» Absolute Polarimetry of Proton Beams at RHIC«

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for the RHIC Polarimetry Group

International Workshop on Polarized Sources, Targets and Polarimetry

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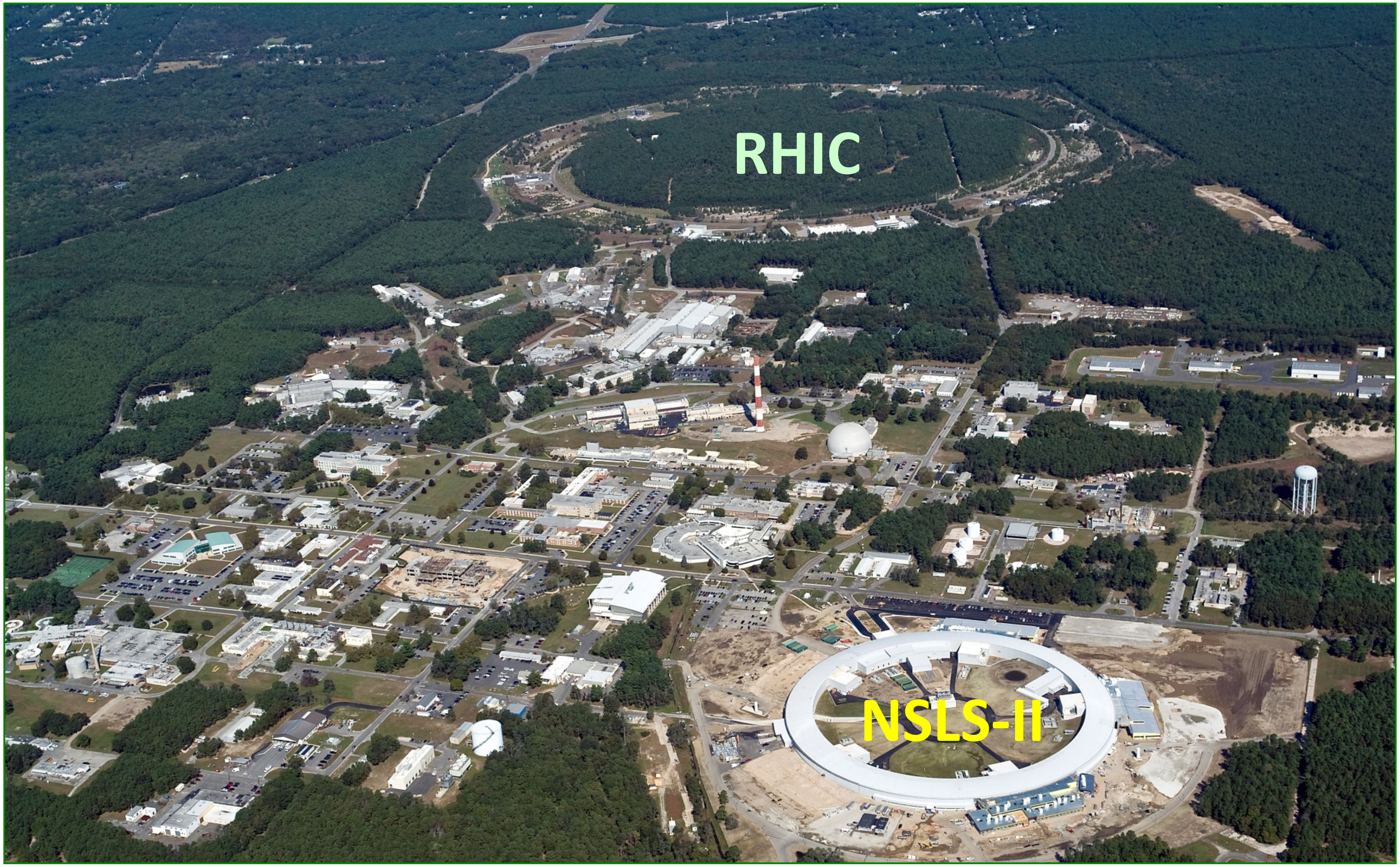
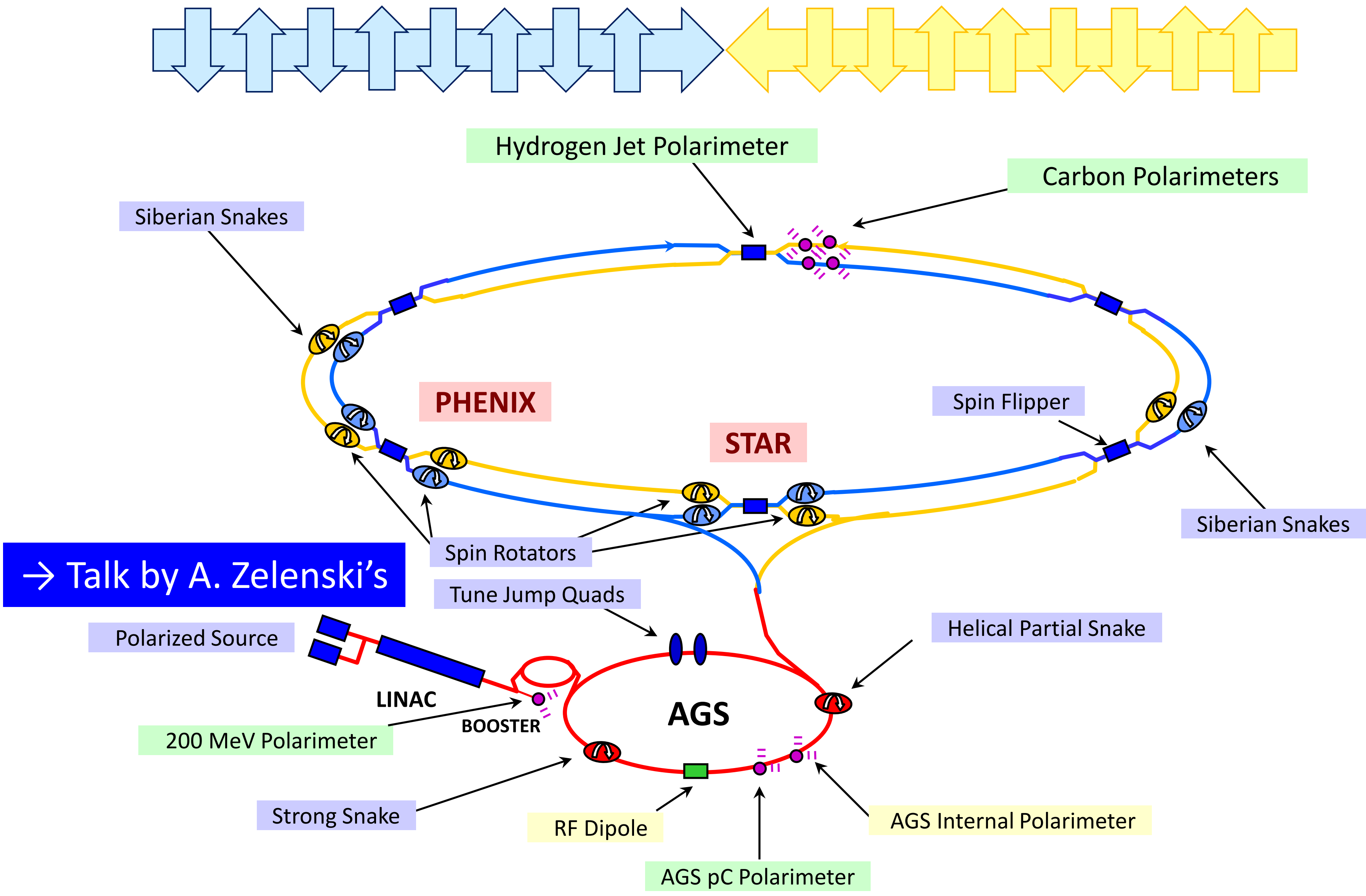
BROOKHAVEN NATIONAL LABORATORY

The Relativistic Heavy Ion Collider

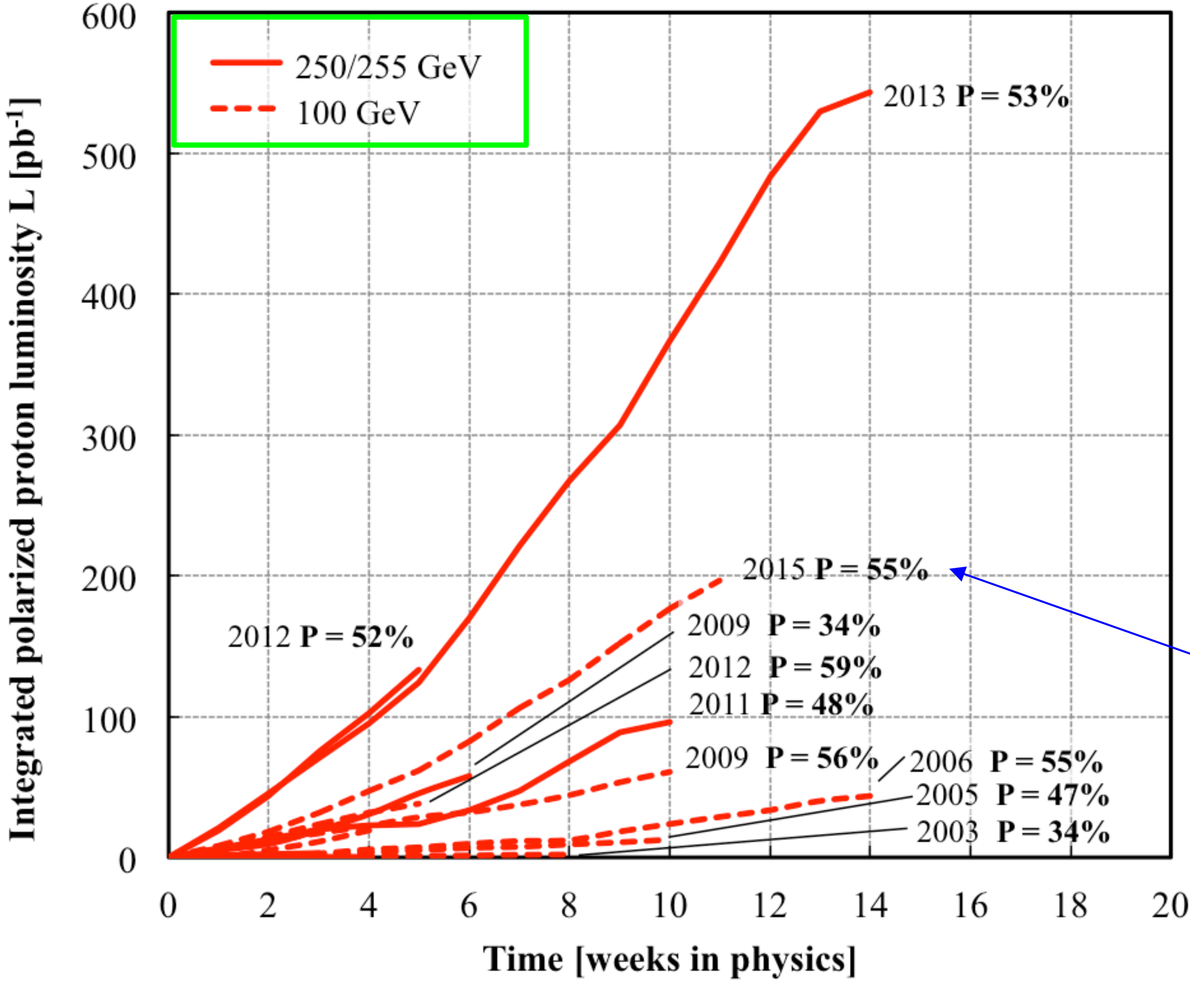
at Brookhaven National Laboratory



Polarized Protons in RHIC



Improvement in Beam Polarization



recent run 2015

Beam energies:

up to 255 GeV

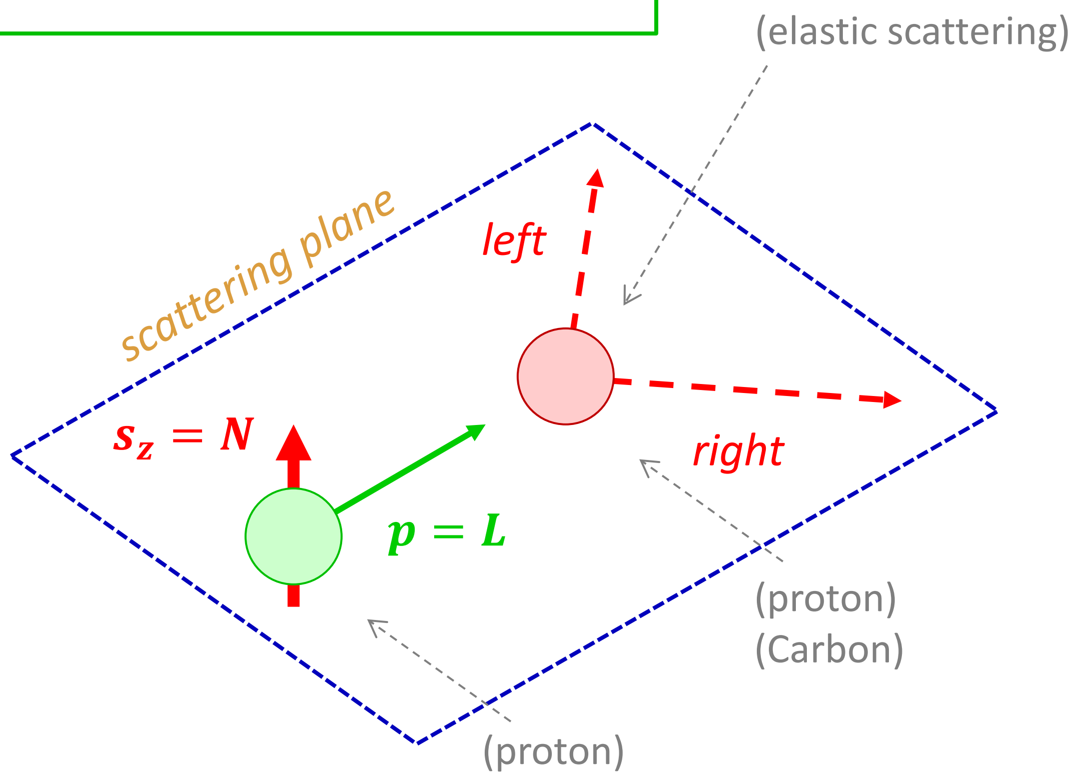
Figure of merit for double helicity measurements:

$$\sim\mathcal{L} \cdot P^4$$

Consistent improvement in delivered luminosity and beam polarization.

