

Physics case for a polarised target for A Fixed Target ExpeRiment @ the LHC (AFTER@LHC)

Jean-Philippe Lansberg

IPN Orsay, CNRS/IN2P3, Université Paris-Sud

The 2015 International Workshop on Polarized Sources, Targets & Polarimetry
(PSTP2015)

14 - 18 September 2015

AFTER@LHC Study group: http://after.in2p3.fr/after/index.php/Current_author_list

Part I

Why a new fixed-target experiment for High-Energy Physics now ?

Decisive advantages of Fixed-target experiments

- Fixed-target experiments offer specific **advantages** that are still nowadays **difficult to challenge by collider experiments**

Decisive advantages of Fixed-target experiments

- Fixed-target experiments offer specific **advantages** that are still nowadays **difficult to challenge by collider experiments**
- They exhibit 4 decisive features,
 - accessing the **high** Feynman $|x_F|$ domain ($x_F \equiv \frac{p_z}{p_{z\max}}$)
 - achieving **high luminosities**,
 - **varying** the atomic mass of the **target** almost at will,
 - **polarising** the target.

Decisive advantages of Fixed-target experiments

- Fixed-target experiments offer specific **advantages** that are still nowadays **difficult to challenge by collider experiments**
- They exhibit 4 decisive features,
 - accessing the **high** Feynman $|x_F|$ domain ($x_F \equiv \frac{p_z}{p_{z \max}}$)
 - achieving **high luminosities**,
 - **varying** the atomic mass of the **target** almost at will,
 - **polarising** the target.
- which are **essential assets** to study
 - rare proton fluctuations at **large x**
 - vector boson production near threshold and other **rare processes**
 - **nuclear dependence** in heavy-ion collisions
 - observables involving **gluons** and the target **proton spin**

Why a fixed-target experiment at the LHC ?

- ADVANCE OUR UNDERSTANDING OF THE LARGE-X GLUON, ANTIQUARK AND HEAVY-QUARK CONTENT IN THE NUCLEON & NUCLEUS

Why a fixed-target experiment at the LHC ?

- ADVANCE OUR UNDERSTANDING OF THE LARGE-X GLUON, ANTIQUARK AND HEAVY-QUARK CONTENT IN THE NUCLEON & NUCLEUS
- Very large PDF uncertainties for $x \gtrsim 0.5$.
[could be crucial to characterise possible BSM discoveries]
- Proton **charm** content important to **high-energy neutrino & cosmic-rays** physics
- **EMC effect** is an open problem; studying a possible **gluon** EMC effect is essential
- Relevance of nuclear PDF to understand the **initial state of heavy-ion collisions**
- Search and study **rare proton fluctuations**
where one gluon carries most of the proton momentum

Why a fixed-target experiment at the LHC ?

· ADVANCE OUR UNDERSTANDING OF THE LARGE-X GLUON, ANTIQUARK AND HEAVY-QUARK CONTENT IN THE NUCLEON & NUCLEUS

- Very large PDF uncertainties for $x \gtrsim 0.5$.

[could be crucial to characterise possible BSM discoveries]

- Proton **charm** content important to **high-energy neutrino & cosmic-rays** physics
- **EMC effect** is an open problem; studying a possible **gluon** EMC effect is essential
- Relevance of nuclear PDF to understand the **initial state of heavy-ion collisions**
- Search and study **rare proton fluctuations**

where one gluon carries most of the proton momentum

· DYNAMICS AND SPIN OF GLUONS INSIDE (UN)POLARISED NUCLEONS

Why a fixed-target experiment at the LHC ?

· ADVANCE OUR UNDERSTANDING OF THE LARGE-X GLUON, ANTIQUARK AND HEAVY-QUARK CONTENT IN THE NUCLEON & NUCLEUS

- Very large PDF uncertainties for $x \gtrsim 0.5$.

[could be crucial to characterise possible BSM discoveries]

- Proton **charm** content important to **high-energy neutrino & cosmic-rays** physics
- **EMC effect** is an open problem; studying a possible **gluon** EMC effect is essential
- Relevance of nuclear PDF to understand the **initial state of heavy-ion collisions**
- Search and study **rare proton fluctuations**

where one gluon carries most of the proton momentum

· DYNAMICS AND SPIN OF GLUONS INSIDE (UN)POLARISED NUCLEONS

- Possible missing contribution to the **proton spin: orbital angular momentum**
- **Test** of the QCD **factorisation** framework
- Determination of the **linearly polarised gluons** in unpolarised protons

Why a fixed-target experiment at the LHC ?

· ADVANCE OUR UNDERSTANDING OF THE LARGE-X GLUON, ANTIQUARK AND HEAVY-QUARK CONTENT IN THE NUCLEON & NUCLEUS

- Very large PDF uncertainties for $x \gtrsim 0.5$.

[could be crucial to characterise possible BSM discoveries]

- Proton **charm** content important to **high-energy neutrino & cosmic-rays** physics
- **EMC effect** is an open problem; studying a possible **gluon** EMC effect is essential
- Relevance of nuclear PDF to understand the **initial state of heavy-ion collisions**
- Search and study **rare proton fluctuations**

where one gluon carries most of the proton momentum

· DYNAMICS AND SPIN OF GLUONS INSIDE (UN)POLARISED NUCLEONS

- Possible missing contribution to the **proton spin: orbital angular momentum**
- **Test** of the QCD **factorisation** framework
- Determination of the **linearly polarised gluons** in unpolarised protons

· HEAVY-ION COLLISIONS TOWARDS LARGE RAPIDITIES

Why a fixed-target experiment at the LHC ?

· ADVANCE OUR UNDERSTANDING OF THE LARGE-X GLUON, ANTIQUARK AND HEAVY-QUARK CONTENT IN THE NUCLEON & NUCLEUS

- Very large PDF uncertainties for $x \gtrsim 0.5$.

[could be crucial to characterise possible BSM discoveries]

- Proton **charm** content important to **high-energy neutrino & cosmic-rays** physics
- **EMC effect** is an open problem; studying a possible **gluon** EMC effect is essential
- Relevance of nuclear PDF to understand the **initial state of heavy-ion collisions**
- Search and study **rare proton fluctuations**

where one gluon carries most of the proton momentum

· DYNAMICS AND SPIN OF GLUONS INSIDE (UN)POLARISED NUCLEONS

- Possible missing contribution to the **proton spin: orbital angular momentum**
- **Test** of the QCD **factorisation** framework
- Determination of the **linearly polarised gluons** in unpolarised protons

· HEAVY-ION COLLISIONS TOWARDS LARGE RAPIDITIES

- Explore the **longitudinal expansion** of QGP formation with **new hard probes**
- Test the **factorisation** of cold nuclear effects **from $p + A$ to $A + B$** collisions
- Test the formation of **azimuthal asymmetries**: hydrodynamics vs. initial-state radiation

Part II

A fixed-target experiment using the LHC beam(s): AFTER@LHC

Generalities

- pp or pA collisions with a 7 TeV p^+ on a fixed target occur at a CM energy

$$\sqrt{s} = \sqrt{2m_N E_p} \simeq 115 \text{ GeV}$$

Generalities

- pp or pA collisions with a 7 TeV p^+ on a fixed target occur at a CM energy

$$\sqrt{s} = \sqrt{2m_N E_p} \simeq 115 \text{ GeV}$$

- In a symmetric collider mode, $\sqrt{s} = 2E_p$, *i.e.* much larger

Generalities

- pp or pA collisions with a 7 TeV p^+ on a fixed target occur at a CM energy

$$\sqrt{s} = \sqrt{2m_N E_p} \simeq 115 \text{ GeV}$$

- In a symmetric collider mode, $\sqrt{s} = 2E_p$, *i.e.* much larger
- Benefit of the fixed target mode : **boost**: $\gamma_{CM}^{Lab} = \frac{\sqrt{s}}{2m_p} \simeq 60$

Generalities

- pp or pA collisions with a 7 TeV p^+ on a fixed target occur at a CM energy

$$\sqrt{s} = \sqrt{2m_N E_p} \simeq 115 \text{ GeV}$$

- In a symmetric collider mode, $\sqrt{s} = 2E_p$, *i.e.* much larger
- Benefit of the fixed target mode : **boost**: $\gamma_{CM}^{Lab} = \frac{\sqrt{s}}{2m_p} \simeq 60$
- Rather soft particles in the CM are in principle detectable

Generalities

- pp or pA collisions with a 7 TeV p^+ on a fixed target occur at a CM energy

$$\sqrt{s} = \sqrt{2m_N E_p} \simeq 115 \text{ GeV}$$

- In a symmetric collider mode, $\sqrt{s} = 2E_p$, *i.e.* much larger
- Benefit of the fixed target mode : **boost**: $\gamma_{CM}^{Lab} = \frac{\sqrt{s}}{2m_p} \simeq 60$
- Rather soft particles in the CM are in principle detectable
- Angle in the Lab. frame: $\tan \theta = \frac{p_T}{p_{z,Lab}} = \frac{1}{\gamma\beta} \Rightarrow \theta \simeq 1^\circ$.

[Rapidity shift: $\Delta y = \tanh^{-1}\beta \simeq 4.8$]

Generalities

- pp or pA collisions with a 7 TeV p^+ on a fixed target occur at a CM energy

$$\sqrt{s} = \sqrt{2m_N E_p} \simeq 115 \text{ GeV}$$

- In a symmetric collider mode, $\sqrt{s} = 2E_p$, *i.e.* much larger

- Benefit of the fixed target mode : **boost**: $\gamma_{CM}^{Lab} = \frac{\sqrt{s}}{2m_p} \simeq 60$

- Rather soft particles in the CM are in principle detectable

- Angle in the Lab. frame: $\tan \theta = \frac{p_T}{p_{z,Lab}} = \frac{1}{\gamma\beta} \Rightarrow \theta \simeq 1^\circ$.

[Rapidity shift: $\Delta y = \tanh^{-1}\beta \simeq 4.8$]

- The entire forward CM hemisphere ($y_{CM} > 0$) within $0^\circ \leq \theta_{Lab} \leq 1^\circ$

Generalities

- pp or pA collisions with a 7 TeV p^+ on a fixed target occur at a CM energy

$$\sqrt{s} = \sqrt{2m_N E_p} \simeq 115 \text{ GeV}$$

- In a symmetric collider mode, $\sqrt{s} = 2E_p$, *i.e.* much larger

- Benefit of the fixed target mode : **boost**: $\gamma_{CM}^{Lab} = \frac{\sqrt{s}}{2m_p} \simeq 60$

- Rather soft particles in the CM are in principle detectable

- Angle in the Lab. frame: $\tan \theta = \frac{p_T}{p_{z,Lab}} = \frac{1}{\gamma\beta} \Rightarrow \theta \simeq 1^\circ$.

[Rapidity shift: $\Delta y = \tanh^{-1}\beta \simeq 4.8$]

- The entire forward CM hemisphere ($y_{CM} > 0$) within $0^\circ \leq \theta_{Lab} \leq 1^\circ$

- **Good thing**: small forward detector \equiv large acceptance

Generalities

- pp or pA collisions with a 7 TeV p^+ on a fixed target occur at a CM energy

$$\sqrt{s} = \sqrt{2m_N E_p} \simeq 115 \text{ GeV}$$

- In a symmetric collider mode, $\sqrt{s} = 2E_p$, *i.e.* much larger

- Benefit of the fixed target mode : **boost**: $\gamma_{CM}^{Lab} = \frac{\sqrt{s}}{2m_p} \simeq 60$

- Rather soft particles in the CM are in principle detectable

- Angle in the Lab. frame: $\tan \theta = \frac{p_T}{p_{z,Lab}} = \frac{1}{\gamma\beta} \Rightarrow \theta \simeq 1^\circ$.

[Rapidity shift: $\Delta y = \tanh^{-1}\beta \simeq 4.8$]

- The entire forward CM hemisphere ($y_{CM} > 0$) within $0^\circ \leq \theta_{Lab} \leq 1^\circ$

- **Good thing**: small forward detector \equiv large acceptance

Generalities

- pp or pA collisions with a 7 TeV p^+ on a fixed target occur at a CM energy

$$\sqrt{s} = \sqrt{2m_N E_p} \simeq 115 \text{ GeV}$$

- In a symmetric collider mode, $\sqrt{s} = 2E_p$, *i.e.* much larger

- Benefit of the fixed target mode : **boost**: $\gamma_{CM}^{Lab} = \frac{\sqrt{s}}{2m_p} \simeq 60$

- Rather soft particles in the CM are in principle detectable

- Angle in the Lab. frame: $\tan \theta = \frac{p_T}{p_{z,Lab}} = \frac{1}{\gamma\beta} \Rightarrow \theta \simeq 1^\circ$.

[Rapidity shift: $\Delta y = \tanh^{-1}\beta \simeq 4.8$]

- The entire forward CM hemisphere ($y_{CM} > 0$) within $0^\circ \leq \theta_{Lab} \leq 1^\circ$

- **Good thing**: small forward detector \equiv large acceptance

- **Bad thing**: high multiplicity \Rightarrow absorber \Rightarrow physics limitation

Boost effect: LHCb becomes a backward detector

- Because of the boost $y_{CM} = 0 \Rightarrow y_{Lab} \simeq 4.8$

Boost effect: LHCb becomes a backward detector

- Because of the boost $y_{CM} = 0 \Rightarrow y_{Lab} \simeq 4.8$
- The pseudo-rapidity coverage of LHCb, $2 \leq \eta \leq 5$, approximately translates to a rapidity coverage in the CM of roughly $-2.8 \leq y_{CM} \leq 0.2$
- LHC muon arm: $2.5 \leq \eta \leq 4 \Rightarrow -2.3 \leq y_{CM} \leq -0.8$

Boost effect: LHCb becomes a backward detector

- Because of the boost $y_{CM} = 0 \Rightarrow y_{Lab} \simeq 4.8$
- The pseudo-rapidity coverage of LHCb, $2 \leq \eta \leq 5$, approximately translates to a rapidity coverage in the *CM* of roughly $-2.8 \leq y_{CM} \leq 0.2$
- LHC muon arm: $2.5 \leq \eta \leq 4 \Rightarrow -2.3 \leq y_{CM} \leq -0.8$
- As a comparison, the PHENIX detector with its forward and backward muons arm *only* goes up to $|y_{CM}| \lesssim 2.2$

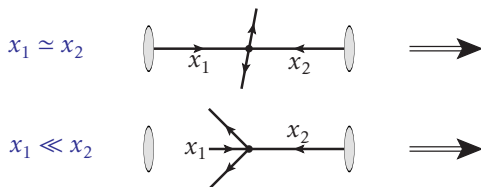
Boost effect: LHCb becomes a backward detector

- Because of the boost $y_{CM} = 0 \Rightarrow y_{Lab} \simeq 4.8$
- The pseudo-rapidity coverage of LHCb, $2 \leq \eta \leq 5$, approximately translates to a rapidity coverage in the CM of roughly $-2.8 \leq y_{CM} \leq 0.2$
- LHC muon arm: $2.5 \leq \eta \leq 4 \Rightarrow -2.3 \leq y_{CM} \leq -0.8$
- As a comparison, the PHENIX detector with its forward and backward muons arm *only* goes up to $|y_{CM}| \lesssim 2.2$
- In addition, there are advantages to go there:
 - reduced multiplicities at large(r) angles
 - **access to partons with momentum fraction $x \rightarrow 1$ in the target**

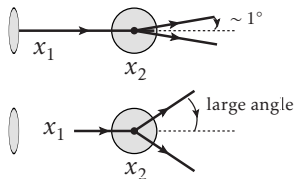
Boost effect: LHCb becomes a backward detector

- Because of the boost $y_{CM} = 0 \Rightarrow y_{Lab} \simeq 4.8$
- The pseudo-rapidity coverage of LHCb, $2 \leq \eta \leq 5$, approximately translates to a rapidity coverage in the CM of roughly $-2.8 \leq y_{CM} \leq 0.2$
- LHC muon arm: $2.5 \leq \eta \leq 4 \Rightarrow -2.3 \leq y_{CM} \leq -0.8$
- As a comparison, the PHENIX detector with its forward and backward muons arm *only* goes up to $|y_{CM}| \lesssim 2.2$
- In addition, there are advantages to go there:
 - reduced multiplicities at large(r) angles
 - **access to partons with momentum fraction $x \rightarrow 1$ in the target**

Hadron center-of-mass system



Target rest frame

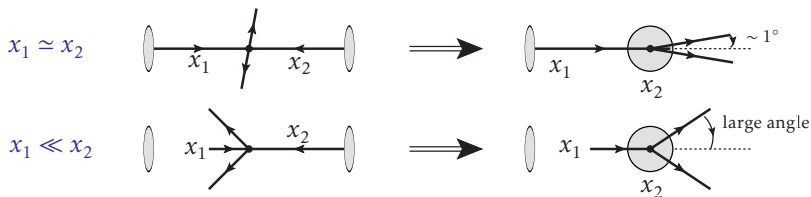


Boost effect: LHCb becomes a backward detector

- Because of the boost $y_{CM} = 0 \Rightarrow y_{Lab} \simeq 4.8$
- The pseudo-rapidity coverage of LHCb, $2 \leq \eta \leq 5$, approximately translates to a rapidity coverage in the CM of roughly $-2.8 \leq y_{CM} \leq 0.2$
- LHC muon arm: $2.5 \leq \eta \leq 4 \Rightarrow -2.3 \leq y_{CM} \leq -0.8$
- As a comparison, the PHENIX detector with its forward and backward muons arm *only* goes up to $|y_{CM}| \lesssim 2.2$
- In addition, there are advantages to go there:
 - reduced multiplicities at large(r) angles
 - **access to partons with momentum fraction $x \rightarrow 1$ in the target**

Hadron center-of-mass system

Target rest frame



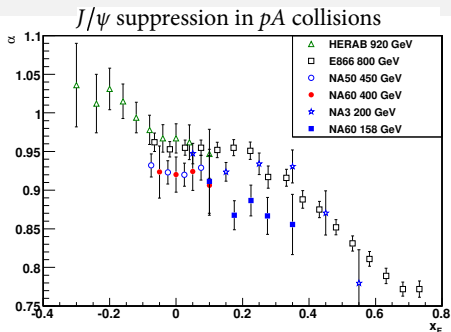
backward physics = large- x_2 physics

First systematic access to the target-rapidity region

($x_F \rightarrow -1$)

First systematic access to the target-rapidity region

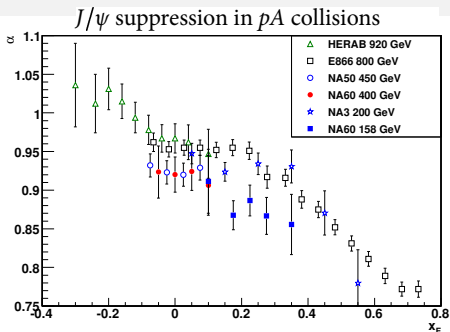
($x_F \rightarrow -1$)



- x_F systematically studied at fixed target experiments up to +1

First systematic access to the target-rapidity region

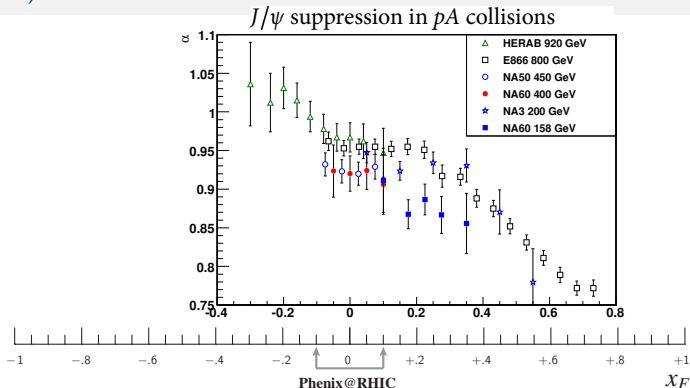
($x_F \rightarrow -1$)



- x_F systematically studied at fixed target experiments **up to +1**
- Hera-B was the only one to really explore $x_F < 0$, up to -0.3

First systematic access to the target-rapidity region

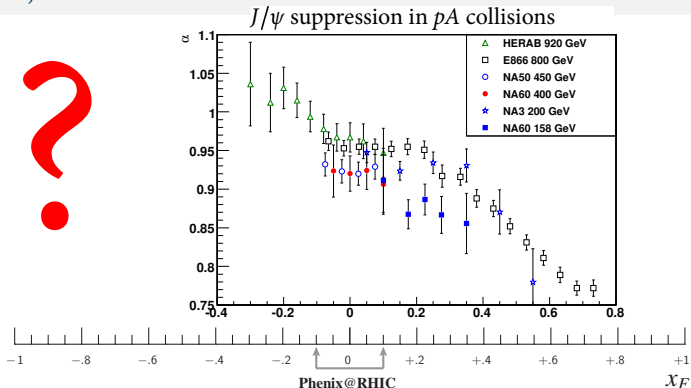
($x_F \rightarrow -1$)



- x_F systematically studied at fixed target experiments up to +1
- Hera-B was the only one to really explore $x_F < 0$, up to -0.3
- PHENIX @ RHIC: $-0.1 < x_F < 0.1$ [could be wider with Υ , but low stat.]
- CMS/ATLAS: $|x_F| < 5 \cdot 10^{-3}$; LHCb-collider: $5 \cdot 10^{-3} < x_F < 4 \cdot 10^{-2}$

First systematic access to the target-rapidity region

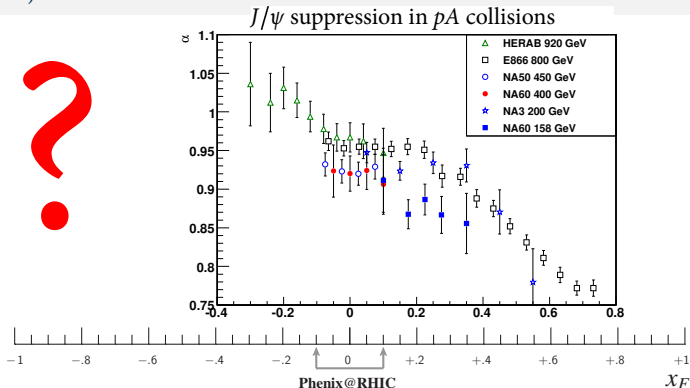
($x_F \rightarrow -1$)



- x_F systematically studied at fixed target experiments up to +1
- Hera-B was the only one to really explore $x_F < 0$, up to -0.3
- PHENIX @ RHIC: $-0.1 < x_F < 0.1$ [could be wider with Υ , but low stat.]
- CMS/ATLAS: $|x_F| < 5 \cdot 10^{-3}$; LHCb-collider: $5 \cdot 10^{-3} < x_F < 4 \cdot 10^{-2}$

First systematic access to the target-rapidity region

($x_F \rightarrow -1$)



- x_F systematically studied at fixed target experiments up to +1
- Hera-B was the only one to really explore $x_F < 0$, up to -0.3
- PHENIX @ RHIC: $-0.1 < x_F < 0.1$ [could be wider with Υ , but low stat.]
- CMS/ATLAS: $|x_F| < 5 \cdot 10^{-3}$; LHCb-collider: $5 \cdot 10^{-3} < x_F < 4 \cdot 10^{-2}$
- If we measure $\Upsilon(b\bar{b})$ at $y_{\text{cms}} \simeq -2.5 \Rightarrow x_F \simeq \frac{2m_Y}{\sqrt{s}} \sinh(y_{\text{cms}}) \simeq -1$

Part III

Colliding the LHC beams on fixed targets: 2 options

The extracted-beam option

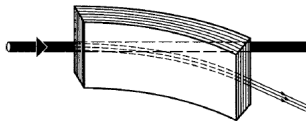
- ★ The LHC beam may be extracted using “Strong crystalline field”
without any decrease in performance of the LHC !

