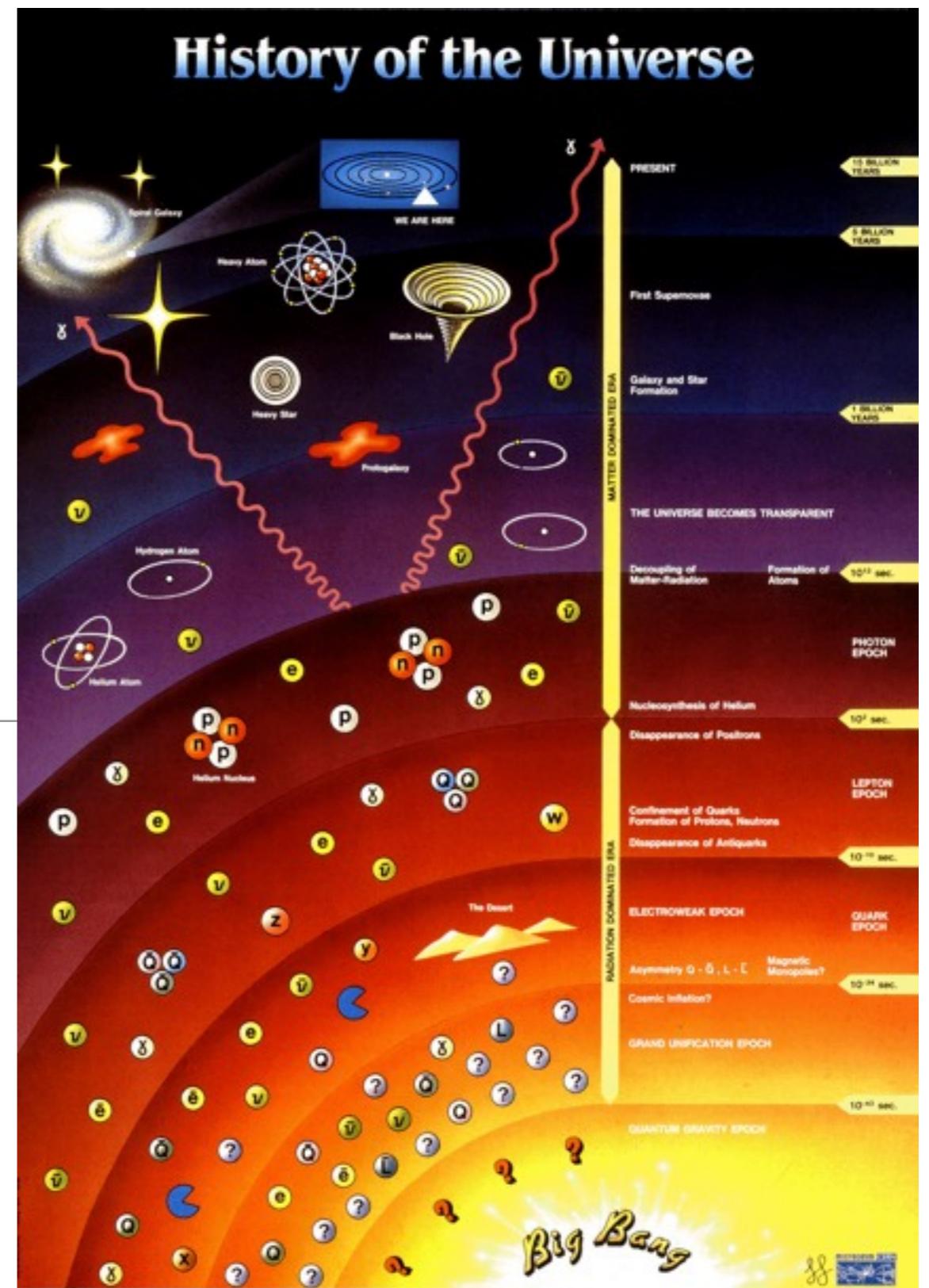


Cosmology

Kerstin Kunze

IUFFyM and Universidad de Salamanca



Overview

- Cosmology 1: The homogeneous universe

★The expansion stages of the universe

- Cosmology 2: The inhomogeneous universe

★The cosmic microwave background

★Large scale structure

- Cosmology 3:

★Dark matter

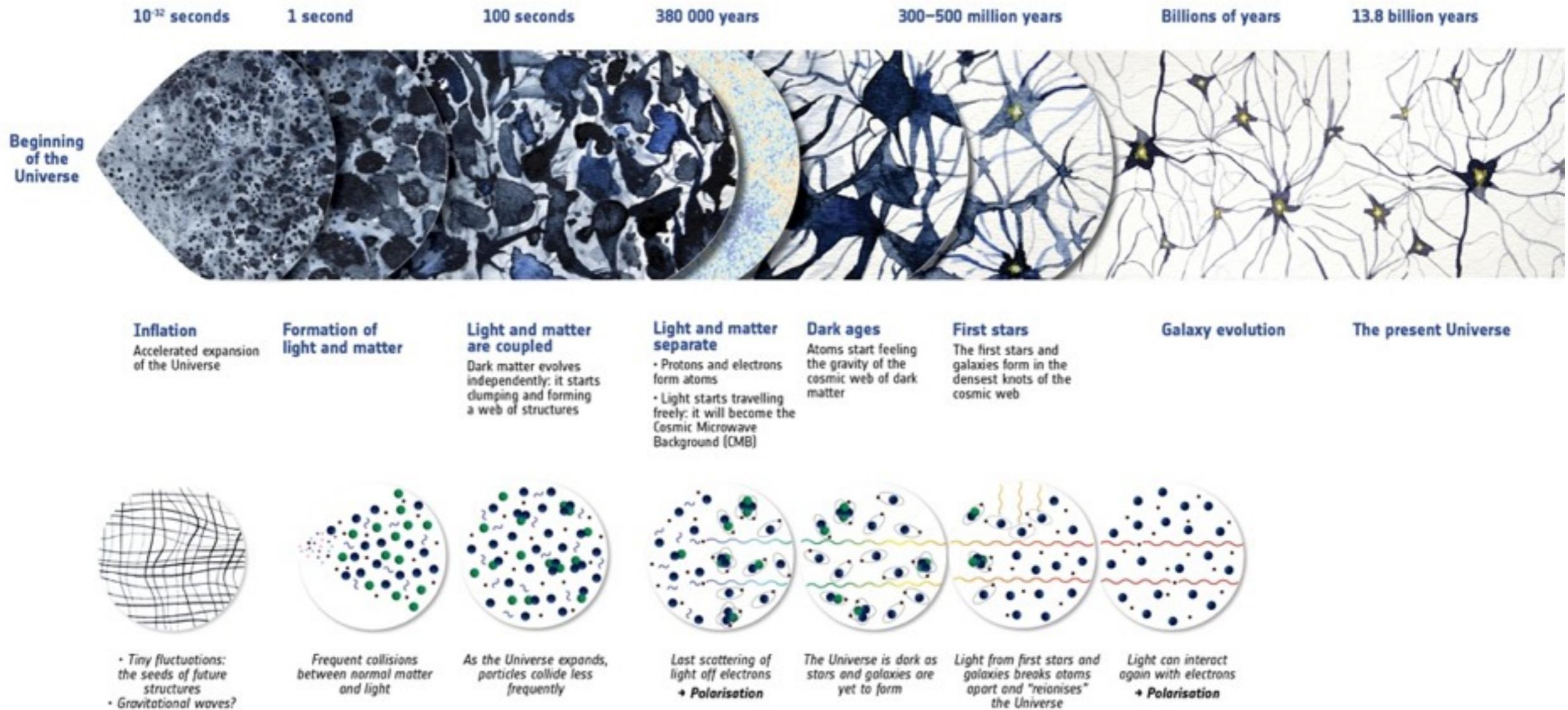
★Dark energy

Cosmology 2

Cosmology 2



→ COSMIC HISTORY

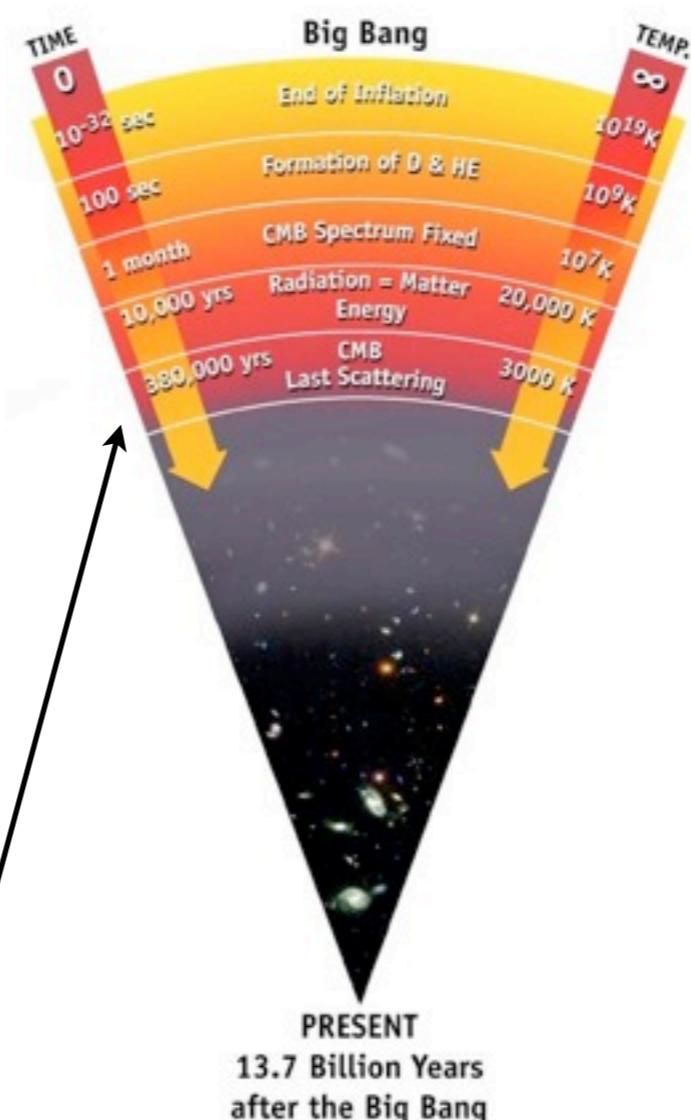
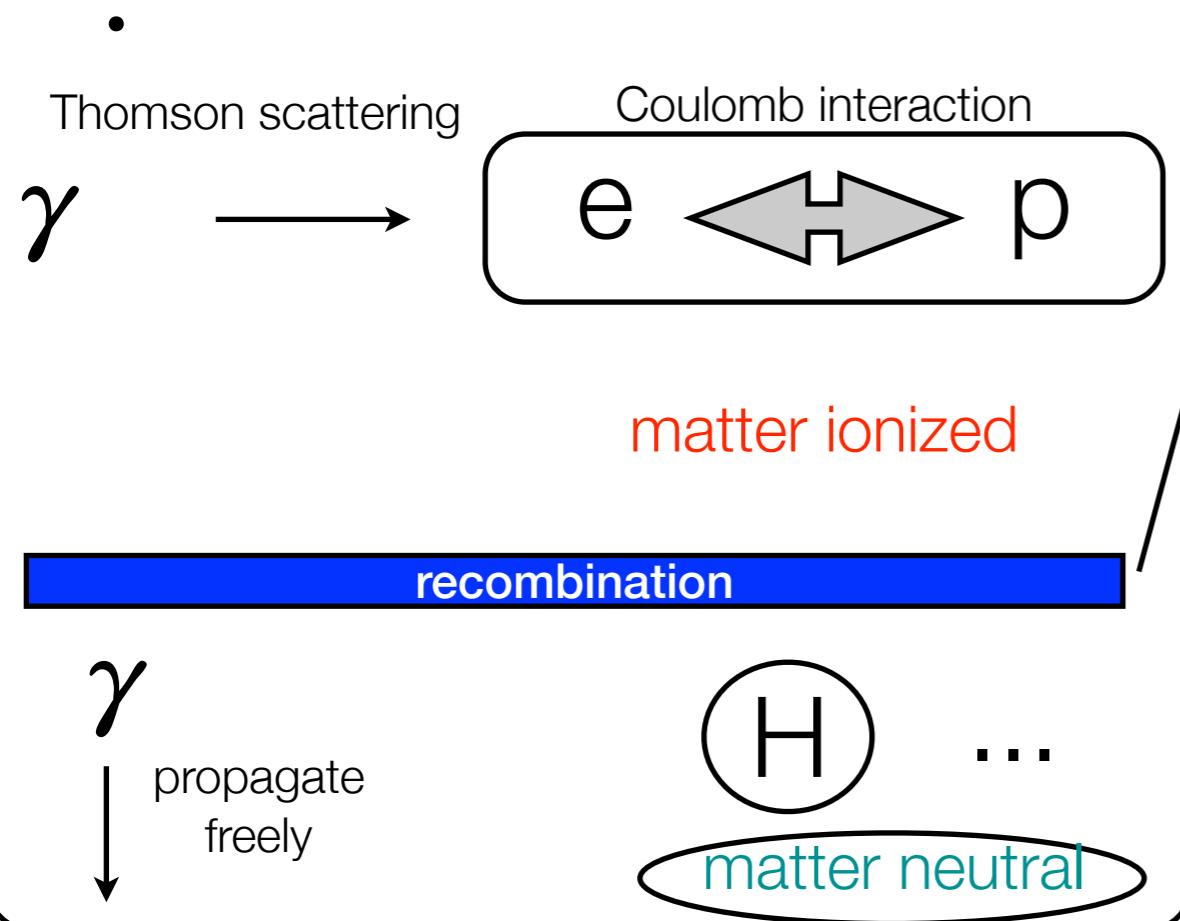


Cosmology 2

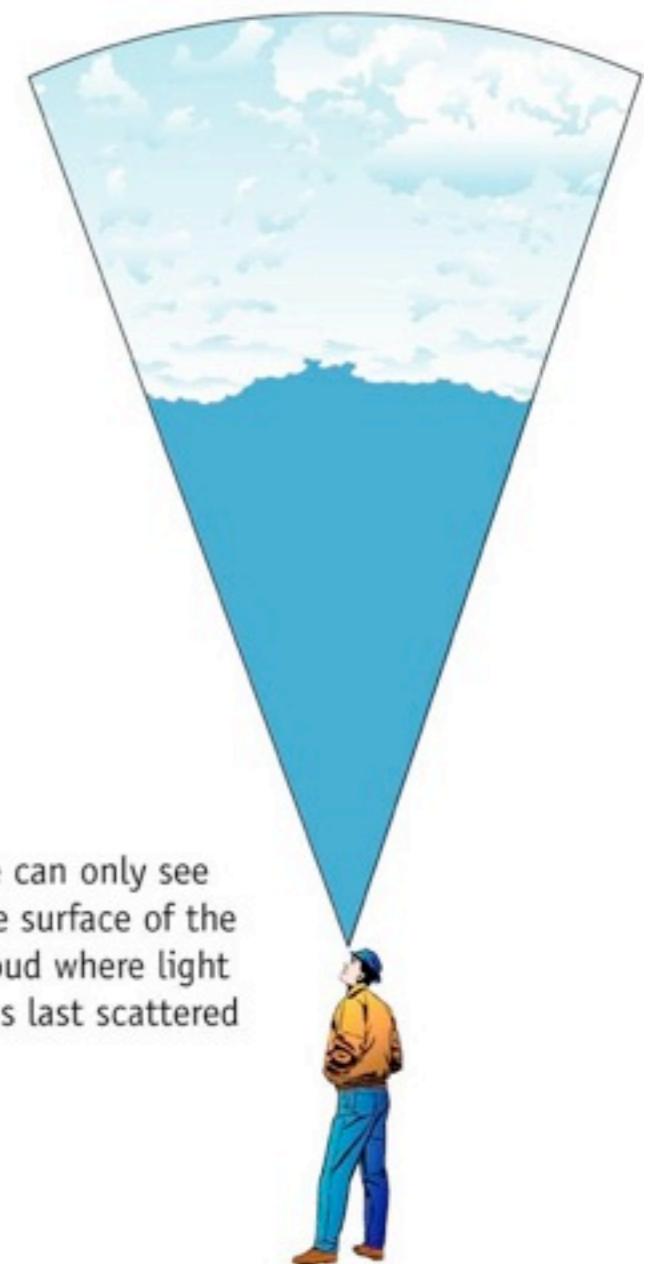
- **The cosmic microwave background**

Before decoupling

tight coupling: baryon-photon fluid

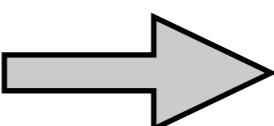


The cosmic microwave background Radiation's "surface of last scatter" is analogous to the light coming through the clouds to our eye on a cloudy day.



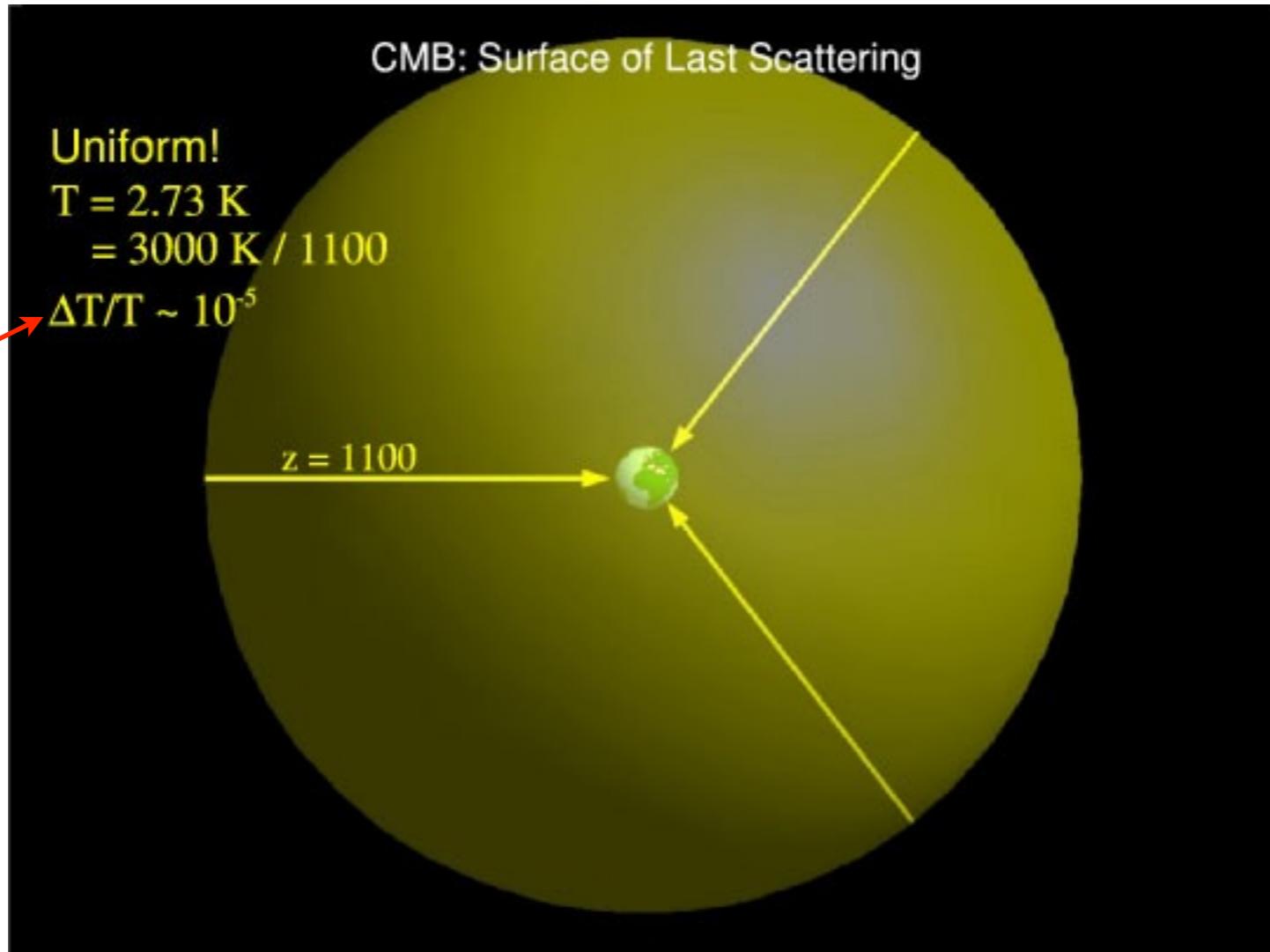
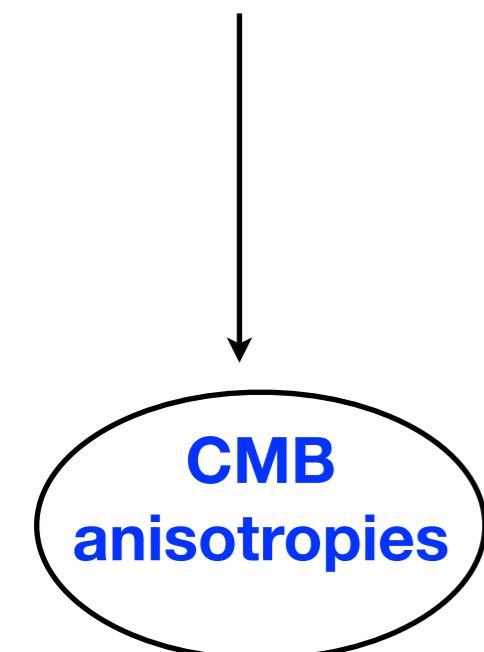
NASA

Cosmology 2

- Recombination would occur much earlier than about 4000K if just depended on $13.6 \text{ eV} = kT$  160000 K.
- Subtleties of recombination are due to two-photon decay process and interactions drop out of Saha equilibrium.
- **Recombination:** epoch when electrons join nuclei. Define at ionization fraction $X_e=0.1$.
 - free electrons become rare
 - Saha equation not applicable anymore
- **Decoupling:** epoch after which photons will not scatter again.
 - residual ionization fraction of order 0.001
 - define as epoch when the duration of photon mean free path equals the age of the universe

Cosmology 2

Small
temperature
perturbations



W. Kinney

Cosmology 2

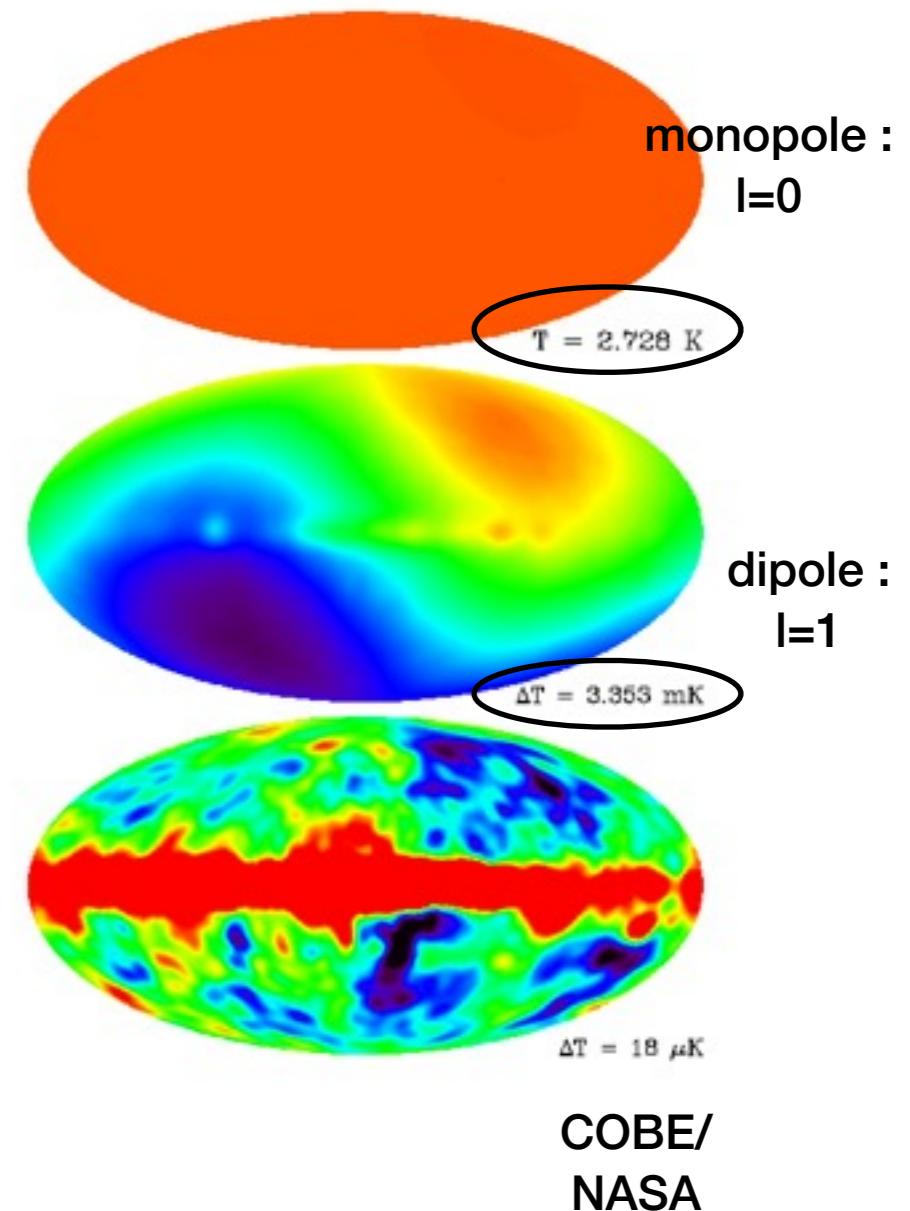
- Expand temperature fluctuations in spherical harmonics:

$$\Theta(\hat{\mathbf{n}}) = \frac{\Delta T}{T} = \sum_{\ell m} \Theta_{\ell m} Y_{\ell m}(\hat{\mathbf{n}})$$

