polarized polarized and target polarized polarized polarized polarized polarized and the HESR with a solid polarized polarized and the HESR with a solid polarized and the HESR with a solid polarized and the HESR with a sol **On the feasibility of using an extracted polarized antiproton beam of the HESR with a solid polarized target**

Presenters: Yu. A. Plis, A.V. Smirnov

PSTP 2015, Bohum, Germany

Authors

- **G.M. Gurevich, Institute for Nuclear Research, RAS, 117312 Moscow, Russia.**
- **A.A. Lukhanin, Kharkiv Institute of Physics and Technology, Kharkiv, Ukraine.**
- **F. Maas, Helmholtz-Institut Mainz, Universität Mainz, D-55099, Germany.**
- **Yu.A. Plis, A.O. Sidorin, A.V. Smirnov, Yu.A. Usov, Joint Institute for Nuclear Research, 141980 Dubna, Russia.**
- **A. Thomas, Institut für Kernphysik, Universität Mainz, D-55099, Germany.**

It seems real to accelerate and store polarized antiprotons in the high-energy storage ring (HESR) of the future GSI Facility for Antiproton and Ion Research (FAIR).

It seems real to accelerate and store polarized antiprotons in the high-energy storage ring (HESR) of the future GSI Facility for Antiproton and Ion Research (FAIR).

Here we do not consider the problems connected with antiproton production and depolarization during acceleration.

```
A. Lehrach et al., SPIN2004, p. 742.
```
It is known that high-energy protons can be efficiently extracted from the beam halo in a storage ring (collider) using channeling in a bent crystal. R. Carrigan et al., Phys. Rev. ST - AC 1 (1998) 022801.

It is known that high-energy protons can be efficiently extracted from the beam halo in a storage ring (collider) using channeling in a bent crystal.

R. Carrigan et al., Phys. Rev. ST - AC 1 (1998) 022801.

The main beam and thus collider experiments would be undisturbed, and the extracted protons are only those which would otherwise be lost.

The experiments at the Tevatron (FNAL) at an energy of 900 GeV confirmed computer simulation and provided the extraction coefficient ≈ 10-6 , which for 10¹²circulating protons resulted in the extracted beam of 10⁶ protons/s.

R. Carrigan et al., Phys. Rev. AB 5 (2002) 043501.

The similar results were received at the CERN SPS. With a 120 GeV circulating beam of 5×10¹¹ protons typically 6×10⁵ protons/s were extracted.

K. Elsener et al., Nucl. Instr. & Meth. B 119 (1996) 215.

But for heavy negative particles (antiprotons) the method of beam deflection by a bent crystal does not work due to fast dechanneling.

V.P. Koshcheev, O.E. Krivosheev., Russian Phys. Journal. 38 (1995) 493; N.F. Shul'ga, I.V. Kirillin, V.I. Truten', Phys. Lett. B 702 (2011) 100.

But for heavy negative particles (antiprotons) the method of beam deflection by a bent crystal does not work due to fast dechanneling.

V.P. Koshcheev, O.E. Krivosheev., Russian Phys. Journal. 38 (1995) 493; N.F. Shul'ga, I.V. Kirillin, V.I. Truten', Phys. Lett. B 702 (2011) 100.

Instead, a stochastic mechanism of multiple scattering of negative particles on the bent chains of crystal atoms was proposed.

Experiments at CERN confirmed the computer simulation results of the stochastic mechanism for negative pion beam deflection by a bent crystal.

W. Scandale et al., Phys. Lett. B 680 (2009) 301.

We propose to extract the halo antiprotons of the HESR, which can not be utilized by the PANDA facility by means of the bent silicon crystal and/or magnetic septum.

The initial direction of polarization is vertical. If necessary, it can be rotated with the spin rotator in the horizontal direction before the polarized target.

