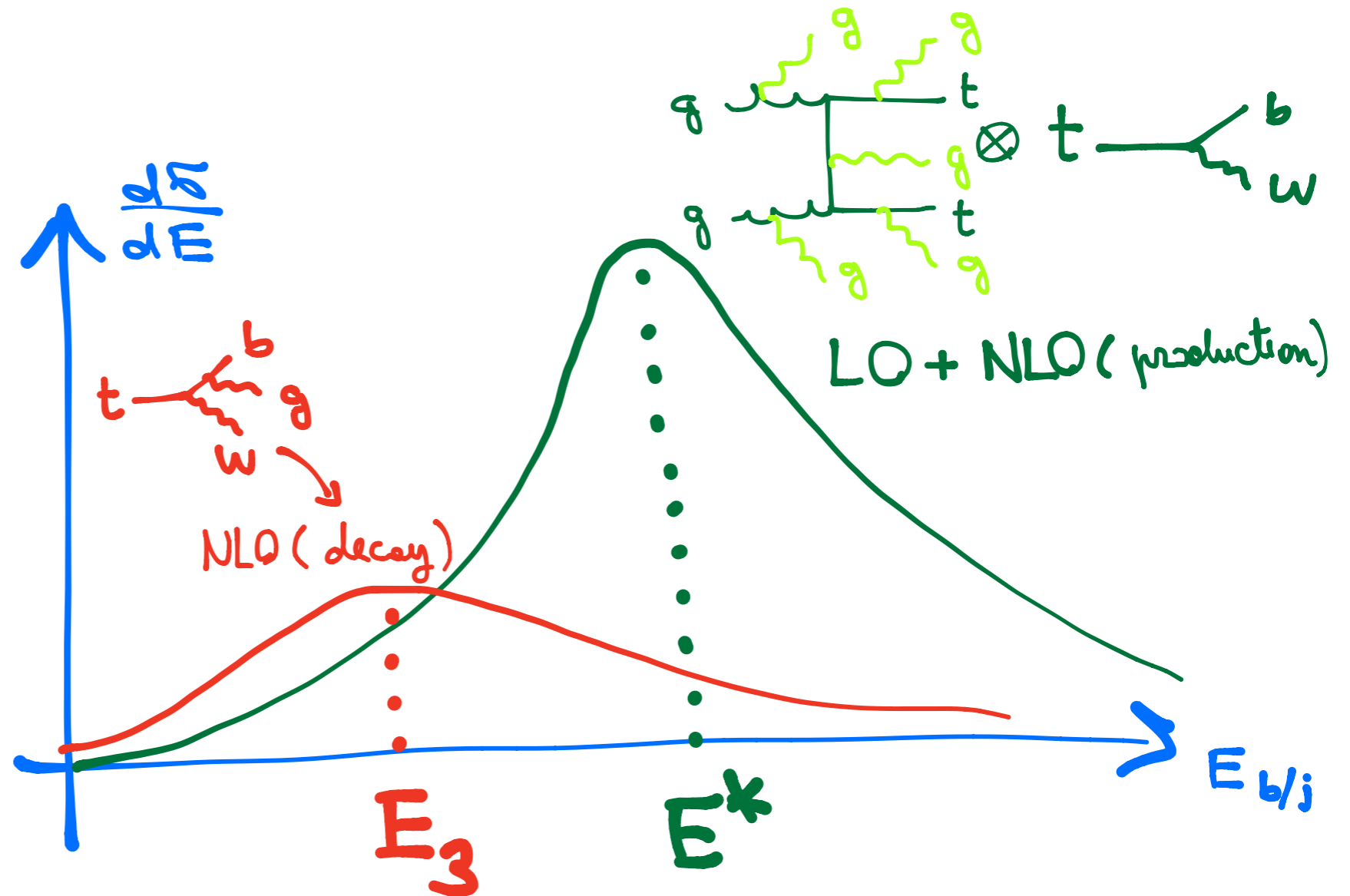


# Decay at NLO

*Agashe, Franceschini, Kim, Schulze - in preparation*



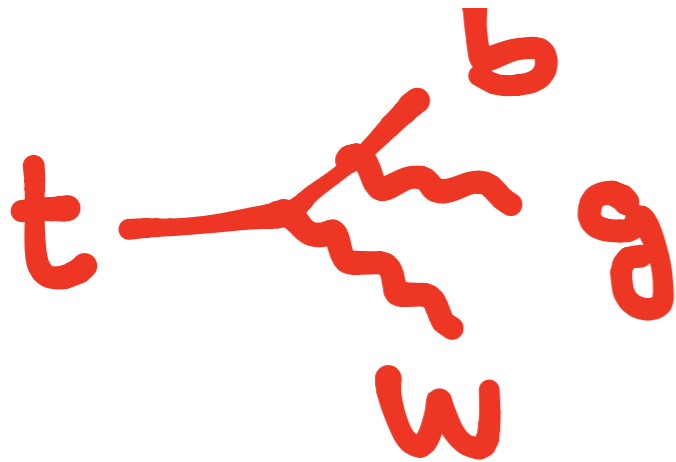
# Peak shift at NLO



$$E^{\text{peak}} = E^* + O(1) \frac{\alpha}{4\pi} E_3$$

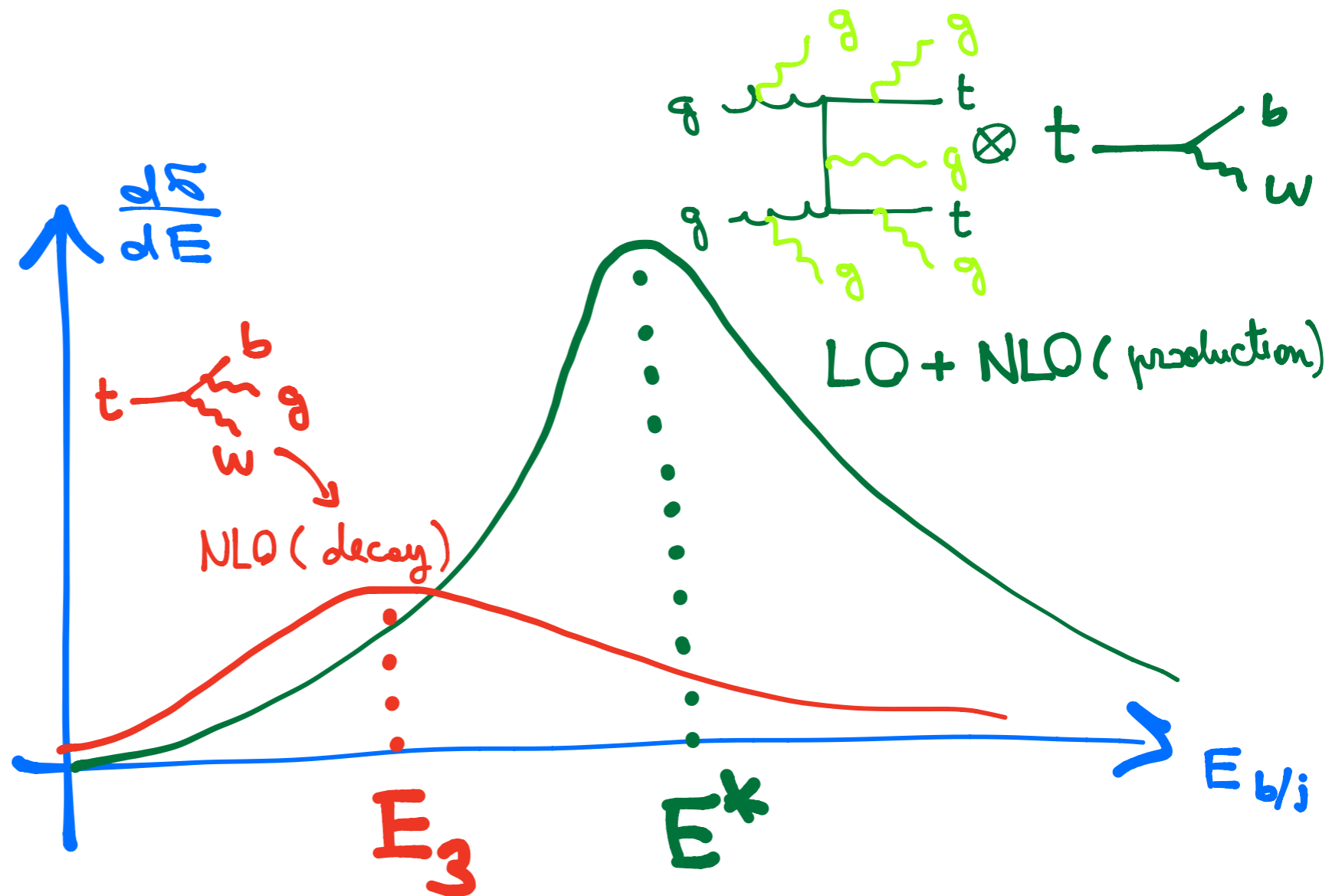


# Peak shift at NLO



BR( $t \rightarrow bWg$ )  
MadGraph5@LO

hard glue	Br
$p_T > 30 \text{ GeV}$ $dR > 0.2$	0.061
$p_T > 30 \text{ GeV}$ $dR > 0.4$	0.043
$p_T > 20 \text{ GeV}$ $dR > 0.2$	0.10
$p_T > 20 \text{ GeV}$ $dR > 0.4$	0.074

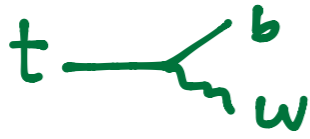


$$E^{\text{peak}} = E^* (1 - \Delta_{\text{TH}}) + \Delta_{\text{TH}} E_3$$

$$\Delta_{\text{TH}} = \text{BR}(t \rightarrow bWg) / \text{BR}(t \rightarrow bW) \approx 0.05$$

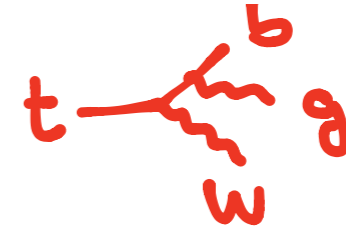
# NLO: production & decay

(MCFM)

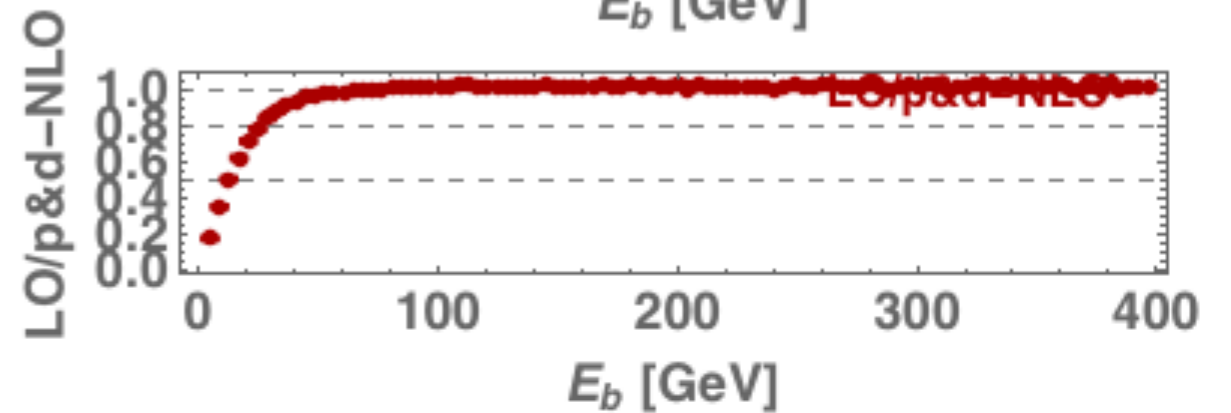
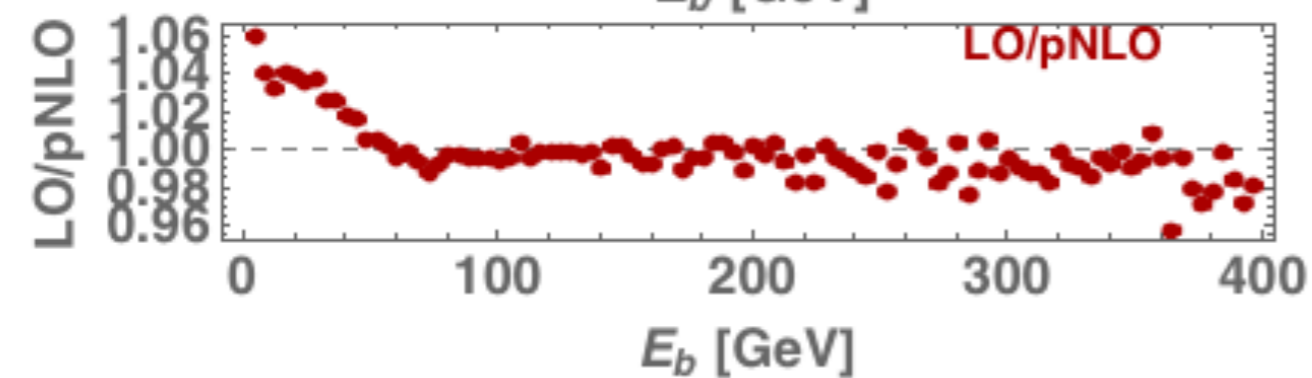
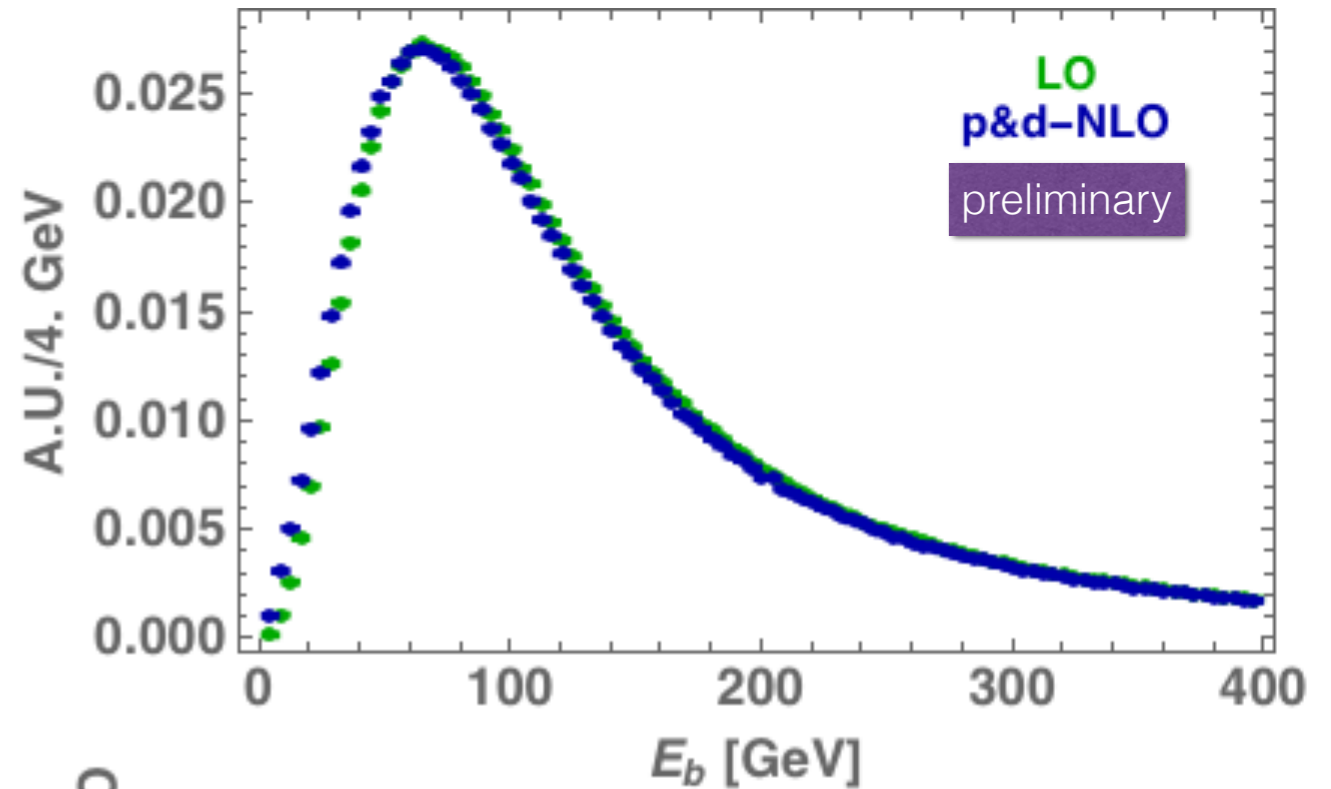
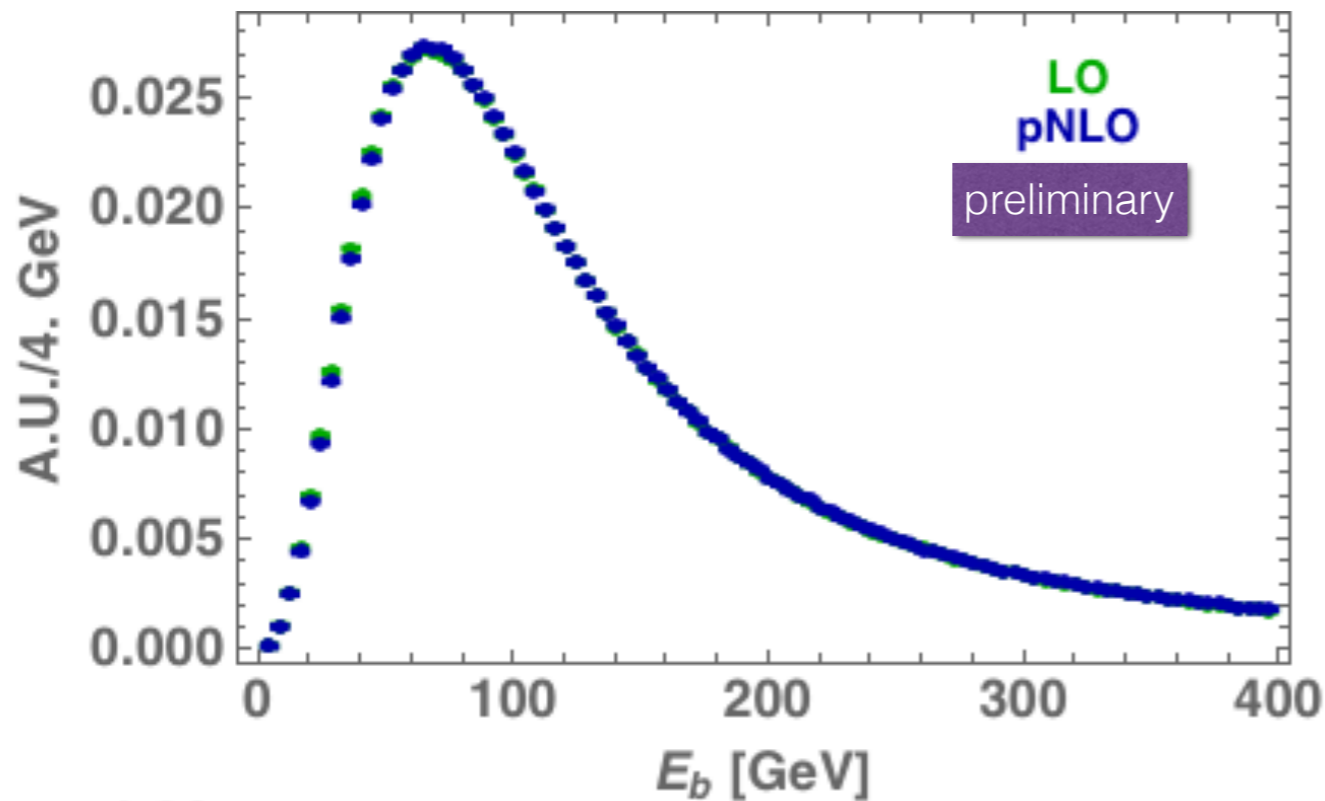


decay at LO

Energy of b



decay at NLO



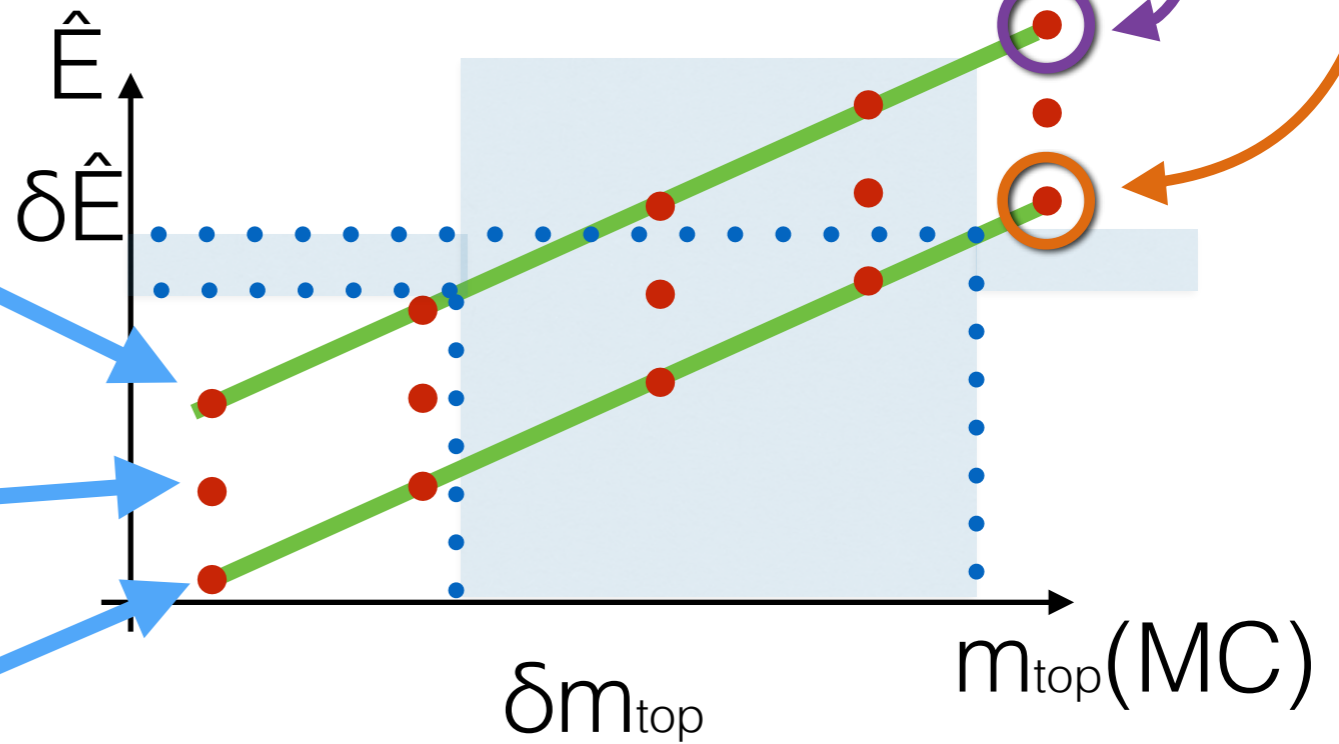
pQCD prediction:  $\hat{E}(m_{\text{top}})$

1. pick top pole mass
2. pick ren./fact. scales
3. energy distribution  $d\sigma/dE_b$
4. peak of the distribution  $\hat{E}$
5.  $\hat{E}(m_{\text{top}})$

pQCD

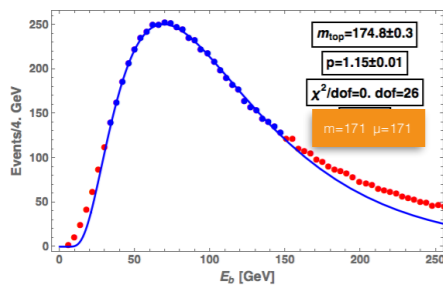
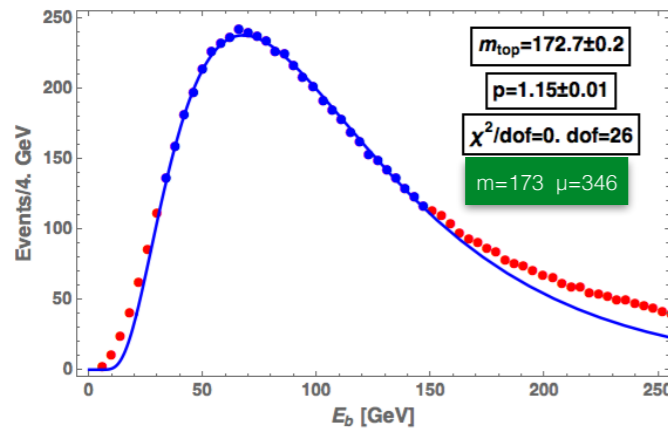
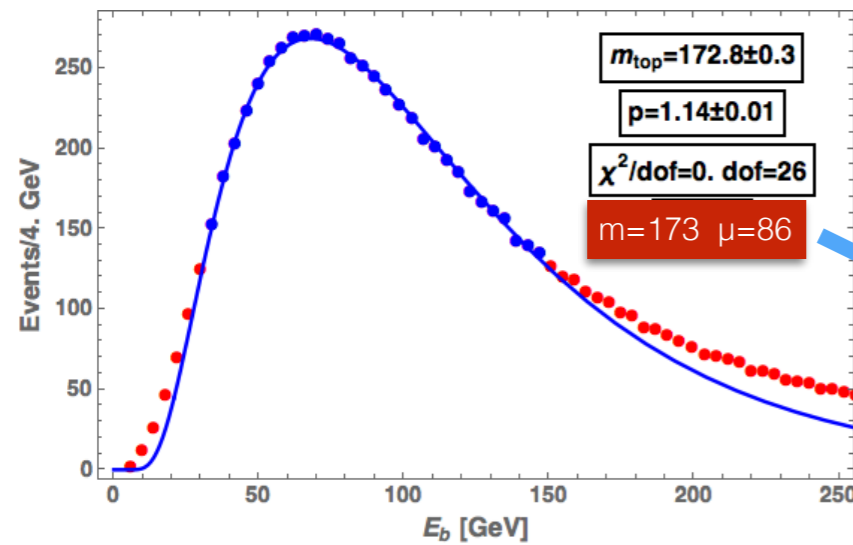
energy peaks

$$\mu \in [\mu_{\text{low}}, \mu_{\text{high}}]$$



Best:

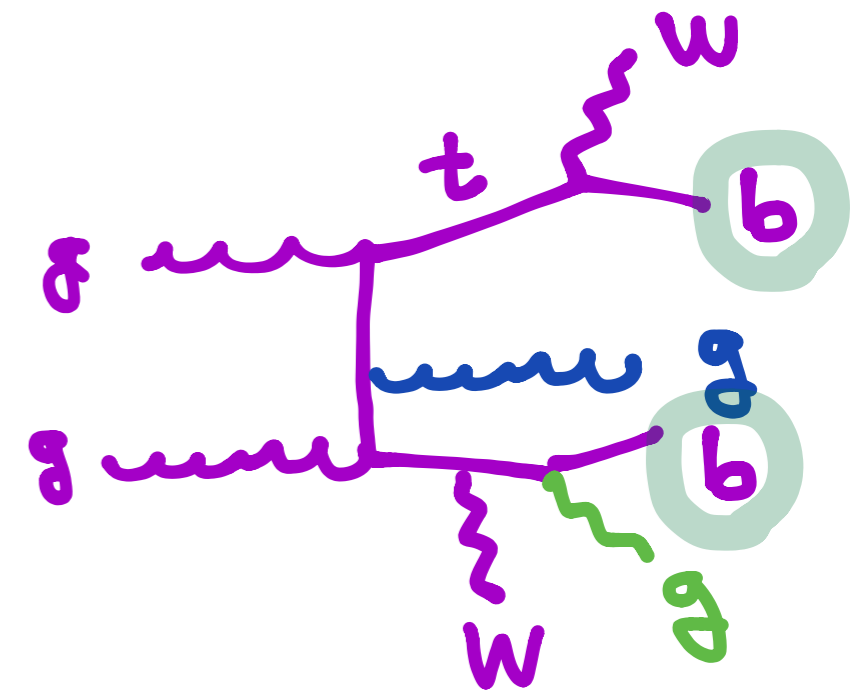
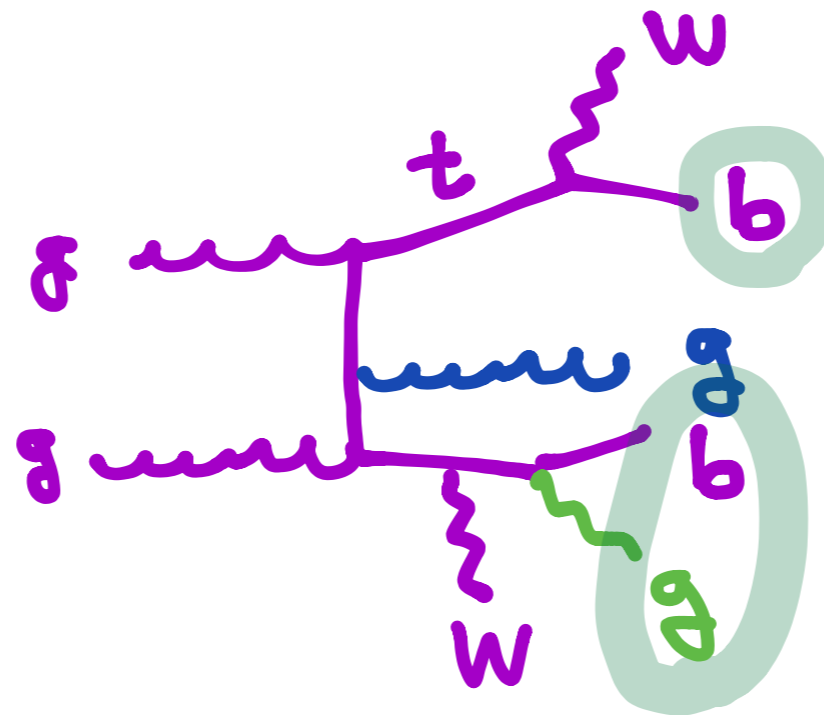
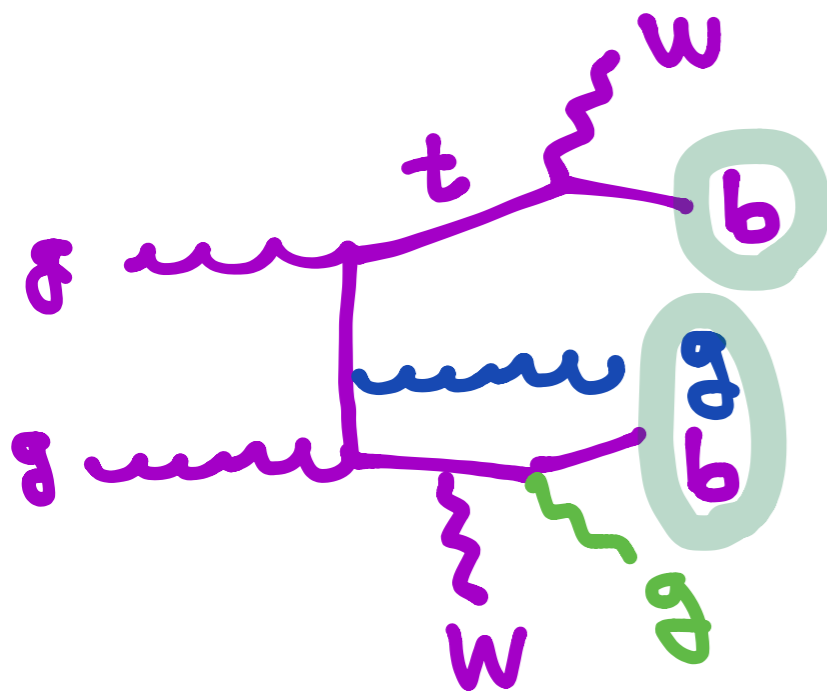
- narrow band between  $\mu_{\text{high}}$  and  $\mu_{\text{low}}$
- steep E vs.  $m_{\text{top}}$  
$$E_b^* = \frac{m_t^2 - m_w^2 + m_b^2}{2m_t}$$



# NLO: production & decay

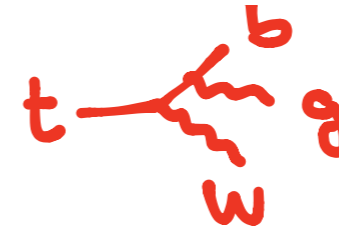
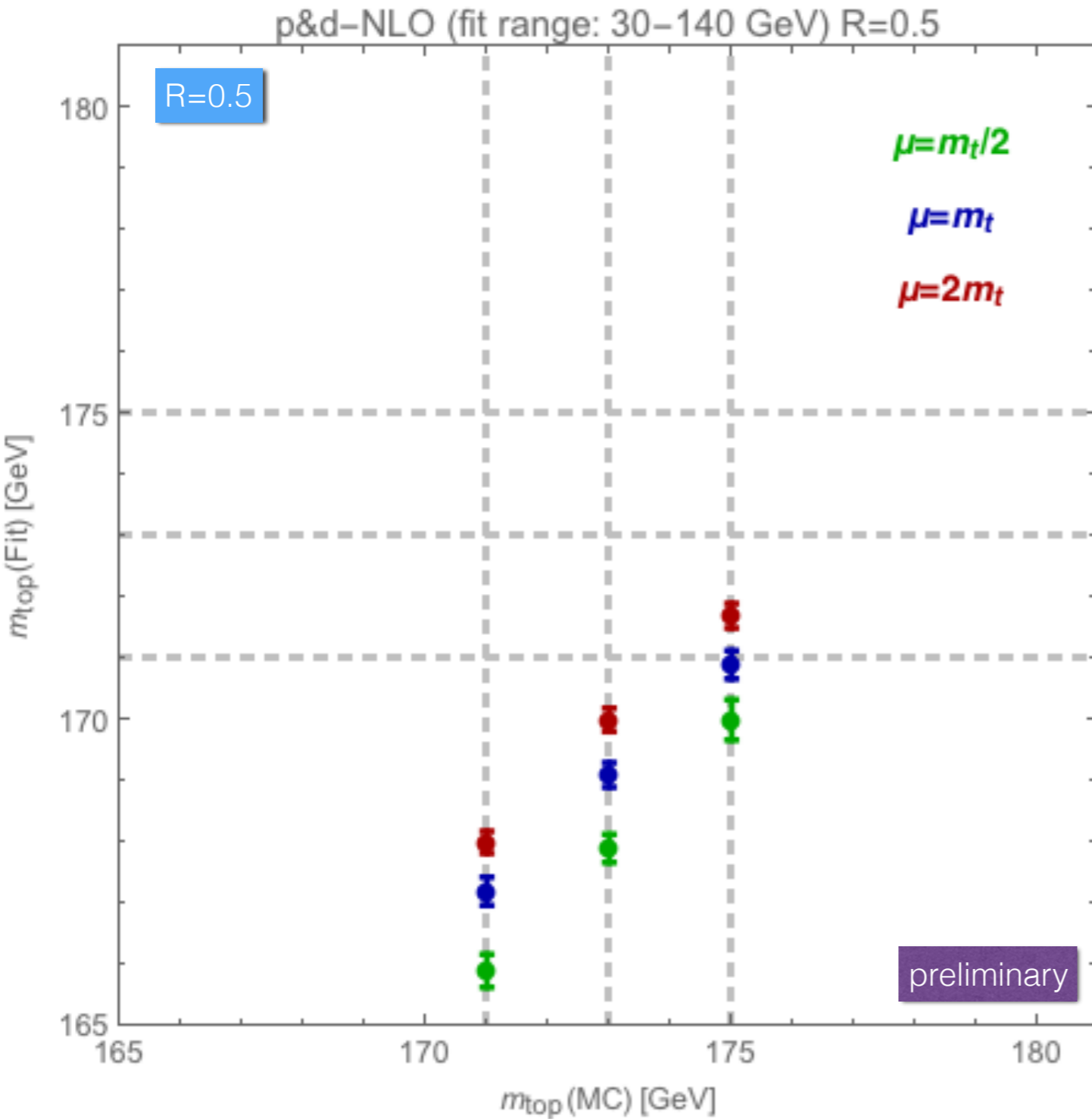
(MCFM)

*Agashe, Franceschini, Kim, Schulze - in preparation*

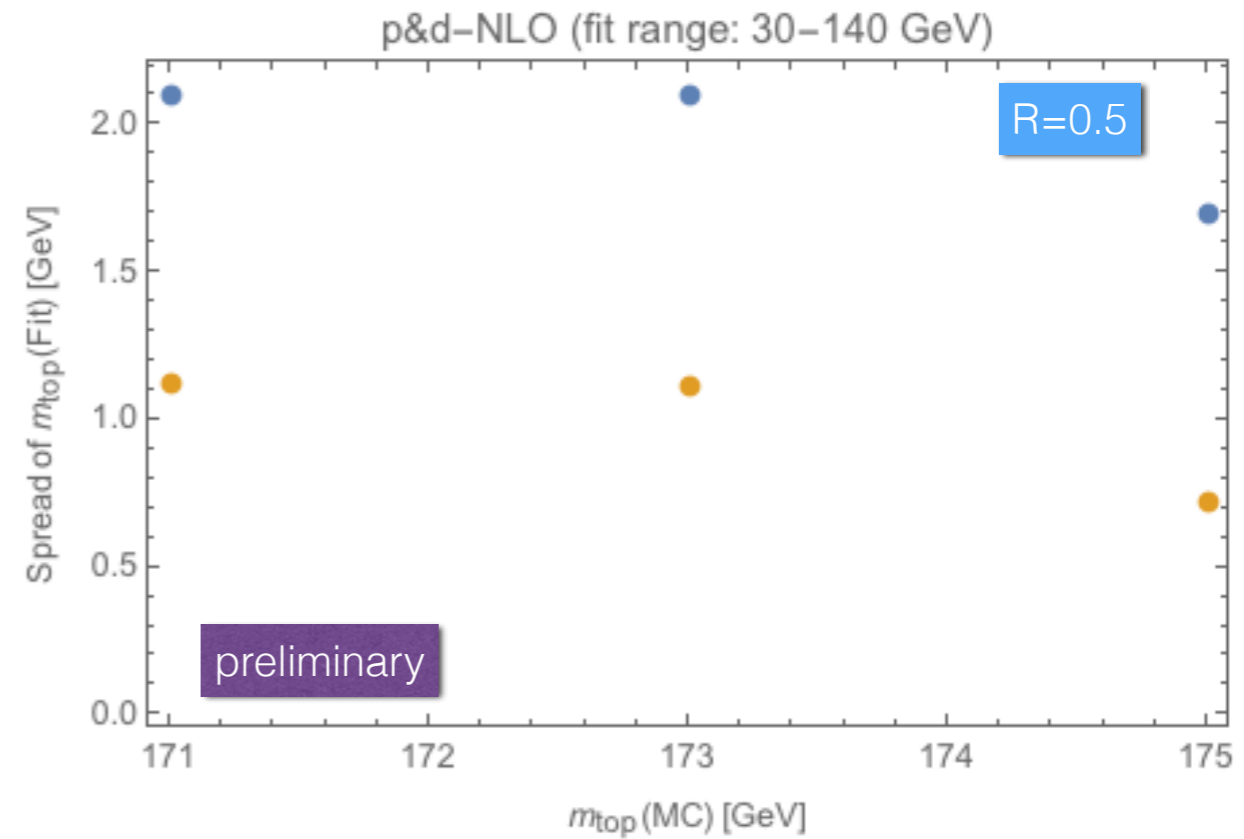


# NLO: production & decay (anti-kT R=0.5)

(MCFM)



$$|\delta| \sim \alpha_3 \sim 1/\mu$$

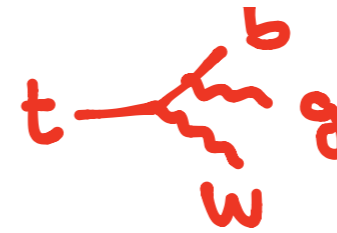
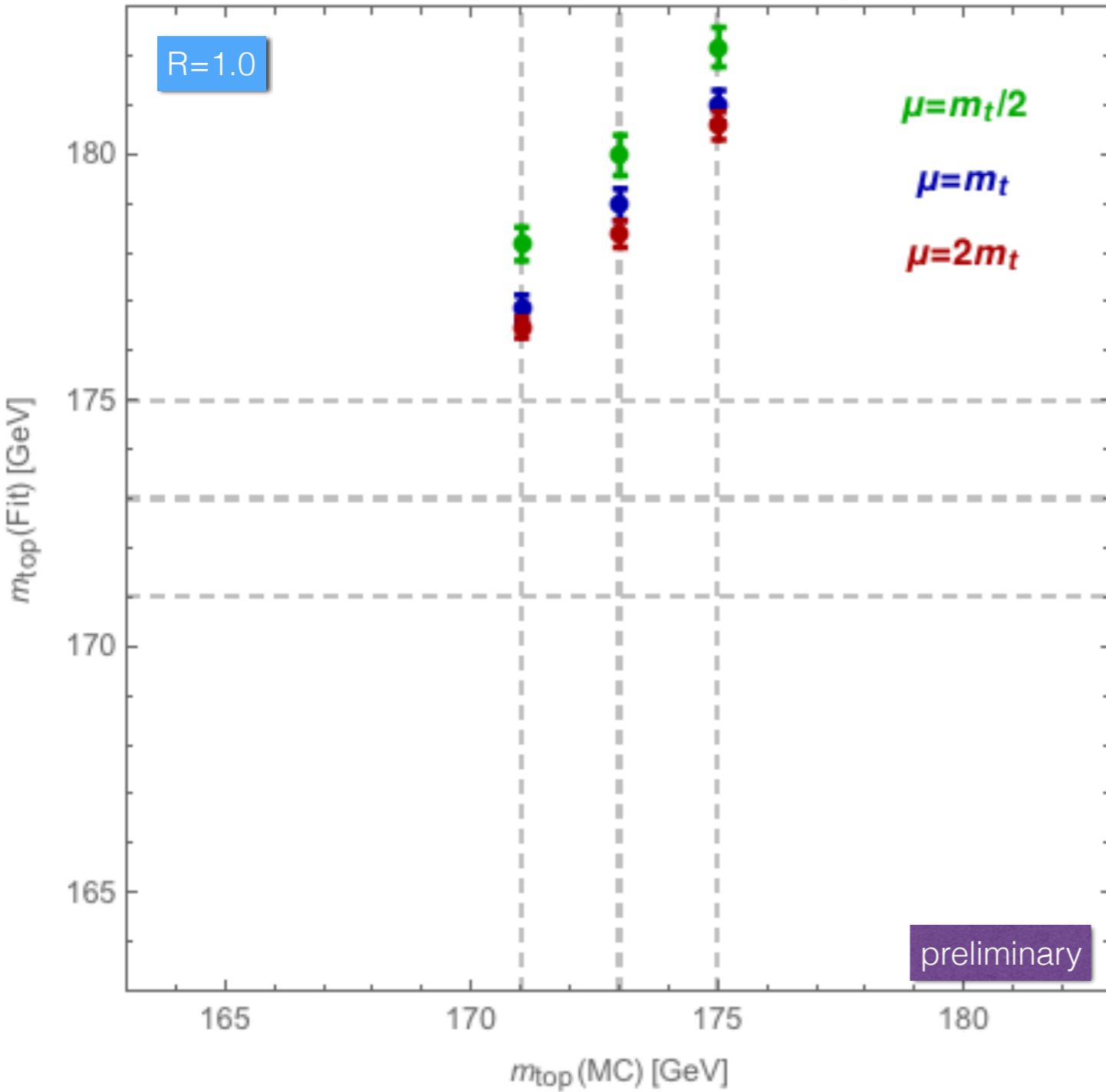


decay NLO sensitive to the scale choice:  $\pm 1$  GeV on  $m_{\text{top}}$

# NLO: production & decay (anti-kT R=1.0)

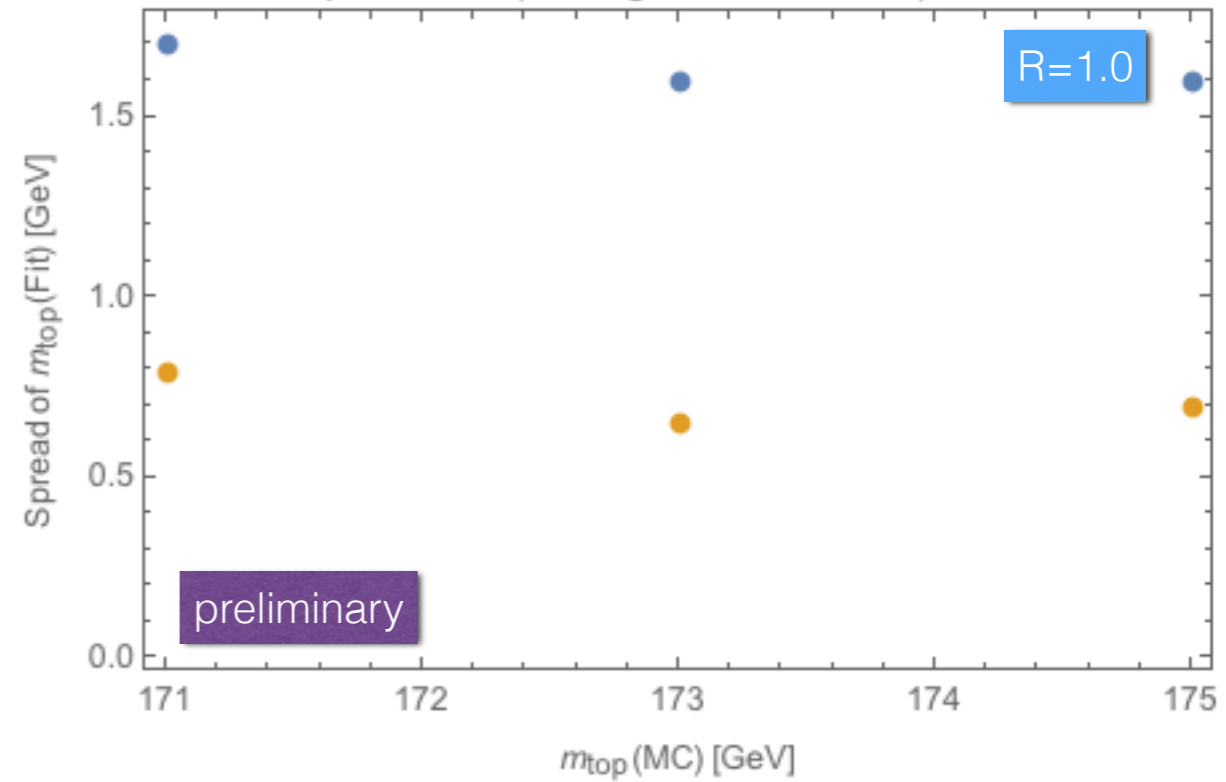
(MCFM)

p&d-NLO (fit range: 30–140 GeV) R=1.