Polarization & Asymmetries

$$s_z = \pm\frac{1}{2}\hbar \Rightarrow P = \frac{n^{\uparrow} - n^{\downarrow}}{n^{\uparrow} + n^{\downarrow}}$$



$$A_N = \frac{d\sigma_{left} - d\sigma_{right}}{d\sigma_{left} + d\sigma_{right}}$$

$$\varepsilon = A_N \cdot P = \frac{N_L - N_R}{N_L + N_R}$$

(*) perpendicular to polarization vector

Elastic Proton-Proton Scattering

$$\varphi(s,t) = \langle \lambda_C \lambda_D | \varphi | \lambda_A \lambda_B \rangle$$

$$\varphi_1(s,t) = \left\langle +\frac{1}{2} + \frac{1}{2} \middle| \varphi \middle| +\frac{1}{2} + \frac{1}{2} \right\rangle$$

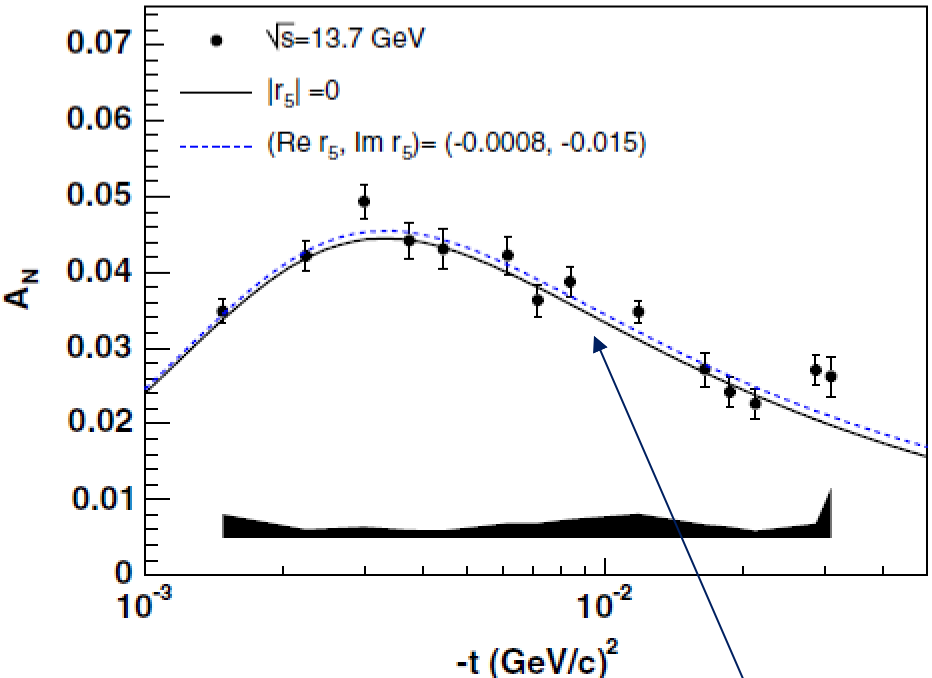
$$\varphi_2(s,t) = \left\langle +\frac{1}{2} + \frac{1}{2} \middle| \varphi \middle| -\frac{1}{2} - \frac{1}{2} \right\rangle$$

$$\varphi_3(s,t) = \left\langle +\frac{1}{2} - \frac{1}{2} \middle| \varphi \middle| +\frac{1}{2} - \frac{1}{2} \right\rangle$$

$$\varphi_4(s,t) = \left\langle +\frac{1}{2} - \frac{1}{2} \middle| \varphi \middle| -\frac{1}{2} + \frac{1}{2} \right\rangle$$

$$\varphi_5(s,t) = \left\langle +\frac{1}{2} + \frac{1}{2} \middle| \varphi \middle| +\frac{1}{2} - \frac{1}{2} \right\rangle$$

Phys. Rev. D 79, 094014 (2009)



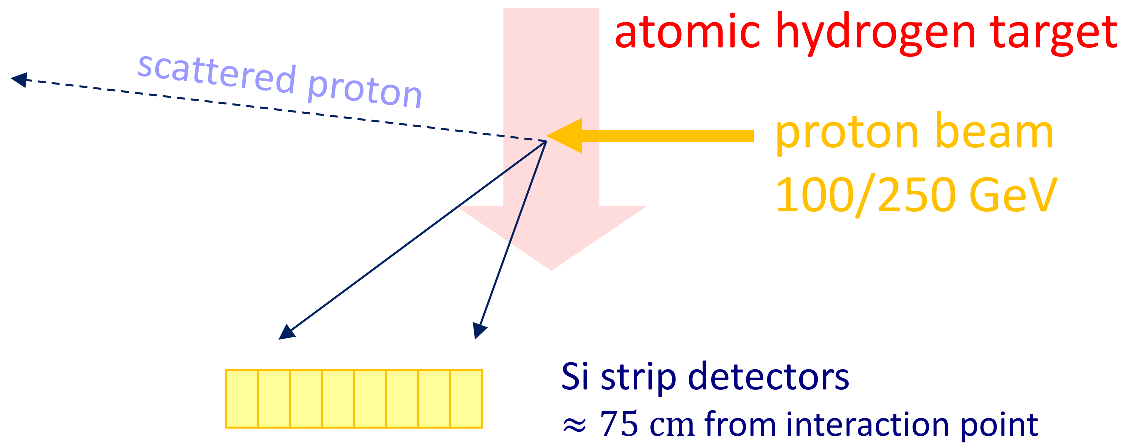
First data from 2004 (100 GeV beam)

$$A_N \frac{ds}{dt} = -\frac{4\pi}{s^2} \mathrm{Im}\left[\varphi_5^{em*}(s,t)\varphi_+^{had}(s,t) + \varphi_5^{had*}(s,t)\varphi_+^{em}(s,t)\right]$$

no-flip amplitude: $\varphi_+(s,t) = \frac{1}{2}[\varphi_1(s,t)+\varphi_3(s,t)]$

Transverse single-spin asymmetries are driven by an interference of amplitudes and can be compared to Regge theory.

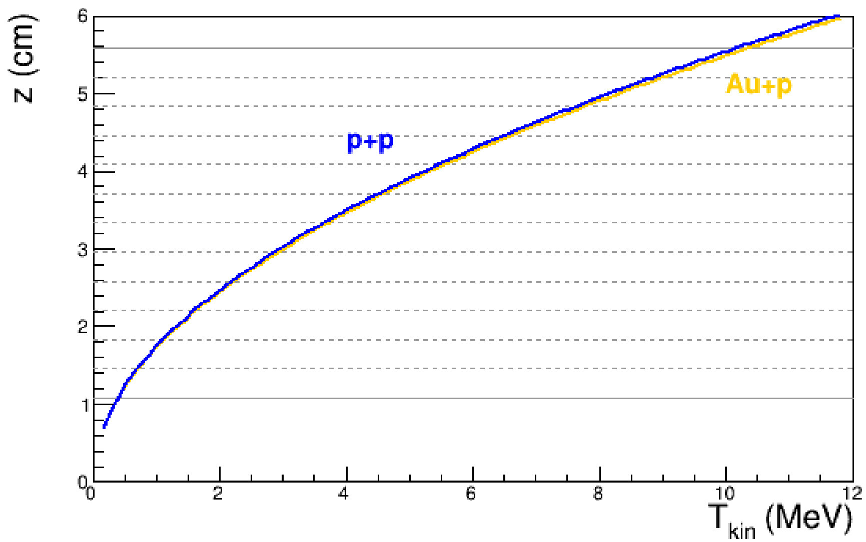
Elastic Recoil Protons

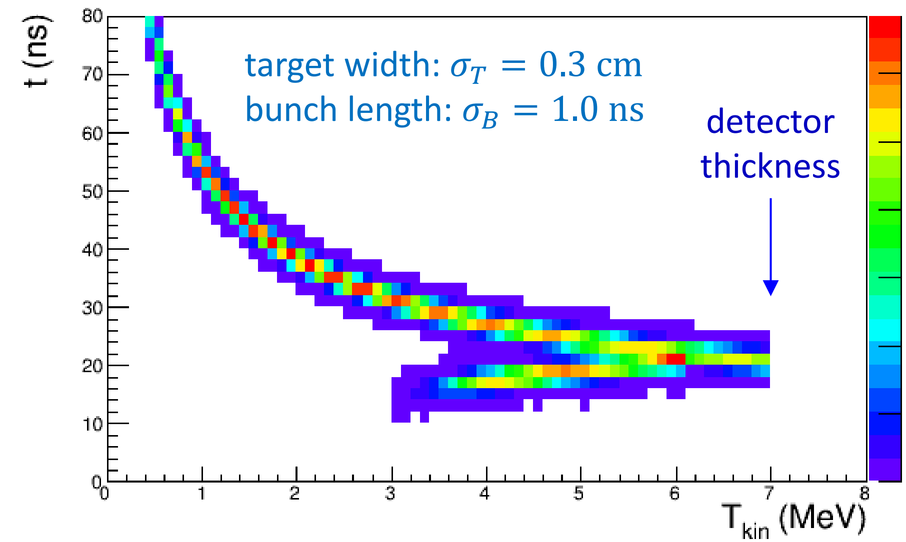


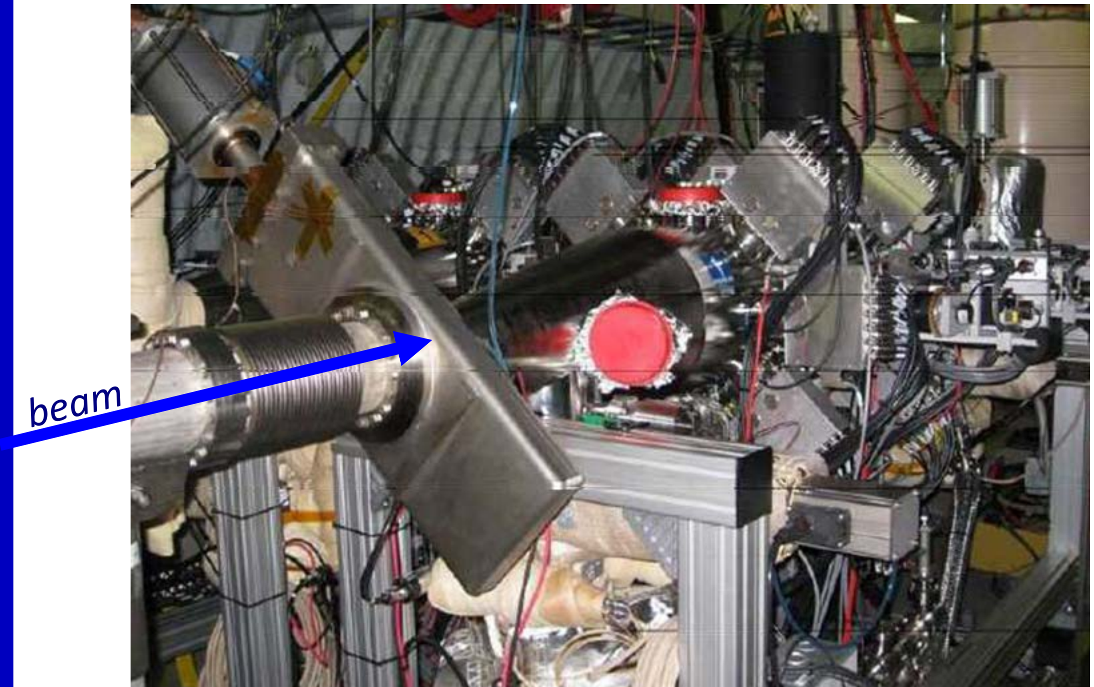
Recoil proton from elastic scattering

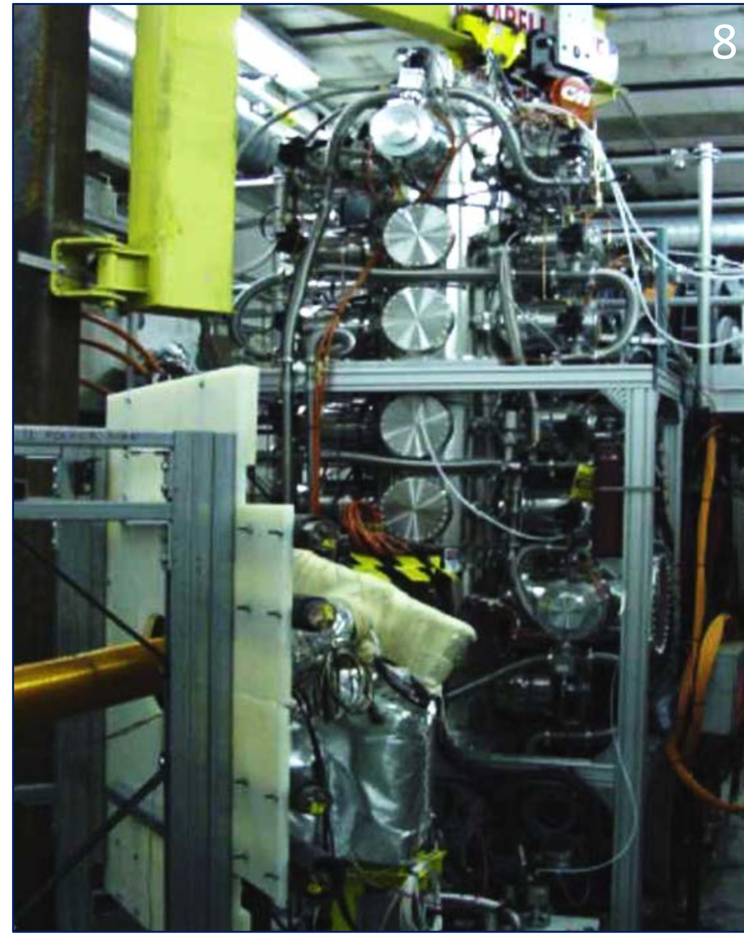
Independent of beam energy, species

Non-relativistic: $T_{kin} = \frac{1}{2}mv^2$









→ Talk by G. Webb

Carbon polarimeters

Two per ring

Fast measurement

$\delta`P/P \approx 4\%$

Beam polarization profile

Polarization decay (time dependence)