We refer to the authors

N.F. Shul'ga, I.V. Kirillin, V.I. Truten',

http://iap.sumy.org/attacments/files/Kharkov2014/12. pdf.

which consider that the stochastic mechanism could be applicable for beam extraction from an antiproton storage ring.

The simulation had been carried out for a beam of 2000 antiprotons with an energy of 10 GeV incident on the silicon crystal in the direction of the <110> axis.

Beam divergence was taken equal to ≈10 μrad, a bending radius R = 1 m, crystal thickness L = 1 mm, bending angle $\alpha = L/R = 1$ **mrad.**

2000 antiprotons, *E*=10 GeV, *L*=1 mm, *R*=1 m, Si
$$
\langle 110 \rangle
$$
, ψ_c =145 prad, $\alpha = \frac{L}{R}$ =1 mrad

Stochastic beam extraction. Eap = 10 GeV, L = 1 mm, R = 1 m, α = L/R = 1 mrad, Si <110>

The simulation had been carried out for a beam of 2000 antiprotons with an energy of 10 GeV incident on the silicon crystal in the direction of the <110> axis.

Beam divergence was taken equal to ≈10 μrad, a bending radius R = 1 m, crystal thickness L = 1 mm, bending angle $\alpha = L/R = 1$ **mrad.**

In this case the maximum of the distribution of the beam deflection reaches 0.3 mrad and the FWHM = 1 mrad.

For the crystal with a thickness L = 4 mm one has two maxima with 0.3 and 1 mrad with FWHM 2 mrad.

In the HESR

(FAIR Project (subproject HESR) Technical Report. http:// www.gsi.de\\/fairreports/btr-e.html)

the planned internal coasting beam could reach ≈10¹¹ antiprotons in the momentum range 1.5-15 GeV/c (energies 0.83-14.1 GeV).

In the HESR

(FAIR Project (subproject HESR) Technical Report. http:// www.gsi.de\\/fairreports/btr-e.html)

the planned internal coasting beam could reach ≈10¹¹ antiprotons in the momentum range 1.5-15 GeV/c (energies 0.83-14.1 GeV).

The beam lifetime in the working mode with the use of hydrogen pellets changes from 0.5 hour at the minimum energy to 2 hours at the maximum energy.

A. Lehrach et al., Nucl. Instr. \& Meth. A 561} (2006) 289.

Furthermore, the mechanism of halo formation in the HESR is different in comparison with the case of the Tevatron where the beam lifetime reached 70 hours. This implies, that the extraction coefficient for the HESR could turn out essentially higher than for the Tevatron.

Furthermore, the mechanism of halo formation in the HESR is different in comparison with the case of the Tevatron where the beam lifetime reached 70 hours. This implies, that the extraction coefficient for the HESR could turn out essentially higher than for the Tevatron.

We consider two possibilities for obtaining an extracted antiproton beam in the HESR with the purpose to use it independently for the study of polarization effects.

1) Since the accumulated antiprotons in the HESR are not utilized entirely, it is feasible to extract them slowly at the end of the cycle and send to the polarized target.

Since the accumulated antiprotons in the HESR are not utilized entirely, it is feasible to extract them slowly at the end of the cycle and send to the polarized target.

For example, at slow variation of the radial betatron oscillation frequency the particles get to the instability region and end up in the septum due to the resonance build-up of the oscillation amplitude.

Scheme of the polarized antiproton extraction.

Let the data acquisition stops when the number of particles drops by an order of value, i.e. down to 10¹⁰ or 10⁹ (depending on the chosen regime). If the particles are extracted for 10 s with the efficiency of 10%, then for the cycle duration equal to 1 hour an average intensity will be 10⁶ – 10⁵ antiproton/s, which, for the typical 10 cm long polarized target, corresponds to the luminosity 10³⁰ - 10²⁹cm-2 s -1 .

2) At the HESR one might try to extract halo antiprotons, which can not be utilised by the PANDA facility, with the assistance of a magnetic septum and, possibly, bent Si crystal.