$$|\delta| \sim \alpha_3 \sim 1/\mu$$

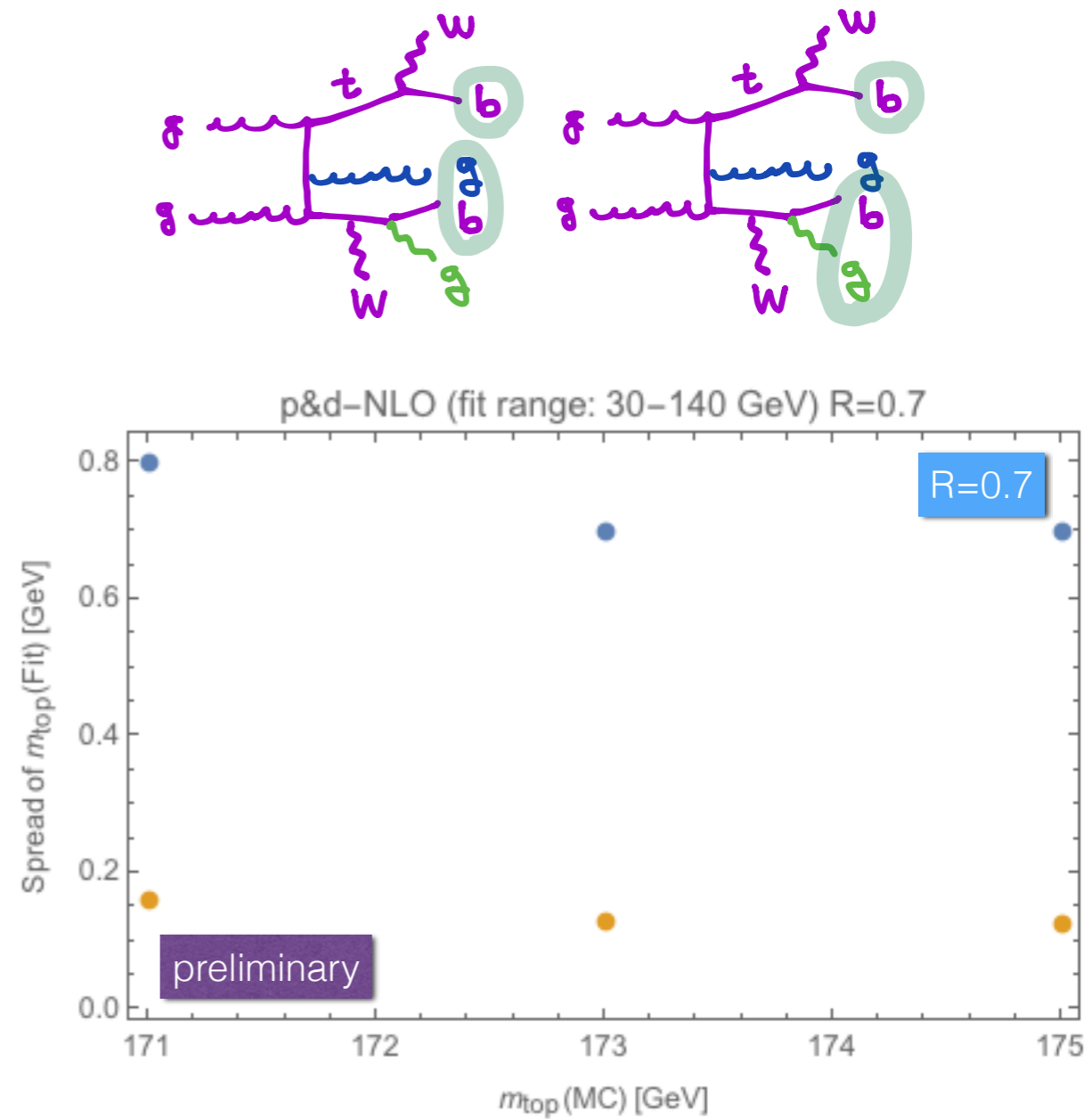
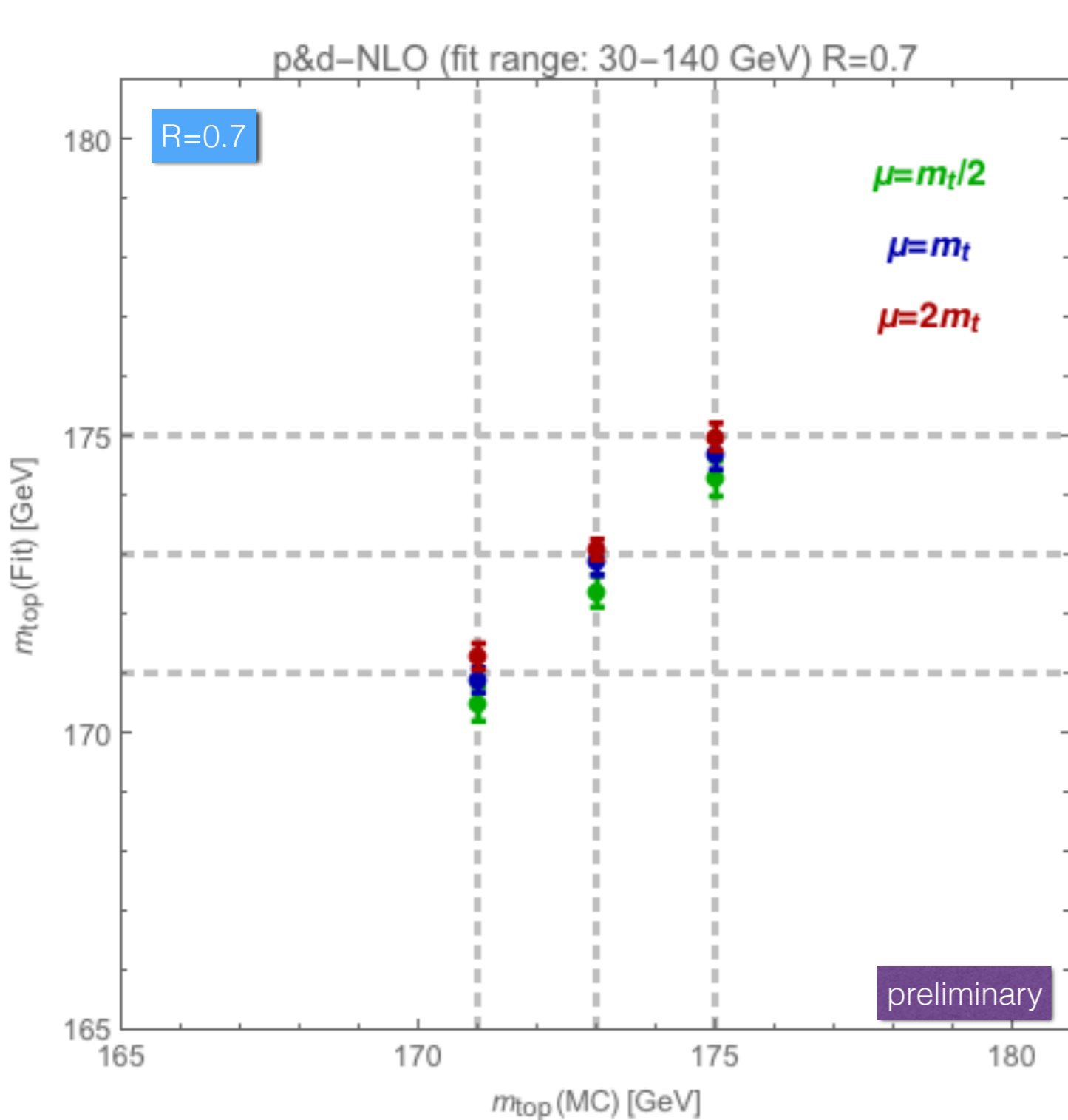
p&d-NLO (fit range: 30–140 GeV) R=1.



decay NLO sensitive to the scale choice:  $\pm 1$  GeV on  $m_{\text{top}}$

# NLO: production & decay (anti-kT R=0.7)

(MCFM)

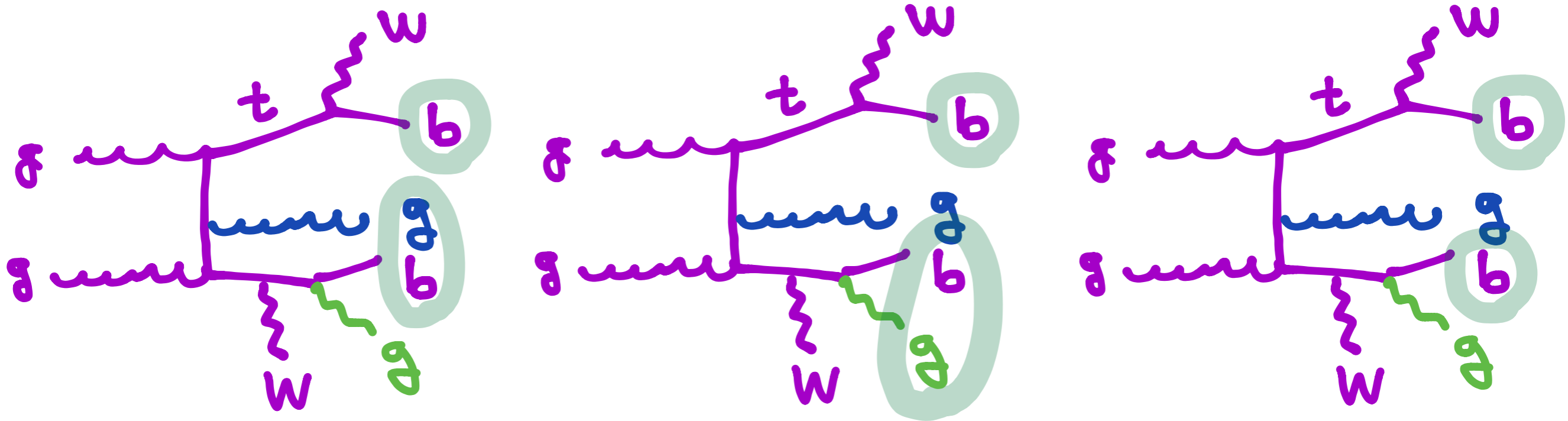


decay NLO sensitive to the scale choice:  $\pm 0.5$  GeV on  $m_{\text{top}}$

# NLO: production & decay

(MCFM)

Agashe, Franceschini, Kim, Schulze - in preparation



$$pp \rightarrow t\bar{t}|_{NLO} \times t \rightarrow b\nu|_{LO}$$

$m_t = 173 \text{ GeV}$	$R = 0.5$	$R = 0.7$	$R = 1.0$
$\mu = 2m_t^{(pole)}$	174.3(1)	175.9(1)	179.4(2)
$\mu = m_t^{(pole)}$	174.5(2)	176.3(2)	180.3(2)
$\mu = m_t^{(pole)}/2$	174.7(2)	176.9(2)	181.5(2)
$\delta_{th}$	$\pm 0.2 \text{ GeV}$	$\pm 0.5 \text{ GeV}$	$\pm 1.0 \text{ GeV}$

$$pp \rightarrow t\bar{t}|_{NLO} \times t \rightarrow b\nu|_{NLO}$$

	$R = 0.5$	$R = 0.7$	$R = 1.0$
$\mu = 2m_t^{(pole)}$	170(2)	173.1(2)	178.4(3)
$\mu = m_t^{(pole)}$	169.1(2)	172.9(2)	179.0(3)
$\mu = m_t^{(pole)}/2$	167.9(2)	172.4(3)	180.0(3)
$\delta_{th}$	$\pm 1 \text{ GeV}$	$\pm 0.4 \text{ GeV}$	$\pm 0.8 \text{ GeV}$



# Mild corrections from NLO

*Agashe, Franceschini, Kim, Schulze - in preparation*

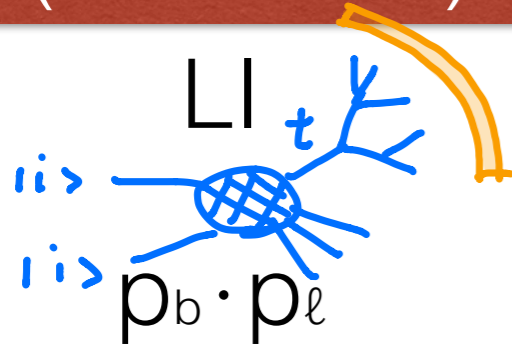
$$\hat{E} = E_{LO}^* \cdot \left[ 1 + f_{pol} + \epsilon_{FSR} \left( C_{bWg} + \underbrace{\delta_{int} + \delta_{PDFs} + \dots}_{\delta_{prod}} \right) \right]$$

$\leq 3 \cdot 10^{-3}$        $\leq 0.1$        $O(1)$

$$O_{NLO} = O_{LO} \cdot \left[ 1 + \underbrace{\delta_{int} + \delta_{PDFs} + \dots}_{\delta_{prod}} \right]$$

# variations around Lorentz Invariance

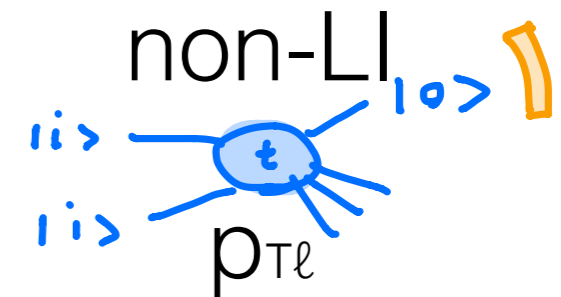
needs two particles  
(combinations)



needs just one particle

“pheno”-LI

$\hat{E}_b$



radiation in decays  
breaks true-LI due to  
reconstruction

radiation in decays  
breaks pheno-LI  
due to 3-body

end-point is safe w.r.t  
radiation in decay

exclusiveness  
breaks pheno-LI

in practice we need the  
tail, which is sensitive to  
radiation

what is the “small parameter”  $\Delta_{TH}$   
that “breaks” (true or effective) LI?

# More (B hadron) peak observables

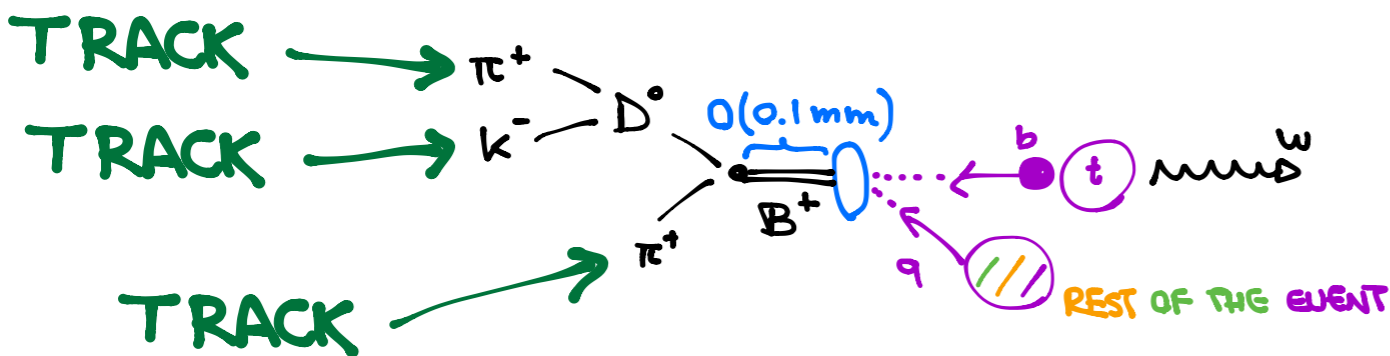
The strength of the future LHC top mass measurement will build on the **diversity of methods**  
⇒ not very useful to talk about “*single best measurement*”

$$\frac{d\sigma}{dE_b} \propto \frac{d\sigma}{d\gamma_b} \propto \frac{d\sigma}{d\lambda}$$

hadron energy peak

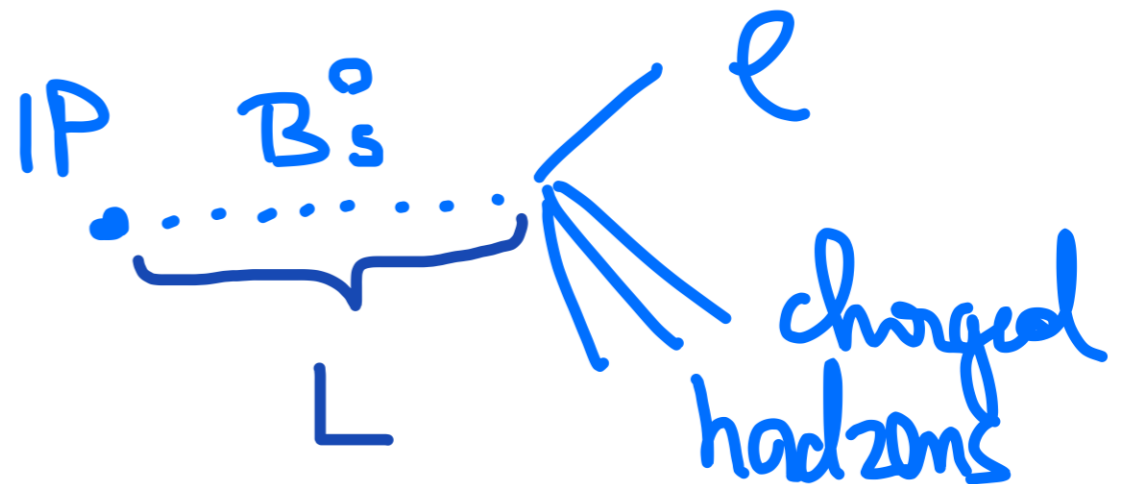
get the hadron energy entirely from tracks

collaboration with  
M. Schulze



mean decay path peak

discussions with  
J. Incandela



exclusive B decays in the top sample

# B hadron observables

B physics in the top sample

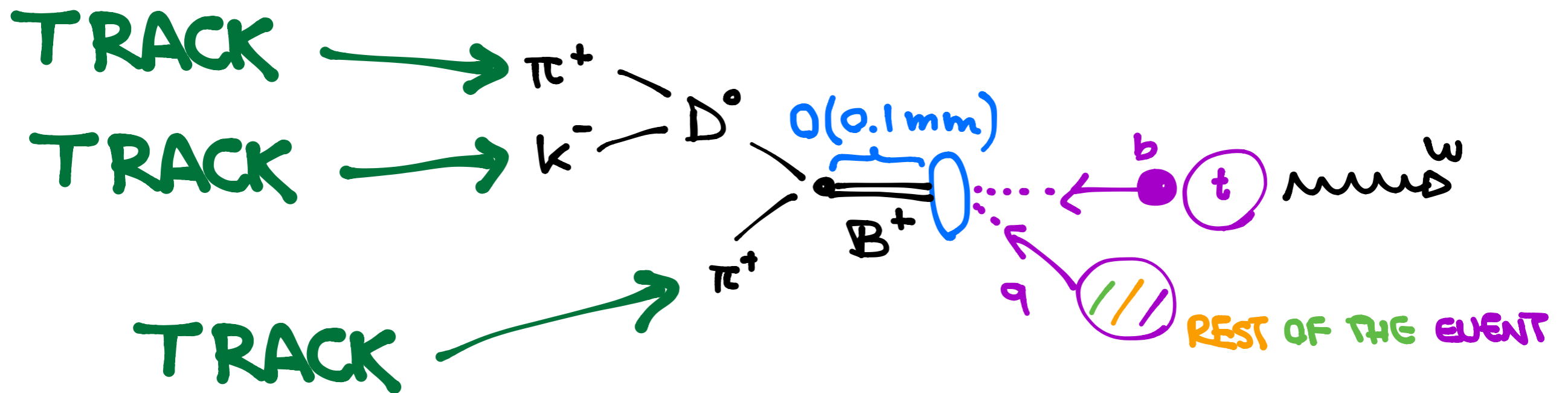
Fragmentation: the b quark energy peak is translated into a (broader) B hadron energy peak

- more exclusive final states
- non-JES uncertainties
- hadronization uncertainties

# B hadron

## energy peak

get the hadron energy entirely from tracks



$B^+ \rightarrow 3 \text{ TRACKS}$

# Exclusive Decay

(Fully reconstructible with tracks)

## J/psi modes

$$b \xrightarrow{\text{few} \cdot 10^{-3}} J/\psi + X \xrightarrow{10^{-1}} \ell\bar{\ell} + X$$

$$B_s^0 \rightarrow J/\psi \phi \rightarrow \mu^- \mu^+ K^+ K^- \quad 1106.4048$$

$$B^0 \rightarrow J/\psi K_S^0 \rightarrow \mu^- \mu^+ \pi^+ \pi^- \quad 1104.2892$$

$$B^+ \rightarrow J/\psi K^+ \rightarrow \mu^+ \mu^- K^+ \quad \begin{matrix} 1101.0131 \\ 1309.6920 \end{matrix}$$

$$\Lambda_b \rightarrow J/\psi \Lambda \rightarrow \mu^+ \mu^- p \pi^- \quad 1205.0594$$

J/psi but no need to require leptonic W decay

## D modes

$$B^0 \xrightarrow{3 \cdot 10^{-3}} D^- \pi^+ \xrightarrow{10^{-2}} K_S^0 \pi^- \pi^+$$

$$B^0 \xrightarrow{3 \cdot 10^{-3}} D^- \pi^+ \xrightarrow{10^{-2}} K^- \pi^+ \pi^- \pi^+$$

$$B^0 \xrightarrow{3 \cdot 10^{-3}} D^- \pi^+ \xrightarrow{3 \cdot 10^{-2}} K_S^0 \pi^+ \pi^- \pi^+$$

$$B^- \xrightarrow{5 \cdot 10^{-3}} D^0 \pi^- \xrightarrow{4 \cdot 10^{-2}} K^- \pi^+ \pi^-$$

$$B^- \xrightarrow{5 \cdot 10^{-3}} D^0 \pi^- \xrightarrow{2 \cdot 10^{-2}} K^{*-}(892) \pi^+ \pi^- \rightarrow K_S^0 \pi^- \pi^+ \pi^-$$

$$B^- \xrightarrow{5 \cdot 10^{-3}} D^0 \pi^- \xrightarrow{6 \cdot 10^{-3}} K_S^0 \rho^0 \pi^-$$

$$B^- \xrightarrow{5 \cdot 10^{-3}} D^0 \pi^- \xrightarrow{5 \cdot 10^{-3}} K^- \pi^+ \rho^0 \pi^-$$

# B hadron $\gamma$ boost factor

$$\frac{d\mathcal{L}}{dE_b} \propto \frac{d\mathcal{L}}{d\gamma_b}$$

hadron energy peak  $\longrightarrow$  hadron boost peak

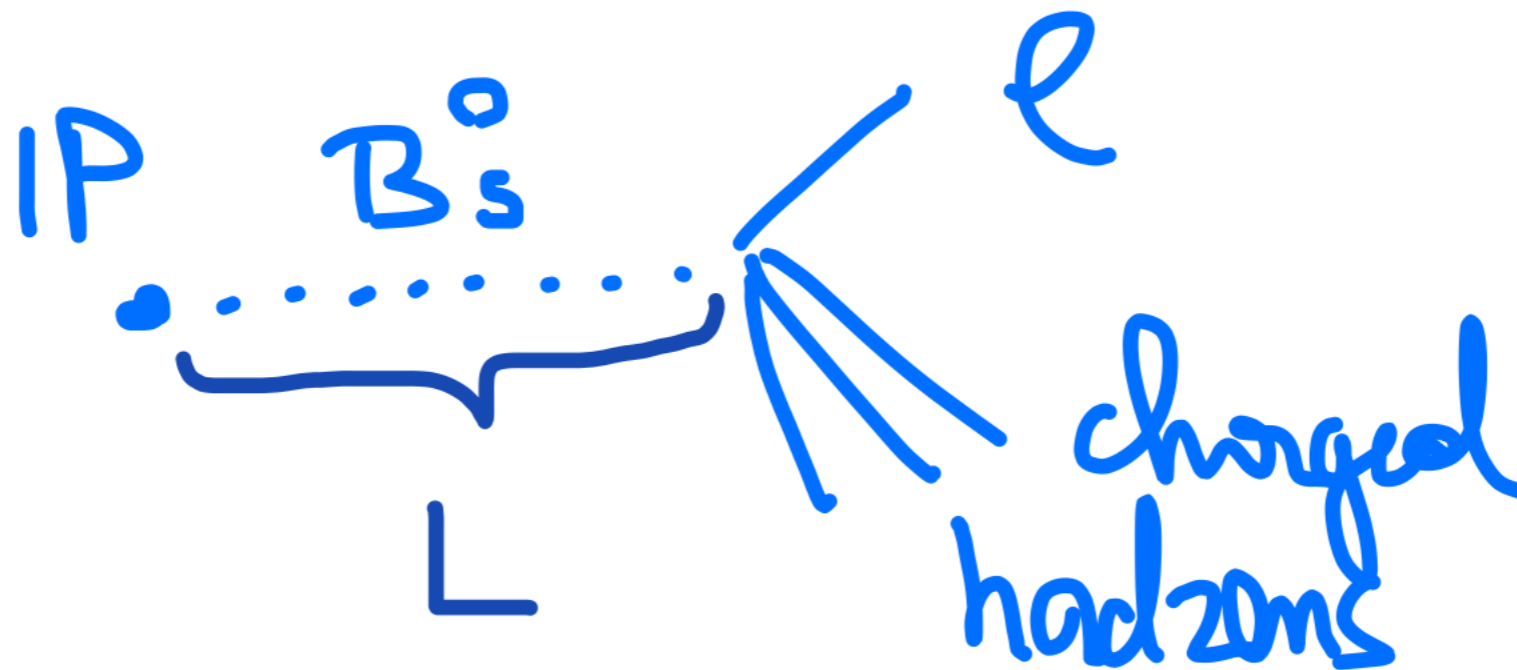
Does the **ratio**  $\gamma = E/m$  help to get rid of exp. uncertainties?

# 3D decay length

*discussion with J. Incandela*

Time of decays is harder to measure than the position

Experiments measure decay length  $L$



**Jet Energy Scale does not affect  $\lambda$ , nor  $L$**



# Mean decay length invariance

$$\gamma = E/m$$

- A peak in the energy distribution of the  $b$  quark implies a peak in the boost factor distribution
- Not so interesting because the boost is not measured directly

However ...

$$\tau'(\text{lab}) = \gamma\tau$$

For  $\beta=1$  is

$$\lambda = c\beta\tau'(\text{lab}) = c\tau E/m$$

$E$  and  $\lambda$   
distributions  
are the same up  
to a rescaling

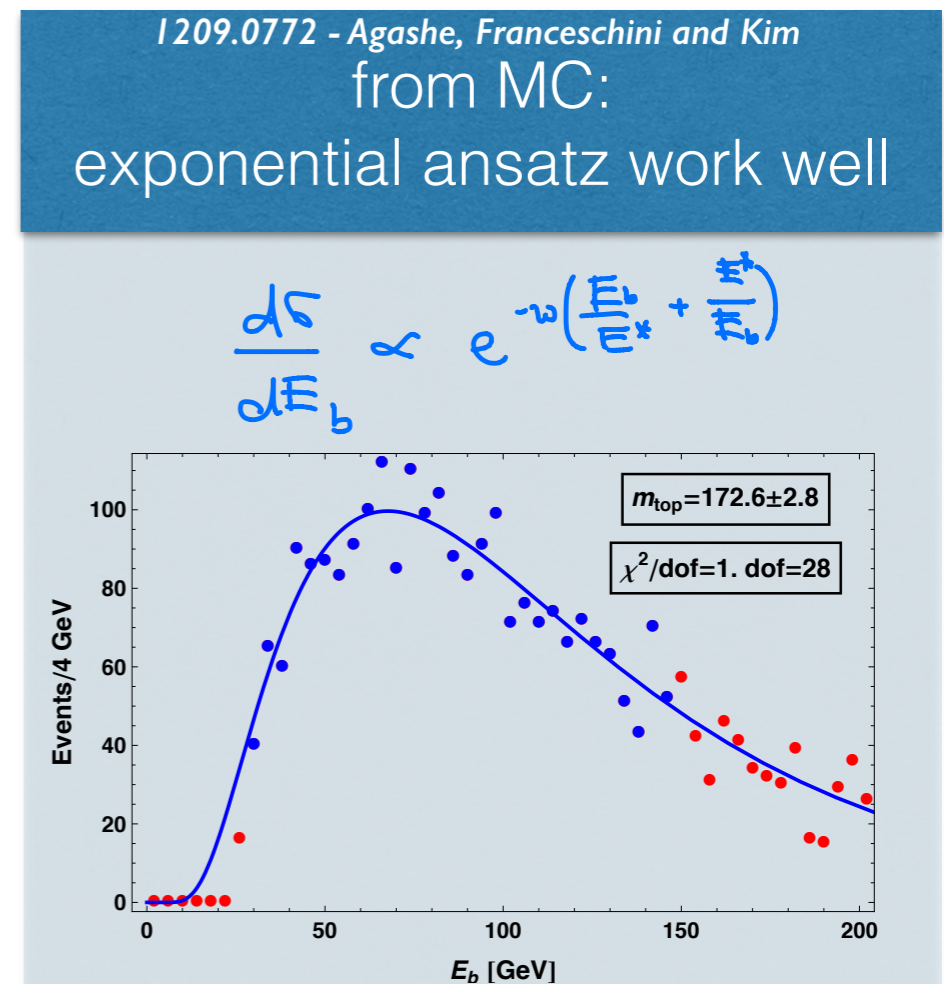
up to  $m^2/E^2$  effects the *mean* decay length of the  $b$  quark has a peak at the top rest frame value

# How to get the distribution of $\lambda$ from the observed $L$ ?

$$\frac{d\mathcal{L}}{dL} = \int e^{-L/\lambda} \otimes \text{pdf}(\lambda) d\lambda$$

For now we just predicted the mode of pdf( $\lambda$ )

$$\frac{d\mathcal{L}}{dE_b} \propto \frac{d\mathcal{L}}{d\gamma_b} \propto \frac{d\mathcal{L}}{d\lambda}$$



How to get the distribution of  $\lambda$  from the observed  $L$ ?

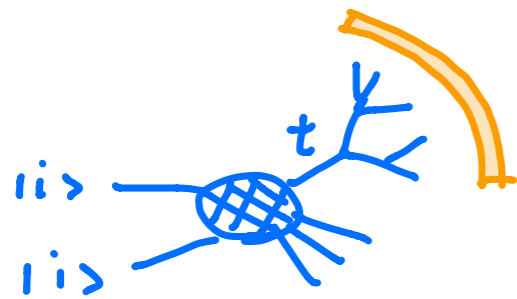
$$\frac{dS}{dL} = \int e^{-L/\lambda} \otimes \text{pdf}(\lambda) d\lambda$$

For now we just predicted the mode of  $\text{pdf}(\lambda)$

$$\text{pdf}(\lambda) = e^{-w \left( \frac{\lambda}{\lambda_0} + \frac{\lambda_0}{\lambda} \right)} ?$$

# (moral) Conclusions

## 1. Energy distributions as Breit-Wigner substitutes

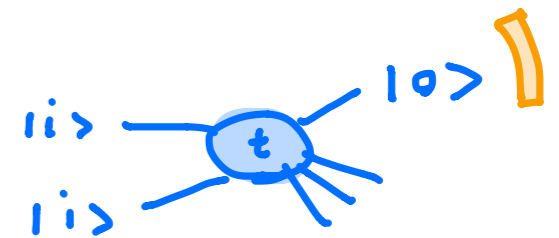


LI

$$\rho_b \cdot \rho_\ell$$

“pheno”-LI

$$\hat{E}_b$$



non-LI

$$\rho_{T\ell}$$

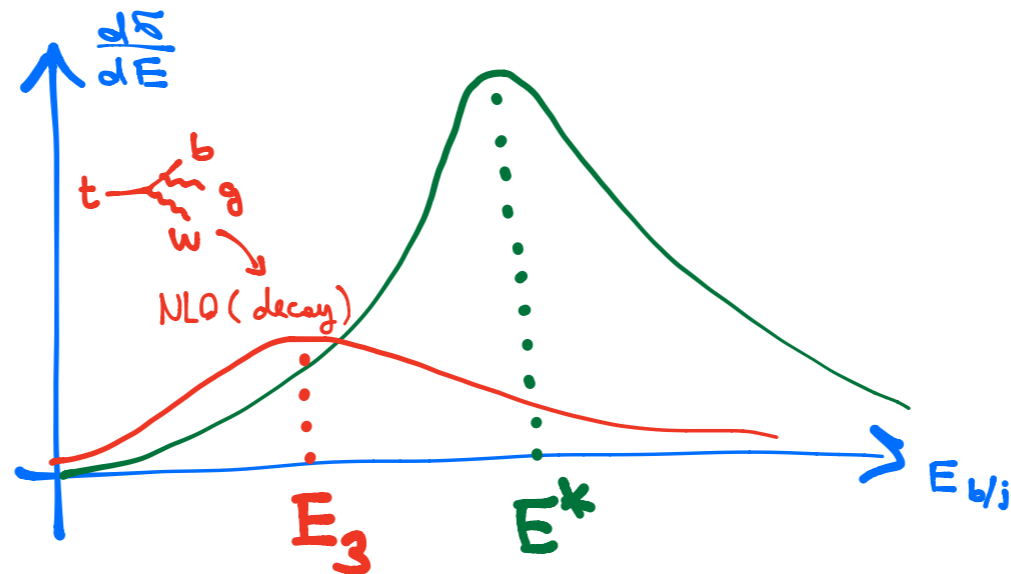


(moral) 

# Conclusions

## 2. Extensive program with b-jets and B-hadrons

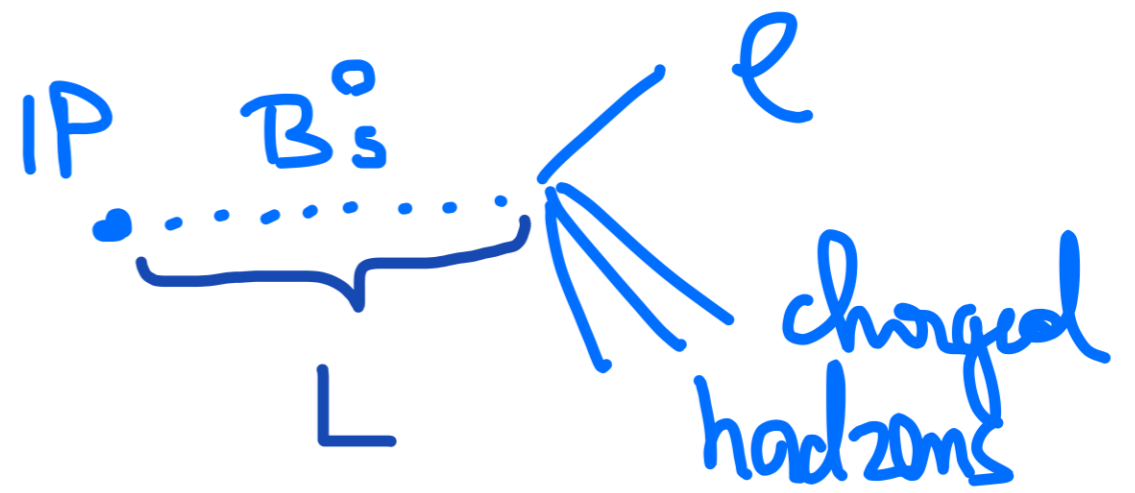
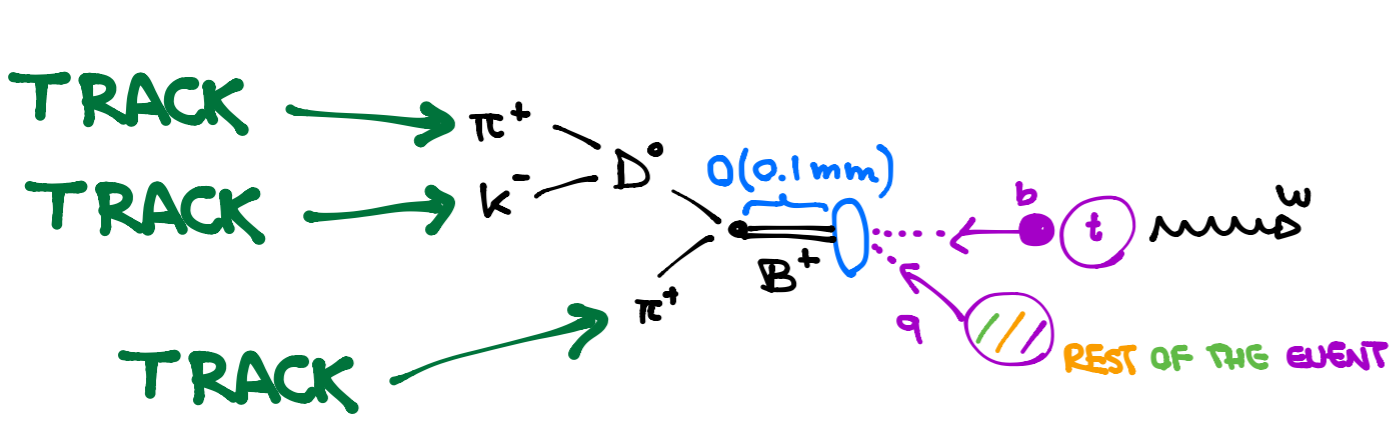
**b-jet energy peak**



**hadron energy peak**

get the hadron energy entirely from tracks

**mean decay path peak**



# (factual) Conclusions

## Peak of b-jet energy distribution

- “**invariance**” holds when only **NLO production** corrections are considered
- full NLO gives  **$\delta m_{\text{top}} \approx \pm 1$  GeV scale sensitivity** for any jet size parameter  $R$

$pp \rightarrow t\bar{t} _{NLO} \times t \rightarrow b\ell\nu _{LO}$				$pp \rightarrow t\bar{t} _{NLO} \times t \rightarrow b\ell\nu _{NLO}$			
$m_t = 173$ GeV	$R = 0.5$	$R = 0.7$	$R = 1.0$		$R = 0.5$	$R = 0.7$	$R = 1.0$
$\mu = 2m_t^{(pole)}$	174.3(1)	175.9(1)	179.4(2)	$\mu = 2m_t^{(pole)}$	170(2)	173.1(2)	178.4(3)
$\mu = m_t^{(pole)}$	174.5(2)	176.3(2)	180.3(2)	$\mu = m_t^{(pole)}$	169.1(2)	172.9(2)	179.0(3)
$\mu = m_t^{(pole)}/2$	174.7(2)	176.9(2)	181.5(2)	$\mu = m_t^{(pole)}/2$	167.9(2)	172.4(3)	180.0(3)
$\delta_{th}$	$\pm 0.2$ GeV	$\pm 0.5$ GeV	$\pm 1.0$ GeV	$\delta_{th}$	$\pm 1$ GeV	$\pm 0.4$ GeV	$\pm 0.8$ GeV

- chances that a **NNLO decay** description would be enough to make a solid prediction at  **$\delta m_{\text{top}} \approx 500$  MeV**

# To Do (in progress)

- check scale sensitivity at  $R \sim 0.82$   
(tt+jet @ NLO)
- check effects of cuts
- compare to moments of  $d\sigma/dE_b$
- B-hadron energy

# To Do (2)

explore:

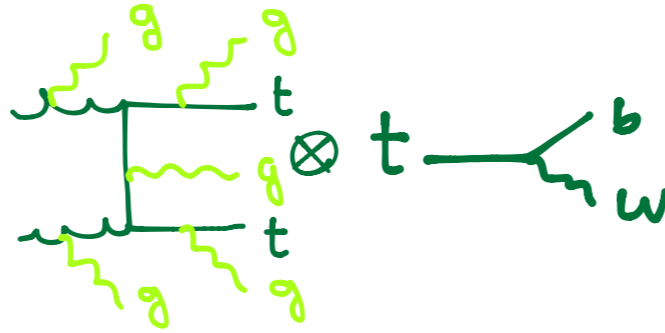
- $t\bar{t}$  vs.  $b\bar{W}bW$
- shower effects (NLO+PS Powheg)
- non-perturbative effects (color re-connection)