normalization

Hydrogen jet polarimeter

Polarized target

Continuous operation

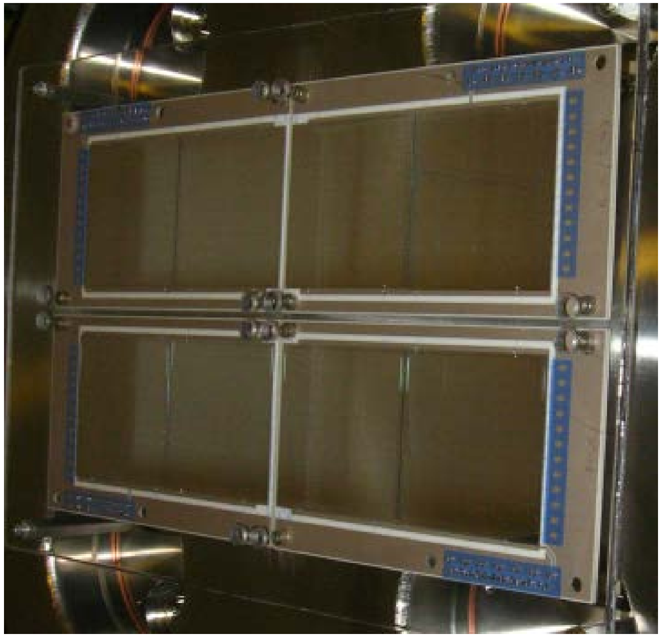
$\delta P/P \approx 5 - 8\%$ per fill

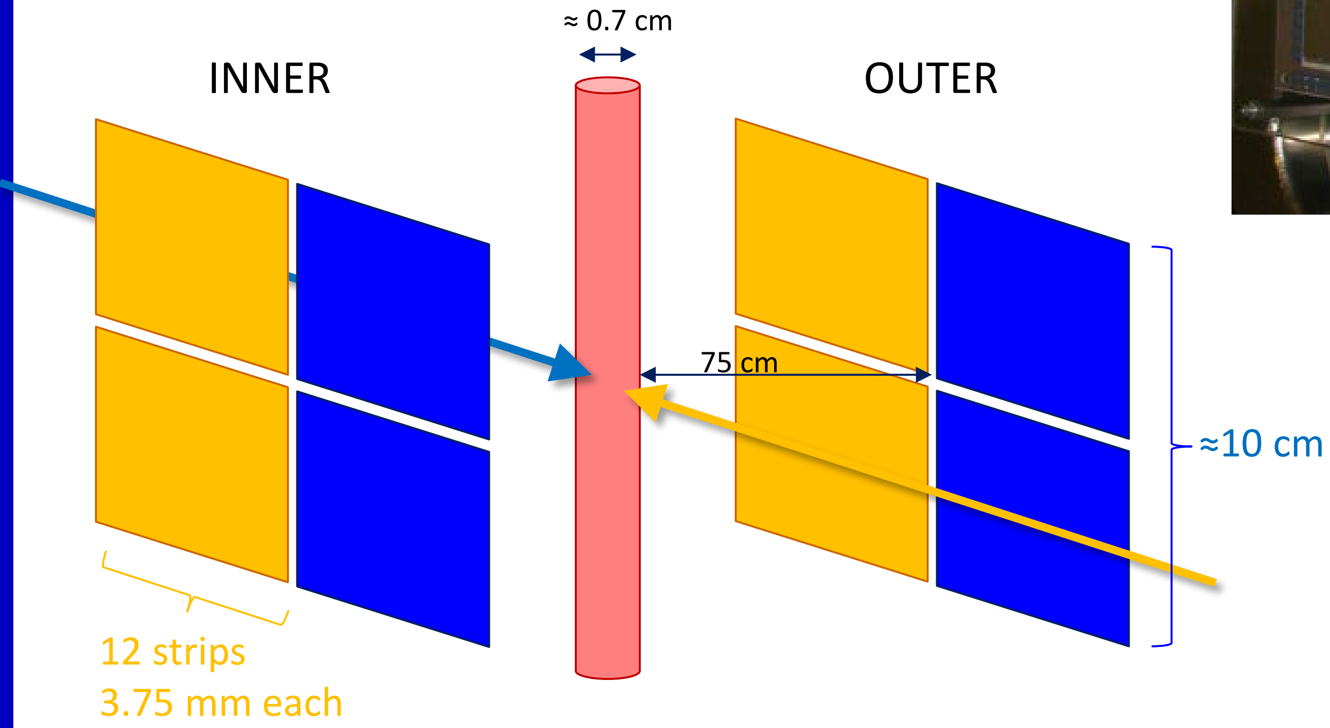
Detector Setup

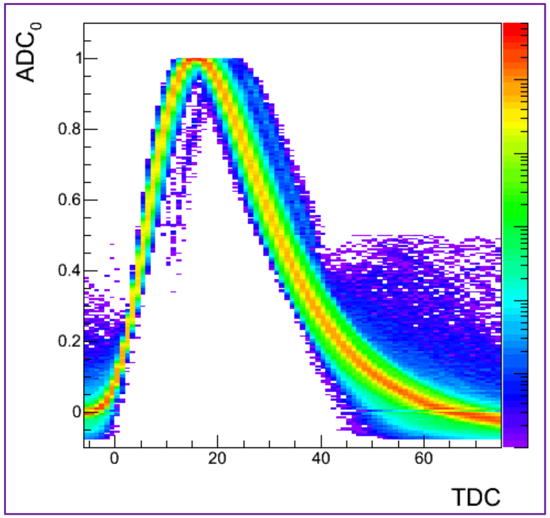
Set of eight Hamamatsu Si strip detectors

12 strips, each 3.75 mm wide, 500 μm thick

Uniform dead layer ≈ 1.5 μm







Energy Calibration

Calibrations are done every few days:

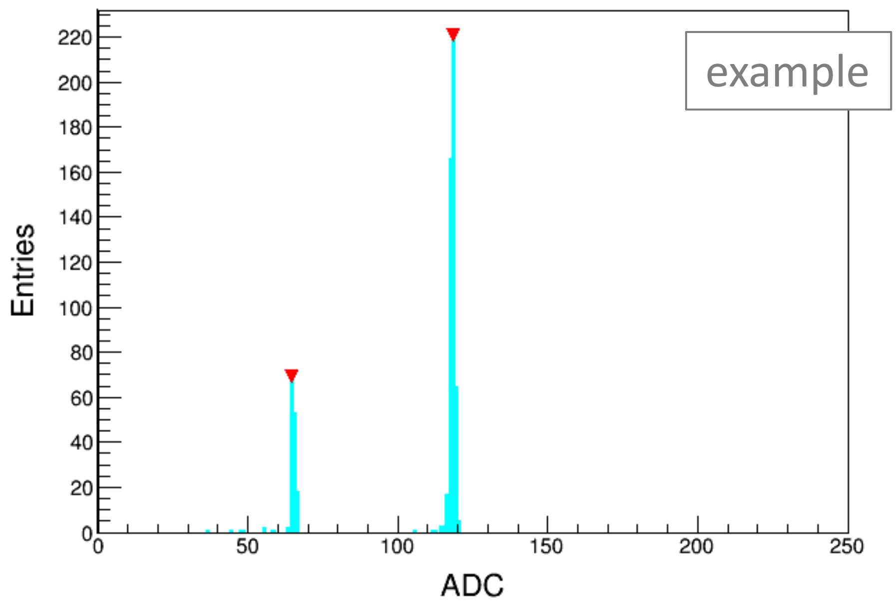
- Gain
- Entrance window (dead layer)

Two different α-sources

$$E_\alpha(Gd) = 3.183 \text{ MeV}$$

$$E_\alpha(Am) = 5.486 \text{ MeV}$$

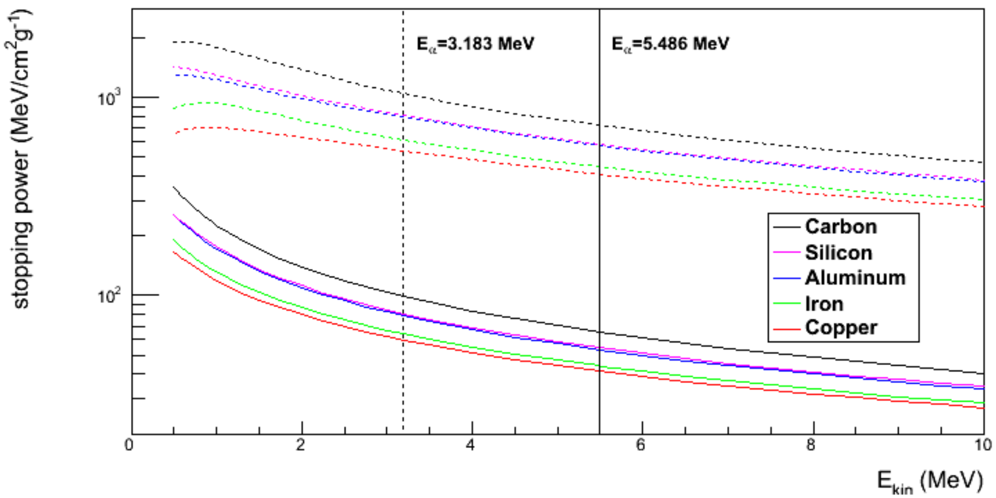
Resolution of peak finding is within 1 ADC count



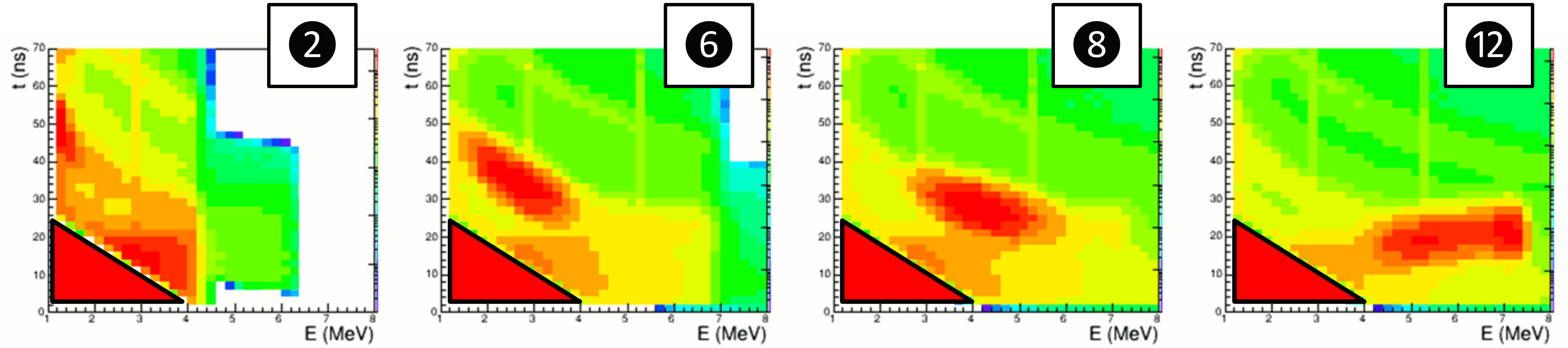
Stopping power for protons and α-particles from NIST database:

$$\Delta E_{\alpha(Am)} = 0.72 \cdot \Delta E_{\alpha(Gd)}$$

$$\Delta E_P = 0.44 \cdot \Delta E_{\alpha(Gd)} \cdot E[MeV]^{-0.64}$$



Kinematics

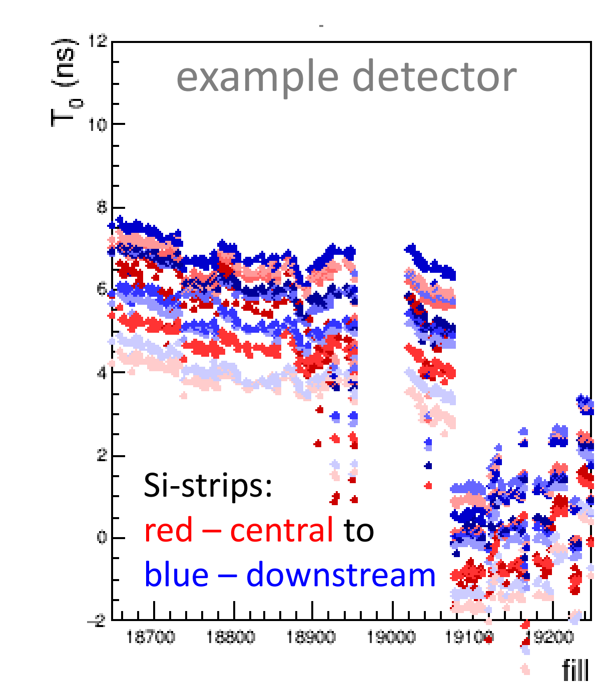


12 strips per detector

Removed peak in prompt hits at low ADC/TDC region

Using elastic p-recoil signature for time-of-flight offset determination

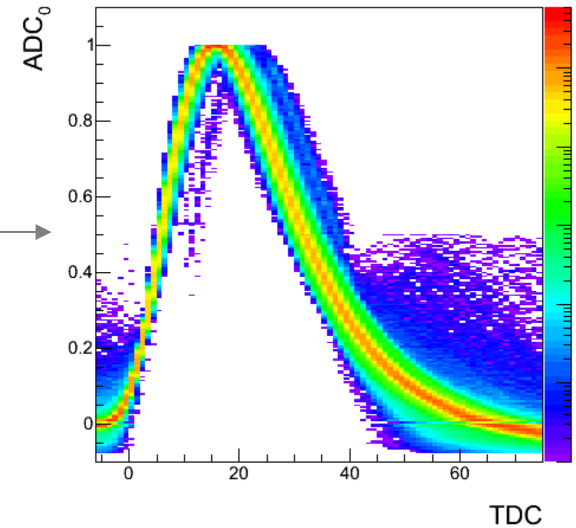
- Slow drift with time (detector/read-out)
- Big jumps when changing the DAQ system