Scheme of the polarized antiproton extraction.

It should be noted that the halo formation mechanism in the HESR differs from that of the FNAL storage ring, where intra-beam scattering was the main process.

At the HESR, at low momenta the main contribution comes from Coulomb scattering, while at high momenta – from nuclear processes

A. Lehrach et al., Nucl. Instr. & Meth. A 561 (2006) 289}).

The beam lifetime for the HESR is about 1 hour, while for the FNAL it was about 70 hours. It can be expected, that in the HESR the substantial share of antiprotons lost due to the interaction with the pellets comes to the halo.

The simulation was done with BETACOOL code ([http://betacool.jinr.ru\)](http://betacool.jinr.ru/) for PANDA parameters with use the barrier bucket and electron cooling system.

A. Smirnov, A. Sidorin, D. Krestnikov., Effective Luminosity Simulation for PANDA Experiment at FAIR. Proceeding of Workshop on Beam Cooling and Related Topics (COOL09). Lanchzhou, China (2009). http://jacow.org.

The main simulation parameters are presented in the following Table.

Simulation parameters

Particle distribution (dots) and barrier buckets (red line) on the longitudinal phase space.

The longitudinal acceptance was chosen equal to the effective barrier bucket height.

It means that particles will be lost after interaction with the target pellet if the energy decreasing is larger than the barrier height.

The longitudinal acceptance was chosen equal to the effective barrier bucket height.

It means that particles will be lost after interaction with the target pellet if the energy decreasing is larger than the barrier height.

If the particle losses are defined by the longitudinal acceptance and by hadronic interaction in the internal target (cross section is (51-100) mbarn, maximum value at 1.5 GeV/c), than the loss rate is about 4×10⁶ particles/s.

Particle number versus time. Particle losses due to the longitudinal acceptance and hadronic interaction in the internal target.

Luminosity evolution due to the longitudinal acceptance and hadronic interaction in the internal target.

Particle number versus time.

Particle losses due to the longitudinal acceptance.

In the case when the particle losses are defined by the longitudinal acceptance only, the particle losses rate can be estimated at the level about 10⁶ particles/s.

In the case when the particle losses are defined by the longitudinal acceptance only, the particle losses rate can be estimated on the level about 10⁶ particles/s.

This value can be used as the simple estimation of the intensity of the slow extracted beam for the chosen parameters. Note, that the HESR horizontal dispersion function has mainly positive values. It means that the particles usually can be extracted to the inner side of the storage ring.

In the case when the particle losses are defined by the longitudinal acceptance only, the particle losses rate can be estimated on the level about 10⁶ particles/s.

This value can be used as the simple estimation of the intensity of the slow extracted beam for the chosen parameters. Note, that the HESR horizontal dispersion function has mainly positive values. It means that the particles usually can be extracted to the inner side of the storage ring.

But the barrier bucket permits to extract the beam to the external side also.

For a 10 cm long polarized target with a nucleon density of 10²⁴ cm-3 the luminosity will be about 10³⁰ cm-2 s -1 .

For a 10 cm long polarized target with a nucleon density of 10²⁴ cm-3 the luminosity will be about 10³⁰ cm-2 s -1 .

For the estimation of the real intensity of the slow extracted beam, the BETACOOL program can be modified in the accordance with necessary physical and geometrical models (beam and target parameters, extraction kicker parameters, etc.).

The particle losses rates (extracted beam intensity) can be controlled with the different parameters of the barrier bucket and electron cooling system which will define the quality of the slow extracted beam also.

In summary, an extracted antiproton beam obtained by the use of one of the described methods could be employed in combination either with an dynamically polarized target (DPT) or frozen spin target (FST), depending on the parameters of the beam which will be reached as a result of joint efforts of scientists from Dubna and Mainz.

Of course, coauthors propose here only the principle way, but the realization of a similar project is a serious technical problem which could be solved by a close co-operation of all relevant experts solely.