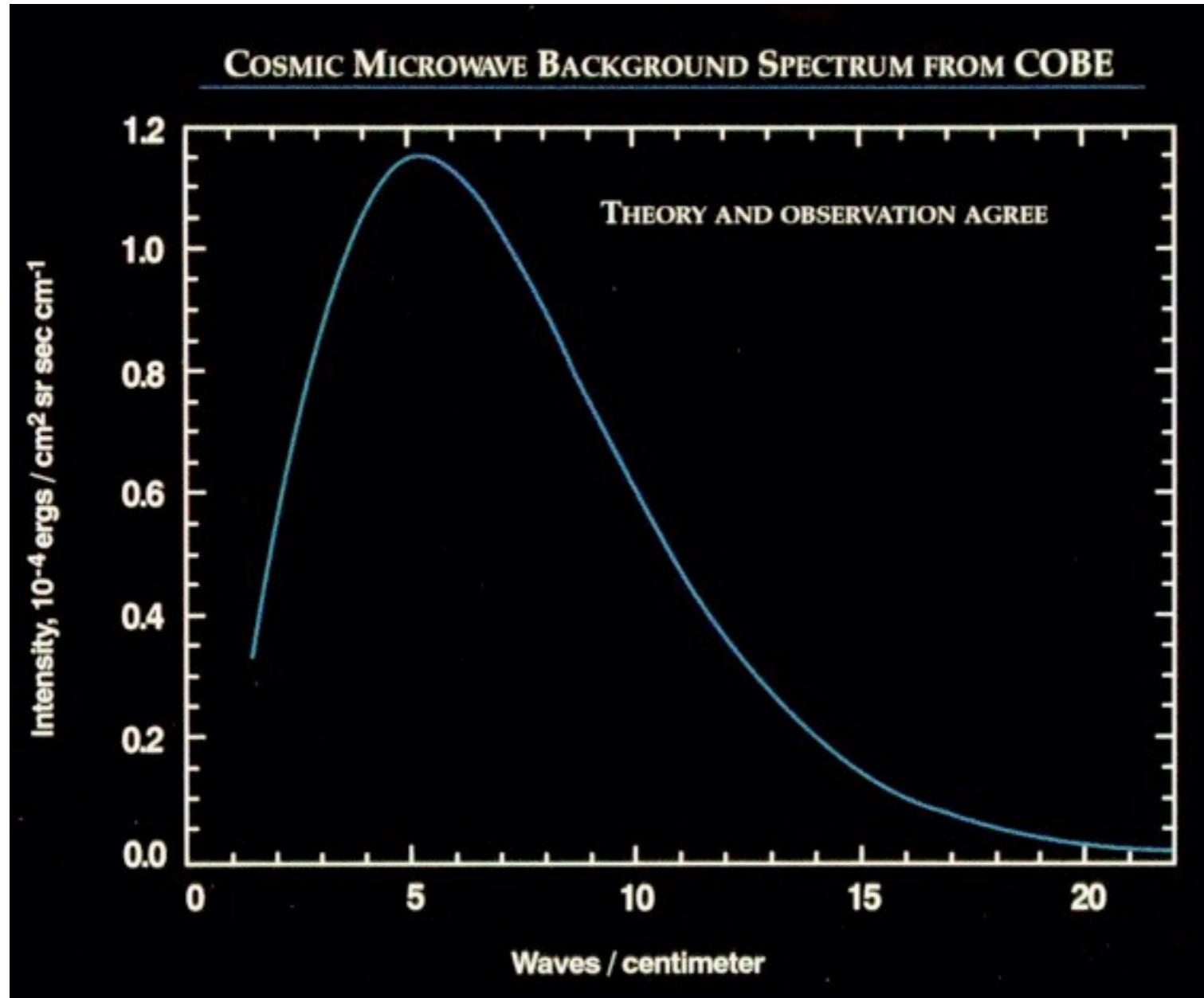
$$\langle \Theta_{\ell m}^* \Theta_{\ell' m'} \rangle = \delta_{\ell' \ell} \delta_{m m'} C_{\ell}$$

dipole component subtracted

**DMR data from the
53 GHz band**



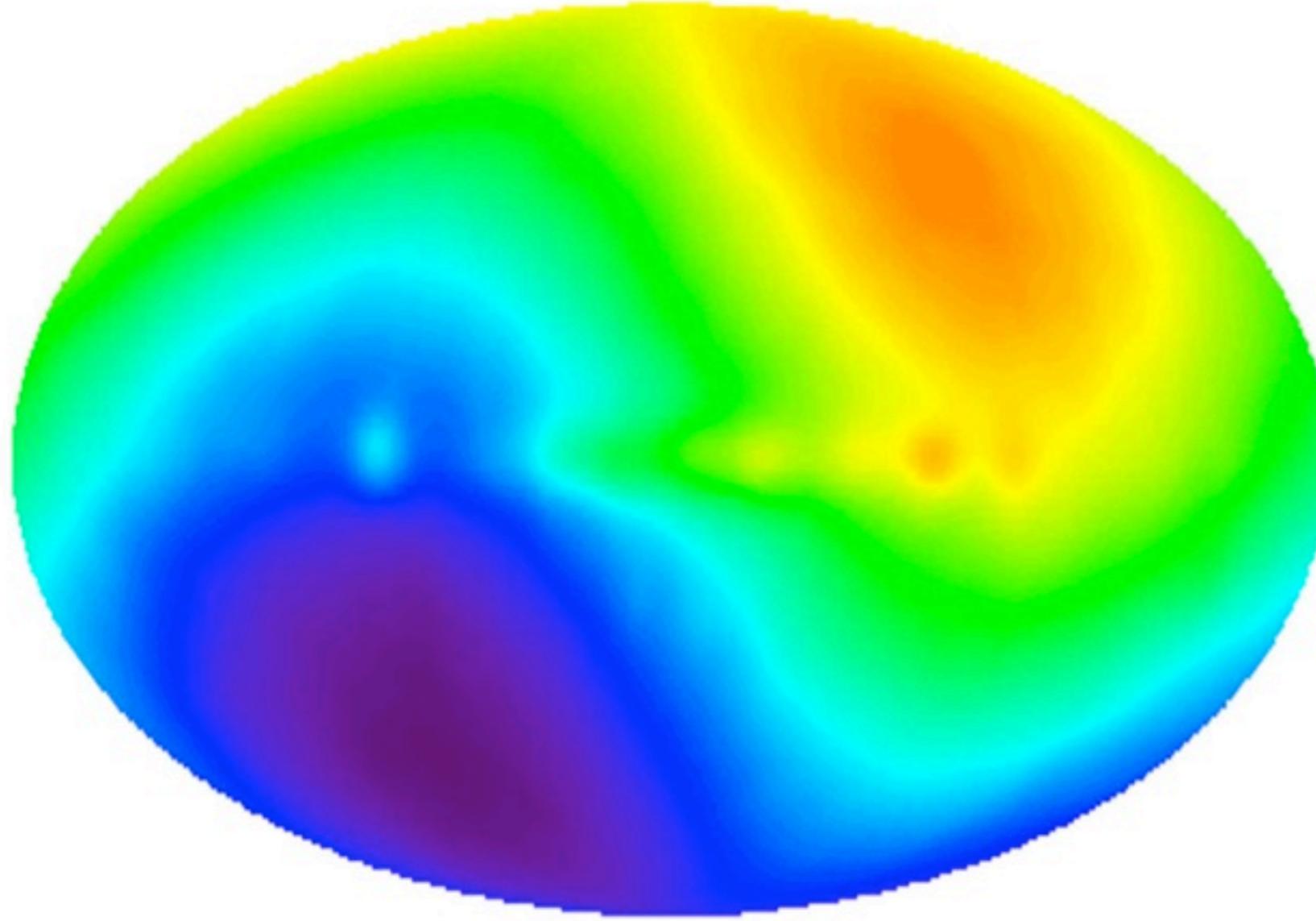
Cosmology 2



monopole

$$T_\gamma = 2.7255 \pm 0.0006 \text{K}$$

Cosmology 2



Dipole

$$T_\gamma = 3.355 \pm 0.008 \text{ mK}$$

Astronomy Picture of the Day

2014 June 15

CMB Dipole: Speeding Through the Universe

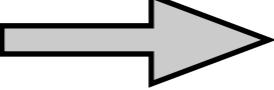
Image Credit: [DMR](#), [COBE](#), [NASA](#), Four-Year Sky Map

Explanation: Our [Earth](#) is not at rest. The Earth moves around the [Sun](#). The Sun orbits the center of the [Milky Way Galaxy](#). The Milky Way Galaxy orbits in the [Local Group of Galaxies](#). The Local Group falls toward the [Virgo Cluster of Galaxies](#). But these speeds are less than the speed that all of these [objects together](#) move relative to the [cosmic microwave background radiation](#) (CMBR). In the [above all-sky map](#) from the [COBE satellite](#), radiation in the Earth's direction of motion appears [blueshifted](#) and hence hotter, while [radiation](#) on the opposite side of the sky is [redshifted](#) and colder. The [map](#) indicates that the [Local Group](#) moves at about 600 kilometers per second relative to this [primordial radiation](#). This [high speed](#) was initially unexpected and its magnitude is still unexplained. [Why are we moving so fast?](#) [What is out there?](#)

Cosmology 2

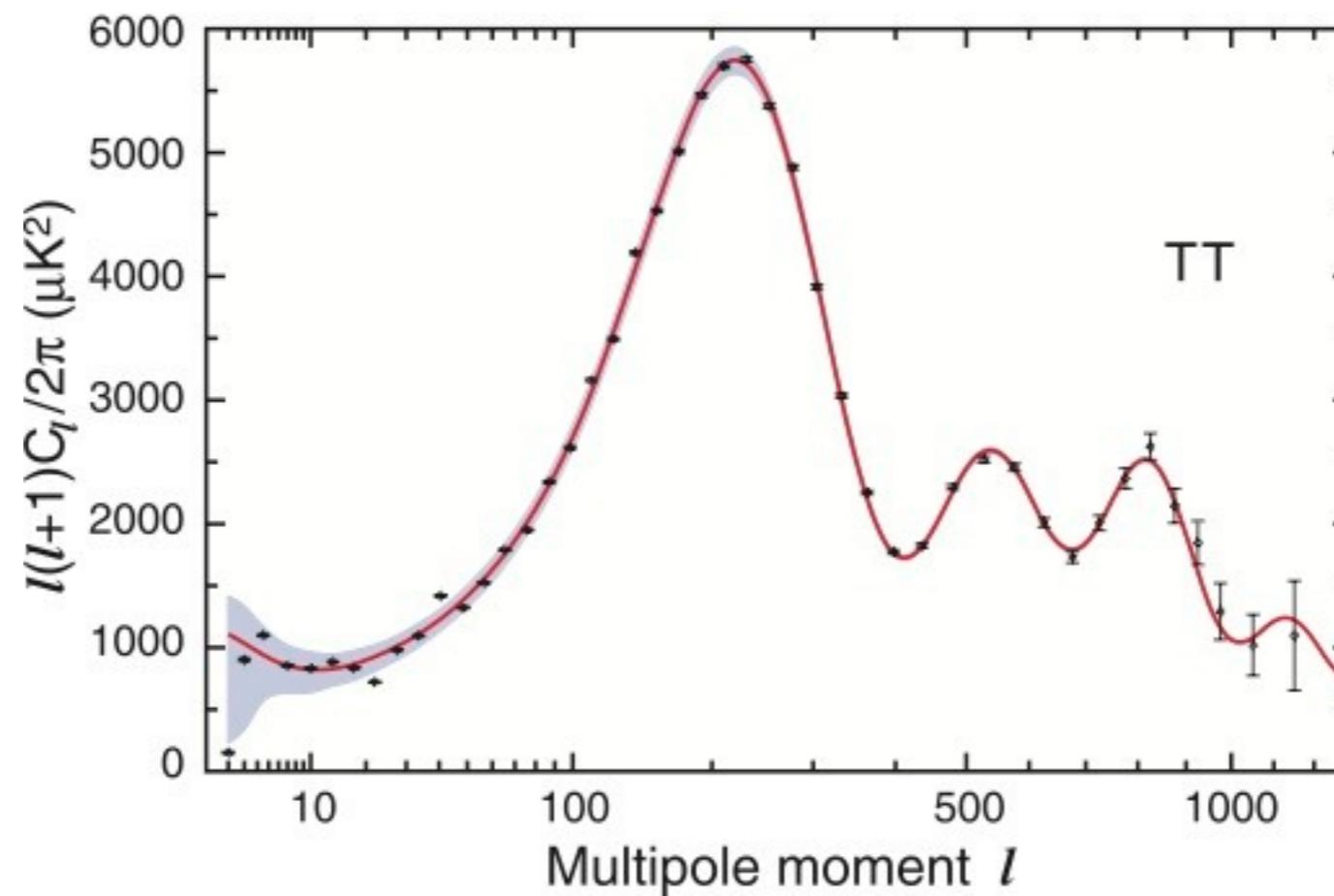
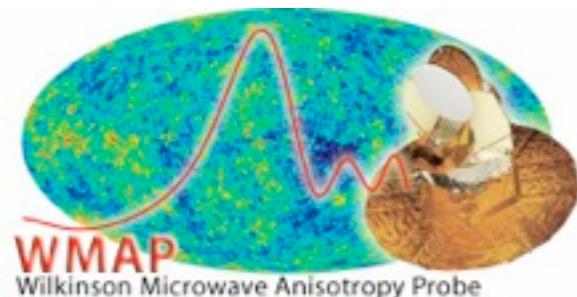
- Dipole (PDG 2014)
- Result of Doppler shift caused by motion of solar system relative to the nearly isotropic black body field. The motion of the observer with velocity $\beta = v/c$ relative to an isotropic Planck radiation field of temperature T_0 produces a Doppler shifted temperature pattern
$$T(\theta) = T_0(1 - \beta^2)^{\frac{1}{2}} / (1 - \beta \cos \theta) \simeq T_0(1 + \beta \cos \theta + \frac{\beta^2}{2} \cos 2\theta + \mathcal{O}(\beta^3))$$
- At every point in the sky a black body spectrum at $T(\theta)$ is observed.
- Velocity of solar system barycenter assuming $T_0 = T_\gamma$

$$v = 369.0 \pm 0.9 \text{ km s}^{-1} \quad \text{towards} \quad (l, b) = (263.99^\circ \pm 0.14^\circ, 48.26^\circ \pm 0.03^\circ)$$

- Dipole is frame-dependent  determine “absolute rest frame” in which CMB dipole would be zero.

Cosmology 2

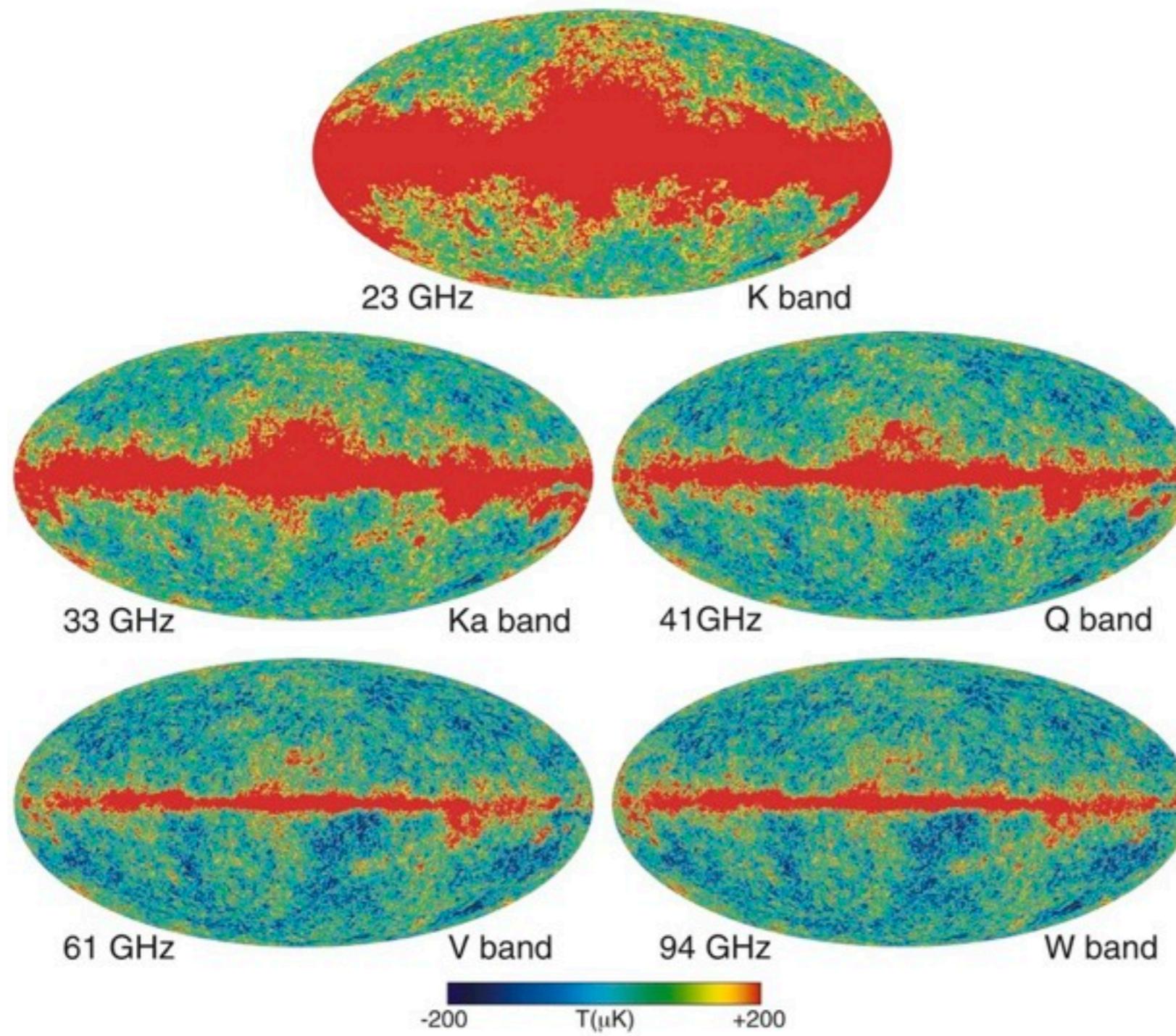
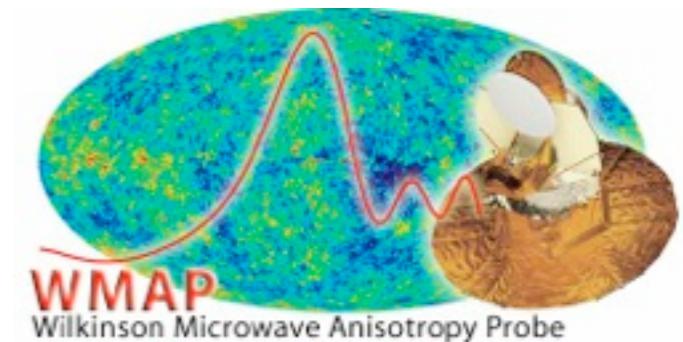
- Higher order multipoles $\ell \geq 2$
- result of perturbations in the density of the early universe, present before last scattering.



WMAP 9:
temperature angular
power spectrum

Bennett et al. 2013

Cosmology 2

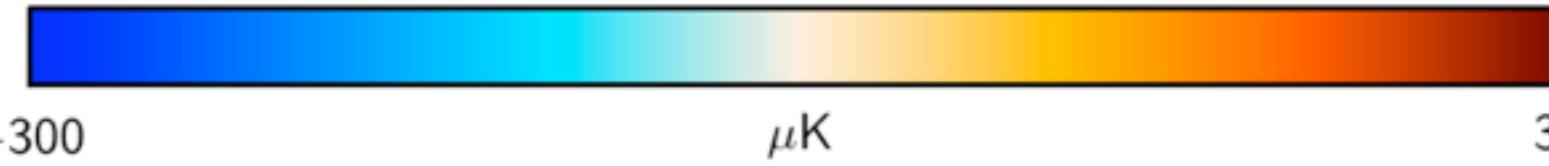
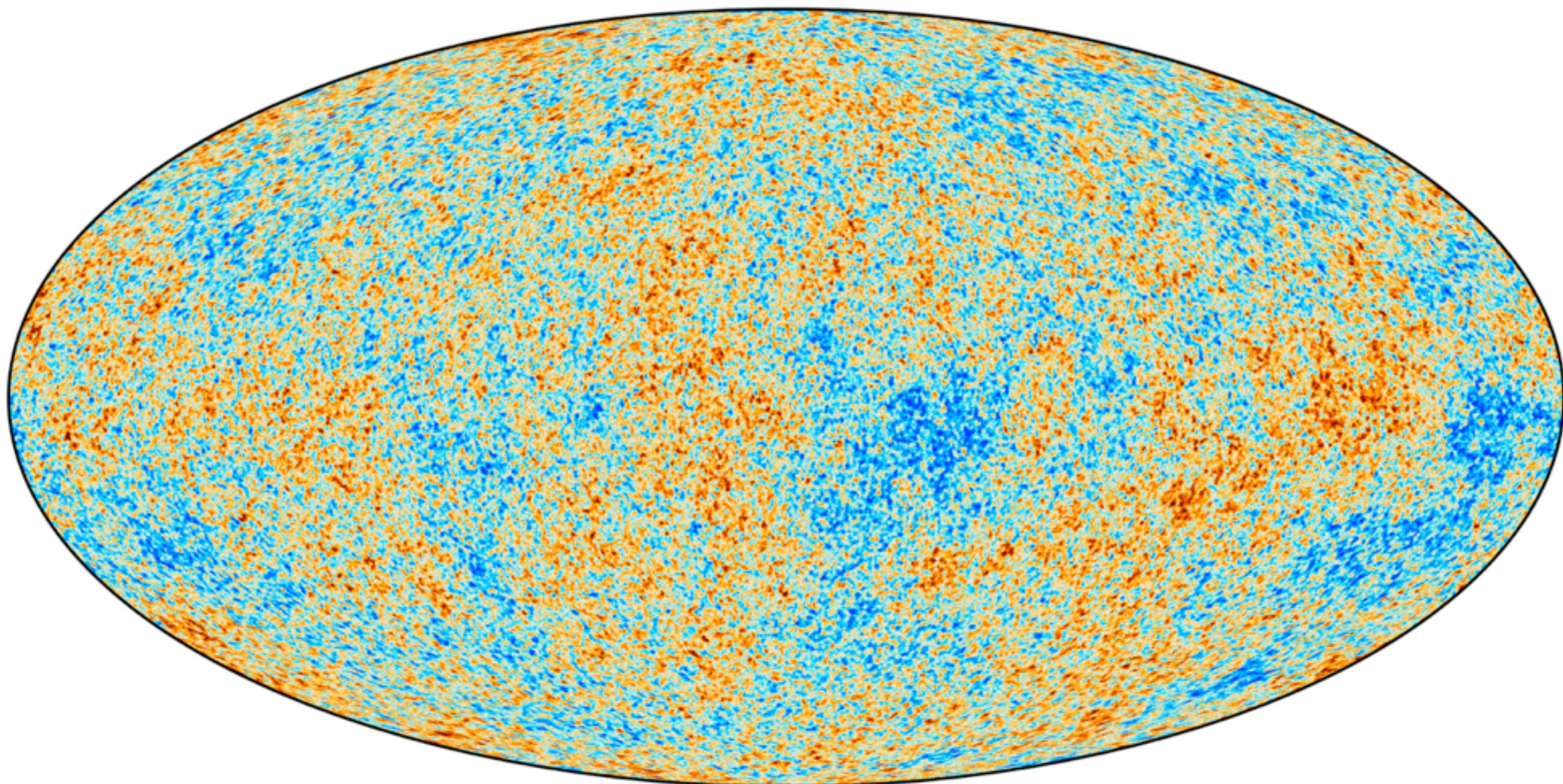


WMAP 9:
temperature sky maps

Bennett et al. 2013

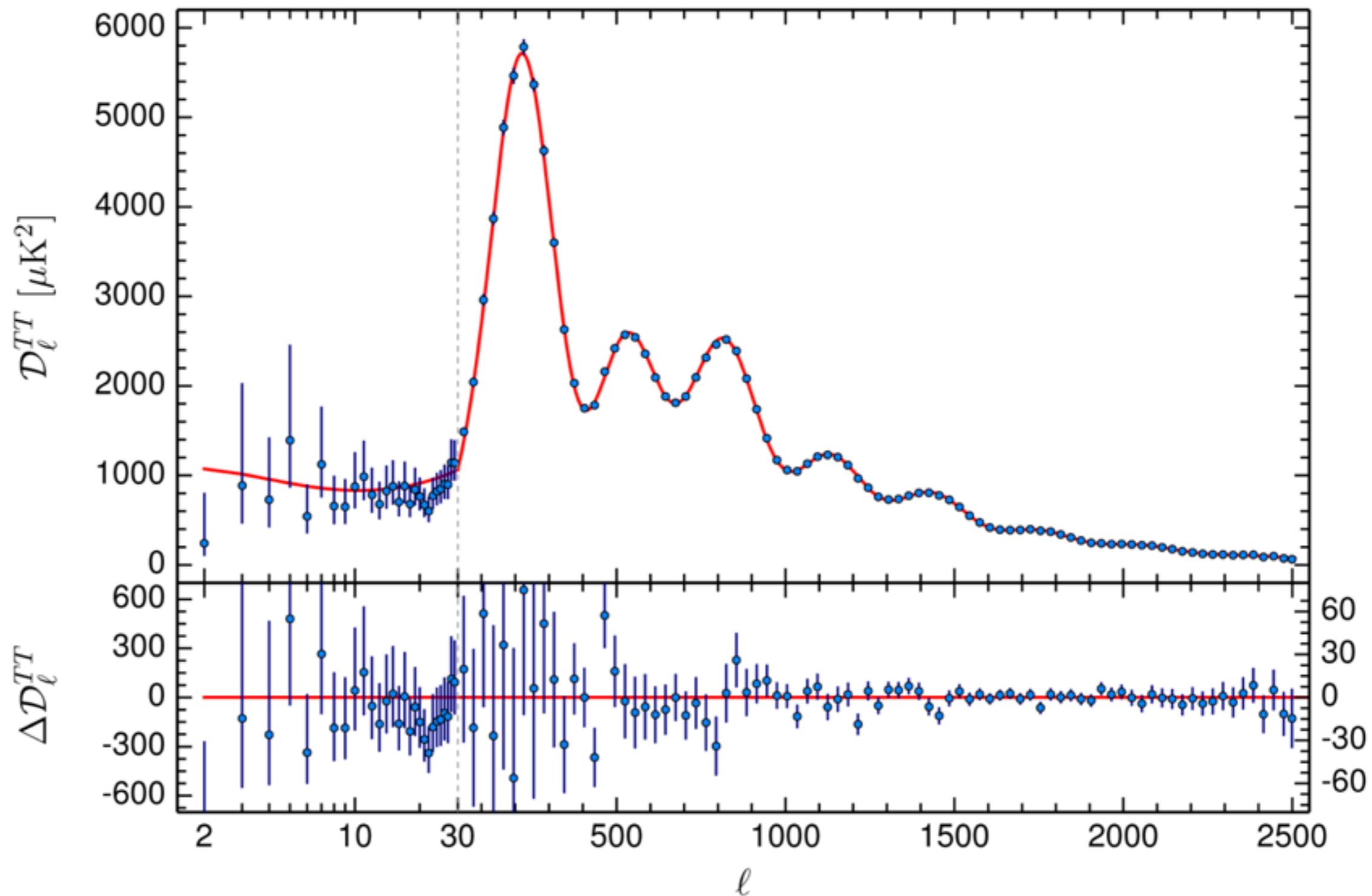
Cosmology 2

Planck 15:
temperature sky maps

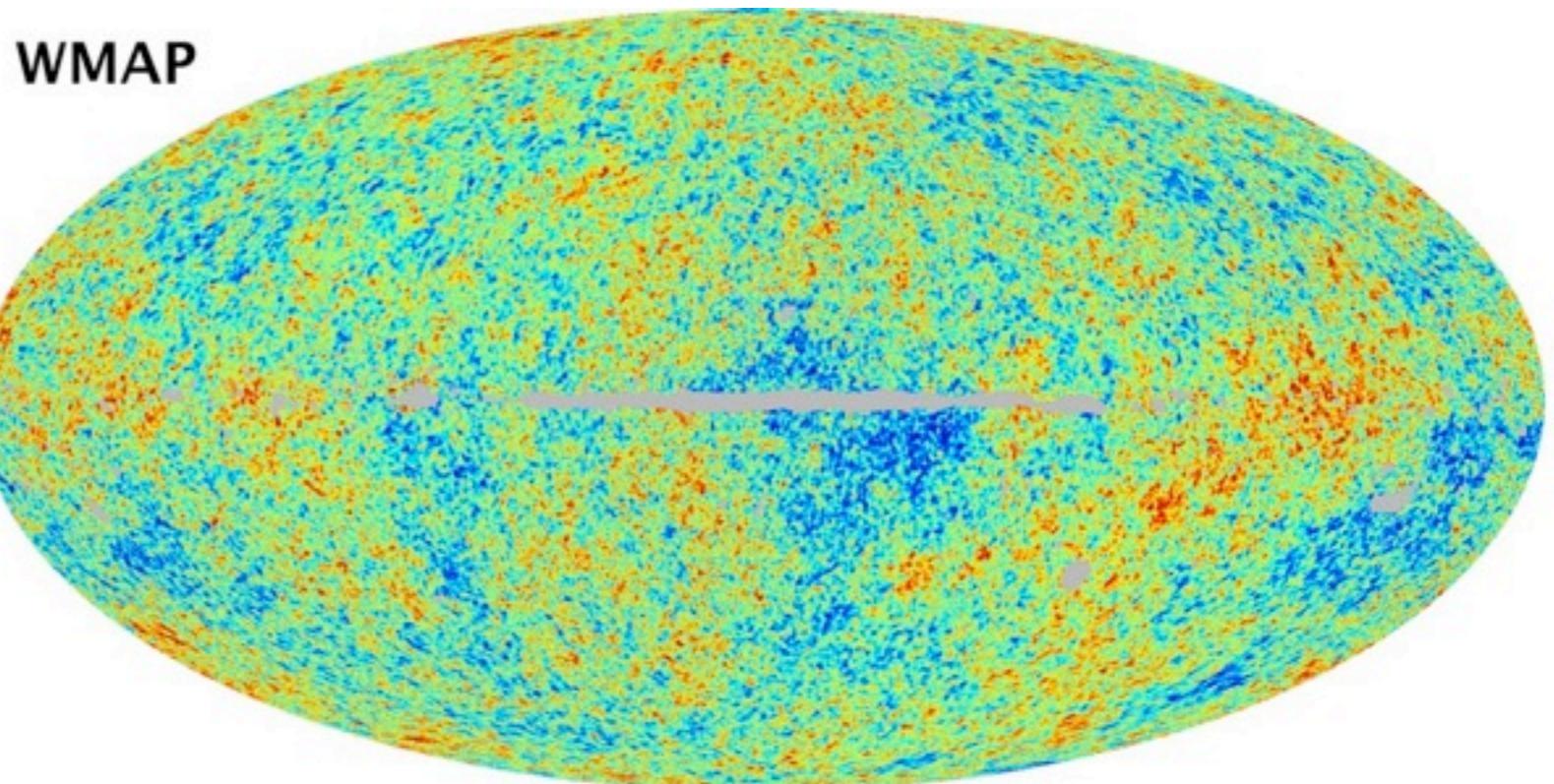


Cosmology 2

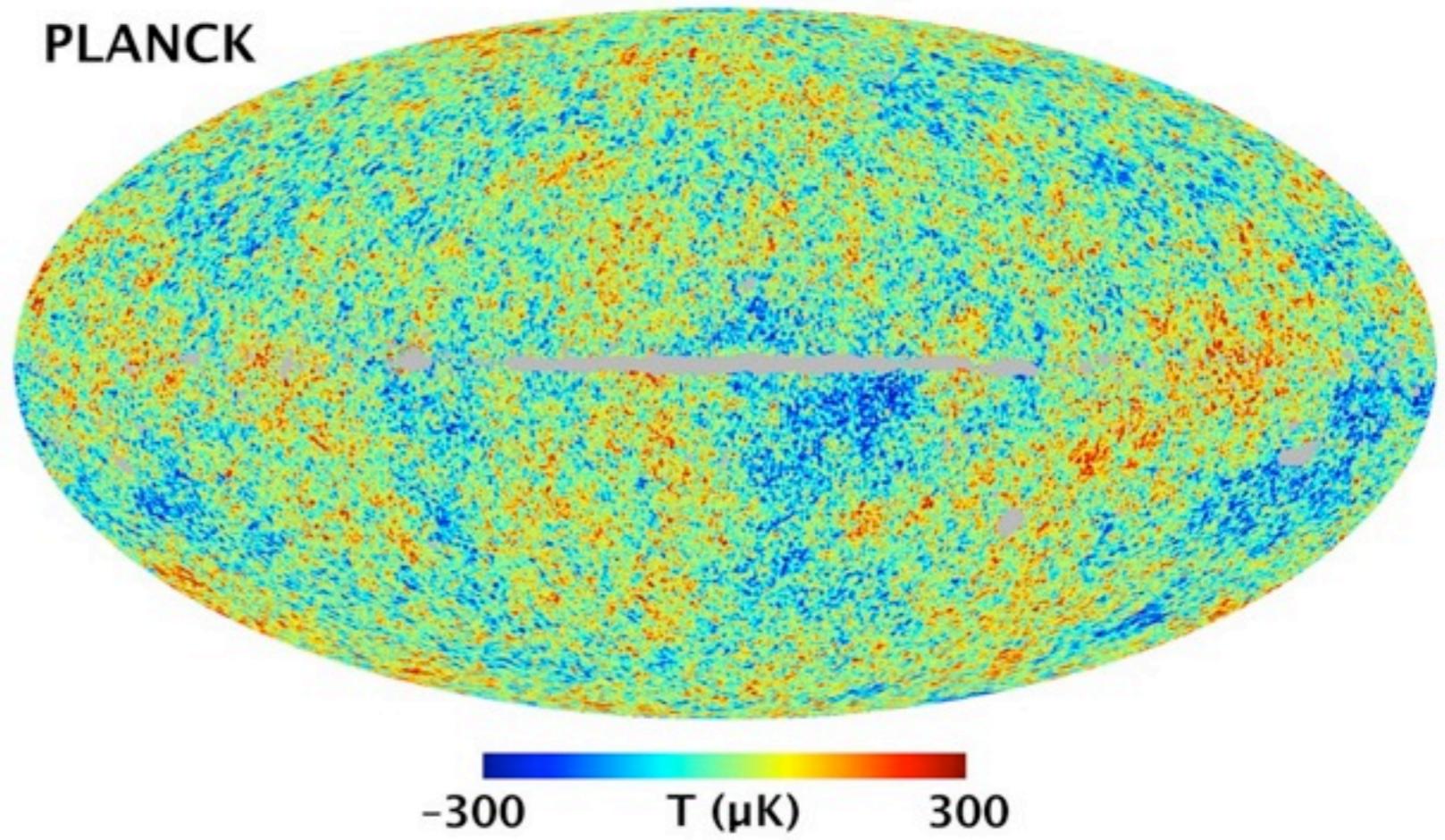
Planck 15:
temperature angular
power spectrum