Stopped Recoil Protons

Normalized to $ADC_{\max}$

Slope δ_{ADC} calculated in six TDC bins around ½ $ADC_{\max}$



→Talk by A. Poblaguev

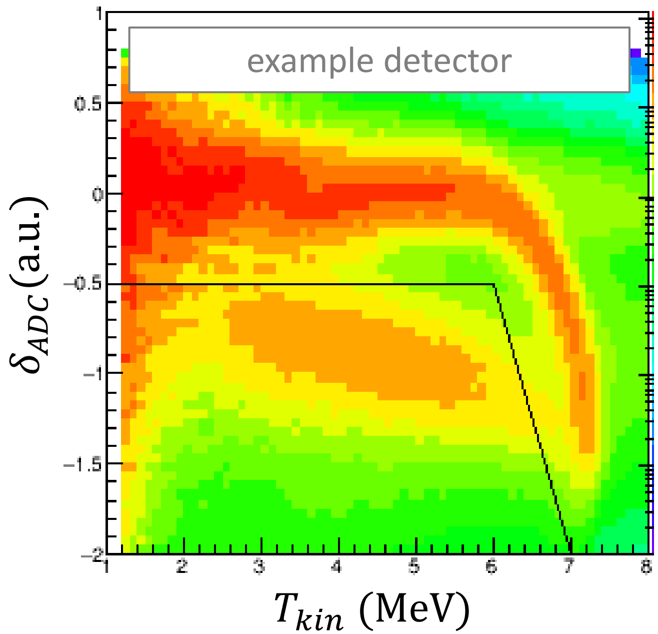
Slope of rise in waveform can be used to identify punch-through particles

Normalized waveform rise ($4.5 < E < 5.5$ MeV) in each detector

Independent of DAQ system (CAMAC/VME)

Remove punch-through particles:

$$(\delta_{\mathrm{ADC}} < -0.5) \wedge (\delta_{ADC} < 8.5 - 1.5 * T_{kin})$$



Detector Alignment

Magnetic holding field for target polarization changes acceptance of detectors on left and right sides

Outer correction field for compensation

For missing proton mass:

$$\sin\theta = \frac{p'}{2 \cdot m_p \cdot p_B}(2 \cdot E + 2 \cdot m_p - T_R)$$

Missing mass:

$$M_{miss}^2 = \begin{pmatrix} E + m_p - E' \\ p_B - p' \end{pmatrix}^2$$

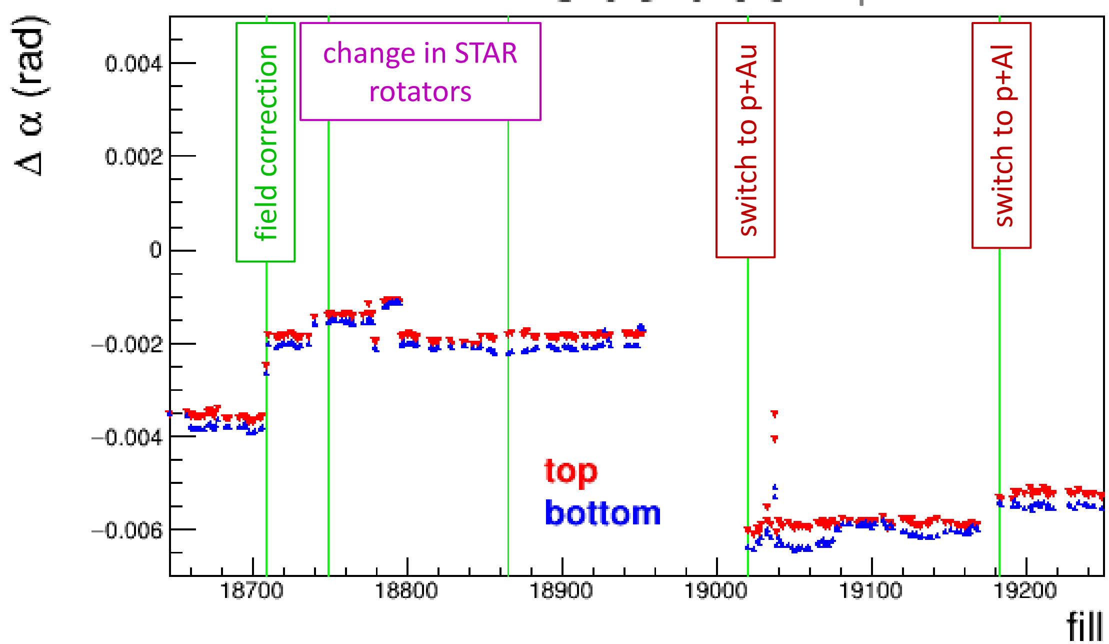
Non-relativistic recoil:

$$p' = \sqrt{2 m_p T_R}$$

Compare with geometry of detector averaged 12 strips

p+Au and p+Al operation had a significant beam angle on the jet target

example detector

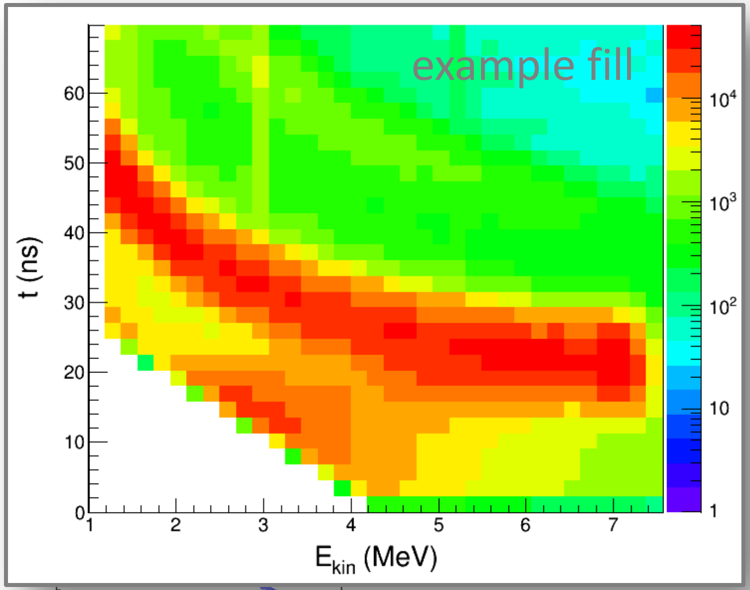


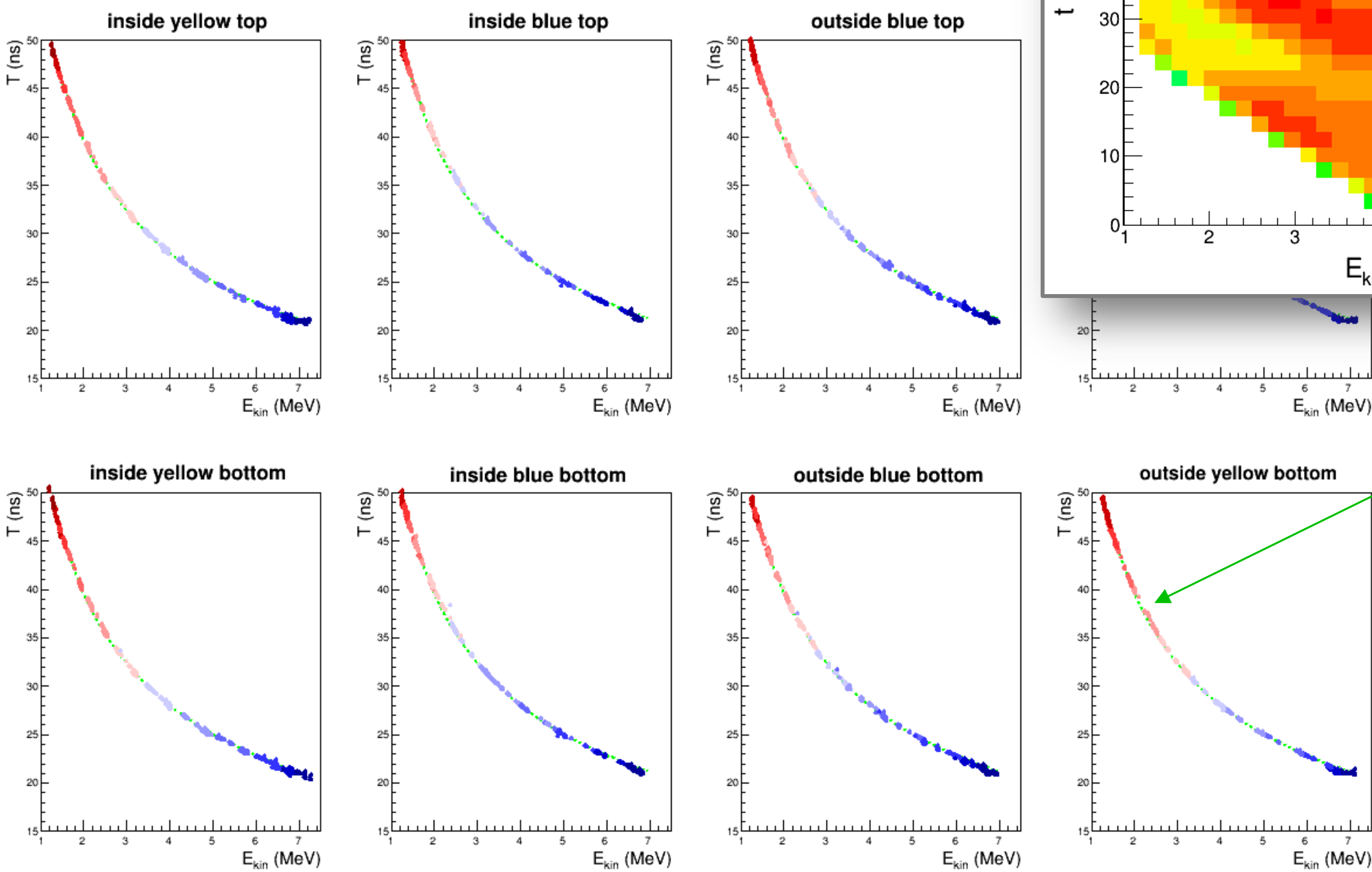
QA: Kinematics

Elastic proton recoil selection:

$$|M_{miss} - m_p| < 100 \text{ MeV}/c^2$$

$$|\Delta t| < 5 \text{ ns}$$





Fit to ALL data, plotted under the distributions in each detector

Si-strips:
red – central to
blue – downstream

Asymmetries & Polarization

$$\varepsilon = A_N \cdot P$$

$$P_{Beam} = -\frac{\varepsilon_{Beam}}{\varepsilon_{Target}} P_{Target}$$

ε_{Target}: measure

P_{Target}: from Breit-Rabi measurement

1

Polarization independent background

$$\varepsilon = \frac{N^{\uparrow} - N^{\downarrow}}{N^{\uparrow} + N^{\downarrow} + 2 \cdot N_{bg}} \Rightarrow \frac{\varepsilon_B}{\varepsilon_T} = \frac{N_B^{\uparrow} - N_B^{\downarrow}}{N_T^{\uparrow} - N_T^{\downarrow}}$$

2

Polarization dependent background

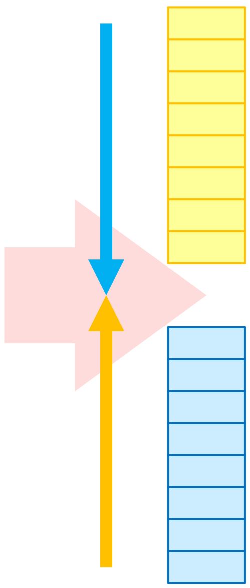
$$\varepsilon = \frac{\varepsilon_{inc} - r \cdot \varepsilon_{bg}}{1 - r}$$