E. Uggerhøj, U.I Uggerhøj, NIM B 234 (2005) 31, Rev. Mod. Phys. 77 (2005) 1131

The extracted-beam option

- ★ The LHC beam may be extracted using “Strong crystalline field”
without any decrease in performance of the LHC !

E. Uggerhøj, U.I Uggerhøj, NIM B 234 (2005) 31, Rev. Mod. Phys. 77 (2005) 1131

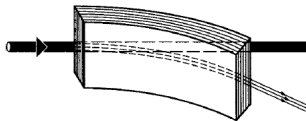


The extracted-beam option

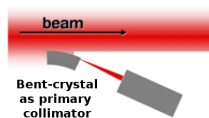
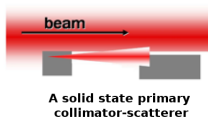
- ★ The LHC beam may be extracted using “Strong crystalline field”

without any decrease in performance of the LHC !

E. Uggerhøj, U.I Uggerhøj, NIM B 234 (2005) 31, Rev. Mod. Phys. 77 (2005) 1131



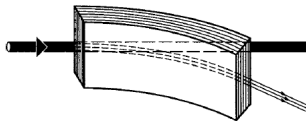
- ★ Illustration for collimation



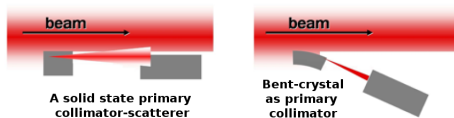
The extracted-beam option

- ★ The LHC beam may be extracted using “Strong crystalline field”
without any decrease in performance of the LHC !

E. Uggerhøj, U.I Uggerhøj, NIM B 234 (2005) 31, Rev. Mod. Phys. 77 (2005) 1131



- ★ **Illustration for collimation**



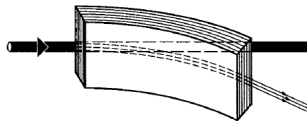
- ★ **Tests** will be performed on the **LHC beam**:

LUA9 proposal approved by the LHCC

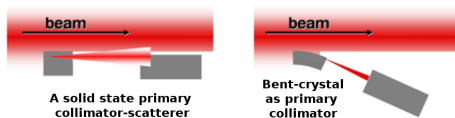
The extracted-beam option

- ★ The LHC beam may be extracted using “Strong crystalline field”
without any decrease in performance of the LHC !

E. Uggerhøj, U.I Uggerhøj, NIM B 234 (2005) 31, Rev. Mod. Phys. 77 (2005) 1131



- ★ **Illustration for collimation**



- ★ **Tests** will be performed on the **LHC beam**:

LUA9 proposal approved by the LHCC

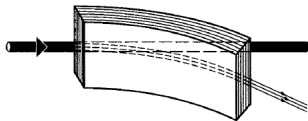
- ★ 2 crystals and 2 goniometers **already installed** in the LHC beampipe

The extracted-beam option

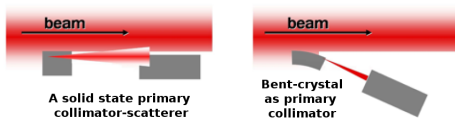
- ★ The LHC beam may be extracted using “Strong crystalline field”

without any decrease in performance of the LHC !

E. Uggerhøj, U.I Uggerhøj, NIM B 234 (2005) 31, Rev. Mod. Phys. 77 (2005) 1131



- ★ **Illustration for collimation**



- ★ **Tests** will be performed on the **LHC beam**:

LUA9 proposal approved by the LHCC

- ★ 2 crystals and 2 goniometers **already installed** in the LHC beampipe

- ★ CRYSBREAM: ERC funded project to extract the LHC beams

with a bent crystal (G. Cavoto - Rome)

Luminosities with extracted-proton beams

- Expected **proton flux** $\Phi_{beam} = 5 \times 10^8 p^+ s^{-1}$

Luminosities with extracted-proton beams

- Expected **proton flux** $\Phi_{beam} = 5 \times 10^8 p^+ s^{-1}$
- Instantaneous **Luminosity**:

$$\mathcal{L} = \Phi_{beam} \times N_{target} = N_{beam} \times (\rho \times \ell \times \mathcal{N}_A) / A$$

[ℓ : target thickness (for instance 1cm)]

Luminosities with extracted-proton beams

- Expected **proton flux** $\Phi_{beam} = 5 \times 10^8 p^+ s^{-1}$
- Instantaneous **Luminosity**:

$$\mathcal{L} = \Phi_{beam} \times N_{target} = N_{beam} \times (\rho \times \ell \times \mathcal{N}_A) / A$$

[ℓ : target thickness (for instance 1cm)]

- Integrated luminosity: $\int dt \mathcal{L}$ over 10^7 s for p^+ and 10^6 for Pb

[the so-called LHC years]

Luminosities with extracted-proton beams

- Expected **proton flux** $\Phi_{beam} = 5 \times 10^8 p^+ s^{-1}$
- Instantaneous **Luminosity**:

$$\mathcal{L} = \Phi_{beam} \times N_{target} = N_{beam} \times (\rho \times \ell \times \mathcal{N}_A) / A$$

[ℓ : target thickness (for instance 1cm)]

- Integrated luminosity: $\int dt \mathcal{L}$ over 10^7 s for p^+ and 10^6 for Pb

[the so-called LHC years]

Target	ρ (g.cm ⁻³)	A	\mathcal{L} ($\mu\text{b}^{-1}.\text{s}^{-1}$)	$\int \mathcal{L}$ (fb ⁻¹ .yr ⁻¹)
1m Liq. H₂	0.07	1	2000	20
1m Liq. D₂	0.16	2	2400	24
1cm Be	1.85	9	.62	.62
1cm Cu	8.96	64	.42	.42
1cm W	19.1	185	.31	.31
1cm Pb	11.35	207	.16	.16

Luminosities with extracted-proton beams

- Expected **proton flux** $\Phi_{beam} = 5 \times 10^8 p^+ s^{-1}$
- Instantaneous **Luminosity**:

$$\mathcal{L} = \Phi_{beam} \times N_{target} = N_{beam} \times (\rho \times \ell \times \mathcal{N}_A) / A$$

[ℓ : target thickness (for instance 1cm)]

- Integrated luminosity: $\int dt \mathcal{L}$ over 10^7 s for p^+ and 10^6 for Pb

[the so-called LHC years]

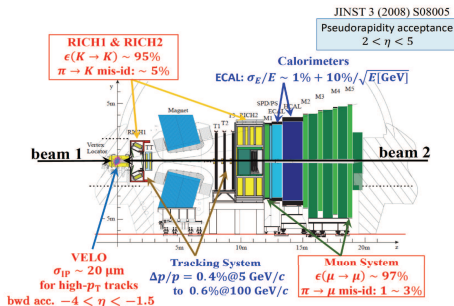
Target	ρ (g.cm ⁻³)	A	\mathcal{L} ($\mu\text{b}^{-1} \cdot \text{s}^{-1}$)	$\int \mathcal{L}$ ($\text{fb}^{-1} \cdot \text{yr}^{-1}$)
1m Liq. H₂	0.07	1	2000	20
1m Liq. D₂	0.16	2	2400	24
1cm Be	1.85	9	62	.62
1cm Cu	8.96	64	42	.42
1cm W	19.1	185	31	.31
1cm Pb	11.35	207	16	.16

- For pp and $p\bar{d}$ collisions : $\mathcal{L}_{H_2/D_2} \simeq 20 \text{ fb}^{-1} \text{ yr}^{-1}$

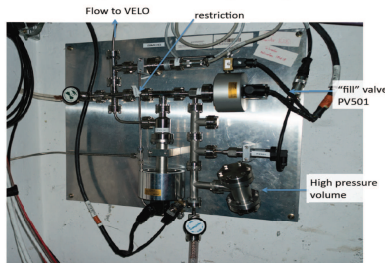
3 orders of magnitude larger than RHIC (200 GeV)

SMOG@LHCb: the first step towards an internal (polarised) target ?

SMOG@LHCb: the first step towards an internal (polarised) target ?

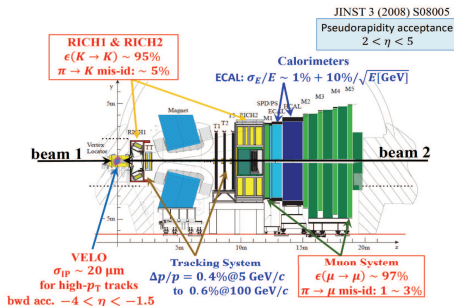


SMOG: System for Measuring Overlap with Gas

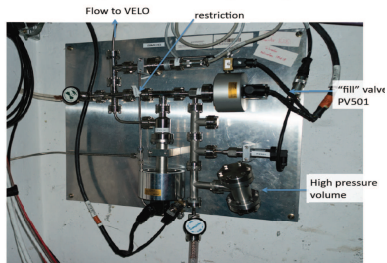


→ injection of Ne-gas into VELO

SMOG@LHCb: the first step towards an internal (polarised) target ?



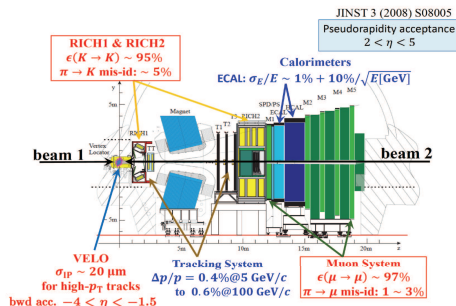
SMOG: System for Measuring Overlap with Gas



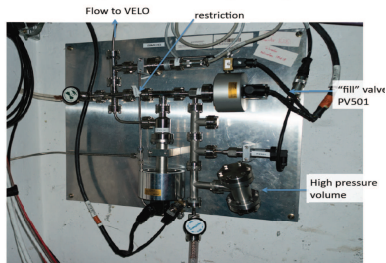
→ injection of Ne-gas into VELO

- Initially: low density Ne-gas injected into LHCb Vertex Locator [LHCb-CONF-2012-034]

SMOG@LHCb: the first step towards an internal (polarised) target ?



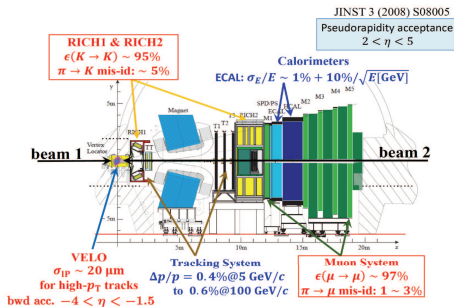
SMOG: System for Measuring Overlap with Gas



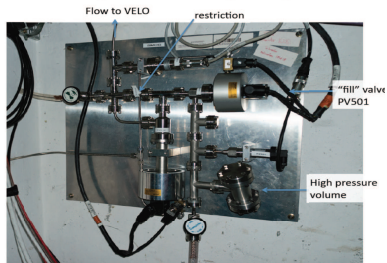
→ injection of Ne-gas into VELO

- Initially: low density Ne-gas injected into LHCb Vertex Locator [LHCb-CONF-2012-034]
- Short pilot runs: 2012 $p\text{Ne}$ at $\sqrt{s_{NN}} = 87 \text{ GeV}$ & 2013 PbNe at $\sqrt{s_{NN}} = 54 \text{ GeV}$

SMOG@LHCb: the first step towards an internal (polarised) target ?



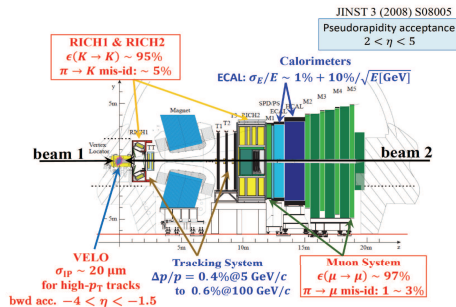
SMOG: System for Measuring Overlap with Gas



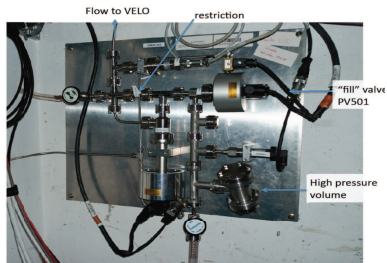
→ injection of Ne-gas into VELO

- Initially: low density Ne-gas injected into LHCb Vertex Locator [LHCb-CONF-2012-034]
- Short pilot runs: 2012 $p\text{Ne}$ at $\sqrt{s_{NN}} = 87 \text{ GeV}$ & 2013 PbNe at $\sqrt{s_{NN}} = 54 \text{ GeV}$
- 12 hours of $p\text{Ne}$ and 8 hours $p\text{He}$ (09/2015); 3 days of $p\text{Ar}$ in (10/2015)
- 3 weeks of PbAr (12/2015)

SMOG@LHCb: the first step towards an internal (polarised) target ?



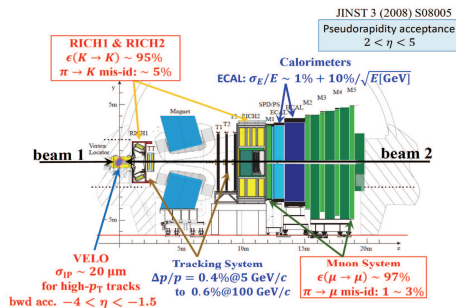
SMOG: System for Measuring Overlap with Gas



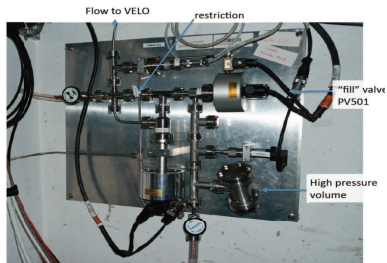
→ injection of Ne-gas into VELO

- Initially: low density Ne-gas injected into LHCb Vertex Locator [LHCb-CONF-2012-034]
- Short pilot runs: 2012 $p\text{Ne}$ at $\sqrt{s_{NN}} = 87 \text{ GeV}$ & 2013 PbNe at $\sqrt{s_{NN}} = 54 \text{ GeV}$
- 12 hours of $p\text{Ne}$ and 8 hours $p\text{He}$ (09/2015); 3 days of $p\text{Ar}$ in (10/2015)
- 3 weeks of PbAr (12/2015)
- Noble gases favoured

SMOG@LHCb: the first step towards an internal (polarised) target ?



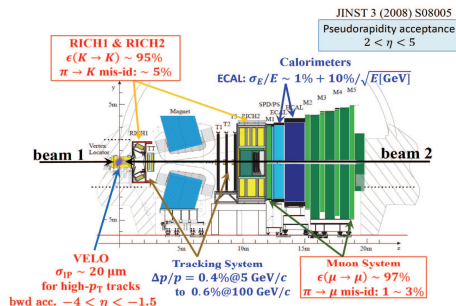
SMOG: System for Measuring Overlap with Gas



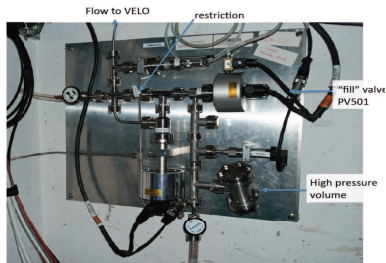
→ injection of Ne-gas into VELO

- Initially: low density Ne-gas injected into LHCb Vertex Locator [LHCb-CONF-2012-034]
- Short pilot runs: 2012 $p\text{Ne}$ at $\sqrt{s_{NN}} = 87 \text{ GeV}$ & 2013 PbNe at $\sqrt{s_{NN}} = 54 \text{ GeV}$
- 12 hours of $p\text{Ne}$ and 8 hours $p\text{He}$ (09/2015); 3 days of $p\text{Ar}$ in (10/2015)
- 3 weeks of PbAr (12/2015)
- Noble gases favoured
- Target unpolarised with the current SMOG system

SMOG@LHCb: the first step towards an internal (polarised) target ?



SMOG: System for Measuring Overlap with Gas



→ injection of Ne-gas into VELO

- Initially: low density Ne-gas injected into LHCb Vertex Locator [LHCb-CONF-2012-034]
- Short pilot runs: 2012 $p\text{Ne}$ at $\sqrt{s_{NN}} = 87 \text{ GeV}$ & 2013 PbNe at $\sqrt{s_{NN}} = 54 \text{ GeV}$
- 12 hours of $p\text{Ne}$ and 8 hours $p\text{He}$ (09/2015); 3 days of $p\text{Ar}$ in (10/2015)
- 3 weeks of PbAr (12/2015)
- Noble gases favoured
- Target unpolarised with the current SMOG system
- SMOG test : no decrease of LHC performances observed

Luminosities with the internal-gas-target option

Luminosities with the internal-gas-target option

- Instantaneous Luminosity: $\mathcal{L} = \Phi_{beam} \times N_{target} = N_{beam} \times (\rho \times \ell \times \mathcal{N}_A) / A$

Luminosities with the internal-gas-target option

- Instantaneous Luminosity: $\mathcal{L} = \Phi_{beam} \times N_{target} = N_{beam} \times (\rho \times \ell \times \mathcal{N}_A) / A$
- $\Phi_{p^+} = 3.2 \times 10^{14} p^+ \times 11000 \text{Hz} = 3.5 \times 10^{18} p^+ \text{ s}^{-1}$
- $\Phi_{Pb} = 4.2 \times 10^{10} p^+ \times 11000 \text{Hz} = 4.6 \times 10^{14} \text{Pb s}^{-1}$

[1/2 Ampère !]

Luminosities with the internal-gas-target option

- Instantaneous Luminosity: $\mathcal{L} = \Phi_{beam} \times N_{target} = N_{beam} \times (\rho \times \ell \times \mathcal{N}_A) / A$
- $\Phi_{p^+} = 3.2 \times 10^{14} p^+ \times 11000 \text{Hz} = 3.5 \times 10^{18} p^+ \text{ s}^{-1}$
- $\Phi_{Pb} = 4.2 \times 10^{10} p^+ \times 11000 \text{Hz} = 4.6 \times 10^{14} \text{Pb s}^{-1}$
- Usable gas zone ℓ , up to 100 cm

[1/2 Ampère !]