Cosmology 2



WMAP 9 vs.
Planck13



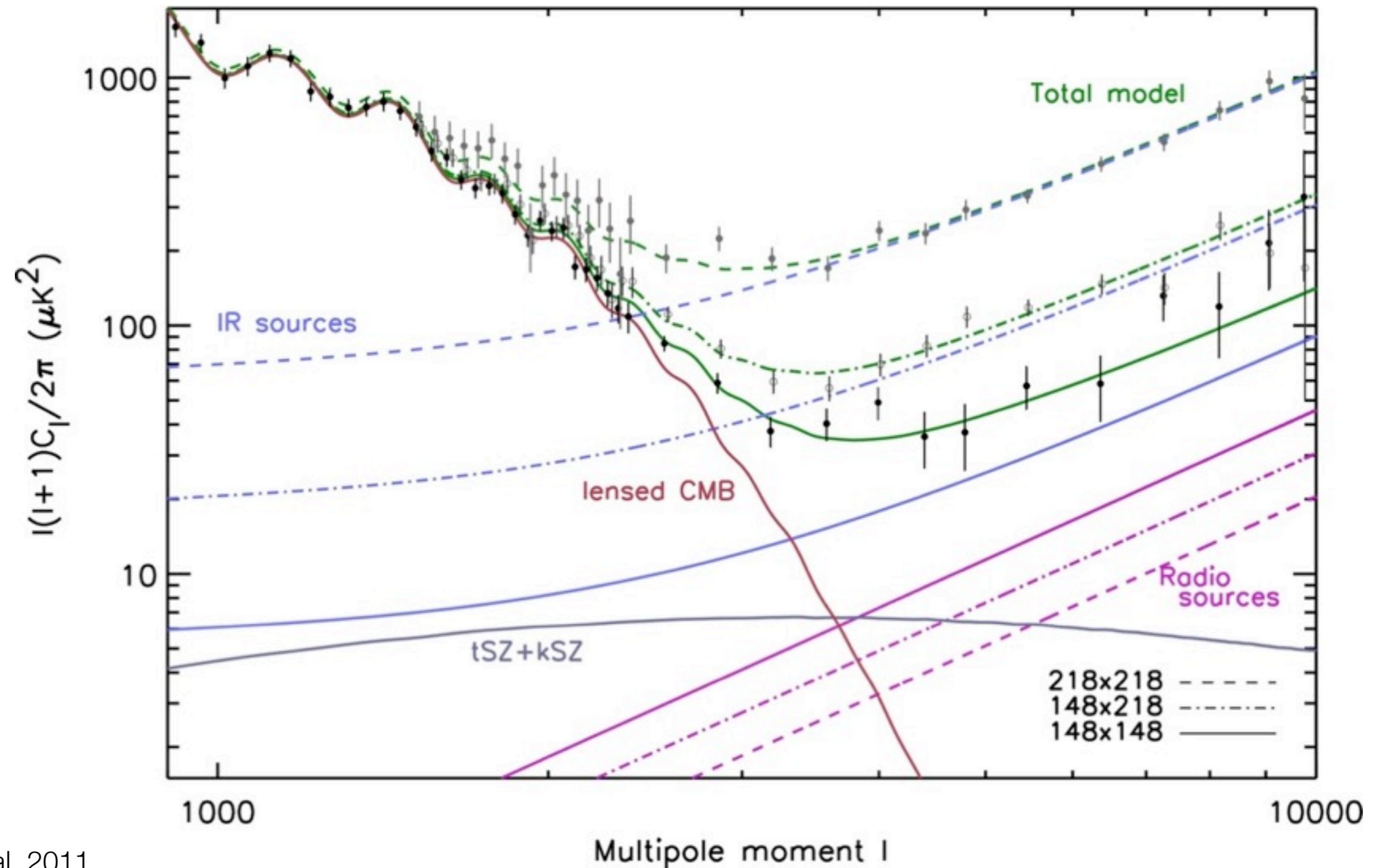
Cosmology 2

The [Atacama Cosmology Telescope](#) ([ACT](#)) is a **6 meter** telescope designed to map the Cosmic Microwave Background (CMB) over a large sky area with an angular resolution of **1 arcmin**. The instrument observes at three frequencies (**148, 218, and 277 GHz**). In order to minimize disturbance from atmospheric effects, the telescope is situated on the Atacama Plateau in northern Chile, at an elevation of **5190 m**.



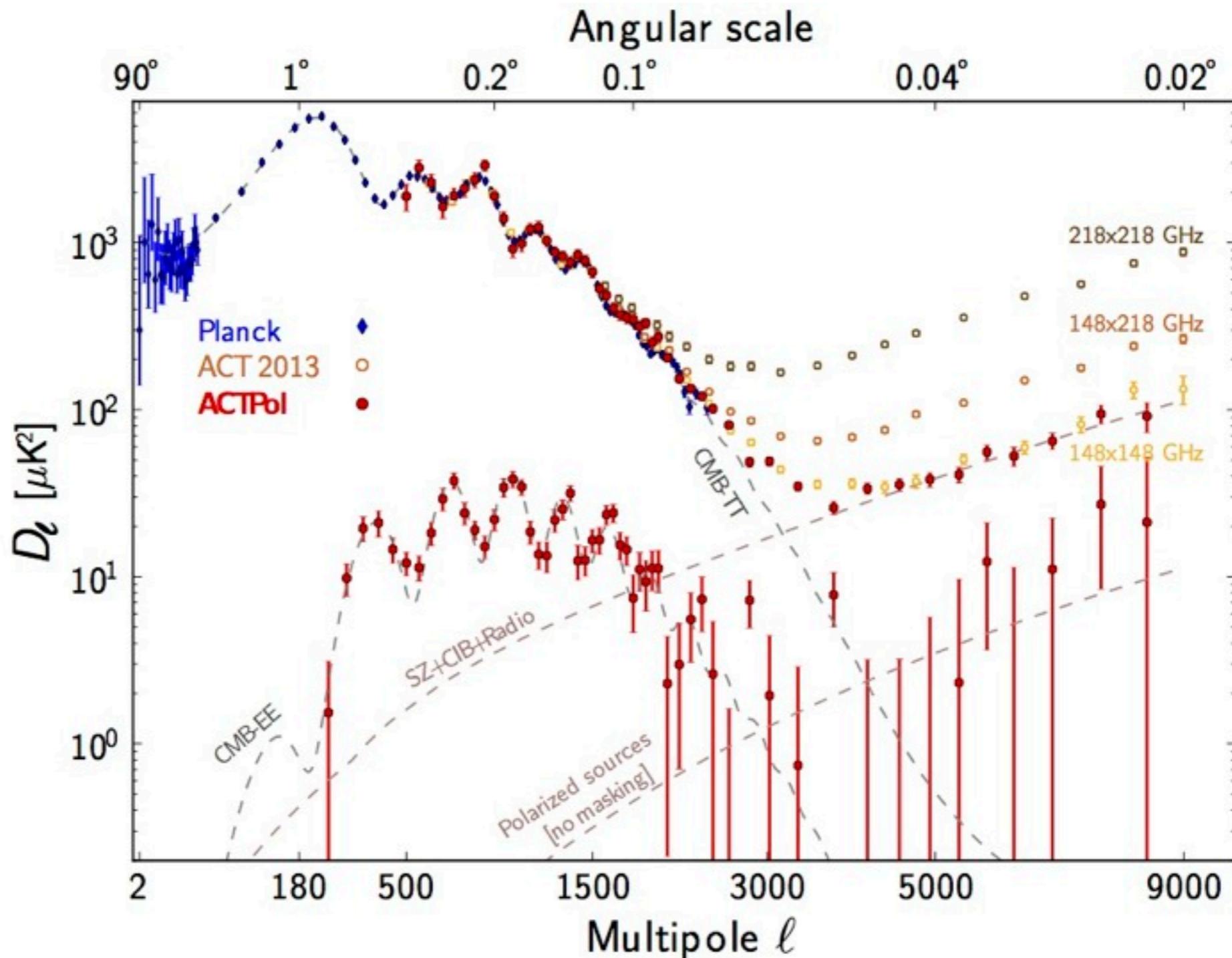
Atacama Cosmology Telescope (ACT)

Cosmology 2



Dunkley et al. 2011

Cosmology 2



Naess et al. 2014

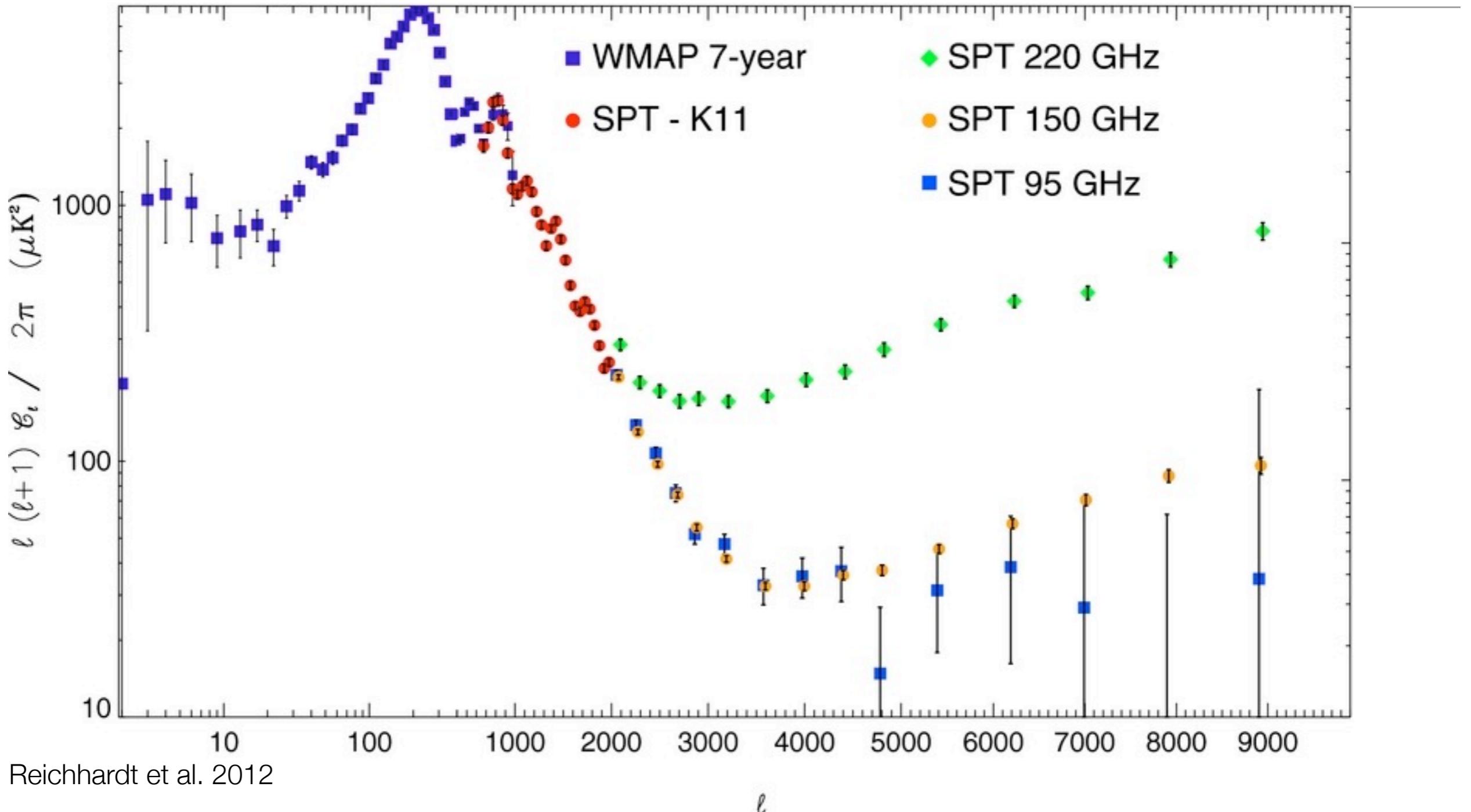
Cosmology 2

The **South Pole Telescope** is a 10 meter diameter telescope operating at the NSF South Pole research station. The telescope is designed for conducting large-area millimeter and sub-millimeter wave surveys of faint, low contrast emission, as required to map primary and secondary anisotropies in the cosmic microwave background.



South Pole Telescope (SPT)

Cosmology 2

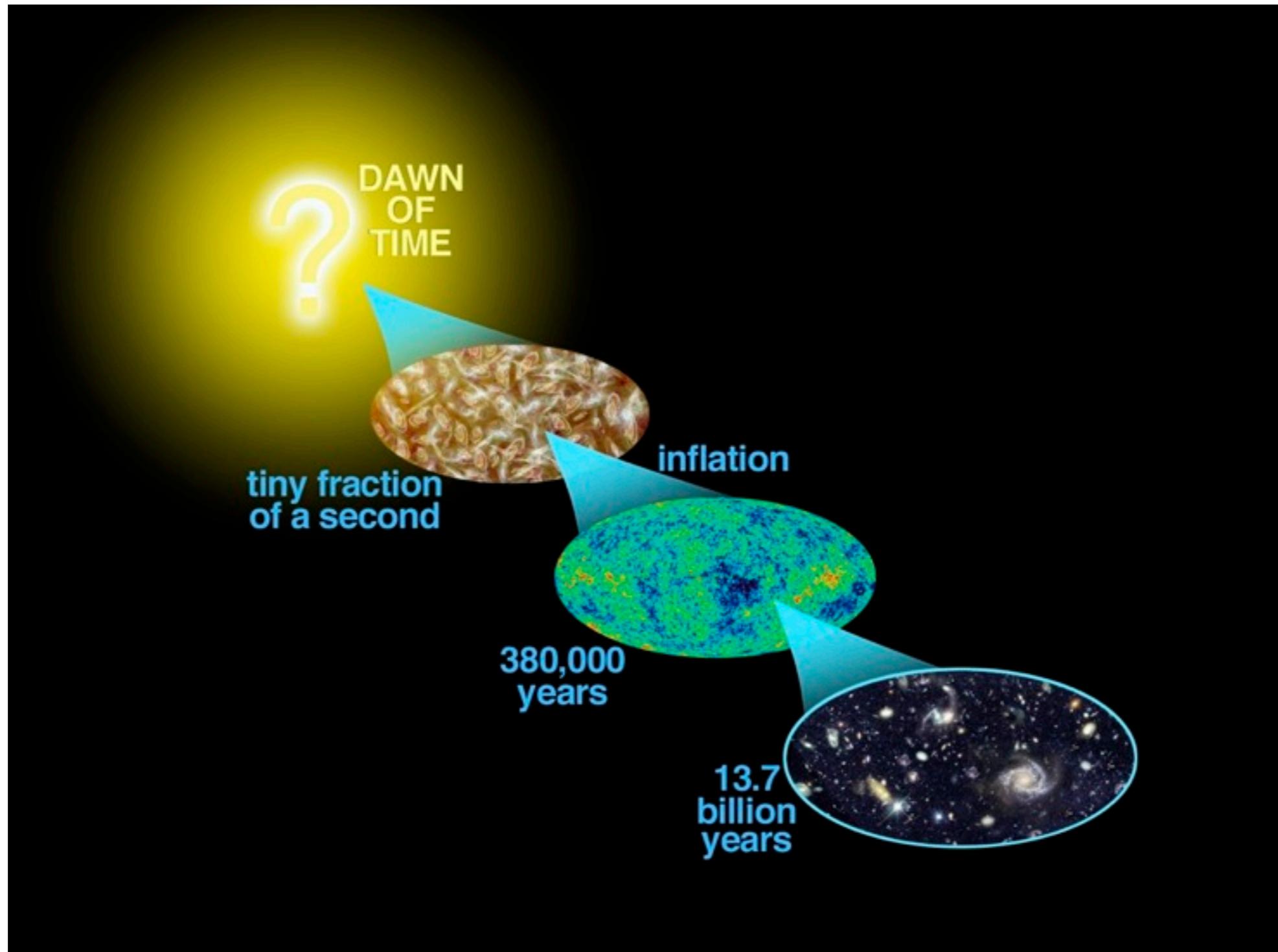


Reichhardt et al. 2012

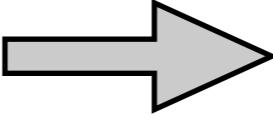
ℓ

WMAP7 and SPT bandpowers. Below $\ell = 2000$, the primary CMB anisotropy is dominant at all frequencies. On smaller scales, the CIB, radio sources, and secondary CMB anisotropies contribute to the signal. With the SPT source masking, the CIB is the largest source of power on sub-arcminute scales at 150 and 220 GHz. Due to the relative spectral behavior of the CIB and synchrotron emission, the 95 GHz bandpowers also have a significant contribution from radio sources.

Cosmology 2



Cosmology 2

- Origin of CMB anisotropies?
- Ripples in space-time and matter momentum tensor  perturb FRW metric and energy
- At linear order: 3 types of perturbations: scalar, vector and tensor

(flat) FRW

$$ds^2 = a^2(\eta)(d\eta^2 - \delta_{ij}dx^i dx^j)$$

most general linear
perturbation

$$ds^2 = a^2(\eta) [(1 - 2A)d\eta^2 + 2B_i d\eta dx^i - [(1 + 2D)\delta_{ij} + 2E_{ij}] dx^i dx^j]$$

expand functions in
scalar, vector and tensor
harmonics

... for details Kodama, Sasaki
ProgTheorPhysSup 78 (1984) 1 or text
book on CMB physics

Cosmology 2

- Perturbed Boltzmann equation
- determines CMB temperature anisotropies and polarization

$$q = a(\eta)p(\eta, \mathbf{x})$$

- Photon phase-space distribution

$$f(\eta, \mathbf{x}, \mathbf{n}, q) = f(q) + \delta f(\eta, \mathbf{x}, \mathbf{n}, q)$$

- define brightness function

$$\Theta \equiv \delta T/T$$

$$f(q) = \frac{1}{e^{q/T_0} - 1}$$

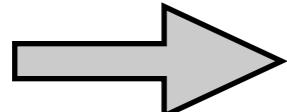
unperturbed universe

$$f(\eta, \mathbf{x}, \mathbf{n}, q) = \frac{1}{e^{q/T_0(1+\Theta)} - 1}$$



$$\delta f(\eta, \mathbf{x}, \mathbf{n}, q) = -q \frac{df(q)}{dq} \Theta(\eta, \mathbf{x}, \mathbf{n})$$

Cosmology 2

-  Boltzmann equation for $\Theta(\eta, \mathbf{k}, \mathbf{n})$

$$\frac{\partial \Theta}{\partial \eta} + ik\mu\Theta - \frac{\partial \Phi}{\partial \eta} + ik\mu\Psi = C[\Theta]$$

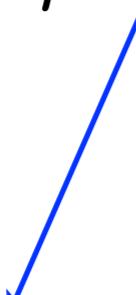
$\mu = \hat{\mathbf{k}} \cdot \mathbf{n}$

collision term

scalar mode perturbation,
conformal Newtonian gauge

Boltzmann hierarchy for $\Theta_\ell(\eta, \mathbf{k})$

Expand $\Theta(\eta, \mathbf{k}, \mathbf{n})$ in spherical harmonics



$$\Theta_0 = \frac{\delta_\gamma}{4}, \quad \Theta_1 = \frac{V_\gamma}{3}, \quad \Theta_2 = \frac{\Pi_\gamma}{12}$$

Lyth, Liddle “The primordial density perturbation”

Cosmology 2

- To solve the Boltzmann equation (hierarchy) together with evolution equations of all matter components of the cosmic fluid (and additional ones you want to put...) there are open source **CMB Boltzmann codes**

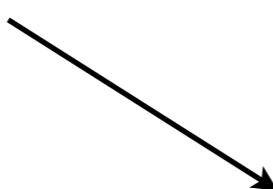
- **COSMOS (Bertschinger)**

- **CMBFAST (Seljak, Zaldarriaga 1996)**

http://lambda.gsfc.nasa.gov/toolbox/tb_cmbfast_ov.cfm

- **CAMB (Lewis et al. 2000)**

<http://camb.info/>



cosmomc:

Markov Chain Monte Carlo

exploration of parameter space

<http://cosmologist.info/cosmomc/>

Cosmology 2

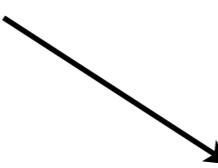
- More Boltzmann codes...