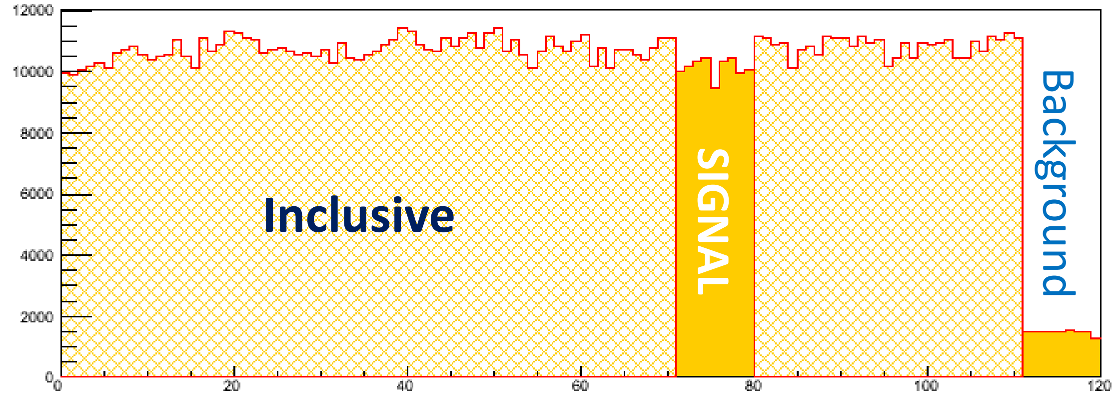
background fraction $r = N_{bg}/N$

Signal & Background I

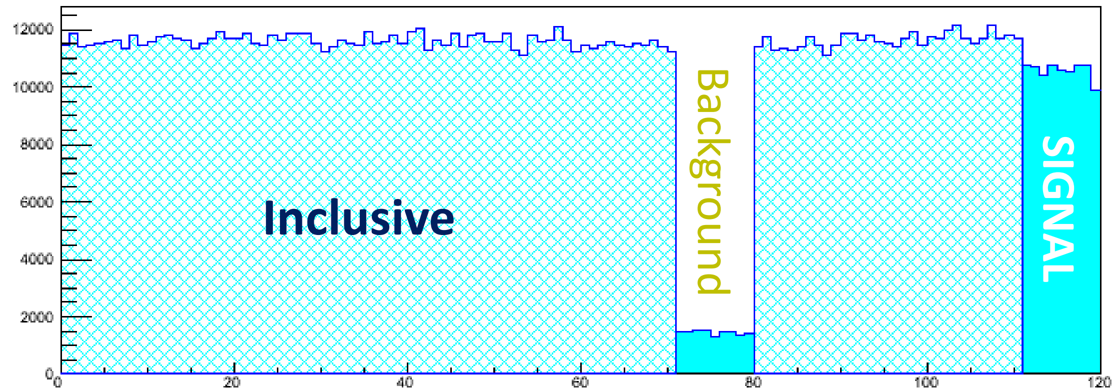
Abort gaps are not aligned at 12 o'clock

Use abort gaps for background and clean signal identification





RHIC bunch



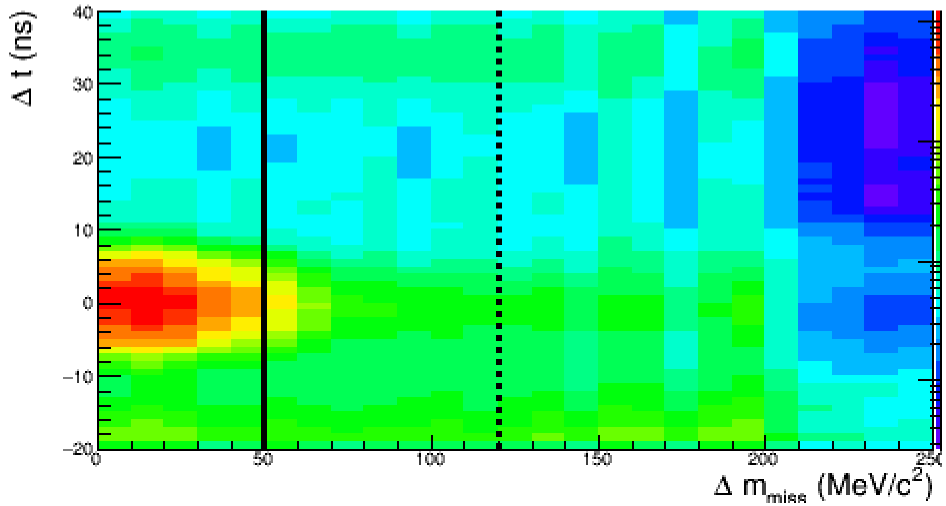
RHIC bunch

Signal & Background II

Δt: difference of time-of-flight to elastic signal (in geometry)

Δm_{miss}: difference of missing mass to scattered proton (in geometry after alignment correction)

Example (logarithmic z-scale)



Position of elastic proton signal is independent of energy and detector

Vertical stripes are a remnant of the spatial detector resolution

Punch through cuts are already applied

Define signal and background regions by missing mass

$$|M_{miss} - M_p| < 50\ \mathrm{MeV}/c^2$$

$$|M_{miss} - M_p| > 120\ \mathrm{MeV}/c^2$$

Signal & Background III

- inclusive (normalized to peak)

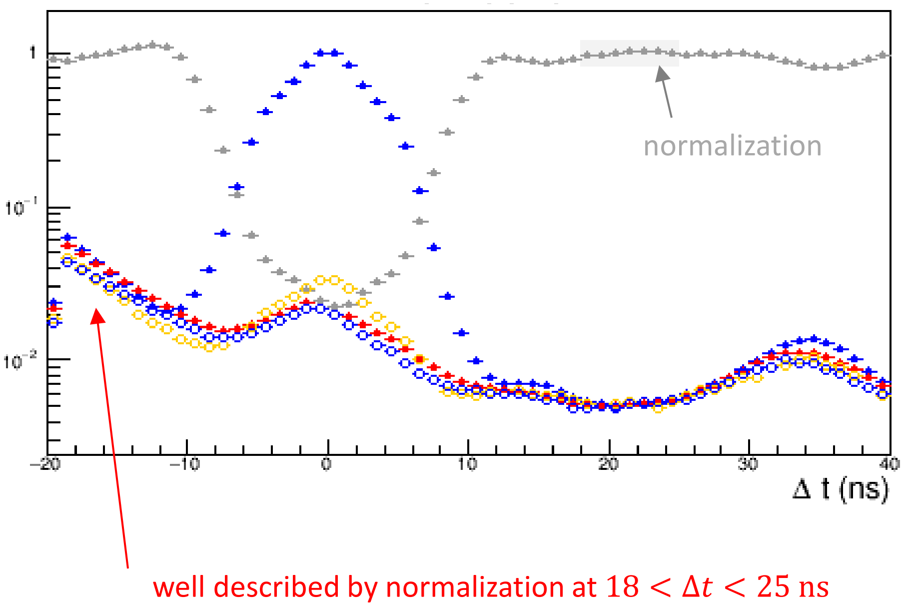
$$|M_{miss} - m_p| < 50 \text{ MeV}/c^2$$

- background (normalized to signal at $18 < \Delta t < 25$ ns)

$$|M_{miss} - m_p| > 120 \text{ MeV}/c^2$$

- background fraction

Example (blue beam, $2 < E_{kin} < 3$ MeV)

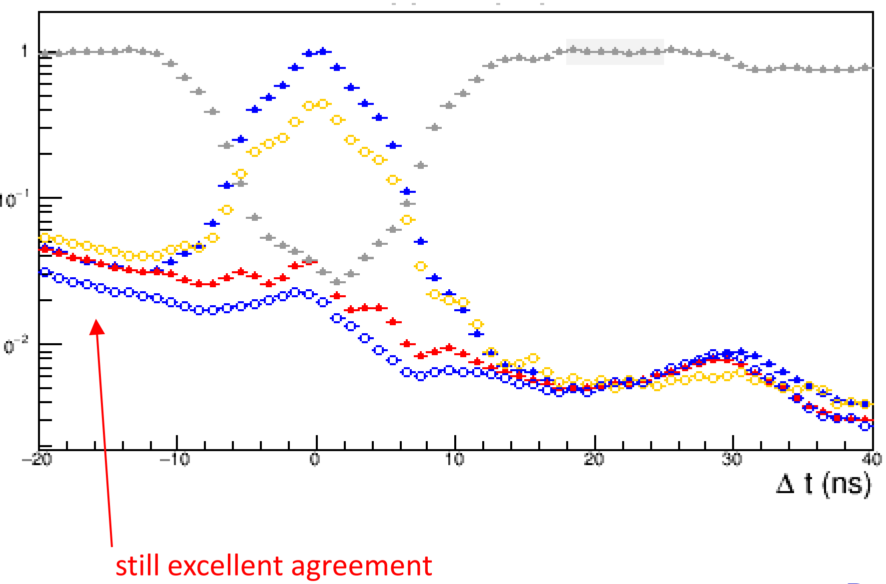


- o Background in yellow abort gap (should be clean blue signal)
- o Signal in blue abort gap (should be only background from yellow beam)

The normalization is same as above → only for comparison of shape and source of background

Background Sources

Example (blue beam, $3 < E_{kin} < 4$ MeV)



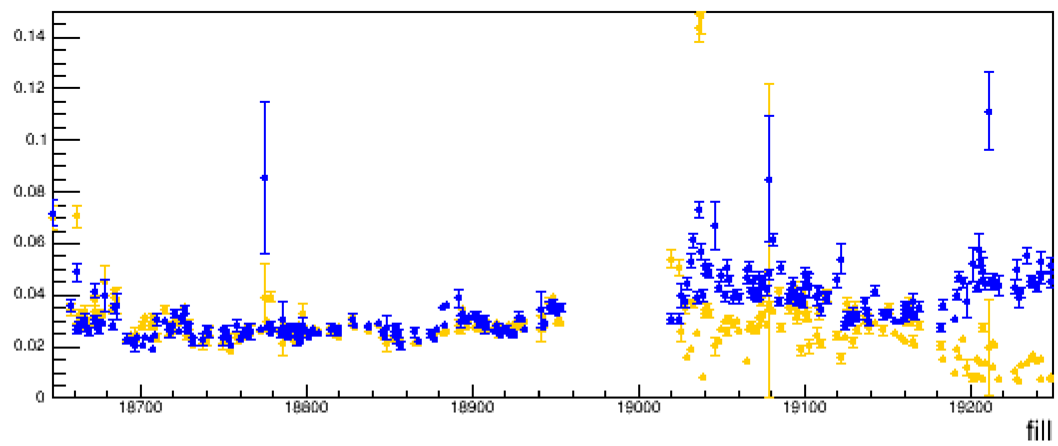
From $p + Au$ operation

Typical bunch shape of Au-beam seen in full background, dominates *early* background

Late background mainly from signal beam

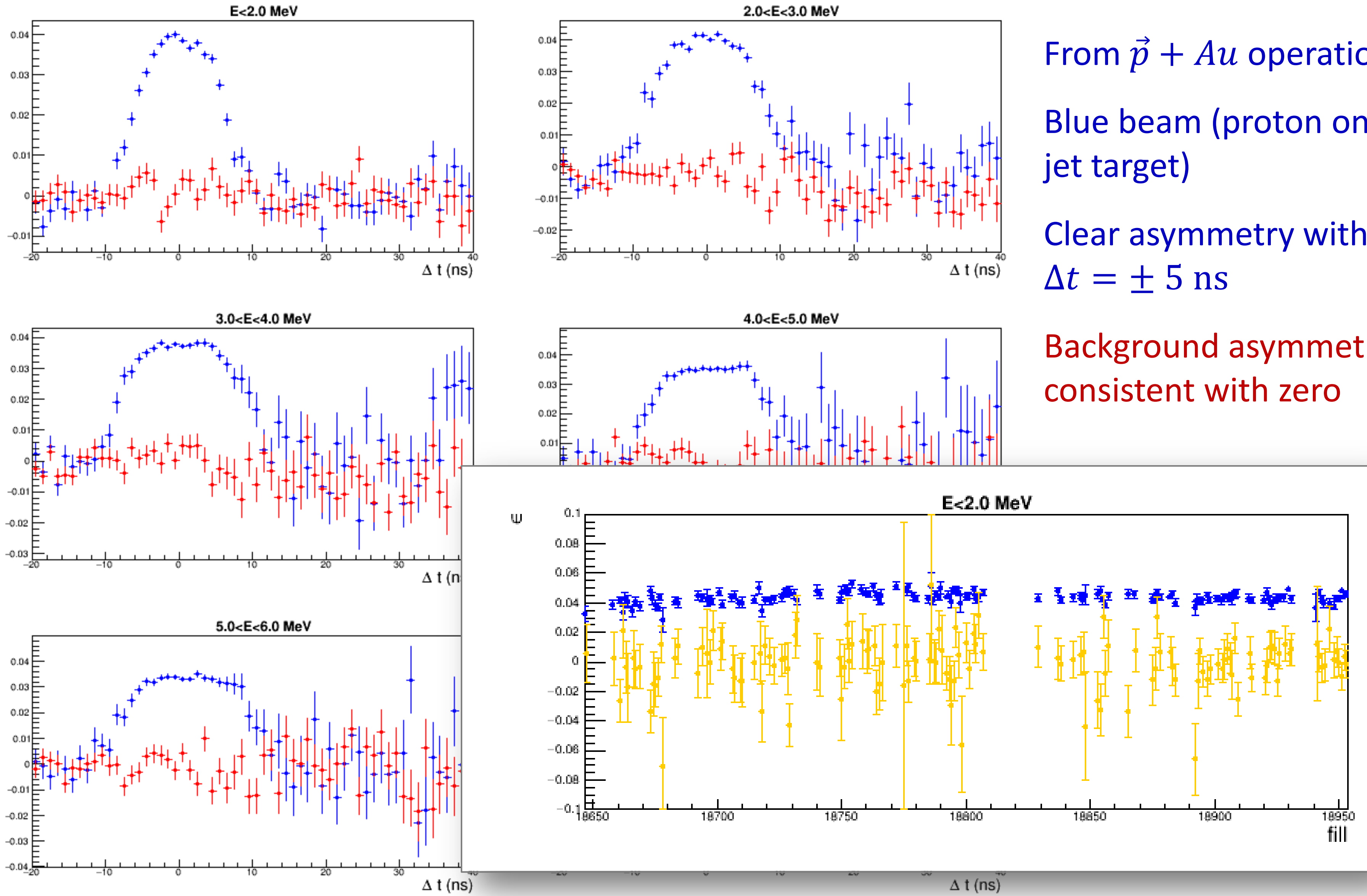
Using signal cuts in blue abort gap: $|M_{miss} - m_p| < 50 \text{ MeV}/c^2$

Background fraction $r = N_{bg}/N$



Fill-by-fill background fraction depends on conditions of both beams → important for beam polarization measurement

Asymmetry Examples



From $\vec{p} + Au$ operation

Blue beam (proton on jet target)