Luminosities with the internal-gas-target option

- Instantaneous Luminosity: $\mathcal{L} = \Phi_{beam} \times N_{target} = N_{beam} \times (\rho \times \ell \times \mathcal{N}_A) / A$
- $\Phi_{p^+} = 3.2 \times 10^{14} p^+ \times 11000 \text{ Hz} = 3.5 \times 10^{18} p^+ \text{ s}^{-1}$ [1/2 Ampère !]
- $\Phi_{Pb} = 4.2 \times 10^{10} p^+ \times 11000 \text{ Hz} = 4.6 \times 10^{14} \text{ Pb s}^{-1}$
- Usable gas zone ℓ , up to 100 cm
- Target density: $\frac{\rho}{P} = c = \frac{A}{22400} \text{ bar}^{-1} \text{ g cm}^{-3} \Rightarrow \mathcal{L} = \Phi_{beam} \times \left(\frac{\mathcal{N}_A}{22400} \times P \times \ell \right)$
[1 mole of a perfect gas occupies 22 400 cm³ at 273 K and 1 bar]

Luminosities with the internal-gas-target option

- Instantaneous Luminosity: $\mathcal{L} = \Phi_{beam} \times N_{target} = N_{beam} \times (\rho \times \ell \times \mathcal{N}_A) / A$
- $\Phi_{p^+} = 3.2 \times 10^{14} p^+ \times 11000 \text{ Hz} = 3.5 \times 10^{18} p^+ \text{ s}^{-1}$ [1/2 Ampère !]
- $\Phi_{Pb} = 4.2 \times 10^{10} p^+ \times 11000 \text{ Hz} = 4.6 \times 10^{14} \text{ Pb s}^{-1}$
- Usable gas zone ℓ , up to 100 cm
- Target density: $\frac{\rho}{P} = c = \frac{A}{22400} \text{ bar}^{-1} \text{ g cm}^{-3} \Rightarrow \mathcal{L} = \Phi_{beam} \times \left(\frac{\mathcal{N}_A}{22400} \times P \times \ell \right)$
[1 mole of a perfect gas occupies 22 400 cm³ at 273 K and 1 bar]
- For $P = 10^{-9} \text{ bar}$ [7x that of SMOG in 2015, the 'vacuum' is 10^{-12} bar], $\mathcal{L}_{pX(\text{PbX})} = 10(10^{-3}) \mu\text{b}^{-1} \text{ s}^{-1}$

Luminosities with the internal-gas-target option

- Instantaneous Luminosity: $\mathcal{L} = \Phi_{beam} \times N_{target} = N_{beam} \times (\rho \times \ell \times \mathcal{N}_A) / A$
- $\Phi_{p^+} = 3.2 \times 10^{14} p^+ \times 11000 \text{ Hz} = 3.5 \times 10^{18} p^+ \text{ s}^{-1}$ [1/2 Ampère !]
- $\Phi_{Pb} = 4.2 \times 10^{10} p^+ \times 11000 \text{ Hz} = 4.6 \times 10^{14} \text{ Pb s}^{-1}$
- Usable gas zone ℓ , up to 100 cm
- Target density: $\frac{\rho}{P} = c = \frac{A}{22400} \text{ bar}^{-1} \text{ g cm}^{-3} \Rightarrow \mathcal{L} = \Phi_{beam} \times \left(\frac{\mathcal{N}_A}{22400} \times P \times \ell \right)$
[1 mole of a perfect gas occupies 22 400 cm³ at 273 K and 1 bar]
- For $P = 10^{-9} \text{ bar}$ [7x that of SMOG in 2015, the 'vacuum' is 10^{-12} bar], $\mathcal{L}_{pX(\text{PbX})} = 10(10^{-3}) \mu\text{b}^{-1} \text{ s}^{-1}$
- Provided that the runs can last as long,
similar luminosities for pA than with the extracted beam options (up to $60 \mu\text{b}^{-1} \text{ s}^{-1}$)

Luminosities with the internal-gas-target option

- Instantaneous Luminosity: $\mathcal{L} = \Phi_{beam} \times N_{target} = N_{beam} \times (\rho \times \ell \times \mathcal{N}_A) / A$
 - $\Phi_{p^+} = 3.2 \times 10^{14} p^+ \times 11000 \text{ Hz} = 3.5 \times 10^{18} p^+ \text{ s}^{-1}$ [1/2 Ampère !]
 - $\Phi_{Pb} = 4.2 \times 10^{10} p^+ \times 11000 \text{ Hz} = 4.6 \times 10^{14} \text{ Pb s}^{-1}$
 - Usable gas zone ℓ , up to 100 cm
 - Target density: $\frac{\rho}{P} = c = \frac{A}{22400} \text{ bar}^{-1} \text{ g cm}^{-3} \Rightarrow \mathcal{L} = \Phi_{beam} \times \left(\frac{\mathcal{N}_A}{22400} \times P \times \ell \right)$
[1 mole of a perfect gas occupies 22 400 cm³ at 273 K and 1 bar]
 - For $P = 10^{-9} \text{ bar}$ [7x that of SMOG in 2015, the 'vacuum' is 10^{-12} bar], $\mathcal{L}_{pX(\text{PbX})} = 10(10^{-3}) \mu\text{b}^{-1} \text{ s}^{-1}$
 - Provided that the runs can last as long,
similar luminosities for pA than with the extracted beam options (up to $60 \mu\text{b}^{-1} \text{ s}^{-1}$)
 - To get $10 \text{ fb}^{-1} \text{ y}^{-1}$ for pp , P should reach $10^{-7} \text{ bar} \leftrightarrow$ target storage cell
which could be polarised
- C. Barschel, P. Lenisa, A. Nass, and E. Steffens, Adv.Hi.En.Phys. (2015) ID:463141; See Ehrard's talk next
- Simply scaled up, this would give for Pbp or PbA $100 \text{ nb}^{-1} \text{ y}^{-1}$.
 \Rightarrow For PbA , limitations would come first from the beam lifetime.

Luminosities with the internal-gas-target option

- Instantaneous Luminosity: $\mathcal{L} = \Phi_{beam} \times N_{target} = N_{beam} \times (\rho \times \ell \times \mathcal{N}_A) / A$
 - $\Phi_{p^+} = 3.2 \times 10^{14} p^+ \times 11000 \text{ Hz} = 3.5 \times 10^{18} p^+ \text{ s}^{-1}$ [1/2 Ampère !]
 - $\Phi_{Pb} = 4.2 \times 10^{10} p^+ \times 11000 \text{ Hz} = 4.6 \times 10^{14} \text{ Pb s}^{-1}$
 - Usable gas zone ℓ , up to 100 cm
 - Target density: $\frac{\rho}{P} = c = \frac{A}{22400} \text{ bar}^{-1} \text{ g cm}^{-3} \Rightarrow \mathcal{L} = \Phi_{beam} \times \left(\frac{\mathcal{N}_A}{22400} \times P \times \ell \right)$
[1 mole of a perfect gas occupies 22 400 cm³ at 273 K and 1 bar]
 - For $P = 10^{-9} \text{ bar}$ [7x that of SMOG in 2015, the 'vacuum' is 10^{-12} bar], $\mathcal{L}_{pX(\text{PbX})} = 10(10^{-3}) \mu\text{b}^{-1} \text{ s}^{-1}$
 - Provided that the runs can last as long,
similar luminosities for pA than with the extracted beam options (up to $60 \mu\text{b}^{-1} \text{ s}^{-1}$)
 - To get $10 \text{ fb}^{-1} \text{ y}^{-1}$ for pp , P should reach $10^{-7} \text{ bar} \leftrightarrow$ target storage cell
which could be polarised
- C. Barschel, P. Lenisa, A. Nass, and E. Steffens, Adv.Hi.En.Phys. (2015) ID:463141; See Ehrard's talk next
- Simply scaled up, this would give for Pbp or PbA $100 \text{ nb}^{-1} \text{ y}^{-1}$.
 \Rightarrow For PbA , limitations would come first from the beam lifetime.
 - A specific gas target could be a competitive alternative to the beam extraction

Part IV

AFTER@LHC: the case of spin physics

The quest for the orbital angular momentum of the quarks and gluons

The quest for the orbital angular momentum of the quarks and gluons

- Quark/Gluon Sivers function: **distortion** in the distribution of an unpolarised partons with momentum fraction x and transverse momentum k_{\perp} **due to the proton transverse polarisation** : $f_{1T}^{\perp}(x, \vec{k}_{\perp}^2)$

The quest for the orbital angular momentum of the quarks and gluons

- Quark/Gluon Sivers function: **distortion** in the distribution of an unpolarised partons with momentum fraction x and transverse momentum k_{\perp} **due to the proton transverse polarisation** : $f_{1T}^{\perp}(x, \vec{k}_{\perp}^2)$
- First suggested by D. Sivers to explain the **large** observed left-right **single transverse spin asymmetries** A_N in $p^{\uparrow}p \rightarrow \pi X$

The quest for the orbital angular momentum of the quarks and gluons

- Quark/Gluon Sivers function: **distortion** in the distribution of an unpolarised partons with momentum fraction x and transverse momentum k_{\perp} **due to the proton transverse polarisation** : $f_{1T}^{\perp}(x, \vec{k}_{\perp}^2)$
- First suggested by D. Sivers to explain the **large** observed left-right **single transverse spin asymmetries** A_N in $p^{\uparrow}p \rightarrow \pi X$
- **non-zero** quark/gluon **Sivers function** \Rightarrow **non-zero** quark/gluon **OAM**

The quest for the orbital angular momentum of the quarks and gluons

- Quark/Gluon Sivers function: **distortion** in the distribution of an unpolarised partons with momentum fraction x and transverse momentum k_{\perp} **due to the proton transverse polarisation** : $f_{1T}^{\perp}(x, \vec{k}_{\perp}^2)$
- First suggested by D. Sivers to explain the **large** observed left-right **single transverse spin asymmetries** A_N in $p^{\uparrow}p \rightarrow \pi X$
- **non-zero** quark/gluon **Sivers function** \Rightarrow **non-zero** quark/gluon **OAM**
- Process dependence predicted: $f_{1T}^{\perp q}(x, \vec{k}_{\perp}^2)_{Drell-Yan} = -f_{1T}^{\perp q}(x, \vec{k}_{\perp}^2)_{Semi-Inclusive DIS}$

The quest for the orbital angular momentum of the quarks and gluons

- Quark/Gluon Sivers function: **distortion** in the distribution of an unpolarised partons with momentum fraction x and transverse momentum k_{\perp} **due to the proton transverse polarisation** : $f_{1T}^{\perp}(x, \vec{k}_{\perp}^2)$
- First suggested by D. Sivers to explain the **large** observed left-right **single transverse spin asymmetries** A_N in $p^{\uparrow}p \rightarrow \pi X$
- **non-zero quark/gluon Sivers function** \Rightarrow **non-zero quark/gluon OAM**
- Process dependence predicted: $f_{1T}^{\perp q}(x, \vec{k}_{\perp}^2)_{Drell-Yan} = -f_{1T}^{\perp q}(x, \vec{k}_{\perp}^2)_{Semi-Inclusive DIS}$
- Several experiments wish to measure $A_N^{Drell-Yan}$ to extract $f_{1T}^{\perp q}(x, \vec{k}_{\perp}^2)$
 - COMPASS: **valence quarks** using a pion beam (160 GeV)
on a polarised proton target
 - P1027: **valence quarks** using a polarised proton beam (120 GeV)
on an unpolarised proton target
 - P1039: **sea quarks** using an unpolarised proton beam (120 GeV)
on a polarised proton target

SSA in Drell-Yan studies with AFTER@LHC

→ Relevant parameters for existing and **proposed polarized DY experiments.**

S.J. Brodsky, F. Fleuret, C. Hadjidakis, JPL, Phys. Rep. 522 (2013) 239

V. Barone, F. Bradamante, A. Martin, Prog. Part. Nucl. Phys. 65 (2010) 267.

Experiment	particles	energy (GeV)	\sqrt{s} (GeV)	x_p^\uparrow	\mathcal{L} (nb ⁻¹ s ⁻¹)
AFTER	$p + p^\uparrow$	7000	115	0.01 ÷ 0.9	1
COMPASS	$\pi^\pm + p^\uparrow$	160	17.4	0.2 ÷ 0.3	2
COMPASS (low mass)	$\pi^\pm + p^\uparrow$	160	17.4	~ 0.05	2
P1039	$p + p^\uparrow$	120	15	0.1 ÷ 0.3	400-1000
P1027	$p^\uparrow + p$	120	15	0.35 ÷ 0.85	400-1000
RHIC	$p^\uparrow + p$	collider	500	0.05 ÷ 0.1	0.2
J-PARC	$p^\uparrow + p$	50	10	0.5 ÷ 0.9	1000
PANDA (low mass)	$\bar{p} + p^\uparrow$	15	5.5	0.2 ÷ 0.4	0.2
PAX	$p^\uparrow + \bar{p}$	collider	14	0.1 ÷ 0.9	0.002
NICA	$p^\uparrow + p$	collider	20	0.1 ÷ 0.8	0.001
RHIC Int. Target (1,2)	$p^\uparrow + p$	250	22	0.2 ÷ 0.5	(2,60)

SSA in Drell-Yan studies with AFTER@LHC

→ Relevant parameters for existing and **proposed polarized DY experiments**.

S.J. Brodsky, F. Fleuret, C. Hadjidakis, JPL, Phys. Rep. 522 (2013) 239

V. Barone, F. Bradamante, A. Martin, Prog. Part. Nucl. Phys. 65 (2010) 267.

Experiment	particles	energy (GeV)	\sqrt{s} (GeV)	x_p^\uparrow	\mathcal{L} (nb ⁻¹ s ⁻¹)
AFTER	$p + p^\uparrow$	7000	115	0.01 ÷ 0.9	1
COMPASS	$\pi^\pm + p^\uparrow$	160	17.4	0.2 ÷ 0.3	2
COMPASS (low mass)	$\pi^\pm + p^\uparrow$	160	17.4	~ 0.05	2
P1039	$p + p^\uparrow$	120	15	0.1 ÷ 0.3	400-1000
P1027	$p^\uparrow + p$	120	15	0.35 ÷ 0.85	400-1000
RHIC	$p^\uparrow + p$	collider	500	0.05 ÷ 0.1	0.2
J-PARC	$p^\uparrow + p$	50	10	0.5 ÷ 0.9	1000
PANDA (low mass)	$\bar{p} + p^\uparrow$	15	5.5	0.2 ÷ 0.4	0.2
PAX	$p^\uparrow + \bar{p}$	collider	14	0.1 ÷ 0.9	0.002
NICA	$p^\uparrow + p$	collider	20	0.1 ÷ 0.8	0.001
RHIC Int. Target (1,2)	$p^\uparrow + p$	250	22	0.2 ÷ 0.5	(2,60)

- For AFTER, \mathcal{L} corresponds to the Barschel *et al.* setup
or an equivalent of 50 cm liquid H target \Rightarrow could yield up to 10 fb⁻¹ per year

SSA in Drell-Yan studies with AFTER@LHC

→ Relevant parameters for existing and **proposed polarized DY experiments**.

S.J. Brodsky, F. Fleuret, C. Hadjidakis, JPL, Phys. Rep. 522 (2013) 239

V. Barone, F. Bradamante, A. Martin, Prog. Part. Nucl. Phys. 65 (2010) 267.

Experiment	particles	energy (GeV)	\sqrt{s} (GeV)	x_p^\uparrow	\mathcal{L} (nb ⁻¹ s ⁻¹)
AFTER	$p + p^\uparrow$	7000	115	0.01 ÷ 0.9	1
COMPASS	$\pi^\pm + p^\uparrow$	160	17.4	0.2 ÷ 0.3	2
COMPASS (low mass)	$\pi^\pm + p^\uparrow$	160	17.4	~ 0.05	2
P1039	$p + p^\uparrow$	120	15	0.1 ÷ 0.3	400-1000
P1027	$p^\uparrow + p$	120	15	0.35 ÷ 0.85	400-1000
RHIC	$p^\uparrow + p$	collider	500	0.05 ÷ 0.1	0.2
J-PARC	$p^\uparrow + p$	50	10	0.5 ÷ 0.9	1000
PANDA (low mass)	$\bar{p} + p^\uparrow$	15	5.5	0.2 ÷ 0.4	0.2
PAX	$p^\uparrow + \bar{p}$	collider	14	0.1 ÷ 0.9	0.002
NICA	$p^\uparrow + p$	collider	20	0.1 ÷ 0.8	0.001
RHIC Int. Target (1,2)	$p^\uparrow + p$	250	22	0.2 ÷ 0.5	(2,60)

- For AFTER, \mathcal{L} corresponds to the Barschel *et al.* setup
or an equivalent of 50 cm liquid H target \Rightarrow could yield up to 10 fb⁻¹ per year
- It is admittedly an apple-to-pear comparison since the precision on A_N
depends on the polarisation of the target/beam and on the cross-sections.

SSA in Drell-Yan studies with AFTER@LHC

→ Relevant parameters for existing and **proposed polarized DY experiments**.

S.J. Brodsky, F. Fleuret, C. Hadjidakis, JPL, Phys. Rep. 522 (2013) 239

V. Barone, F. Bradamante, A. Martin, Prog. Part. Nucl. Phys. 65 (2010) 267.

Experiment	particles	energy (GeV)	\sqrt{s} (GeV)	x_p^\dagger	\mathcal{L} (nb ⁻¹ s ⁻¹)
AFTER	$p + p^\dagger$	7000	115	0.01 ÷ 0.9	1
COMPASS	$\pi^\pm + p^\dagger$	160	17.4	0.2 ÷ 0.3	2
COMPASS (low mass)	$\pi^\pm + p^\dagger$	160	17.4	~ 0.05	2
P1039	$p + p^\dagger$	120	15	0.1 ÷ 0.3	400-1000
P1027	$p^\dagger + p$	120	15	0.35 ÷ 0.85	400-1000
RHIC	$p^\dagger + p$	collider	500	0.05 ÷ 0.1	0.2
J-PARC	$p^\dagger + p$	50	10	0.5 ÷ 0.9	1000
PANDA (low mass)	$\bar{p} + p^\dagger$	15	5.5	0.2 ÷ 0.4	0.2
PAX	$p^\dagger + \bar{p}$	collider	14	0.1 ÷ 0.9	0.002
NICA	$p^\dagger + p$	collider	20	0.1 ÷ 0.8	0.001
RHIC Int. Target (1,2)	$p^\dagger + p$	250	22	0.2 ÷ 0.5	(2,60)

- For AFTER, \mathcal{L} corresponds to the Barschel *et al.* setup
or an equivalent of 50 cm liquid H target \Rightarrow could yield up to 10 fb⁻¹ per year
- It is admittedly an apple-to-pear comparison since the precision on A_N
depends on the polarisation of the target/beam and on the cross-sections.
- Nota: At RHIC energy, Drell-Yan studies are very delicate (see later)

[not yet done for unpolarised pp collisions]

SSA in Drell-Yan studies with AFTER@LHC

→ Relevant parameters for existing and **proposed polarized DY experiments.**

S.J. Brodsky, F. Fleuret, C. Hadjidakis, JPL, Phys. Rep. 522 (2013) 239

V. Barone, F. Bradamante, A. Martin, Prog. Part. Nucl. Phys. 65 (2010) 267.

Experiment	particles	energy (GeV)	\sqrt{s} (GeV)	x_p^\uparrow	\mathcal{L} (nb ⁻¹ s ⁻¹)
AFTER	$p + p^\uparrow$	7000	115	0.01 ÷ 0.9	1
COMPASS	$\pi^\pm + p^\uparrow$	160	17.4	0.2 ÷ 0.3	2
COMPASS (low mass)	$\pi^\pm + p^\uparrow$	160	17.4	~ 0.05	2
P1039	$p + p^\uparrow$	120	15	0.1 ÷ 0.3	400-1000
P1027	$p^\uparrow + p$	120	15	0.35 ÷ 0.85	400-1000
RHIC	$p^\uparrow + p$	collider	500	0.05 ÷ 0.1	0.2
J-PARC	$p^\uparrow + p$	50	10	0.5 ÷ 0.9	1000
PANDA (low mass)	$\bar{p} + p^\uparrow$	15	5.5	0.2 ÷ 0.4	0.2
PAX	$p^\uparrow + \bar{p}$	collider	14	0.1 ÷ 0.9	0.002
NICA	$p^\uparrow + p$	collider	20	0.1 ÷ 0.8	0.001
RHIC Int. Target (1,2)	$p^\uparrow + p$	250	22	0.2 ÷ 0.5	(2,60)

- For AFTER, \mathcal{L} corresponds to the Barschel *et al.* setup
or an equivalent of 50 cm liquid H target \Rightarrow could yield up to 10 fb⁻¹ per year
- It is admittedly an apple-to-pear comparison since the precision on A_N
depends on the polarisation of the target/beam and on the cross-sections.
- Nota: At RHIC energy, Drell-Yan studies are very delicate (see later)

[not yet done for unpolarised pp collisions]

- **AFTER could be the only project able to reach $x^\uparrow = 10^{-2}$ and $x^\uparrow > 0.4$**

SSA in Drell-Yan studies with AFTER@LHC

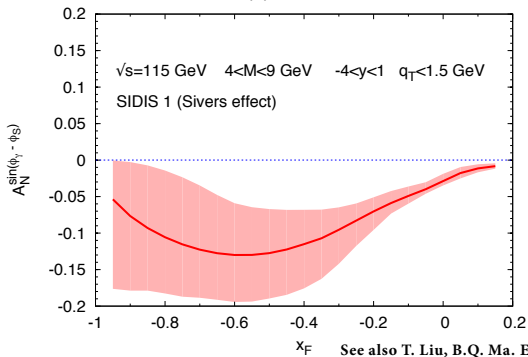
Expected asymmetries

The target-rapidity region (negative x_F) corresponds to **high x^\uparrow**
where the **k_T -spin correlation is the largest**

SSA in Drell-Yan studies with AFTER@LHC

Expected asymmetries

The target-rapidity region (negative x_F) corresponds to **high x^\uparrow**
where the **k_T -spin correlation is the largest**
 $p p^\uparrow \rightarrow l^+ \Gamma + X$

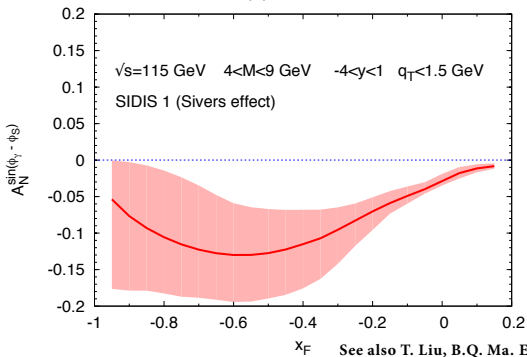


SSA in Drell-Yan studies with AFTER@LHC

Expected asymmetries

The target-rapidity region (negative x_F) corresponds to **high x^\uparrow**

where the **k_T -spin correlation is the largest**
 $p p^\uparrow \rightarrow l^+ \Gamma + X$



Experimental goal: to measure asymmetries on the order of 5-10 % at $x_F < 0$

With 10 fb^{-1} , one can expect up to 10^6 DY events in $4 < M < 9 \text{ GeV}$ (see later)

The gluon OAM contribution to the proton spin



- **Gluon Sivers effect** essentially **unconstrained**

D. Boer, C. Lorcé, C. Pisano, J. Zhou. *Adv.Hi.En.Phys.* (2015) ID:371396

The gluon OAM contribution to the proton spin



- **Gluon Sivers effect** essentially **unconstrained**

D. Boer, C. Lorcé, C. Pisano, J. Zhou. *Adv.Hi.En.Phys.* (2015) ID:371396

- It can be measured via A_N of **gluon sensitive probes** [as opposed to DY for quarks]



The gluon OAM contribution to the proton spin

- **Gluon Sivers effect** essentially **unconstrained**

D. Boer, C. Lorcé, C. Pisano, J. Zhou. *Adv.Hi.En.Phys.* (2015) ID:371396

- It can be measured via A_N of **gluon sensitive probes** [as opposed to DY for quarks]
- Theoretical complications suggest to analyse **multiple probes**



The gluon OAM contribution to the proton spin

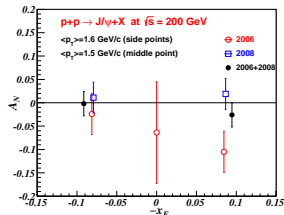
- **Gluon Sivers effect** essentially **unconstrained**

D. Boer, C. Lorcé, C. Pisano, J. Zhou. Adv.Hi.En.Phys. (2015) ID:371396

- It can be measured via A_N of **gluon sensitive probes** [as opposed to DY for quarks]
- Theoretical complications suggest to analyse **multiple probes**
- quarkonia (J/ψ , Υ , χ_C , η_C , ...)