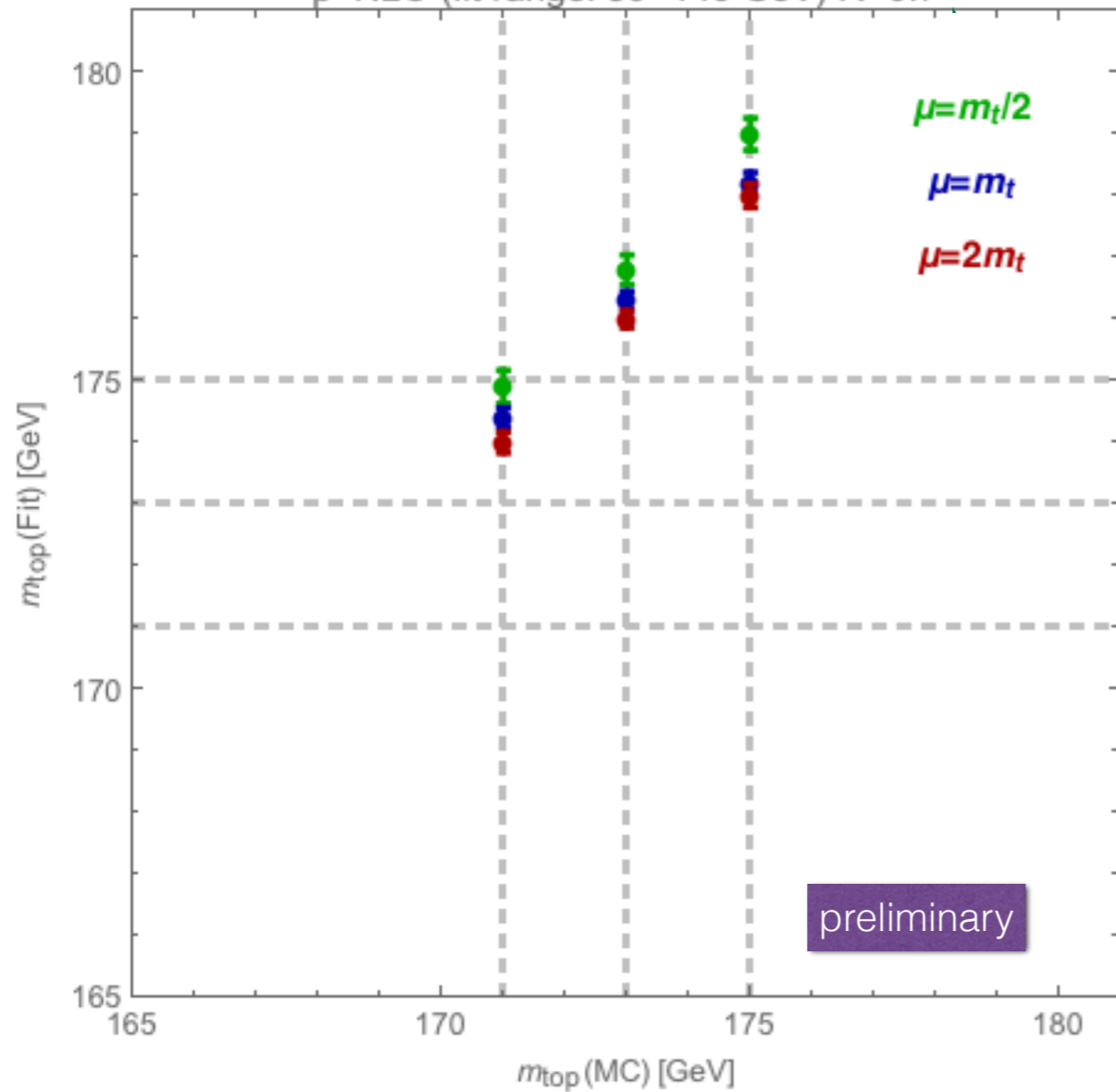
Extra

# NLO: production

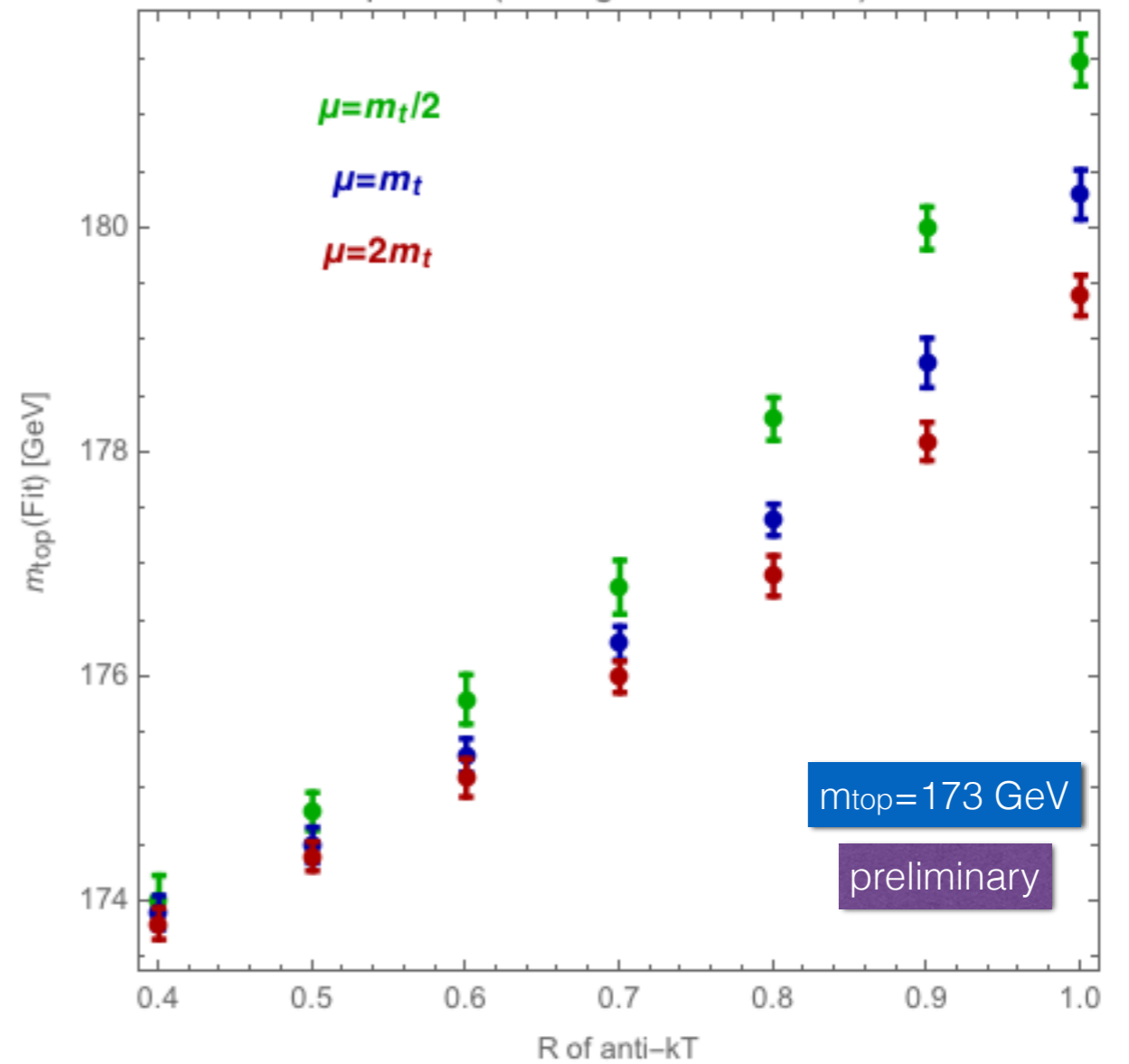
(MCFM)



p-NLO (fit range: 30–140 GeV)  $R=0.7$



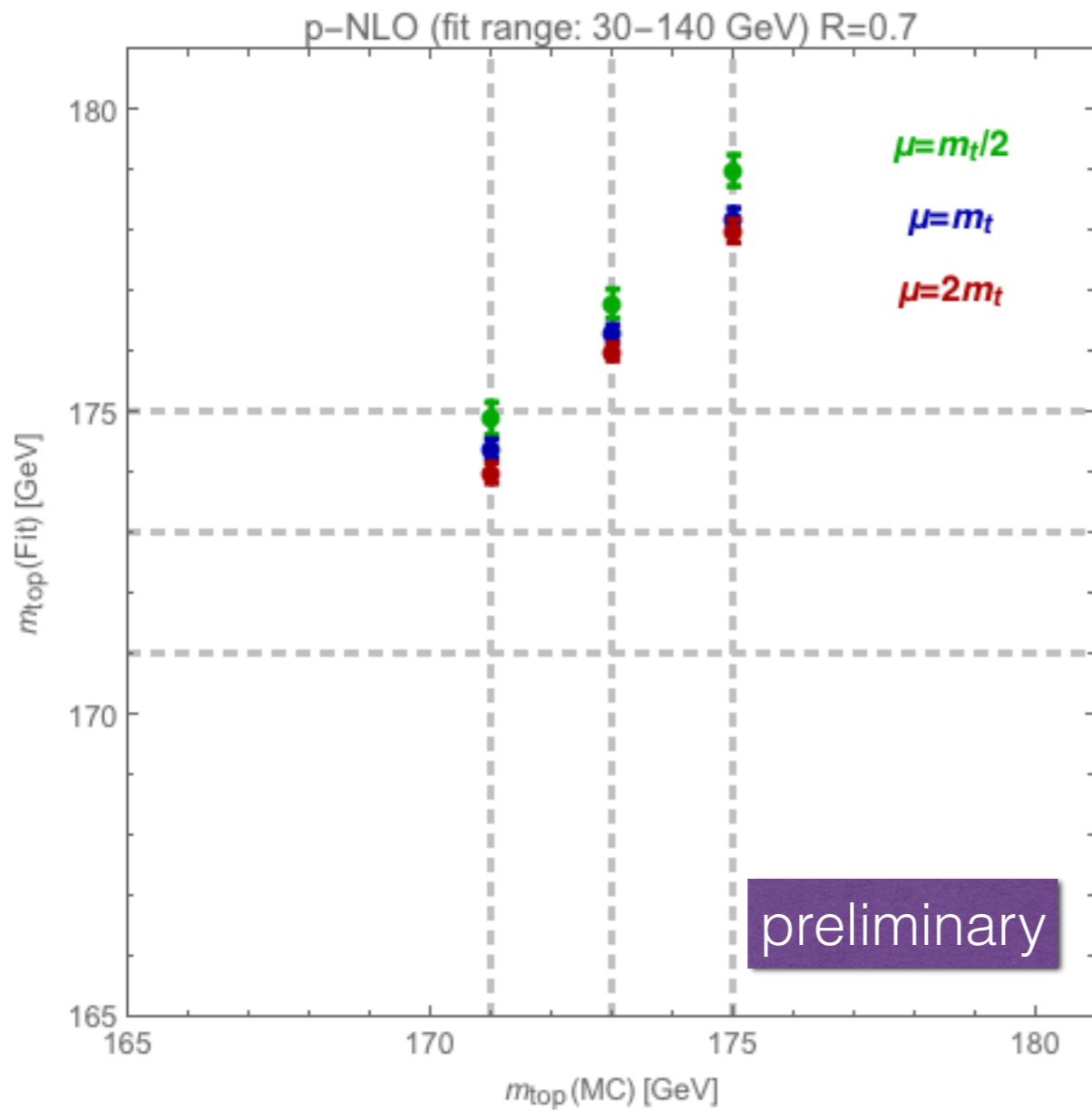
p-NLO (fit range: 30–140 GeV)



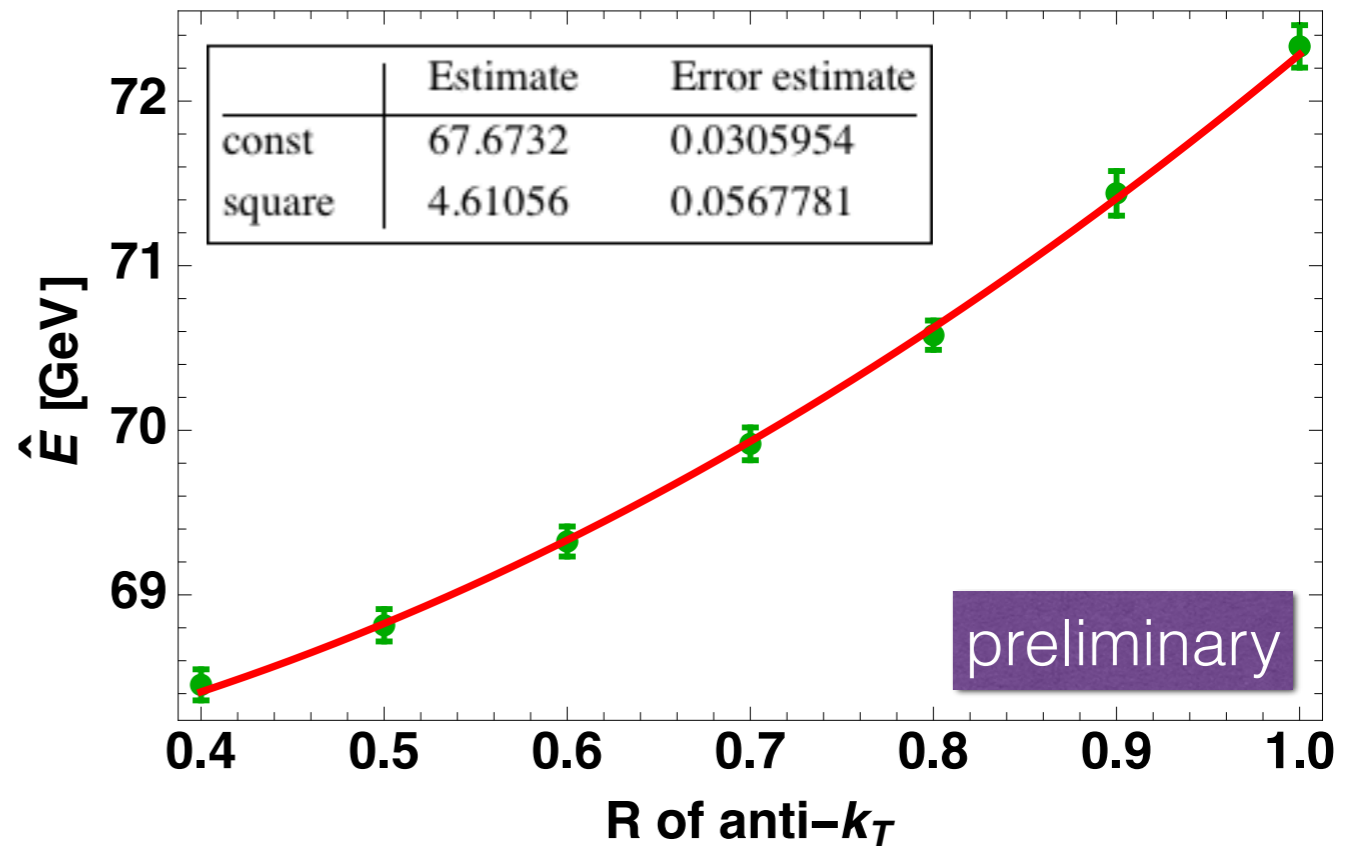
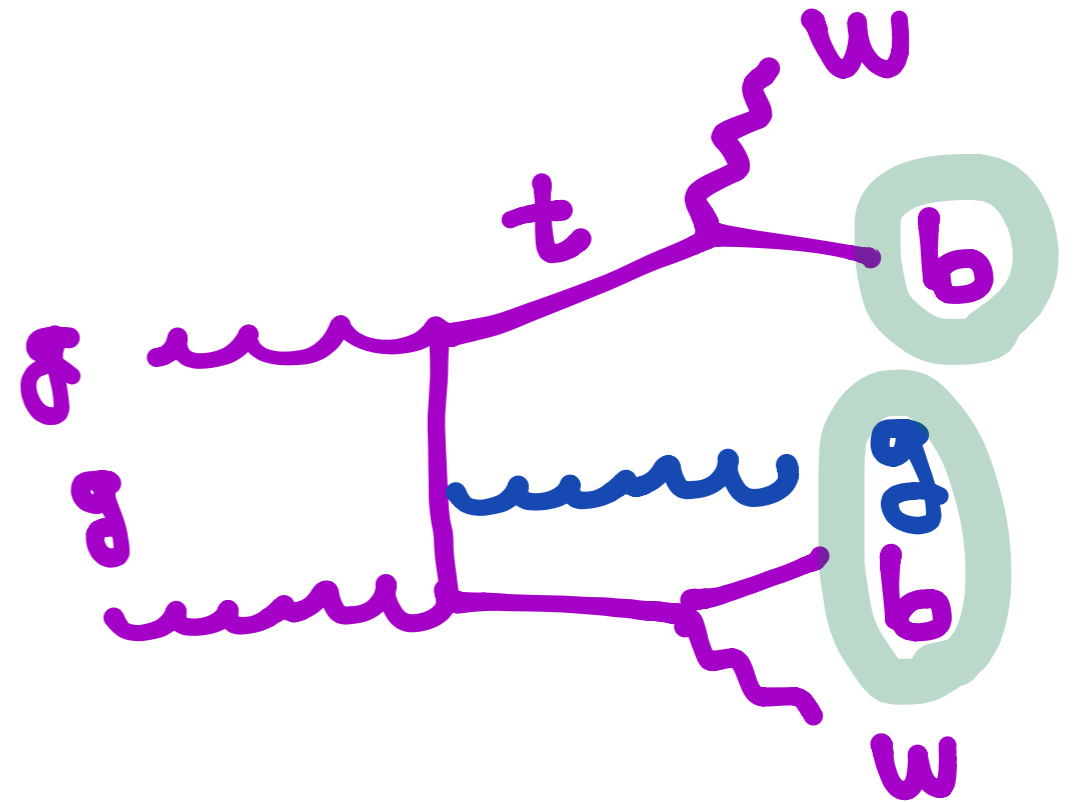
very little sensitive to the scale choice (less than 400 MeV on  $m_{\text{top}}$ )

# NLO: production

(MCFM)

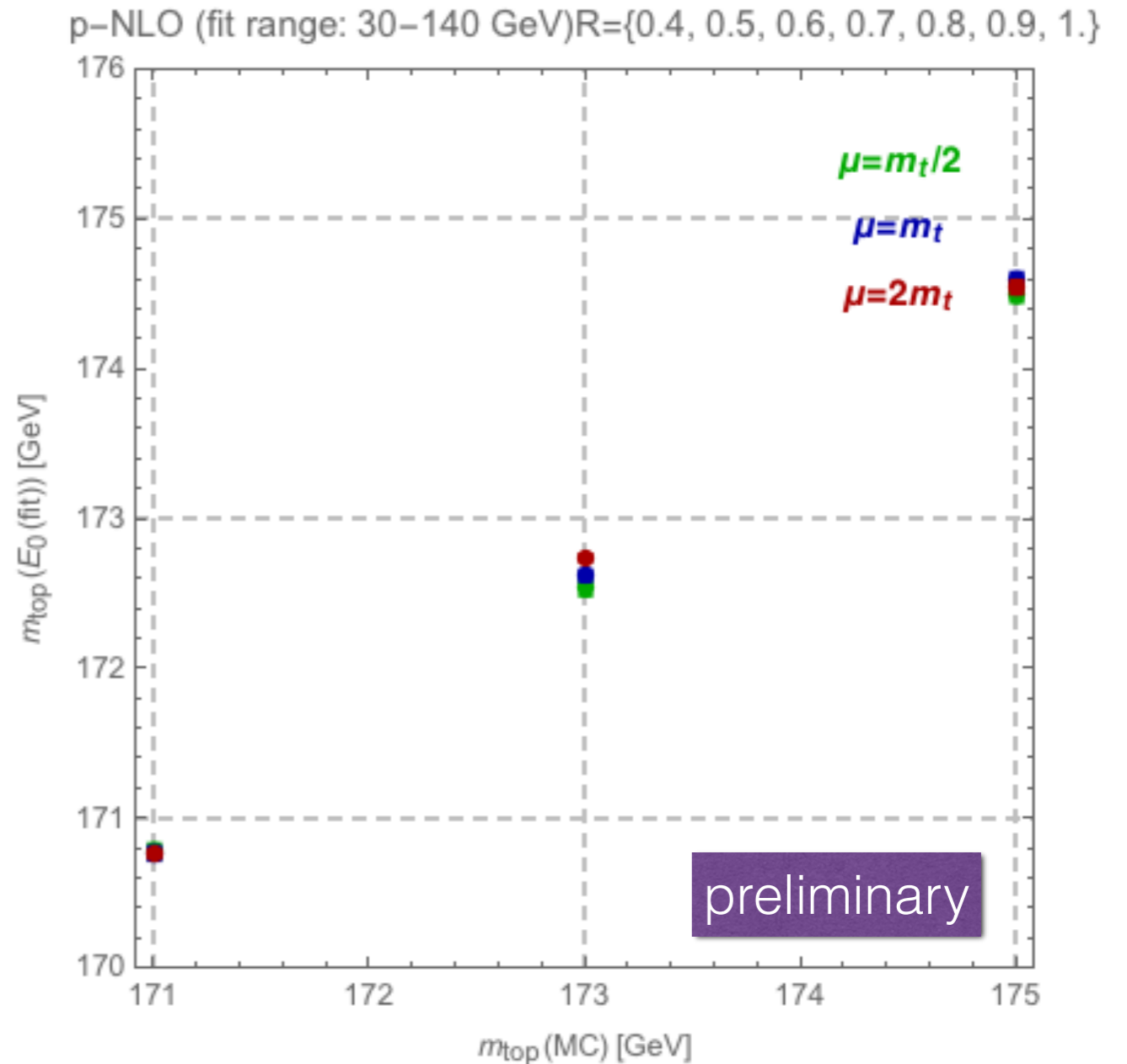
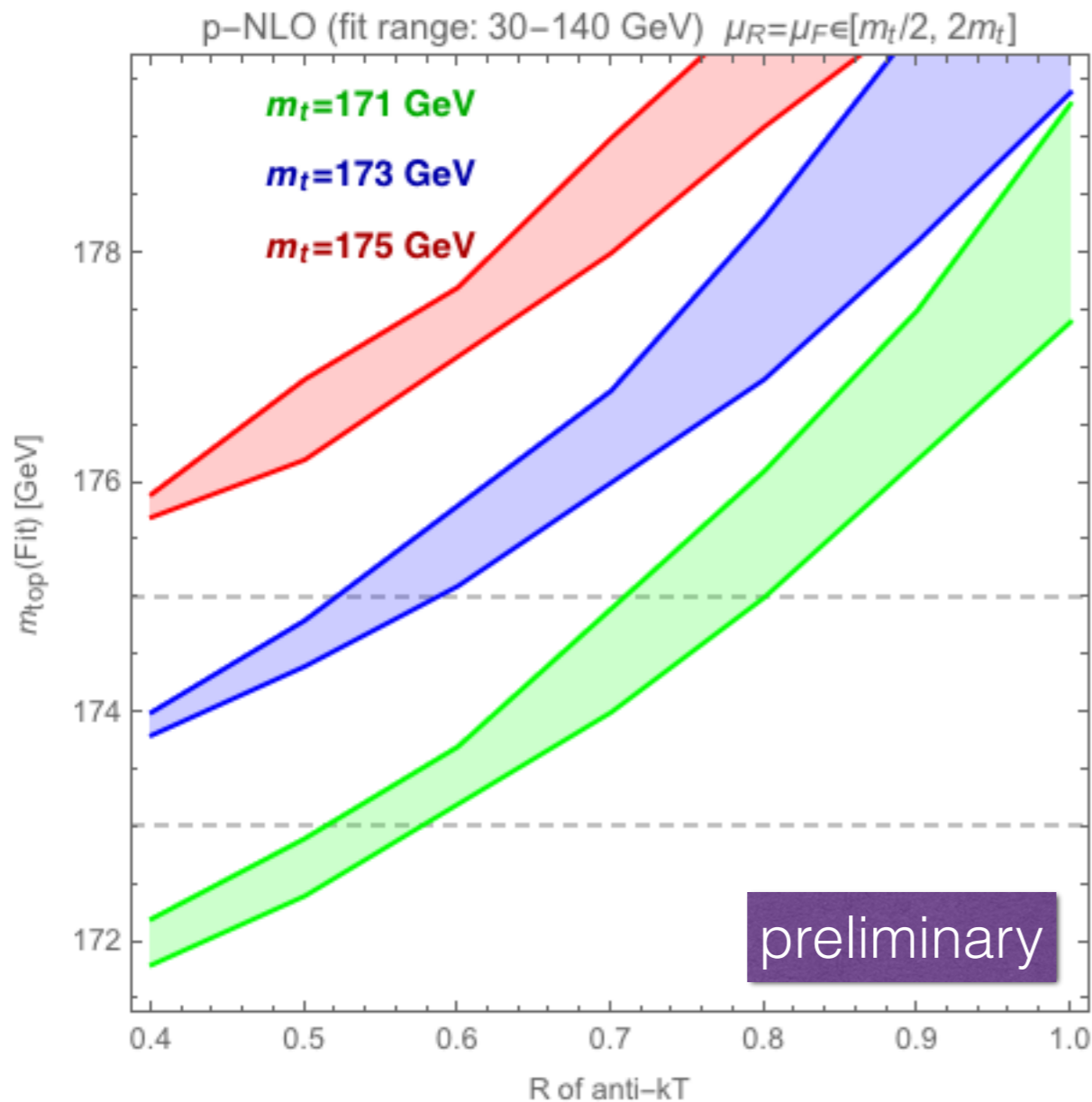


shift  $\sim R^p$  (p~2 jet area)  
 shift  $\sim 1/\mu$  (real radiation)



# NLO: production

(MCFM)

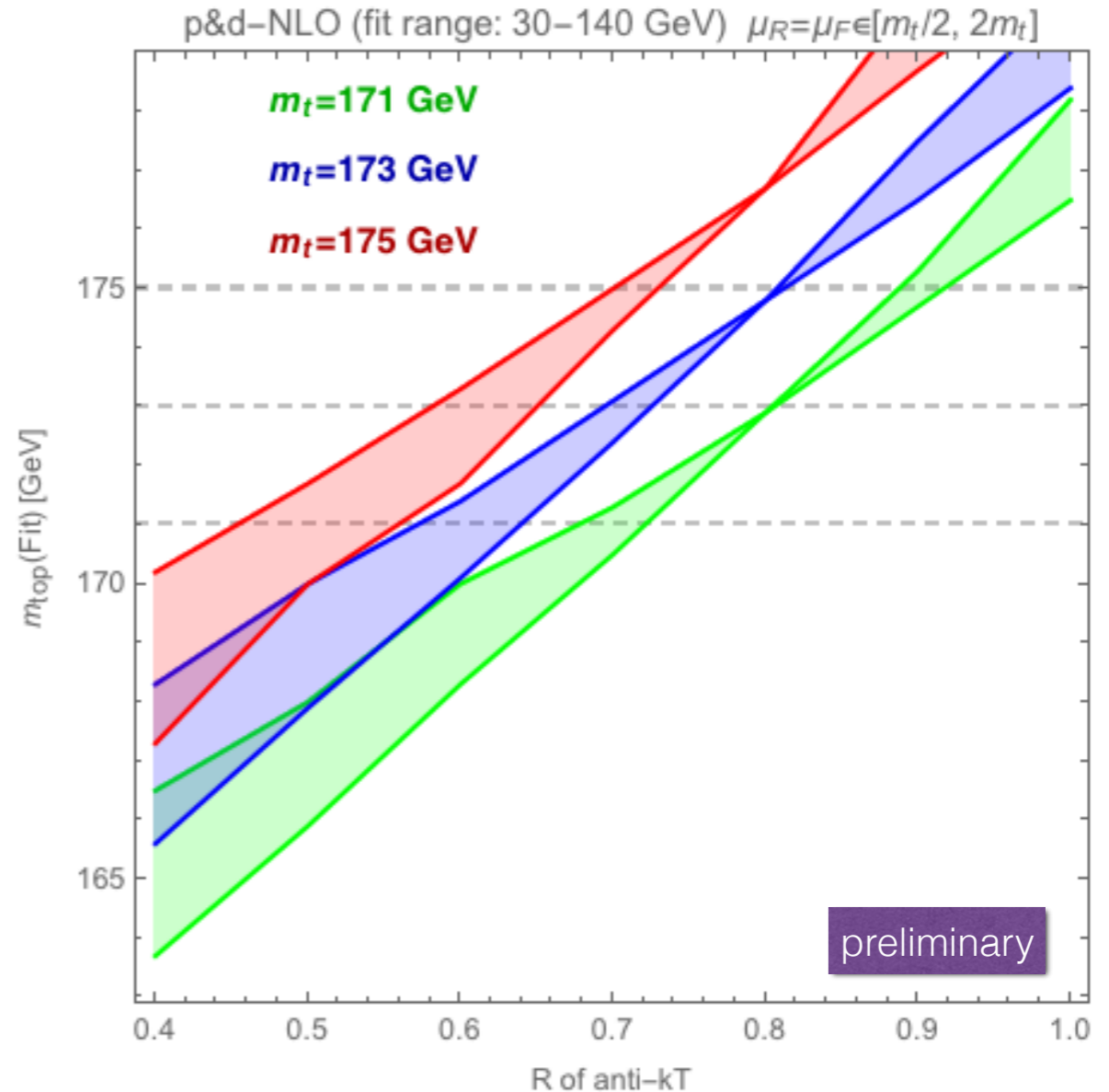
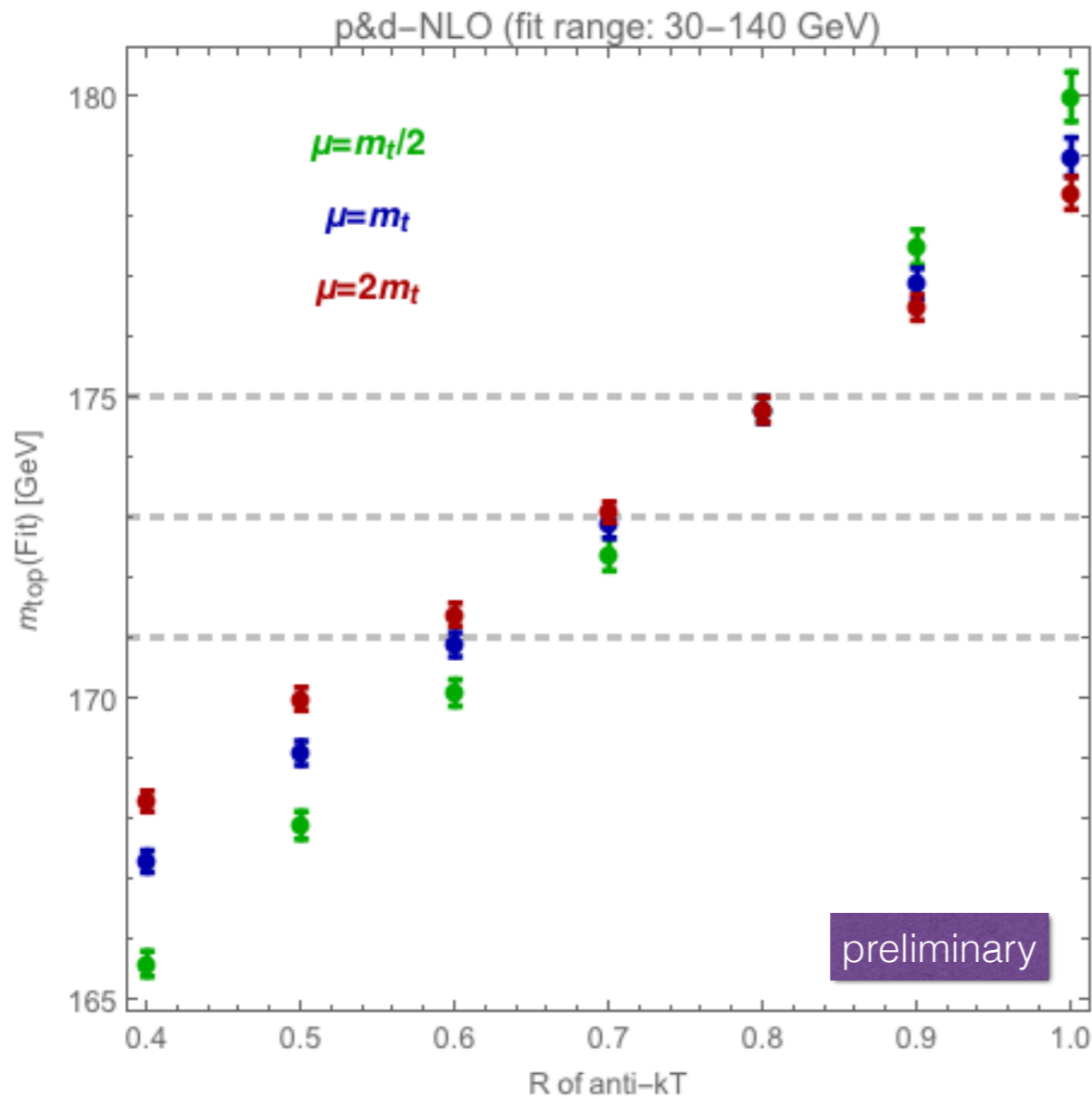
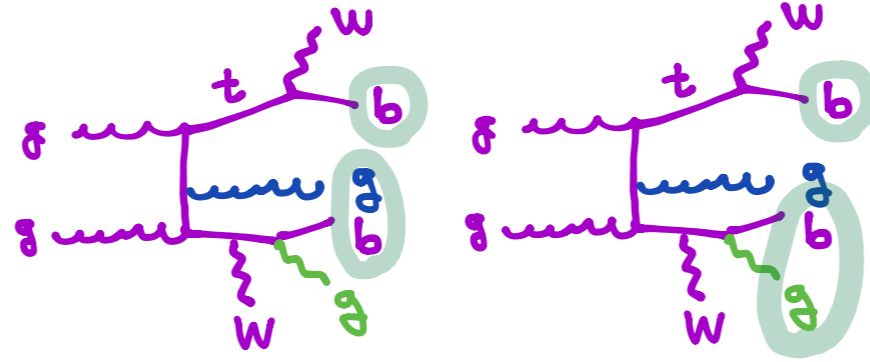


shift  $\sim R^p$  ( $p \sim 2$  jet area)  
 shift  $\sim 1/\mu$  (real radiation)

$$\hat{E} = E_0 + \alpha(\mu) \cdot p \cdot R^2 + \dots$$

$E_b \times = \frac{m_t^2 - m_w^2 + m_b^2}{2m_t}$

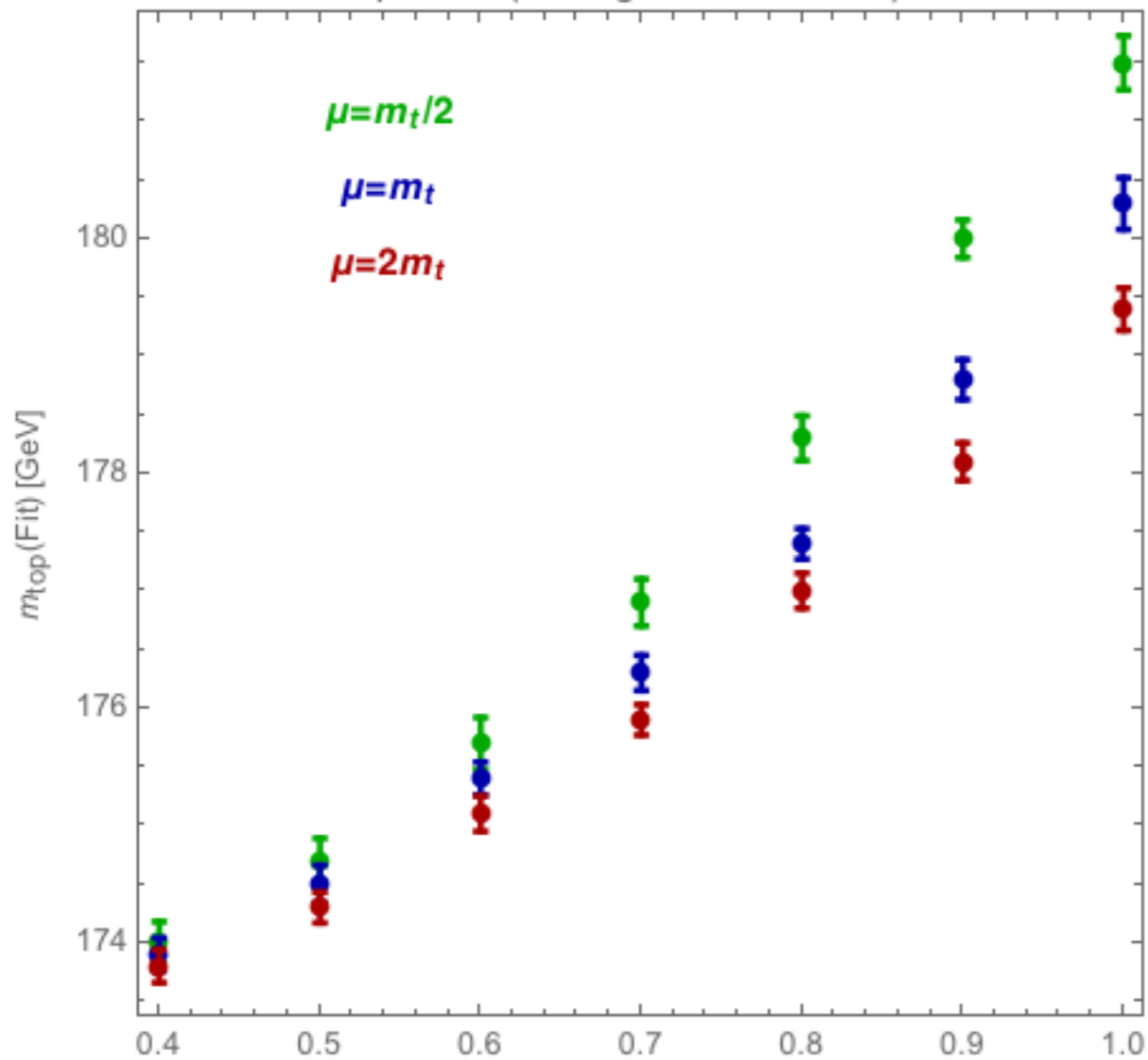
# NLO: production & decay



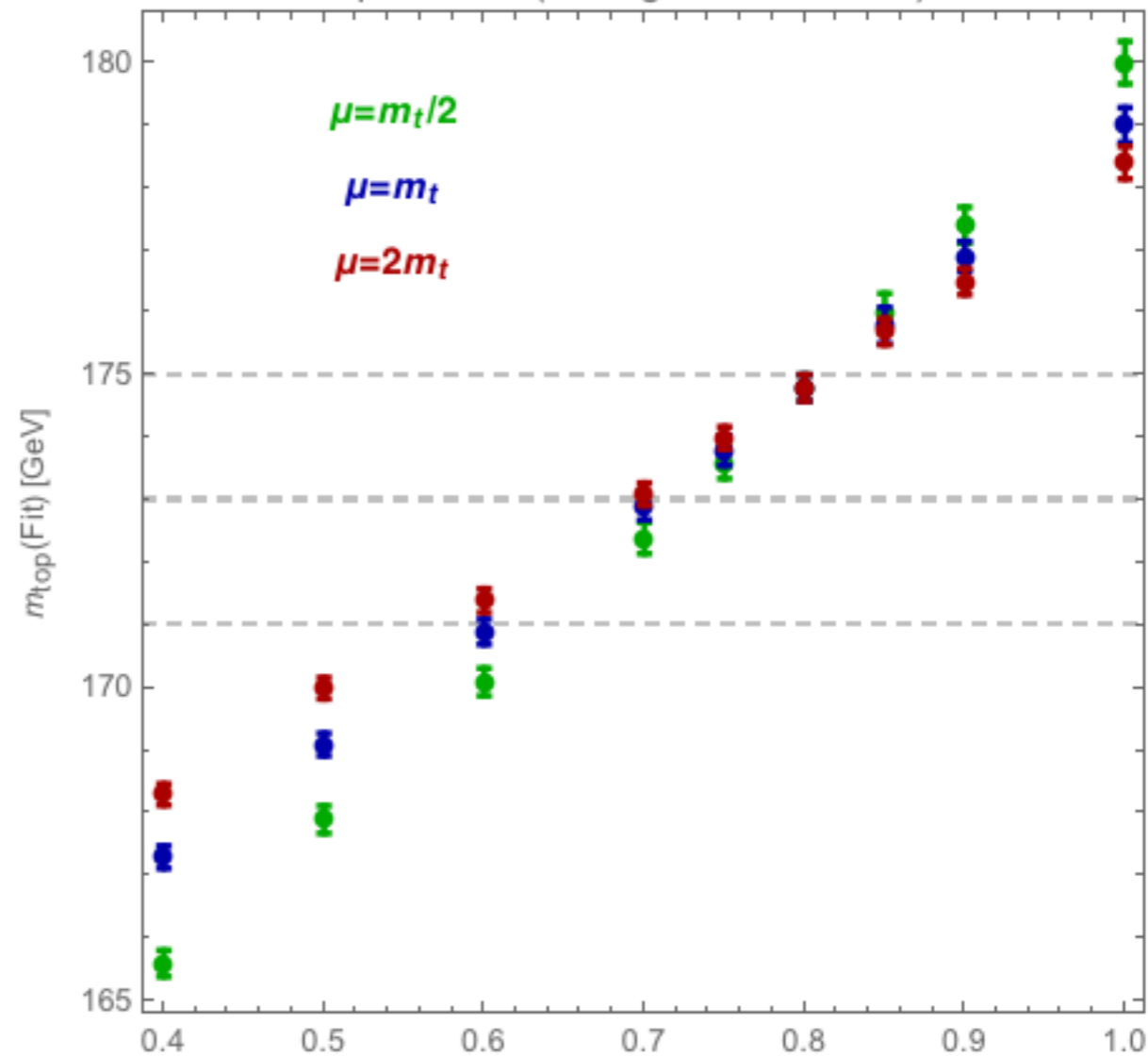
$$\hat{E} = E_0 + \alpha(\mu) \cdot [ p \cdot R^2 + p \cdot \log R ] + \dots$$

decay NLO sensitive to the scale choice:  $\pm 1$  GeV on  $m_{\text{top}}$

p-NLO (fit range: 30–140 GeV)

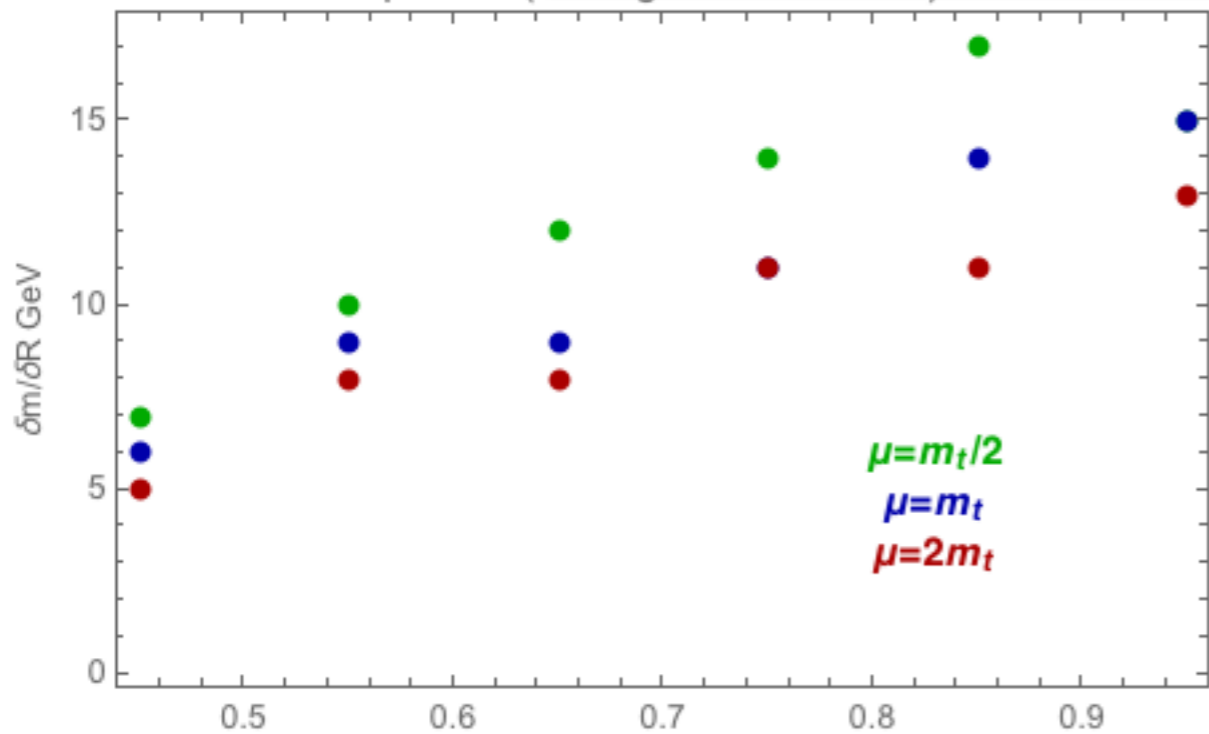


p&amp;d-NLO (fit range: 30–140 GeV)

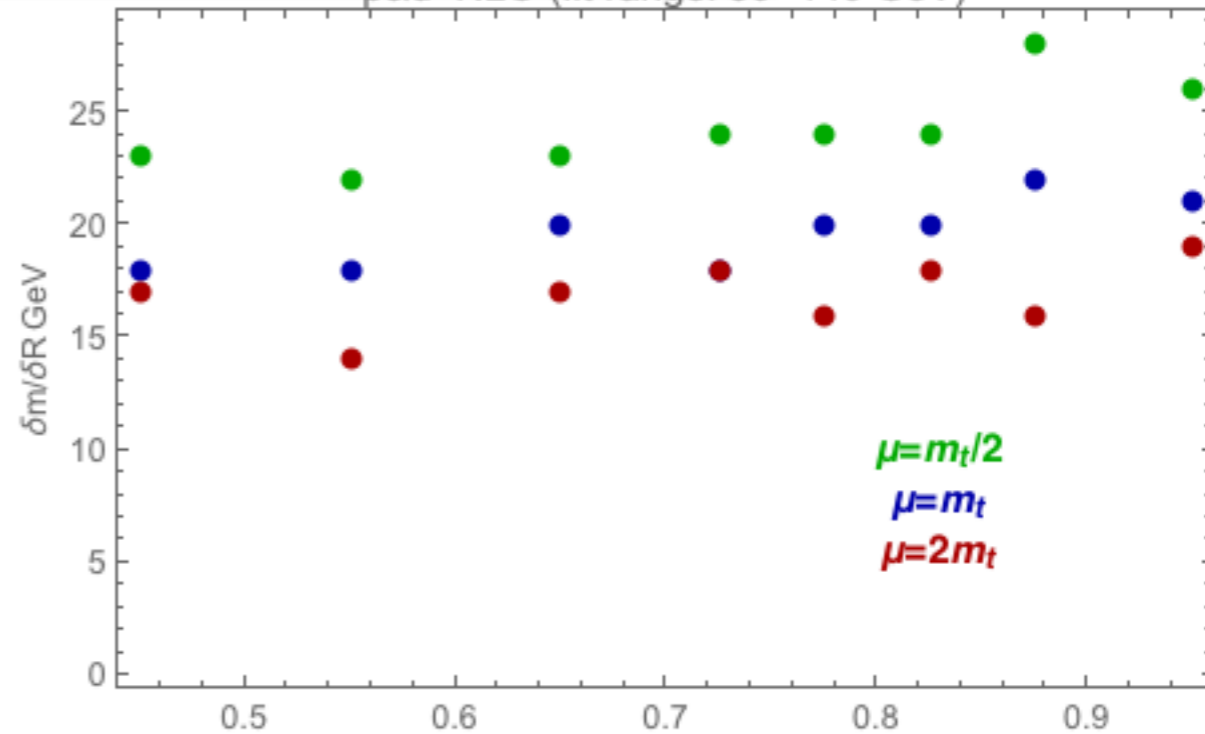


$$dE/dR = 2pR + d/R \quad d \rightarrow 2p \quad dE/dR = 4p + 2p(1-R)^2 + \dots$$

p-NLO (fit range: 30–140 GeV)



p&amp;d-NLO (fit range: 30–140 GeV)



# NLO virtues

*Agashe, Franceschini, Kim, Schulze - in preparation*

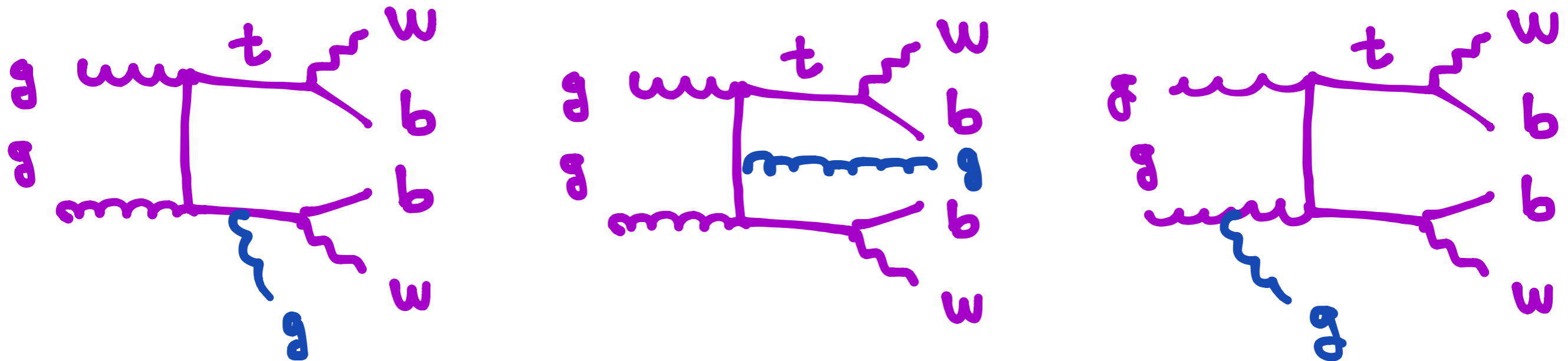
- **Invariance holds for  $pp \rightarrow tt$  @ NLO**
- Not sensitive to Initial State Radiation
- Not sensitive to Parton Distribution Functions
- Not sensitive to the exact energy of the collider