Clear asymmetry within $\Delta t = \pm 5$ ns

Background asymmetry consistent with zero

Correlated Background

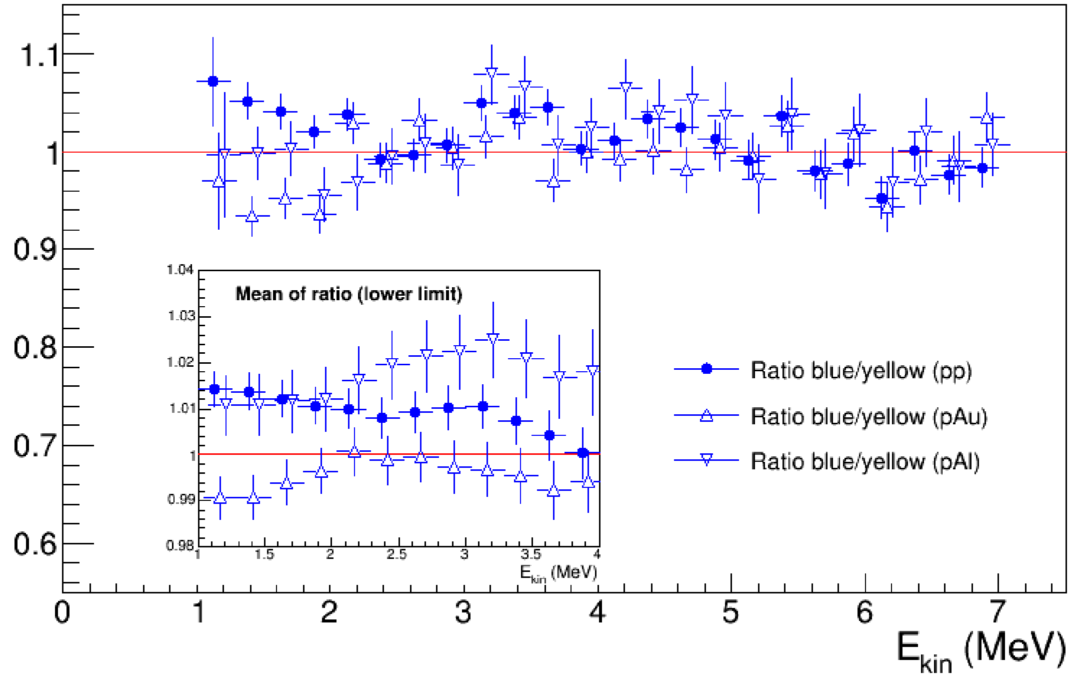
<u>Inelastic background:</u> expected to be polarization independent (or small compared to signal)

<u>Elastic background:</u> opposite beam can affect low energy signal

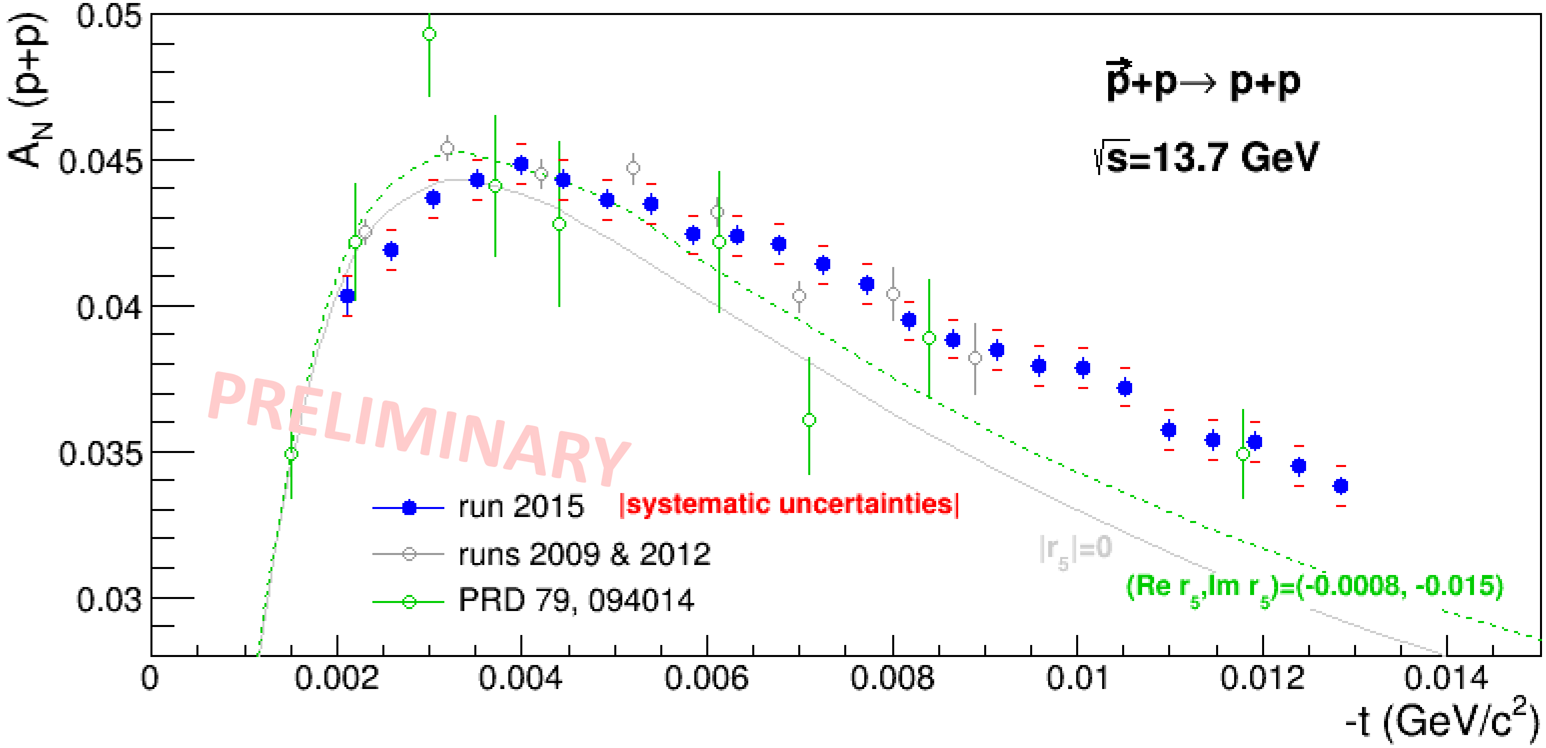
- jet target size ≈ 0.7 cm
- would have opposite sign for target polarization, no effect on beam polarization
- consistent with zero

Increased asymmetries have been observed at low energies, attributed to parts of the blue-side detectors (masked out)

Possibly due to correlated background that *suppresses* the signal in one polarization state



Analyzing Power: $A_N(\vec{p}+p)$



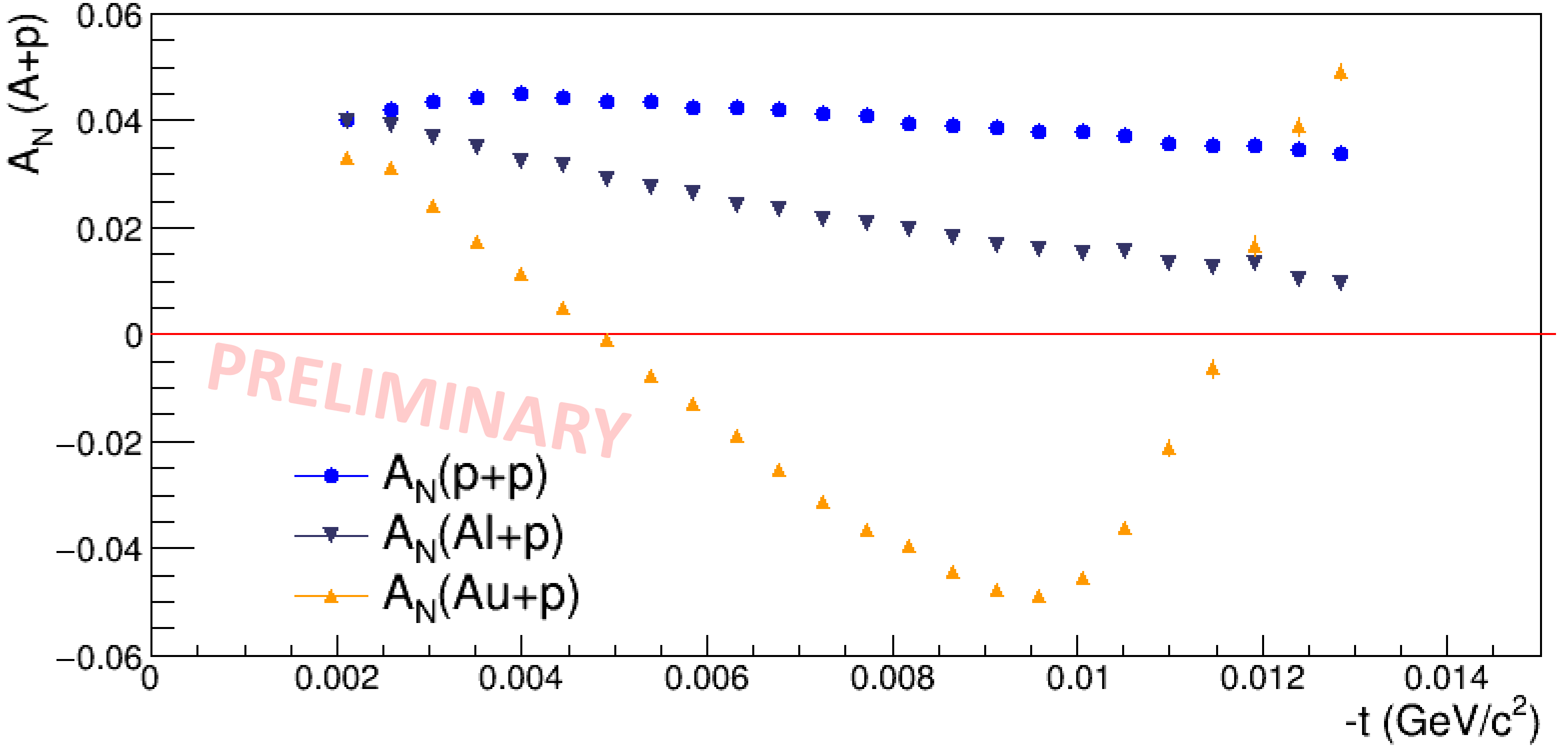
Atomic hydrogen polarization $P = 96\%$

Molecular component $R_{H_2} = 3\%$ (by mass)

Global uncertainty from target polarization not included

$-t$-range can be extended with punch-through protons

Analyzing Power: $A_N(\vec{p} + A)$



Atomic hydrogen polarization $P = 96\%$

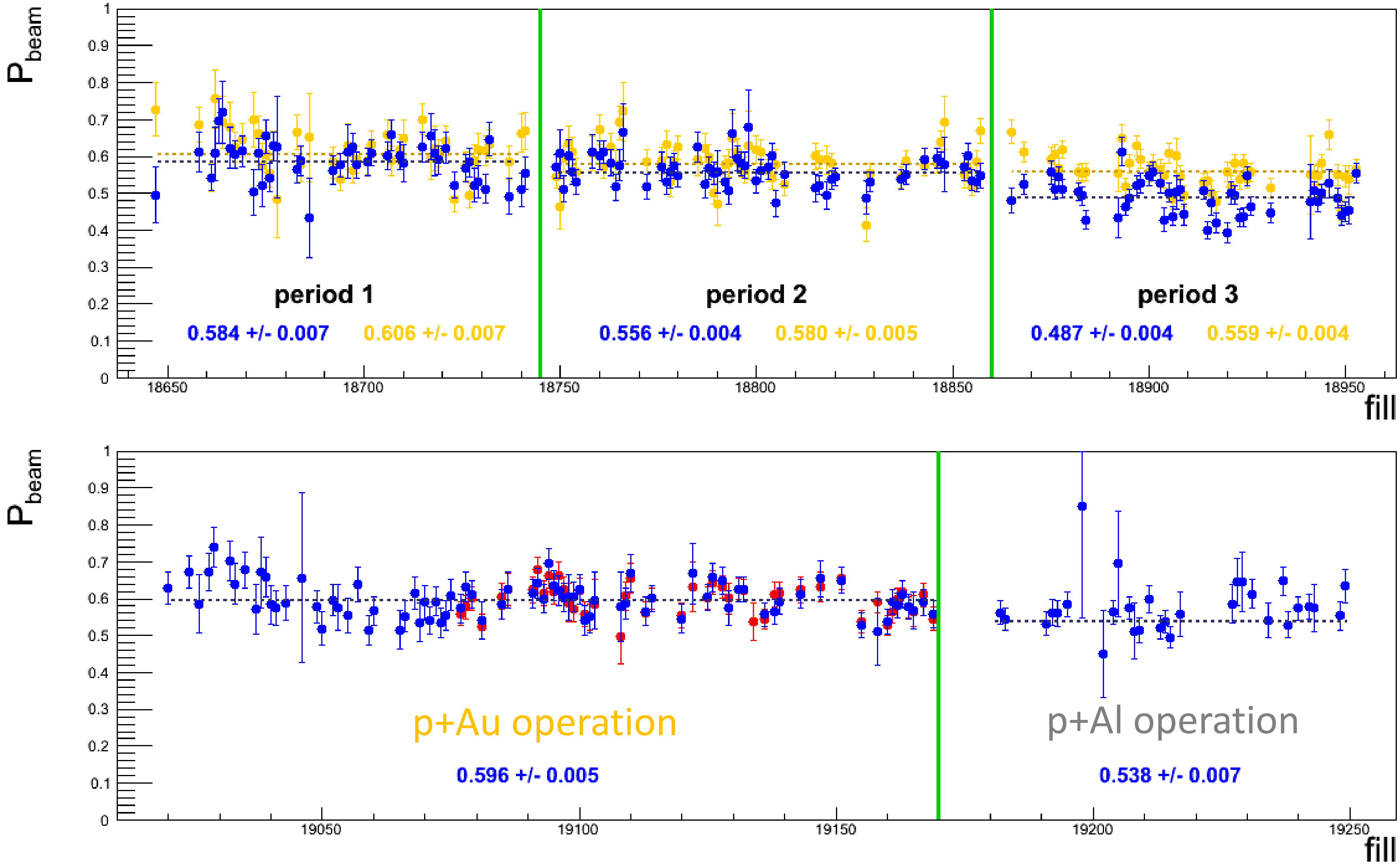
Molecular component $R_{H_2} = 3\%$ (by mass)