F. Yuan, PRD 78 (2008) 014024; A. Schaefer, J. Zhou, PRD (2013)

PHENIX Phys.Rev. D86 (2012) 099904





The gluon OAM contribution to the proton spin

- **Gluon Sivers effect** essentially **unconstrained**

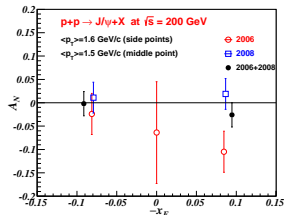
D. Boer, C. Lorcé, C. Pisano, J. Zhou. Adv.Hi.En.Phys. (2015) ID:371396

- It can be measured via A_N of **gluon sensitive probes** [as opposed to DY for quarks]
- Theoretical complications suggest to analyse **multiple probes**
- quarkonia (J/ψ , Υ , χ_C , η_C , ...)

F. Yuan, PRD 78 (2008) 014024; A. Schaefer, J. Zhou, PRD (2013)
PHENIX Phys.Rev. D86 (2012) 099904

- B & D meson production

M. Anselmino, *et al.* PRD 70 (2004) 074025.





The gluon OAM contribution to the proton spin

- **Gluon Sivers effect** essentially **unconstrained**

D. Boer, C. Lorcé, C. Pisano, J. Zhou. *Adv.Hi.En.Phys.* (2015) ID:371396

- It can be measured via A_N of **gluon sensitive probes** [as opposed to DY for quarks]
- Theoretical complications suggest to analyse **multiple probes**

- quarkonia (J/ψ , Υ , χ_C , η_C , ...)

F. Yuan, *PRD* 78 (2008) 014024; A. Schaefer, J. Zhou, *PRD* (2013)

PHENIX *Phys.Rev.* D86 (2012) 099904

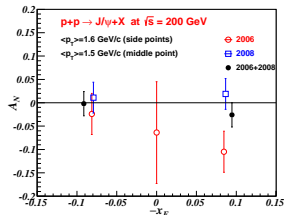
- B & D meson production

M. Anselmino, *et al.* *PRD* 70 (2004) 074025.

- γ , γ -jet, $\gamma - \gamma$

A. Bacchetta, *et al.*, *PRL* 99 (2007) 212002

J.W. Qiu, *et al.*, *PRL* 107 (2011) 062001





The gluon OAM contribution to the proton spin

- **Gluon Sivers effect** essentially **unconstrained**

D. Boer, C. Lorcé, C. Pisano, J. Zhou. *Adv.Hi.En.Phys.* (2015) ID:371396

- It can be measured via A_N of **gluon sensitive probes** [as opposed to DY for quarks]
- Theoretical complications suggest to analyse **multiple probes**

- quarkonia (J/ψ , Υ , χ_C , η_C , ...)

F. Yuan, *PRD* 78 (2008) 014024; A. Schaefer, J. Zhou, *PRD* (2013)

PHENIX *Phys.Rev.* D86 (2012) 099904

- B & D meson production

M. Anselmino, *et al.* *PRD* 70 (2004) 074025.

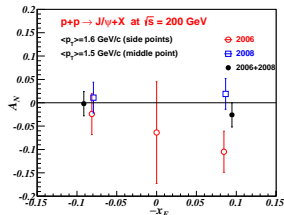
- γ , γ -jet, $\gamma - \gamma$

A. Bacchetta, *et al.*, *PRL* 99 (2007) 212002

J.W. Qiu, *et al.*, *PRL* 107 (2011) 062001

- $J/\psi + \gamma$: the cleanest; sensitive to gluons up to $x^\uparrow \simeq 0.5$

W. den Dunnen, J.P.L., C. Pisano, M. Schlegel, *PRL* 112, 212001 (2014); J.P.L., C. Pisano, M. Schlegel (work in progress)





The gluon OAM contribution to the proton spin

- **Gluon Sivers effect** essentially **unconstrained**

D. Boer, C. Lorcé, C. Pisano, J. Zhou. Adv.Hi.En.Phys. (2015) ID:371396

- It can be measured via A_N of **gluon sensitive probes** [as opposed to DY for quarks]
- Theoretical complications suggest to analyse **multiple probes**

- quarkonia (J/ψ , Υ , χ_C , η_C , ...)

F. Yuan, PRD 78 (2008) 014024; A. Schaefer, J. Zhou, PRD (2013)

PHENIX Phys.Rev. D86 (2012) 099904

- B & D meson production

M. Anselmino, *et al.* PRD 70 (2004) 074025.

- γ , γ -jet, $\gamma - \gamma$

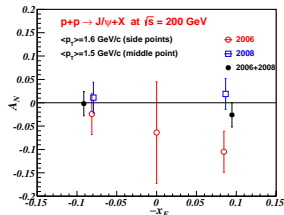
A. Bacchetta, *et al.*, PRL 99 (2007) 212002

J.W. Qiu, *et al.*, PRL 107 (2011) 062001

- $J/\psi + \gamma$: the cleanest; sensitive to gluons up to $x^\uparrow \simeq 0.5$

W. den Dunnen, J.P.L., C. Pisano, M. Schlegel, PRL 112, 212001 (2014); J.P.L., C. Pisano, M. Schlegel (work in progress)

- **All these measurements can be done with AFTER@LHC with the required precision: $10^9 J/\psi$, $10^6 \Upsilon$, $10^8 B$, etc ...**



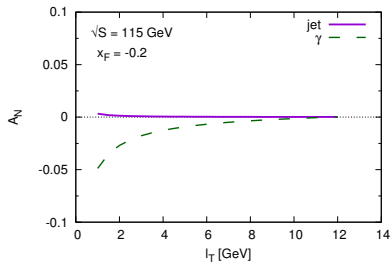
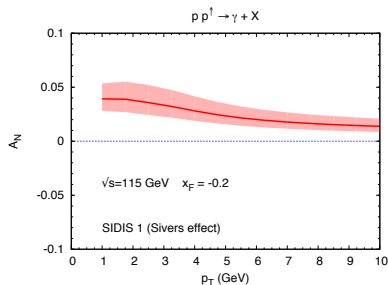


Further studies of the Sivers effect

- A_N^γ is predicted to have an **opposite sign** between the Generalised Parton Model (GPM) and the Collinear-Twist 3 (CT3) approach

GPM: M. Anselmino, U. D'Alesio, S. Melis. Adv.Hi.En.Phys. (2015) ID:475040

CT3: K. Kanazawa, Y. Koike, A. Metz, and D. Pitonyak. Adv.Hi.En.Phys. (2015) ID:257934.



- A_N^π : sign mismatch issue with $f_{1T}^{\perp,q}(x, \vec{k}_\perp^2)$ extracted from SIDIS
 - A_N^{jet} : complementary since no “contamination” (fragmentation Collins effect)
 - A_N^π should be measured at larger p_T

Part V

First simulation results

First simulation: is the boost an issue ?

B. Trzeciak, L. Massacrier *et al.*, 1504.05145 [hep-ex], to appear in Adv.Hi.En.Phys.

First simulation: is the boost an issue ?

B. Trzeciak, L. Massacrier *et al.*, 1504.05145 [hep-ex], to appear in Adv.Hi.En.Phys.

- LHCb has successfully carried out $p\text{Pb}$ and $\text{Pb}p$ analyses at 5 TeV

First simulation: is the boost an issue ?

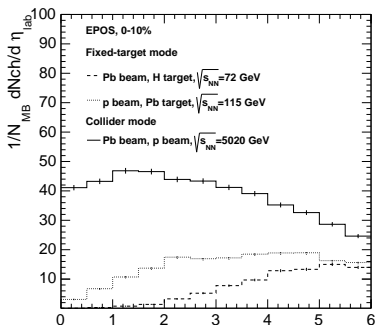
B. Trzeciak, L. Massacrier *et al.*, 1504.05145 [hep-ex], to appear in *Adv.Hi.En.Phys.*

- LHCb has successfully carried out p Pb and Pb p analyses at 5 TeV
- We have compared the multiplicity as function of η in the collider mode ($\sqrt{s} = 5$ TeV) vs. that in fixed target mode ($\sqrt{s} = 115$ TeV) using EPOS

First simulation: is the boost an issue ?

B. Trzeciak, L. Massacrier *et al.*, 1504.05145 [hep-ex], to appear in Adv.Hi.En.Phys.

- LHCb has successfully carried out p Pb and Pb p analyses at 5 TeV
- We have compared the multiplicity as function of η in the collider mode ($\sqrt{s} = 5$ TeV) vs. that in fixed target mode ($\sqrt{s} = 115$ TeV) using EPOS

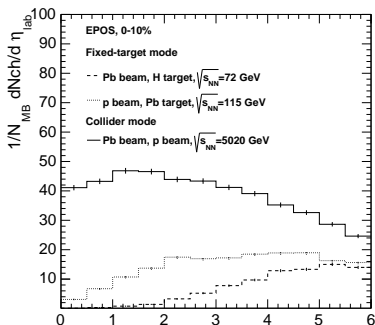


- Despite the boost, the multiplicity in the LHCb acceptance [forward η] is **lower** in the fixed mode than in the collider mode (at higher \sqrt{s})

First simulation: is the boost an issue ?

B. Trzeciak, L. Massacrier *et al.*, 1504.05145 [hep-ex], to appear in Adv.Hi.En.Phys.

- LHCb has successfully carried out p Pb and Pb p analyses at 5 TeV
- We have compared the multiplicity as function of η in the collider mode ($\sqrt{s} = 5$ TeV) vs. that in fixed target mode ($\sqrt{s} = 115$ TeV) using EPOS



- Despite the boost, the multiplicity in the LHCb acceptance [forward η] is **lower** in the fixed mode than in the collider mode (at higher \sqrt{s})
- Simulation backed-up with a comparison of the number-of-track distribution between **simulations at the detector level and data**

Z. Yang, private comm.

Fast simulation using LHCb reconstruction parameters

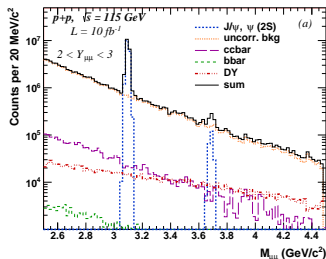
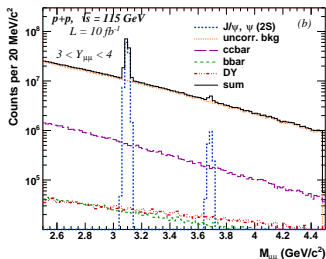
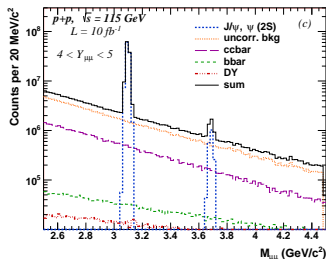
Projection for a LHCb-like detector

L. Massacrier, B. Trzeciak, *et al.*, Adv.Hi.En.Phys. (2015) ID:986348

- Simulations with Pythia 8.185
- the LHCb detector is NOT simulated but LHCb reconstruction parameters are introduced in the fast simulation (resolution, analysis cuts, efficiencies,...)
- Requirements:
 - Momentum resolution : $\Delta p/p = 0.5\%$
 - Muon identification efficiency: 98%
- Cuts at the single muon level
 - $2 < \eta_\mu < 5$
 - $p_{T\mu} > 0.7 \text{ GeV}$
- Muon misidentification:
 - If π and K decay before the calorimeters (12m), they are rejected by the tracking
 - otherwise a misidentification probability is applied following: F. Achilli et al, arXiv:1306.0249

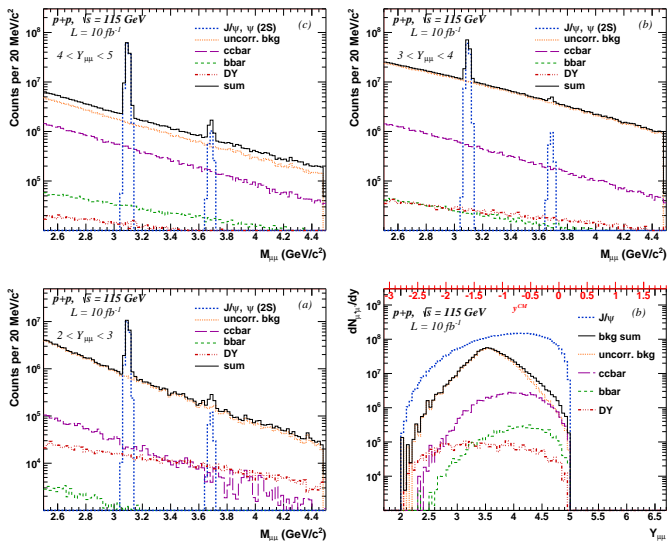
Charmonium background & its rapidity dependence

B. Trzeciak, L. Massacrier *et al.*, 1504.05145 [hep-ex], Adv.Hi.En.Phys. (2015) ID:986348



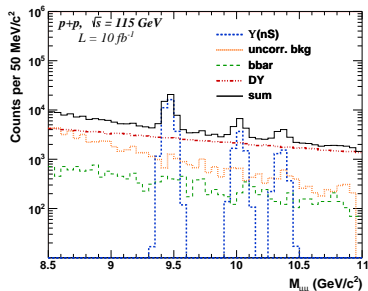
Charmonium background & its rapidity dependence

B. Trzeciak, L. Massacrier *et al.*, 1504.05145 [hep-ex], Adv.Hi.En.Phys. (2015) ID:986348



Bottomonium background & signal reach

B. Trzeciak, L. Massacrier *et al.*, 1504.05145 [hep-ex], Adv.Hi.En.Phys. (2015) ID:986348

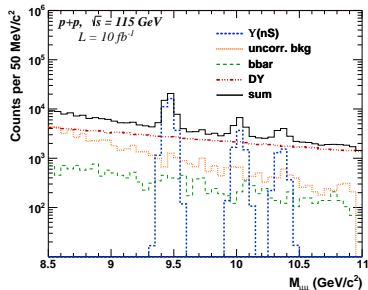


The dominant background is Drell-Yan

3 peaks well resolved

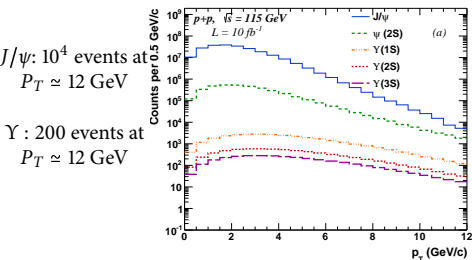
Bottomonium background & signal reach

B. Trzeciak, L. Massacrier *et al.*, 1504.05145 [hep-ex], Adv.Hi.En.Phys. (2015) ID:986348



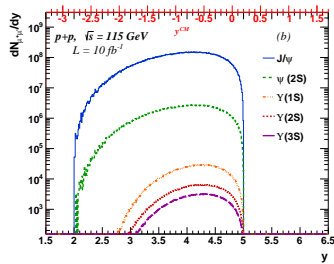
The dominant background is Drell-Yan

3 peaks well resolved



J/ψ : 10^4 events at
 $P_T \approx 12 \text{ GeV}$

Y: 200 events at
 $P_T \approx 12 \text{ GeV}$

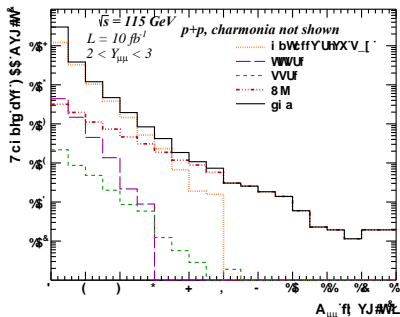


J/ψ : reach cut by the detector acceptance

Y: 200 events at
 $y_{c.m.s.}^Y \approx -2.1$, i.e.
 $x_2 \approx 0.7$

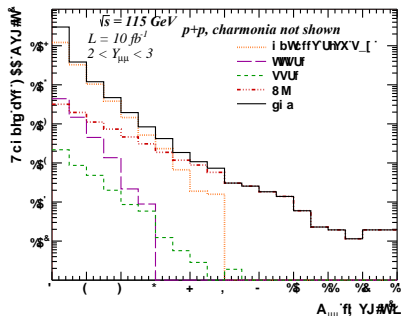
Drell-Yan background & signal reach

- At backward rapidities, quark-induced processes are favoured \Rightarrow Bkgd get smaller



Drell-Yan background & signal reach

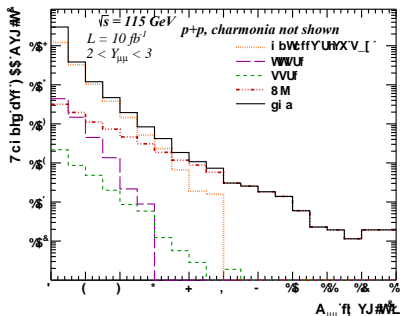
- At backward rapidities, quark-induced processes are favoured \Rightarrow Bkgd get smaller



- Charm and beauty background can be cut (2nd vertex) but interesting on their own

Drell-Yan background & signal reach

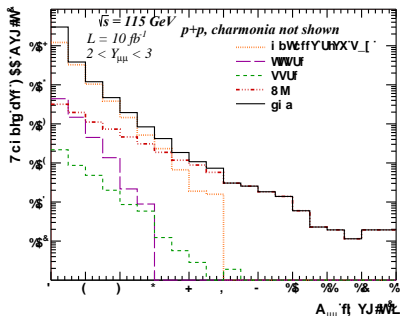
- At backward rapidities, quark-induced processes are favoured \Rightarrow Bkgd get smaller



- Charm and beauty background can be cut (2nd vertex) but interesting on their own
- Uncorrelated background can be subtracted by the mixing-event method
[up to which S/B depends on the systematics of the subtraction]

Drell-Yan background & signal reach

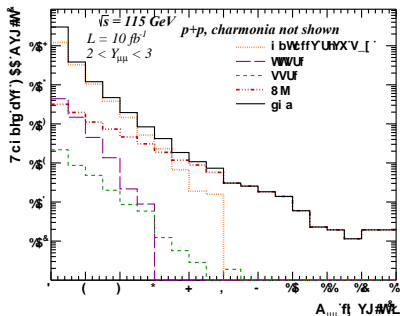
- At backward rapidities, quark-induced processes are favoured \Rightarrow Bkgd get smaller



- Charm and beauty background can be cut (2nd vertex) but interesting on their own
- Uncorrelated background can be subtracted by the mixing-event method
[up to which S/B depends on the systematics of the subtraction]
- Still **4000+ DY events** left in $2 < Y < 3$ for $8 < M < 9$ GeV, i.e. at $x^\uparrow \simeq 0.7$

Drell-Yan background & signal reach

- At backward rapidities, quark-induced processes are favoured \Rightarrow Bkgd get smaller



- Charm and beauty background can be cut (2nd vertex) but interesting on their own
- Uncorrelated background can be subtracted by the mixing-event method [up to which S/B depends on the systematics of the subtraction]
- Still **4000+ DY events** left in $2 < Y < 3$ for $8 < M < 9$ GeV, i.e. at $x^\dagger \simeq 0.7$
- Should yield to **precise measurements of A_N^{DY} at large x**

Part VI

Further readings

Further readings

Heavy-Ion Physics

- *Gluon shadowing effects on J/ψ and Υ production in $p+Pb$ collisions at $\sqrt{s_{NN}} = 115$ GeV and $Pb+p$ collisions at $\sqrt{s_{NN}} = 72$ GeV at AFTER@LHC* by R. Vogt. Adv.Hi.En.Phys. (2015) ID:492302.
- *Prospects for open heavy flavor measurements in heavy-ion and $p+A$ collisions in a fixed-target experiment at the LHC* by D. Kikola. Adv.Hi.En.Phys. (2015) ID:783134
- *Quarkonium suppression from coherent energy loss in fixed-target experiments using LHC beams* by F. Arleo, S.Peigné. [arXiv:1504.07428 [hep-ph]]. Adv.Hi.En.Phys. (2015) ID:961951
- *Anti-shadowing Effect on Charmonium Production at a Fixed-target Experiment Using LHC Beams* by K. Zhou, Z. Chen, P. Zhuang. arXiv:1507.05413 [nucl-th].
- *Lepton-pair production in ultraperipheral collisions at AFTER@LHC*
By J.P. Lansberg, L. Szymanowski, J. Wagner. arXiv:1504.02733 [hep-ph]. To appear in JHEP
- *Quarkonium Physics at a Fixed-Target Experiment using the LHC Beams.* By J.P. Lansberg, S.J. Brodsky, F. Fleuret, C. Hadjidakis. [arXiv:1204.5793 [hep-ph]]. Few Body Syst. 53 (2012) 11.