**only sensitive to the NLO decay  $t \rightarrow bWg$**

# Insenstive to production at NLO

Agashe, Franceschini, Kim, Schulze - in preparation

Production NLO only affects the boost distribution of top



**The energy peak position is unchanged**

$$E_b^{\text{peak}} = \frac{m_t^2 - m_W^2 + m_{b/j}^2}{2m_t} = E_b^*$$

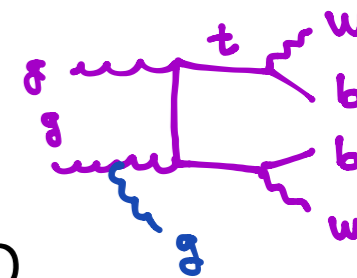


# NLO virtues

- Invariance holds for  $pp \rightarrow tt$  @ NLO
- **Not sensitive to Initial State Radiation**
- Not sensitive to Parton Distribution Functions
- Not sensitive to the exact energy of the collider

**only sensitive to the NLO decay  $t \rightarrow bWg$**

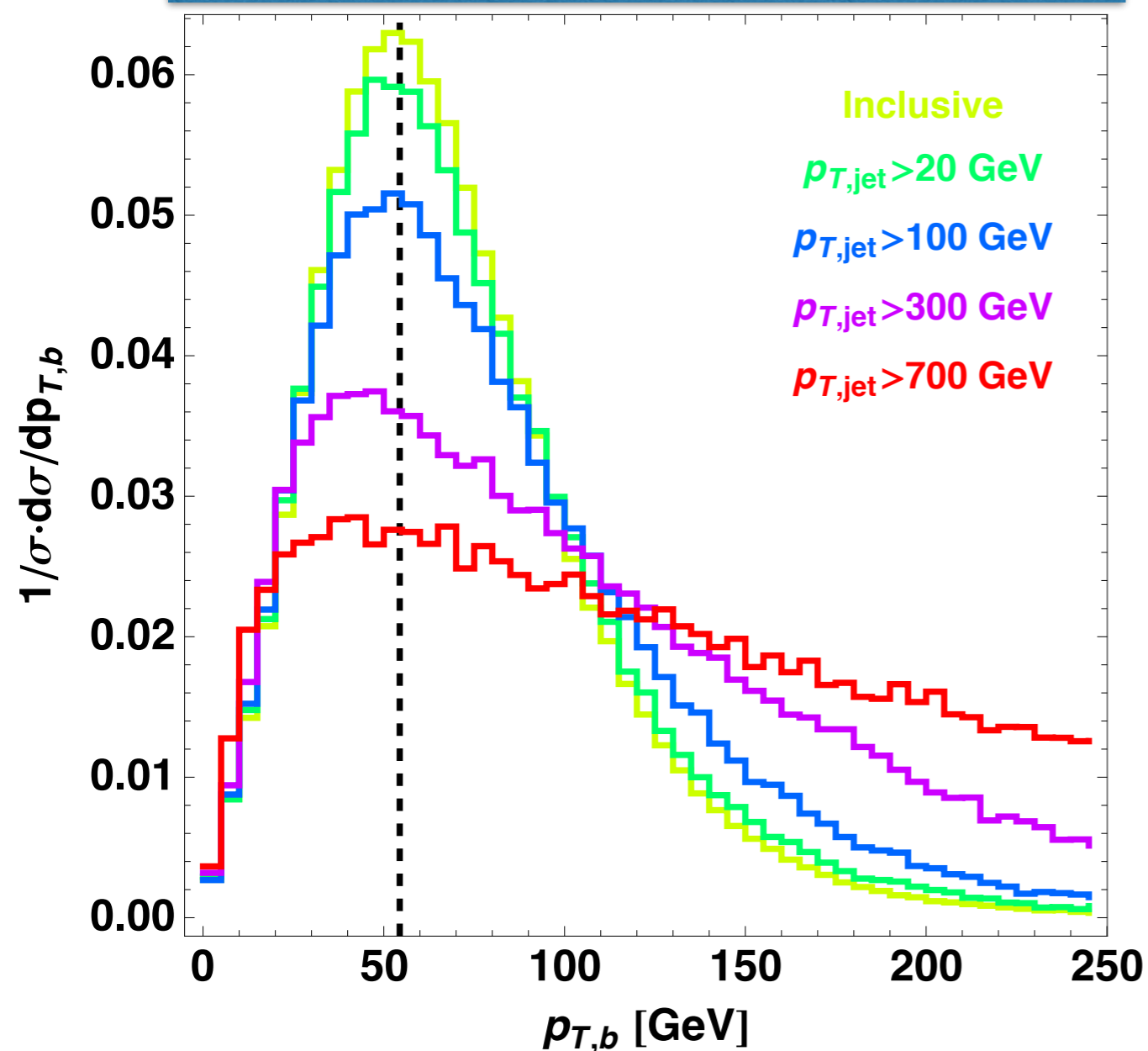
# Effect of initial state radiation



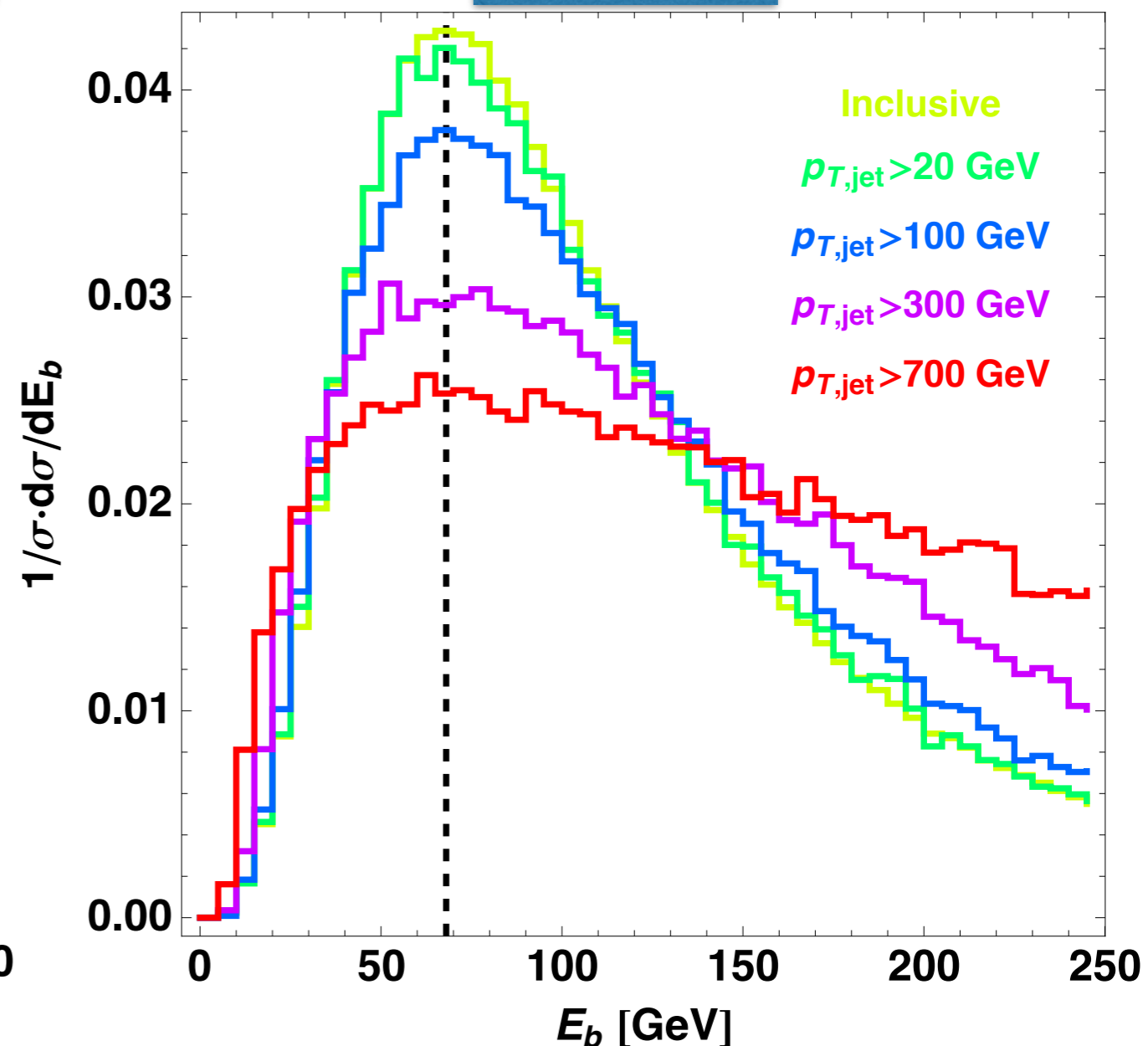
ISR only affects the boost distribution of top

*Agashe, Franceschini, Kim, Schulze - in preparation*

## Transverse Momentum



## Energy



peak stability →

# NLO virtues

- Invariance holds for  $pp \rightarrow tt$  @ NLO
- Not sensitive to Initial State Radiation
- **Not sensitive to Parton Distribution Functions**
- **Not sensitive to the exact energy of the collider**

**only sensitive to the NLO decay  $t \rightarrow bWg$**

# Top mass combination

1403.4427 - First combination of Tevatron and LHC measurements of the top-quark mass

## LHC/Tevatron NOTE

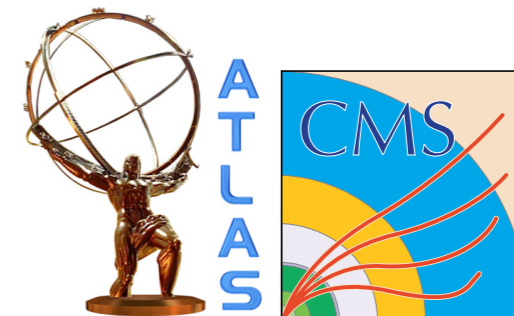
ATLAS-CONF-2014-008

CDF Note 11071

CMS PAS TOP-13-014

D0 Note 6416

March 17, 2014



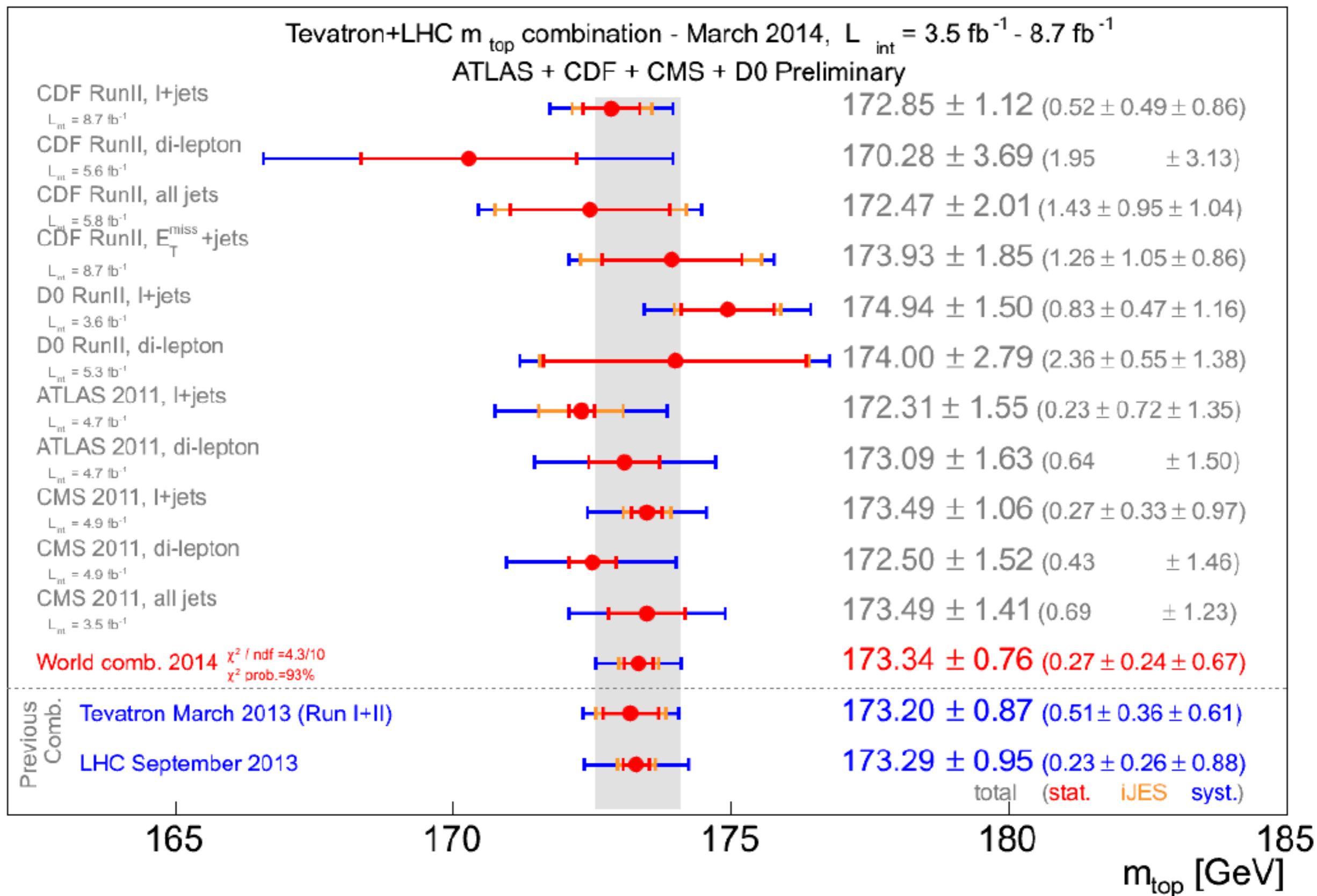
Experiment	$t\bar{t}$ final state	$\mathcal{L}_{int}$ [ $\text{fb}^{-1}$ ]	$m_{top} \pm (\text{stat.}) \pm (\text{syst.})$ [GeV]	Total uncertainty on $m_{top}$ [GeV] ([%])	Reference
CDF	$l+jets$	8.7	$172.85 \pm 0.52 \pm 0.99$	<u>1.12</u> (0.65)	[8]
	dilepton	5.6	$170.28 \pm 1.95 \pm 3.13$	3.69 (2.17)	[9]
	all jets	5.8	$172.47 \pm 1.43 \pm 1.41$	2.01 (1.16)	[10]
	$E_T^{miss}+jets$	8.7	$173.93 \pm 1.26 \pm 1.36$	1.85 (1.07)	[11]
D0	$l+jets$	3.6	$174.94 \pm 0.83 \pm 1.25$	1.50 (0.86)	[12]
	dilepton	5.3	$174.00 \pm 2.36 \pm 1.49$	2.79 (1.60)	[13]
ATLAS	$l+jets$	4.7	$172.31 \pm 0.23 \pm 1.53$	1.55 (0.90)	[14]
	dilepton	4.7	$173.09 \pm 0.64 \pm 1.50$	1.63 (0.94)	[15]
CMS	$l+jets$	4.9	$173.49 \pm 0.27 \pm 1.03$	<u>1.06</u> (0.61)	[16]
	dilepton	4.9	$172.50 \pm 0.43 \pm 1.46$	1.52 (0.88)	[17]
	all jets	3.5	$173.49 \pm 0.69 \pm 1.23$	1.41 (0.81)	[18]

**LHC-7 is on par with Tevatron**

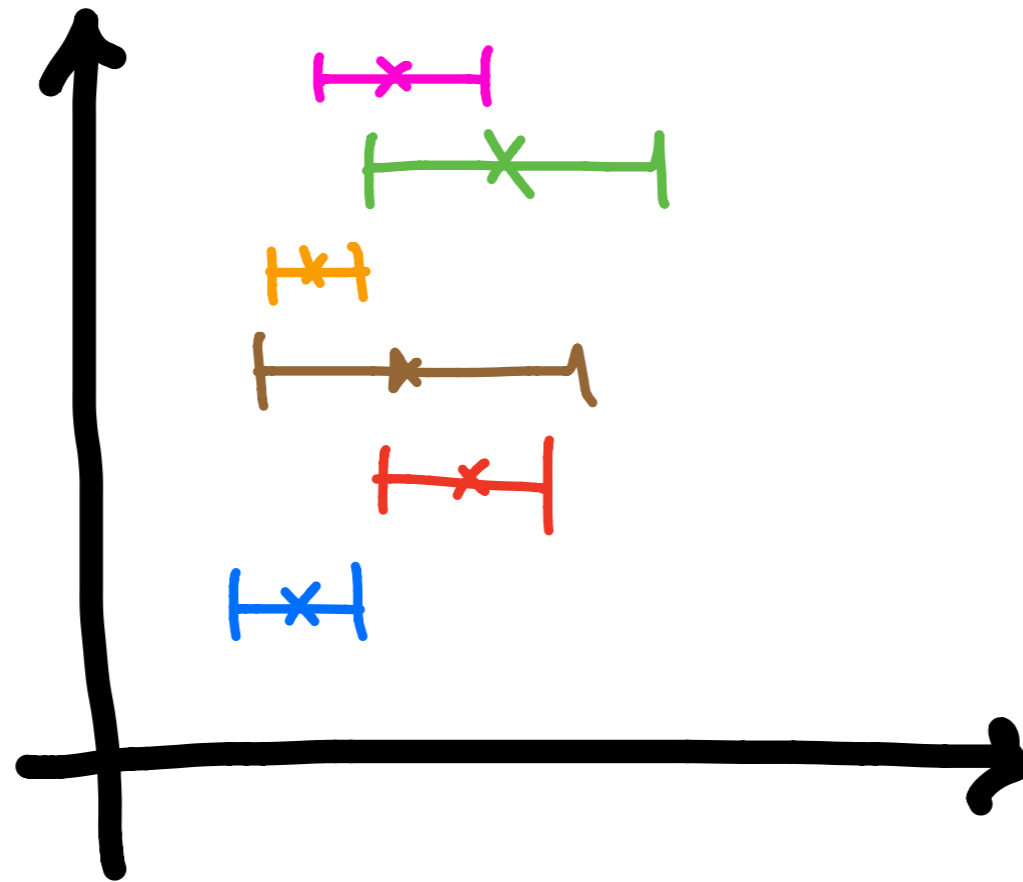
**$173.34 \pm 0.27(\text{stat}) \pm 0.71(\text{syst})$  GeV  
dominated by systematics**

$l+jets$   
dilepton  
all jets

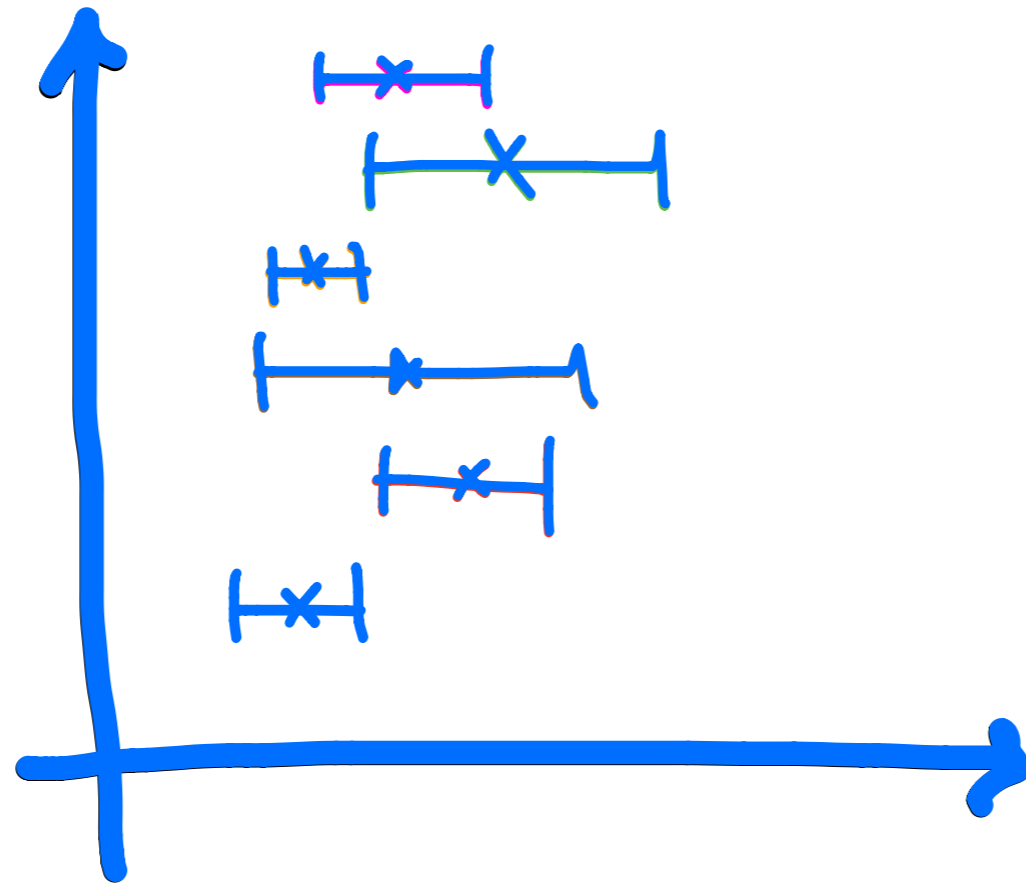
# Many measurements



# Many measurements?



# Many measurements?

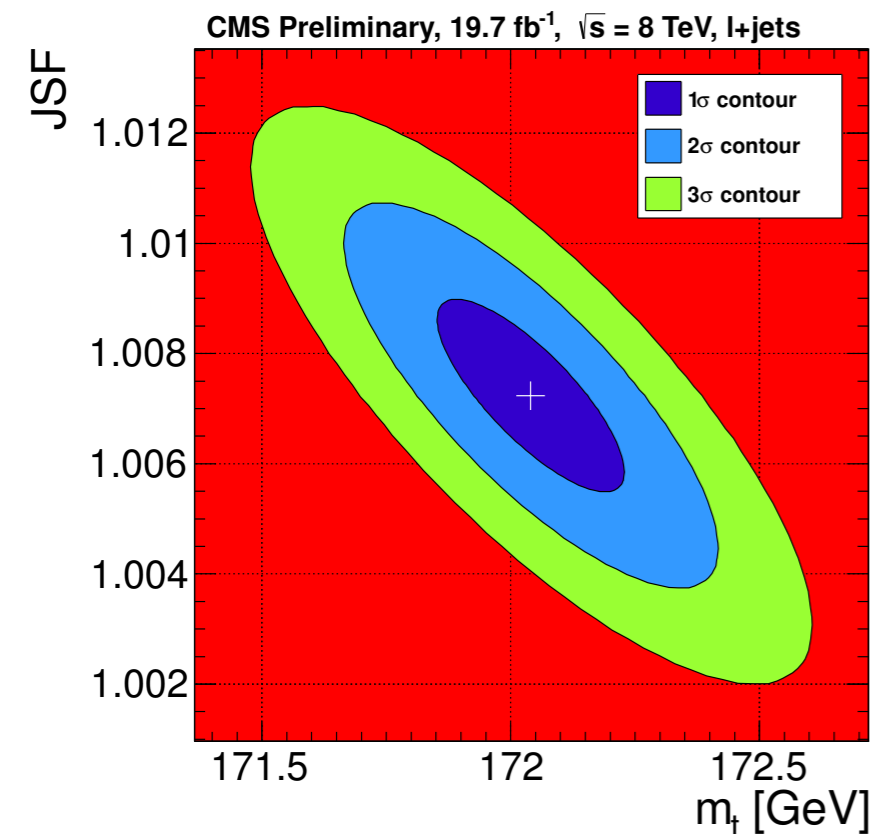


# CMS PAS TOP-14-001

$172.04 \pm 0.19$  (stat.+JSF)  $\pm 0.75$  (syst.) GeV

## Ideogram Method (Kinematic fit)

**MG5+Py6 or POWHEG**



	$\delta m_t^{2D}$ (GeV)	$\delta$ JSF	$\delta m_t^{1D}$ (GeV)
<b>Experimental uncertainties</b>			
Fit calibration	0.10	0.001	0.06
$p_T$ - and $\eta$ -dependent JES	0.18	0.007	1.17
Lepton energy scale	0.03	<0.001	0.03
MET	0.09	0.001	0.01
Jet energy resolution	0.26	0.004	0.07
b tagging	0.02	<0.001	0.01
Pileup	0.27	0.005	0.17
Non- $t\bar{t}$ background	0.11	0.001	0.01
<b>Modeling of hadronization</b>			
Flavor-dependent JSF	0.41	0.004	0.32
b fragmentation	0.06	0.001	0.04
Semi-leptonic B hadron decays	0.16	<0.001	0.15
<b>Modeling of the hard scattering process</b>			
PDF	0.09	0.001	0.05
Renormalization and factorization scales	$0.12 \pm 0.13$	$0.004 \pm 0.001$	$0.25 \pm 0.08$
ME-PS matching threshold	$0.15 \pm 0.13$	$0.003 \pm 0.001$	$0.07 \pm 0.08$
ME generator	$0.23 \pm 0.14$	$0.003 \pm 0.001$	$0.20 \pm 0.08$
<b>Modeling of non-perturbative QCD</b>			
Underlying event	$0.14 \pm 0.17$	$0.002 \pm 0.002$	$0.06 \pm 0.10$
Color reconnection modeling	$0.08 \pm 0.15$	$0.002 \pm 0.001$	$0.07 \pm 0.09$
<b>Total</b>	<b>0.75</b>	<b>0.012</b>	<b>1.29</b>



# ATLAS-CONF-2013-046

$$m_{\text{top}} = 172.31 \pm 0.23 \text{ (stat)} \pm 0.27 \text{ (JSF)} \pm 0.67 \text{ (bJSF)} \pm 1.35 \text{ (syst)} \text{ GeV}$$

## 3D Method (Kinematic Fit)

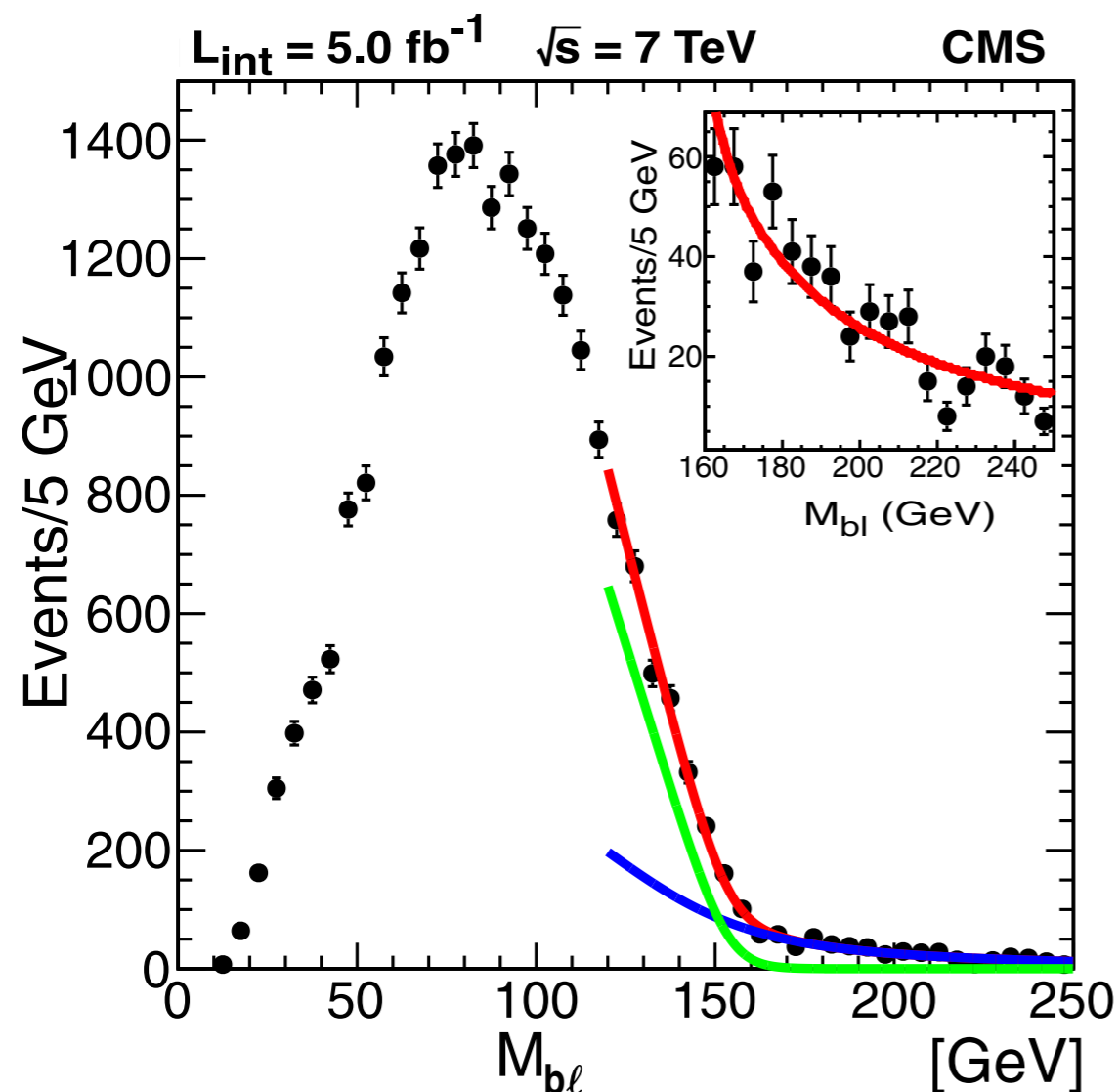
MC@NLO or POWHEG	2d-analysis		3d-analysis		
	$m_{\text{top}}$ [GeV]	JSF	$m_{\text{top}}$ [GeV]	JSF	bJSF
Measured value	172.80	1.014	172.31	1.014	1.006
Data statistics	0.23	0.003	0.23	0.003	0.008
Jet energy scale factor (stat. comp.)	0.27	n/a	0.27	n/a	n/a
bJet energy scale factor (stat. comp.)	n/a	n/a	0.67	n/a	n/a
Method calibration	0.13	0.002	0.13	0.002	0.003
Signal MC generator	0.36	0.005	0.19	0.005	0.002
Hadronisation	1.30	0.008	0.27	0.008	0.013
Underlying event	0.02	0.001	0.12	0.001	0.002
Colour reconnection	0.03	0.001	0.32	0.001	0.004
ISR and FSR (signal only)	0.96	0.017	0.45	0.017	0.006
Proton PDF	0.09	0.000	0.17	0.000	0.001
single top normalisation	0.00	0.000	0.00	0.000	0.000
$W$ +jets background	0.02	0.000	0.03	0.000	0.000
QCD multijet background	0.04	0.000	0.10	0.000	0.001
Jet energy scale	0.60	0.005	0.79	0.004	0.007
$b$ -jet energy scale	0.92	0.000	0.08	0.000	0.002
Jet energy resolution	0.22	0.006	0.22	0.006	0.000
Jet reconstruction efficiency	0.03	0.000	0.05	0.000	0.000
$b$ -tagging efficiency and mistag rate	0.17	0.001	0.81	0.001	0.011
Lepton energy scale	0.03	0.000	0.04	0.000	0.000
Missing transverse momentum	0.01	0.000	0.03	0.000	0.000
Pile-up	0.03	0.000	0.03	0.000	0.001
Total systematic uncertainty	2.02	0.021	1.35	0.021	0.020
Total uncertainty	2.05	0.021	1.55	0.021	0.022

# CMS-TOP-11-027

$$M_t = 173.9 \pm 0.9 \text{ (stat.)}_{-2.1}^{+1.7} \text{ (syst.) GeV}$$

$m(b,l)$  end-point

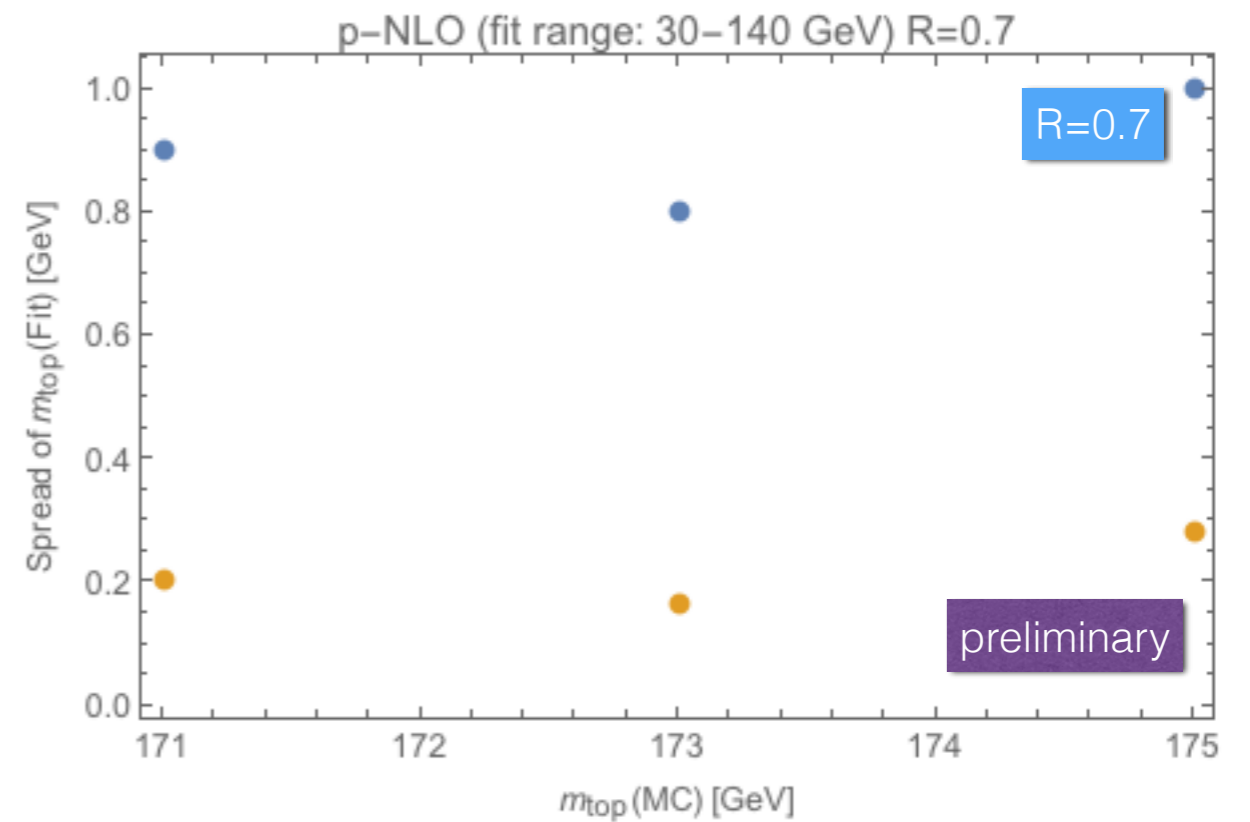
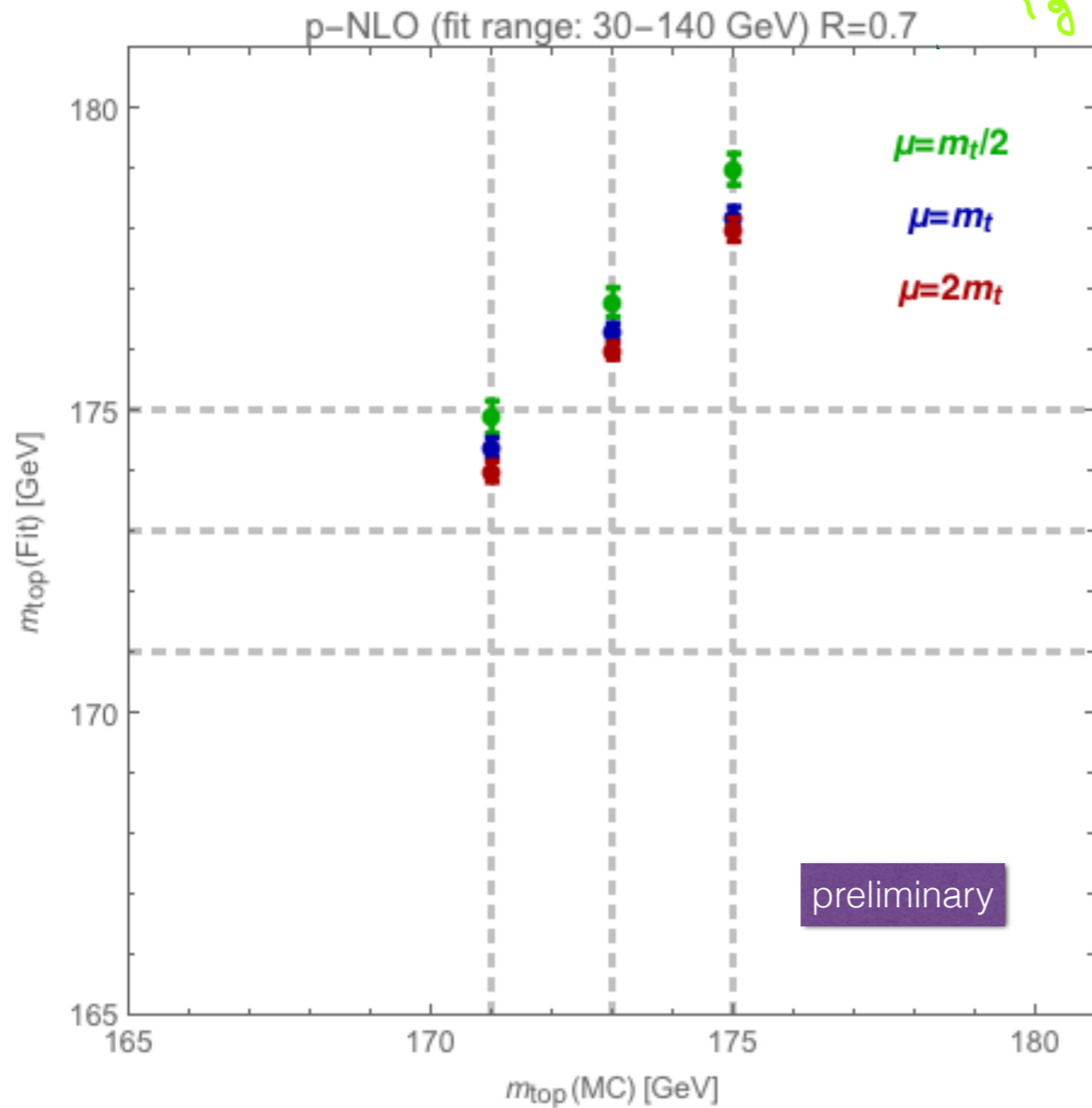
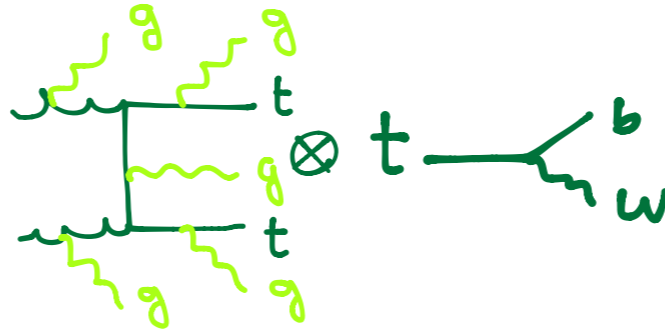
- robust to NLO
- robust to combinatorics
- robust to hadronization



Source	$\delta M_t$ (GeV)
Jet Energy Scale	+1.3 -1.8
Jet Energy Resolution	$\pm 0.5$
Lepton Energy Scale	+0.3 -0.4
Fit Range	$\pm 0.6$
Background Shape	$\pm 0.5$
Jet and Lepton Efficiencies	+0.1 -0.2
Pileup	$< 0.1$
QCD effects	$\pm 0.6$
Total	+1.7 -2.1

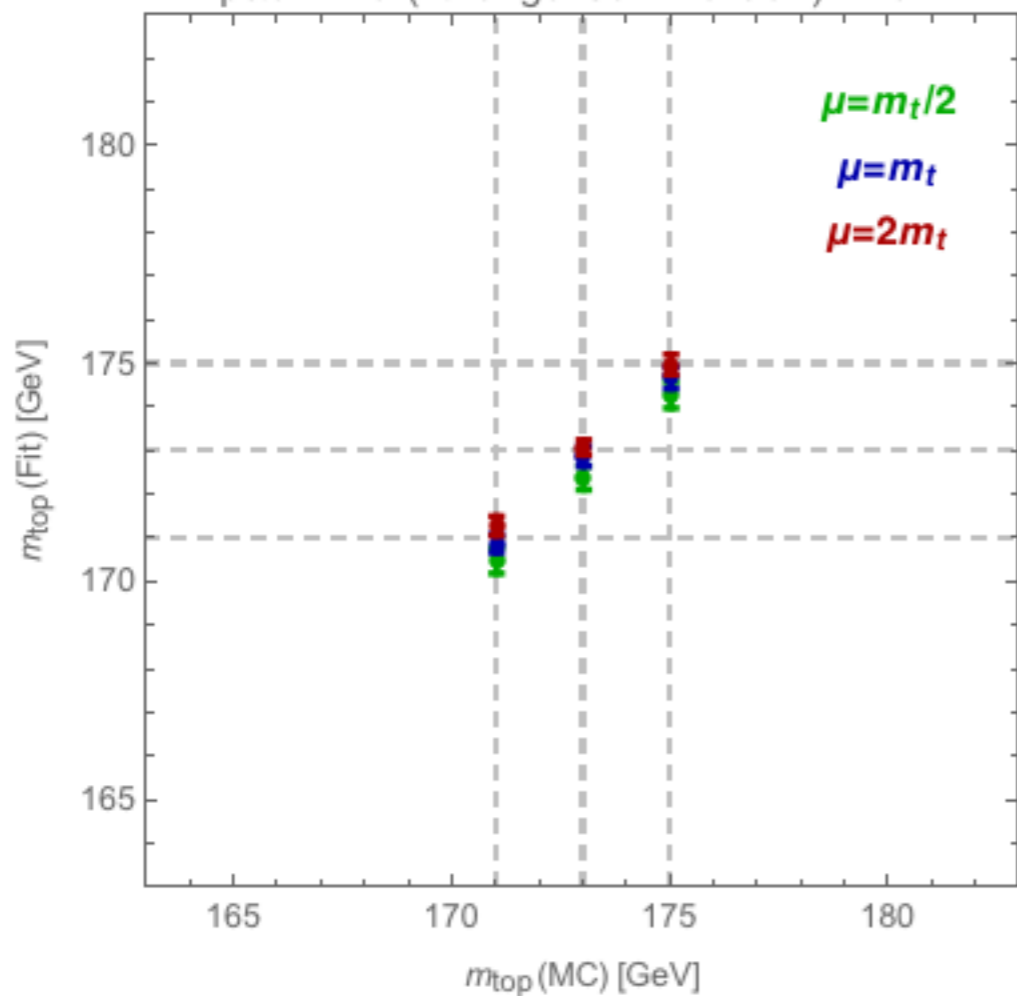
# NLO: production

(MCFM)

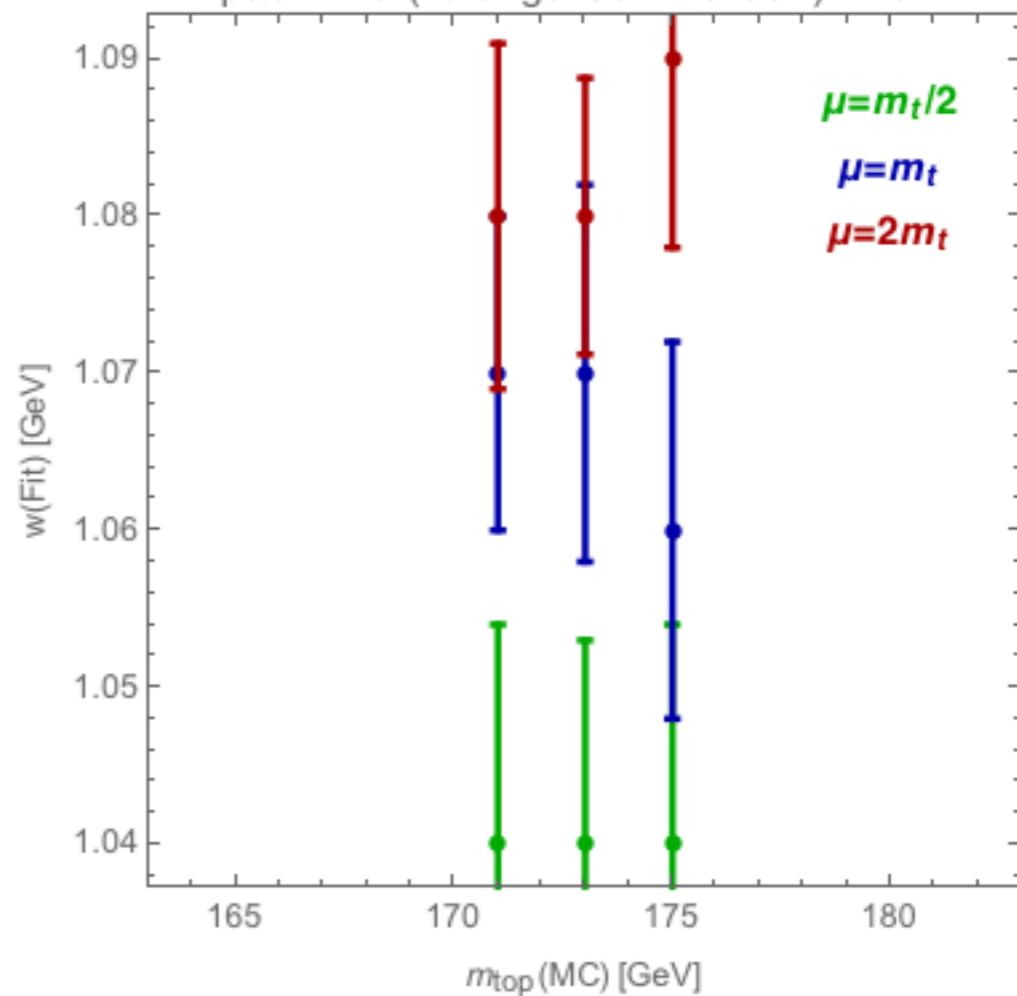


very little sensitive to the scale choice (less than 400 MeV on  $m_{\text{top}}$ )

