

Higgs Boson Pair Production at NNLO in the EFT

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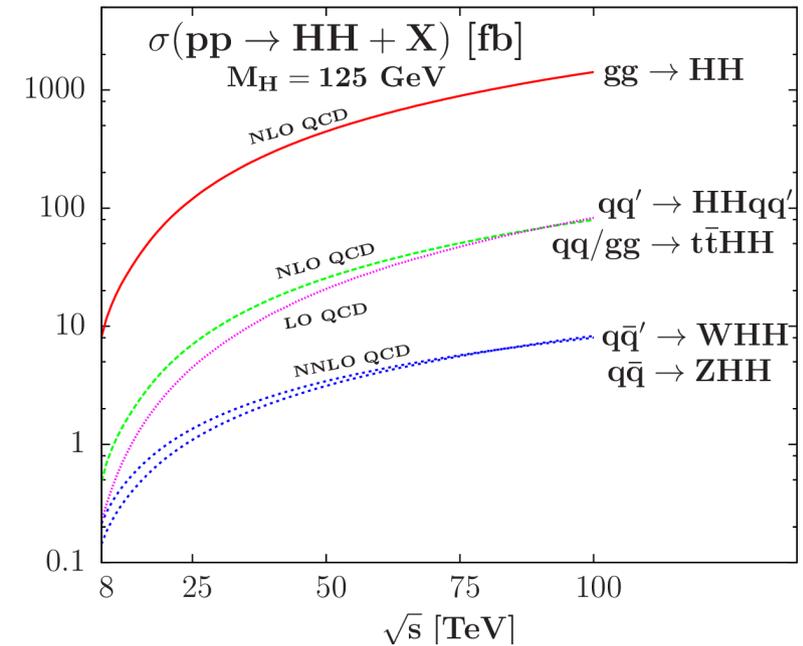
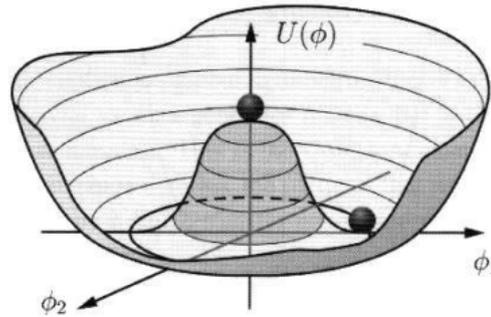
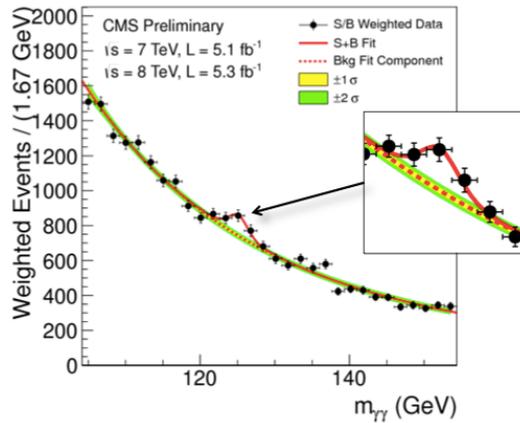
in collaboration with **Daniel de Florian**

