

How well can we constrain primordial non-Gaussianity with Large Scale Structure ?

Vincent Desjacques

Early Universe workshop, CERN, Jan 8, 2015

Local quadratic primordial NG:

$$
\Phi(\mathbf{x}) = \phi(\mathbf{x}) + f_{\rm NL} \phi^2(\mathbf{x})
$$

$$
|\phi| \sim 10^{-5}
$$

$$
\langle \Phi(\mathbf{k}_1) \Phi(\mathbf{k}_2) \Phi(\mathbf{k}_3) \rangle_c = (2\pi)^3 \delta_D(\mathbf{k}_1 + \mathbf{k}_2 + \mathbf{k}_3) \xi_{\Phi}^{(3)}(\mathbf{k}_1, \mathbf{k}_2, \mathbf{k}_3)
$$

$$
\xi_{\Phi}^{(3)}(\mathbf{k}_1, \mathbf{k}_2, \mathbf{k}_3) = 2f_{\rm NL} [P_{\phi}(k_1) P_{\phi}(k_2) + 2 \text{ cyc.}]
$$

$$
P_{\phi}(k) \propto k^{n_s - 4} \sim k^{-3}
$$

Outline

- *• Non-Gaussian bias as a probe of PNG*
- *• 3 universal truths about non-Gaussian bias*
- *• Peak theory: the not so universal truths*
- *• Implications for fNL constraints*

Galaxy clustering probes:

• Cluster counts

$$
S_3 \sim \int d^3 \mathbf{k}_1 \int d^3 \mathbf{k}_2 \, \xi_{\Phi}^{(3)}(\mathbf{k}_1, \mathbf{k}_2, -\mathbf{k}_1 - \mathbf{k}_2)
$$

• Galaxy power spectrum

$$
\Delta b_1(k) \sim \int d^3 \mathbf{k} \, \xi_{\Phi}^{(3)}(\mathbf{k}_1, -\mathbf{k}_1, \mathbf{k}), \quad S_3
$$

• Galaxy bispectrum

$$
\xi_{\Phi}^{(3)}(\mathbf{k}_1,\mathbf{k}_2,-\mathbf{k}_1-\mathbf{k}_2)\,,\cdots
$$

The non-Gaussian bias for local quadratic NG:

$$
P_{gg}(k) = (b_1 + \Delta b_1^{NG}(k))^2 P_{mm}(k)
$$

$$
\Delta b_1^{NG}(k) = \frac{2f_{NL}b_{NG}}{\mathcal{M}(k)} \sim \frac{2f_{NL}b_{NG}}{k^2}
$$

Most recent constraints:

$-40 \le f_{\rm NL} \le +40$ (95\% C.L.)

 Giannantonio et al. '12, Ho et al. '13, Leistedt et al. '14

Forecast for a Euclid-like survey :

 $\Delta f_{\rm NL} \sim 3$

 Giannantonio et al. '12

Peak-background split (PBS):

 $\delta(\mathbf{x}) = \delta_l(\mathbf{x}) + \delta_s(\mathbf{x})$

Kaiser '84; Bardeen et al. '86; Cole & Kaiser '89; Mo & White '96; Sheth & Tormen '99; ...

Thursday, 8 January 15

$$
\delta_{\rm g}(\mathbf{x}) \equiv \frac{\bar{n}_{\rm g}(\mathbf{x})}{\bar{n}_{\rm g}} - 1 = \frac{\bar{n}_{\rm g}(\delta_c - \delta_l(\mathbf{x}), \sigma)}{\bar{n}_{\rm g}(\delta_c, \sigma)} - 1 = \left(-\frac{1}{\bar{n}_{\rm g}}\frac{d\bar{n}_{\rm g}}{d\delta_c}\right)\delta_l(\mathbf{x}) + \dots
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PBS: non-Gaussian bias

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\Phi(\mathbf{x}) = \phi(\mathbf{x}) + f_{\text{NL}}\phi(\mathbf{x})^2
$$

$$
\phi(\mathbf{x}) = \phi_l(\mathbf{x}) + \phi_s(\mathbf{x})
$$

$$
\Phi = (\phi_l + f_{\text{NL}}\phi_l^2) + \phi_s(1 + 2f_{\text{NL}}\phi_l) + f_{\text{NL}}\phi_s^2
$$

$$
\downarrow
$$

$$
\sigma_s \to \sigma_s(1 + 2f_{\text{NL}}\phi_l)
$$

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on-Gaussian bias

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\delta_{l}(\mathbf{x}) = b_{1}
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PBS: non-Gaussian bias