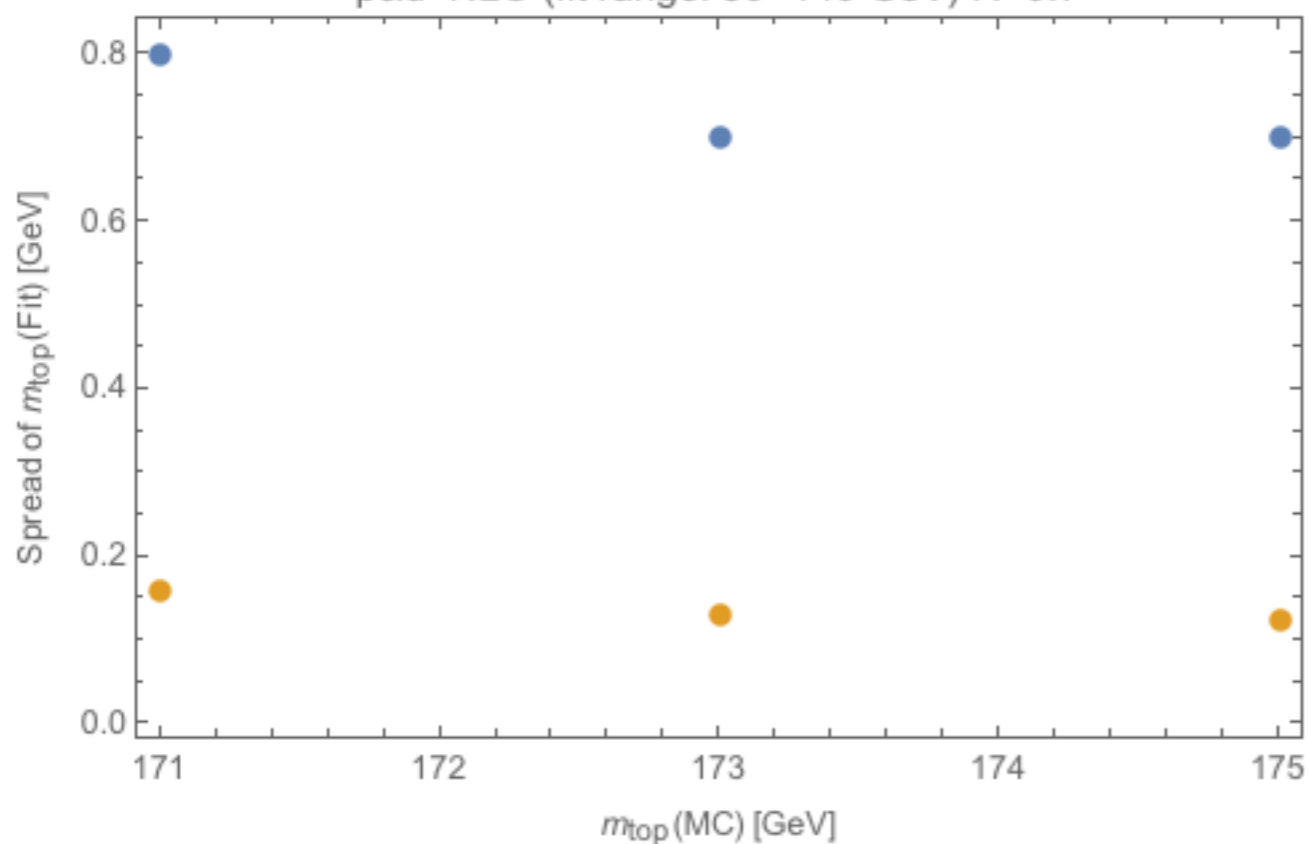
p&amp;d-NLO (fit range: 30–140 GeV) R=0.7



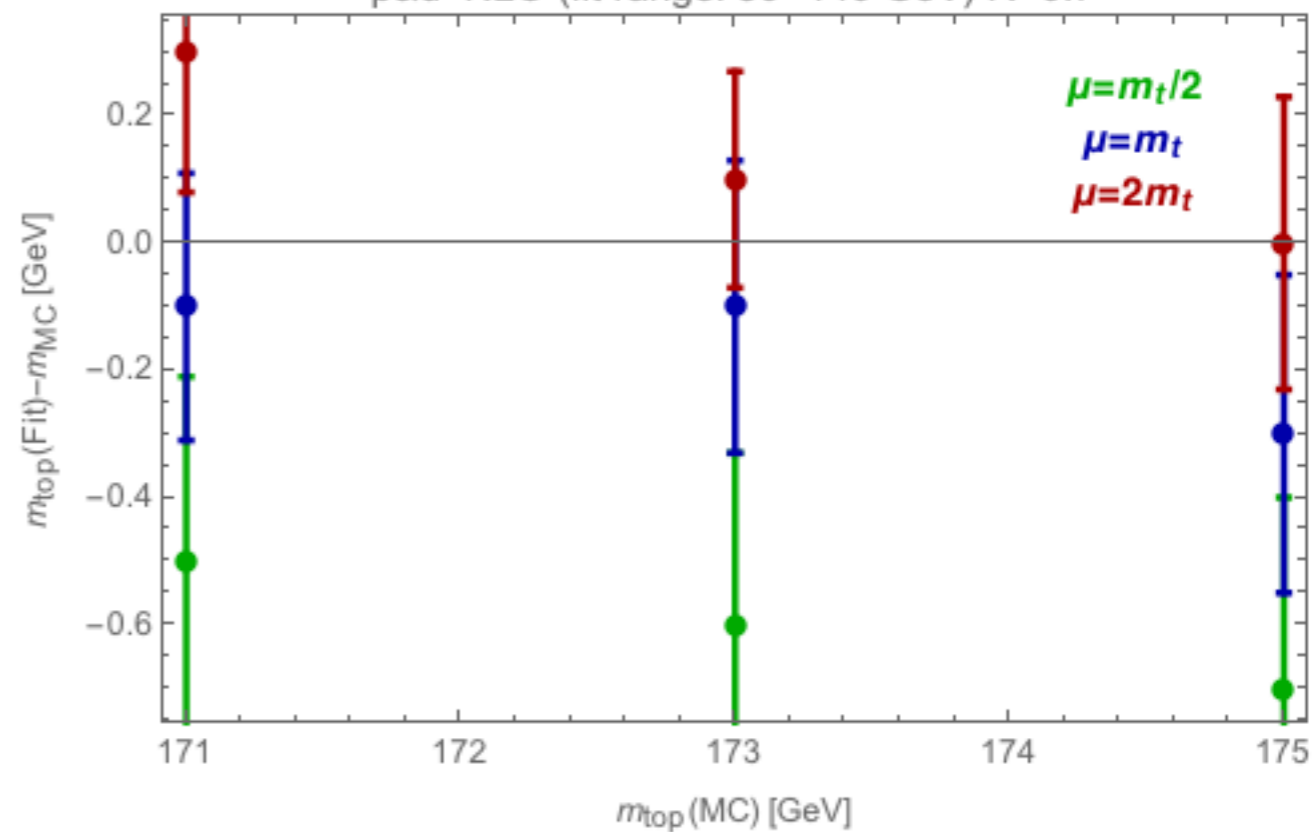
p&amp;d-NLO (fit range: 30–140 GeV) R=0.7



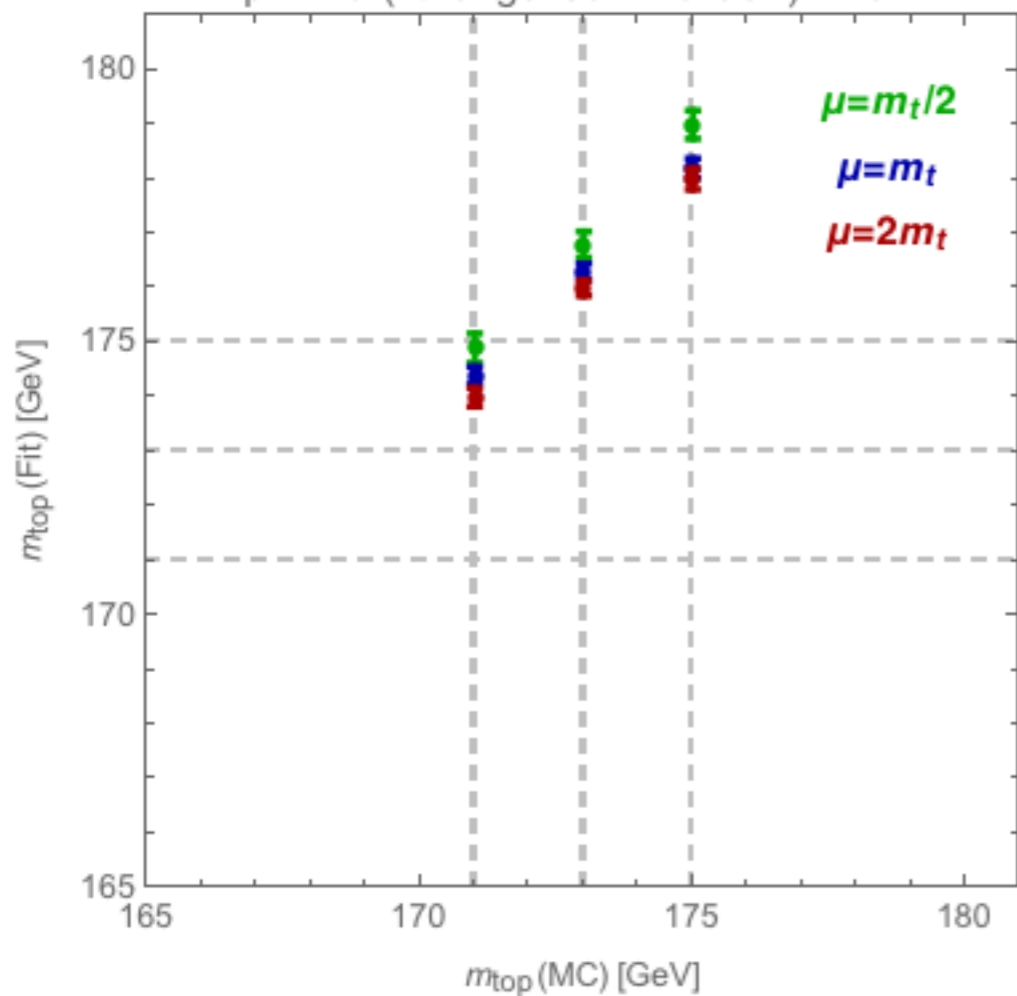
p&amp;d-NLO (fit range: 30–140 GeV) R=0.7



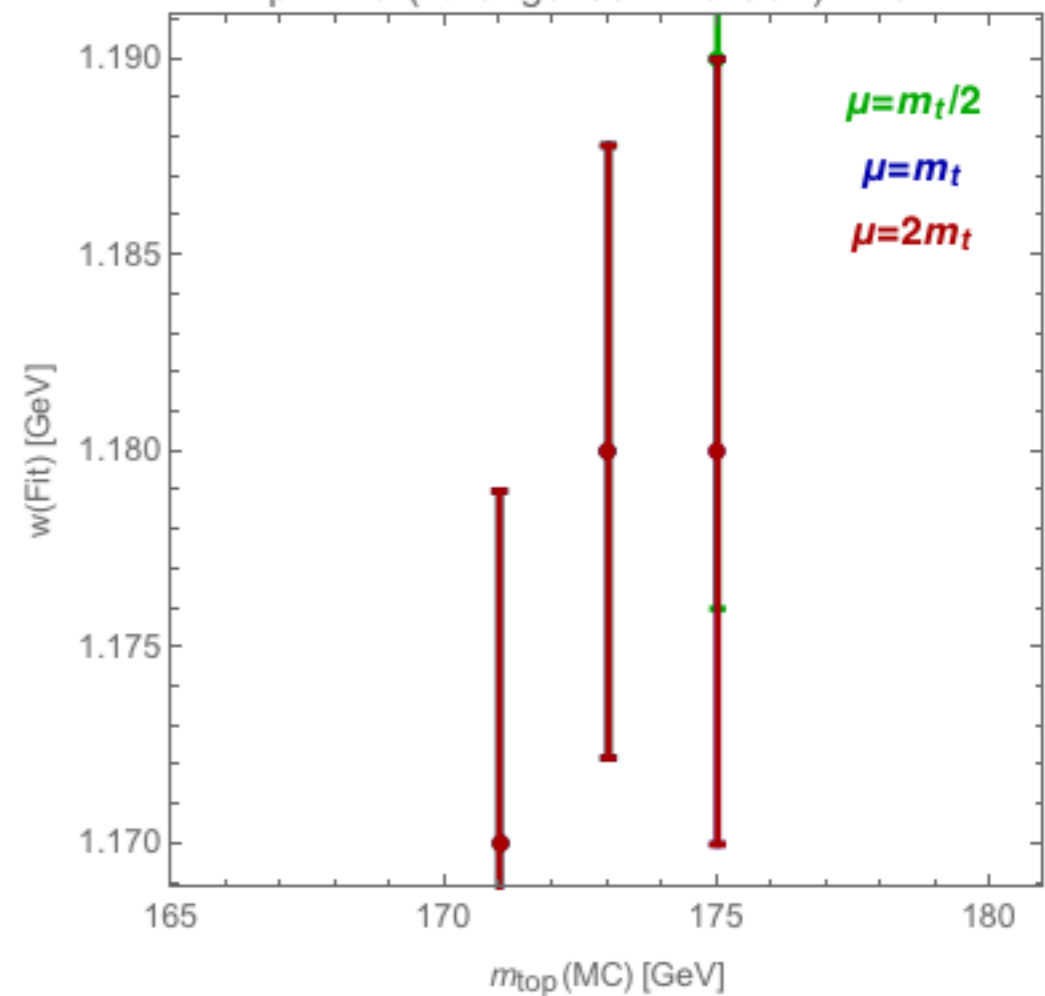
p&amp;d-NLO (fit range: 30–140 GeV) R=0.7



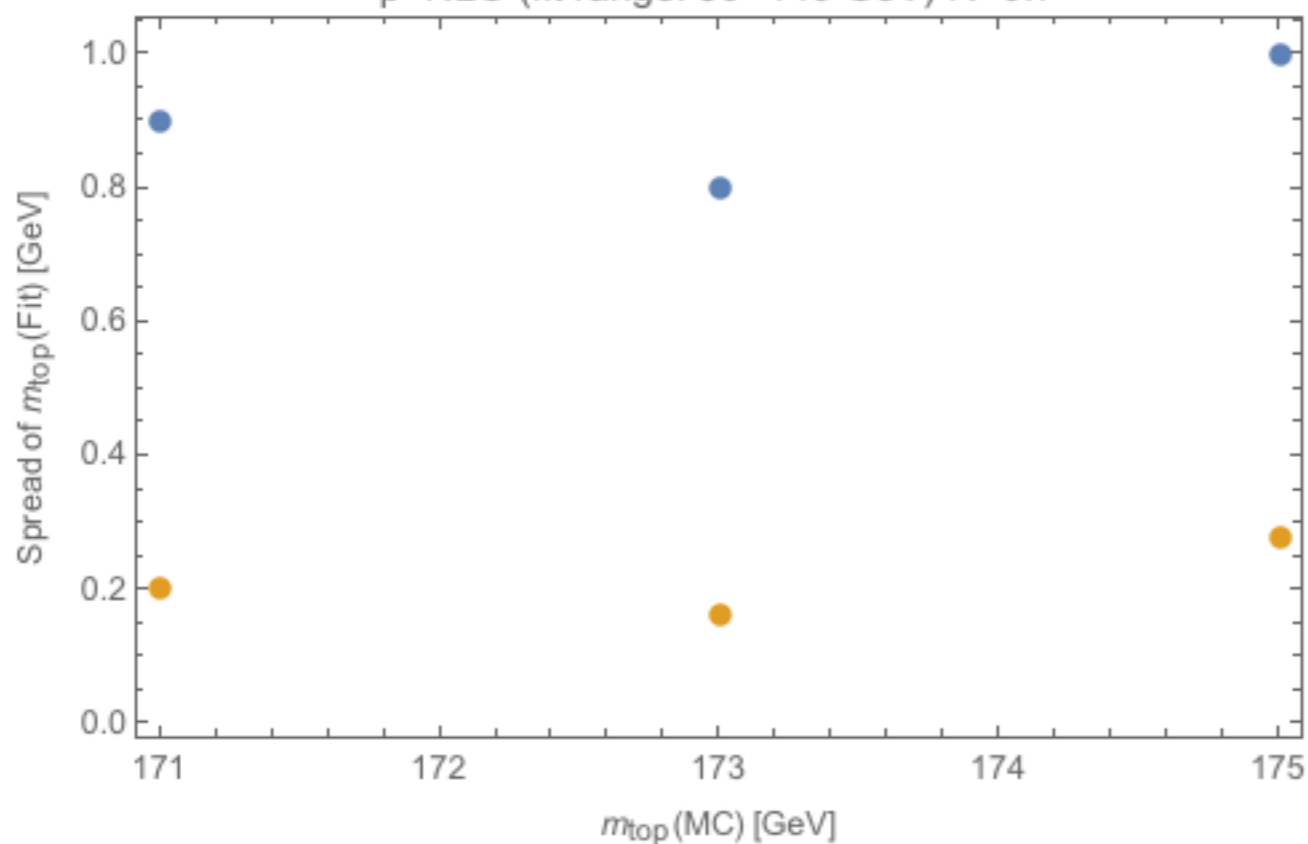
p-NLO (fit range: 30–140 GeV) R=0.7



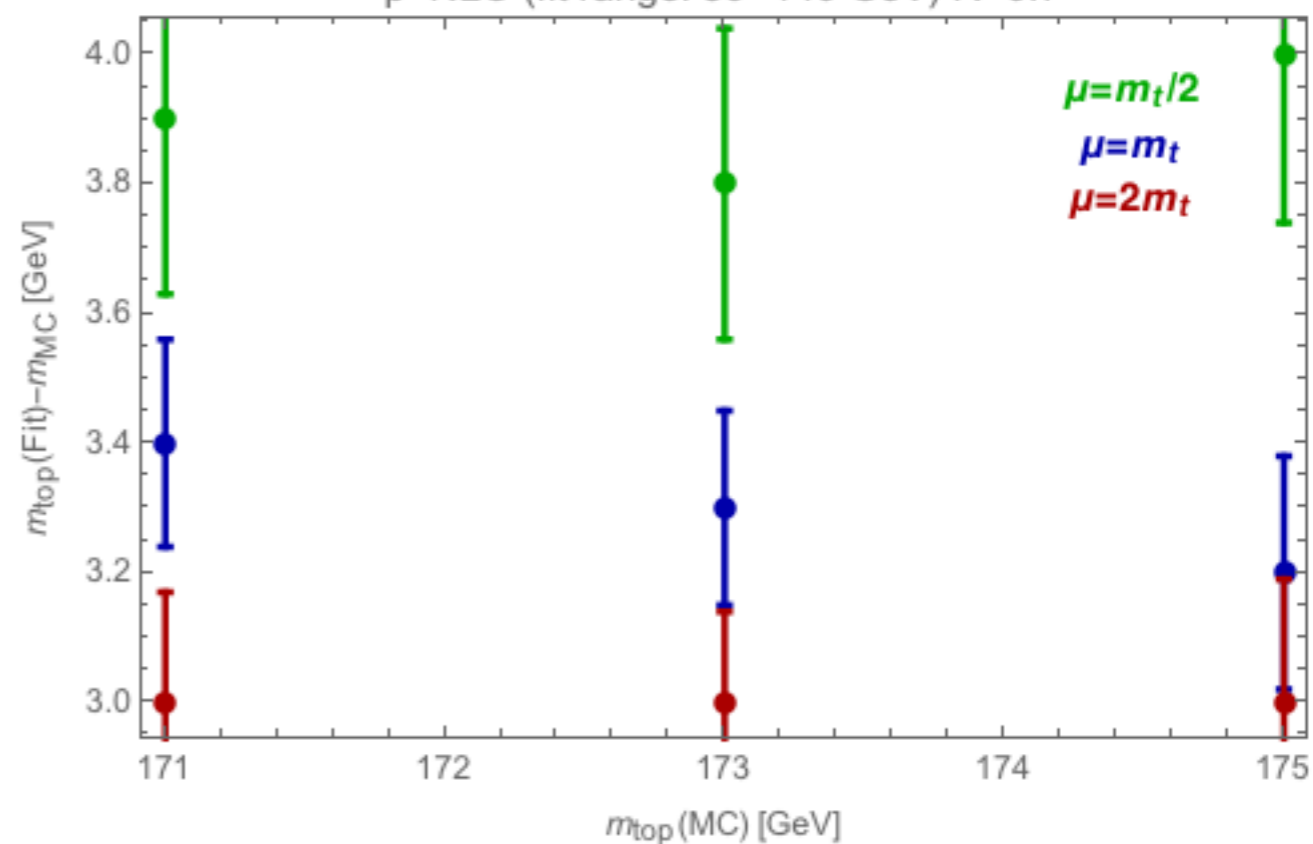
p-NLO (fit range: 30–140 GeV) R=0.7



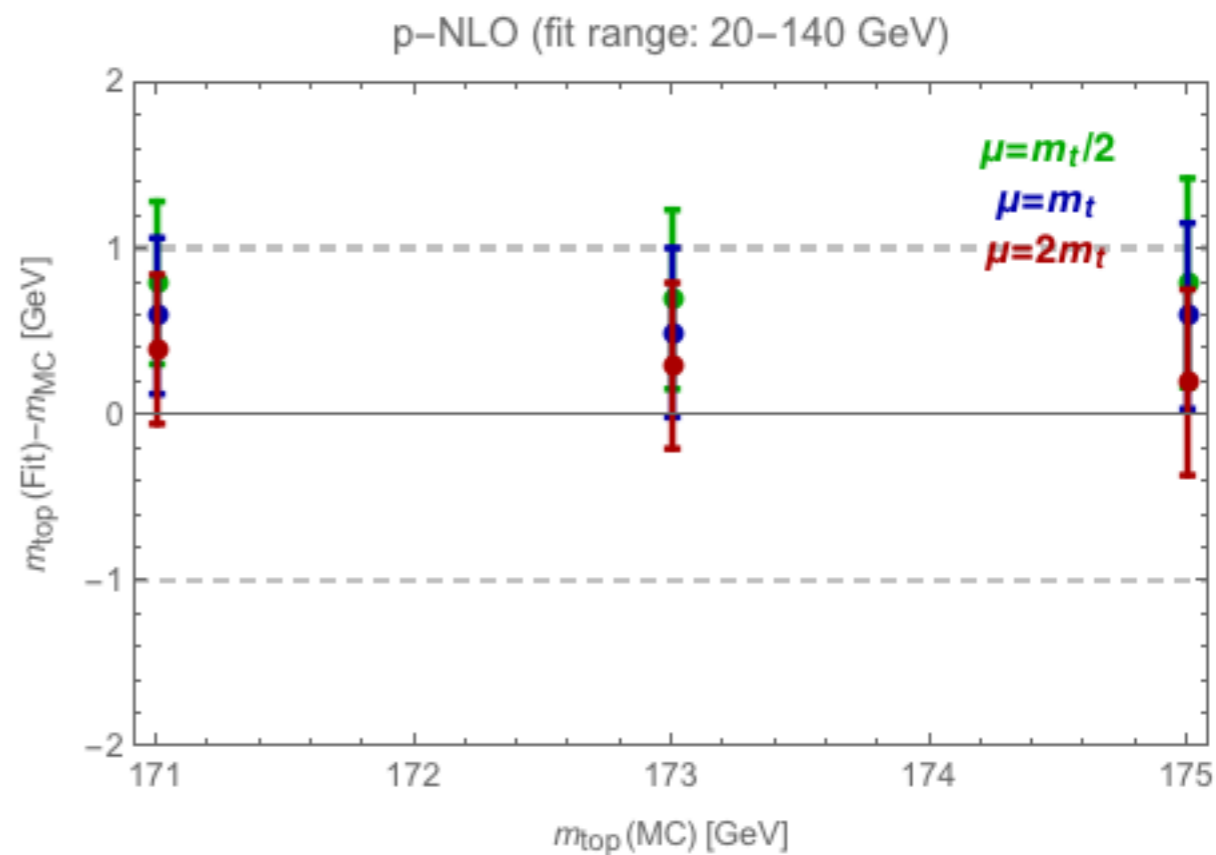
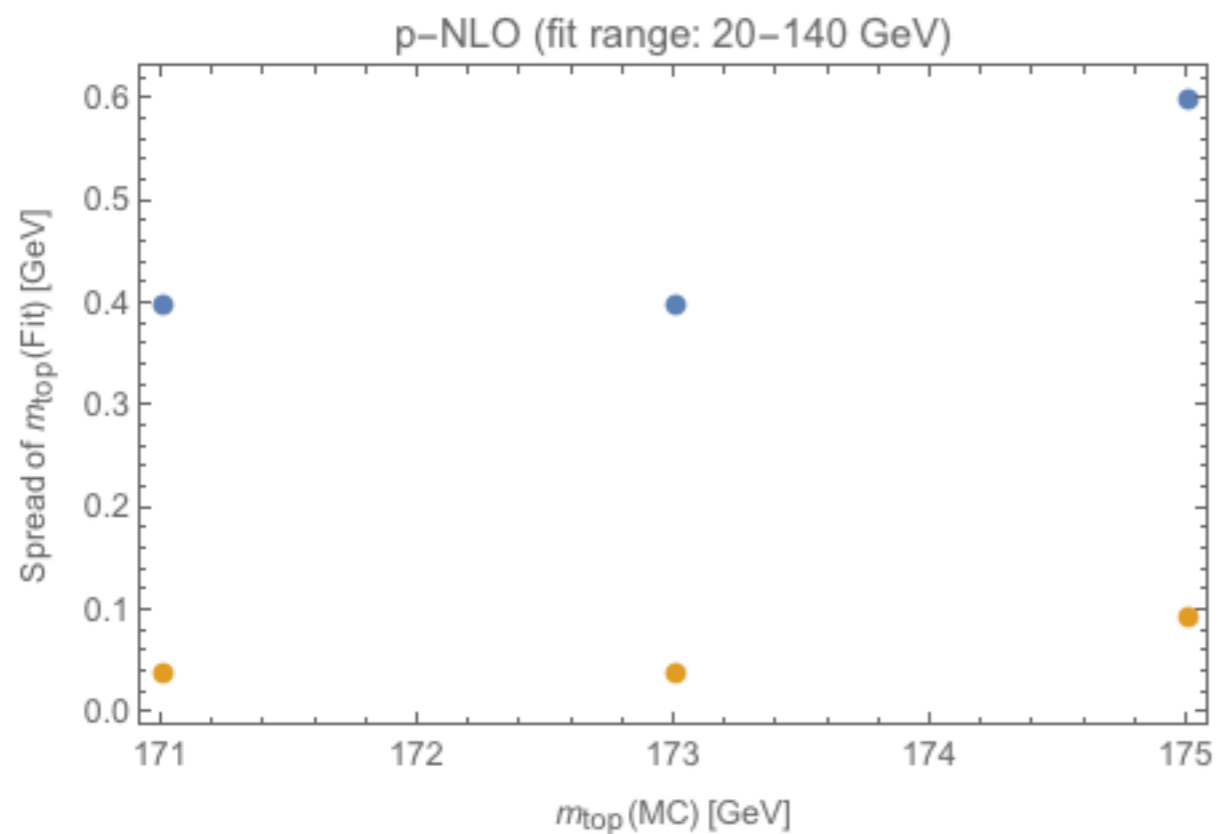
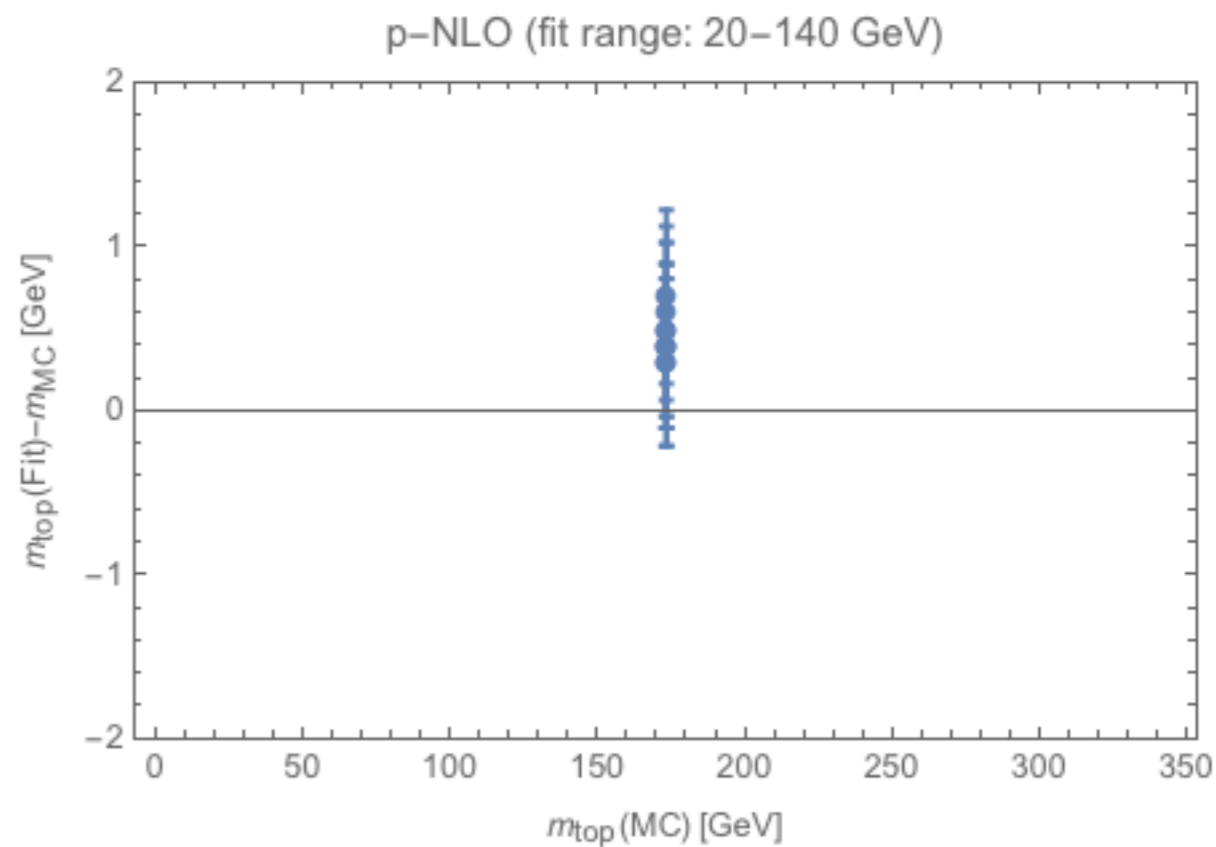
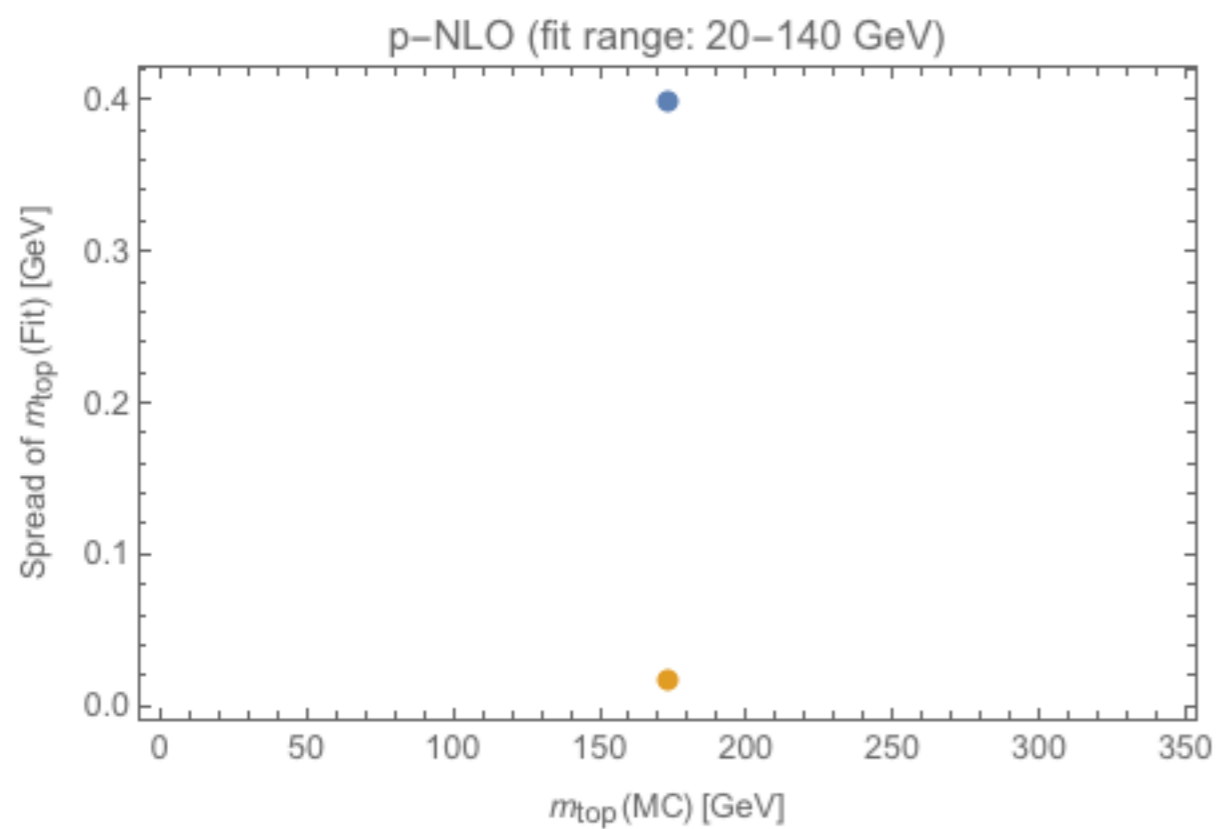
p-NLO (fit range: 30–140 GeV) R=0.7



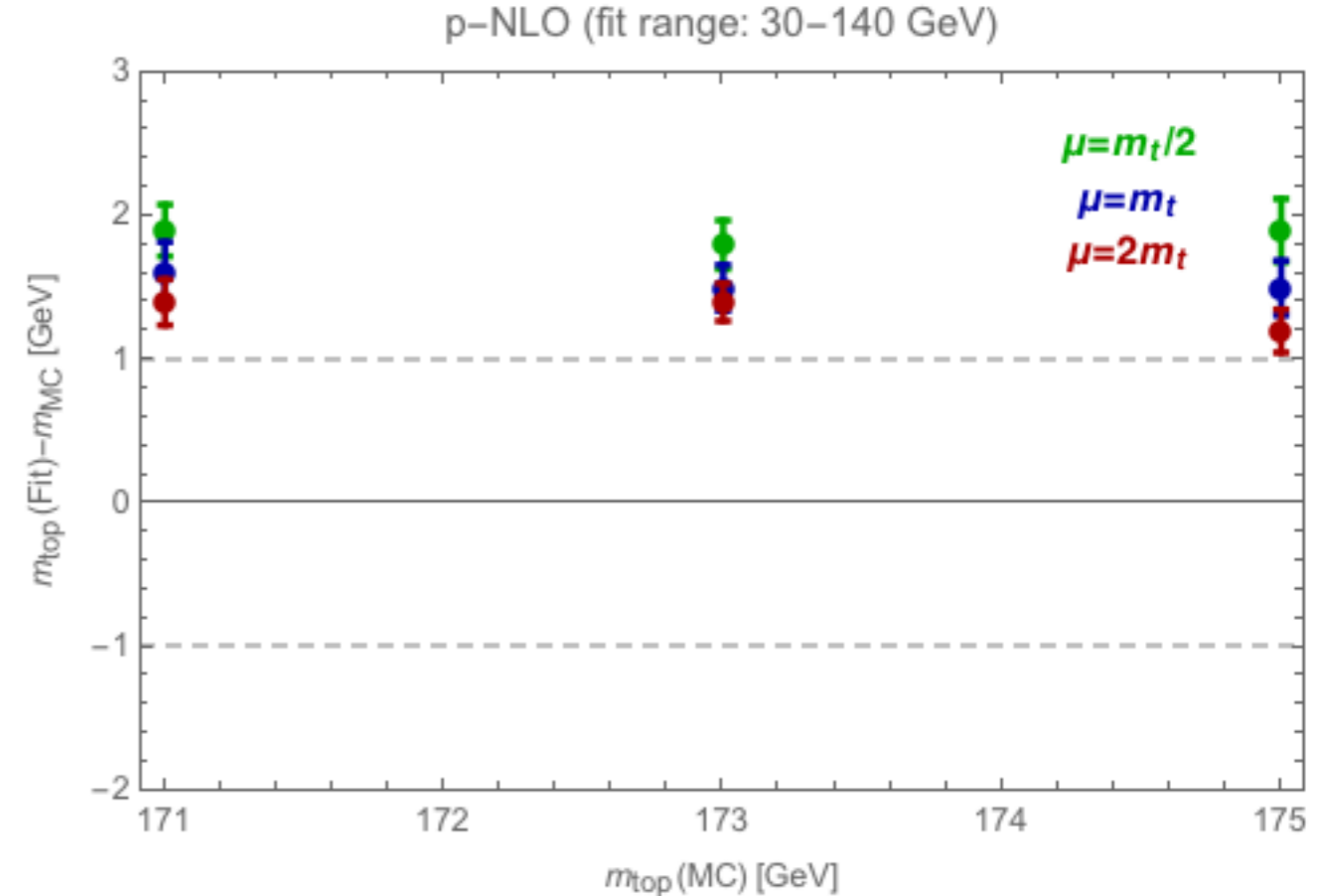
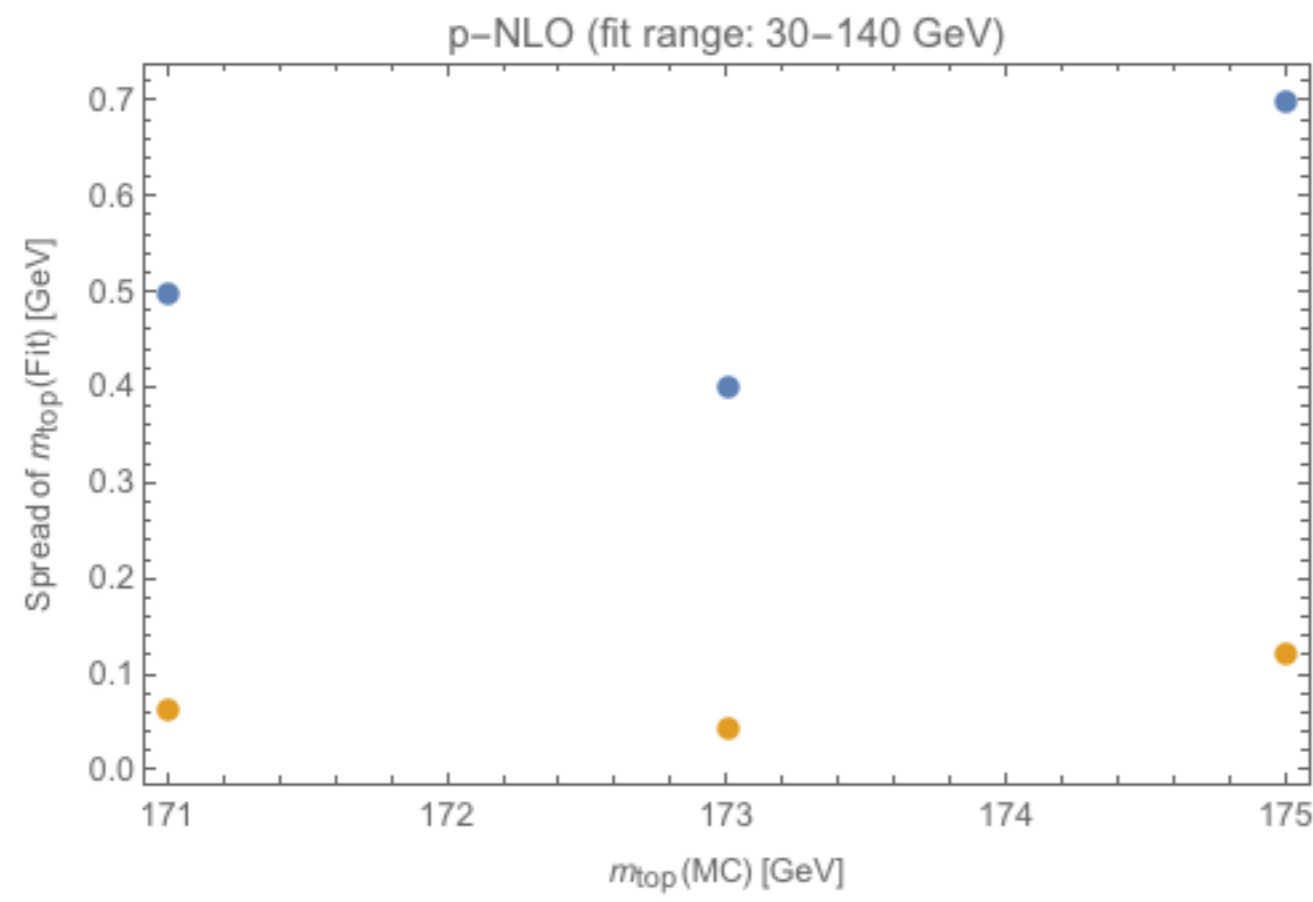
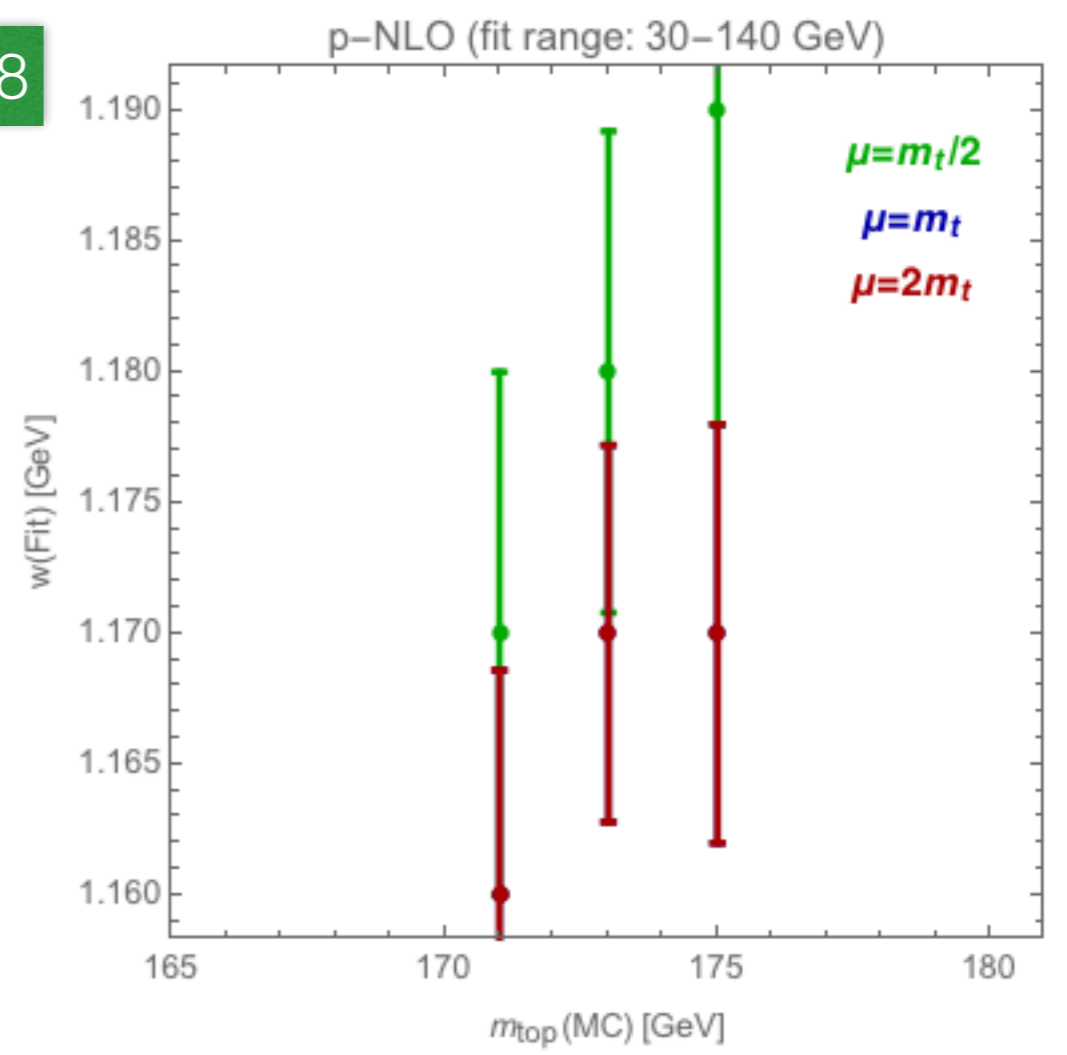
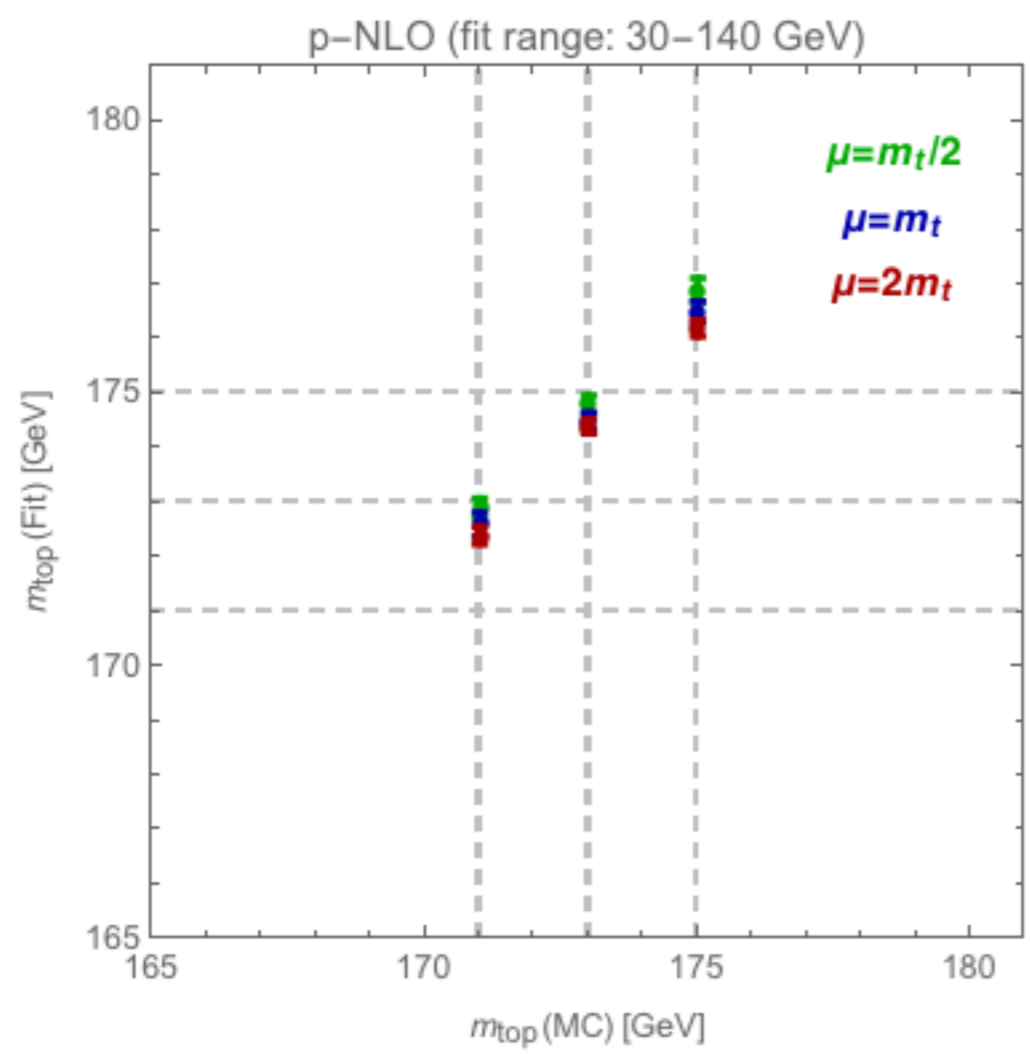
p-NLO (fit range: 30–140 GeV) R=0.7