Further readings

Spin physics

- *Transverse single-spin asymmetries in proton-proton collisions at the AFTER@LHC experiment* by K. Kanazawa, Y. Koike, A. Metz, and D. Pitonyak. [arXiv:1502.04021 [hep-ph]]. Adv.Hi.En.Phys. (2015) ID:257934.
- *Transverse single-spin asymmetries in proton-proton collisions at the AFTER@LHC experiment in a TMD factorisation scheme* by M. Anselmino, U. D'Alesio, and S. Melis. [arXiv:1504.03791 [hep-ph]]. Adv.Hi.En.Phys. (2015) ID:475040.
- *The gluon Sivers distribution: status and future prospects* by D. Boer, C. Lorcé, C. Pisano, and J. Zhou. [arXiv:1504.04332 [hep-ph]]. Adv.Hi.En.Phys. (2015) ID:371396
- *Azimuthal asymmetries in lepton-pair production at a fixed-target experiment using the LHC beams (AFTER)* By T. Liu, B.Q. Ma. [arXiv:1203.5579 [hep-ph]]. Eur.Phys.J. C72 (2012) 2037.
- *Polarized gluon studies with charmonium and bottomonium at LHCb and AFTER* By D. Boer, C. Pisano. [arXiv:1208.3642 [hep-ph]]. Phys.Rev. D86 (2012) 094007.

Further readings

Hadron structure

- *Double-quarkonium production at a fixed-target experiment at the LHC (AFTER@LHC).*
by J.P. Lansberg, H.S. Shao. [arXiv:1504.06531 [hep-ph]]. To appear in Nucl. Phys. B
- *Next-To-Leading Order Differential Cross-Sections for Jpsi, psi(2S) and Upsilon Production in Proton-Proton Collisions at a Fixed-Target Experiment using the LHC Beams (AFTER@LHC)*
by Y. Feng, and J.X. Wang. Adv.Hi.En.Phys. (2015) ID:726393, in press.
- *η_c production in photon-induced interactions at a fixed target experiment at LHC as a probe of the odderon*
By V.P. Goncalves, W.K. Sauter. arXiv:1503.05112 [hep-ph].Phys.Rev. D91 (2015) 9, 094014.
- *A review of the intrinsic heavy quark content of the nucleon*
by S. J. Brodsky, A. Kusina, F. Lyonnet, I. Schienbein, H. Spiesberger, and R. Vogt. Adv.Hi.En.Phys. (2015) ID:231547, in press.
- *Hadronic production of Ξ_{cc} at a fixed-target experiment at the LHC*
By G. Chen *et al.* [arXiv:1401.6269 [hep-ph]]. Phys.Rev. D89 (2014) 074020.

Further readings

Feasibility study and technical ideas

- *Feasibility studies for quarkonium production at a fixed-target experiment using the LHC proton and lead beams (AFTER@LHC)* by L. Massacrier, B. Trzeciak, F. Fleuret, C. Hadjidakis, D. Kikola, J.P.Lansberg, and H.S. Shao arXiv:1504.05145 [hep-ex]. Adv.Hi.En.Phys. (2015) ID:986348
- *A Gas Target Internal to the LHC for the Study of pp Single-Spin Asymmetries and Heavy Ion Collisions* by C. Barschel, P. Lenisa, A. Nass, and E. Steffens. Adv.Hi.En.Phys. (2015) ID:463141
- *Quarkonium production and proposal of the new experiments on fixed target at LHC* by N.S. Topilskaya, and A.B. Kurepin. Adv.Hi.En.Phys. (2015) ID:760840

Generalities

- *Physics Opportunities of a Fixed-Target Experiment using the LHC Beams* By S.J. Brodsky, F. Fleuret, C. Hadjidakis, J.P. Lansberg. [arXiv:1202.6585 [hep-ph]]. Phys.Rept. 522 (2013) 239.

Part VII

Conclusion and outlooks

Conclusion

- **THREE MAIN THEMES PUSH FOR A FIXED-TARGET PROGRAM AT THE LHC**
[without interfering with the other experiments]

Conclusion

- **THREE MAIN THEMES PUSH FOR A FIXED-TARGET PROGRAM AT THE LHC**
[without interfering with the other experiments]
 - **The large x frontier:** new probes of the confinement
and connections with astroparticles

Conclusion

- **THREE MAIN THEMES PUSH FOR A FIXED-TARGET PROGRAM AT THE LHC**
[without interfering with the other experiments]
 - **The large x frontier:** new probes of the confinement
and connections with astroparticles
 - **The nucleon spin and the transverse dynamics of the partons**

Conclusion

- **THREE MAIN THEMES PUSH FOR A FIXED-TARGET PROGRAM AT THE LHC**
[without interfering with the other experiments]
 - **The large x frontier:** new probes of the confinement
and connections with astroparticles
 - **The nucleon spin and the transverse dynamics of the partons**
 - **The approach to the deconfinement phase transition:**
new energy, new rapidity domain and new probes

Conclusion

- **THREE MAIN THEMES PUSH FOR A FIXED-TARGET PROGRAM AT THE LHC**
[without interfering with the other experiments]
 - **The large x frontier:** new probes of the confinement
and connections with astroparticles
 - **The nucleon spin and the transverse dynamics of the partons**
 - **The approach to the deconfinement phase transition:**
new energy, new rapidity domain and new probes
- **2 WAYS TOWARDS FIXED-TARGET COLLISIONS WITH THE LHC BEAMS**

Conclusion

- **THREE MAIN THEMES PUSH FOR A FIXED-TARGET PROGRAM AT THE LHC**
[without interfering with the other experiments]
 - **The large x frontier:** new probes of the confinement
and connections with astroparticles
 - **The nucleon spin and the transverse dynamics of the partons**
 - **The approach to the deconfinement phase transition:**
new energy, new rapidity domain and new probes
- **2 WAYS TOWARDS FIXED-TARGET COLLISIONS WITH THE LHC BEAMS**
 - A slow extraction with a **bent crystal**
 - An internal **gas target** inspired from SMOG@LHCb/Hermes/H-jet@RHIC, ...

Conclusion

- **THREE MAIN THEMES PUSH FOR A FIXED-TARGET PROGRAM AT THE LHC**
[without interfering with the other experiments]
 - **The large x frontier:** new probes of the confinement
and connections with astroparticles
 - **The nucleon spin and the transverse dynamics of the partons**
 - **The approach to the deconfinement phase transition:**
new energy, new rapidity domain and new probes
- **2 WAYS TOWARDS FIXED-TARGET COLLISIONS WITH THE LHC BEAMS**
 - A slow extraction with a **bent crystal**
 - An internal **gas target** inspired from SMOG@LHCb/Hermes/H-jet@RHIC, ...
- **LARGE POTENTIAL FOR SPIN PHYSICS** with

Conclusion

- **THREE MAIN THEMES PUSH FOR A FIXED-TARGET PROGRAM AT THE LHC**
[without interfering with the other experiments]
 - **The large x frontier:** new probes of the confinement
and connections with astroparticles
 - **The nucleon spin and the transverse dynamics of the partons**
 - **The approach to the deconfinement phase transition:**
new energy, new rapidity domain and new probes
- **2 WAYS TOWARDS FIXED-TARGET COLLISIONS WITH THE LHC BEAMS**
 - A slow extraction with a **bent crystal**
 - An internal **gas target** inspired from SMOG@LHCb/Hermes/H-jet@RHIC, ...
- **LARGE POTENTIAL FOR SPIN PHYSICS** with
 - extracted beam + polarised target as COMPASS, P1039, JLab ...
 - polarised gas target

Conclusion

- **THREE MAIN THEMES PUSH FOR A FIXED-TARGET PROGRAM AT THE LHC**
[without interfering with the other experiments]
 - **The large x frontier:** new probes of the confinement
and connections with astroparticles
 - **The nucleon spin and the transverse dynamics of the partons**
 - **The approach to the deconfinement phase transition:**
new energy, new rapidity domain and new probes
- **2 WAYS TOWARDS FIXED-TARGET COLLISIONS WITH THE LHC BEAMS**
 - A slow extraction with a **bent crystal**
 - An internal **gas target** inspired from SMOG@LHCb/Hermes/H-jet@RHIC, ...
- **LARGE POTENTIAL FOR SPIN PHYSICS** with
 - extracted beam + polarised target as COMPASS, P1039, JLab ...
 - polarised gas target
- We have started to draft an Expression of Interest to be submitted to the LHCC

Conclusion

- **THREE MAIN THEMES PUSH FOR A FIXED-TARGET PROGRAM AT THE LHC**
[without interfering with the other experiments]
 - **The large x frontier:** new probes of the confinement
and connections with astroparticles
 - **The nucleon spin and the transverse dynamics of the partons**
 - **The approach to the deconfinement phase transition:**
new energy, new rapidity domain and new probes
- **2 WAYS TOWARDS FIXED-TARGET COLLISIONS WITH THE LHC BEAMS**
 - A slow extraction with a **bent crystal**
 - An internal **gas target** inspired from SMOG@LHCb/Hermes/H-jet@RHIC, ...
- **LARGE POTENTIAL FOR SPIN PHYSICS** with
 - extracted beam + polarised target as COMPASS, P1039, JLab ...
 - polarised gas target
- We have started to draft an Expression of Interest to be submitted to the LHCC
- Your contribution is welcome especially on the polarised target

Conclusion

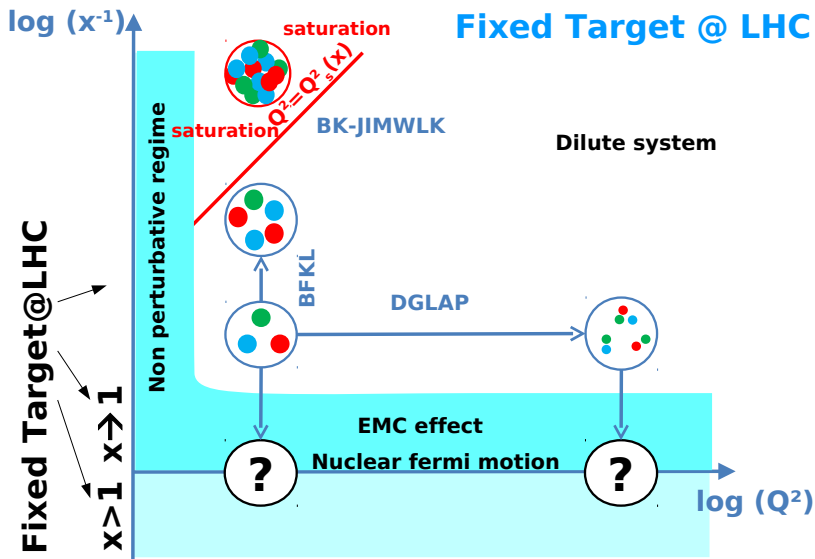
- **THREE MAIN THEMES PUSH FOR A FIXED-TARGET PROGRAM AT THE LHC**
[without interfering with the other experiments]
 - **The large x frontier:** new probes of the confinement
and connections with astroparticles
 - **The nucleon spin and the transverse dynamics of the partons**
 - **The approach to the deconfinement phase transition:**
new energy, new rapidity domain and new probes
- **2 WAYS TOWARDS FIXED-TARGET COLLISIONS WITH THE LHC BEAMS**
 - A slow extraction with a **bent crystal**
 - An internal **gas target** inspired from SMOG@LHCb/Hermes/H-jet@RHIC, ...
- **LARGE POTENTIAL FOR SPIN PHYSICS** with
 - extracted beam + polarised target as COMPASS, P1039, JLab ...
 - polarised gas target
- We have started to draft an Expression of Interest to be submitted to the LHCC
- Your contribution is welcome especially on the polarised target
- Webpage: <http://after.in2p3.fr>



Part VIII

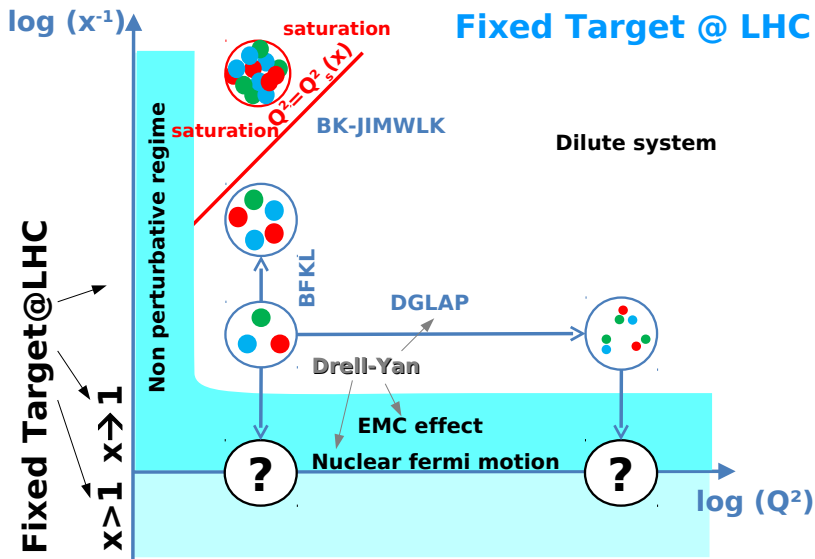
Backup slides

Overall



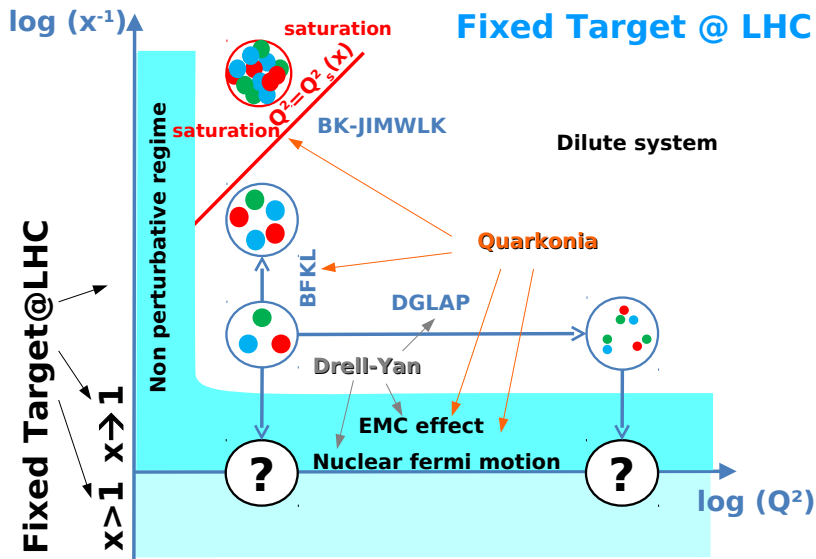
Overall

Fixed Target @ LHC



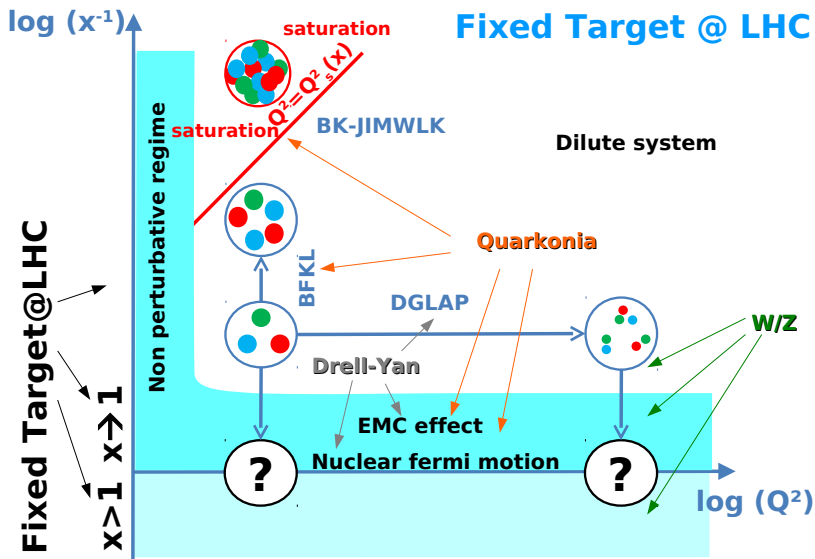
Overall

Fixed Target @ LHC



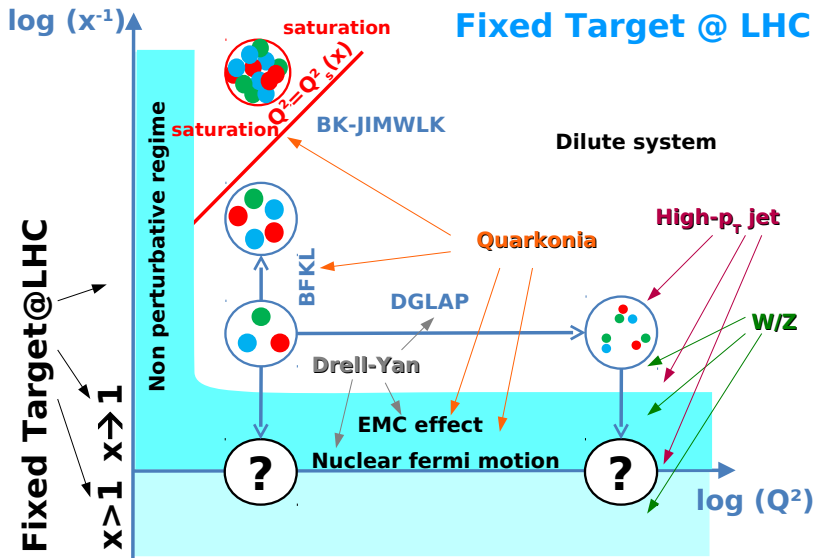
Overall

Fixed Target @ LHC



Overall

Fixed Target @ LHC



Gas target

C. Barschel, P. Lenisa, A. Nass, and E. Steffens, Adv.Hi.En.Phys. (2015) ID:463141

TABLE 1: Comparison of gas targets in storage rings with a hypothetical target for the proposed AFTER@LHC initiative [1, 2]. The target gas ^1H , ^2D , or ^3He is assumed to be spin polarized.

Storage ring	Particle	E_{max} [GeV]	Target type	L [m]	T [K]	L_{max} [1/cm ³ s]	Remarks	Reference
HERA-e DESY (term. 2007)	e^{\pm} pol.	27.6	Cell ^1H , ^2D , ^3He	0.4	100 25	$2.5 \cdot 10^{31}$ $2.5 \cdot 10^{32}$	HERMES exp. 1995–2007	[9]
RHIC-p BNL	p pol.	250	Jet	—	—	$1.7 \cdot 10^{30}$	Absolute p polarimeter	[10]
COSY FZ Jülich	p, d pol.	3.77 $T = 49.3$ MeV	Cell ^1H , ^2D Cell ^1H	0.4	300	10^{29} $2.75 \cdot 10^{29}$	ANKE exp. PAX exp.	[4, 5] [11]
LHC CERN (proposed)	p unpol. heavy ions	7,000 $2,760 \cdot A$	Cell ^1H , ^2D Xe $M \approx 131$	1.0	100 ≥ 100	10^{33} $10^{27} - 10^{28}$	Based on techn. of HERMES target	this paper

→ beam lifetime with $\mathcal{L}_{\text{pp}} = 10^{33} \text{ cm}^{-2} \text{ s}^{-1} = 10 \text{ nb}^{-1} \text{ s}^{-1}$ of 2×10^6 s (or 23 days).