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\delta_{g}(\mathbf{x}) = \left(-\frac{1}{\bar{n}_{g}}\frac{d\bar{n}_{g}}{d\delta_{c}}\right)\delta_{l}(\mathbf{x}) + \left(\frac{1}{\bar{n}_{g}}\frac{d\bar{n}_{g}}{d\sigma}\right)2f_{\rm NL}\sigma\phi_{l}(\mathbf{x}) + \dots
$$

$$
= b_{1}\delta_{l}(\mathbf{x}) + 2f_{\rm NL}\left(\frac{\partial\ln\bar{n}_{g}}{\partial\ln\sigma}\right)\phi_{l}(\mathbf{x}) + \dots
$$

Slosar et al. '08

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 iff $\bar{n}_{\rm g} \equiv \bar{n}_{\rm g} (\delta_c / \sigma)$

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Slosar et al. '08

$$
\text{Local bias approach:} \qquad \delta_{\mathrm{g}}(\mathbf{x}) = b_1 \delta_l(\mathbf{x}) + \frac{1}{2} b_2 \delta_l^2(\mathbf{x}) + \dots
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$$

g(x) = *^b*1*l*(x)+2*f*NL²*b*2*l*(x) + *... Taruya et al. '08*

The following three statements are certainly true:

• The amplitude of the non-Gaussian bias is given by

$$
\left(\frac{\partial\ln\bar{n}_\mathrm{g}}{\partial\ln\sigma_8}\right)
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• $\delta_c b_1$ is equal to $\left(\frac{\partial \ln \bar{n}_{\rm g}}{\partial \ln \sigma_8}\right)$ only if the halo mass function is universal ◆ $\delta_c b_1$

• Local bias expansions cannot correctly predict the amplitude of non-Gaussian bias

VD, Seljak & Iliev '09

Hamaus, Seljak & VD '11

Toy model: maxima of the linear density field as a proxy for the formation sites *of dark matter haloes (Peacock & Heavens '85; BBKS '86)*

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\nu(\mathbf{x}) \equiv \frac{1}{\sigma_0} \delta_s(\mathbf{x}), \qquad \eta_i(\mathbf{x}) = \frac{1}{\sigma_1} \partial_i \delta_s(\mathbf{x}), \qquad \zeta_{ij} \equiv \frac{1}{\sigma_2} \partial_i \partial_j \delta_s(\mathbf{x})
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\sigma_n^2 = \frac{1}{2\pi^2} \int_0^\infty dk \, k^{2(n+1)} P_s(k)
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Peak constraint:

(i)
$$
\nu(\mathbf{x}_p) = \frac{\delta_c}{\sigma_0} = \frac{1.68}{\sigma_0} \equiv \nu_c
$$

\n(ii)
$$
\eta_i(\mathbf{x}_p) = 0
$$

\n(iii)
$$
\lambda_1(\mathbf{x}_p) \ge \lambda_2(\mathbf{x}_p) \ge \lambda_3(\mathbf{x}_p) > 0
$$

 $\lambda_a(\mathbf{x}) =$ eigenvalues of $-\zeta_{ij}(\mathbf{x})$

"Localized" number density:

$$
n_{\rm pk}(\mathbf{x}) = \sum_{\mathbf{x}_p} \delta_D(\mathbf{x} - \mathbf{x}_p) = \frac{(6\pi)^{3/2}}{V_{\star}} \left| \det \zeta(\mathbf{x}) \right| \delta_D[\boldsymbol{\eta}(\mathbf{x})] \theta_H[\lambda_3(\mathbf{x})] \delta_D(\nu(\mathbf{x}) - \nu_c)
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Kac '43; Rice '51; Bardeen et al.'86

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So, one ends up computing sth like

$$
\int d^{10} \mathbf{y}_1 \dots \int d^{10} \mathbf{y}_N \, n_{\text{pk}}(\mathbf{y}_1) \times \dots \times n_{\text{pk}}(\mathbf{y}_N) P_N(\mathbf{y}_1, \dots, \mathbf{y}_N)
$$

$$
\mathbf{y}_\alpha = \left(\nu(\mathbf{x}_\alpha), \eta_i(\mathbf{x}_\alpha), \zeta_A(\mathbf{x}_\alpha) \right)
$$

Bardeen et al. '86; Regos & Szalay '95; VD '08; VD et al. '10

• Find all rotational invariants: $\nu(\mathbf{x})$

 $u(\mathbf{x}) \equiv -\mathrm{tr}\zeta(\mathbf{x})$

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\n
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\n
$$
\zeta^2(\mathbf{x}) \equiv \frac{3}{2} \text{tr}(\tilde{\zeta}^2)(\mathbf{x}), \qquad \tilde{\zeta}_{ij} = \zeta_{ij} + \frac{1}{3} u \,\delta_{ij}
$$