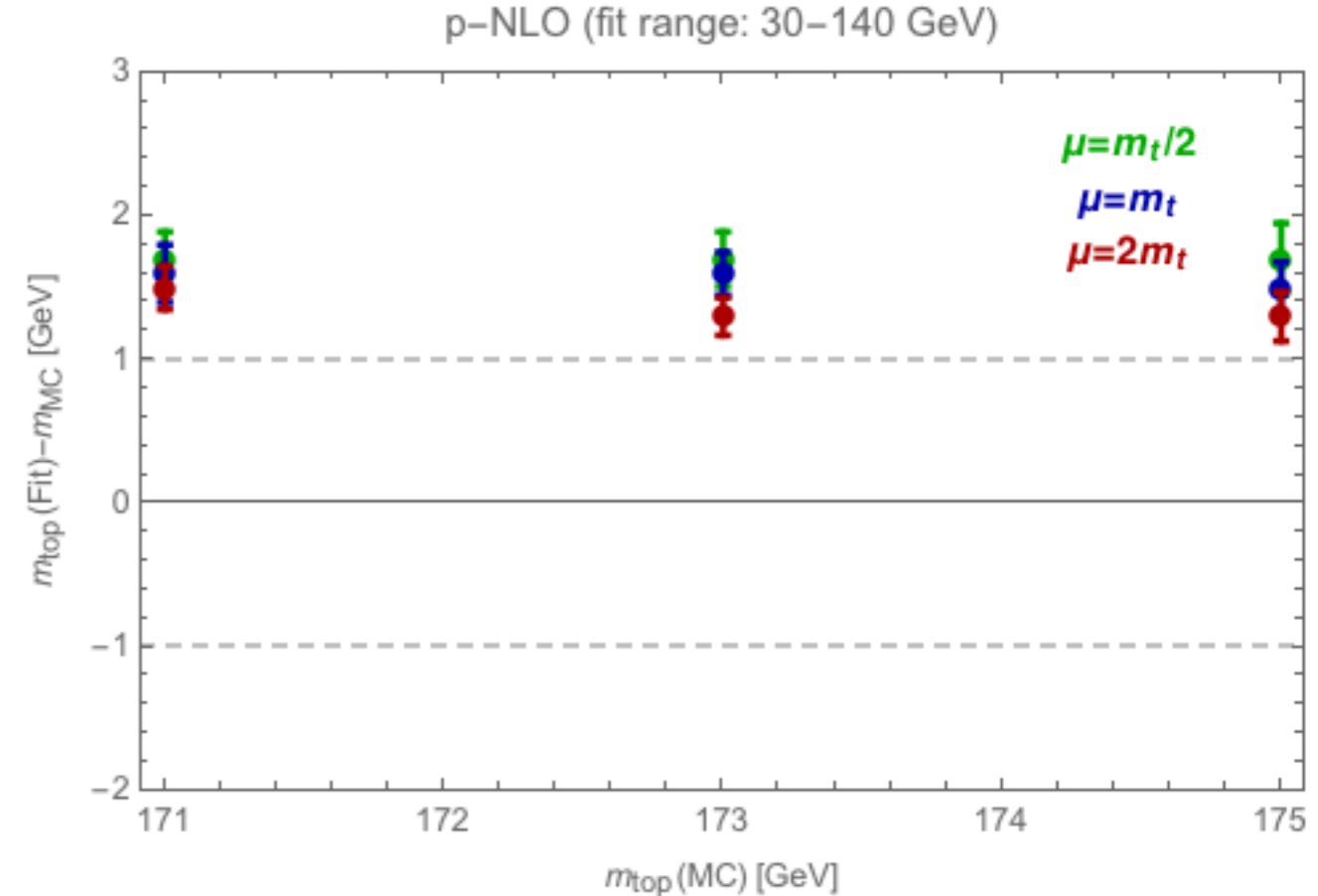
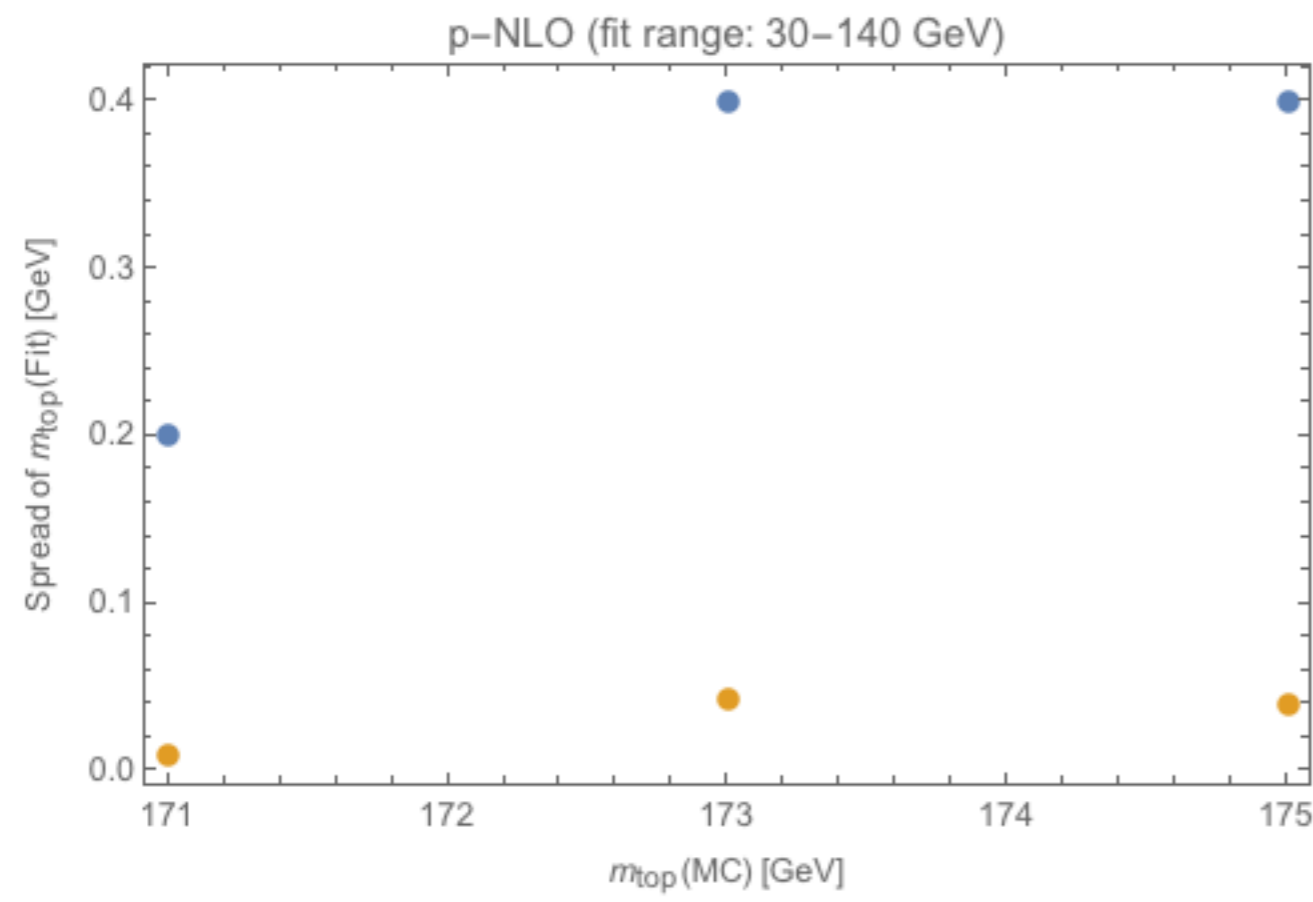
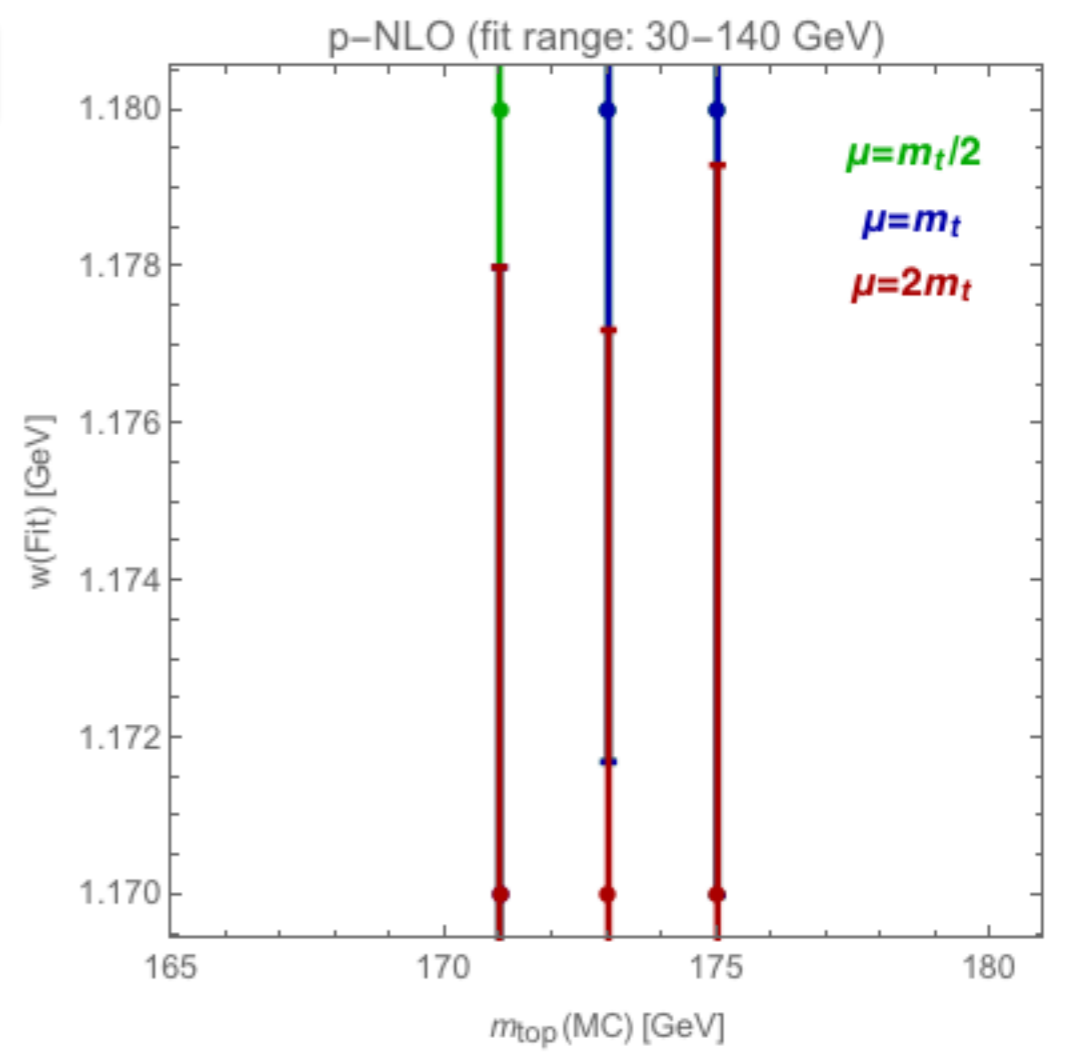
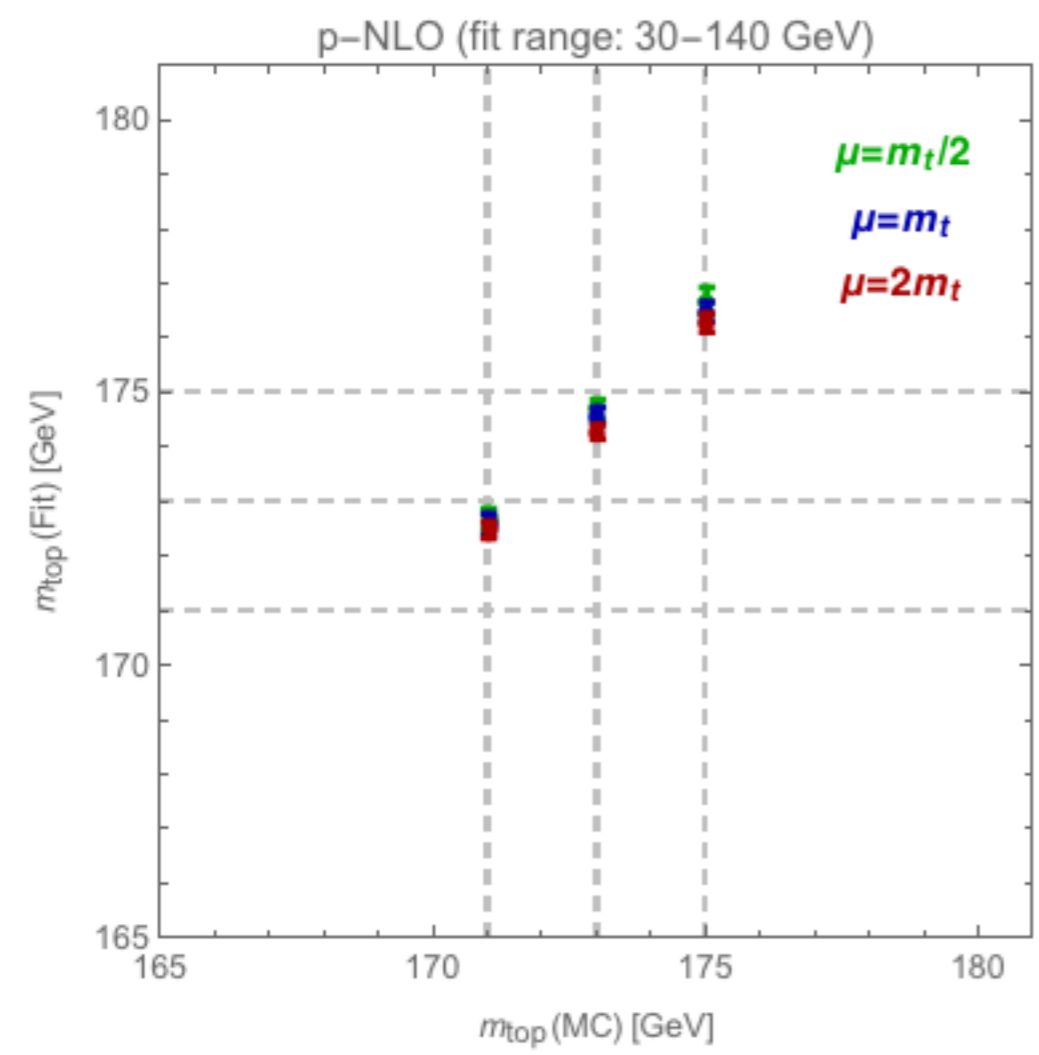
$$\mu_F \neq \mu_R$$



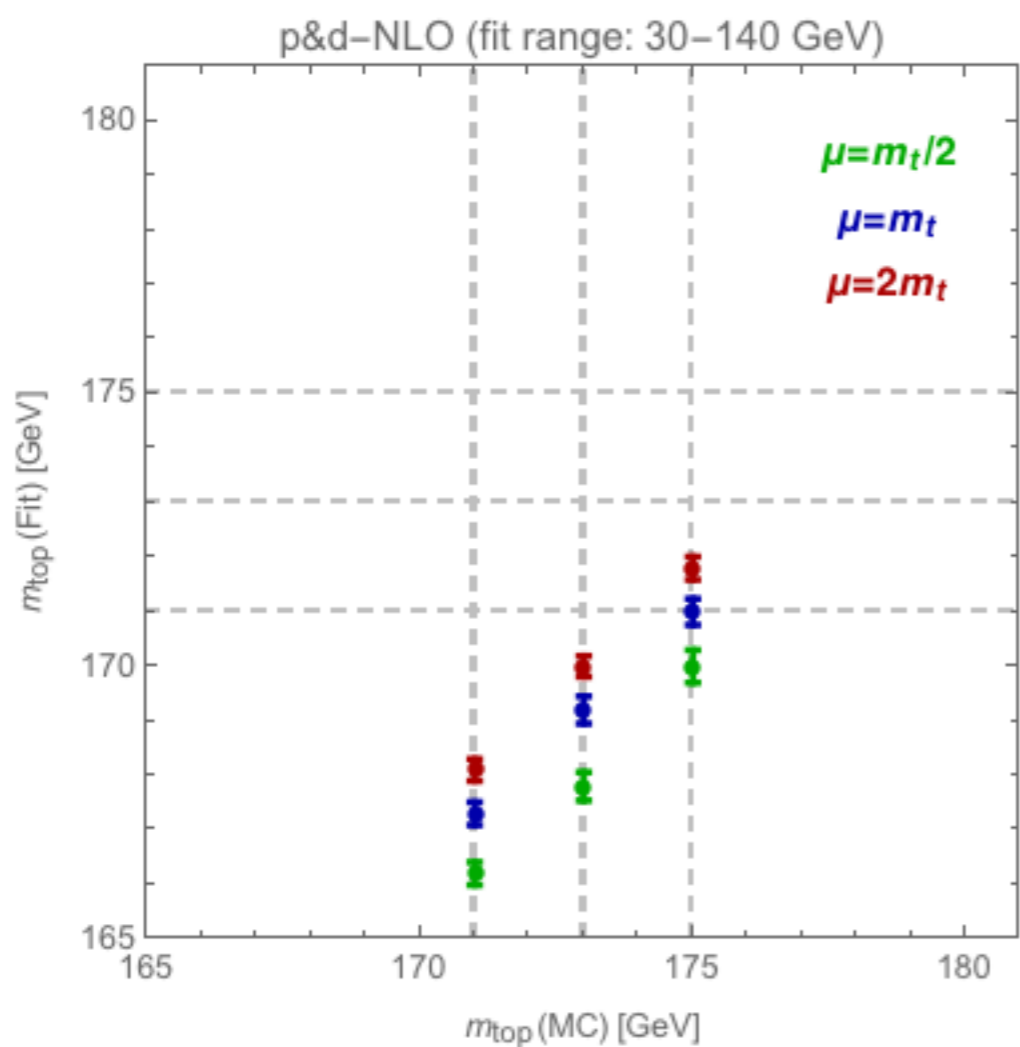
MSTW08



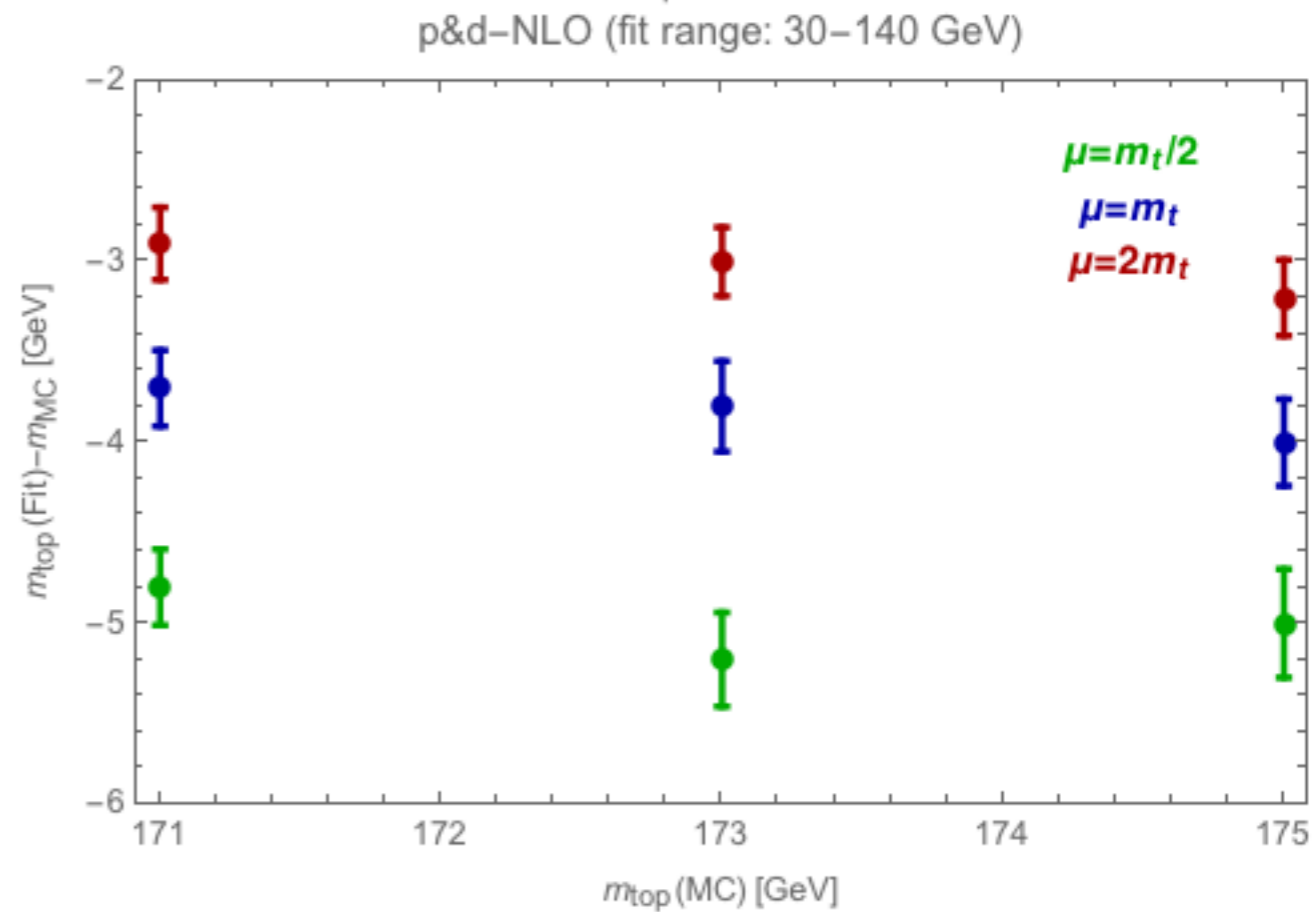
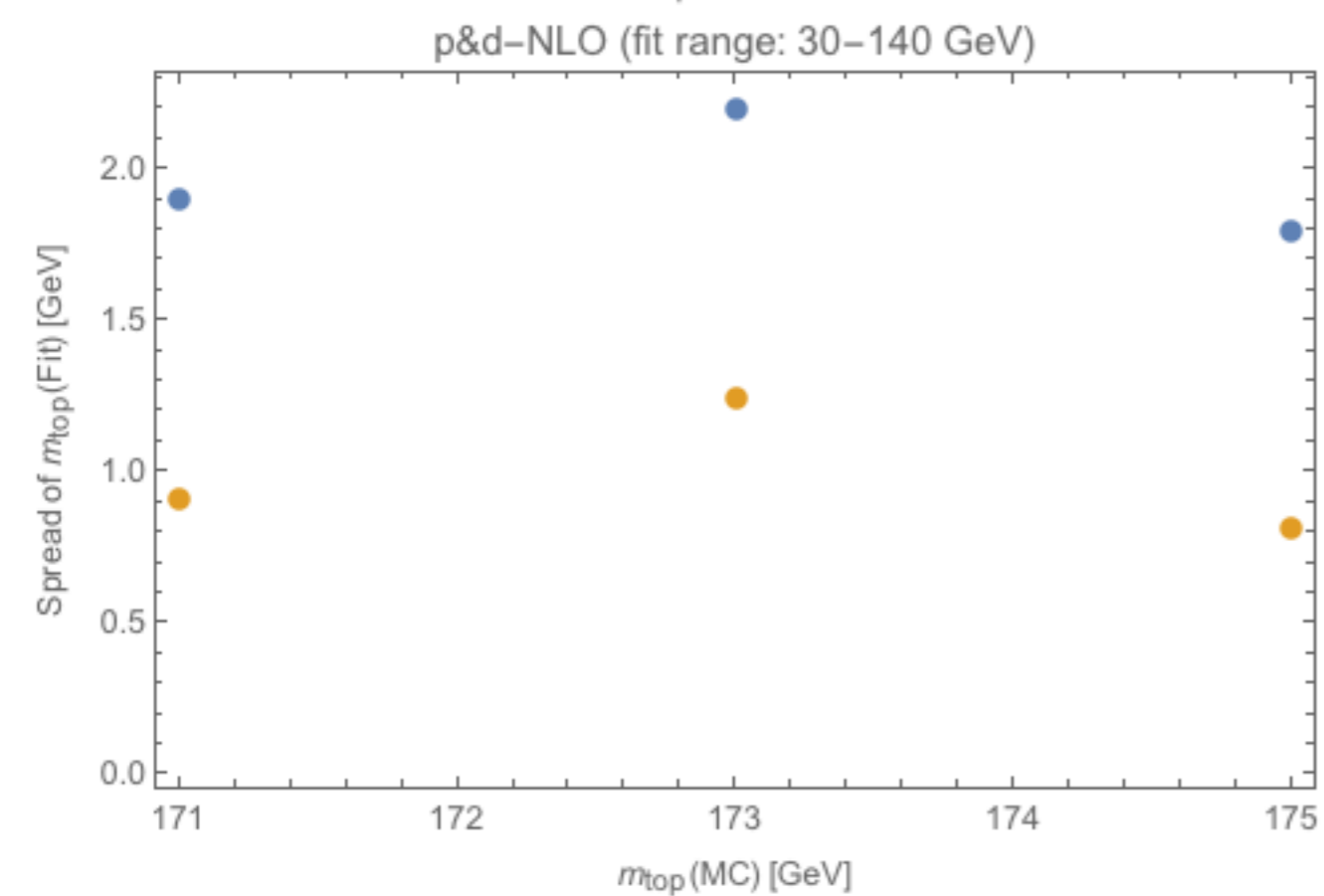
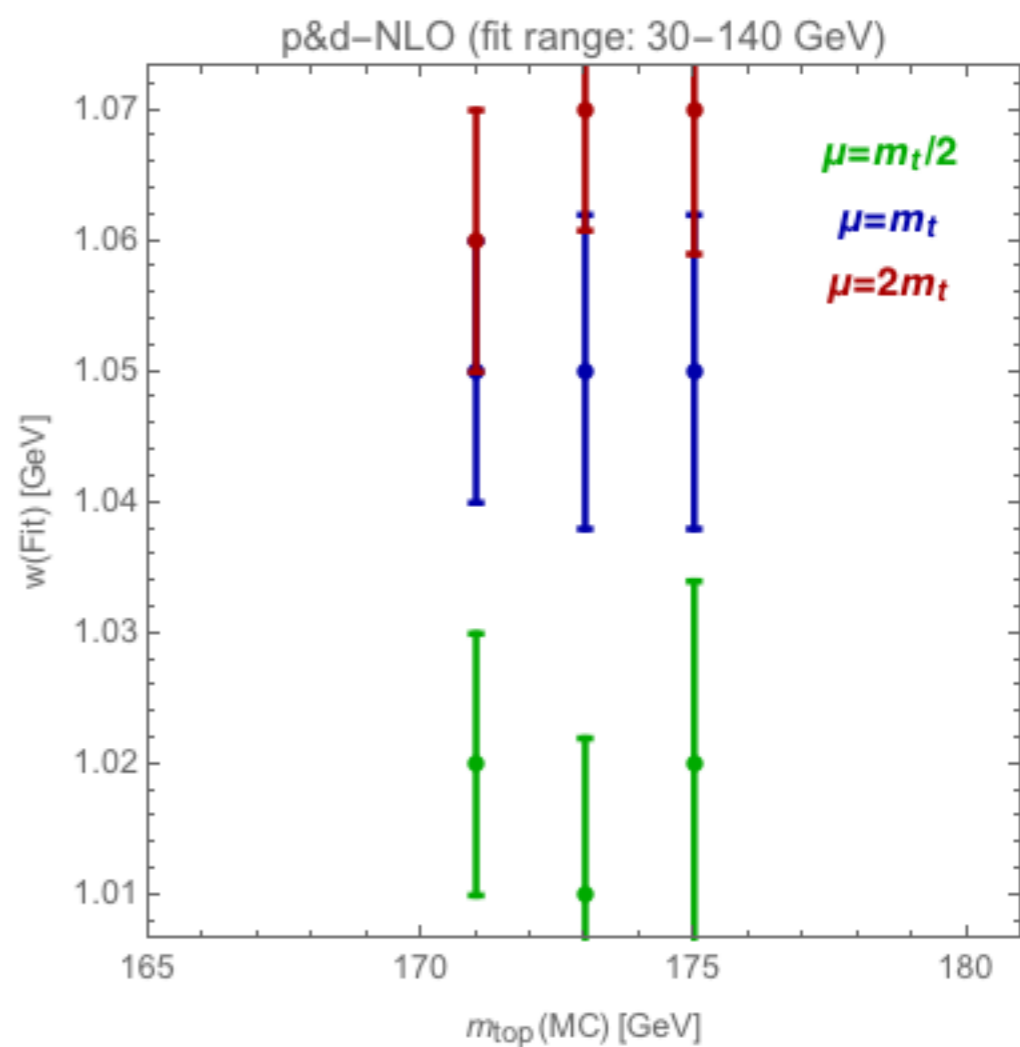
CT10

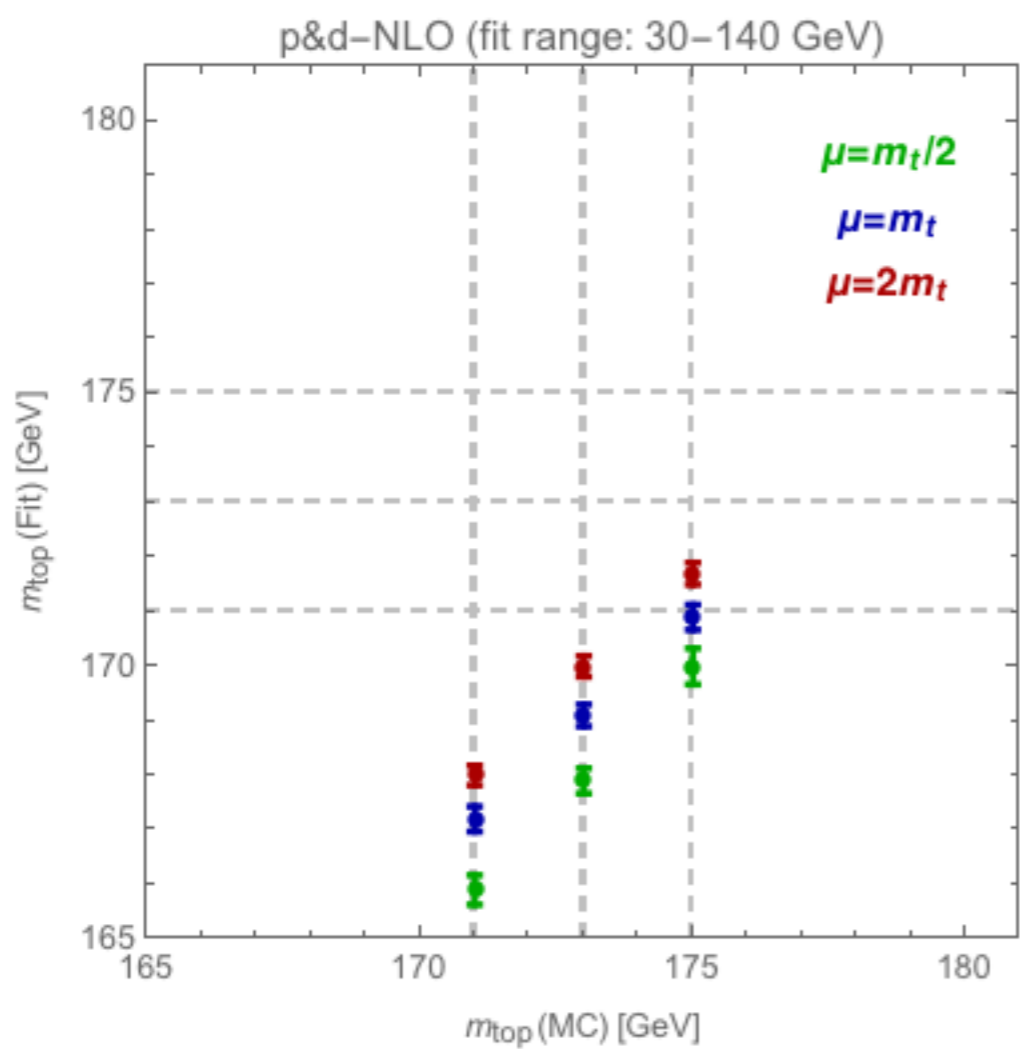




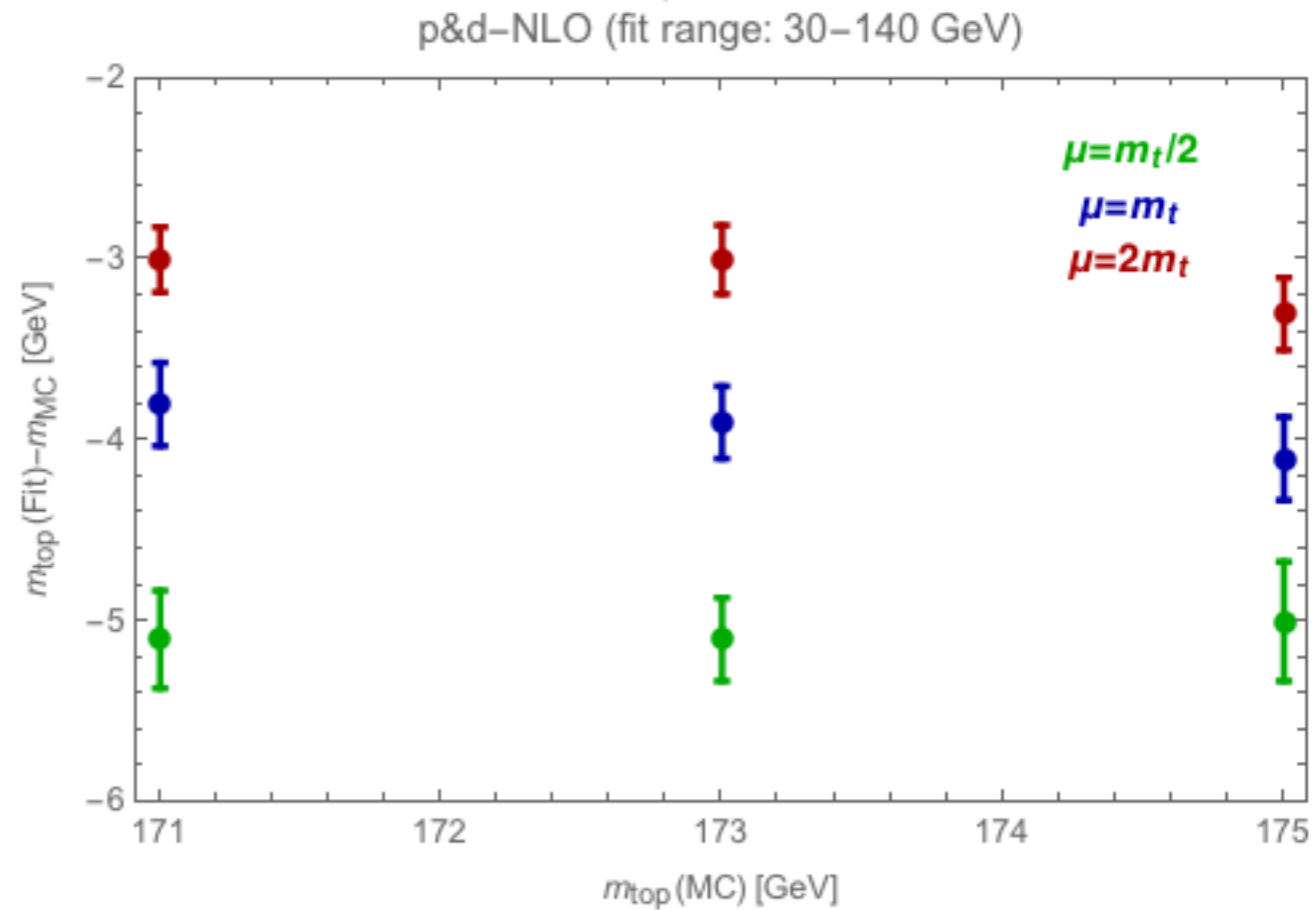
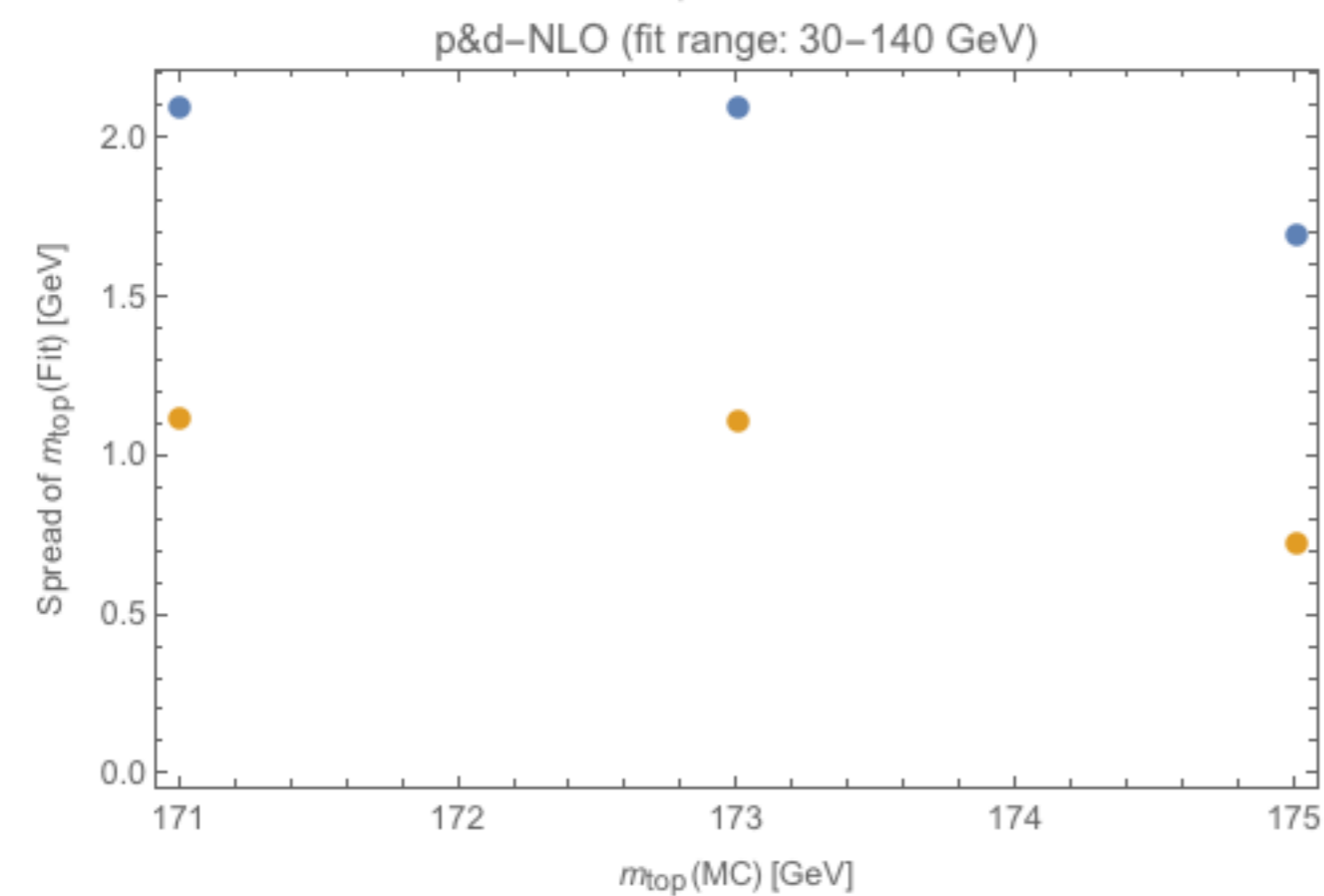
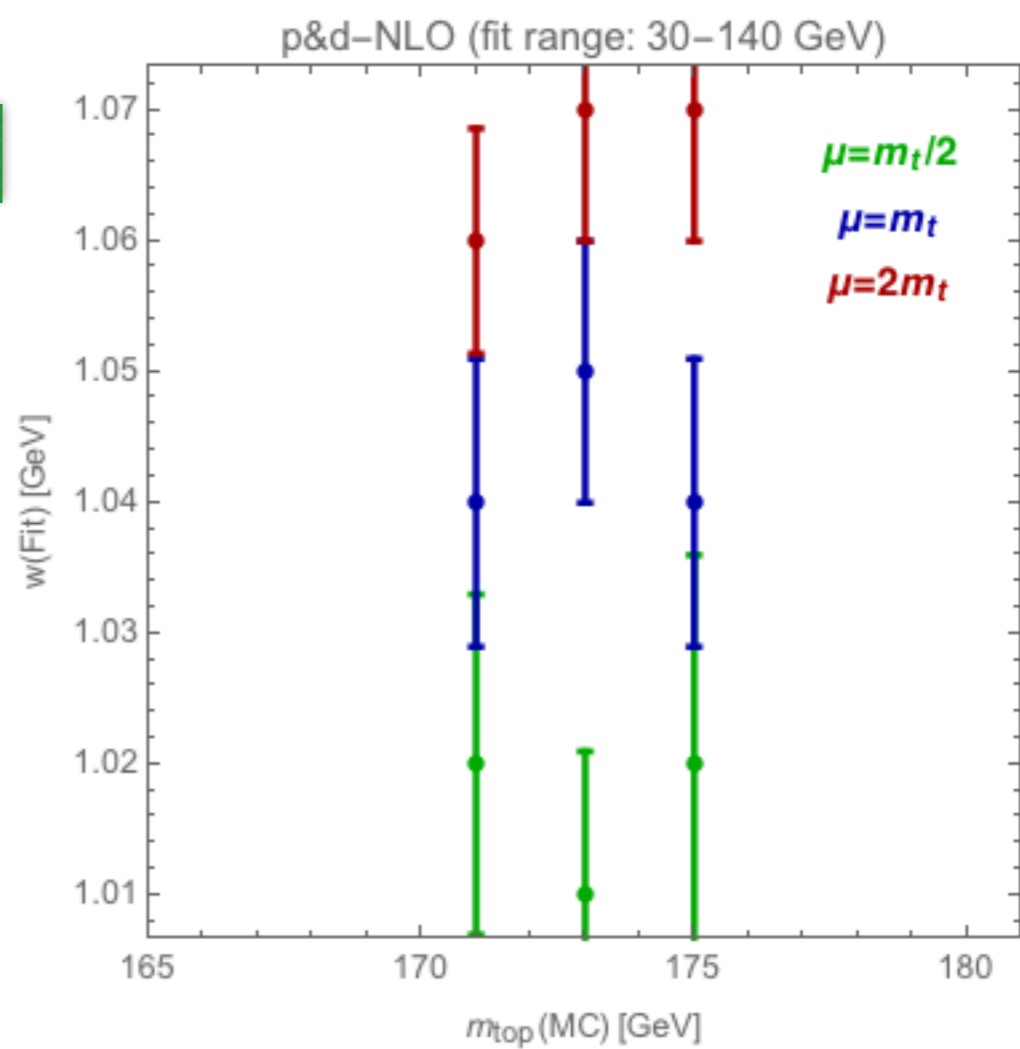