- **CMBEASY (Doran 2005)**

<http://www.thphys.uni-heidelberg.de/~robbers/cmbeasy/>

- **CLASS (Blas et al. 2011)**

<http://class-code.net/>



montepython:
Markov Chain Monte Carlo

exploration of parameter space

https://github.com/baudren/montepython_public

Cosmology 2

- Below $T < 1$ MeV the cosmic fluid has four components:

* baryons

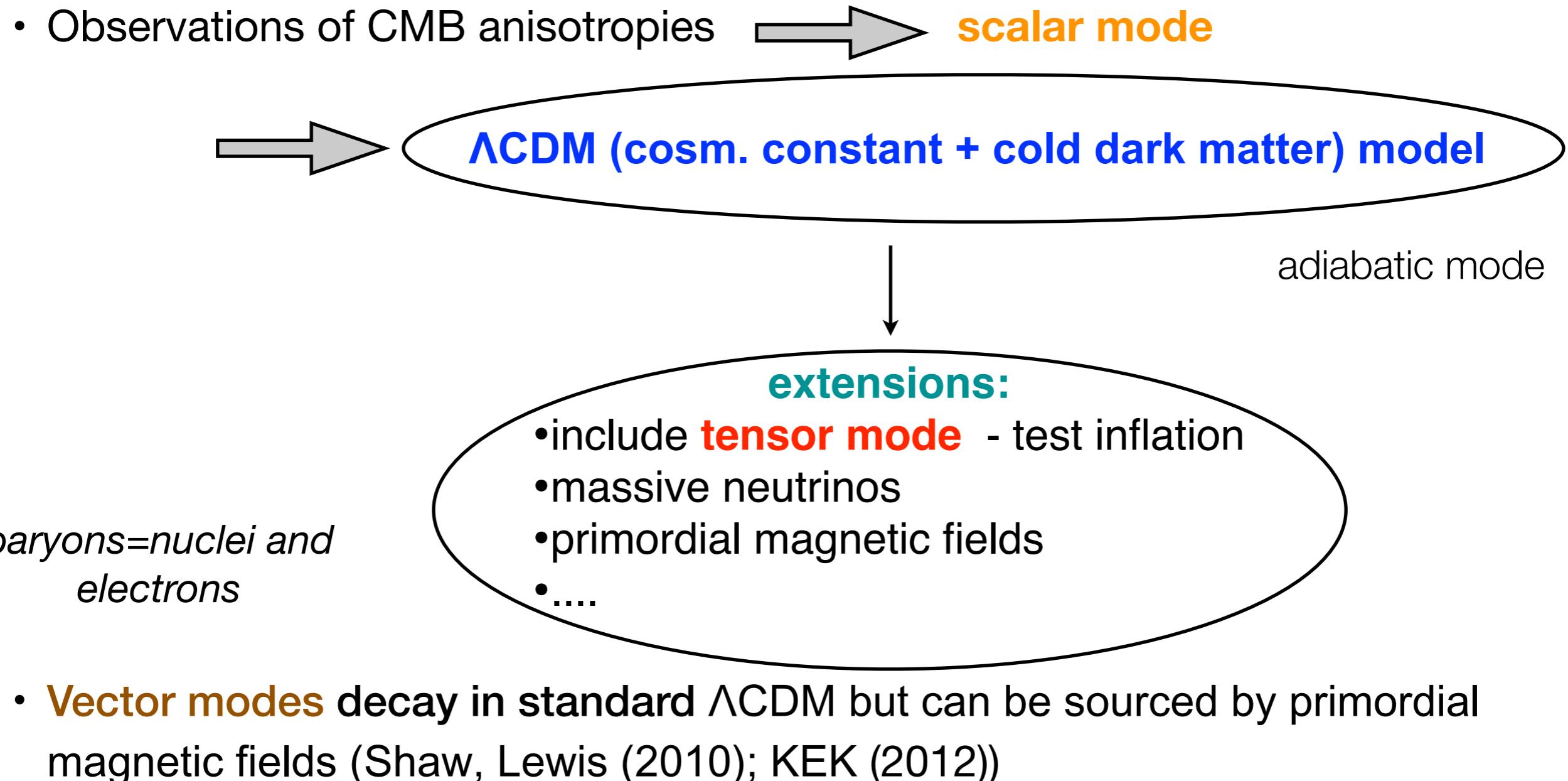
* cold dark matter (CDM)

* photons

*(light) neutrinos

baryons=nuclei and electrons since their number densities are basically equal at each point in space because of Coulomb interaction

Cosmology 2



Cosmology 2

- **Scalar mode**

- just focus on density perturbation and velocity of photons and baryons

gauge-invariant formalism

Kodama, Sasaki (1984),
Doran et al. (2003),
KEK (2011)

$$\dot{\Delta}_\gamma = -\frac{4}{3}kV_\gamma,$$

$$\dot{V}_\gamma = k(\Psi - \Phi) + \frac{k}{4}\Delta_\gamma - \frac{k}{6}\pi_\gamma + \tau_c^{-1}(V_b - V_\gamma),$$

$$\dot{\Delta}_b = -kV_b - 3c_s^2\mathcal{H}\Delta_b.$$

$$\begin{aligned}\dot{V}_b = & (3c_s^2 - 1)\mathcal{H}V_b + k(\Psi - 3c_s^2\Phi) + kc_s^2\Delta_b \\ & + R\tau_c^{-1}(V_\gamma - V_b)\end{aligned}$$

$$\tau_c^{-1} = an_e\sigma_T$$

mean free path of photons between scatterings

$$R \equiv \frac{4}{3} \frac{\rho_\gamma}{\rho_b}.$$

$c_s^2 = \frac{\partial \bar{p}}{\partial \bar{\rho}}$ is the adiabatic sound speed

Cosmology 2

- Initial conditions
- set on superhorizon scales $k\eta \ll 1$
- **adiabatic initial conditions**
- Radiation-matter fluid specific entropy: $S_{\gamma b} = S_\gamma / n_b$

$$\frac{\delta S_{\gamma b}}{S_{\gamma b}} = \frac{\delta S_\gamma}{S_\gamma} - \frac{\delta n_b}{n_b} = 3 \frac{\delta T_\gamma}{T_\gamma} - \frac{\delta n_b}{n_b} = \frac{\delta \rho_\gamma}{\rho_\gamma + p_\gamma} - \frac{\delta n_b}{n_b}$$

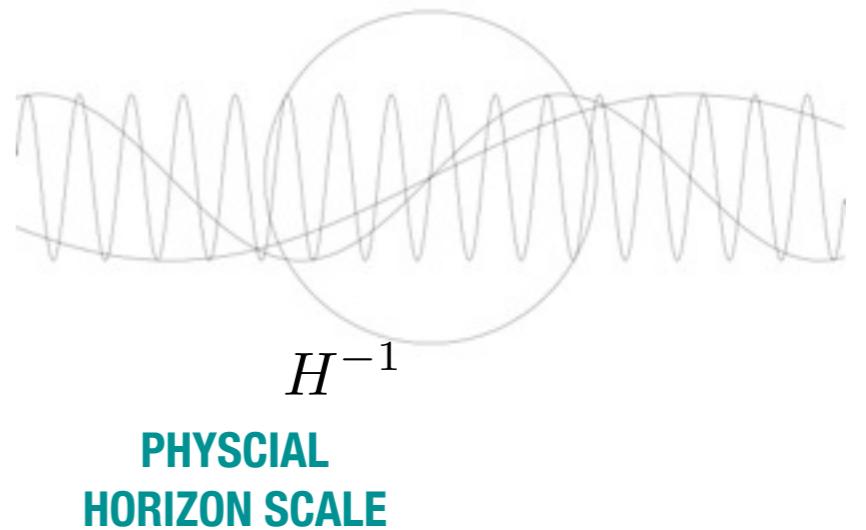
$\frac{\delta S_{\gamma b}}{S_{\gamma b}} = 0 \Rightarrow \frac{3}{4} \delta_\gamma = \delta_b$

adiabatic i.c.

generalize to all components j

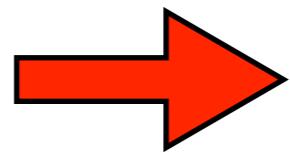
$p_j = w_j \rho_j$

$\boxed{\frac{S_{\gamma j}}{S_{\gamma j}} = \frac{\delta_\gamma}{w_\gamma + 1} - \frac{\delta_j}{1 + w_j}}$



Cosmology 2

$$\frac{\delta S_{\gamma j}}{S_{\gamma j}} = 0$$



Adiabatic i.c.

$$\frac{1}{4}\delta_\gamma = \frac{1}{4}\delta_\nu = \frac{1}{3}\delta_b = \frac{1}{3}\delta_c$$



curvature perturbation ζ

$$\zeta = \frac{\Delta_\gamma}{4}$$

- Gauge-invariant formalism:
- Single field inflation generates adiabatic initial conditions. Inflaton decays into the usual species, their overall ratios are fixed

$$\delta(n_A/n_B) = 0 \Rightarrow \delta S_{A,B}/S_{A,B} = 0$$

Cosmology 2

- Curvature perturbation from inflation

$$\zeta = -H \frac{\delta \rho}{\dot{\rho}} \Rightarrow \zeta = -H \frac{\delta \phi}{\dot{\phi}}$$

$$\langle \zeta_{\mathbf{k}} \zeta_{\mathbf{k}'} \rangle = \frac{2\pi^2}{k^3} \mathcal{P}_\zeta \delta_{\mathbf{k}, \mathbf{k}'}$$

$$\mathcal{P}_\zeta(k) = \frac{1}{4\pi} \left(\frac{H^2}{\dot{\phi}} \right)^2 \Big|_k$$

slow roll inflation

$$\mathcal{P}_\zeta(k) = \frac{1}{24\pi^2 M_P^4} \left. \frac{V}{\epsilon} \right|_{k=aH}$$

Spectral index

$$n - 1 \equiv \frac{d \ln \mathcal{P}_\zeta(k)}{d \ln k}$$

spectral index

$$n(k) - 1 = -6\epsilon + 2\eta$$

running of spectral index

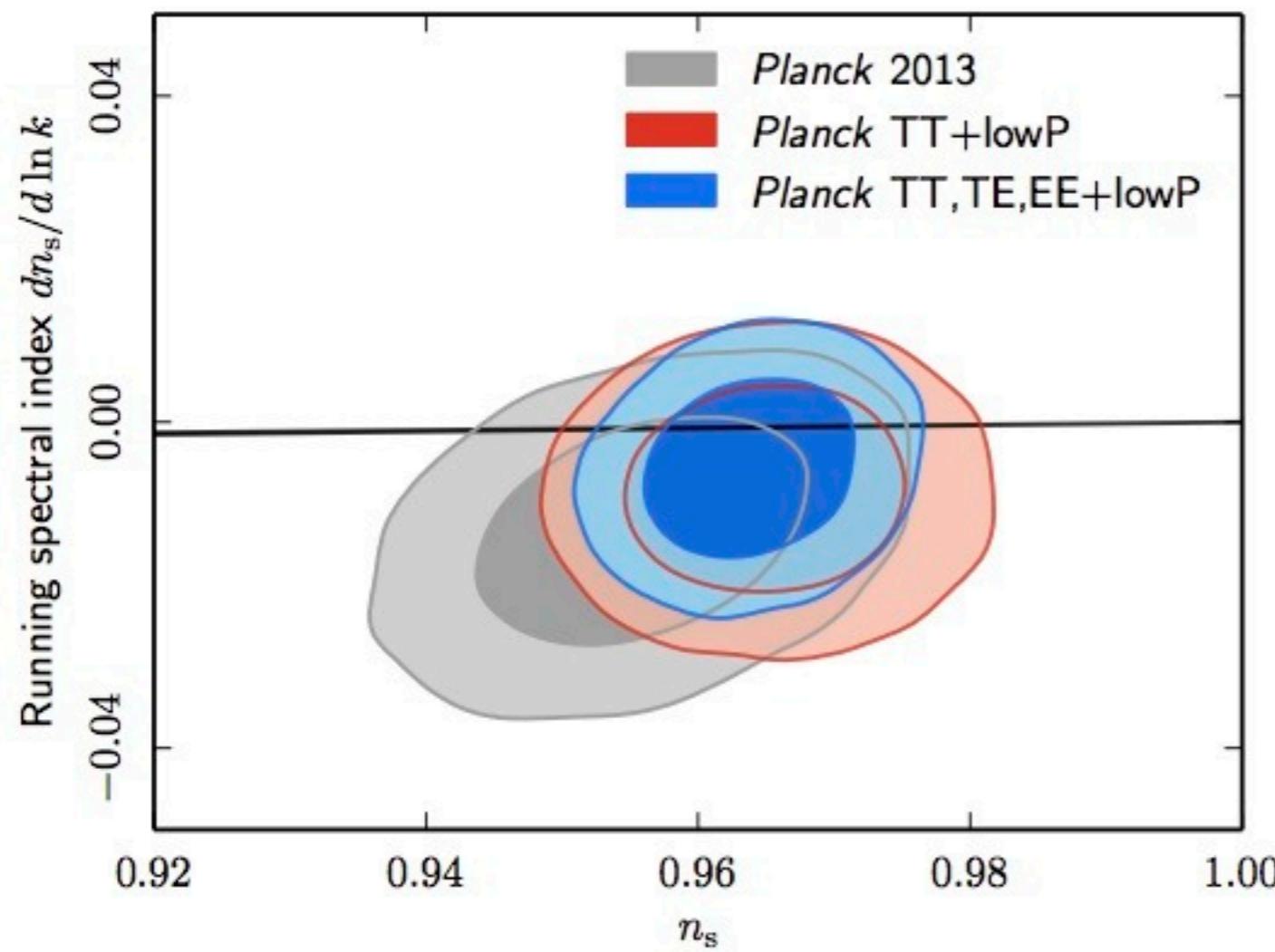
$$dn/d \ln k = -16\epsilon\eta + 24\epsilon^2 + 2\xi$$

slow roll parameters

$$\epsilon(\phi) = \frac{M_P^2}{16\pi} \left(\frac{V'}{V} \right)^2 \quad \eta = \frac{M_P^2}{8\pi} \frac{V''}{V} \quad \xi \equiv \frac{M_P^4}{64\pi^2} \frac{V' V'''}{V^2}$$

Cosmology 2

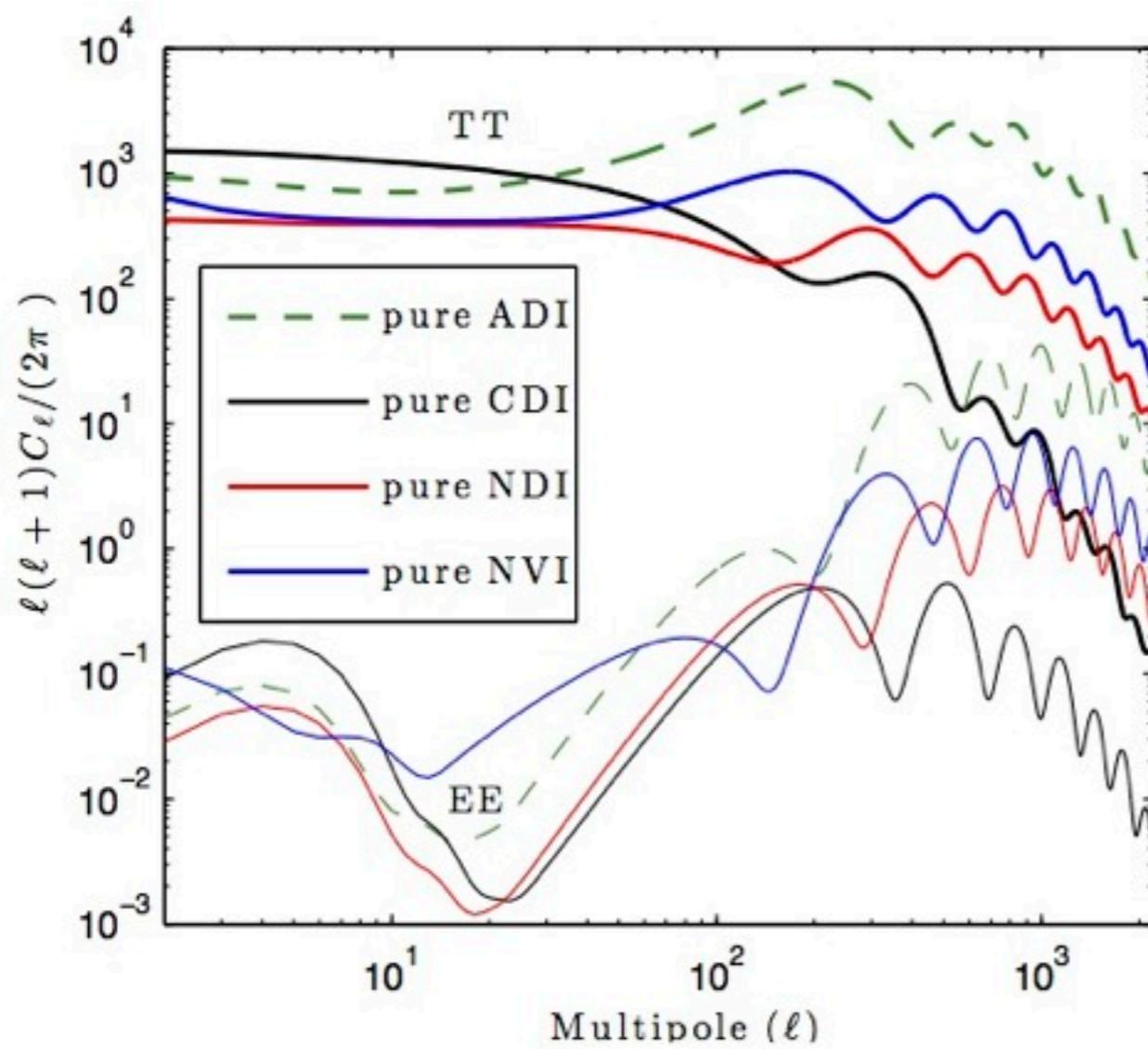
- Constraints from Planck 15



Ade et al.2015 Planck 2015.XX.

Cosmology 2

- Another type of initial condition: isocurvature i.c. $\zeta = 0$



$$\frac{\delta S_{ij}}{S_i} = \text{const.} \neq 0$$

ADI - adiabatic mode
CDI - CDM density isocurvature mode
NDI - neutrino density isocurvature mode
NVI - neutrino velocity isocurvature mode

Ade et al. 2015 Planck 2015.XX.

Cosmology 2

- Now going back to the **baryon-photon fluid**. In the **tight-coupling limit** the equations can be combined to the equation of a forced harmonic oscillator:

$$\ddot{\Delta}_\gamma + \frac{\dot{R}_b}{1+R_b} \dot{\Delta}_\gamma + c_{sb\gamma}^2 k^2 \Delta_\gamma \simeq \frac{4k^2}{3} \frac{2+R_b}{1+R_b} \Phi,$$

which is solved by

$$\begin{aligned} \Delta_\gamma(\tau) = & \frac{1}{(1+R_b)^{1/4}} \left[\Delta_\gamma(0) \cos(kr_s(\tau)) \right. \\ & + \frac{\sqrt{3}}{k} \left[\dot{\Delta}_\gamma(0) + \frac{1}{4} \dot{R}_b(0) \Delta_\gamma(0) \right] \sin(kr_s(\tau)) \\ & + \frac{\sqrt{3}}{k} \int_0^\tau d\tau' (1+R_b(\tau'))^{3/4} \right. \\ & \times \left. \sin[kr_s(\tau) - kr_s(\tau')] F(\tau') \right], \end{aligned}$$

where $R_b \equiv \frac{1}{R}$ and $c_{sb\gamma}^2 \equiv \frac{1}{3} \frac{1}{R_b+1}$



sound speed of baryon-photon fluid

$$F(\tau) \equiv \frac{4k^2}{3} \frac{2+R_b}{1+R_b} \Phi$$

sound horizon

$$r_s(\tau) \equiv \int_0^\tau c_{sb\gamma} d\tau'$$

Hu, Sugiyama

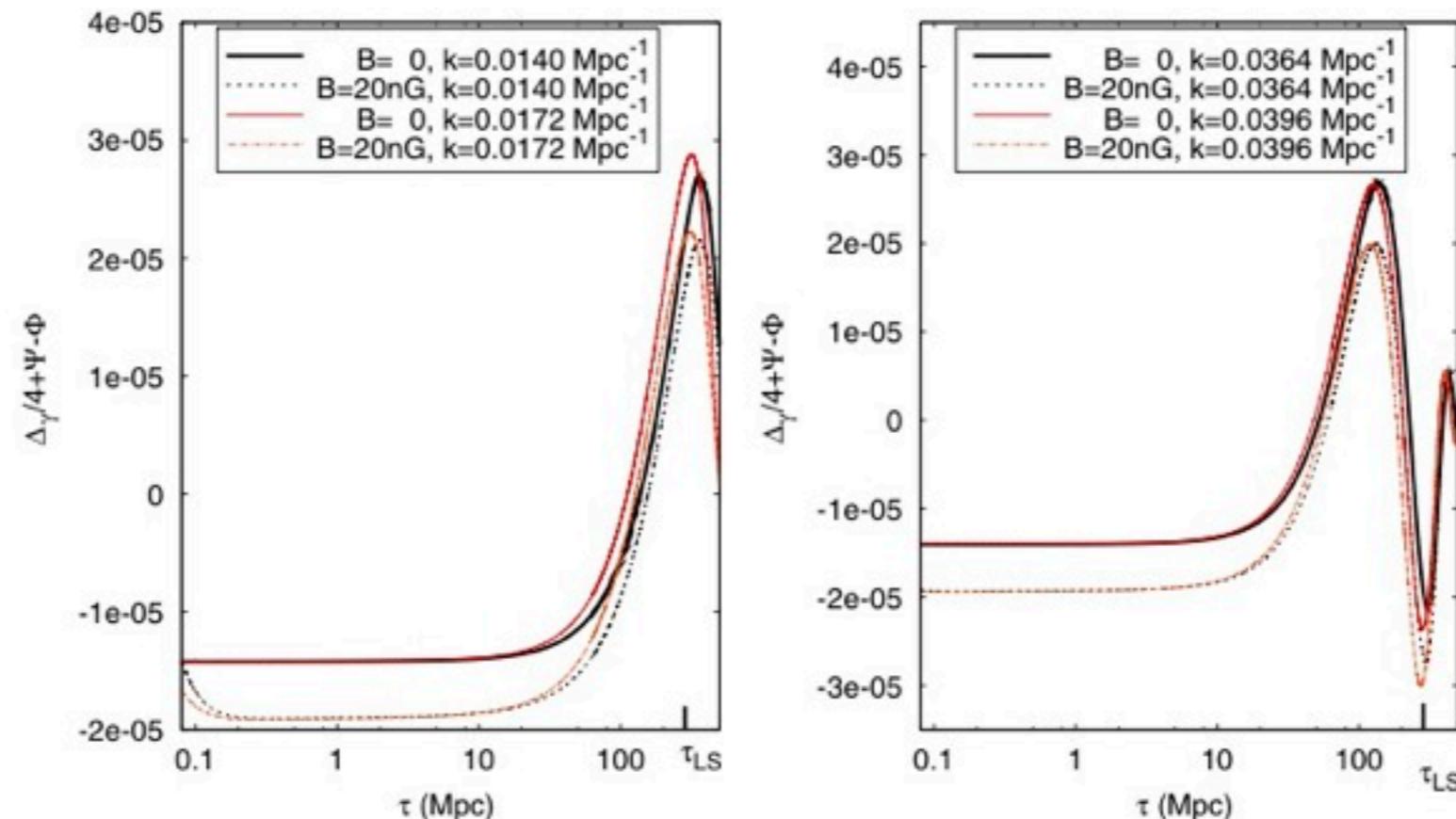
Cosmology 2

**Effective
temperature
perturbation** for

different wave
numbers

$$\frac{\Delta_\gamma}{4} + \Psi - \Phi$$

(ignore magnetic
field...)



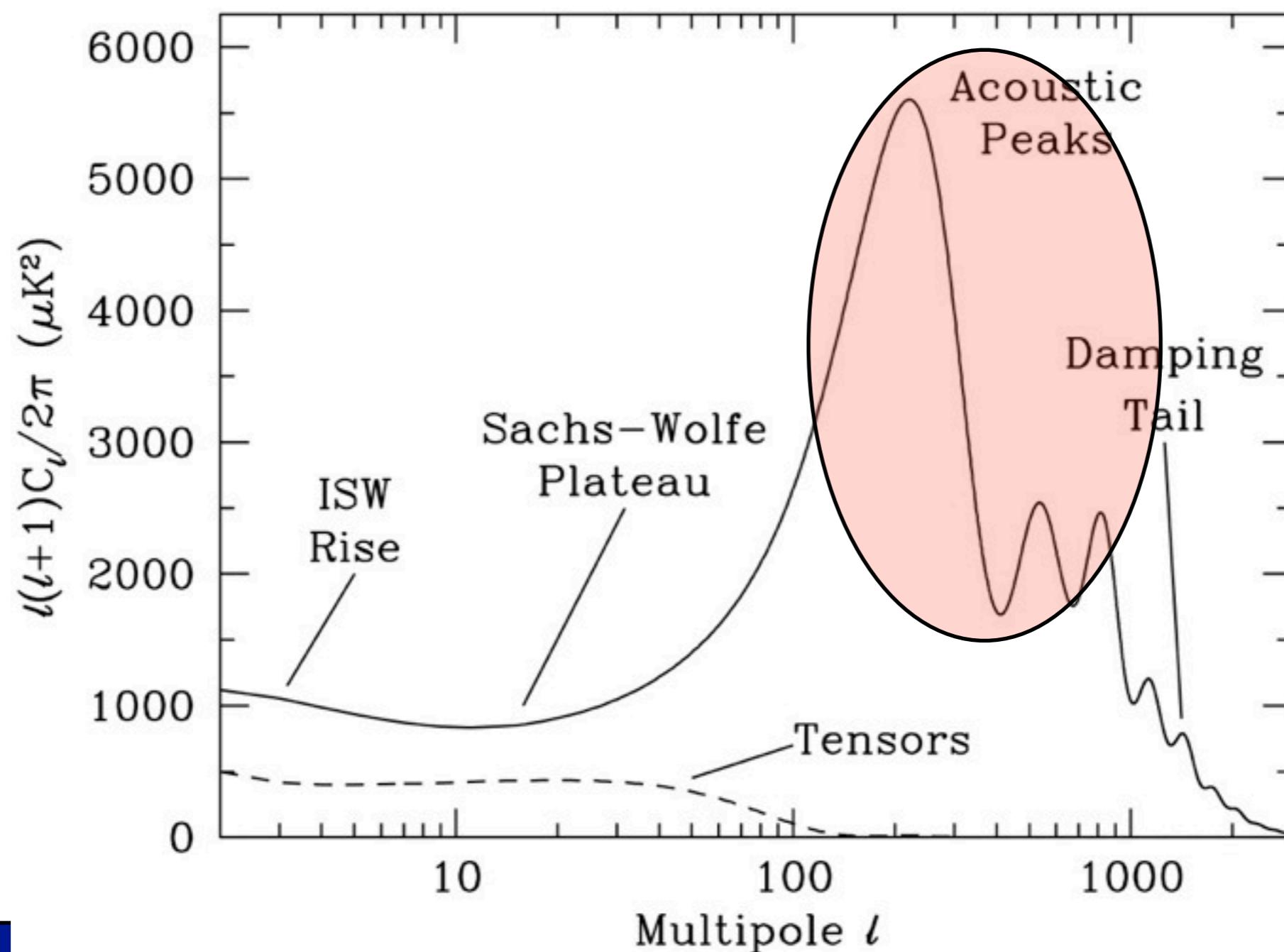
KEK (2011)

FIG. 6 (color online). The numerical solution of the evolution of the effective temperature perturbation is shown for different wave numbers. The WMAP7 best fit solution is shown in comparison with the solution in presence of a stochastic magnetic field with $B = 20 \text{ nG}$ and magnetic spectral index $n_B = -2.9$. Indicated is the time of last scattering, which in this case is $\tau_{\text{LS}} = 285 \text{ Mpc}$. *Left:* The solutions for the wave numbers $k = 0.0140 \text{ Mpc}^{-1}$ and $k = 0.0172 \text{ Mpc}^{-1}$ have a maximum at the time of last scattering. These wave numbers correspond to the region of the first acoustic peak, whereas the former corresponds to $\ell = 200$, the latter corresponds to $\ell = 244$. *Right:* The solutions for the wave numbers $k = 0.0364 \text{ Mpc}^{-1}$ and $k = 0.0396 \text{ Mpc}^{-1}$ have a minimum at the time of last scattering. These wave numbers correspond to the region of the second acoustic peak, whereas the former corresponds to $\ell = 523$, the latter corresponds to $\ell = 569$.

Cosmology 2

- Features in the angular power spectrum

(PDG 2014)



Cosmology 2

- Damping tail
- Close to last scattering photons and baryons are no longer strongly coupled leading to diffusion damping (Silk damping) of density perturbations.
- Amplitudes are multiplied by a factor

$$\exp(-k^2/k_D^2(\eta))$$

Silk scale (photon diffusion scale) k_D^{-1}
is roughly the comoving distance that
a photon has time to travel since
some initial time.