Accessing the large x glue with quarkonia:

PYTHIA simulation
 $\sigma(y) / \sigma(y=0.4)$
 statistics for one month
 5% acceptance considered

Statistical relative uncertainty
 Large statistics allow to access
 very backward region

Gluon uncertainty from
 MSTWPDF
 - only for the gluon content of
 the target
 - assuming

$$x_g = M_{J/\psi} / \sqrt{s} e^{-y_{CM}}$$

J/ψ

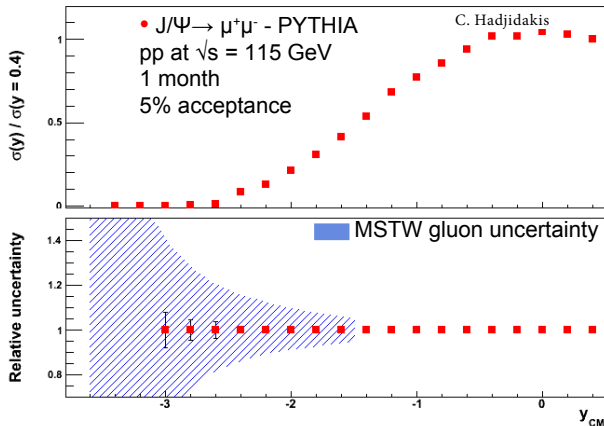
$y_{CM} \sim 0 \rightarrow x_g = 0.03$

$y_{CM} \sim -3.6 \rightarrow x_g = 1$

Y : larger x_g for same y_{CM}

$y_{CM} \sim 0 \rightarrow x_g = 0.08$

$y_{CM} \sim -2.4 \rightarrow x_g = 1$



⇒ Backward measurements allow to access large x gluon pdf

Assuming that we understand the
 quarkonium-production mechanisms

Distribution of linearly polarised gluons in unpolarised protons

Distribution of linearly polarised gluons in unpolarised protons

PHYSICAL REVIEW D **86**, 094007 (2012)

Polarized gluon studies with charmonium and bottomonium at LHCb and AFTER

Daniël Boer^{*}

Theory Group, KVI, University of Groningen, Zernikelaan 25, NL-9747 AA Groningen, The Netherlands

Cristian Pisano[†]

Istituto Nazionale di Fisica Nucleare, Sezione di Cagliari, C.P. 170, I-09042 Monserrato (CA), Italy

Distribution of linearly polarised gluons in unpolarised protons

PHYSICAL REVIEW D **86**, 094007 (2012)

Polarized gluon studies with charmonium and bottomonium at LHCb and AFTER

Daniël Boer*

Theory Group, KVI, University of Groningen, Zernikelaan 25, NL-9747 AA Groningen, The Netherlands

Cristian Pisano†

Istituto Nazionale di Fisica Nucleare, Sezione di Cagliari, C.P. 170, I-09042 Monserrato (CA), Italy

- **Low P_T C-even quarkonium** production is a good probe of the distribution of linearly polarised gluons in unpolarised protons: $h_1^{\perp g}$

Distribution of linearly polarised gluons in unpolarised protons

PHYSICAL REVIEW D 86, 094007 (2012)

Polarized gluon studies with charmonium and bottomonium at LHCb and AFTER

Daniël Boer*

Theory Group, KVI, University of Groningen, Zernikelaan 25, NL-9747 AA Groningen, The Netherlands

Cristian Pisano†

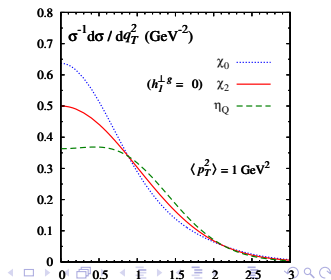
Istituto Nazionale di Fisica Nucleare, Sezione di Cagliari, C.P. 170, I-09042 Monserrato (CA), Italy

- **Low P_T C-even quarkonium** production is a good probe of the distribution of linearly polarised gluons in unpolarised protons: $h_1^{\perp g}$

- Affect the **low P_T spectra**:

$$\frac{1}{\sigma} \frac{d\sigma(\eta_Q)}{dq_T^2} \propto 1 - R(\mathbf{q}_T^2) \quad \& \quad \frac{1}{\sigma} \frac{d\sigma(\chi_{0,Q})}{dq_T^2} \propto 1 + R(\mathbf{q}_T^2)$$

(R involves $f_1^g(x, k_T, \mu)$ and $h_1^{\perp g}(x, k_T, \mu)$)



Distribution of linearly polarised gluons in unpolarised protons

PHYSICAL REVIEW D 86, 094007 (2012)

Polarized gluon studies with charmonium and bottomonium at LHCb and AFTER

Daniël Boer*

Theory Group, KVI, University of Groningen, Zernikelaan 25, NL-9747 AA Groningen, The Netherlands

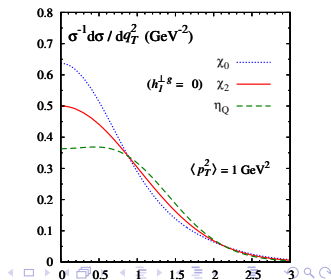
Cristian Pisano†

Istituto Nazionale di Fisica Nucleare, Sezione di Cagliari, C.P. 170, I-09042 Monserrato (CA), Italy

- **Low P_T C-even quarkonium** production is a good probe of the distribution of linearly polarised gluons in unpolarised protons: $h_1^{\perp g}$
- Affect the **low P_T spectra**:

$$\frac{1}{\sigma} \frac{d\sigma(\eta_Q)}{dq_T^2} \propto 1 - R(\mathbf{q}_T^2) \quad \& \quad \frac{1}{\sigma} \frac{d\sigma(\chi_{0,Q})}{dq_T^2} \propto 1 + R(\mathbf{q}_T^2)$$

(R involves $f_1^g(x, k_T, \mu)$ and $h_1^{\perp g}(x, k_T, \mu)$)
- The boost is of great help to access low P_T P -wave quarkonia



Distribution of linearly polarised gluons in unpolarised protons

PHYSICAL REVIEW D 86, 094007 (2012)

Polarized gluon studies with charmonium and bottomonium at LHCb and AFTER

Daniël Boer*

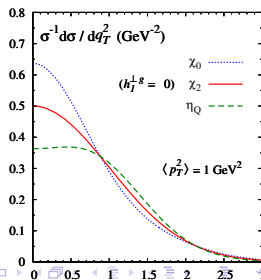
Theory Group, KVI, University of Groningen, Zernikelaan 25, NL-9747 AA Groningen, The Netherlands

Cristian Pisano†

Istituto Nazionale di Fisica Nucleare, Sezione di Cagliari, C.P. 170, I-09042 Monserrato (CA), Italy

- **Low P_T C-even quarkonium** production is a good probe of the distribution of linearly polarised gluons in unpolarised protons: $h_1^{\perp g}$
- Affect the **low P_T spectra**:

$$\frac{1}{\sigma} \frac{d\sigma(\eta_Q)}{dq_T^2} \propto 1 - R(\mathbf{q}_T^2) \quad \& \quad \frac{1}{\sigma} \frac{d\sigma(\chi_{0,Q})}{dq_T^2} \propto 1 + R(\mathbf{q}_T^2)$$
(R involves $f_1^g(x, k_T, \mu)$ and $h_1^{\perp g}(x, k_T, \mu)$)
- The boost is of great help to access low P_T P -wave quarkonia
- $h_1^{\perp g}$ is connected to the Higgs transverse-momentum distribution D. Boer, *et al.* PRL 108 (2012) 032002



Access to h_1^{1g} : II

Access to h_1^{1g} : II

PRL 112, 212001 (2014)

PHYSICAL REVIEW LETTERS

week ending
30 MAY 2014

**Accessing the Transverse Dynamics and Polarization of Gluons inside
the Proton at the LHC**

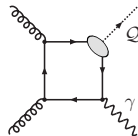
Wilco J. den Dunnen,^{1,*} Jean-Philippe Lansberg,^{2,†} Cristian Pisano,^{3,‡} and Marc Schlegel^{1,§}

¹*Institute for Theoretical Physics, Universität Tübingen, Auf der Morgenstelle 14, D-72076 Tübingen, Germany*

²*IPNO, Université Paris-Sud, CNRS/IN2P3, F-91406, Orsay, France*

³*Nikhef and Department of Physics and Astronomy, VU University Amsterdam,*

De Boelelaan 1081, NL-1081 HV Amsterdam, The Netherlands



Access to h_1^{1g} : II

PRL 112, 212001 (2014)

PHYSICAL REVIEW LETTERS

week ending
30 MAY 2014

Accessing the Transverse Dynamics and Polarization of Gluons inside the Proton at the LHC

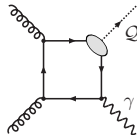
Wilco J. den Dunnen,^{1,*} Jean-Philippe Lansberg,^{2,†} Cristian Pisano,^{3,‡} and Marc Schlegel^{1,§}

¹*Institute for Theoretical Physics, Universität Tübingen, Auf der Morgenstelle 14, D-72076 Tübingen, Germany*

²*IPNO, Université Paris-Sud, CNRS/IN2P3, F-91406, Orsay, France*

³*Nikhef and Department of Physics and Astronomy, VU University Amsterdam,*

De Boelelaan 1081, NL-1081 HV Amsterdam, The Netherlands



- Gluon B-M can also be accessed via back-to-back $\psi/\Upsilon + \gamma$ associated production at the LHC. Also true at AFTER@LHC !

Access to $h_1^{\perp g}$: II

PRL 112, 212001 (2014)

PHYSICAL REVIEW LETTERS

week ending
30 MAY 2014

Accessing the Transverse Dynamics and Polarization of Gluons inside the Proton at the LHC

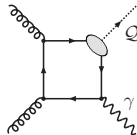
Wilco J. den Dunnen,^{1,*} Jean-Philippe Lansberg,^{2,†} Cristian Pisano,^{3,‡} and Marc Schlegel^{1,§}

¹Institute for Theoretical Physics, Universität Tübingen, Auf der Morgenstelle 14, D-72076 Tübingen, Germany

²IPNO, Université Paris-Sud, CNRS/IN2P3, F-91406, Orsay, France

³Nikhef and Department of Physics and Astronomy, VU University Amsterdam,

De Boelelaan 1081, NL-1081 HV Amsterdam, The Netherlands



- Gluon B-M can also be accessed via back-to-back $\psi/\Upsilon + \gamma$ associated production at the LHC. Also true at AFTER@LHC !
- Smaller yield (14 TeV \rightarrow 115 GeV) compensated by an access to lower P_T

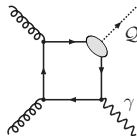
Access to $h_1^{\perp g}$: II

PRL 112, 212001 (2014)

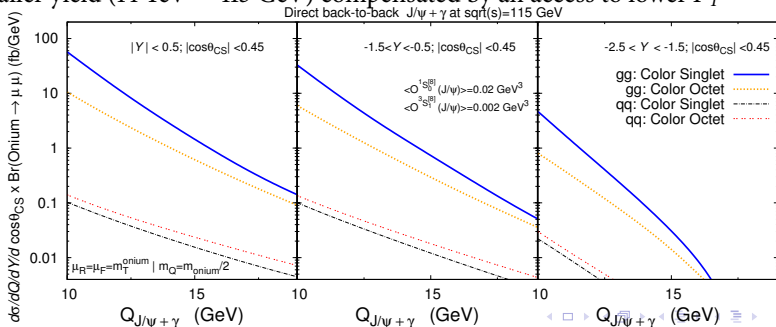
PHYSICAL REVIEW LETTERS

week ending
30 MAY 2014Accessing the Transverse Dynamics and Polarization of Gluons inside
the Proton at the LHCWilco J. den Dunnen,^{1,*} Jean-Philippe Lansberg,^{2,†} Cristian Pisano,^{3,‡} and Marc Schlegel^{1,§}¹Institute for Theoretical Physics, Universität Tübingen, Auf der Morgenstelle 14, D-72076 Tübingen, Germany²IPNO, Université Paris-Sud, CNRS/IN2P3, F-91406, Orsay, France³Nikhef and Department of Physics and Astronomy, VU University Amsterdam,

De Boelelaan 1081, NL-1081 HV Amsterdam, The Netherlands



- Gluon B-M can also be accessed via back-to-back $\psi/\Upsilon + \gamma$ associated production at the LHC. Also true at AFTER@LHC !
- Smaller yield (14 TeV \rightarrow 115 GeV) compensated by an access to lower P_T



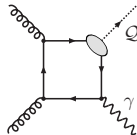
Access to h_1^{1g} : II

PRL 112, 212001 (2014)

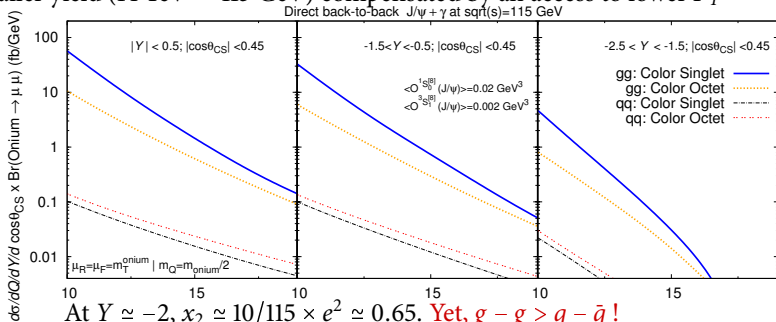
PHYSICAL REVIEW LETTERS

week ending
30 MAY 2014Accessing the Transverse Dynamics and Polarization of Gluons inside
the Proton at the LHCWilco J. den Dunnen,^{1,*} Jean-Philippe Lansberg,^{2,†} Cristian Pisano,^{3,‡} and Marc Schlegel^{1,§}¹Institute for Theoretical Physics, Universität Tübingen, Auf der Morgenstelle 14, D-72076 Tübingen, Germany²IPNO, Université Paris-Sud, CNRS/IN2P3, F-91406, Orsay, France³Nikhef and Department of Physics and Astronomy, VU University Amsterdam,

De Boelelaan 1081, NL-1081 HV Amsterdam, The Netherlands



- Gluon B-M can also be accessed via back-to-back $\psi/\Upsilon + \gamma$ associated production at the LHC. Also true at AFTER@LHC !
- Smaller yield (14 TeV \rightarrow 115 GeV) compensated by an access to lower P_T



Luminosities with extracted-lead beams

Luminosities with extracted-lead beams

- Instantaneous Luminosity:

$$\mathcal{L} = \Phi_{beam} \times N_{target} = N_{beam} \times (\rho \times \ell \times \mathcal{N}_A) / A$$

$$\Phi_{beam} = 2 \times 10^5 \text{ Pb s}^{-1}, \quad \ell = 1 \text{ cm (target thickness)}$$

Luminosities with extracted-lead beams

- Instantaneous Luminosity:

$$\mathcal{L} = \Phi_{beam} \times N_{target} = N_{beam} \times (\rho \times \ell \times \mathcal{N}_A) / A$$

$$\Phi_{beam} = 2 \times 10^5 \text{ Pb s}^{-1}, \quad \ell = 1 \text{ cm (target thickness)}$$

- Integrated luminosity $\int dt \mathcal{L} = \mathcal{L} \times 10^6 \text{ s}$ for Pb

Luminosities with extracted-lead beams

- Instantaneous Luminosity:

$$\mathcal{L} = \Phi_{beam} \times N_{target} = N_{beam} \times (\rho \times \ell \times \mathcal{N}_A) / A$$

$$\Phi_{beam} = 2 \times 10^5 \text{ Pb s}^{-1}, \quad \ell = 1 \text{ cm (target thickness)}$$

- Integrated luminosity $\int dt \mathcal{L} = \mathcal{L} \times 10^6 \text{ s}$ for Pb
- Expected luminosities with $2 \times 10^5 \text{ Pb s}^{-1}$ extracted (1cm-long target)

Target	ρ (g.cm ⁻³)	A	\mathcal{L} (mb ⁻¹ .s ⁻¹) = $\int \mathcal{L}$ (nb ⁻¹ .yr ⁻¹)
1m Liq. H₂	0.07	1	800
1m Liq. D₂	0.16	2	1000
1cm Be	1.85	9	25
1cm Cu	8.96	64	17
1cm W	19.1	185	13
1cm Pb	11.35	207	7

Luminosities with extracted-lead beams

- Instantaneous Luminosity:

$$\mathcal{L} = \Phi_{beam} \times N_{target} = N_{beam} \times (\rho \times \ell \times \mathcal{N}_A) / A$$

$$\Phi_{beam} = 2 \times 10^5 \text{ Pb s}^{-1}, \quad \ell = 1 \text{ cm (target thickness)}$$

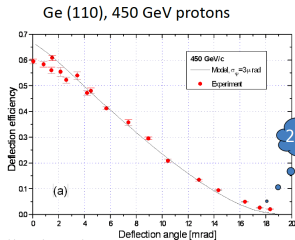
- Integrated luminosity $\int dt \mathcal{L} = \mathcal{L} \times 10^6 \text{ s}$ for Pb
- Expected luminosities with $2 \times 10^5 \text{ Pb s}^{-1}$ extracted (1cm-long target)

Target	ρ (g.cm ⁻³)	A	\mathcal{L} (mb ⁻¹ .s ⁻¹) = $\int \mathcal{L}$ (nb ⁻¹ .yr ⁻¹)
1m Liq. H₂	0.07	1	800
1m Liq. D₂	0.16	2	1000
1cm Be	1.85	9	25
1cm Cu	8.96	64	17
1cm W	19.1	185	13
1cm Pb	11.35	207	7

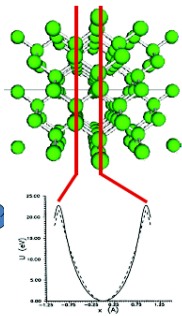
- Planned lumi for PHENIX Run15AuAu 2.8 nb^{-1} (0.13 nb^{-1} at 62 GeV)
- Nominal LHC lumi for PbPb 0.5 nb^{-1}

The beam extraction with a bent crystal

- Inter-crystalline fields are huge

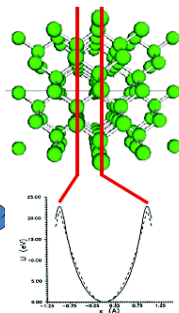
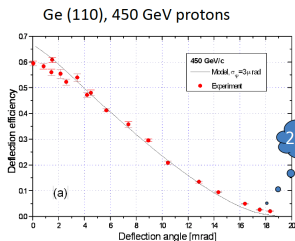


2000 T !



The beam extraction with a bent crystal

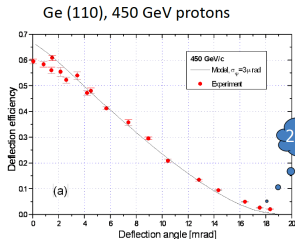
- Inter-crystalline fields are huge



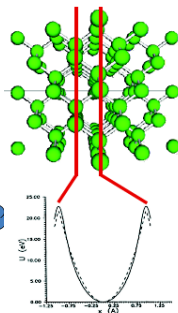
- The channeling efficiency is high for a deflection of a few mrad

The beam extraction with a bent crystal

- Inter-crystalline fields are huge



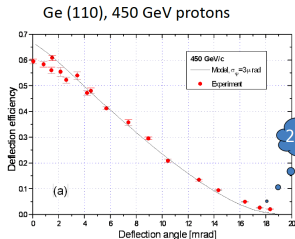
2000 T !



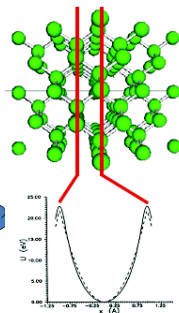
- The **channeling efficiency** is high for a deflection of a few mrad
- One can **extract** a significant part of the **beam loss** ($10^9 p^+ s^{-1}$)

The beam extraction with a bent crystal

- Inter-crystalline **fields are huge**



2000 T !



- The **channeling efficiency** is high for a deflection of a few mrad
- One can **extract** a significant part of the **beam loss** ($10^9 p^+ s^{-1}$)
- Simple and robust way to extract the most energetic beam ever:



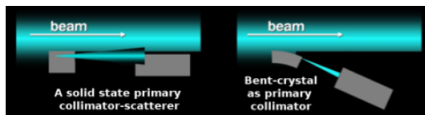
The beam extraction: news

[S. Montesano, *Physics at AFTER using LHC beams, ICT* Trento, Feb. 2013*]



UA9 installation in the SPS

Goal : assess the possibility to use bent crystals as primary collimators in hadronic accelerators and colliders



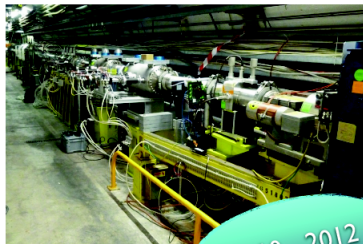
Prototype crystal collimation system at SPS :

- local beam loss reduction (5÷20x reduction for proton beam)
- beam loss map show average loss reduction in the entire SPS ring
- halo extraction efficiency
70÷80% for protons (50÷70% for Pb)

The beam extraction: news

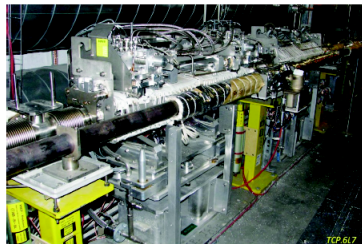
[S. Montesano, *Physics at AFTER using LHC beams, ECT* Trento, Feb. 2013*]

Goal : assess the possibility to use bent crystals as primary collimators in hadronic accelerators and colliders



UA9 installation in the SPS

2010 - 2012



LUA9 future installation in LHC

Prototype crystal collimation system at SPS :

- local beam loss reduction (5+20x reduction for proton beam)
- beam loss map show average loss reduction in the entire SPS ring
- halo extraction efficiency
70+80% for protons (50+70% for Pb)

The beam extraction: news

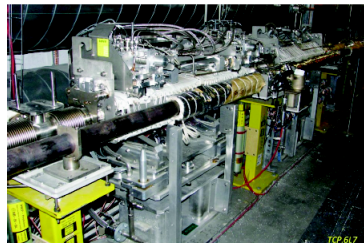
[S. Montesano, *Physics at AFTER using LHC beams*, TCT* Trento, Feb. 2013]

Goal : assess the possibility to use bent crystals as primary collimators in hadronic accelerators and colliders



UA9 installation in the SPS

2010 - 2012



LUA9 future installation in LHC

Prototype crystal collimation system at SPS :

- local beam loss reduction (5+20x reduction for proton beam)
- beam loss map show average loss reduction in the entire SPS ring
- halo extraction efficiency
70+80% for protons (50+70% for Pb)

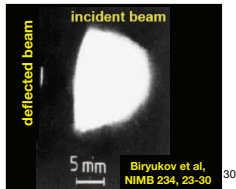
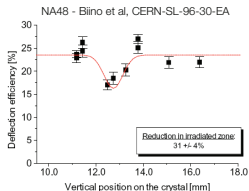
Towards an installation in the LHC : propose and install during LSI a min. number of devices

- 2 crystals

Long term plan is ambitious : propose a collimation system based on bent crystals for the upgrade of the current LHC collimation system

Crystal resistance to irradiation

- **IHEP U-70** (Biryukov et al, NIMB 234, 23-30):
 - 70 GeV protons, 50 ms spills of 10^{14} protons every 9.6 s, several minutes irradiation
 - equivalent to 2 nominal LHC bunches for 500 turns every 10 s
 - 5 mm silicon crystal, **channeling efficiency unchanged**
- **SPS North Area - NA48** (Biino et al, CERN-SL-96-30-EA):
 - 450 GeV protons, 2.4 s spill of 5×10^{12} protons every 14.4 s, one year irradiation, 2.4×10^{20} protons/cm² in total,
 - equivalent to several year of operation for a primary collimator in LHC
 - $10 \times 50 \times 0.9$ mm³ silicon crystal, 0.8×0.3 mm² area irradiated, **channeling efficiency reduced by 30%**.
- **HRMT16-UA9CRY** (HiRadMat facility, November 2012):
 - 440 GeV protons, up to 288 bunches in 7.2 μ s, 1.1×10^{11} protons per bunch (3×10^{13} protons in total)
 - energy deposition comparable to an asynchronous beam dump in LHC
 - 3 mm long silicon crystal, **no damage to the crystal after accurate visual inspection**, more tests planned to assess possible crystal lattice damage
 - **accurate FLUKA simulation of energy deposition** and residual dose