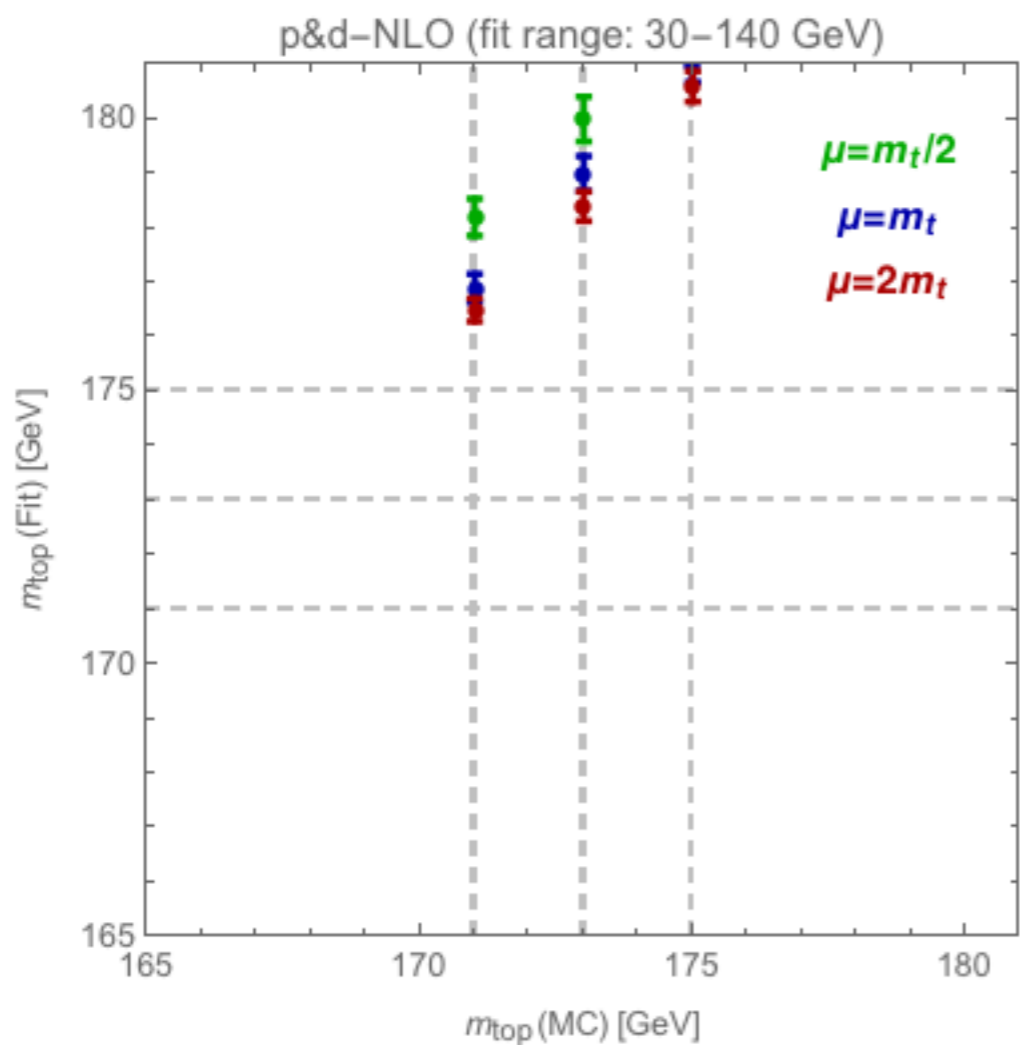
R=0.5  
 CT10



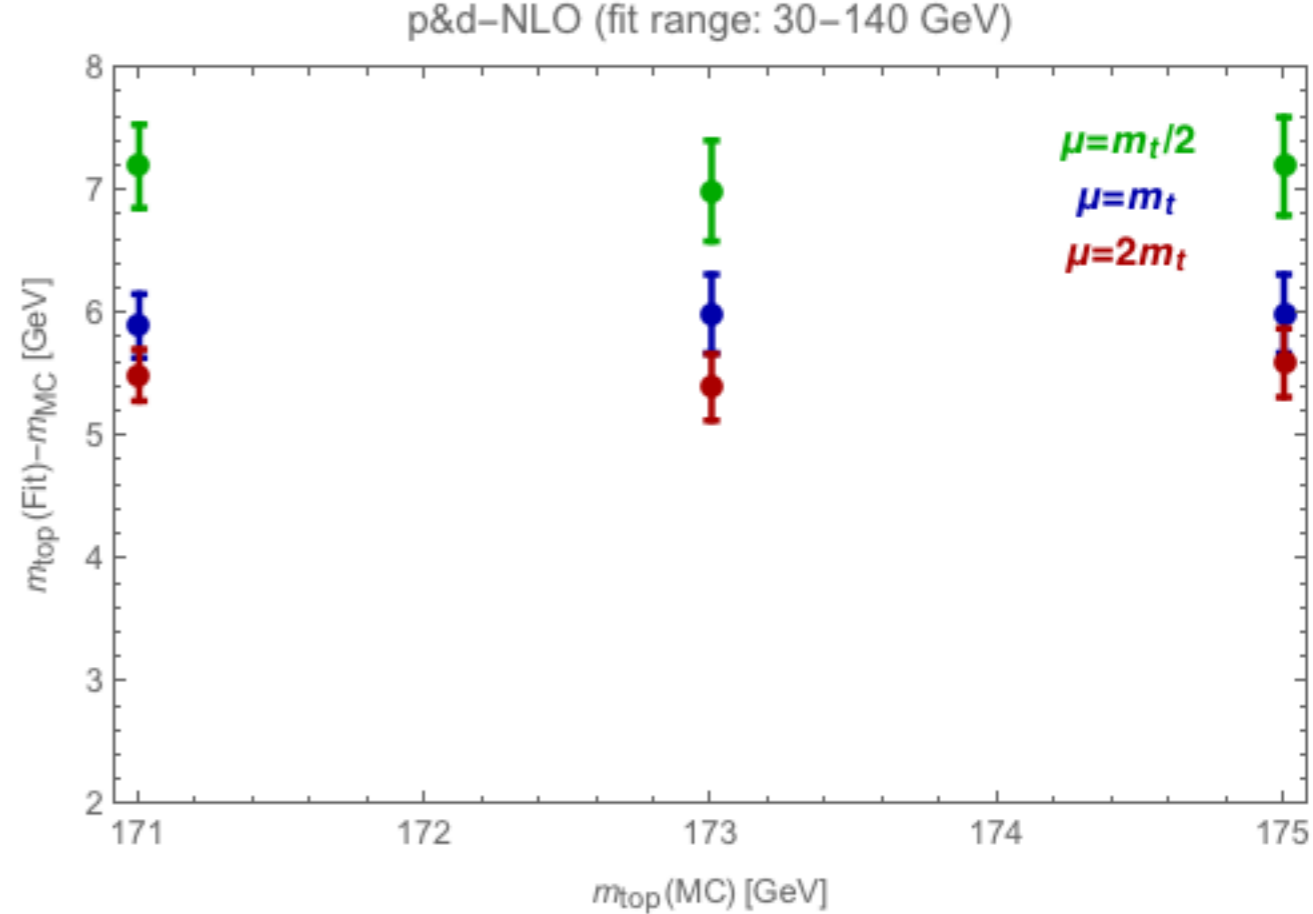
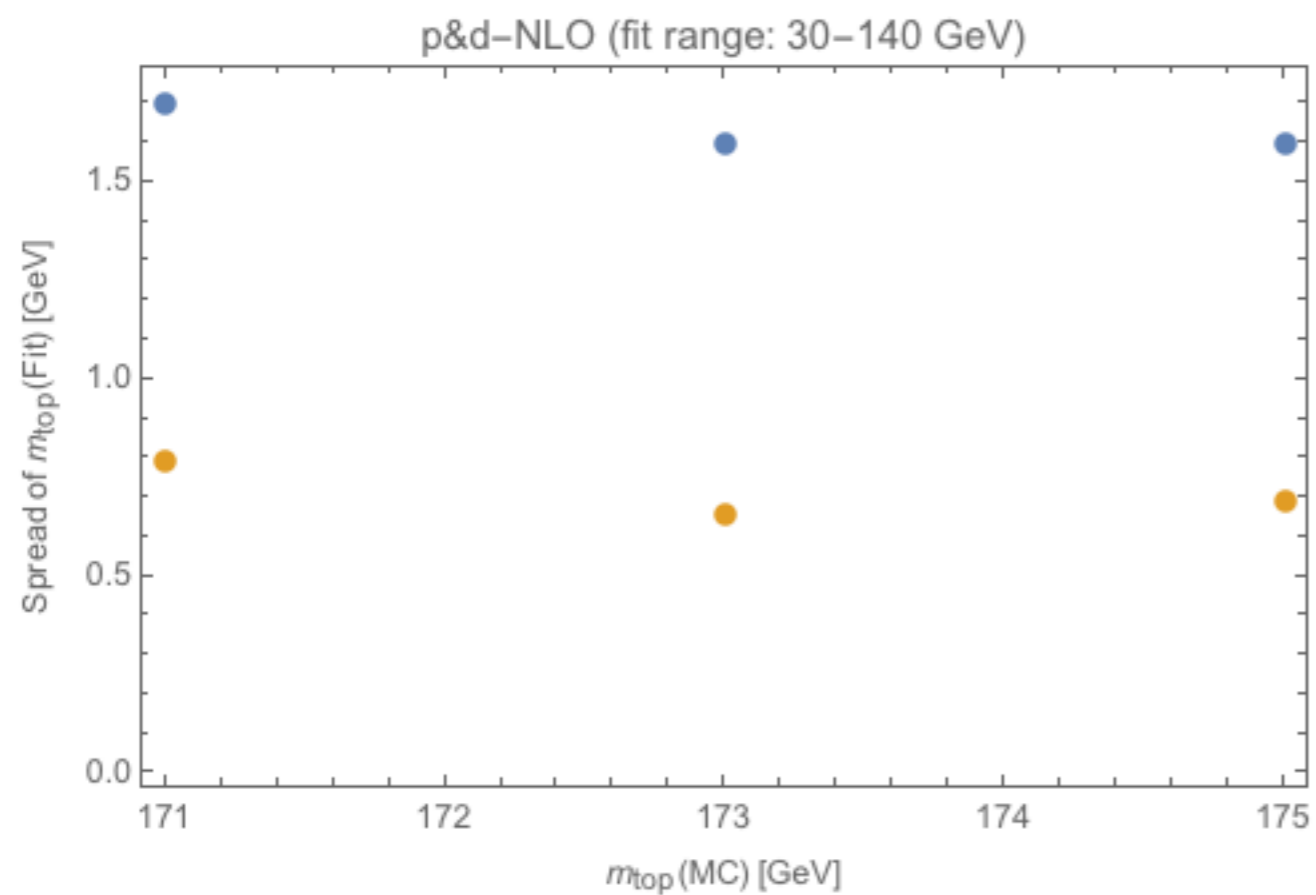
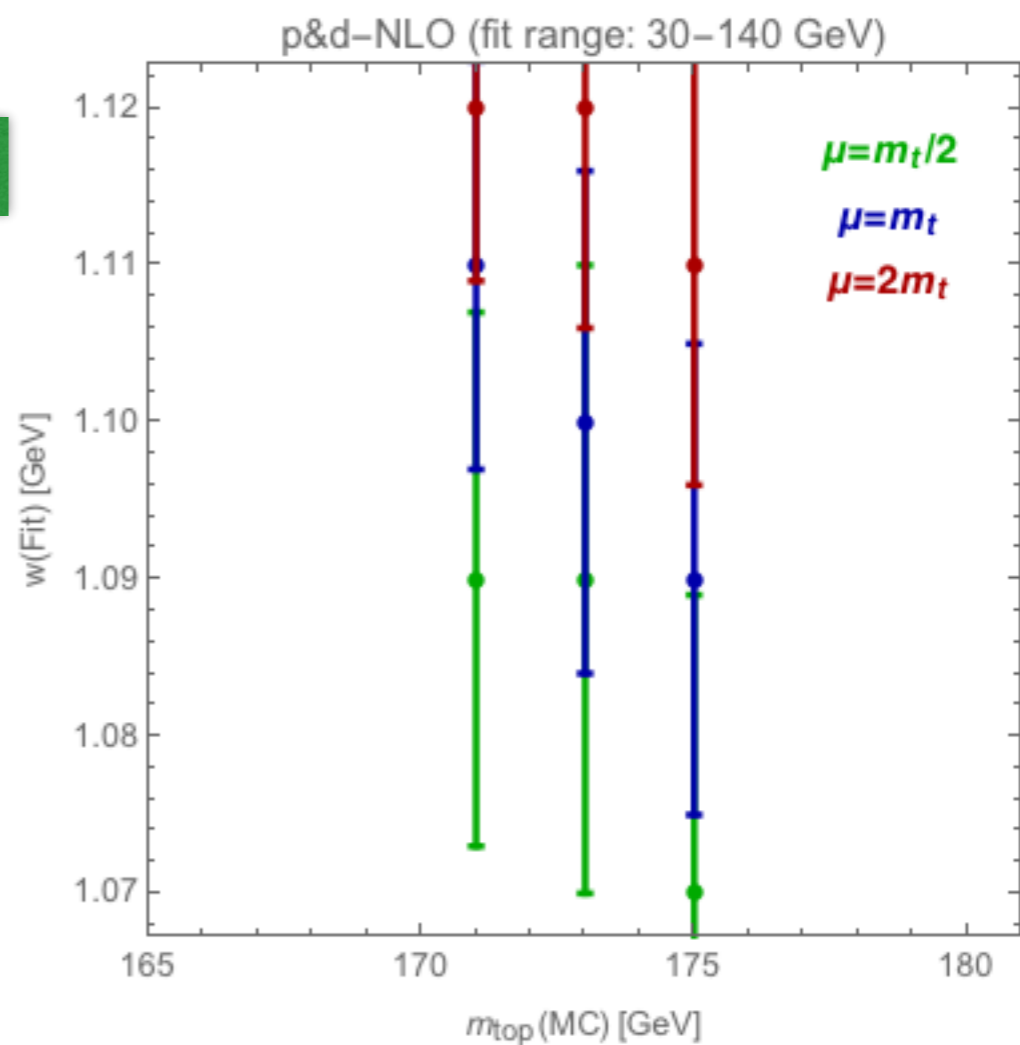


R=0.5  
MSTW08

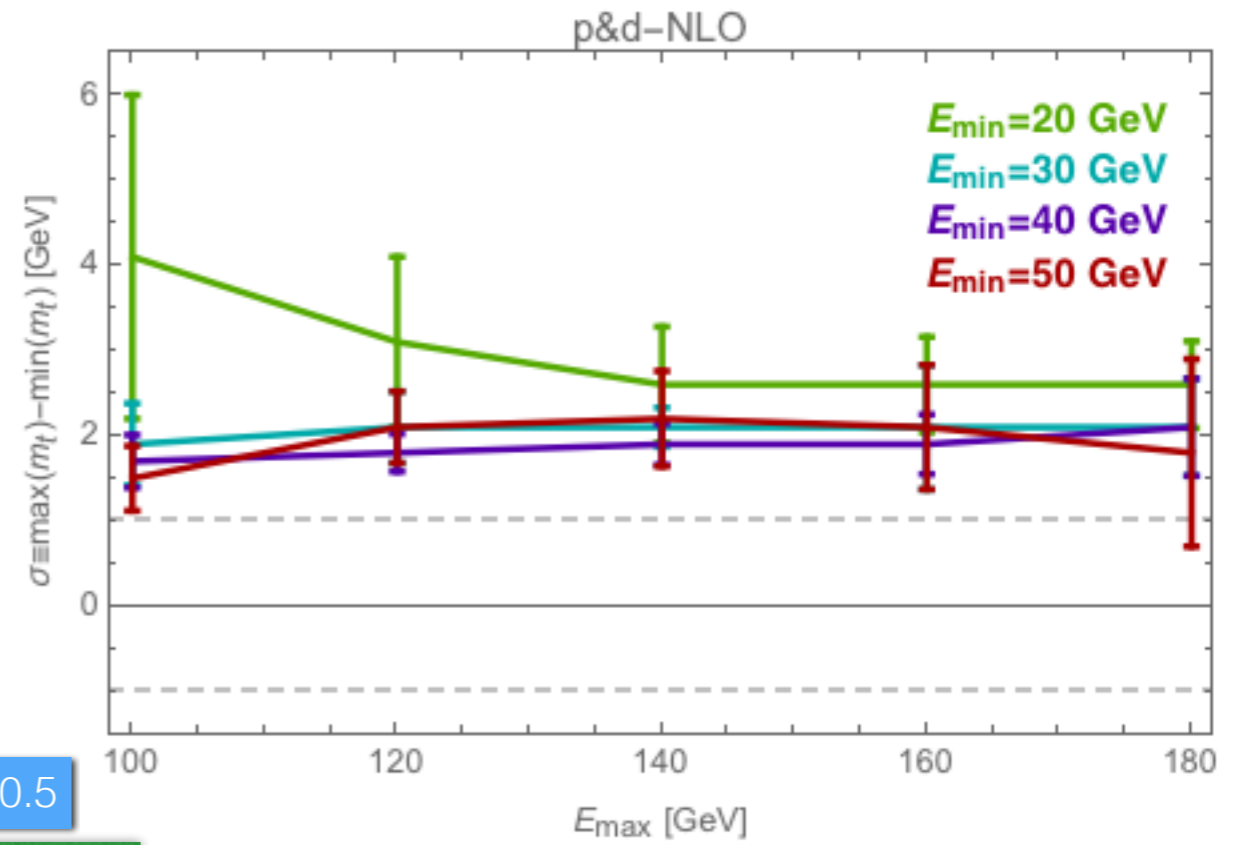
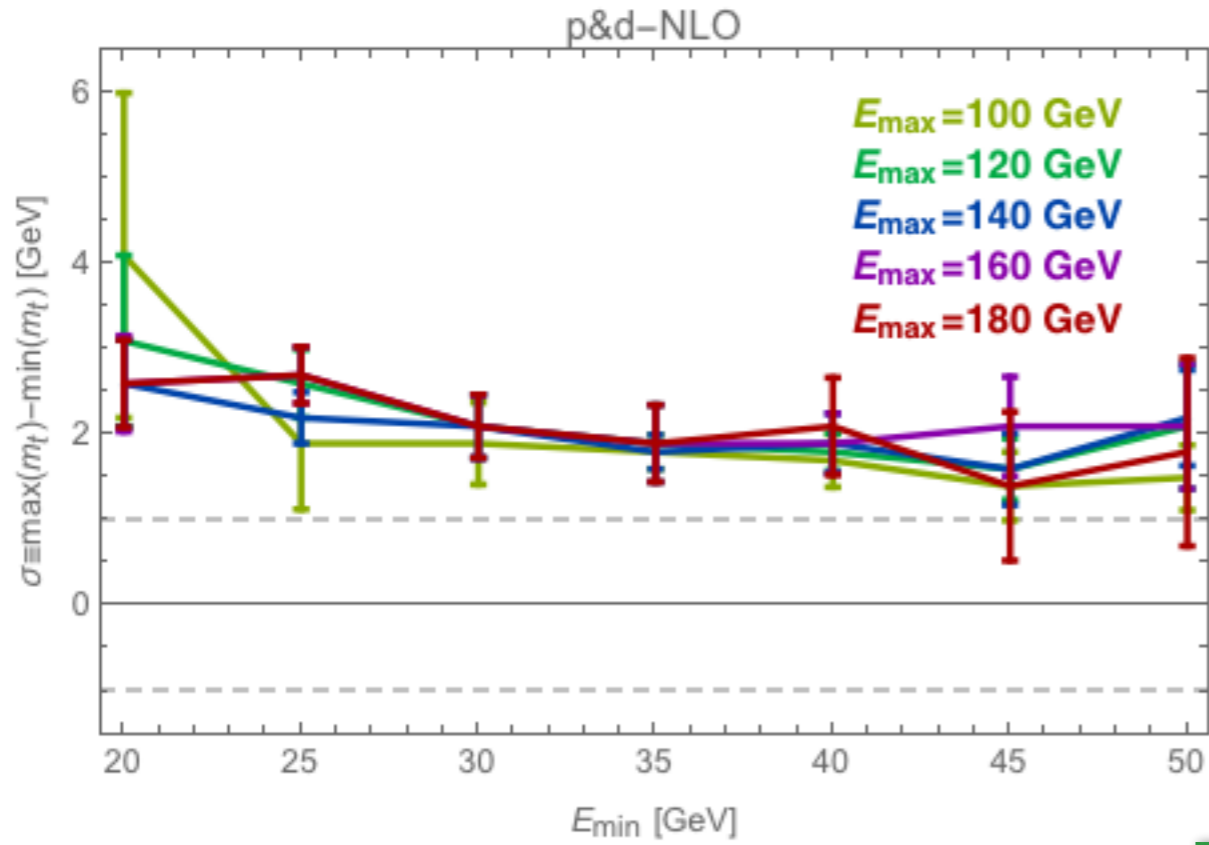




R=1.0  
MSTW08

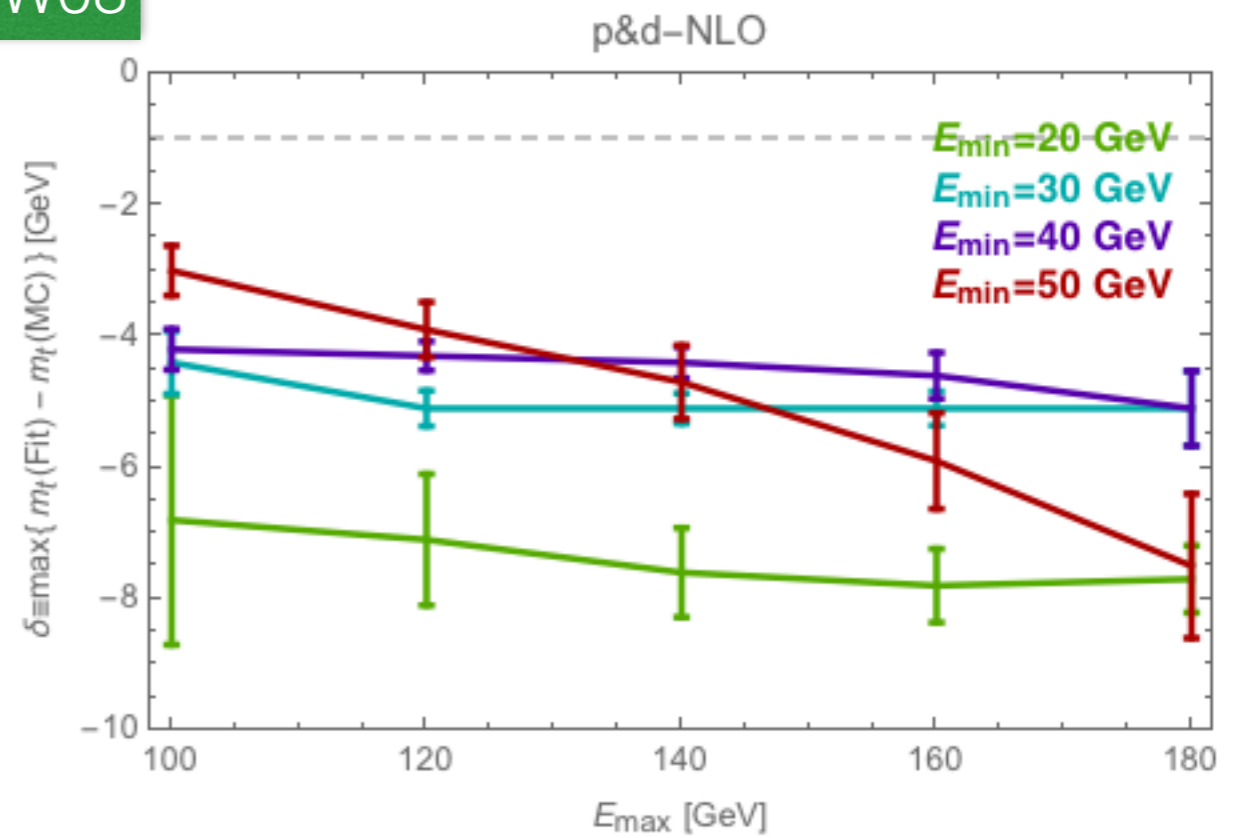
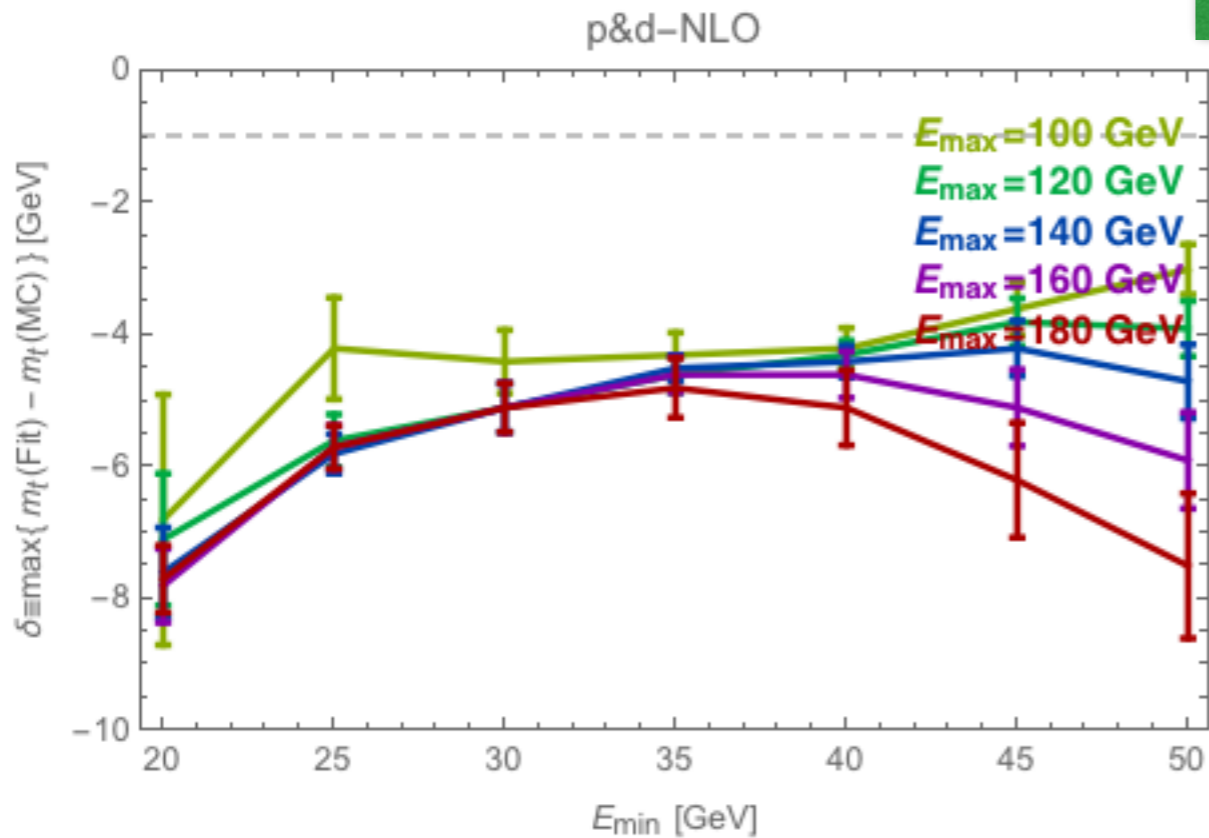


# Fit Variations p&d-NLO

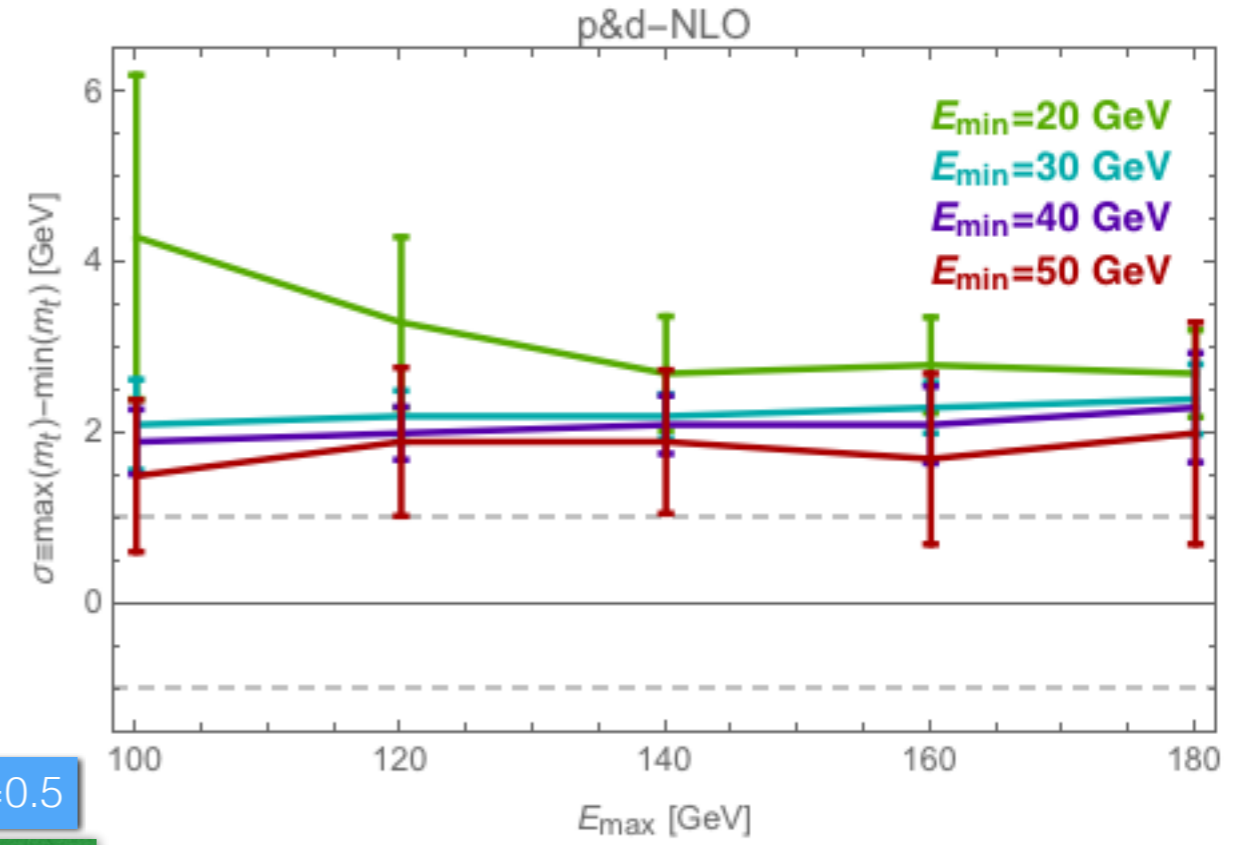
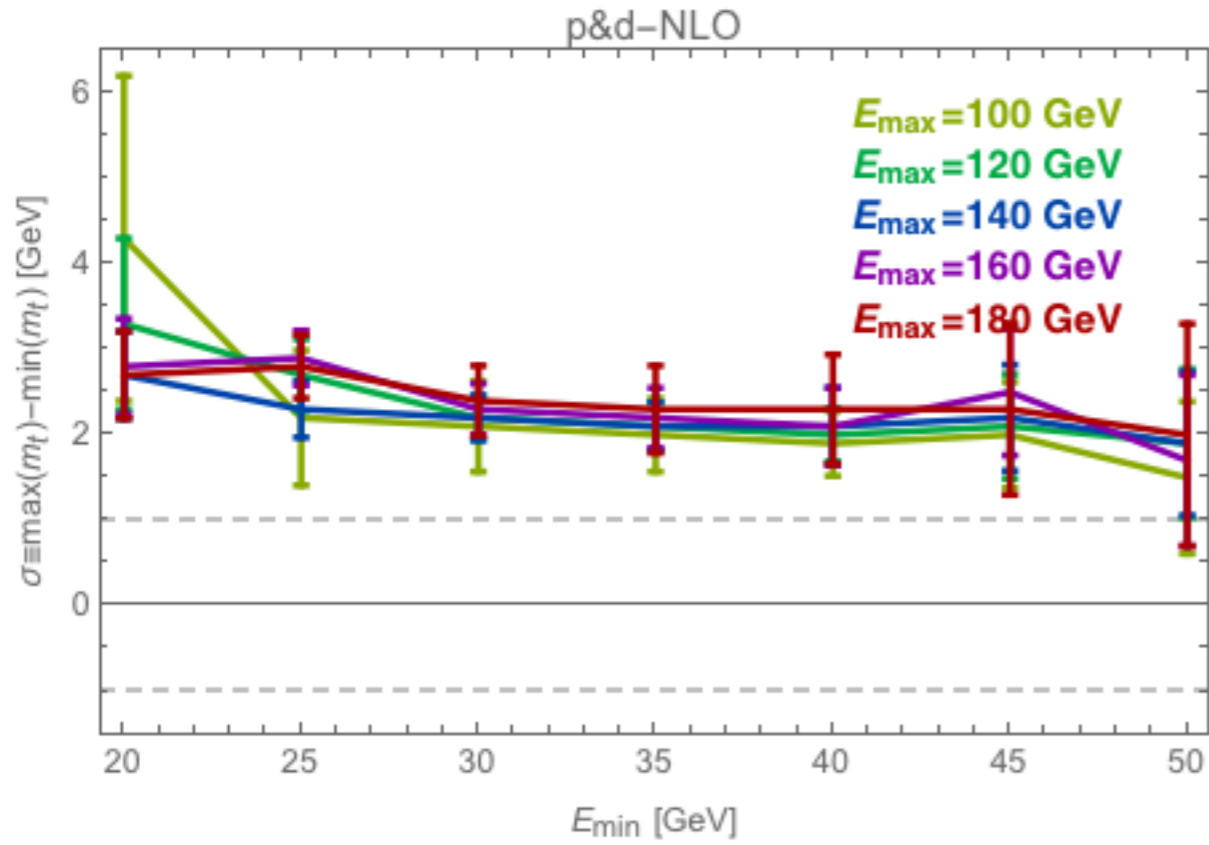


R=0.5

MSTW08

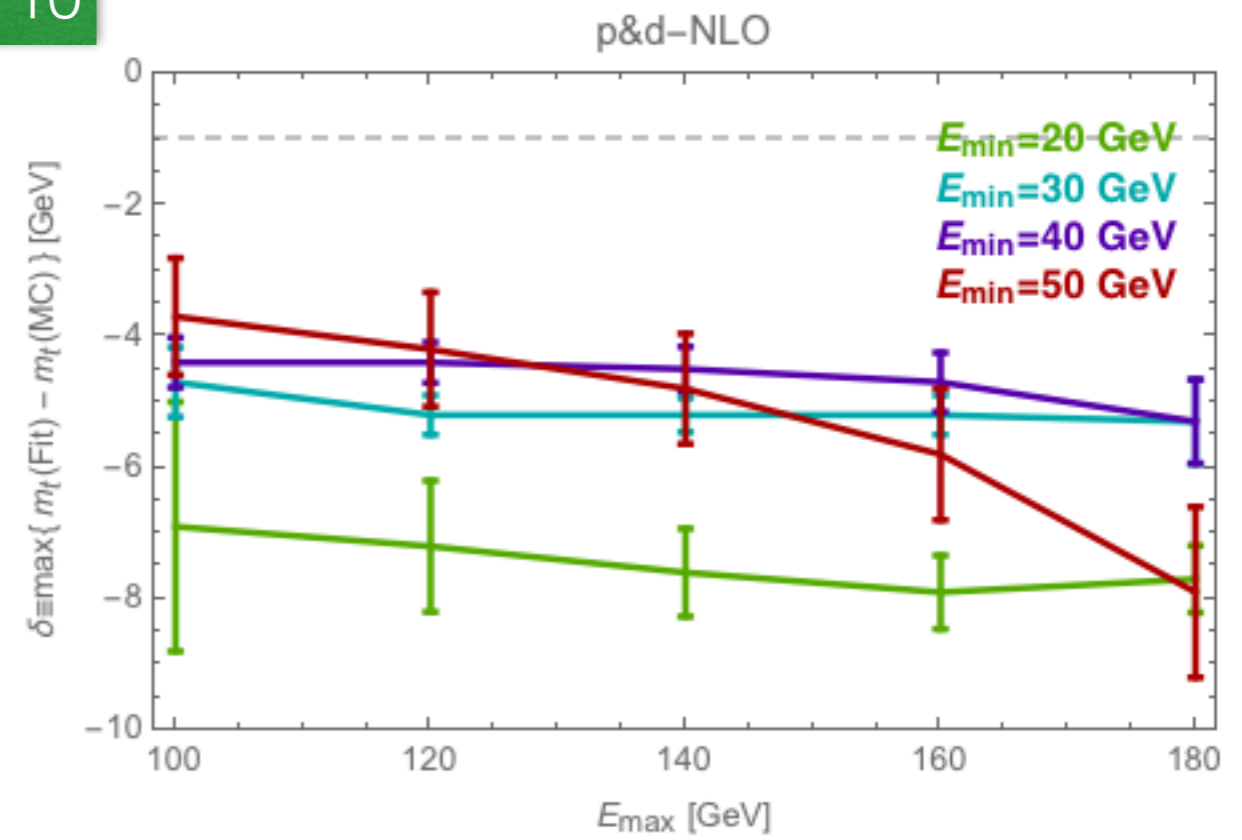
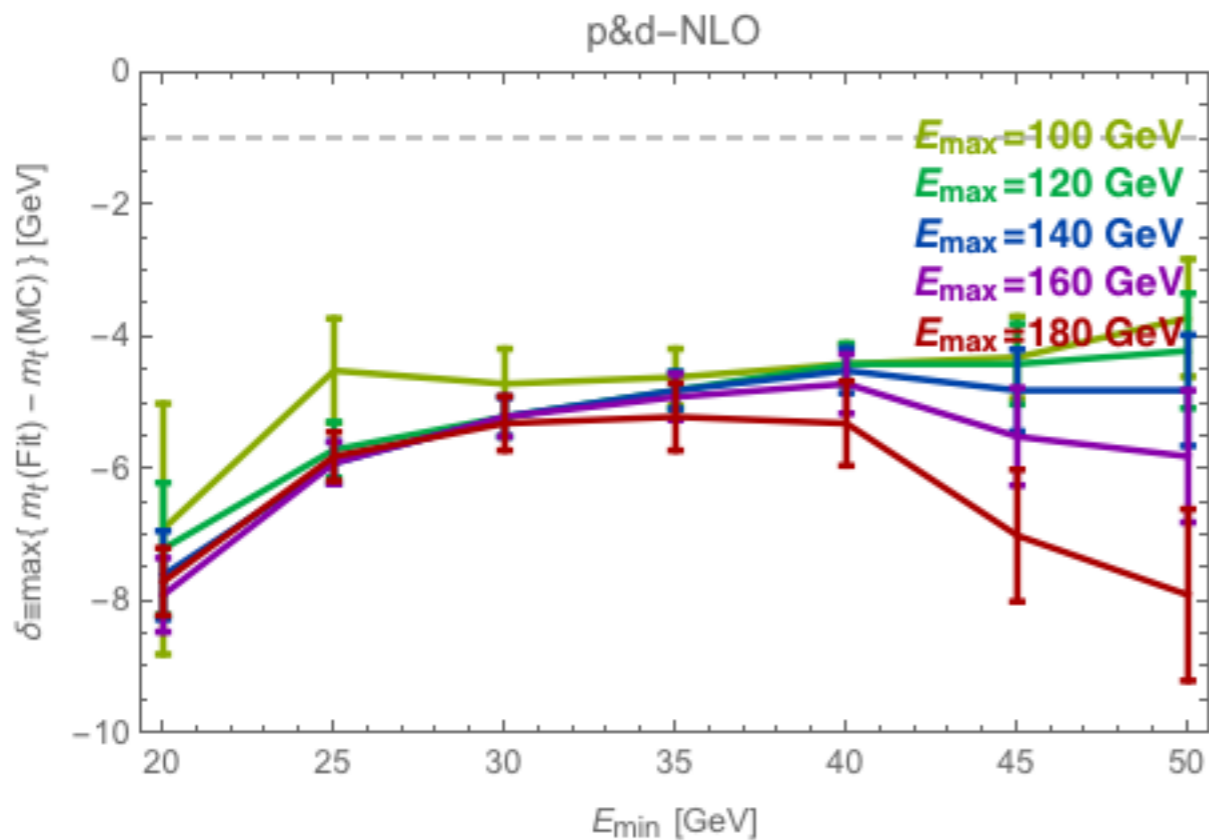


# Fit Variations p&d-NLO



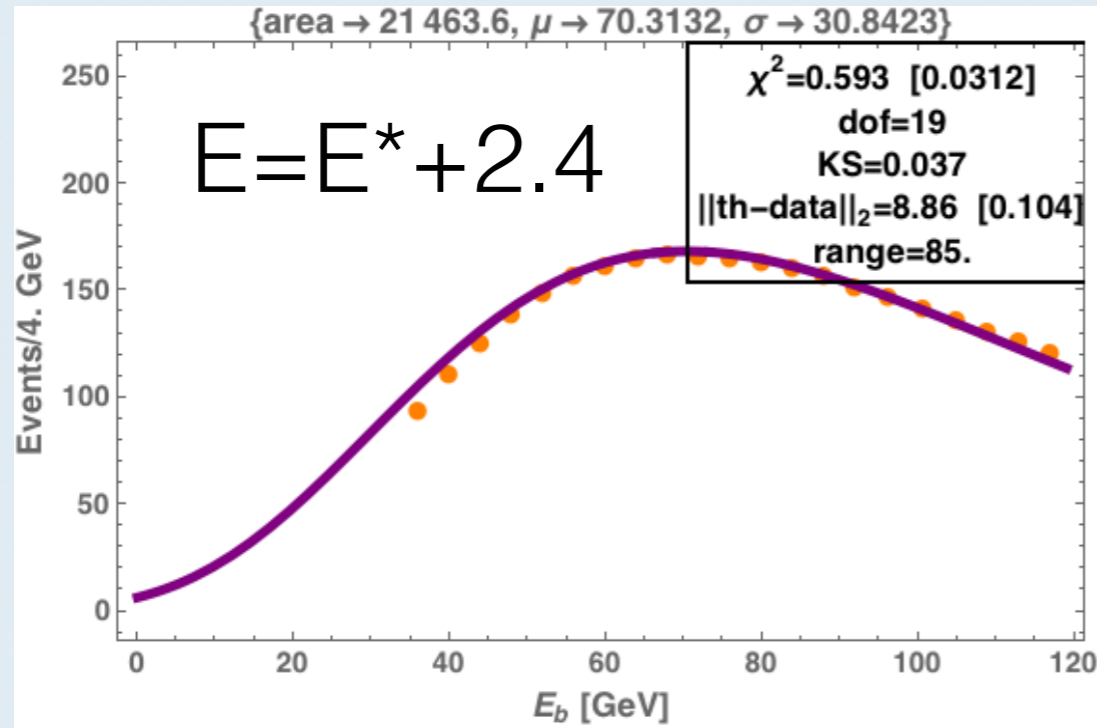
R=0.5

CT10

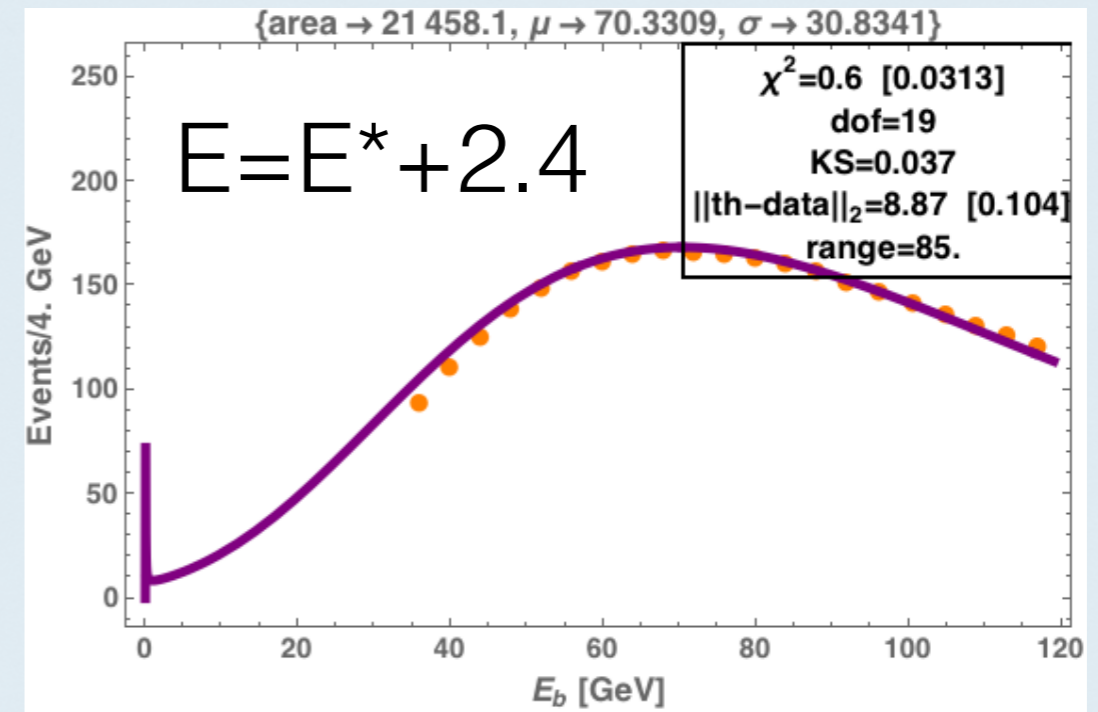


# LO MCFM fixed $\mu=m_{\text{top}}$ ( $E=67.9$ GeV)

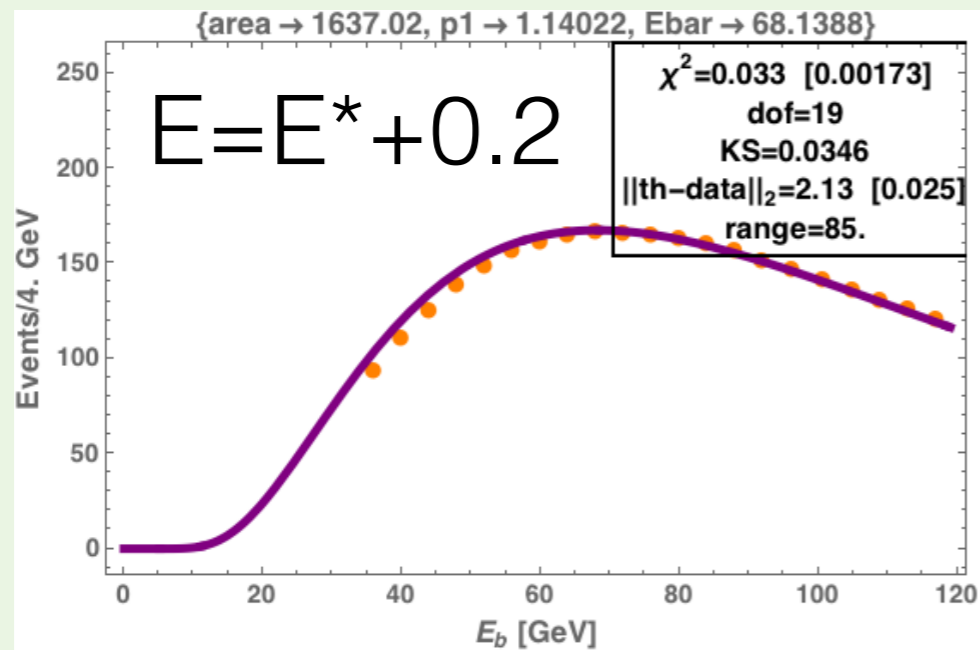
Moyal(x)