Global uncertainty from target polarization not included

$-t$-range can be extended with punch-through protons

Beam Polarizations

Online results from 2015, no background correction included



Summary

- Polarimetry at RHIC
  - Essential input for experiments
  - Fast feedback during collider operation

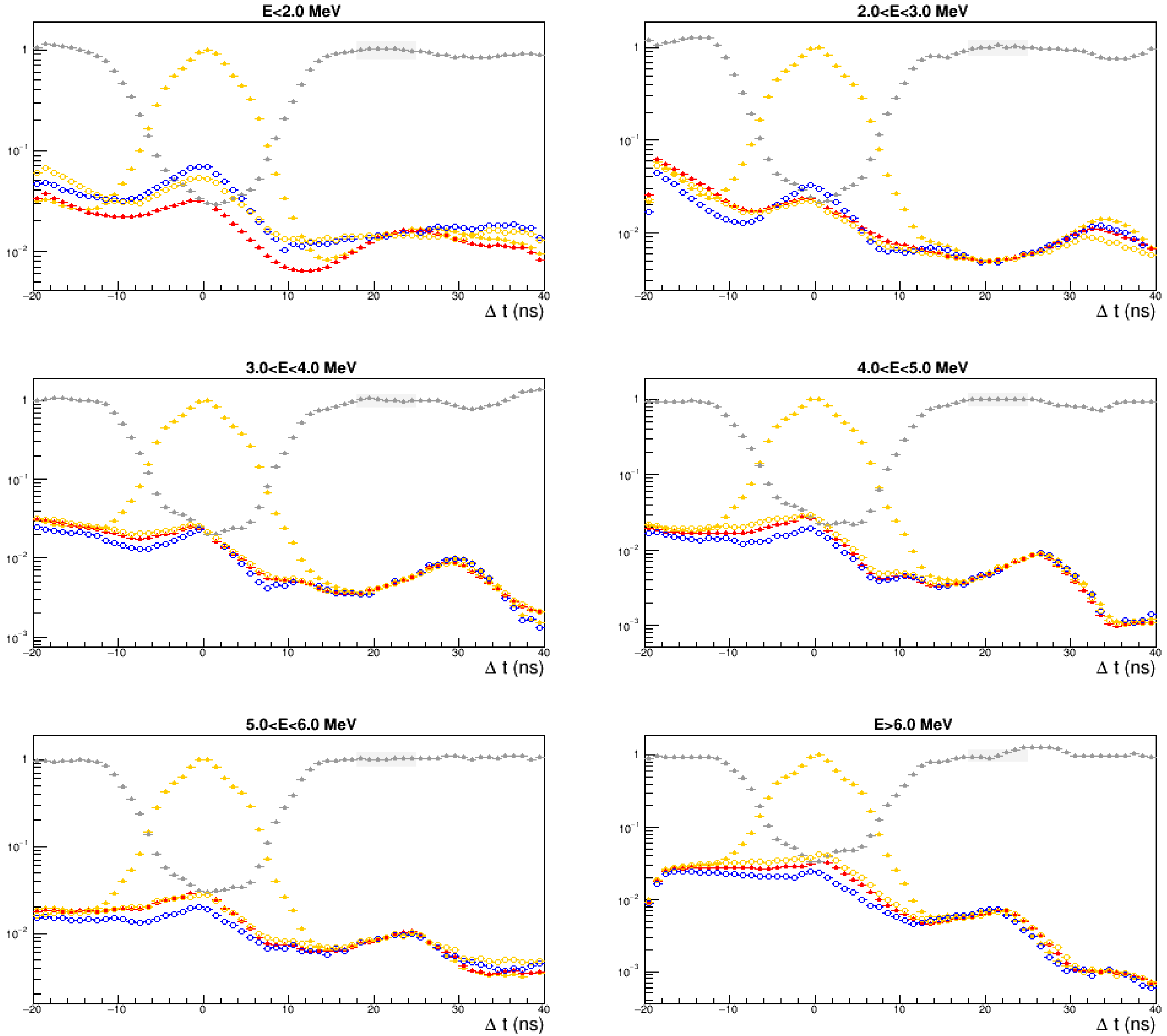
Fast polarization measurement with Carbon targets

Absolute normalization with polarized hydrogen jet target

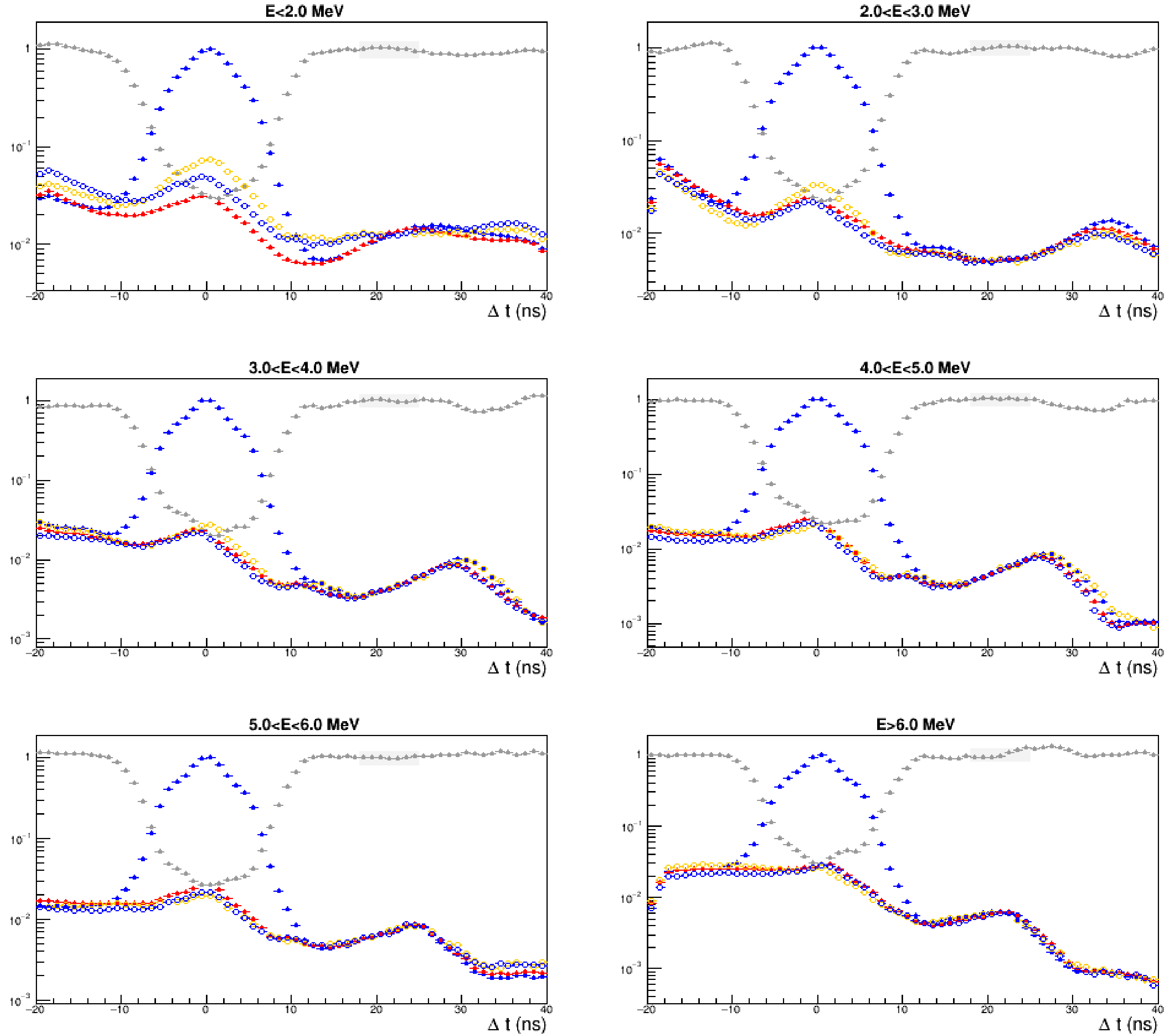
- Analyzing power with new detectors in 2015 → improved precision
- New asymmetries from elastic proton-heavy-ion scattering
- Determination of beam polarizations with background correction expected soon