A few figures on the (extracted) proton beam

- Beam loss: $10^9 p^+ s^{-1}$
- Extracted intensity: $5 \times 10^8 p^+ s^{-1}$ (1/2 the beam loss)

E. Uggerhøj, U.I Uggerhøj, NIM B 234 (2005) 31

A few figures on the (extracted) proton beam

- Beam loss: $10^9 p^+ s^{-1}$
- Extracted intensity: $5 \times 10^8 p^+ s^{-1}$ (1/2 the beam loss)
- Number of p^+ : 2808 bunches of $1.15 \times 10^{11} p^+ = 3.2 \times 10^{14} p^+$

E. Uggerhøj, U.I Uggerhøj, NIM B 234 (2005) 31

A few figures on the (extracted) proton beam

- Beam loss: $10^9 p^+ s^{-1}$
- Extracted intensity: $5 \times 10^8 p^+ s^{-1}$ (1/2 the beam loss)
- Number of p^+ : 2808 bunches of $1.15 \times 10^{11} p^+ = 3.2 \times 10^{14} p^+$
- Revolution frequency: Each bunch passes the extraction point at a rate of $3.10^5 \text{ km}\cdot\text{s}^{-1} / 27 \text{ km} \simeq 11 \text{ kHz}$

E. Uggerhøj, U.I Uggerhøj, NIM B 234 (2005) 31

A few figures on the (extracted) proton beam

- Beam loss: $10^9 p^+ s^{-1}$
- Extracted intensity: $5 \times 10^8 p^+ s^{-1}$ (1/2 the beam loss)
- Number of p^+ : 2808 bunches of $1.15 \times 10^{11} p^+ = 3.2 \times 10^{14} p^+$
- Revolution frequency: Each bunch passes the extraction point at a rate of $3.10^5 \text{ km}\cdot\text{s}^{-1}/27 \text{ km} \simeq 11 \text{ kHz}$
- Extracted “mini” bunches:
 - the crystal sees $2808 \times 11000 \text{ s}^{-1} \simeq 3.10^7 \text{ bunches s}^{-1}$
 - one extracts $5.10^8/3.10^7 \simeq 15p^+$ from each bunch at each pass
 - Provided that the probability of interaction with the target is below 5%,

E. Uggerhøj, U.I Uggerhøj, NIM B 234 (2005) 31

pile-up is not an issue

A few figures on the (extracted) proton beam

- Beam loss: $10^9 p^+ s^{-1}$
- Extracted intensity: $5 \times 10^8 p^+ s^{-1}$ (1/2 the beam loss)
- Number of p^+ : 2808 bunches of $1.15 \times 10^{11} p^+ = 3.2 \times 10^{14} p^+$
- Revolution frequency: Each bunch passes the extraction point at a rate of $3.10^5 \text{ km} \cdot \text{s}^{-1} / 27 \text{ km} \simeq 11 \text{ kHz}$
- Extracted “mini” bunches:
 - the crystal sees $2808 \times 11000 \text{ s}^{-1} \simeq 3.10^7 \text{ bunches s}^{-1}$
 - one extracts $5.10^8 / 3.10^7 \simeq 15 p^+$ from each bunch at each pass
 - Provided that the probability of interaction with the target is below 5%,
- Extraction over a 10h fill:
 - $5 \times 10^8 p^+ \times 3600 \text{ s h}^{-1} \times 10 \text{ h} = 1.8 \times 10^{13} p^+ \text{ fill}^{-1}$
 - This means $1.8 \times 10^{13} / 3.2 \times 10^{14} \simeq 5.6\%$ of the p^+ in the beam

E. Uggerhøj, U.I Uggerhøj, NIM B 234 (2005) 31

pile-up is not an issue

These protons are lost anyway !

A few figures on the (extracted) proton beam

- Beam loss: $10^9 p^+ s^{-1}$
- Extracted intensity: $5 \times 10^8 p^+ s^{-1}$ (1/2 the beam loss)
- Number of p^+ : 2808 bunches of $1.15 \times 10^{11} p^+ = 3.2 \times 10^{14} p^+$
- Revolution frequency: Each bunch passes the extraction point at a rate of $3.10^5 \text{ km} \cdot \text{s}^{-1} / 27 \text{ km} \simeq 11 \text{ kHz}$
- Extracted “mini” bunches:
 - the crystal sees $2808 \times 11000 \text{ s}^{-1} \simeq 3.10^7 \text{ bunches s}^{-1}$
 - one extracts $5.10^8 / 3.10^7 \simeq 15 p^+$ from each bunch at each pass
 - Provided that the probability of interaction with the target is below 5%,
- Extraction over a 10h fill:
 - $5 \times 10^8 p^+ \times 3600 \text{ s h}^{-1} \times 10 \text{ h} = 1.8 \times 10^{13} p^+ \text{ fill}^{-1}$
 - This means $1.8 \times 10^{13} / 3.2 \times 10^{14} \simeq 5.6\%$ of the p^+ in the beam
- similar figures for the Pb-beam extraction

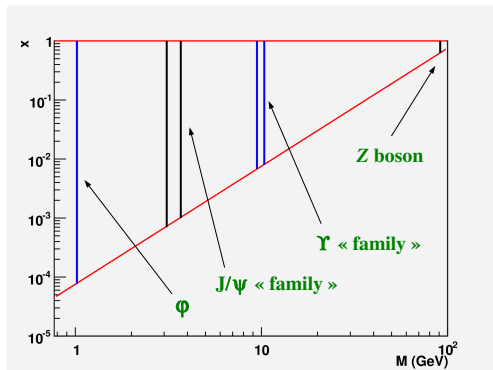
E. Uggerhøj, U.I Uggerhøj, NIM B 234 (2005) 31

pile-up is not an issue

These protons are lost anyway !

AFTER@LHC: A dilepton observatory ?

→ Region in x probed by dilepton production as function of $M_{\ell\ell}$

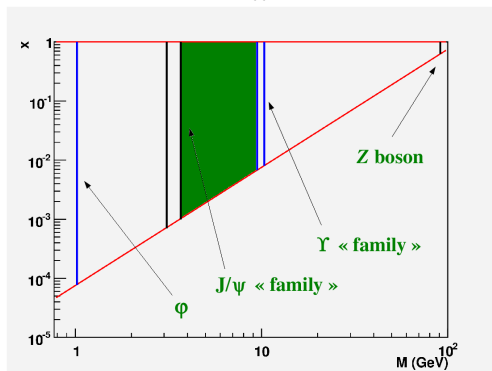


AFTER@LHC: A dilepton observatory ?

→ Region in x probed by dilepton production as function of $M_{\ell\ell}$

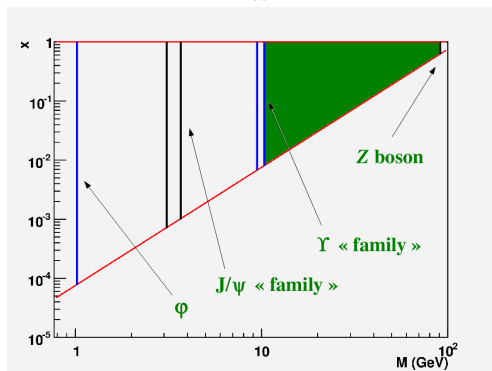
→ Above $c\bar{c}$: $x \in [10^{-3}, 1]$

→ Above $b\bar{b}$: $x \in [9 \times 10^{-3}, 1]$



AFTER@LHC: A dilepton observatory ?

- Region in x probed by dilepton production as function of $M_{\ell\ell}$
- Above $c\bar{c}$: $x \in [10^{-3}, 1]$
- Above $b\bar{b}$: $x \in [9 \times 10^{-3}, 1]$



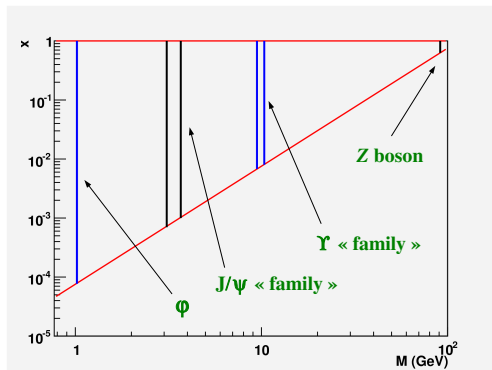
AFTER@LHC: A dilepton observatory ?

→ Region in x probed by dilepton production as function of $M_{\ell\ell}$

→ Above $c\bar{c}$: $x \in [10^{-3}, 1]$

→ Above $b\bar{b}$: $x \in [9 \times 10^{-3}, 1]$

Note: $x_{target} (\equiv x_2) > x_{projectile} (\equiv x_1)$
 “backward” region



AFTER@LHC: A dilepton observatory ?

→ Region in x probed by dilepton production as function of $M_{\ell\ell}$

→ Above $c\bar{c}$: $x \in [10^{-3}, 1]$

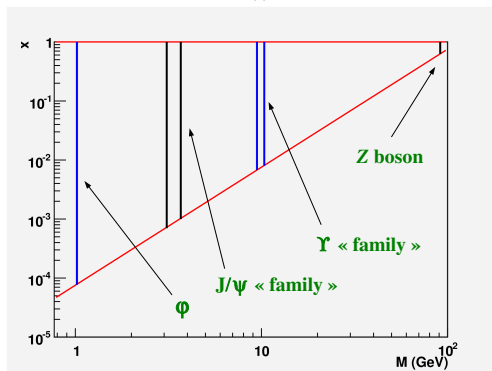
→ Above $b\bar{b}$: $x \in [9 \times 10^{-3}, 1]$

Note: $x_{target} (\equiv x_2) > x_{projectile} (\equiv x_1)$
“backward” region

→ **sea-quark asymmetries**
via p and d studies

- at large(est) x : backward (“easy”)

- at small(est) x : forward (need to stop the (extracted) beam)



AFTER@LHC: A dilepton observatory ?

⇒ Region in x probed by dilepton production as function of $M_{\ell\ell}$

→ Above $c\bar{c}$: $x \in [10^{-3}, 1]$

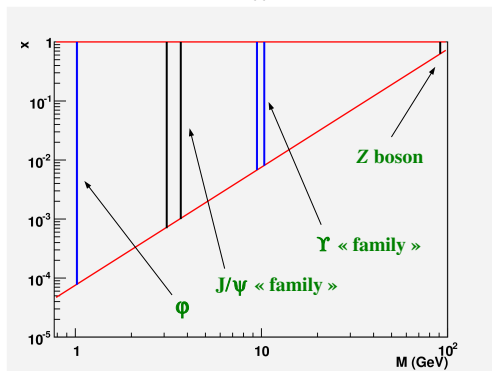
→ Above $b\bar{b}$: $x \in [9 \times 10^{-3}, 1]$

Note: $x_{target} (\equiv x_2) > x_{projectile} (\equiv x_1)$
“backward” region

→ sea-quark asymmetries
via p and d studies

- at large(est) x : backward (“easy”)

- at small(est) x : forward (need to stop the (extracted) beam)



⇒ To do: to look at the rates to see how competitive this will be

AFTER, among other things, a quarkonium observatory in pp

- Interpolating the world data set:

Target	$\int \mathcal{L} \text{ (fb}^{-1}\cdot\text{yr}^{-1}\text{)}$	$N(\text{J}/\Psi) \text{ yr}^{-1}$ $= A\mathcal{L}B\sigma_{\Psi}$	$N(\Upsilon) \text{ yr}^{-1}$ $= A\mathcal{L}B\sigma_{\Upsilon}$
1 m Liq. H₂	20	4.0 10⁸	8.0 10⁵
1 m Liq. D₂	24	9.6 10⁸	1.9 10⁶
LHC pp 14 Tev (low pT)	0.05 (ALICE) 2 LHCb	3.6 10⁷ 1.4 10⁹	1.8 10⁵ 7.2 10⁶
RHIC pp 200GeV	1.2 10⁻²	4.8 10⁵	1.2 10³

AFTER, among other things, a quarkonium observatory in pp

- Interpolating the world data set:

Target	$\int \mathcal{L} \text{ (fb}^{-1}\cdot\text{yr}^{-1}\text{)}$	$N(\text{J}/\Psi) \text{ yr}^{-1}$ $= A\mathcal{L}B\sigma_{\Psi}$	$N(\Upsilon) \text{ yr}^{-1}$ $= A\mathcal{L}B\sigma_{\Upsilon}$
1 m Liq. H₂	20	4.0 10⁸	8.0 10⁵
1 m Liq. D₂	24	9.6 10⁸	1.9 10⁶
LHC pp 14 Tev (low pT)	0.05 (ALICE) 2 LHCb	3.6 10⁷ 1.4 10⁹	1.8 10⁵ 7.2 10⁶
RHIC pp 200GeV	1.2 10⁻²	4.8 10⁵	1.2 10³

- 1000 times higher than at RHIC; comparable to ALICE/LHCb at the LHC

AFTER, among other things, a quarkonium observatory in pp

- Interpolating the world data set:

Target	$\int \mathcal{L} \text{ (fb}^{-1}\cdot\text{yr}^{-1}\text{)}$	$N(\text{J}/\Psi) \text{ yr}^{-1}$ $= A\mathcal{L}B\sigma_{\Psi}$	$N(\Upsilon) \text{ yr}^{-1}$ $= A\mathcal{L}B\sigma_{\Upsilon}$
1 m Liq. H₂	20	4.0 10⁸	8.0 10⁵
1 m Liq. D₂	24	9.6 10⁸	1.9 10⁶
LHC pp 14 Tev (low pT)	0.05 (ALICE) 2 LHCb	3.6 10⁷ 1.4 10⁹	1.8 10⁵ 7.2 10⁶
RHIC pp 200GeV	1.2 10⁻²	4.8 10⁵	1.2 10³

- 1000 times higher than at RHIC; comparable to ALICE/LHCb at the LHC
- Numbers are for only one unit of rapidity about 0

AFTER, among other things, a quarkonium observatory in pp

- Interpolating the world data set:

Target	$\int \mathcal{L} \text{ (fb}^{-1}\cdot\text{yr}^{-1}\text{)}$	$N(\text{J}/\Psi) \text{ yr}^{-1}$ $= A\mathcal{L}B\sigma_{\Psi}$	$N(\Upsilon) \text{ yr}^{-1}$ $= A\mathcal{L}B\sigma_{\Upsilon}$
1 m Liq. H₂	20	4.0 10⁸	8.0 10⁵
1 m Liq. D₂	24	9.6 10⁸	1.9 10⁶
LHC pp 14 Tev (low pT)	0.05 (ALICE) 2 LHCb	3.6 10⁷ 1.4 10⁹	1.8 10⁵ 7.2 10⁶
RHIC pp 200GeV	1.2 10⁻²	4.8 10⁵	1.2 10³

- 1000 times higher than at RHIC; comparable to ALICE/LHCb at the LHC
- Numbers are for only one unit of rapidity about 0
- Unique access in the backward region

AFTER, among other things, a quarkonium observatory in pp

- Interpolating the world data set:

Target	$\int \mathcal{L} \text{ (fb}^{-1}\cdot\text{yr}^{-1}\text{)}$	$N(\text{J}/\Psi) \text{ yr}^{-1}$ $= A\mathcal{L}B\sigma_{\Psi}$	$N(\Upsilon) \text{ yr}^{-1}$ $= A\mathcal{L}B\sigma_{\Upsilon}$
1 m Liq. H₂	20	4.0 10⁸	8.0 10⁵
1 m Liq. D₂	24	9.6 10⁸	1.9 10⁶
LHC pp 14 Tev (low pT)	0.05 (ALICE) 2 LHCb	3.6 10⁷ 1.4 10⁹	1.8 10⁵ 7.2 10⁶
RHIC pp 200GeV	1.2 10⁻²	4.8 10⁵	1.2 10³

- 1000 times higher than at RHIC; comparable to ALICE/LHCb at the LHC
- Numbers are for only one unit of rapidity about 0
- Unique access in the backward region
- Probe of the (very) large x in the target

Need for a quarkonium observatory

- Many **hopes** were put in **quarkonium studies** to extract **gluon PDF**

Need for a quarkonium observatory

- Many **hopes** were put in **quarkonium studies** to extract **gluon PDF**
 - in photo/lepto production (DIS)
 - but also pp collisions in gg -fusion process
 - mainly because of the presence of a natural “hard” scale: m_Q
 - and the good detectability of a dimuon pair

Need for a quarkonium observatory

- Many **hopes** were put in **quarkonium studies** to extract **gluon PDF**
 - in photo/lepto production (DIS)
 - but also pp collisions in gg -fusion process
 - mainly because of the presence of a natural “hard” scale: m_Q
 - and the good detectability of a dimuon pair

PHYSICAL REVIEW D

VOLUME 37, NUMBER 5

1 MARCH 1988

Structure-function analysis and ψ , jet, W , and Z production: Determining the gluon distribution

A. D. Martin

Department of Physics, University of Durham, Durham, England

R. G. Roberts

Rutherford Appleton Laboratory, Didcot, Oxon, England

W. J. Stirling

Department of Physics, University of Durham, Durham, England

(Received 27 July 1987)

We perform a next-to-leading-order structure-function analysis of deep-inelastic μN and νN scattering data and find acceptable fits for a range of input gluon distributions. We show three equally acceptable sets of parton distributions which correspond to gluon distributions which are (1) “soft,” (2) “hard,” and (3) which behave as $xG(x) \sim 1/\sqrt{x}$ at small x . J/ψ and prompt photon hadroproduction data are used to discriminate between the three sets. Set 1, with the “soft”-gluon distribution, is favored. W , Z , and jet production data from the CERN collider are well described but do not distinguish between the sets of structure functions. The precision of the predictions for σ_W and σ_Z allow the collider measurements to yield information on the number of light neutrinos and the mass of the top quark. Finally we discuss how the gluon distribution at very small x may be directly measured at DESY HERA.

Need for a quarkonium observatory

- Many **hopes** were put in **quarkonium studies** to extract **gluon PDF**
 - in photo/lepto production (DIS)
 - but also pp collisions in gg -fusion process
 - mainly because of the presence of a natural “hard” scale: m_Q
 - and the good detectability of a dimuon pair

PHYSICAL REVIEW D

VOLUME 37, NUMBER 5

1 MARCH 1988

Structure-function analysis and ψ , jet, W , and Z production: Determining the gluon distribution

A. D. Martin

Department of Physics, University of Durham, Durham, England

R. G. Roberts

Rutherford Appleton Laboratory, Didcot, Oxon, England

W. J. Stirling

Department of Physics, University of Durham, Durham, England

(Received 27 July 1987)

We perform a next-to-leading-order structure-function analysis of deep-inelastic μN and νN scattering data and find acceptable fits for a range of input gluon distributions. We show three equally acceptable sets of parton distributions which correspond to gluon distributions which are (1) “soft,” (2) “hard,” and (3) which behave as $xG(x) \sim 1/\sqrt{x}$ at small x . J/ψ and prompt photon hadroproduction data are used to discriminate between the three sets. Set 1, with the “soft”-gluon distribution, is favored. W , Z , and jet production data from the CERN collider are well described but do not distinguish between the sets of structure functions. The precision of the predictions for σ_W and σ_Z allow the collider measurements to yield information on the number of light neutrinos and the mass of the top quark. Finally we discuss how the gluon distribution at very small x may be directly measured at DESY HERA.

- Production **puzzle** → quarkonium not used anymore in global fits

Need for a quarkonium observatory

- Many **hopes** were put in **quarkonium studies** to extract **gluon PDF**
 - in photo/lepto production (DIS)
 - but also pp collisions in gg -fusion process
 - mainly because of the presence of a natural “hard” scale: m_Q
 - and the good detectability of a dimuon pair

PHYSICAL REVIEW D

VOLUME 37, NUMBER 5

1 MARCH 1988

Structure-function analysis and ψ , jet, W , and Z production: Determining the gluon distribution

A. D. Martin

Department of Physics, University of Durham, Durham, England

R. G. Roberts

Rutherford Appleton Laboratory, Didcot, Oxon, England

W. J. Stirling

Department of Physics, University of Durham, Durham, England

(Received 27 July 1987)

We perform a next-to-leading-order structure-function analysis of deep-inelastic μN and νN scattering data and find acceptable fits for a range of input gluon distributions. We show three equally acceptable sets of parton distributions which correspond to gluon distributions which are (1) “soft,” (2) “hard,” and (3) which behave as $xG(x) \sim 1/\sqrt{x}$ at small x . J/ψ and prompt photon hadroproduction data are used to discriminate between the three sets. Set 1, with the “soft”-gluon distribution, is favored. W , Z , and jet production data from the CERN collider are well described but do not distinguish between the sets of structure functions. The precision of the predictions for σ_W and σ_Z allow the collider measurements to yield information on the number of light neutrinos and the mass of the top quark. Finally we discuss how the gluon distribution at very small x may be directly measured at DESY HERA.