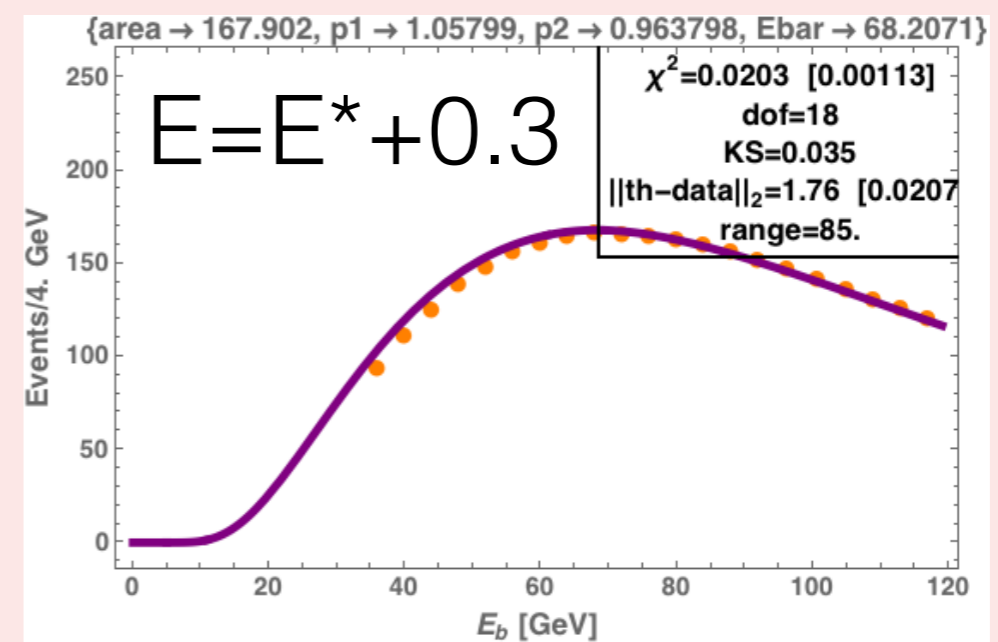
Moyal(1/x+x)



1par Exp(x+1/x)



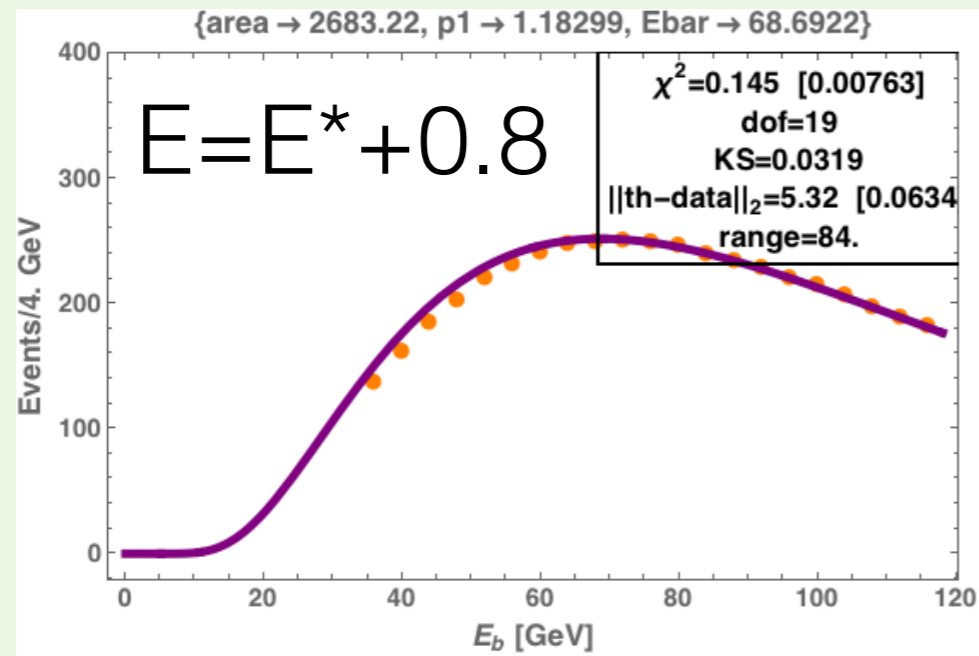
2pars Exp(x+1/x)



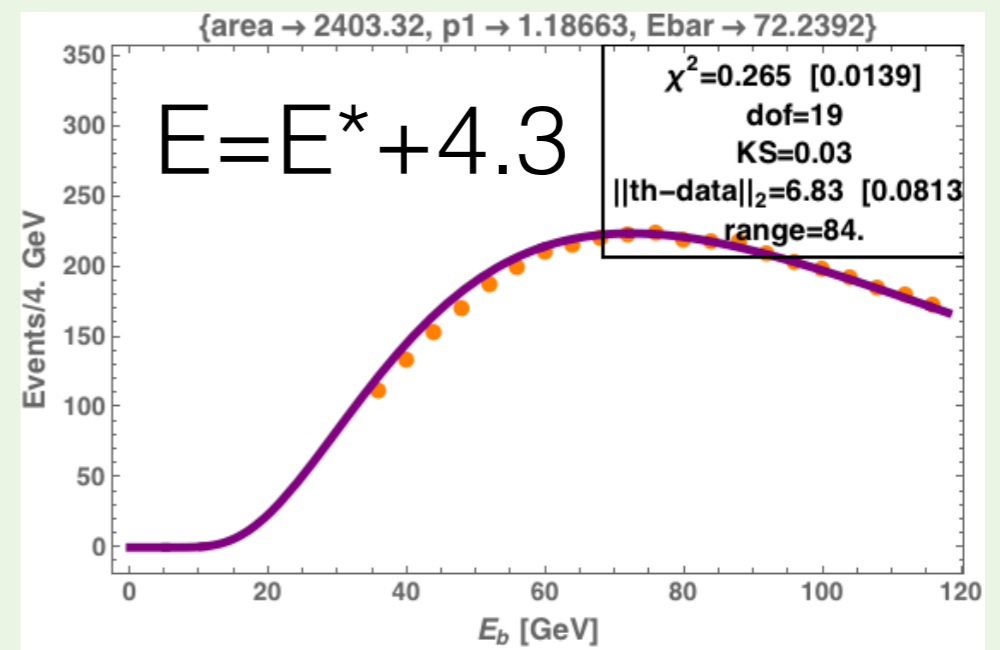
# pNLO MCFM fixed $\mu=m_{\text{top}}$ (E=67.9 GeV)

1par Exp(x+1/x)

R=0.5



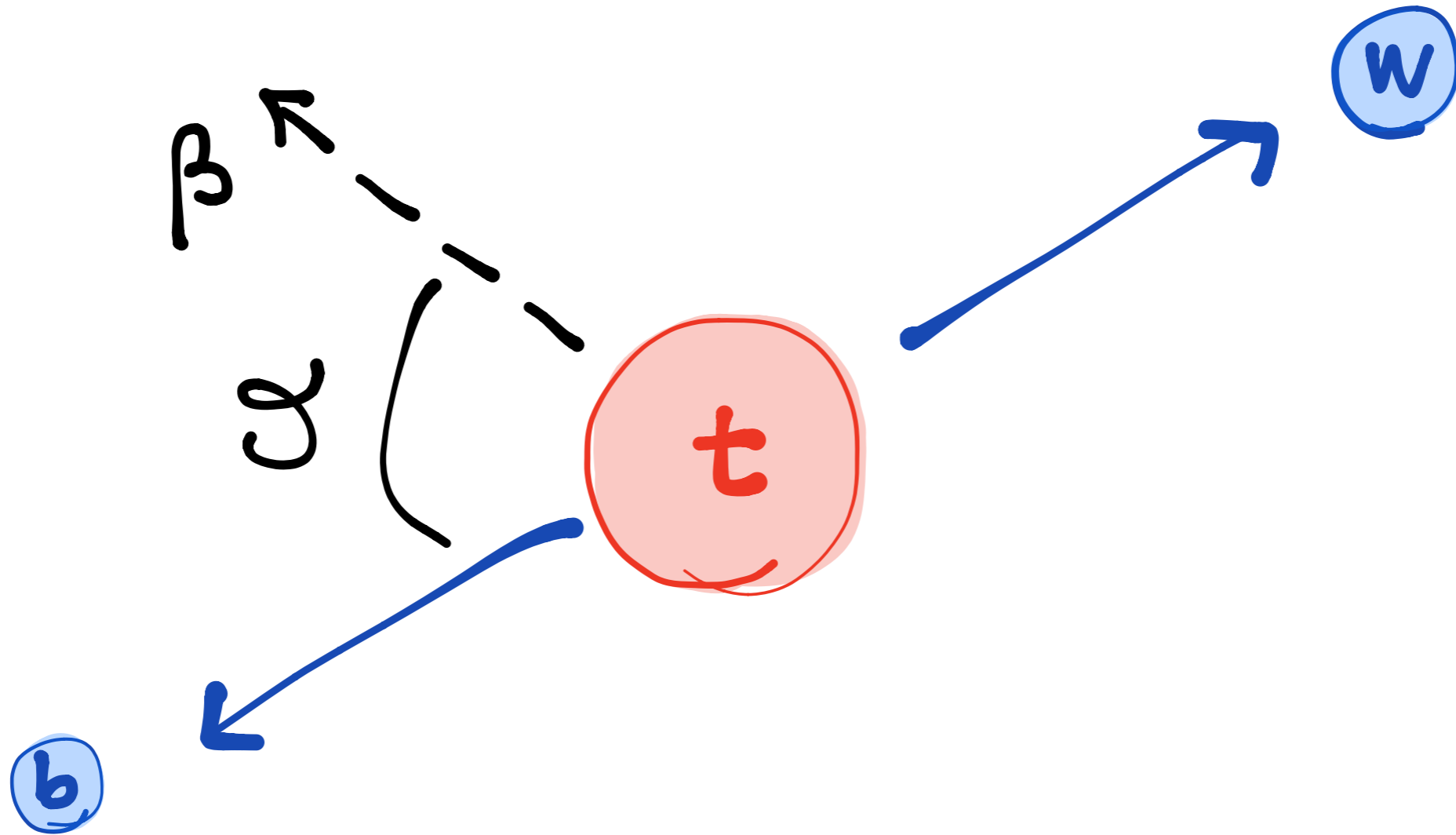
R=1.0



# A simple, yet subtle, invariance of the two body decay

*1209.0772 - Agashe, Franceschini and Kim*

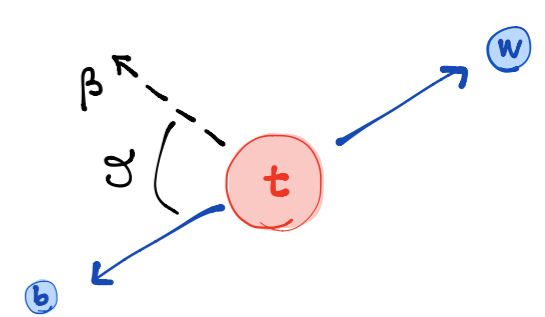




$$E_{\text{lab},b} = E_b^* \gamma + p_b^* \gamma \beta \cos \vartheta$$

Event-by-event we cannot tell anything

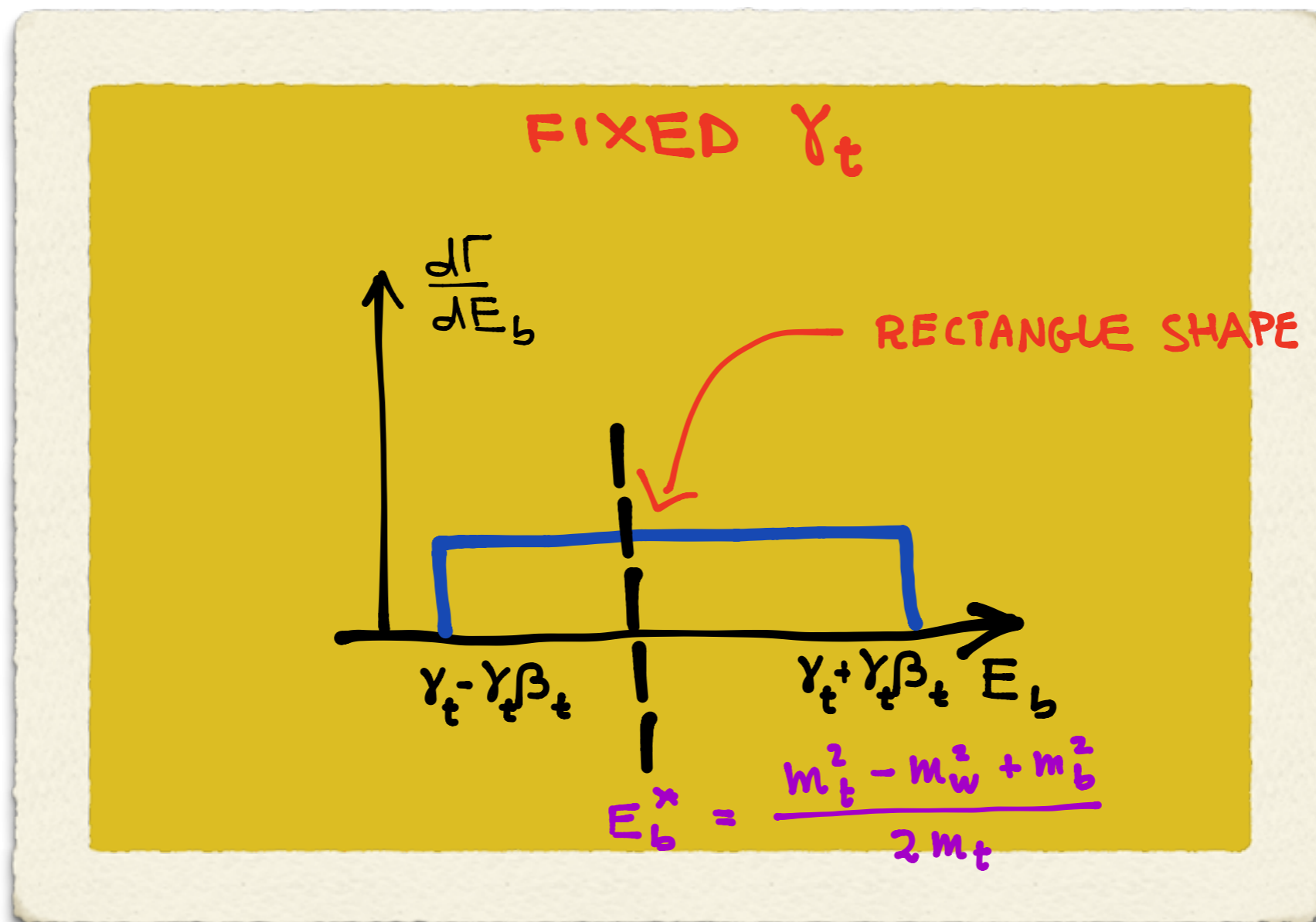
# Fixed top boost decay



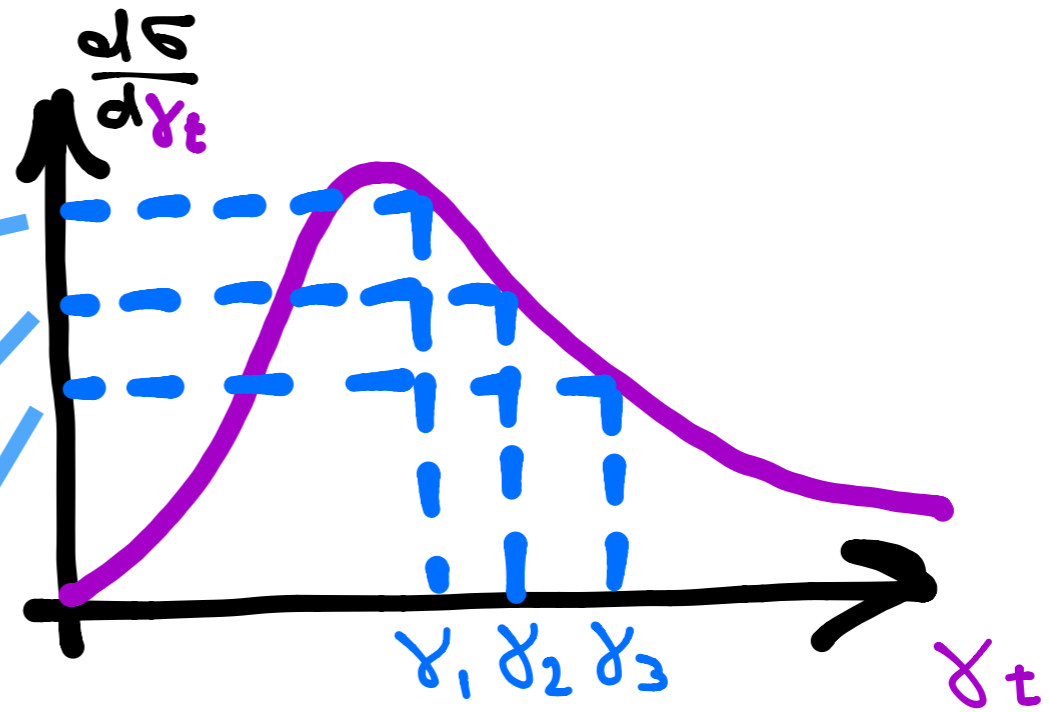
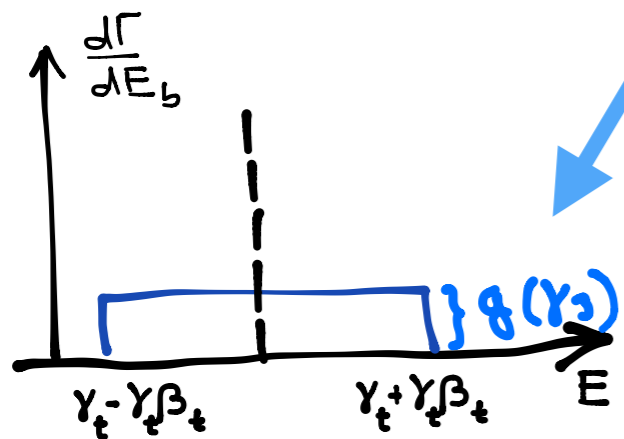
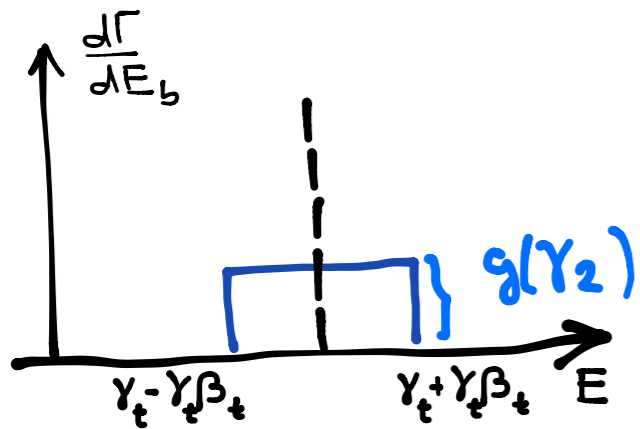
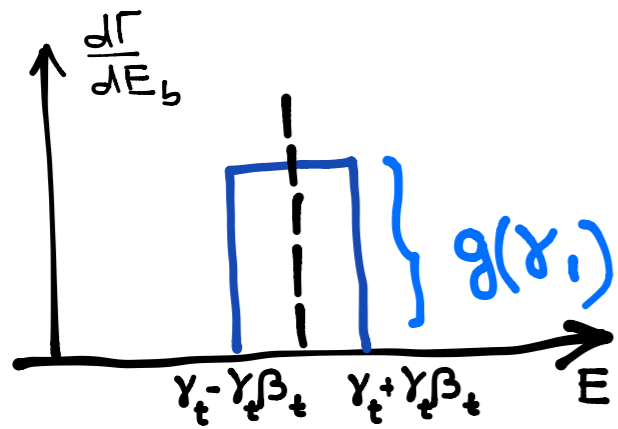
Massless  $b$ -quark (for now)

$$E_{lab,b} = E_b^* (\gamma + \gamma\beta \cos\vartheta)$$

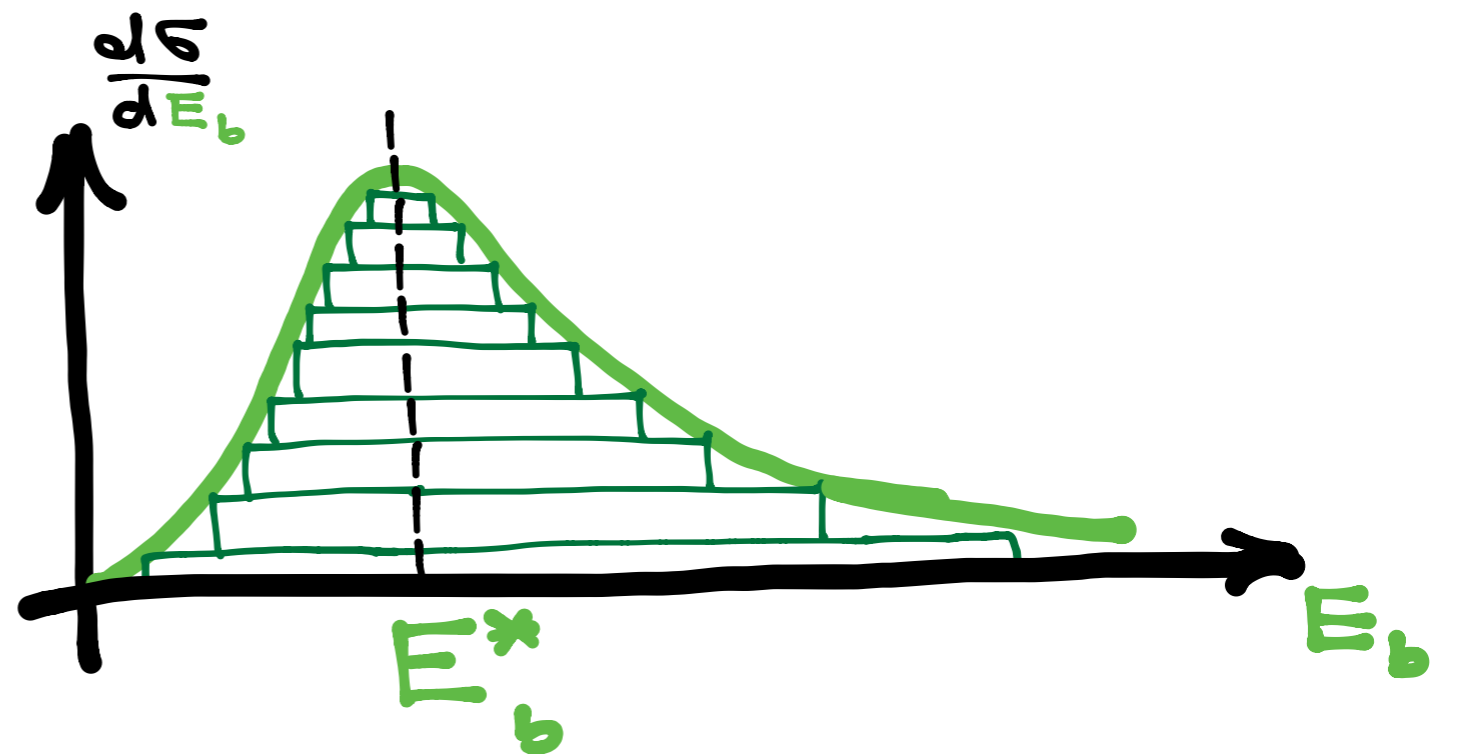
unpolarized top sample  $\rightarrow$   $\cos\theta$  is flat



# Summing over the top boosts



THE ENERGY DISTRIBUTION IN THE LAB IS THE SUM OF ALL THE RECTANGLES

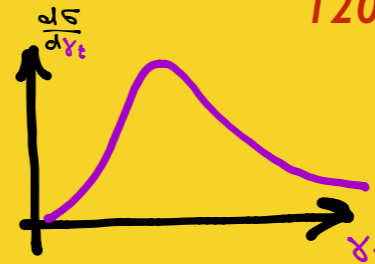


# Lab-frame energy distribution

1209.0772 - Agashe, Franceschini and Kim

also Stecker 1971

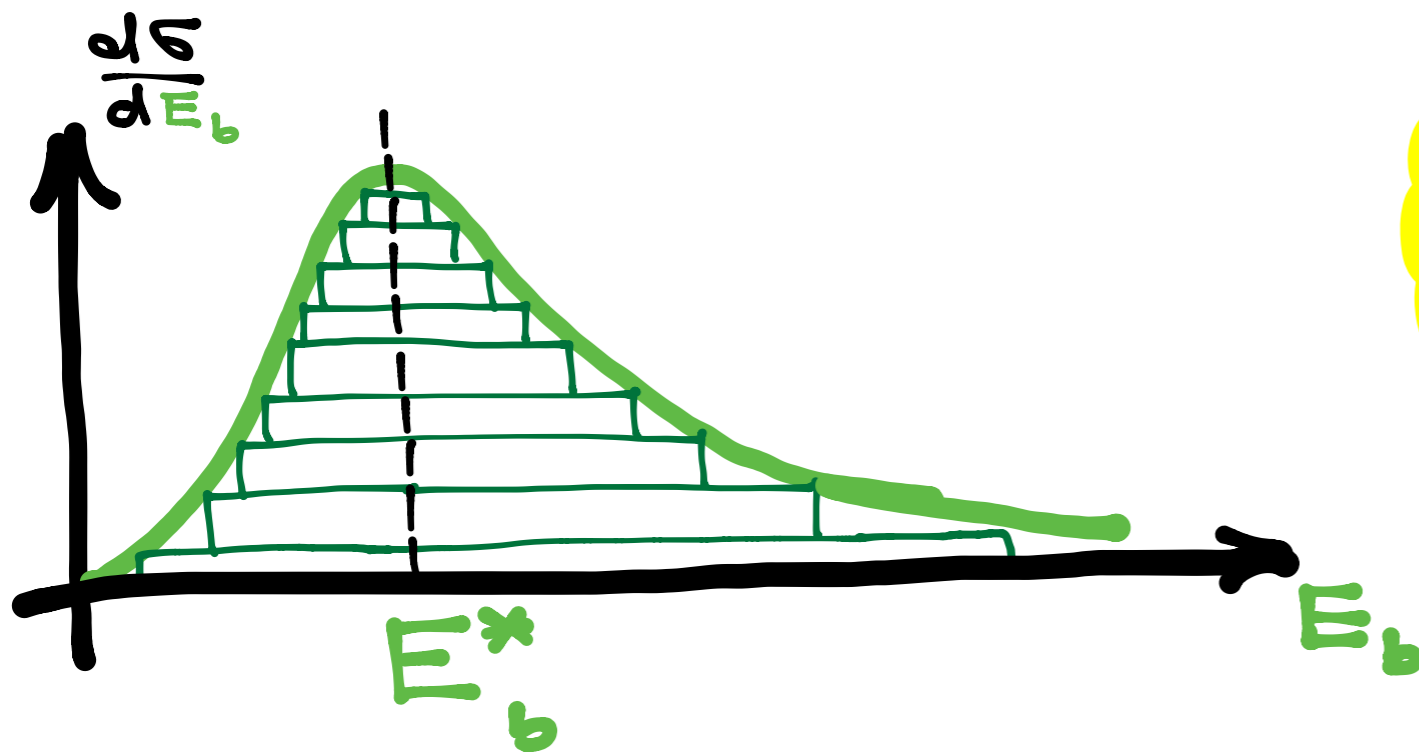
for any top boost distribution



the peak:

- is the same as in the rest frame
- encodes invariant

$$E_b^* = \frac{m_t^2 - m_w^2 + m_b^2}{2m_t}$$



THE FRAME-DEPENDENT  
ENERGY DISTRIBUTION ENCODES  
THE INVARIANT  $E_b^*$  IN A  
VERY SIMPLE WAY

There is no difference when the b-mass is taken into account provided  $\gamma_{top} < 500$

# On mass measurements

- Lorentz invariants
- resonance reconstruction

# Ideal mass measurements



Lorentz invariant

insensitive to:

- Parton Distribution Functions
- Production Mode (qq or gg, SM or BSM, ISR, ...)

beware of radiation for precision measurement

# Less ideal mass measurements

One particle is just lost



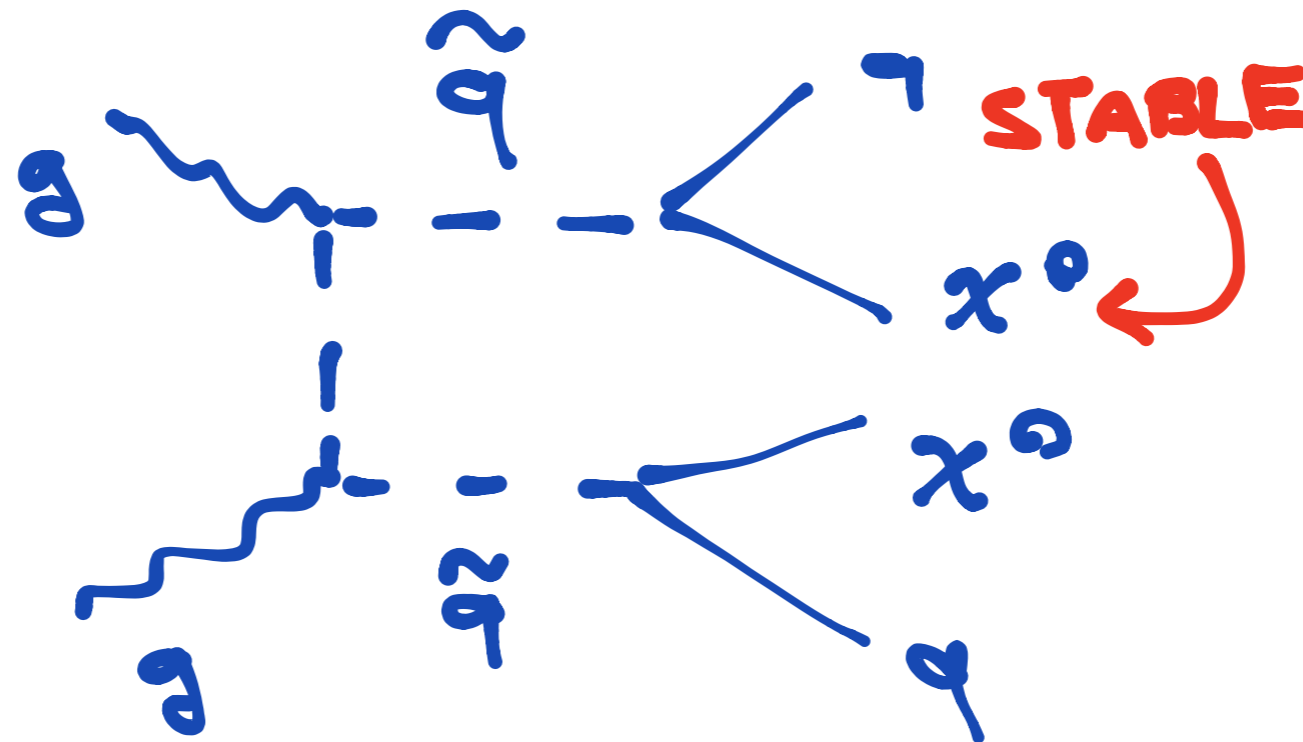
Need to come up with a trick

for example:

- Transverse Mass (use  $m_{ET}$ )
- $p_T$  (nuisances are back:  $qq$  or  $gg$ , SM or BSM, ISR, ...)

# ... and it can get worse

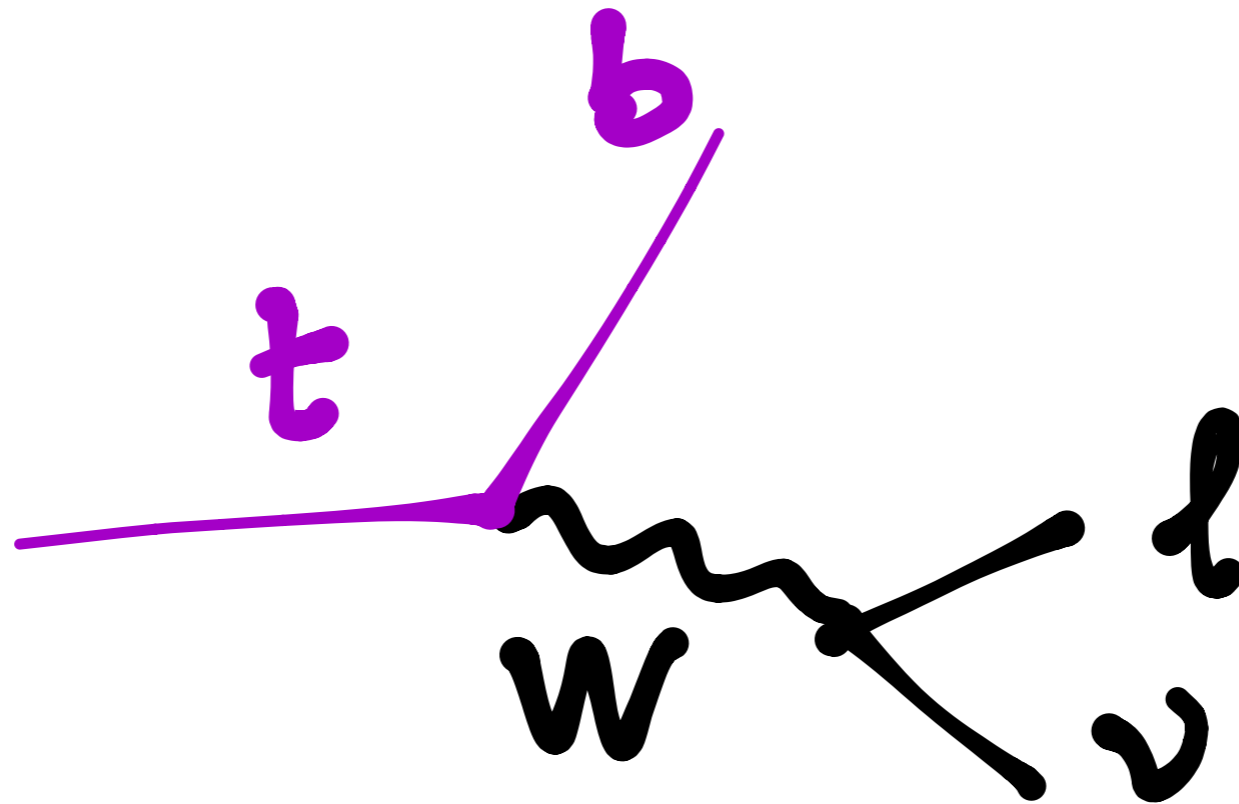
any BSM with some sort of Matter Parity (e.g. RPC SUSY)



can we make a mass measurement without ever mentioning the unobservable particle  $\chi$ ?

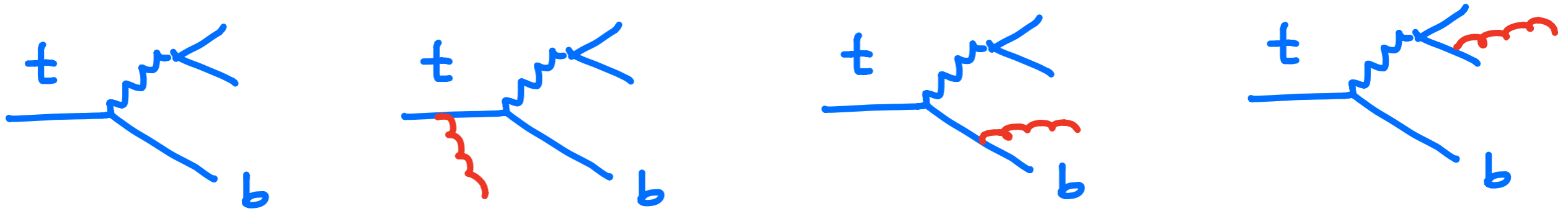


“useful” top is semi-invisible

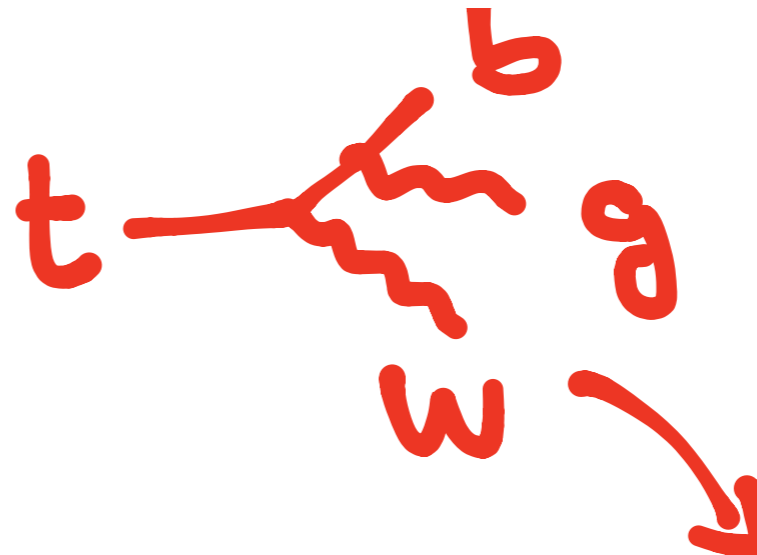
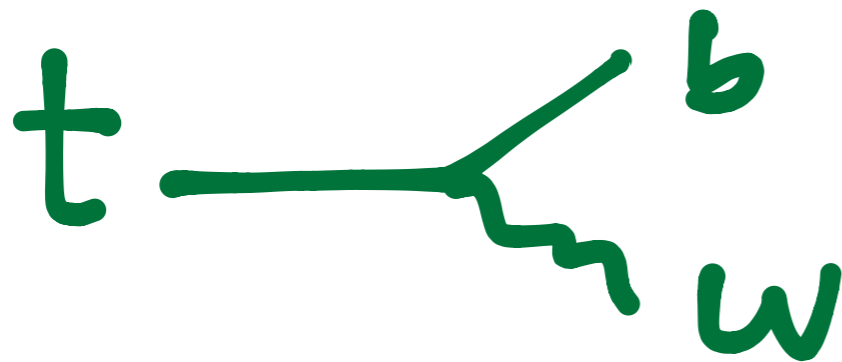


can we make a mass measurement without ever mentioning the unobservable particle  $W$ ?

# To reconstruct or not to reconstruct?



top quark reconstruction is entangled with some picture of the kinematics (fixed order?)



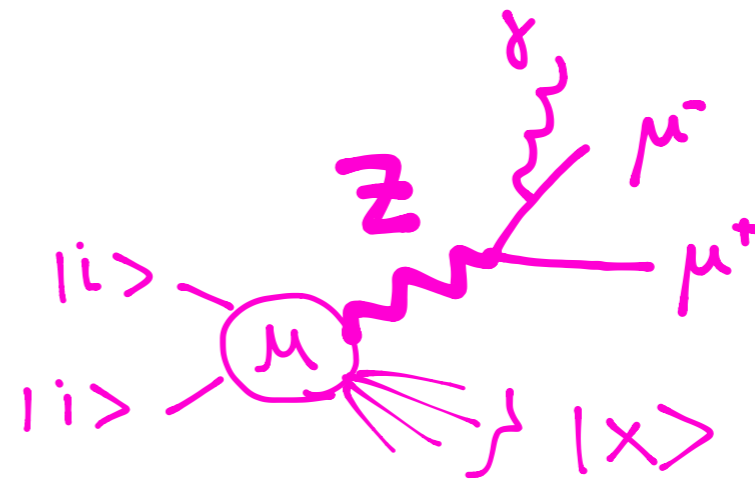
NLO (decay)

NLO+PS in 1412.1828

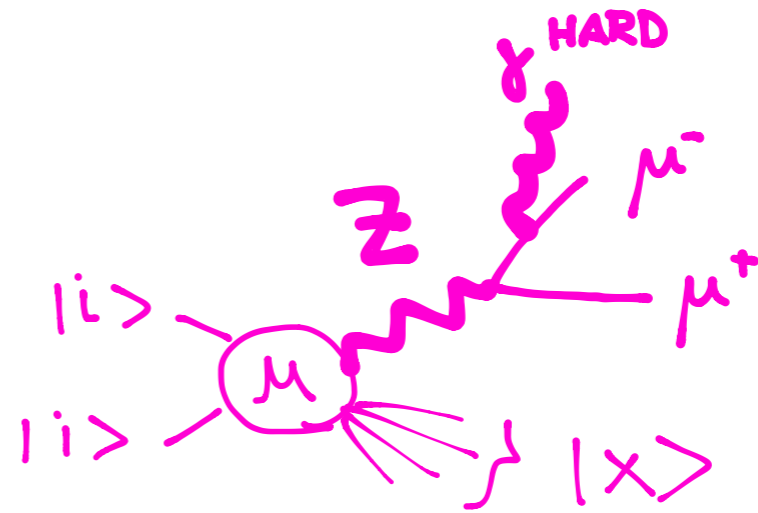
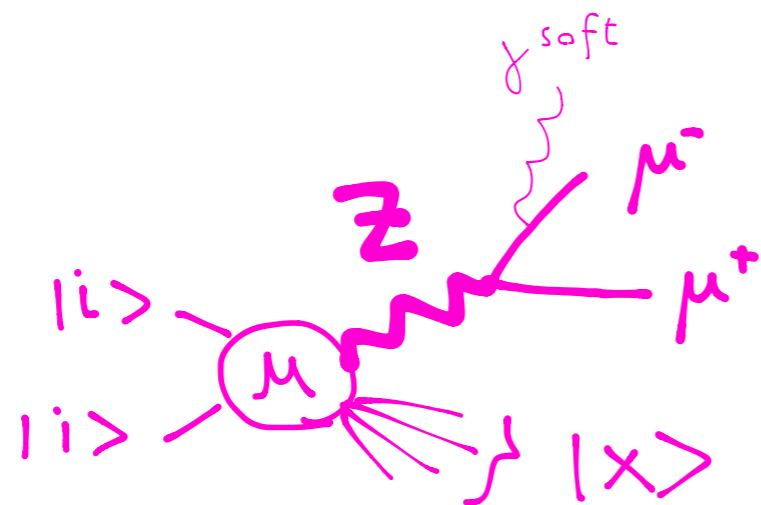
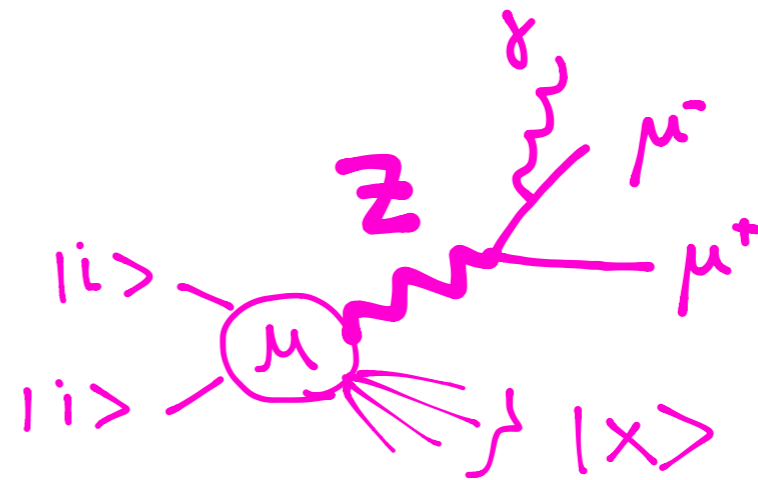
# To reconstruct or not to reconstruct?



# To reconstruct or not to reconstruct?



# To reconstruct or not to reconstruct?



# To reconstruct or not to reconstruct?



does (not) distinguish where  
the final state came from (t, t\*, bW, bWg, bqqg)

need (not) to define the top

might (not) depend on the production mechanism

...