Cosmology 2

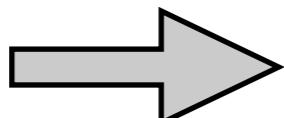
- To calculate photon diffusion scale: Model movement of photon in local baryon rest frame as random walk. Mean time between collisions:

$$t_c \sim (n_e \sigma_T)^{-1}$$

- Average number of steps in time t :

$$N = t/t_c$$

- During t the photon diffuses a distance: $d \sim \sqrt{N}t_c \sim (tt_c)^{1/2}$



$$ak_D^{-1} \simeq \left(\frac{t}{n_e \sigma_T} \right)^{\frac{1}{2}}$$

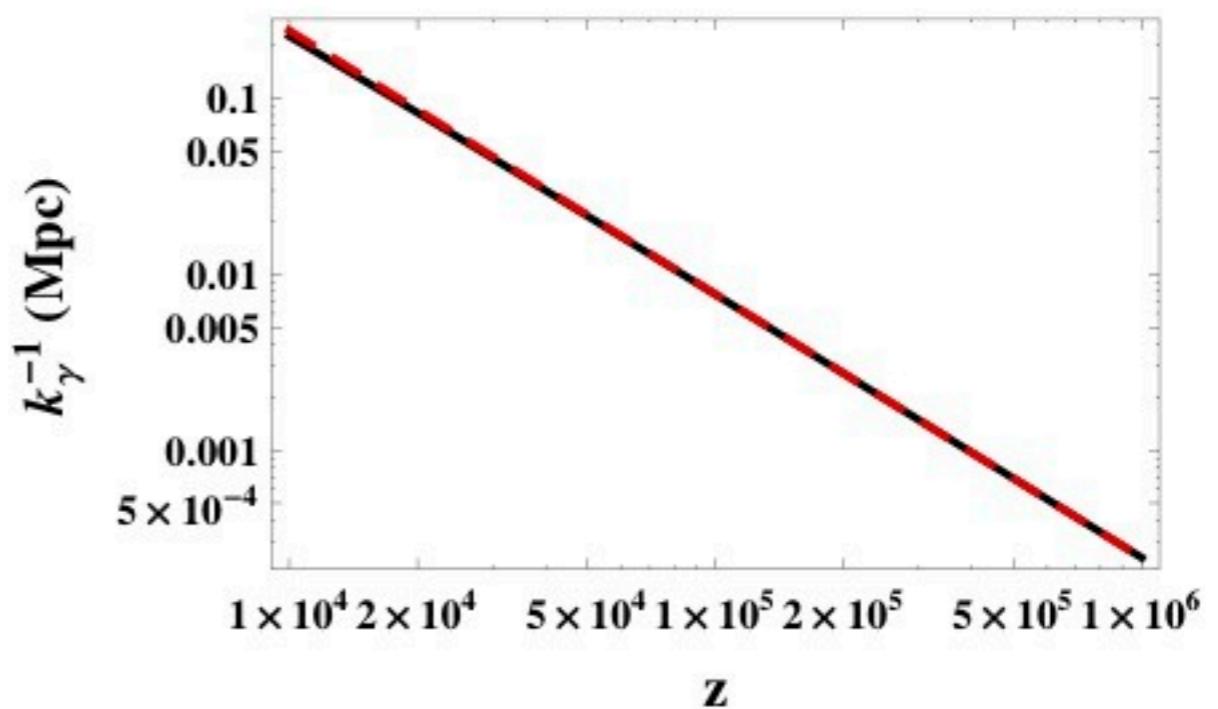
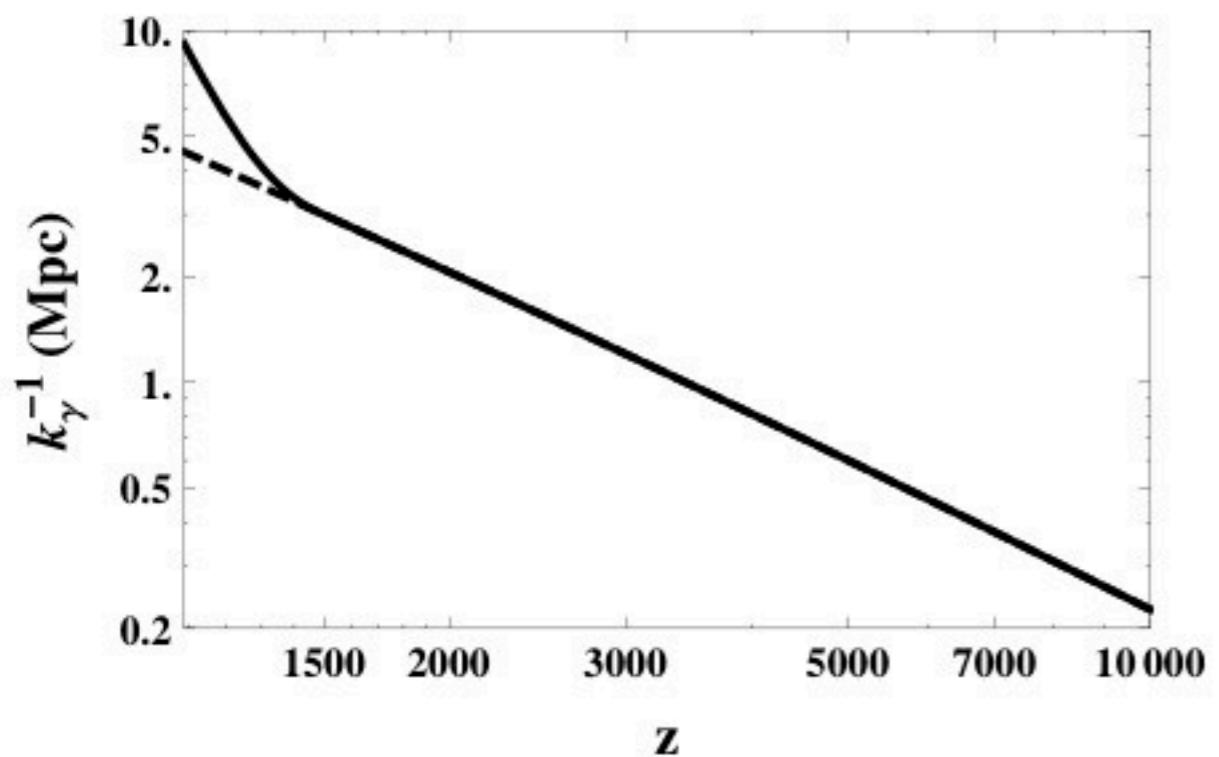
- during matter domination: $k_D^{-1} \sim a^{5/4}$

Each scale k^{-1} starts out bigger than the Silk scale at horizon entry with photon diffusion damping setting in when

$$k_D(\eta) = k$$

- during radiation domination: $k_D^{-1} \sim a^{3/2}$

Cosmology 2

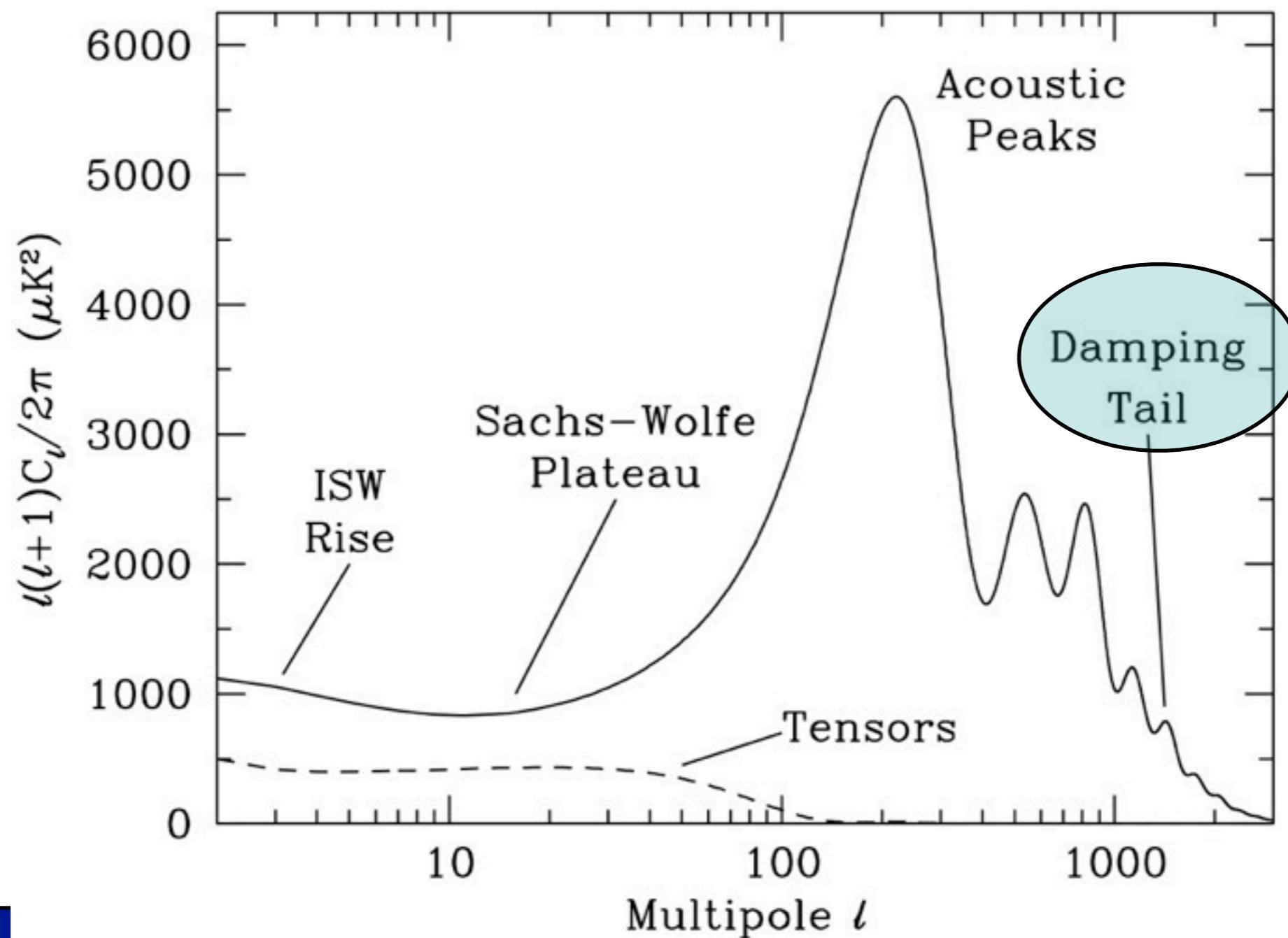


Evolution of photon diffusion scale as function of redshift z
(KEK, Komatsu (2014) based on model of Hu, Sugiyama (1995))

Cosmology 2

- Features in the angular power spectrum

(PDG 2014)



Cosmology 2

- **Sachs-Wolfe contribution**

- For all scales the solution of the Boltzmann hierarchy is formally given by a **line-of-sight integral** (Seljak, Zaldarriaga):

$$\Delta_\ell^{(s)}(q) = \int_{\eta_{in}}^{\eta_0} d\eta S(k, \eta) j_\ell[k(\eta_0 - \eta)]$$

↑
source function

- Split source function into contributions from the Sachs-Wolfe term (SW), integrated Sachs-Wolfe term (ISW), the Doppler term (dop) and a term related to polarisation (pol), in the newtonian gauge:

Cosmology 2

$$\Delta_\ell^{(s)}(q) = \int_{\eta_{in}}^{\eta_0} d\eta S(k, \eta) j_\ell[k(\eta_0 - \eta)]$$

$$S = \sum_i S_i$$

visibility function
 $g = \dot{\tau} e^{-\tau}$

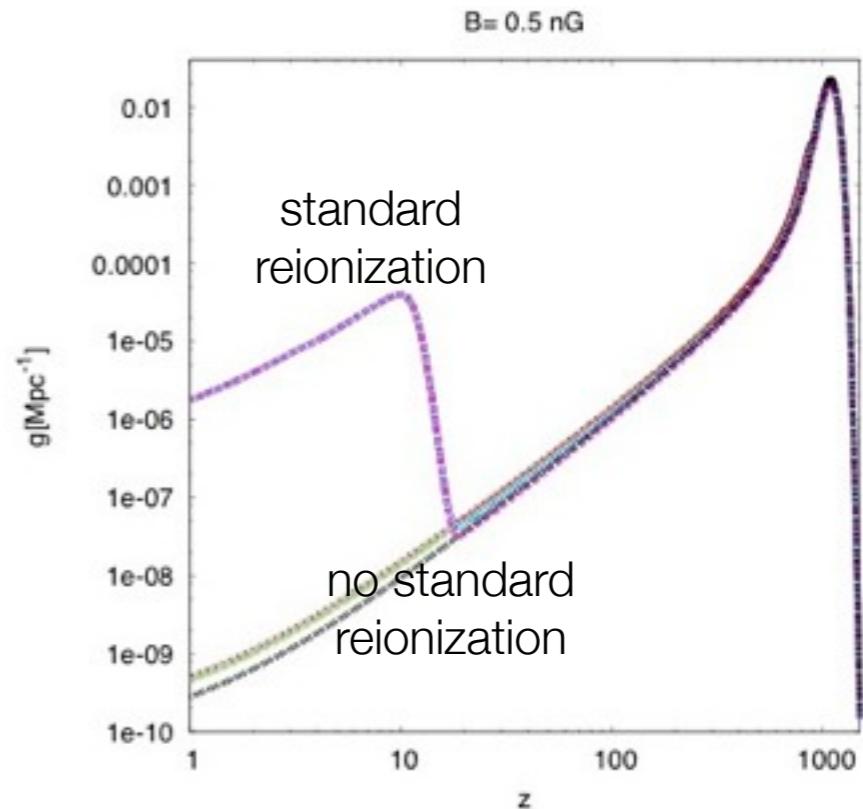
$$\boxed{S_{SW} = g \left(\frac{\delta_\gamma}{4} + \psi \right)}$$

optical depth

$$S_{ISW} = g (\phi - \psi) + e^{-\tau} 2\dot{\phi}$$

$$S_{dop} = k^{-2} \left(g \dot{\theta}_b + \dot{g} \theta_b \right)$$

$$S_{pol} = g P^{(0)}$$

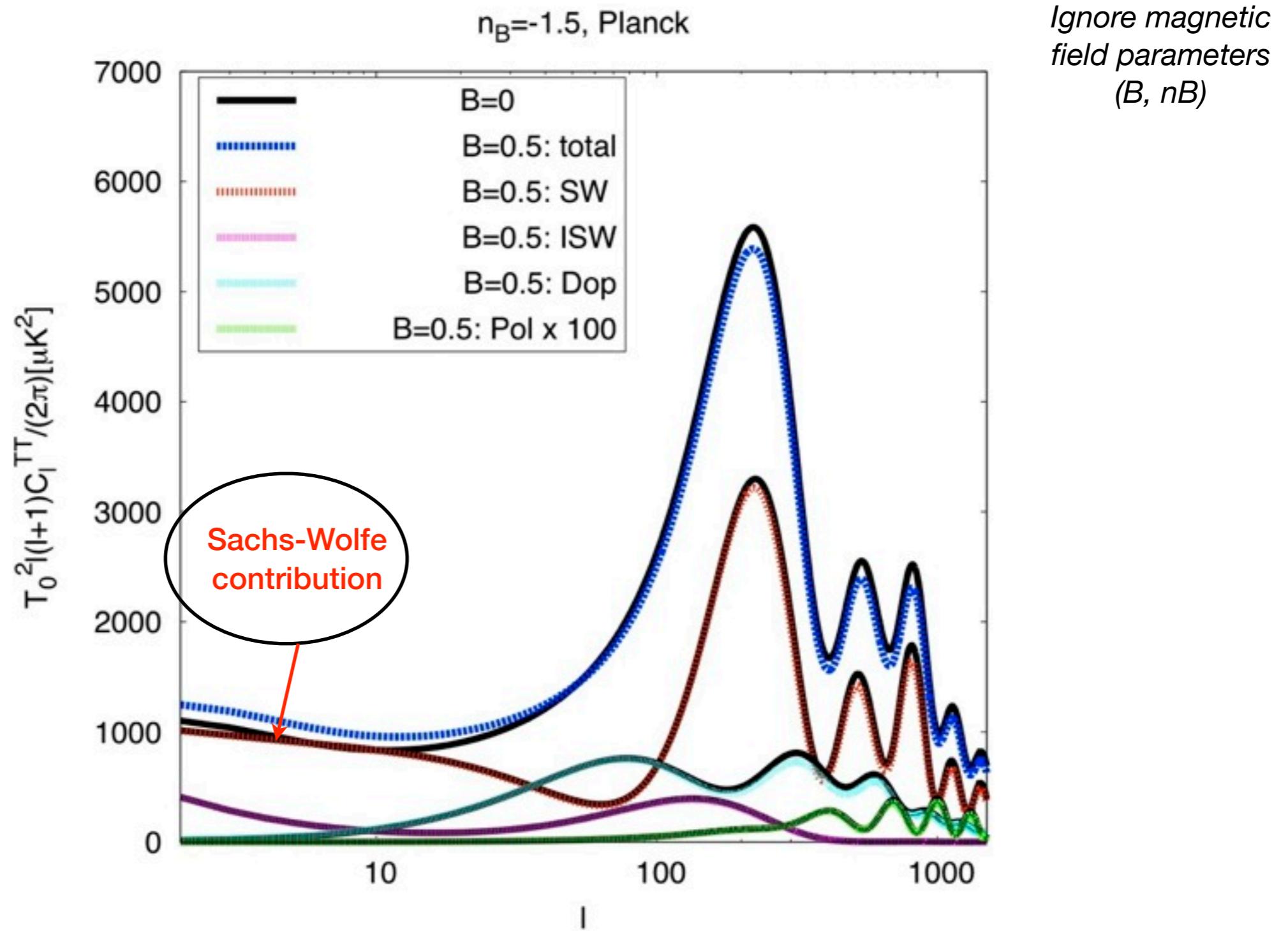


$$\tau(t) = \sigma_T \int_t^{t_0} n_e(t) dt$$

number density of
free electrons

$e^{-\tau(t)}$: probability that a CMB photon observed today has not scattered since time t.

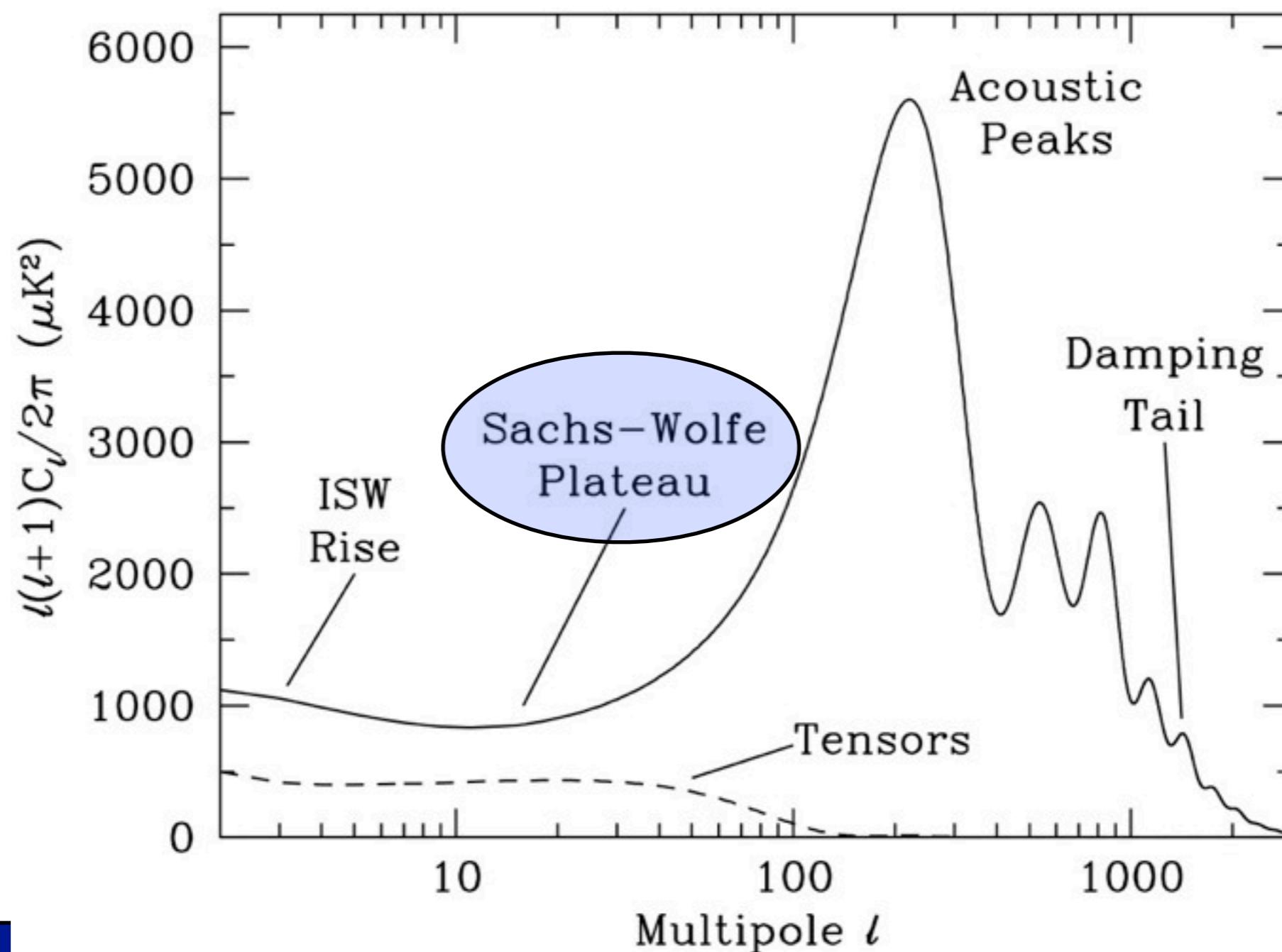
Cosmology 2



Cosmology 2

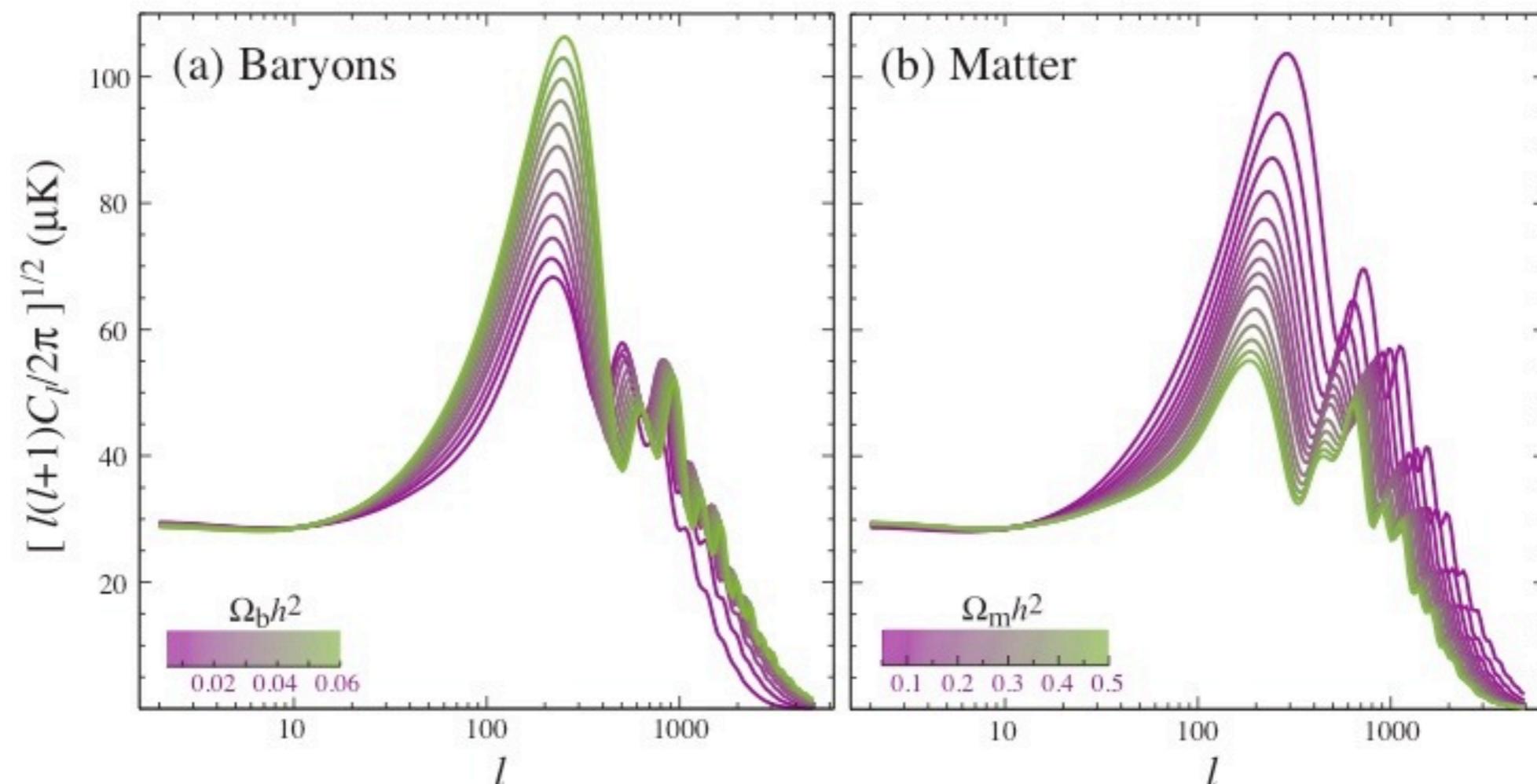
- Features in the angular power spectrum

(PDG 2014)



Cosmology 2

- Changing cosmological parameters



Hu (2008)

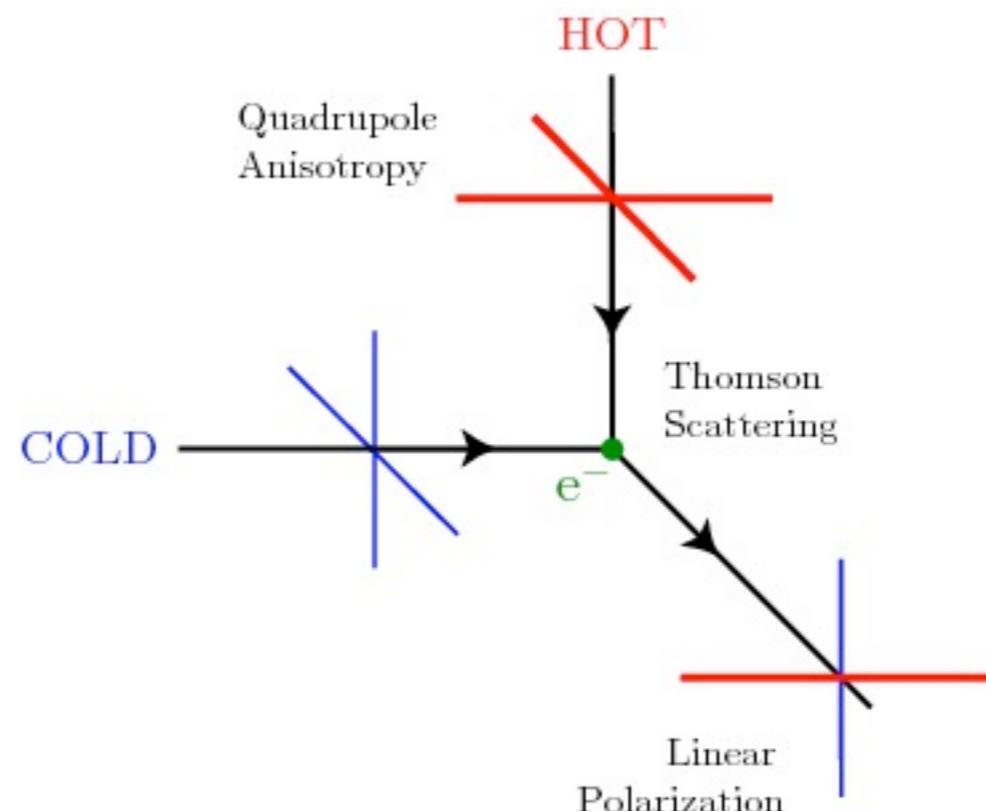
Cosmology 2

- What else do we see?

→ Polarization

Thomson scattering

- Isotropic radiation scatters into unpolarized radiation
- Radiation with quadrupole anisotropy scatters into linearly polarized radiation.