- Production **puzzle** → quarkonium not used anymore in global fits
- With systematic studies, one would **restore its status as gluon probe**

AFTER: also a quarkonium observatory in pA

Target	A	$\int \mathcal{L} \text{ (fb}^{-1}\cdot\text{yr}^{-1})$	$N(J/\Psi) \text{ yr}^{-1}$ $= A\mathcal{L}B\sigma_{\Psi}$	$N(\Upsilon) \text{ yr}^{-1}$ $= A\mathcal{L}B\sigma_{\Upsilon}$
1cm Be	9	0.62	1.1 10^8	2.2 10^5
1cm Cu	64	0.42	5.3 10^8	1.1 10^6
1cm W	185	0.31	1.1 10^9	2.3 10^6
1cm Pb	207	0.16	6.7 10^8	1.3 10^6
LHC pPb 8.8 TeV	207	10^{-4}	1.0 10^7	7.5 10^4
RHIC dAu 200GeV	198	1.5 10^{-4}	2.4 10^6	5.9 10^3
RHIC dAu 62GeV	198	3.8 10^{-6}	1.2 10^4	18

- In principle, one can get **300 times more J/ψ** –not counting the likely wider y coverage– than at RHIC, allowing for

AFTER: also a quarkonium observatory in pA

Target	A	$\int \mathcal{L} \text{ (fb}^{-1}\cdot\text{yr}^{-1})$	$N(J/\Psi) \text{ yr}^{-1}$ $= A\mathcal{L}B\sigma_{\Psi}$	$N(\Upsilon) \text{ yr}^{-1}$ $= A\mathcal{L}B\sigma_{\Upsilon}$
1cm Be	9	0.62	1.1 10^8	2.2 10^5
1cm Cu	64	0.42	5.3 10^8	1.1 10^6
1cm W	185	0.31	1.1 10^9	2.3 10^6
1cm Pb	207	0.16	6.7 10^8	1.3 10^6
LHC pPb 8.8 TeV	207	10^{-4}	1.0 10^7	7.5 10^4
RHIC dAu 200GeV	198	$1.5 \cdot 10^{-4}$	2.4 10^6	5.9 10^3
RHIC dAu 62GeV	198	$3.8 \cdot 10^{-6}$	1.2 10^4	18

- In principle, one can get **300 times more J/ψ** –not counting the likely wider y coverage– than at RHIC, allowing for
 - χ_c measurement in pA via $J/\psi + \gamma$ (extending Hera-B studies)

AFTER: also a quarkonium observatory in pA

Target	A	$\int \mathcal{L} \text{ (fb}^{-1}\cdot\text{yr}^{-1})$	$N(J/\Psi) \text{ yr}^{-1}$ $= A\mathcal{L}B\sigma_{\Psi}$	$N(\Upsilon) \text{ yr}^{-1}$ $= A\mathcal{L}B\sigma_{\Upsilon}$
1cm Be	9	0.62	1.1 10^8	2.2 10^5
1cm Cu	64	0.42	5.3 10^8	1.1 10^6
1cm W	185	0.31	1.1 10^9	2.3 10^6
1cm Pb	207	0.16	6.7 10^8	1.3 10^6
LHC pPb 8.8 TeV	207	10^{-4}	1.0 10^7	7.5 10^4
RHIC dAu 200GeV	198	$1.5 \cdot 10^{-4}$	2.4 10^6	5.9 10^3
RHIC dAu 62GeV	198	$3.8 \cdot 10^{-6}$	1.2 10^4	18

- In principle, one can get **300 times more J/ψ** –not counting the likely wider y coverage– than at RHIC, allowing for
 - χ_c measurement in pA via $J/\psi + \gamma$ (extending Hera-B studies)
 - **Polarisation** measurement as **the centrality, y or P_T**

AFTER: also a quarkonium observatory in pA

Target	A	$\int \mathcal{L} \text{ (fb}^{-1}\cdot\text{yr}^{-1})$	$N(J/\Psi) \text{ yr}^{-1}$ $= A\mathcal{L}B\sigma_{\Psi}$	$N(\Upsilon) \text{ yr}^{-1}$ $= A\mathcal{L}B\sigma_{\Upsilon}$
1cm Be	9	0.62	1.1 10^8	2.2 10^5
1cm Cu	64	0.42	5.3 10^8	1.1 10^6
1cm W	185	0.31	1.1 10^9	2.3 10^6
1cm Pb	207	0.16	6.7 10^8	1.3 10^6
LHC pPb 8.8 TeV	207	10^{-4}	1.0 10^7	7.5 10^4
RHIC dAu 200GeV	198	$1.5 \cdot 10^{-4}$	2.4 10^6	5.9 10^3
RHIC dAu 62GeV	198	$3.8 \cdot 10^{-6}$	1.2 10^4	18

- In principle, one can get **300 times more J/ψ** –not counting the likely wider y coverage– than at RHIC, allowing for
 - χ_c measurement in pA via $J/\psi + \gamma$ (extending Hera-B studies)
 - **Polarisation** measurement as **the centrality, y or P_T**
 - Ratio ψ' over **direct J/ψ** measurement in pA

AFTER: also a quarkonium observatory in pA

Target	A	$\int \mathcal{L} \text{ (fb}^{-1}\cdot\text{yr}^{-1})$	$N(J/\Psi) \text{ yr}^{-1}$ $= A\mathcal{L}B\sigma_{\Psi}$	$N(\Upsilon) \text{ yr}^{-1}$ $= A\mathcal{L}B\sigma_{\Upsilon}$
1cm Be	9	0.62	1.1 10^8	2.2 10^5
1cm Cu	64	0.42	5.3 10^8	1.1 10^6
1cm W	185	0.31	1.1 10^9	2.3 10^6
1cm Pb	207	0.16	6.7 10^8	1.3 10^6
LHC pPb 8.8 TeV	207	10^{-4}	1.0 10^7	7.5 10^4
RHIC dAu 200GeV	198	$1.5 \cdot 10^{-4}$	2.4 10^6	5.9 10^3
RHIC dAu 62GeV	198	$3.8 \cdot 10^{-6}$	1.2 10^4	18

- In principle, one can get **300 times more J/ψ** –not counting the likely wider y coverage– than at RHIC, allowing for
 - χ_c measurement in pA via $J/\psi + \gamma$ (extending Hera-B studies)
 - **Polarisation** measurement as **the centrality, y or P_T**
 - Ratio ψ' over **direct J/ψ** measurement in pA
 - not to mention ratio with **open charm, Drell-Yan**, etc ...

What for ?

- The **target versatility** of a fixed-target experiment is undisputable

What for ?

- The **target versatility** of a fixed-target experiment is undisputable
- A **wide rapidity coverage** is needed for:
 - a precise analysis of **gluon nuclear PDF**: $y, p_T \leftrightarrow x_2$
 - a handle on **formation time effects**

What for ?

- The **target versatility** of a fixed-target experiment is undisputable
- A **wide rapidity coverage** is needed for:
 - a precise analysis of **gluon nuclear PDF**: $y, p_T \leftrightarrow x_2$
 - a handle on **formation time effects**
- Strong need for **cross checks** from **various** measurements

What for ?

- The **target versatility** of a fixed-target experiment is undisputable
- A **wide rapidity coverage** is needed for:
 - a precise analysis of **gluon nuclear PDF**: $y, p_T \leftrightarrow x_2$
 - a handle on **formation time effects**
- Strong need for **cross checks** from **various** measurements
- The **backward kinematics** is very useful for large- x_{target} studies

What for ?

- The **target versatility** of a fixed-target experiment is undisputable
- A **wide rapidity coverage** is needed for:
 - a precise analysis of **gluon nuclear PDF**: $y, p_T \leftrightarrow x_2$
 - a handle on **formation time effects**
- Strong need for **cross checks** from **various** measurements
- The **backward kinematics** is very useful for large- x_{target} studies
 - What is the amount of Intrinsic charm ? Is it color filtered ?

What for ?

- The **target versatility** of a fixed-target experiment is undisputable
- A **wide rapidity coverage** is needed for:
 - a precise analysis of **gluon nuclear PDF**: $y, p_T \leftrightarrow x_2$
 - a handle on **formation time effects**
- Strong need for **cross checks** from **various** measurements
- The **backward kinematics** is very useful for large- x_{target} studies
 - What is the amount of Intrinsic charm ? Is it color filtered ?
 - **Is there an EMC effect for gluon ?** (reminder: EMC region $0.3 < x < 0.7$)

What for ?

- The **target versatility** of a fixed-target experiment is undisputable
- A **wide rapidity coverage** is needed for:
 - a precise analysis of **gluon nuclear PDF**: $\gamma, p_T \leftrightarrow x_2$
 - a handle on **formation time effects**
- Strong need for **cross checks** from **various** measurements
- The **backward kinematics** is very useful for large- x_{target} studies
 - What is the amount of Intrinsic charm ? Is it color filtered ?
 - **Is there an EMC effect for gluon ?** (reminder: EMC region $0.3 < x < 0.7$)
- One should be careful with factorization breaking effects:

This calls for multiple measurements to (in)validate factorisation

AFTER: also an heavy-flavour observatory in PbA

- Luminosities and yields with the extracted 2.76 TeV Pb beam

$$(\sqrt{s_{NN}} = 72 \text{ GeV})$$

Target	A.B	$\int \mathcal{L} \text{ (nb}^{-1}\cdot\text{yr}^{-1}\text{)}$	$N(J/\Psi) \text{ yr}^{-1}$ $= AB\mathcal{L}B\sigma_{\Psi}$	$N(\Upsilon) \text{ yr}^{-1}$ $= AB\mathcal{L}B\sigma_{\Upsilon}$
1 m Liq. H₂	207.1	800	3.4 10⁶	6.9 10³
1cm Be	207.9	25	9.1 10⁵	1.9 10³
1cm Cu	207.64	17	4.3 10⁶	0.9 10³
1cm W	207.185	13	9.7 10⁶	1.9 10⁴
1cm Pb	207.207	7	5.7 10⁶	1.1 10⁴
LHC PbPb 5.5 TeV	207.207	0.5	7.3 10⁶	3.6 10⁴
RHIC AuAu 200GeV	198.198	2.8	4.4 10⁶	1.1 10⁴
RHIC AuAu 62GeV	198.198	0.13	4.0 10⁴	61

AFTER: also an heavy-flavour observatory in PbA

- Luminosities and yields with the extracted 2.76 TeV Pb beam

$$(\sqrt{s_{NN}} = 72 \text{ GeV})$$

Target	A.B	$\int \mathcal{L} \text{ (nb}^{-1}\cdot\text{yr}^{-1}\text{)}$	$N(J/\Psi) \text{ yr}^{-1}$ = $AB\mathcal{L}B\sigma_{\Psi}$	$N(\Upsilon) \text{ yr}^{-1}$ = $AB\mathcal{L}B\sigma_{\Upsilon}$
1 m Liq. H₂	207.1	800	3.4 10⁶	6.9 10³
1cm Be	207.9	25	9.1 10⁵	1.9 10³
1cm Cu	207.64	17	4.3 10⁶	0.9 10³
1cm W	207.185	13	9.7 10⁶	1.9 10⁴
1cm Pb	207.207	7	5.7 10⁶	1.1 10⁴
LHC PbPb 5.5 TeV	207.207	0.5	7.3 10⁶	3.6 10⁴
RHIC AuAu 200GeV	198.198	2.8	4.4 10⁶	1.1 10⁴
RHIC AuAu 62GeV	198.198	0.13	4.0 10⁴	61

- Yields **similar** to those of RHIC at 200 GeV,
100 times those of RHIC at 62 GeV

AFTER: also an heavy-flavour observatory in PbA

- Luminosities and yields with the extracted 2.76 TeV Pb beam

$$(\sqrt{s_{NN}} = 72 \text{ GeV})$$

Target	A.B	$\int \mathcal{L} \text{ (nb}^{-1}\cdot\text{yr}^{-1}\text{)}$	$N(J/\Psi) \text{ yr}^{-1}$ = $AB\mathcal{L}B\sigma_{\Psi}$	$N(\Upsilon) \text{ yr}^{-1}$ = $AB\mathcal{L}B\sigma_{\Upsilon}$
1 m Liq. H₂	207.1	800	3.4 10⁶	6.9 10³
1cm Be	207.9	25	9.1 10⁵	1.9 10³
1cm Cu	207.64	17	4.3 10⁶	0.9 10³
1cm W	207.185	13	9.7 10⁶	1.9 10⁴
1cm Pb	207.207	7	5.7 10⁶	1.1 10⁴
LHC PbPb 5.5 TeV	207.207	0.5	7.3 10⁶	3.6 10⁴
RHIC AuAu 200GeV	198.198	2.8	4.4 10⁶	1.1 10⁴
RHIC AuAu 62GeV	198.198	0.13	4.0 10⁴	61

- Yields **similar** to those of RHIC at 200 GeV,
100 times those of RHIC at 62 GeV
- Also **very competitive** compared to the **LHC**.

AFTER: also an heavy-flavour observatory in PbA

- Luminosities and yields with the extracted 2.76 TeV Pb beam

$$(\sqrt{s_{NN}} = 72 \text{ GeV})$$

Target	A.B	$\int \mathcal{L} \text{ (nb}^{-1}\cdot\text{yr}^{-1}\text{)}$	$N(J/\Psi) \text{ yr}^{-1}$ = $AB\mathcal{L}B\sigma_{\Psi}$	$N(\Upsilon) \text{ yr}^{-1}$ = $AB\mathcal{L}B\sigma_{\Upsilon}$
1 m Liq. H₂	207.1	800	3.4 10⁶	6.9 10³
1cm Be	207.9	25	9.1 10⁵	1.9 10³
1cm Cu	207.64	17	4.3 10⁶	0.9 10³
1cm W	207.185	13	9.7 10⁶	1.9 10⁴
1cm Pb	207.207	7	5.7 10⁶	1.1 10⁴
LHC PbPb 5.5 TeV	207.207	0.5	7.3 10⁶	3.6 10⁴
RHIC AuAu 200GeV	198.198	2.8	4.4 10⁶	1.1 10⁴
RHIC AuAu 62GeV	198.198	0.13	4.0 10⁴	61

- Yields **similar** to those of RHIC at 200 GeV,
100 times those of RHIC at 62 GeV
- Also **very competitive** compared to the **LHC**.

The same picture also holds for **open heavy flavour**

What for ?

Observation of J/ψ sequential suppression **seems to be hindered** by

- the **Cold Nuclear Matter effects**: non trivial and

... not well understood

What for ?

Observation of J/ψ sequential suppression **seems to be hindered** by

- the **Cold Nuclear Matter effects**: non trivial and ... not well understood
- the difficulty to observe directly the **excited states** which would melt before the ground states
 - χ_c **never studied** in AA collisions
 - $\psi(2S)$ **not yet** studied in AA collisions **at RHIC**

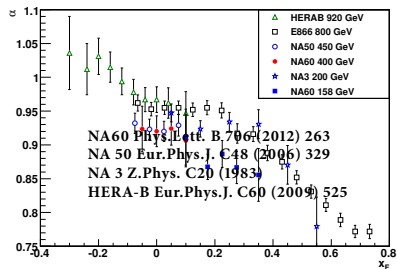
What for ?

Observation of J/ψ sequential suppression **seems to be hindered** by

- the **Cold Nuclear Matter effects**: non trivial and ... not well understood
- the difficulty to observe directly the **excited states** which would melt before the ground states
 - χ_c **never studied** in AA collisions
 - $\psi(2S)$ **not yet studied** in AA collisions **at RHIC**
- the possibilities for **$c\bar{c}$ recombination**
 - **Open charm** studies are **difficult** where recombination matters most i.e. at **low P_T**
 - Only indirect indications –from the γ and P_T dependence of R_{AA} – that recombination may be at work
 - CNM effects may show a non-trivial γ and P_T dependence ...

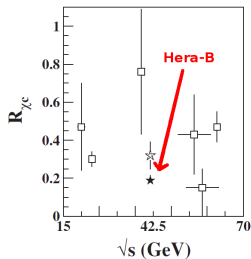
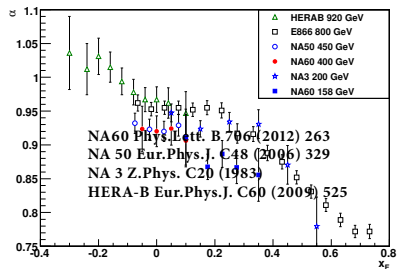
SPS and Hera-B

– J/ψ data in pA collisions



SPS and Hera-B

- J/ψ data in pA collisions
- χ_c data in pA collisions



HERA-B PRD 79 (2009) 012001, and ref. therein

LHB

Our idea is not completely new

Nuclear Instruments and Methods in Physics Research A 333 (1993) 125–135
North-Holland

**NUCLEAR
INSTRUMENTS
& METHODS
IN PHYSICS
RESEARCH**
Section A

LHB, a fixed target experiment at LHC to measure CP violation in B mesons

Flavio Costantini

University of Pisa and INFN, Italy

A fixed target experiment at LHC to measure CP violation in B mesons is presented. A description of the proposed apparatus is given together with its sensitivity on the CP violation asymmetry measurement for the two benchmark decay channels $B^0 \rightarrow J/\psi + K_s^0$, $B^0 \rightarrow \pi^+ \pi^-$. The possibility of obtaining an extracted LHC beam hinges on channeling in a bent silicon crystal. Recent results on beam extraction efficiencies measured at CERN SPS based on this technique are presented.

LHB

Our idea is not completely new

1. Introduction

...

This paper presents a fixed target experiment to measure CP violation in the B system based on the possibility of extracting the 8 TeV LHC proton beam using a bent silicon crystal [4]. A 10% extraction efficiency of the LHC beam halo will give an extracted beam intensity of about 10^8 protons/s allowing the production of as many as 10^{10} $B\bar{B}$ pairs per year, i.e. about two orders of magnitude more than what could be produced by an e^+e^- asymmetric B factory with 10^{34} $\text{cm}^{-2}\text{s}^{-1}$ luminosity [5].



LHB

Our idea is not completely new

1. Introduction

...

This paper presents a fixed target experiment to measure CP violation in the B system based on the possibility of extracting the 8 TeV LHC proton beam using a bent silicon crystal [4]. A 10% extraction efficiency of the LHC beam halo will give an extracted beam intensity of about 10^8 protons/s allowing the production of as many as 10^{10} $B\bar{B}$ pairs per year, i.e. about two orders of magnitude more than what could be produced by an e^+e^- asymmetric B factory with 10^{34} $\text{cm}^{-2}\text{s}^{-1}$ luminosity [5].

10^{10} $B\bar{B}$ pairs per year

- B-factories: 1 ab^{-1} means $10^9 B\bar{B}$ pairs



LHB

Our idea is not completely new

1. Introduction

...

This paper presents a fixed target experiment to measure CP violation in the B system based on the possibility of extracting the 8 TeV LHC proton beam using a bent silicon crystal [4]. A 10% extraction efficiency of the LHC beam halo will give an extracted beam intensity of about 10^8 protons/s allowing the production of as many as 10^{10} $B\bar{B}$ pairs per year, i.e. about two orders of magnitude more than what could be produced by an e^+e^- asymmetric B factory with 10^{34} $\text{cm}^{-2}\text{s}^{-1}$ luminosity [5].

10^{10} $B\bar{B}$ pairs per year

- B-factories: 1 ab^{-1} means $10^9 B\bar{B}$ pairs
- For LHCb, typically 1 fb^{-1} means $\simeq 2 \times 10^{11} B\bar{B}$ pairs at 14 TeV



LHB

Our idea is not completely new

1. Introduction

...

This paper presents a fixed target experiment to measure CP violation in the B system based on the possibility of extracting the 8 TeV LHC proton beam using a bent silicon crystal [4]. A 10% extraction efficiency of the LHC beam halo will give an extracted beam intensity of about 10^8 protons/s allowing the production of as many as 10^{10} $B\bar{B}$ pairs per year, i.e. about two orders of magnitude more than what could be produced by an e^+e^- asymmetric B factory with 10^{34} $\text{cm}^{-2}\text{s}^{-1}$ luminosity [5].

10^{10} $B\bar{B}$ pairs per year

- B-factories: 1 ab^{-1} means $10^9 B\bar{B}$ pairs
- For LHCb, typically 1 fb^{-1} means $\simeq 2 \times 10^{11} B\bar{B}$ pairs at 14 TeV
- LHB turned down in favour of LHCb mainly because of the **fear of a premature degradation of the bent crystal** due to radiation damages.



LHB

Our idea is not completely new

1. Introduction

...

This paper presents a fixed target experiment to measure CP violation in the B system based on the possibility of extracting the 8 TeV LHC proton beam using a bent silicon crystal [4]. A 10% extraction efficiency of the LHC beam halo will give an extracted beam intensity of about 10^8 protons/s allowing the production of as many as 10^{10} $B\bar{B}$ pairs per year, i.e. about two orders of magnitude more than what could be produced by an c^+e^- asymmetric B factory with 10^{34} $\text{cm}^{-2}\text{s}^{-1}$ luminosity [5].

10^{10} $B\bar{B}$ pairs per year



- B-factories: 1 ab^{-1} means $10^9 B\bar{B}$ pairs
- For LHCb, typically 1 fb^{-1} means $\simeq 2 \times 10^{11} B\bar{B}$ pairs at 14 TeV
- LHB turned down in favour of LHCb mainly because of the **fear of a premature degradation of the bent crystal** due to radiation damages.
- Nowadays, degradation is known to be $\simeq 6\%$ per 10^{20} particles/ cm^2
- 10^{20} particles/ cm^2 : one year of operation for realistic conditions

LHB

Our idea is not completely new

1. Introduction

...

This paper presents a fixed target experiment to measure CP violation in the B system based on the possibility of extracting the 8 TeV LHC proton beam using a bent silicon crystal [4]. A 10% extraction efficiency of the LHC beam halo will give an extracted beam intensity of about 10^8 protons/s allowing the production of as many as 10^{10} $B\bar{B}$ pairs per year, i.e. about two orders of magnitude more than what could be produced by an c^+e^- asymmetric B factory with 10^{34} $\text{cm}^{-2}\text{s}^{-1}$ luminosity [5].

10^{10} $B\bar{B}$ pairs per year



- B-factories: 1 ab^{-1} means $10^9 B\bar{B}$ pairs
- For LHCb, typically 1 fb^{-1} means $\simeq 2 \times 10^{11} B\bar{B}$ pairs at 14 TeV
- LHB turned down in favour of LHCb mainly because of the **fear of a premature degradation of the bent crystal** due to radiation damages.
- Nowadays, degradation is known to be $\simeq 6\%$ per 10^{20} particles/ cm^2
- 10^{20} particles/ cm^2 : one year of operation for realistic conditions
- After a year, one simply moves the crystal by less than one mm ...