HH subgroup meeting – October 2014, CERN

CONICET



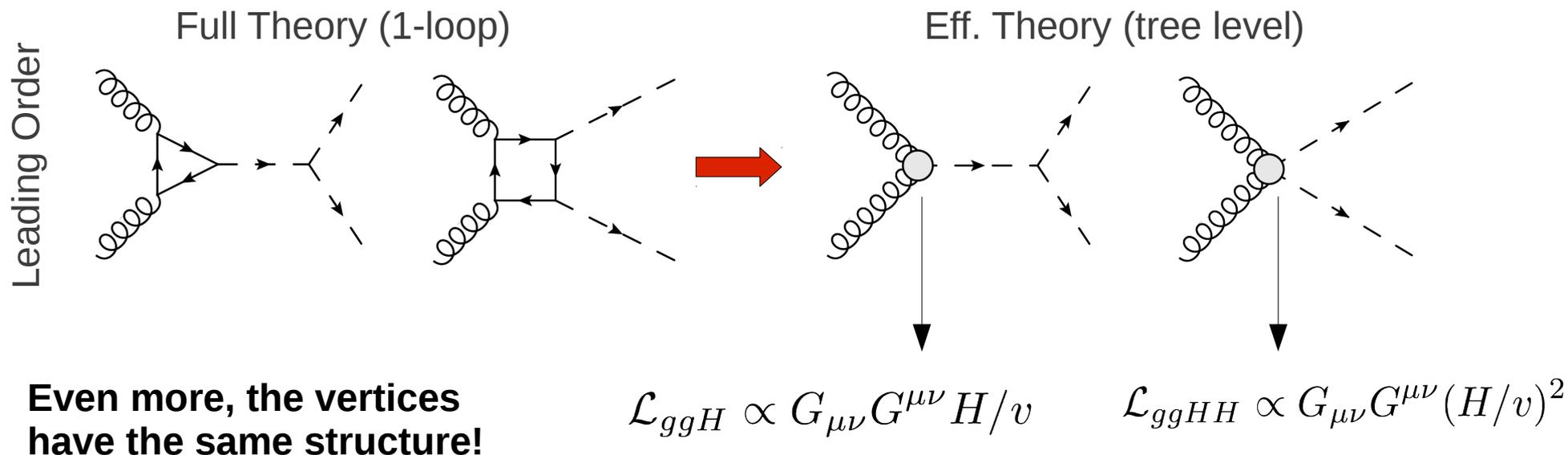
Motivation



- Why computing HH production via gluon fusion at NNLO in the EFT?
 - Gluon fusion is the main HH production channel
 - NLO corrections in the EFT are large
K factor close to 2
Large theoretical uncertainties
 - We need to improve both precision and accuracy of the prediction
Finite top mass effects at NLO
NNLO in the EFT ←

Calculation in the EFT

- Calculation is much simpler! Loop induced in the full theory, tree level in the EFT



- Even more, the vertices have the same structure!

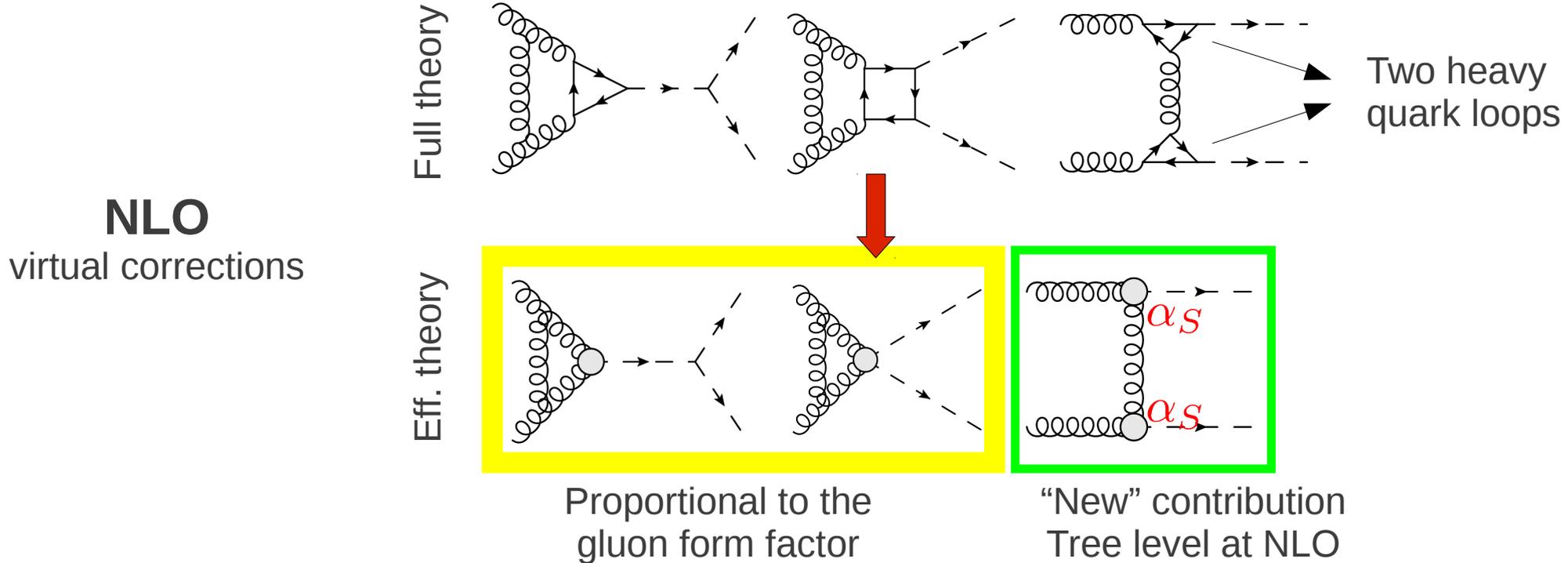
- We don't need to calculate those diagrams again: use the single Higgs case