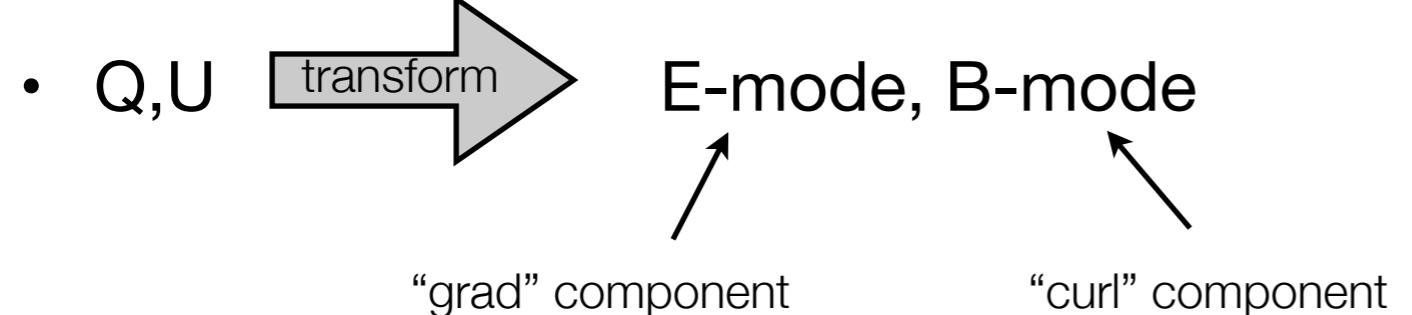


Hu,White

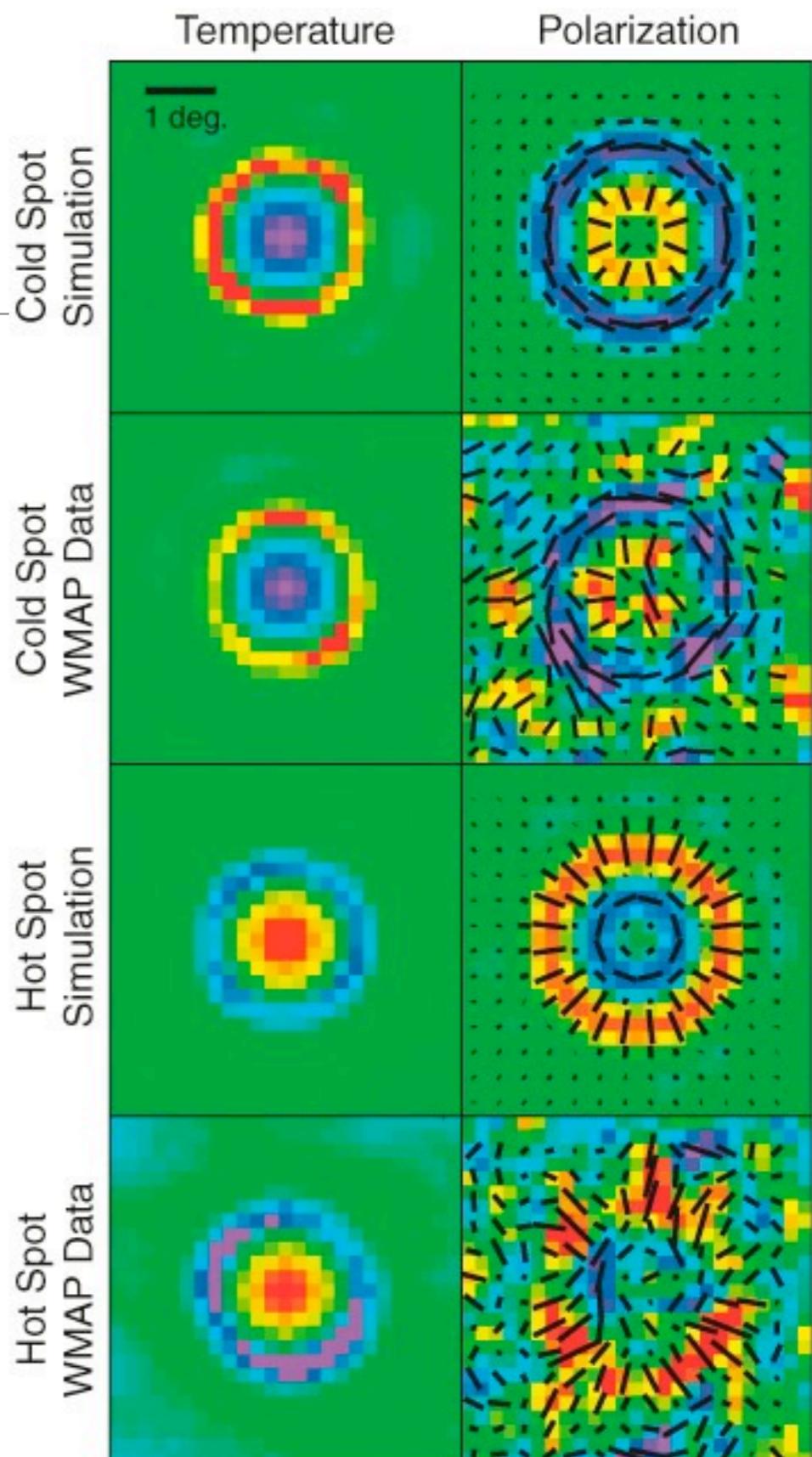
Cosmology 2

- To describe polarization use polarization tensor
- For linearly polarized radiation it is parametrized by the Stokes parameters:

I: intensity
Q, U : specify plane polarization

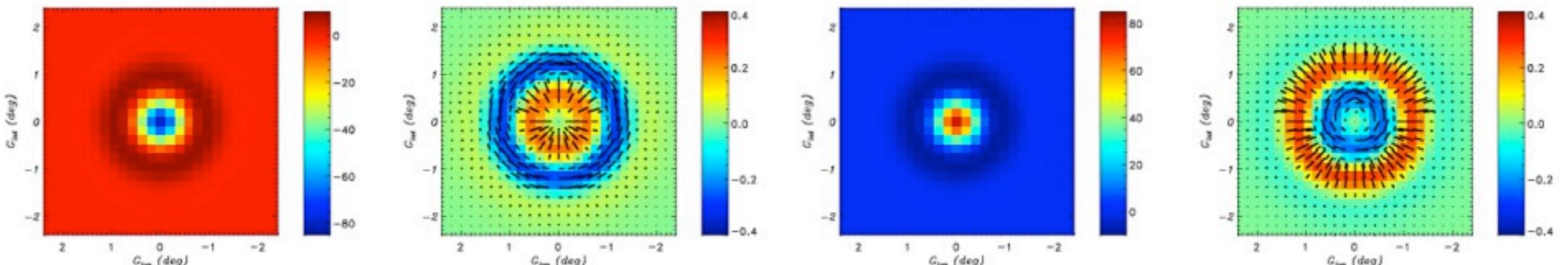
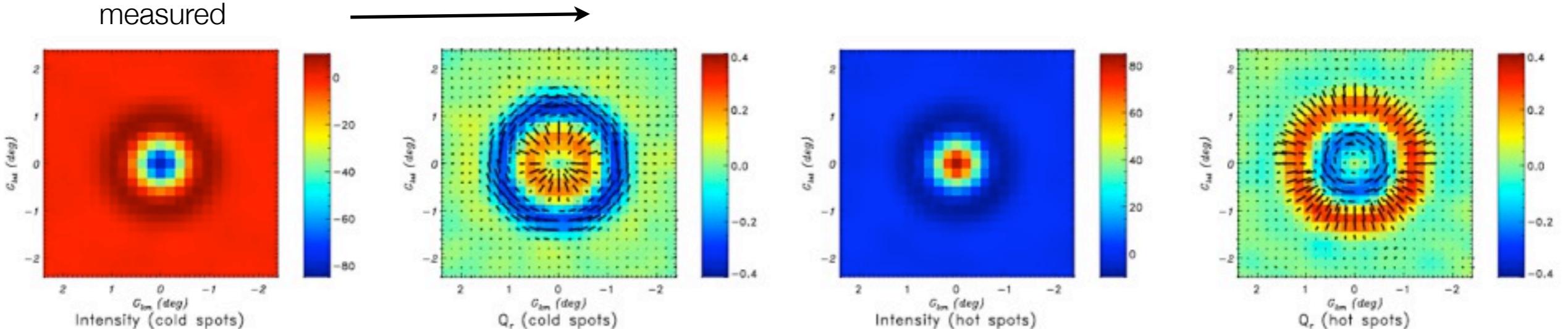


WMAP 7



Cosmology 2

measured

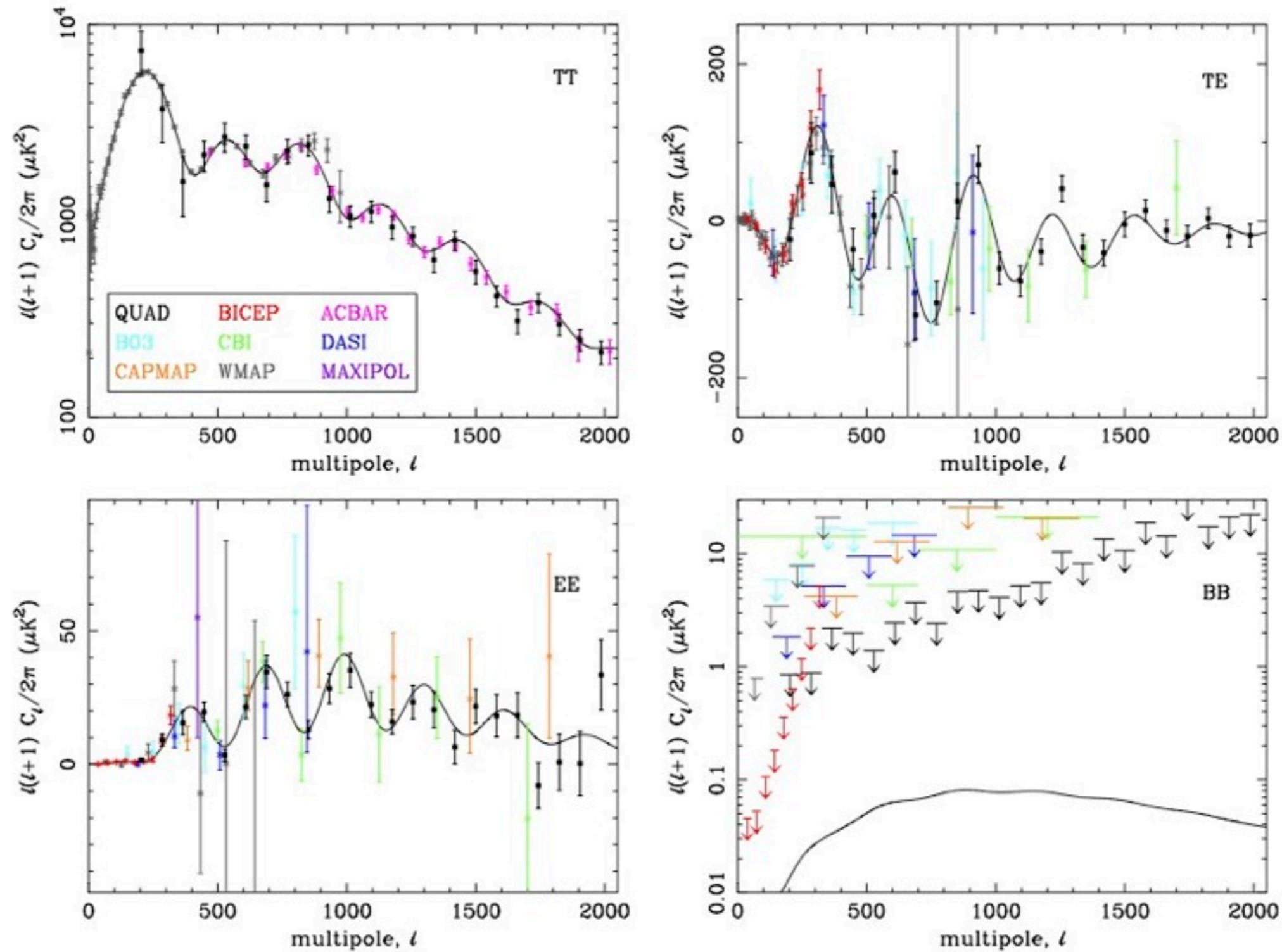


simulation



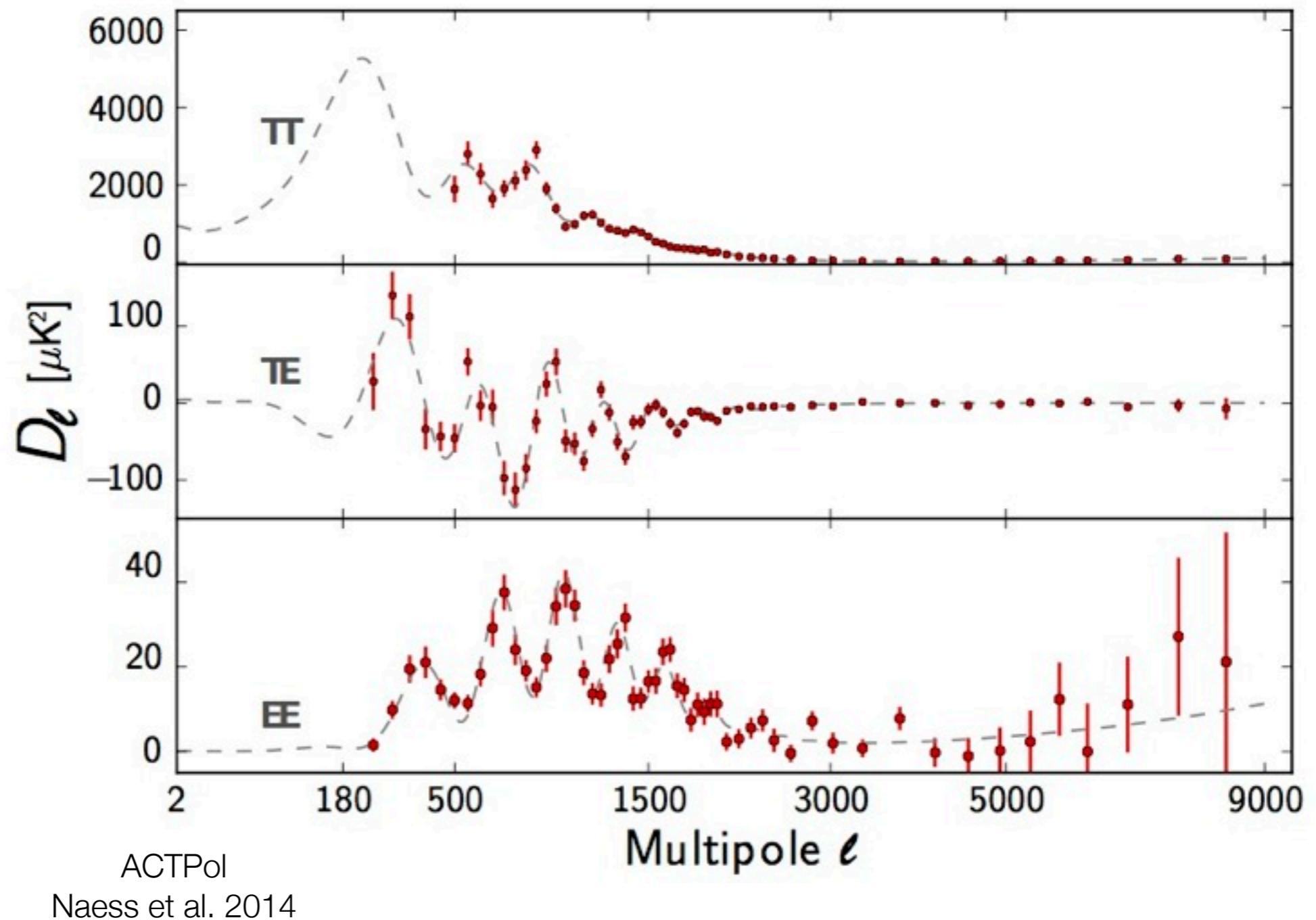
Planck13

Cosmology 2



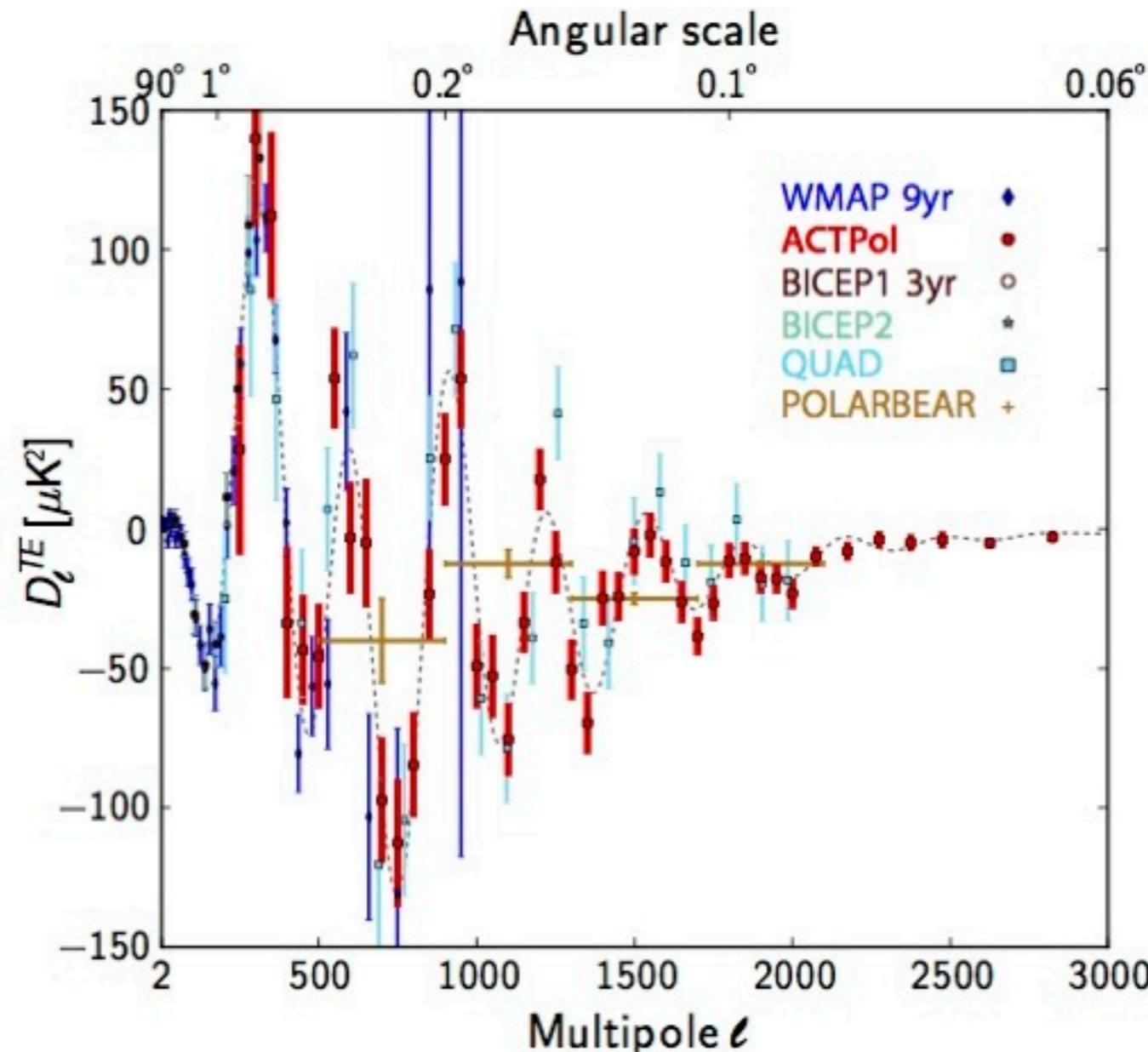
Cosmology 2

The **ACTPol** TT, TE, and EE power spectra, together with the best-fitting Λ CDM cosmological model and foreground components.



Cosmology 2

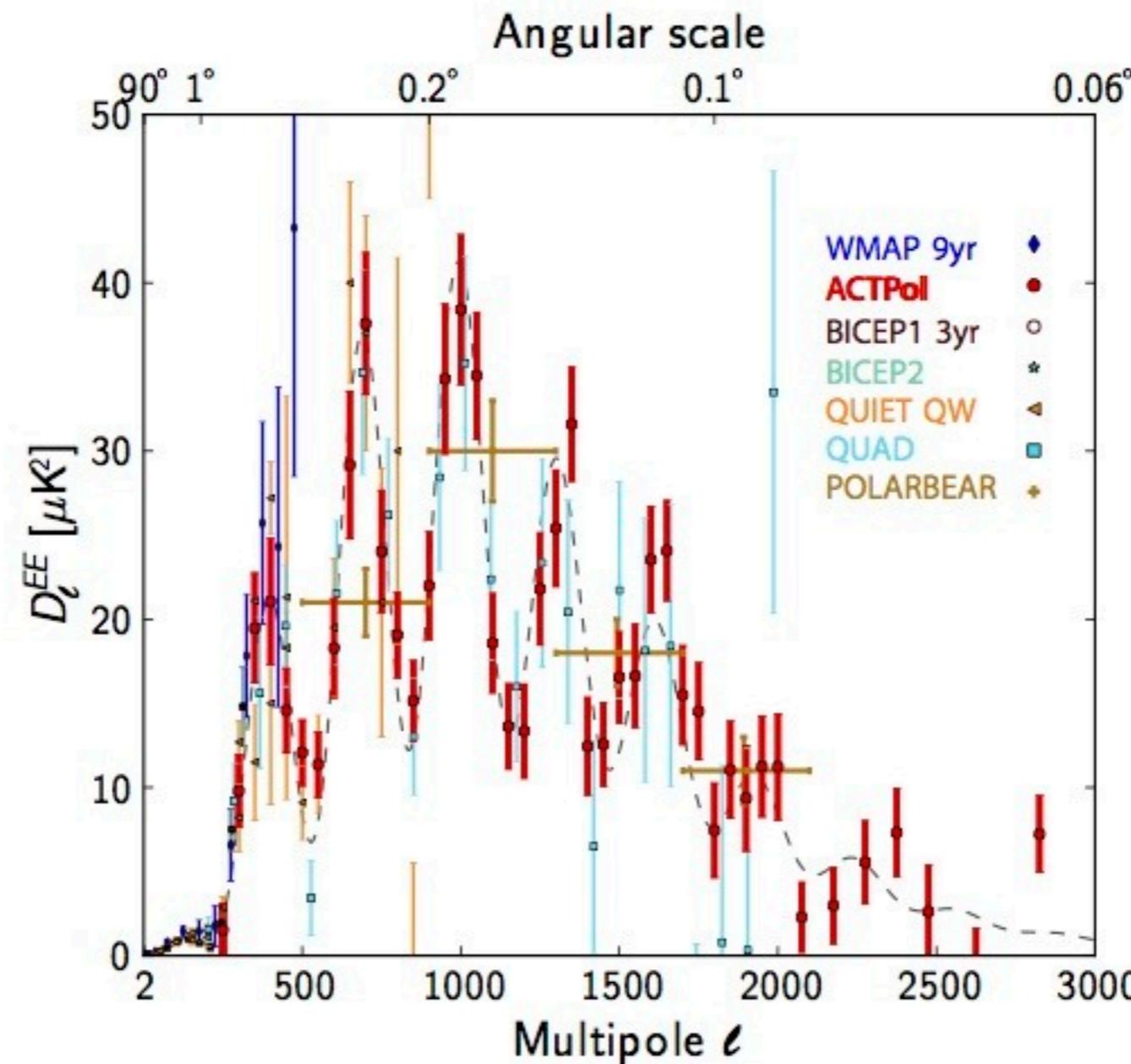
TE cross
correlation



ACTPol
Naess et al. 2014

Cosmology 2

EE auto
correlation

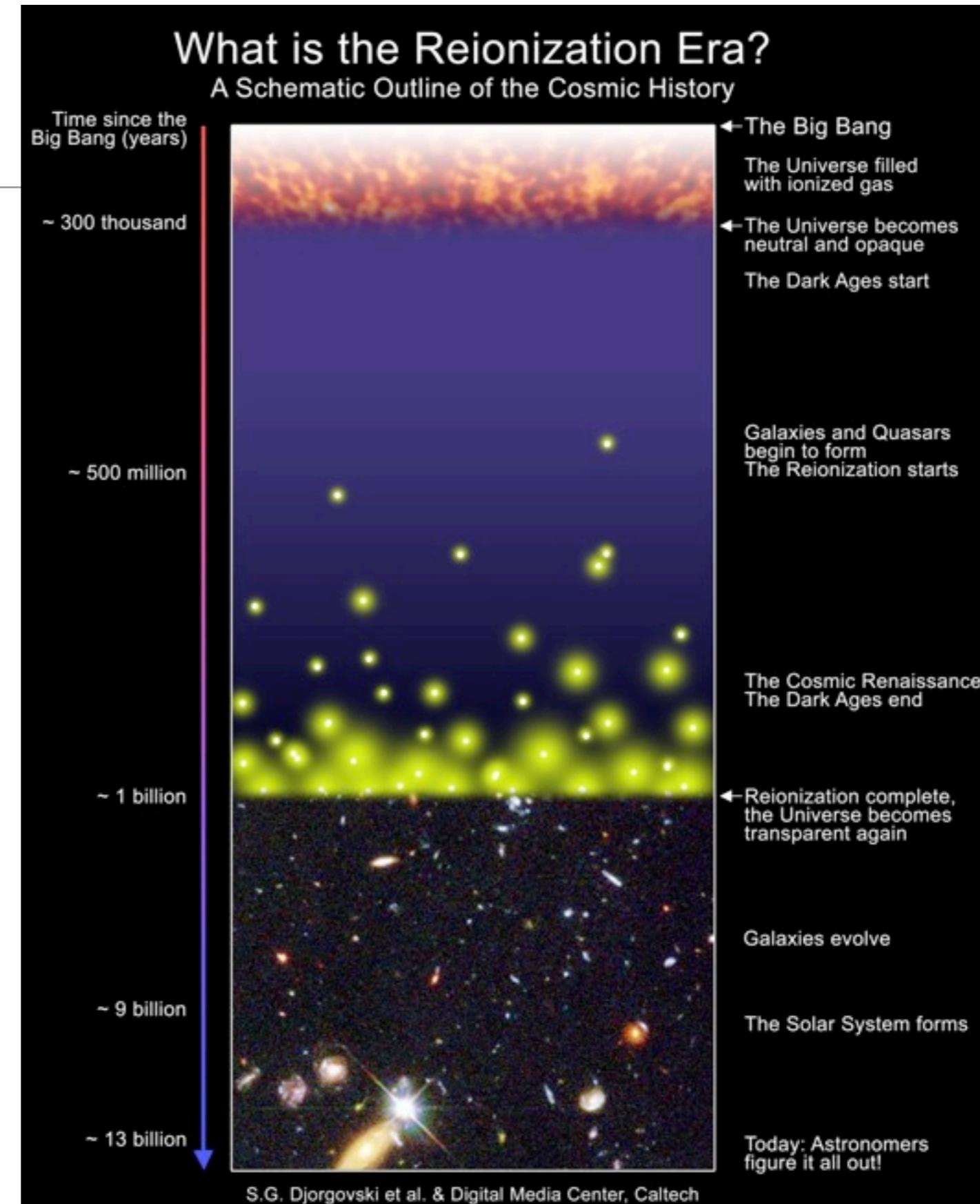


ACTPol
Naess et al. 2014

Cosmology 2

• Reionization

An interesting feature in the CMB polarization signal has been detected which is related to the more recent history of the universe: the universe is completely reionized today.



Caltech

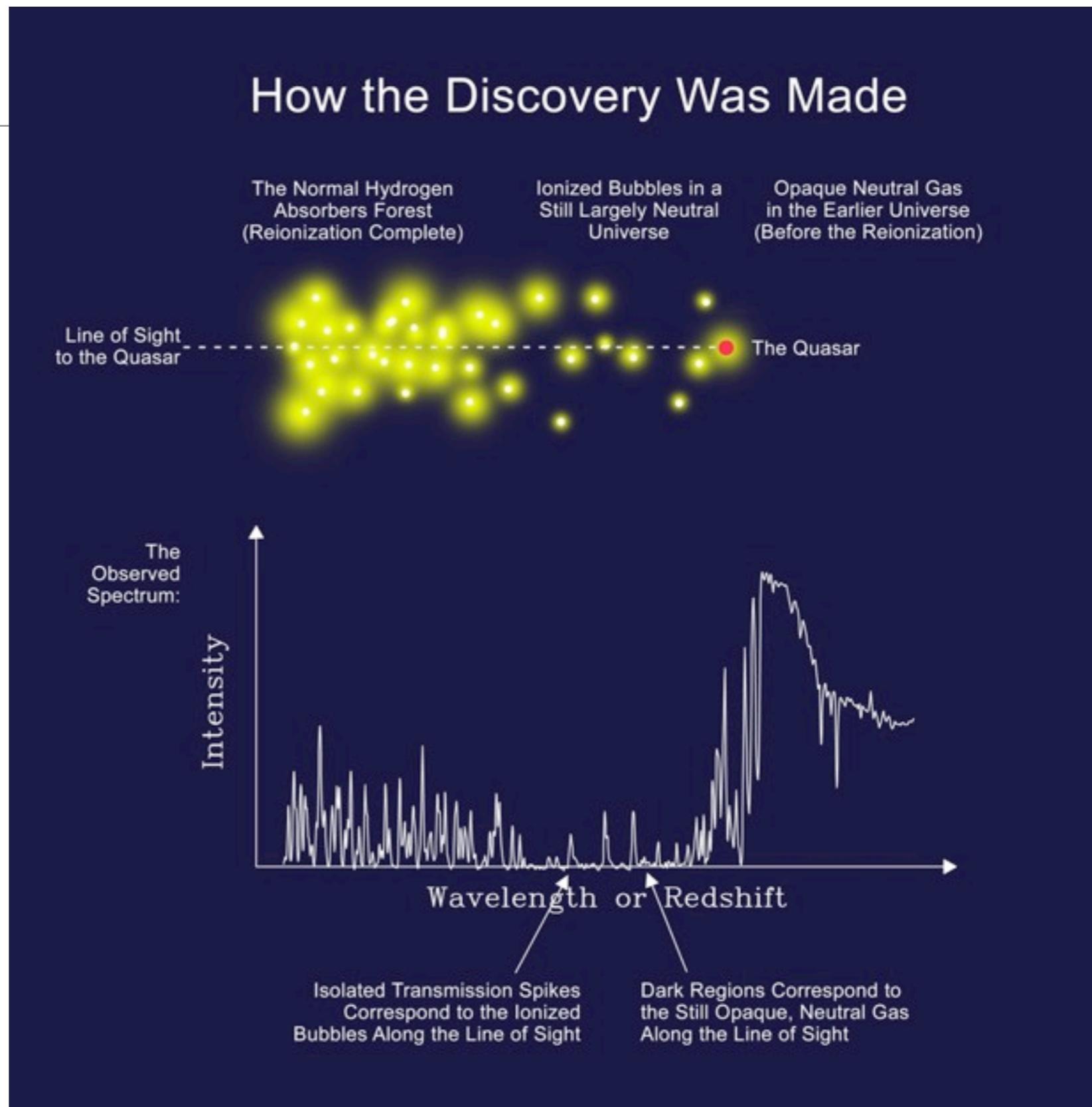
Cosmology 2

- Reionization

Evidence from
quasar spectrum

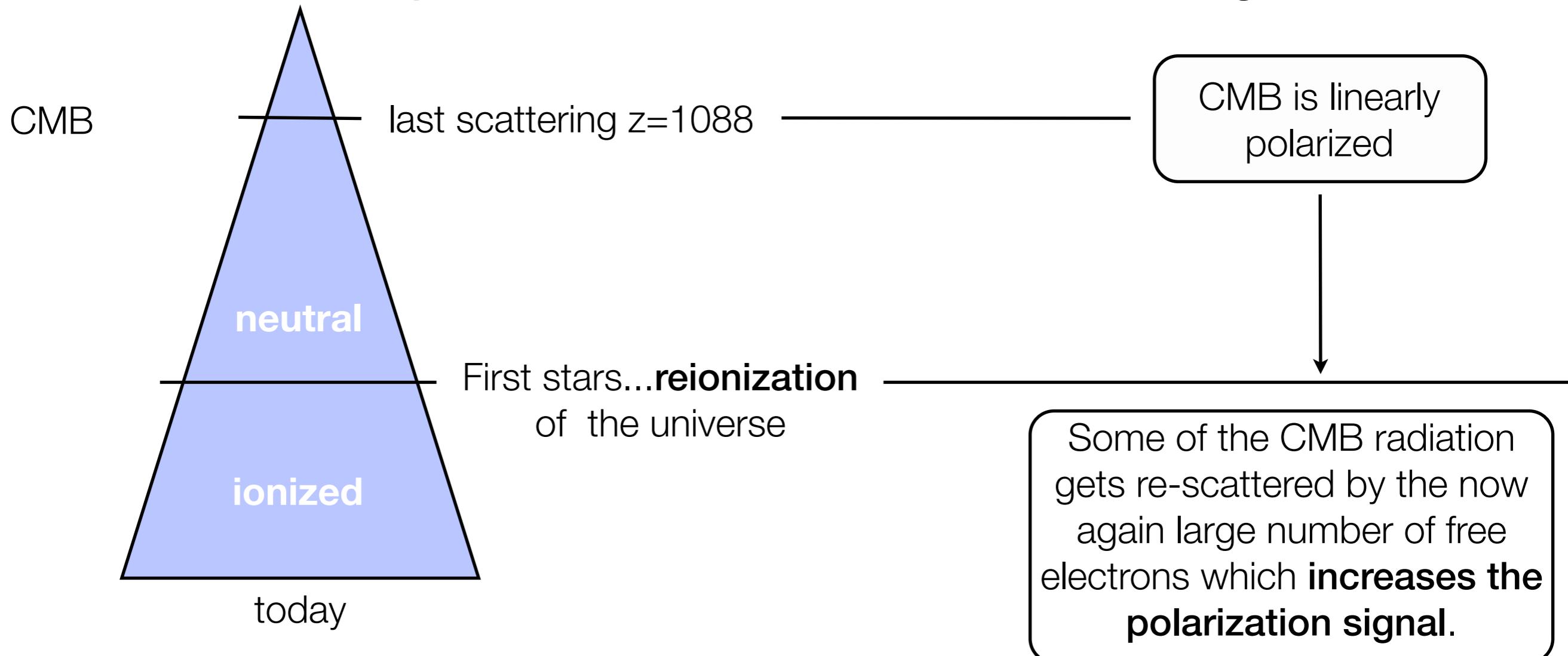
If neutral hydrogen
present, then
absorption at Ly α
transition. So if
there is
transmission,
hydrogen is
ionized.
(Gunn-Peterson)