\n
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VD '13

• Write down the "effective" bias expansion:

$$
\delta_{\rm pk}(\mathbf{x}) = \sigma_0 b_{10} \nu(\mathbf{x}) + \sigma_2 b_{01} u(\mathbf{x}) + \frac{1}{2} \sigma_0^2 b_{20} \nu^2(\mathbf{x}) + \sigma_0 \sigma_2 b_{11} \nu(\mathbf{x}) u(\mathbf{x}) + \frac{1}{2} \sigma_2^2 b_{02} u^2(\mathbf{x}) + \sigma_1^2 \chi_{10} \eta^2(\mathbf{x}) + \sigma_2^2 \chi_{01} \zeta^2(\mathbf{x}) + \dots
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$$

• Compute the bias factors: $\sigma^i_0\sigma^j_2b_{ij} =$ $\frac{1}{\bar{n}_{\rm pk}}$ \int $d^{10}\mathbf{y}$ $n_\mathrm{pk}(\mathbf{y}) H_{ij}(\nu,u) P_1(\mathbf{y})$ $\sigma_1^{2k}\chi_{k0} =$ $\frac{1}{\bar{n}_{\rm pk}}$ \int d^{10} y $n_{\rm pk}$ (y) $L_k^{(1/2)}$ $\sqrt{3\eta^2}$ 2 ◆ $P_1(\mathbf{y})$ $\sigma_2^{2k}\chi_{0k} =$ $\frac{1}{\bar{n}_{\rm pk}}$ \int d^{10} y $n_{\rm pk}$ (y) $L_k^{(3/2)}$ $\sqrt{5\zeta^2}$ 2 ◆ $P_1(\mathbf{y})$

⁰¹ : *bias induced by peak profile asphericity*

FIG. 7—The 95%, 90%, and 50% contours of the conditional probability for ellipticity $e = y/x$ and prolateness $p = z/x$ subject to the constraint of given x for peaks (eq. [7.6]). (The x and x_{*} used here are 1.58 $\approx y^{-1}$ ti

Fig. from Bardeen et al '86

⁰¹ : *bias induced by peak profile asphericity*

• N-point connected correlations can be perturbatively computed from (2×7)

$$
\xi_{\rm pk}^{(N)}(\mathbf{x}_1,\ldots,\mathbf{x}_N)\equiv\langle\delta_{\rm pk}(\mathbf{x}_1)\times\cdots\times\delta_{\rm pk}(\mathbf{x}_N)\rangle
$$

"Excursion set peaks": combine peak constraint with first-crossing condition

$$
\mu(\mathbf{x}) \equiv -\frac{d\delta_s}{dR_s}(\mathbf{x})
$$

The "localized" number density of excursion set peaks becomes

$$
n_{\text{ESP}}(\mu, \mathbf{y}) = \left(\frac{\mu}{\gamma_{\nu\mu}\nu_c}\right) \theta_H(\mu) n_{\text{pk}}(\mathbf{y})
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Appel & Jones '91; Paranjape & Sheth '12; Paranjape, Sheth & VD '13

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and the "effective" bias expansion generalizes to

 $\delta_{\text{ESP}}(\mathbf{x}) = \sigma_0 b_{100} \nu(\mathbf{x}) + \sigma_2 b_{010} u(\mathbf{x}) + b_{001} \mu(\mathbf{x}) + \dots$

VD, Gong & Riotto '13

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f_{\rm NL}\phi^2: \quad \Delta b_1^{\rm NG}(k) = \frac{2f_{\rm NL}b_{\rm NG}}{\mathcal{M}(k)}
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b_{\rm NG} = \sigma_0^2 b_2
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 \nTaruya et al '08

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Excursion set peaks:

VD, Gong & Riotto '13

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b_{\rm NG}=\sigma_0^2b_{200}+2\sigma_1^2b_{110}+\sigma_2^2b_{020}+2\sigma_1^2\chi_{10}+2\sigma_2^2\chi_{01}+\Delta_0^2b_{002}+2\gamma_{\nu\mu}\sigma_0b_{101}+2\gamma_{u\mu}\sigma_2b_{011}
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The following three statements are certainly true:

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• Local bias expansions cannot correctly predict the amplitude of non-Gaussian bias

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All the results discussed so far have a common denominator: the collapse threshold is assumed to be constant and deterministic while it is in fact moving (e.g. halo mass-dependent) and stochastic

$$
B(\sigma)=\delta_c
$$

Paranjape, Sheth & VD '13; Biagetti et al '14

Preliminary ! work in progress with Matteo Biagetti

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non-Gaussian bias

Hamaus, Seljak & VD '11

Beyond galaxy power spectrum: galaxy bispectrum and matter statistics

Sefusatti, Crocce & VD '12

Conclusion

- *• Non-Gaussian bias is subtle: peak-background split may not work*
- *• If confirmed, many of the current constraints and forecasts are in need of revision*
- Peak theory is a powerful approach to understand the nitty-gritty *details of galaxy bias*