$$\text{Diagram 1} = \text{Diagram 2} \times \text{Diagram 3} \quad (1)$$

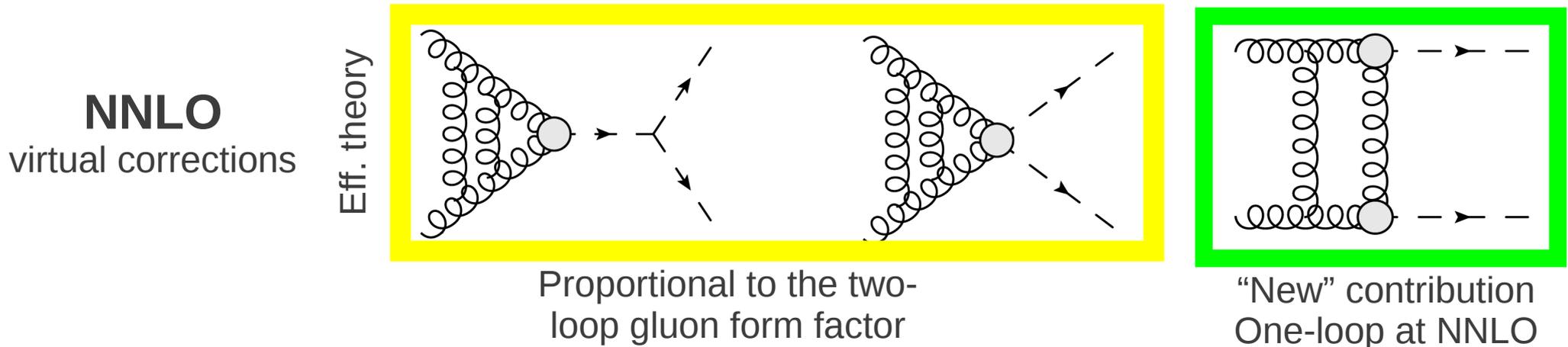
$$\text{Diagram 1} = \text{Diagram 2} \times \left(-\frac{1}{v}\right) \quad (2)$$

- These relations also hold at higher orders, in a diagram by diagram basis

Calculation in the EFT – virtual corrections



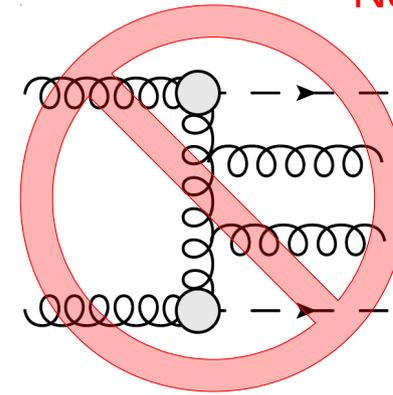
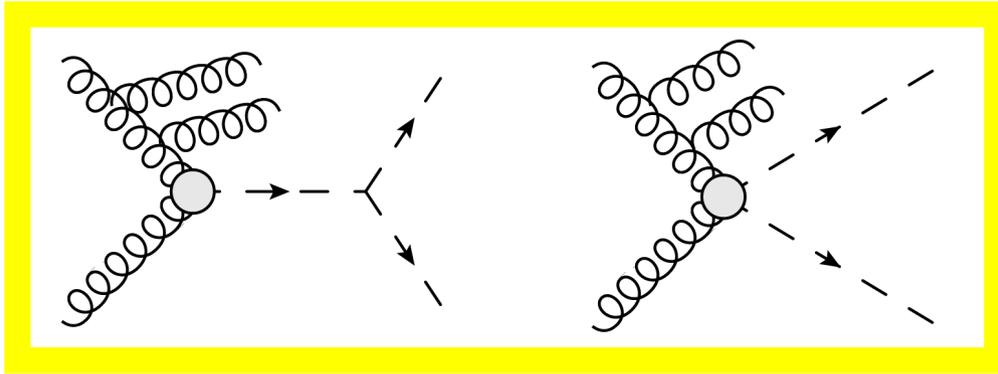
- For the NNLO virtual corrections we only need to compute one-loop diagrams