How the Discovery Was Made



Cosmology 2

- Reionization
- Evidence from the *polarization of the cosmic microwave background*



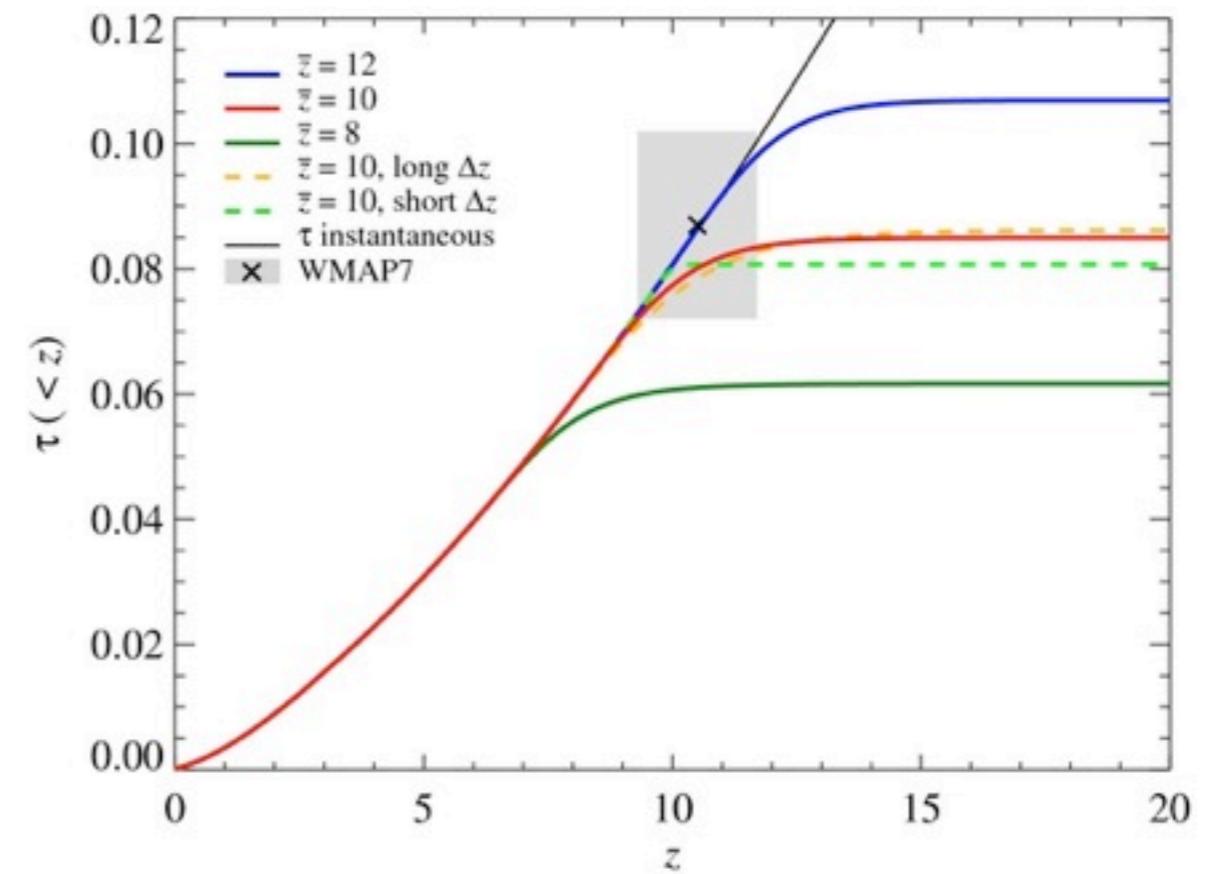
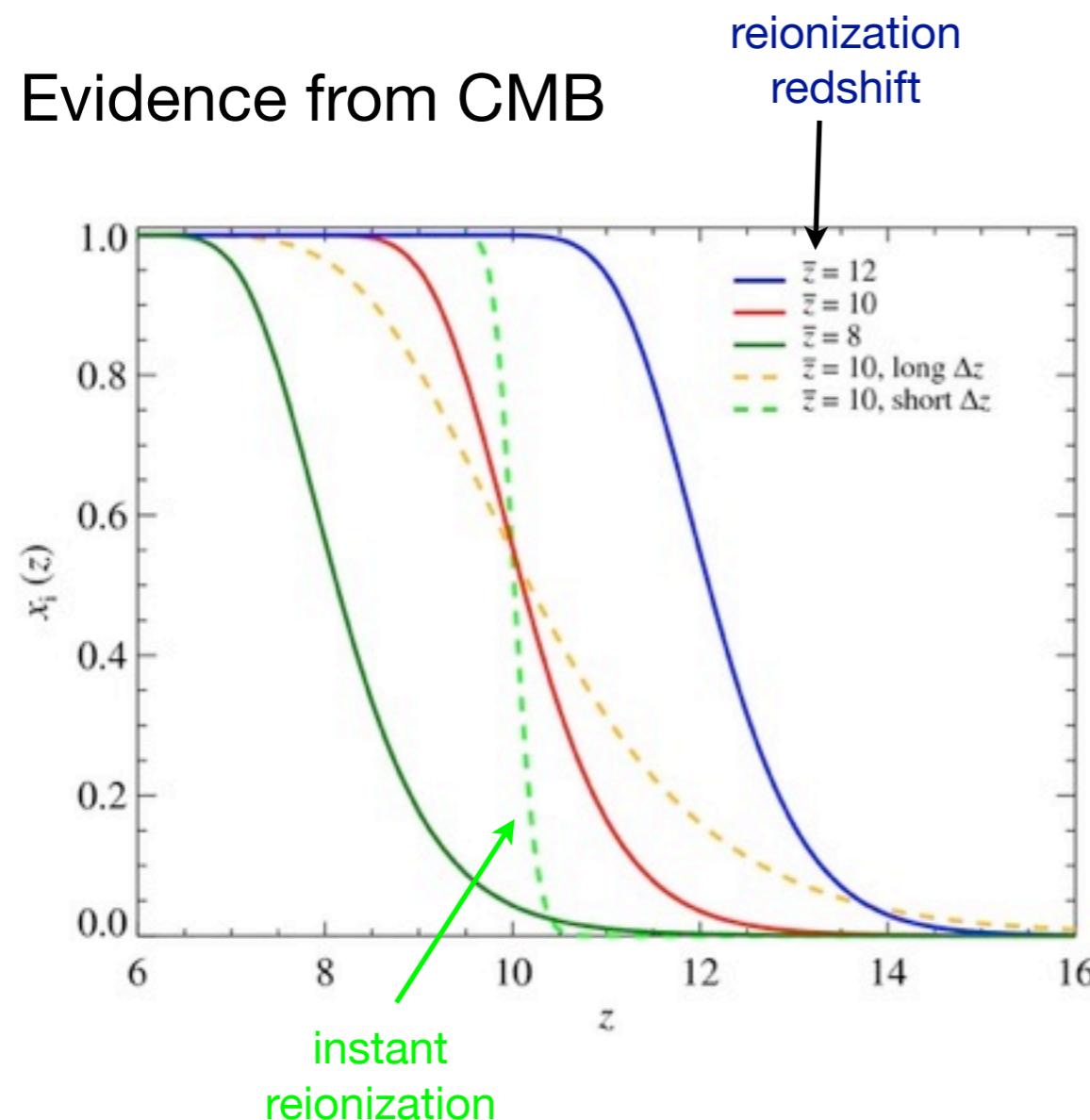
Cosmology 2

- The reionization generated polarization signal is redshift dependent:
- allows to constrain ionization history
- allows to determine *when* reionization took place

Cosmology 2

- Reionization
- Evidence from CMB

ionization history



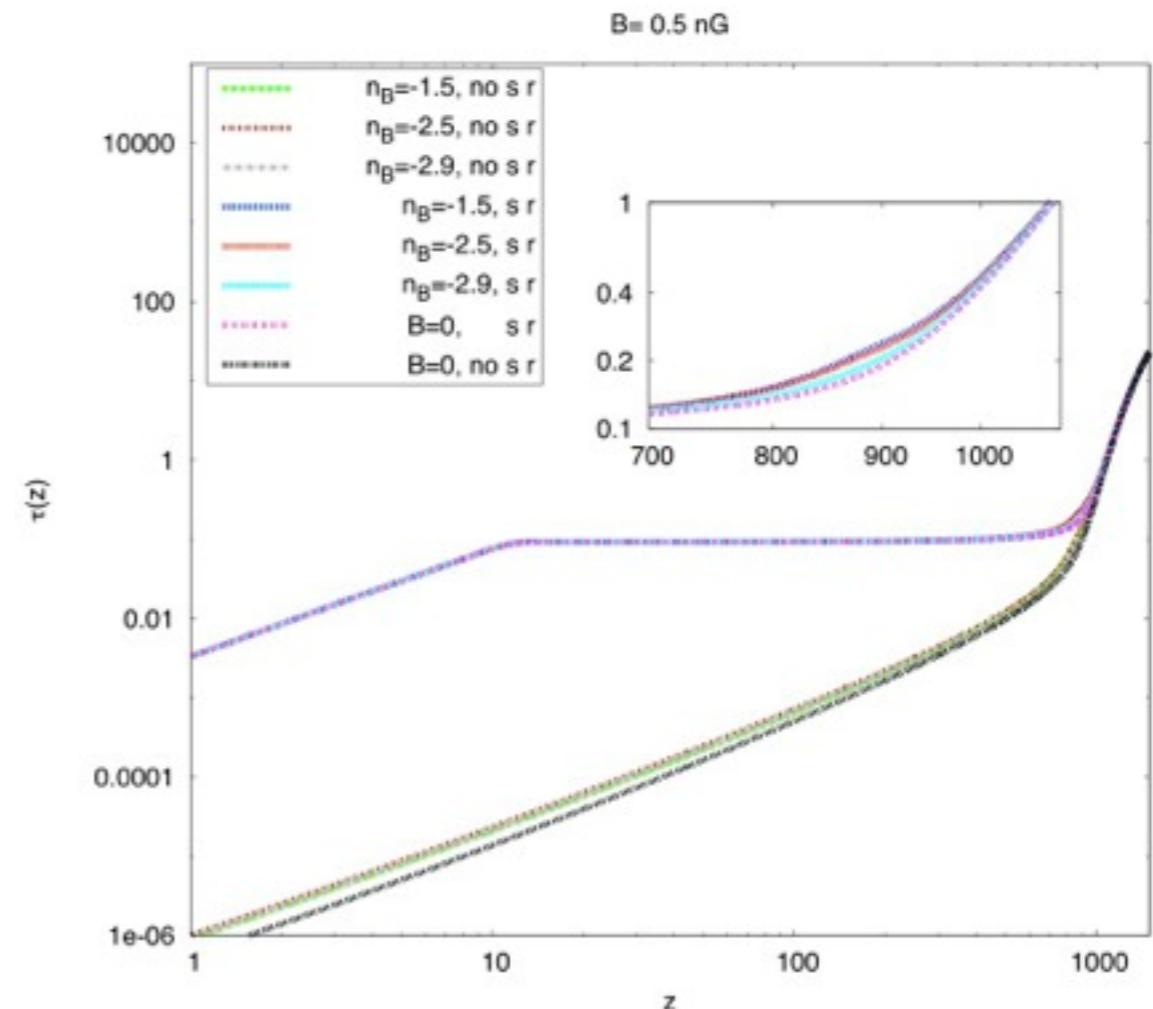
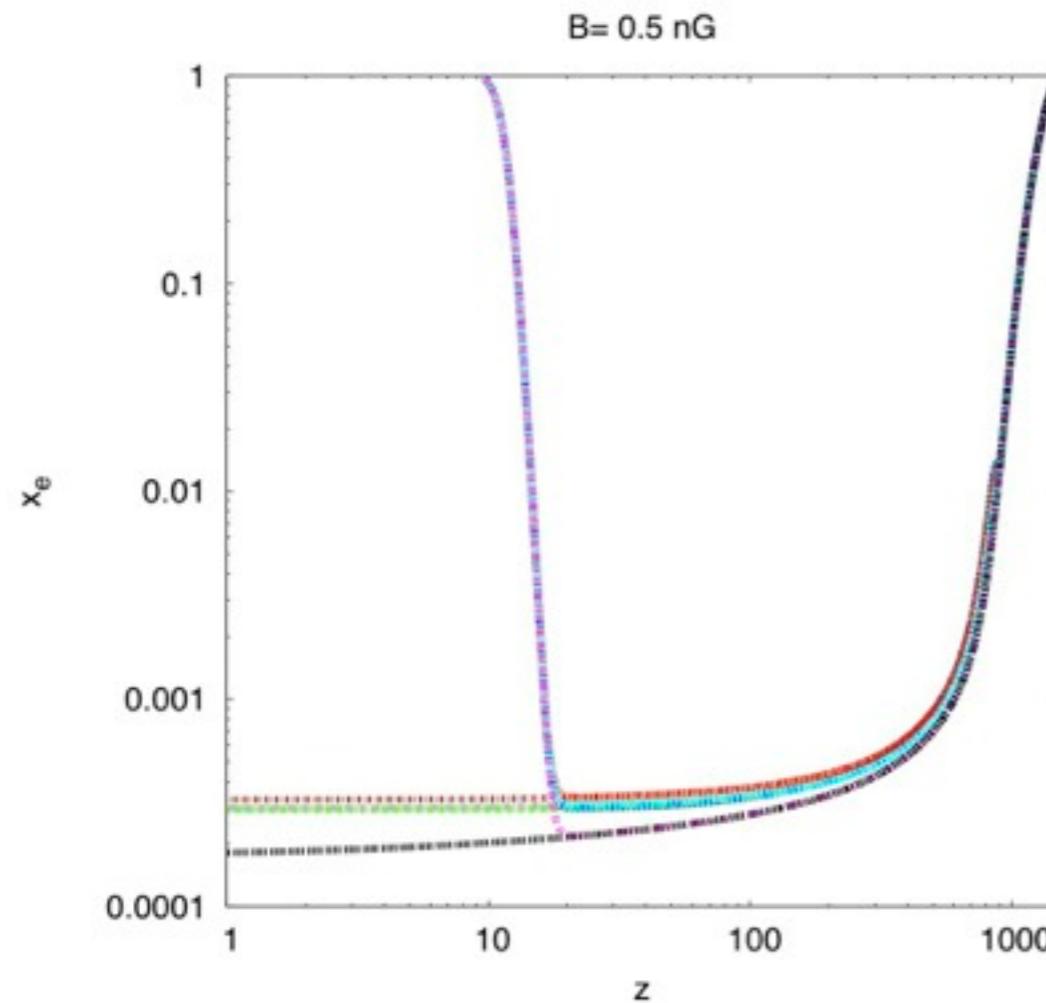
Battaglia et al. 2013

Cosmology 2

- Reionization

ionization history: from decoupling to close to the present

to put it in context (ignore magnetic field contributions...)

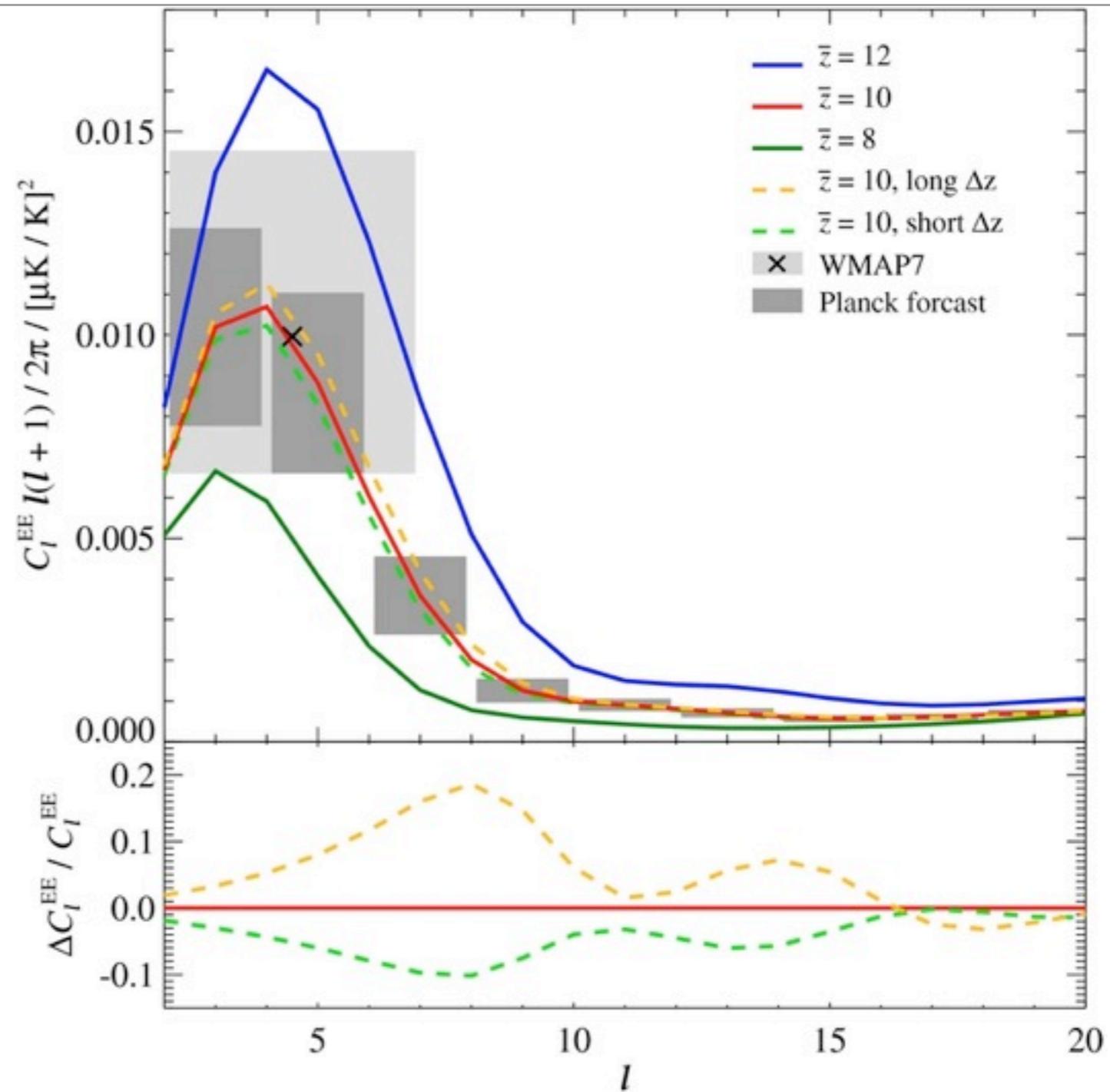


KEK, Komatsu 2015

Cosmology 2

- Reionization

low l EE polarization power spectrum



Battaglia et al. 2013

Cosmology 2

- CMB spectral distortions
- Spectrum well fitted by Planck black body spectrum

$$n_\nu = \left[\exp\left(\frac{h\nu}{kT}\right) - 1 \right]^{-1}$$

$$B_\nu(T) = \frac{2h\nu^3/c^2}{\exp\left(\frac{h\nu}{kT}\right) - 1}$$

FIXSEN ET AL. 1996

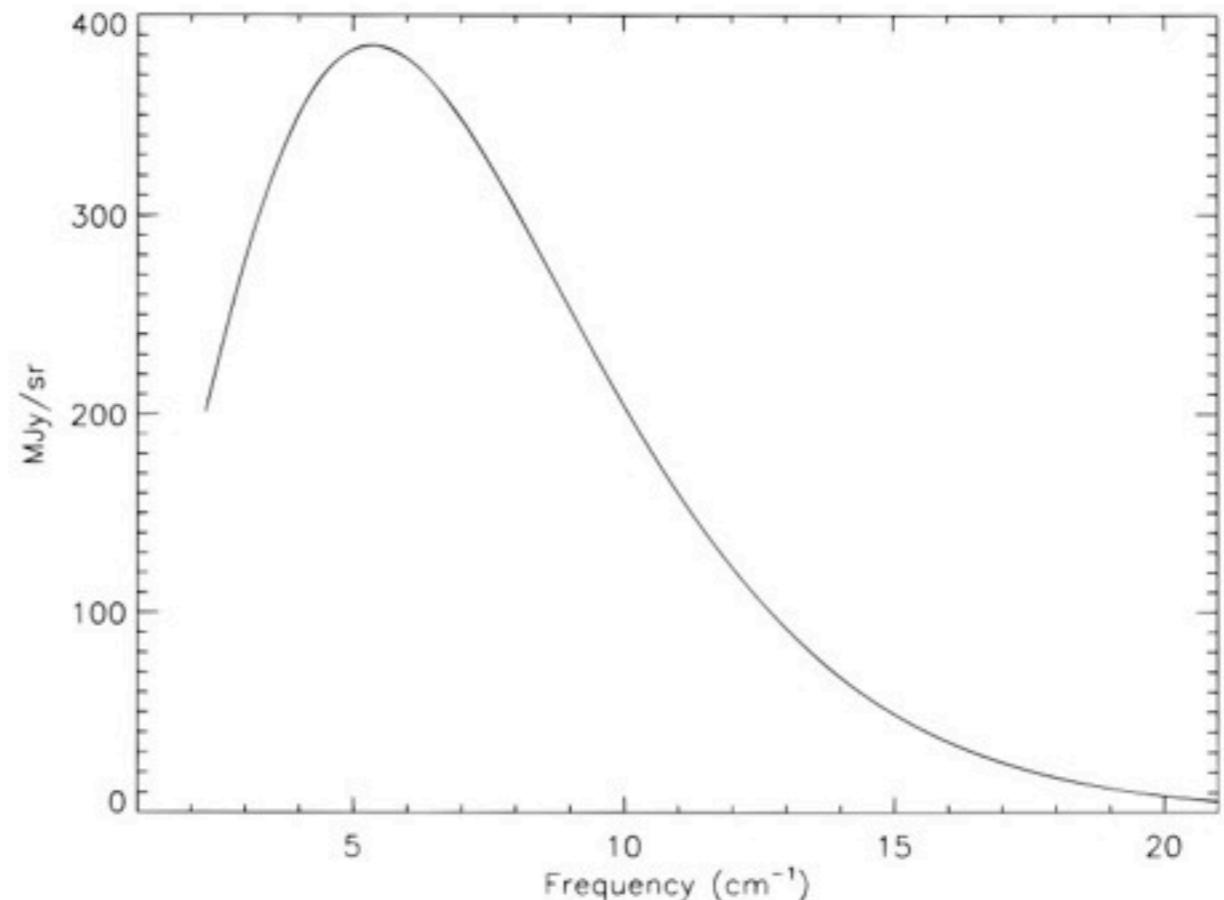
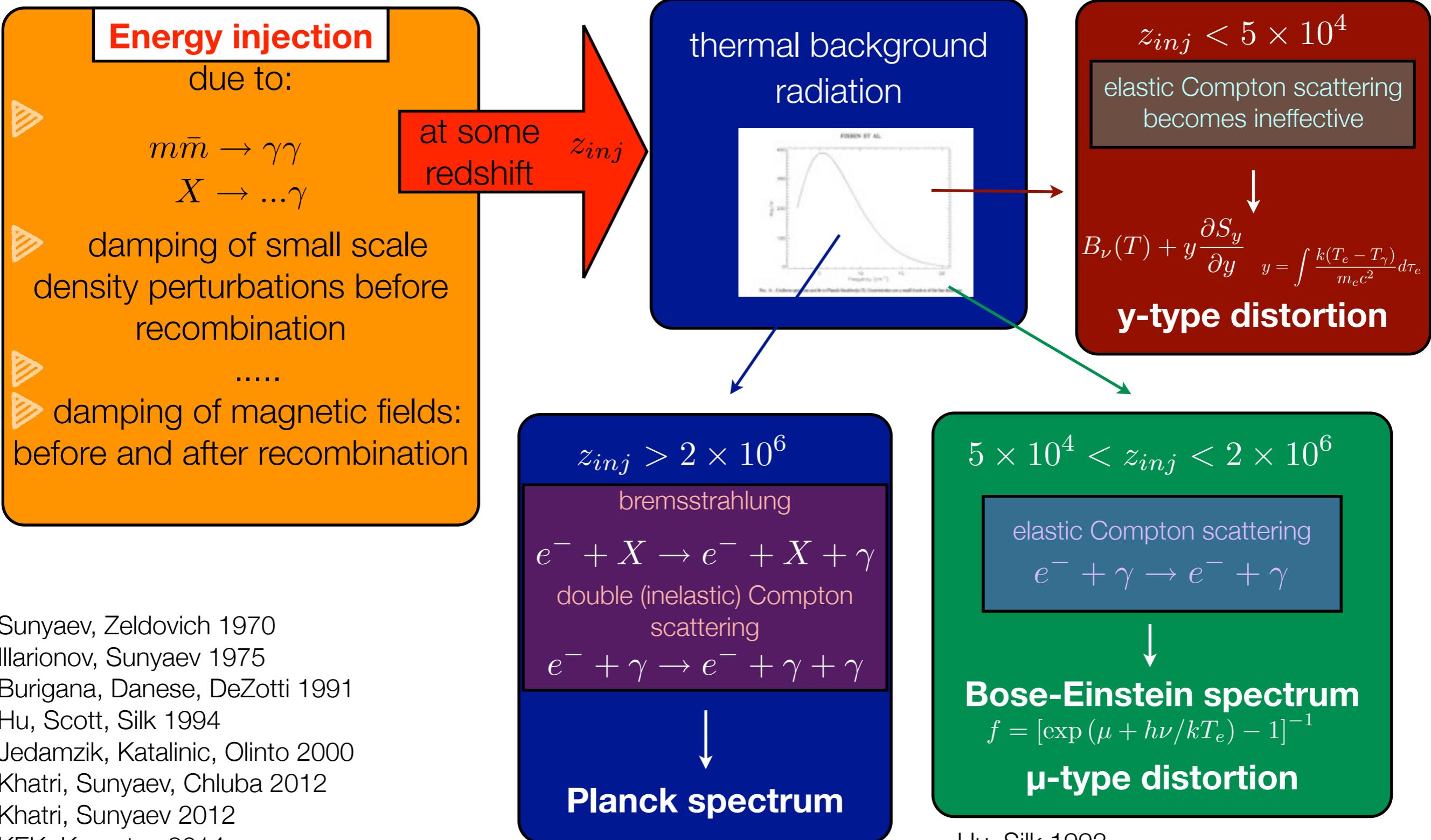


FIG. 4.—Uniform spectrum and fit to Planck blackbody (T). Uncertainties are a small fraction of the line thickness.