BACK UP

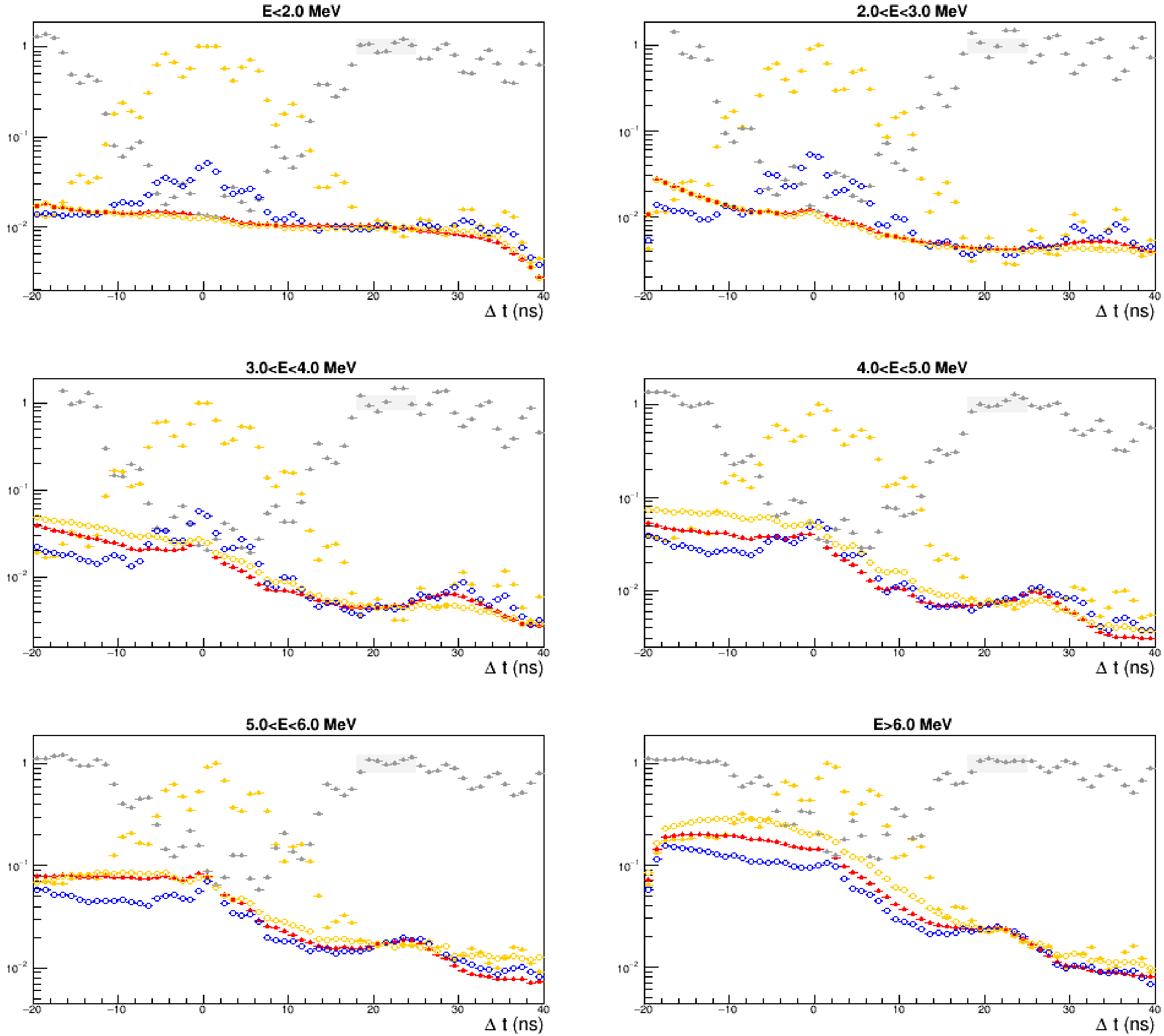
$\vec{p} + p$, yellow beam



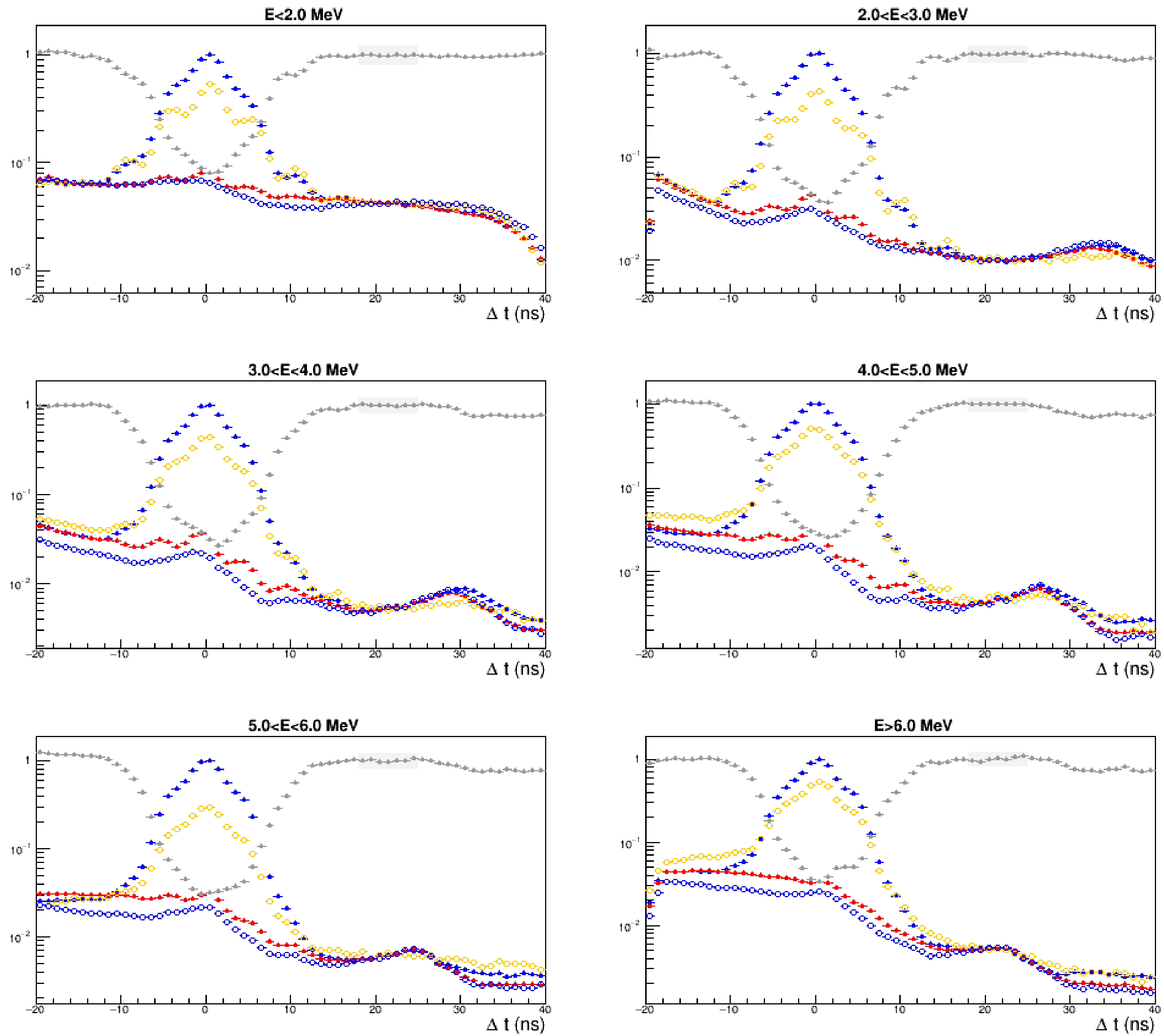
$\vec{p}+p$, blue beam



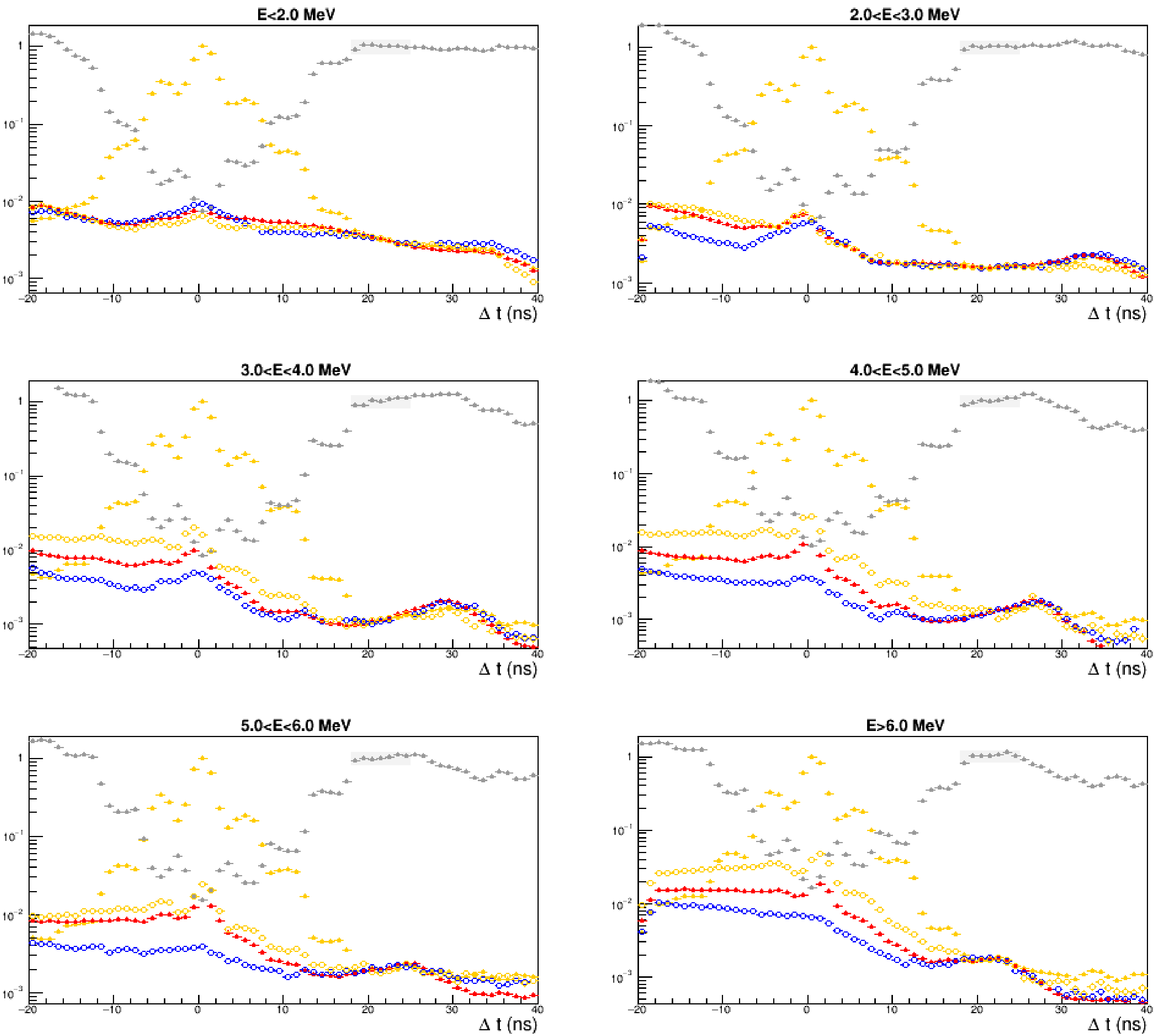
$\vec{p} + Au$, yellow beam



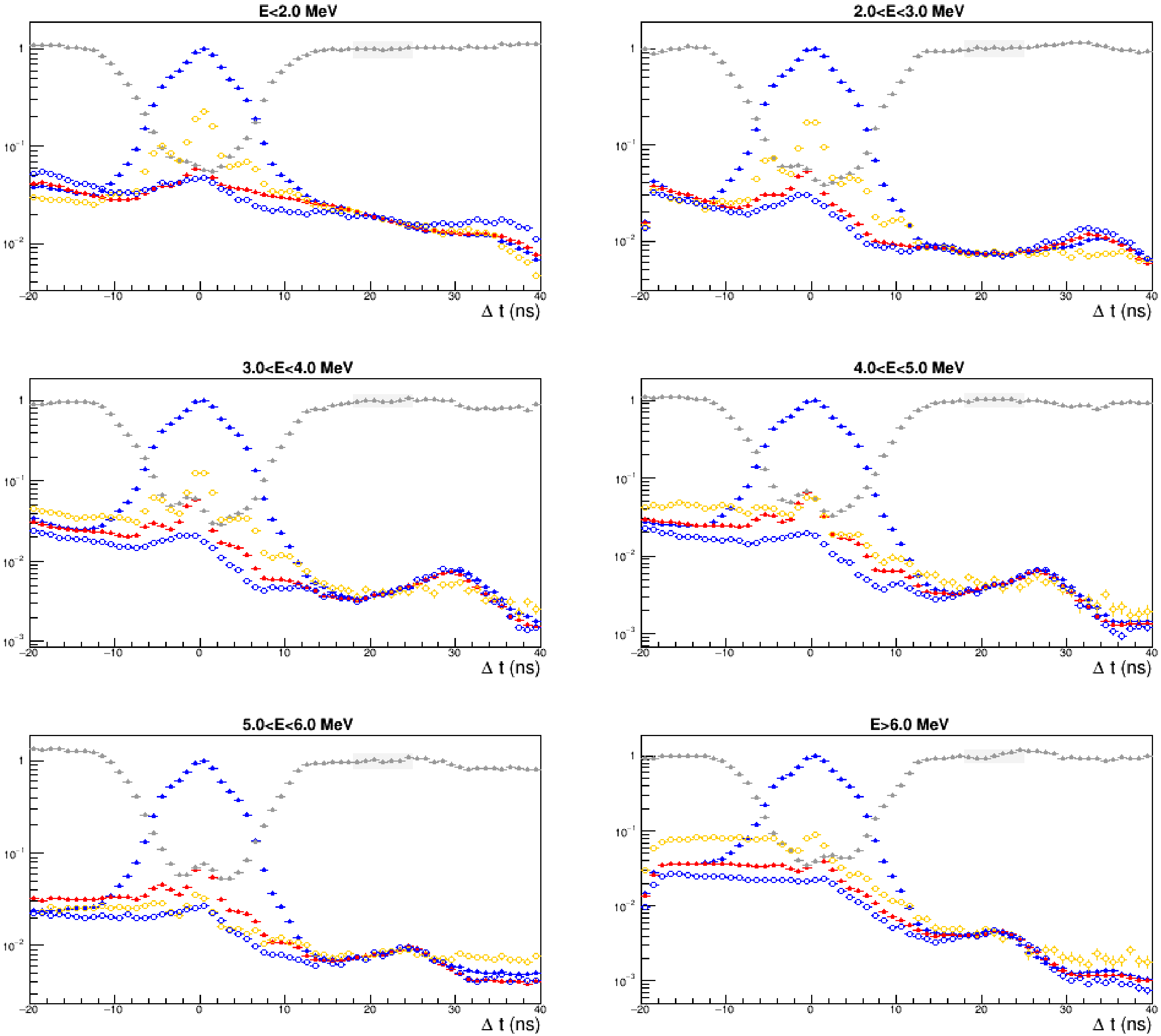
$\vec{p} + Au$, blue beam



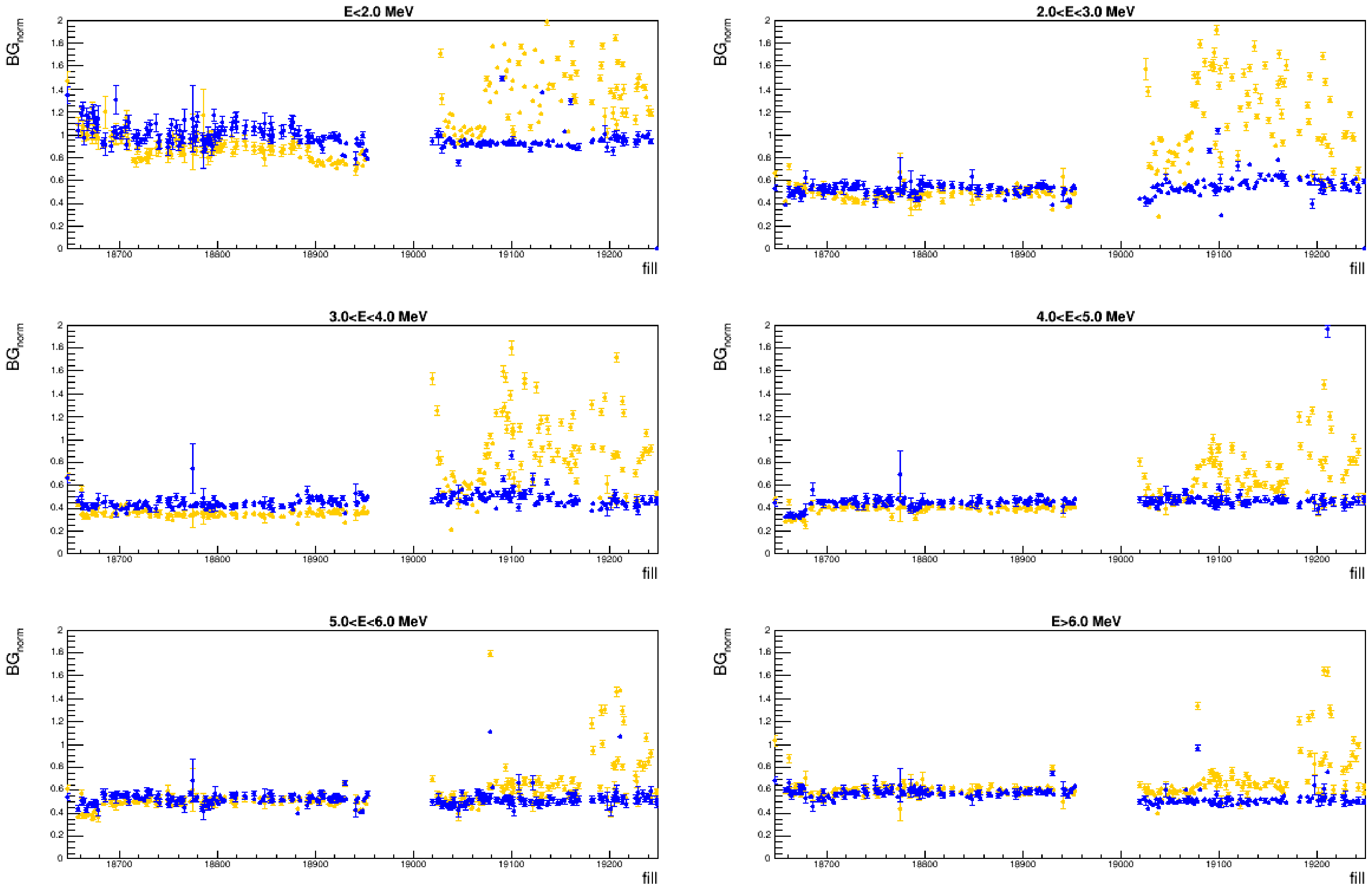
$\vec{p} + Al$, yellow beam



$\vec{p} + Al$, blue beam



Background Normalization ($18 < \Delta t < 25$ ns)



Background Fraction ($|\Delta t| < 5$ ns)