- Spectral distortions: γ - and μ -type are small

Cosmology 2



Sunyaev, Zeldovich 1970

Illarionov, Sunyaev 1975

Burigana, Danese, DeZotti 1991

Hu, Scott, Silk 1994

Jedamzik, Katalinic, Olinto 2000

Khatri, Sunyaev, Chluba 2012

Khatri, Sunyaev 2012

KEK, Komatsu 2014

Cosmology 2

- μ -type distortion

Hu, Silk 1993

$$\frac{d\mu}{dt} = -\frac{\mu}{t_{DC}(z)} + \frac{1.4}{3} \frac{dQ}{\rho_\gamma}$$


time scale for double
Compton scattering

$$t_{DC} = 2.06 \times 10^{33} \left(1 - \frac{Y_P}{2}\right)^{-1} (\Omega_b h^2)^{-1} z^{-\frac{9}{2}} \text{ s}$$

Pre-decoupling era

mixture of black body spectra:
1/3 of injected energy leads to
spectral distortions, 2/3 to raise
average temperature (Khatri,
Sunyaev, Chluba 2012)

Cosmology 2

- Observational limits

COBE FIRAS

Fixsen et al. 1996

$$|y| < 1.5 \times 10^{-5}$$

$$|\mu| < 9 \times 10^{-5}$$

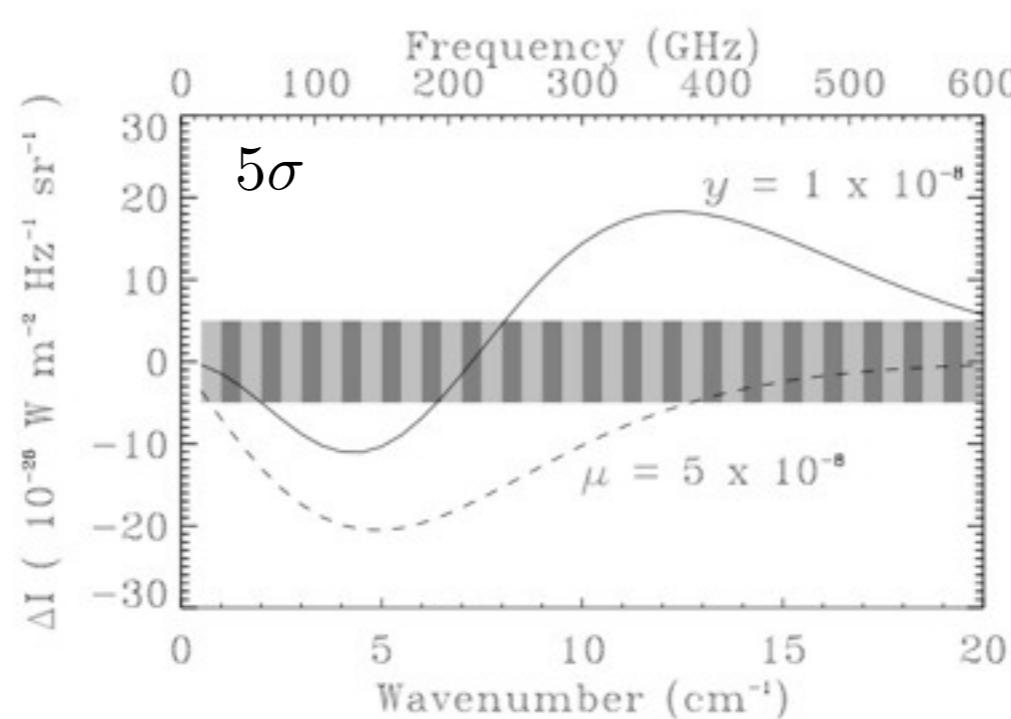
(95%CL)

- Future experiments

PIXIE

(Primordial Inflation Explorer)

Kogut et al. 2011

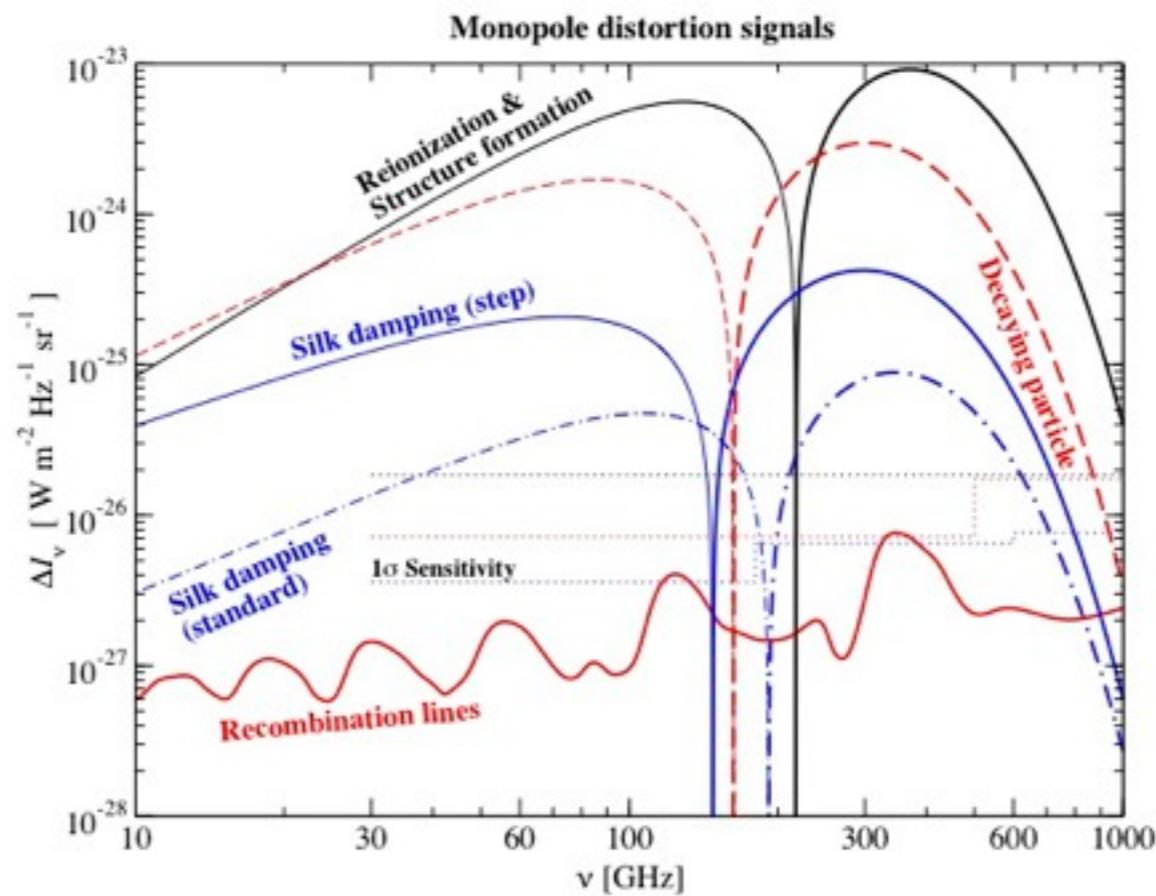


Cosmology 2

- Future experiments

PRISM (Polarized Radiation Imaging and Spectroscopy Mission)

Andre et al. 2013



Cosmology 2

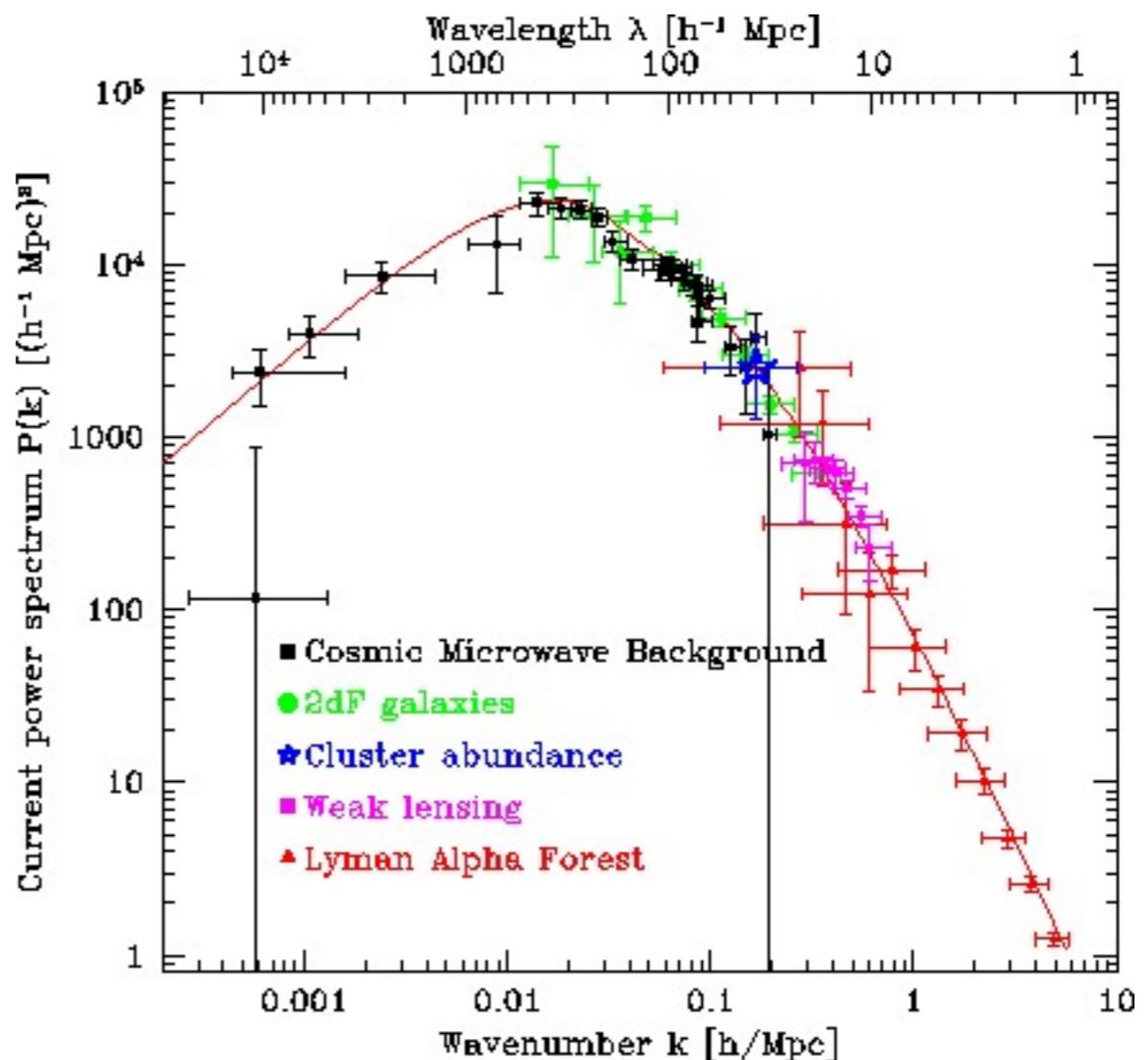
- The matter power spectrum

total matter perturbation

$$\Delta_m \equiv R_c \Delta_c + R_b \Delta_b$$
$$R_i \equiv \frac{\rho_i}{\rho_{matter}}$$

(during matter domination)

$$P_m(k) = \frac{k^3}{2\pi^2} |\Delta_m|^2$$



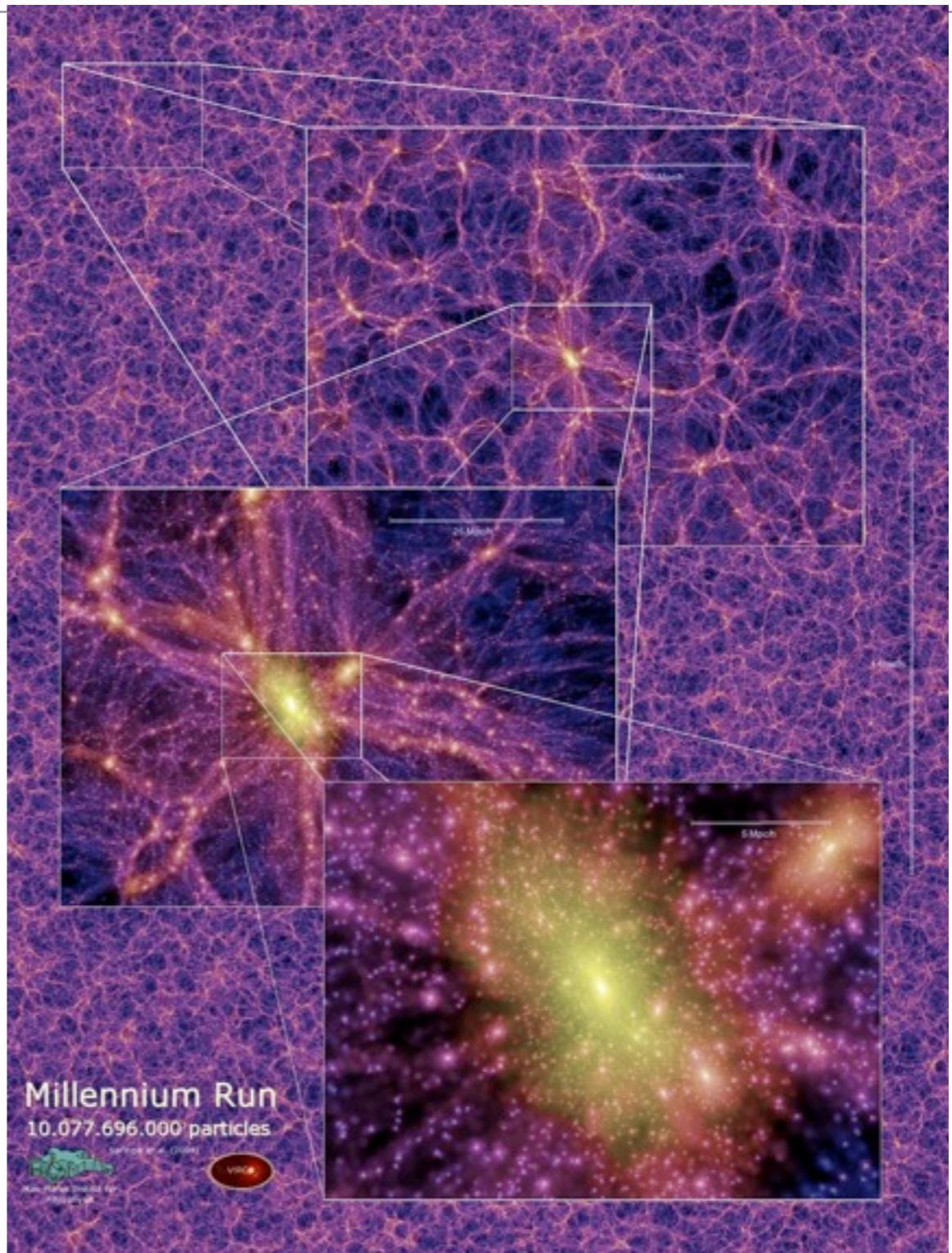
Tegmark (2002)

Cosmology 2

- On small scales: nonlinear effects become important
- numerical simulations
- analytical techniques....renormalization group approach

Cosmology 2

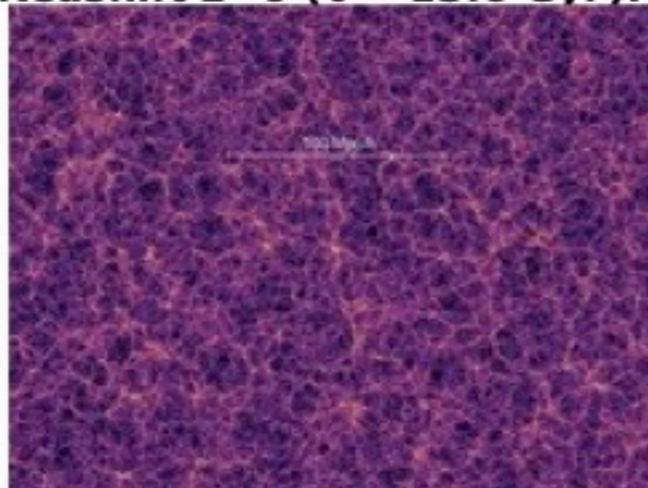
The *Millennium Run* used more than 10 billion particles to trace the evolution of the matter distribution in a cubic region of the Universe over 2 billion light-years on a side.



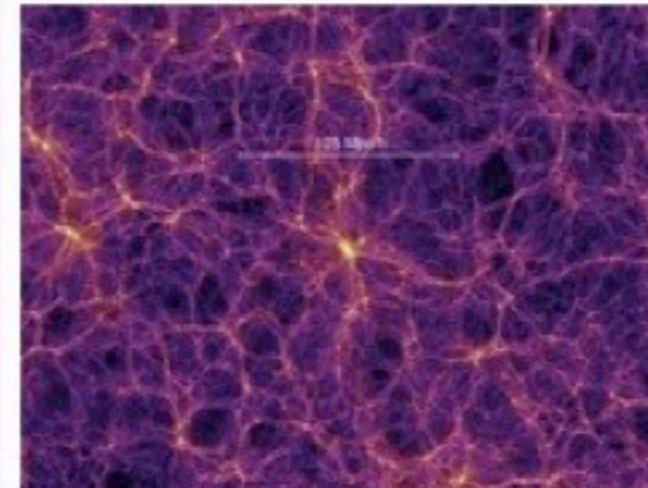
Cosmology 2

The following slices through the density field are all 15 Mpc/h thick. For each redshift, we show three panels. Subsequent panels zoom in by a factor of four with respect to the previous ones.

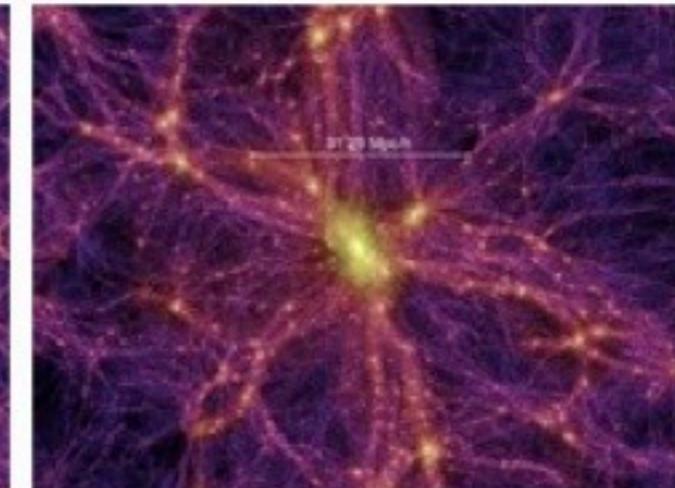
Redshift z=0 (t = 13.6 Gyr):



► 1024x768 ► 2048x1536

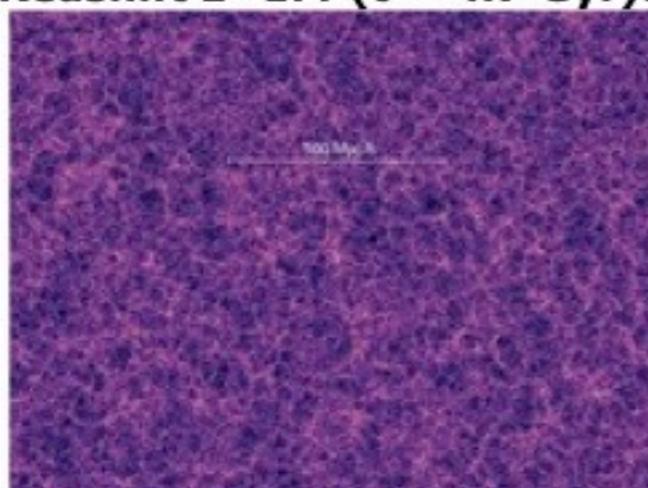


► 1024x768 ► 2048x1536

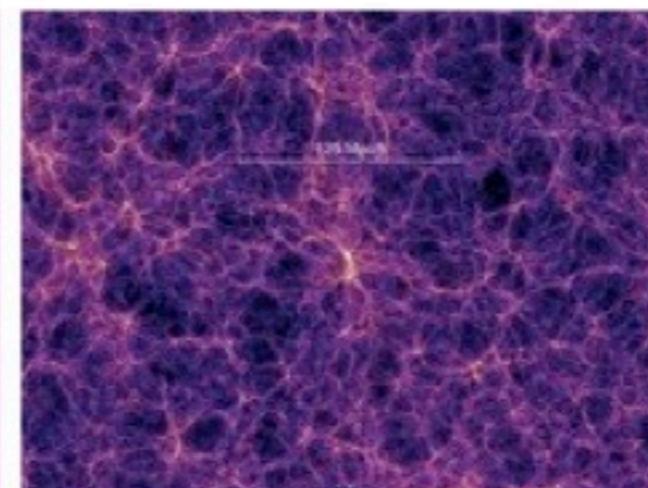


► 1024x768 ► 2048x1536

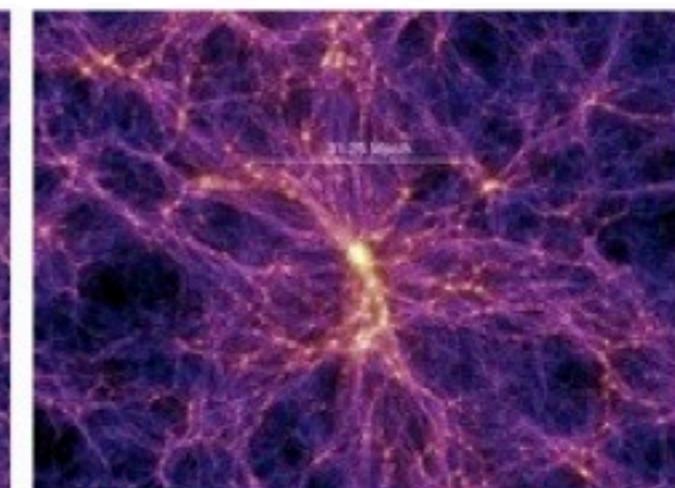
Redshift z=1.4 (t = 4.7 Gyr):



► 1024x768 ► 2048x1536



► 1024x768 ► 2048x1536



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The Millennium Simulation Project

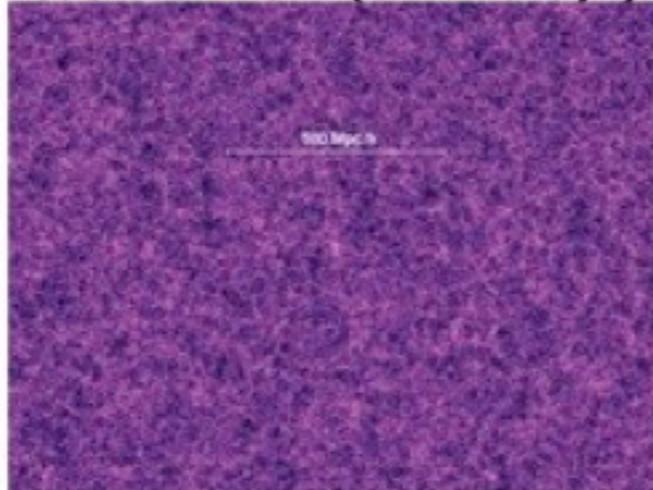


THE VIRGO CONSORTIUM

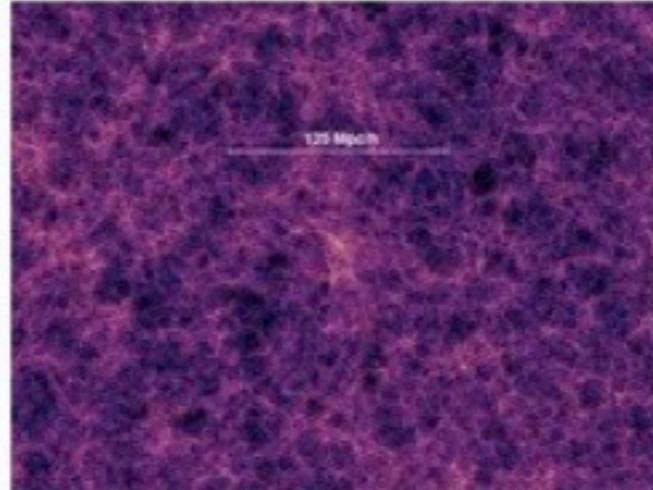
Cosmology 2

The Millennium Simulation Project

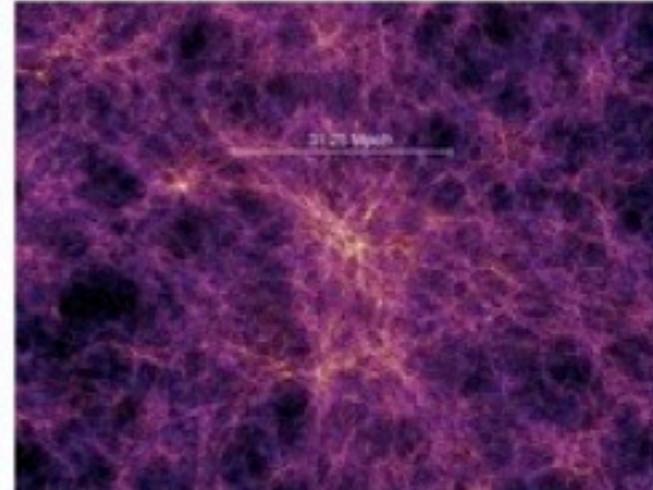
Redshift $z=5.7$ ($t = 1.0$ Gyr):



► 1024x768 ► 2048x1536

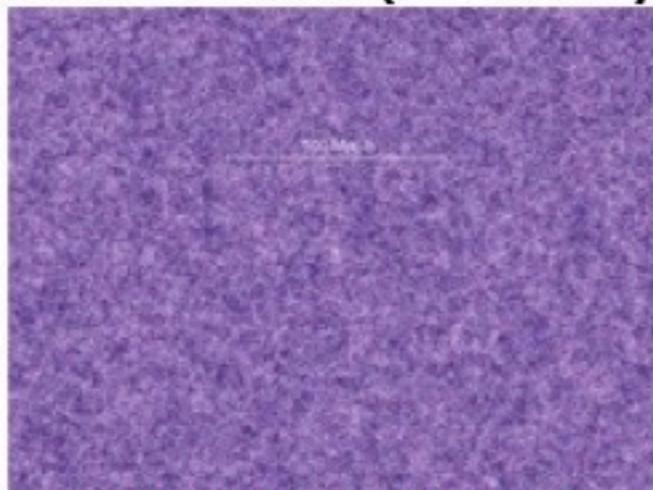


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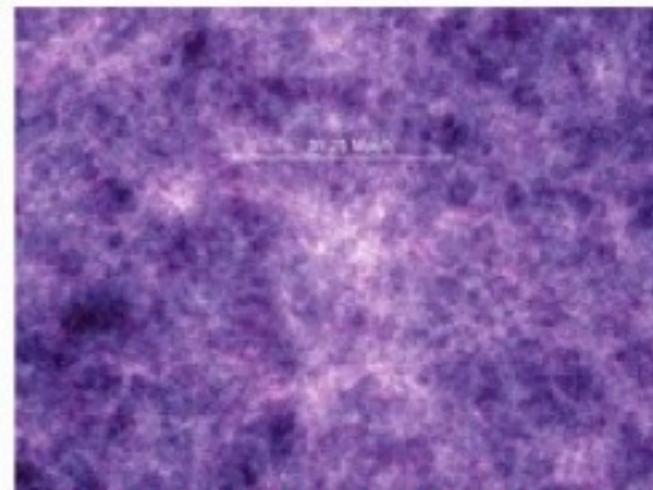
Redshift $z=18.3$ ($t = 0.21$ Gyr):



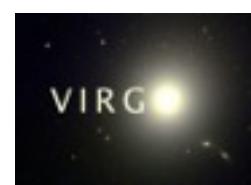
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THE VIRGO CONSORTIUM

<http://www.mpa-garching.mpg.de/galform/virgo/millennium/>