- Reduction into master integrals using FIRE

Calculation in the EFT – real corrections

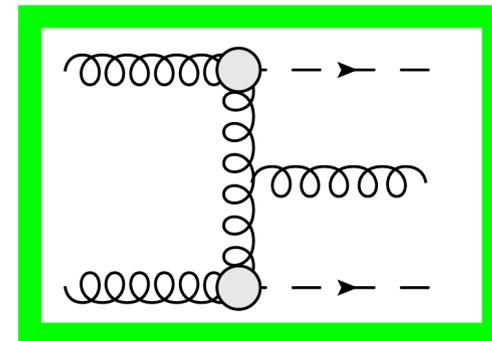
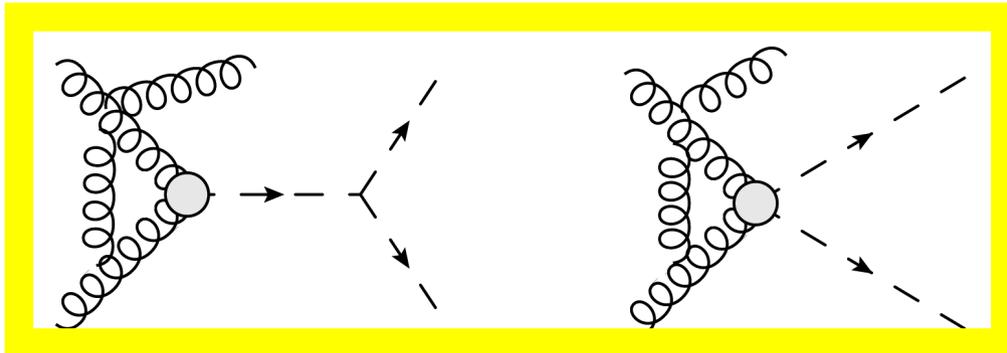
- **Double real emission:**



N3LO contribution

- Everything can be taken from the single Higgs results

- **Single real emission at “one-loop”:**



- Single soft and collinear divergencies at tree level: we use FKS subtraction

Calculation in the EFT – in short:

- We split the calculation: $Q^2 \frac{d\hat{\sigma}}{dQ^2} = \hat{\sigma}^a + \hat{\sigma}^b$

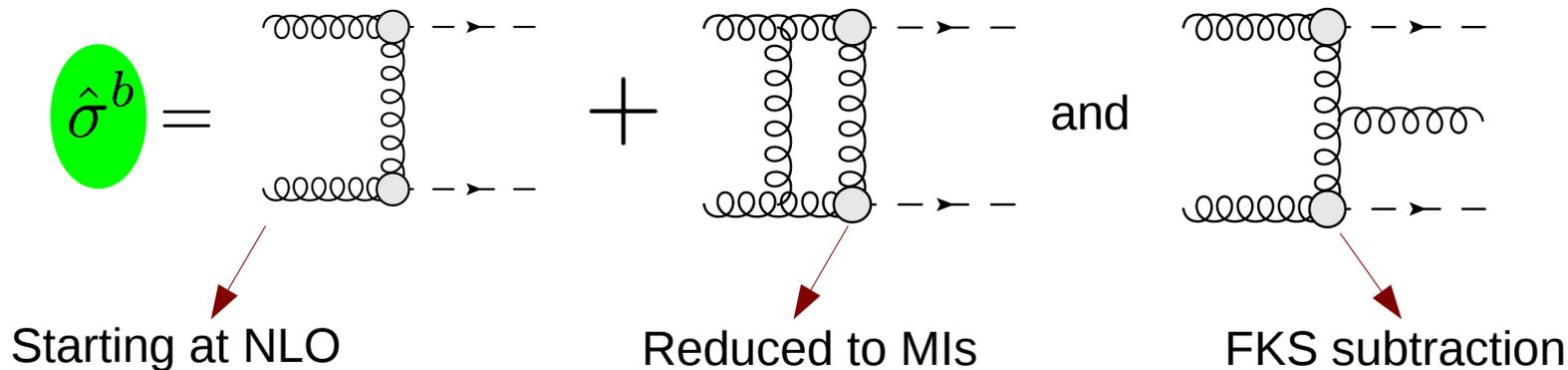
Single-Higgs like

New topologies with two effective vertices

$$\hat{\sigma}_{ij}^a = \hat{\sigma}_{LO} \left\{ \eta_{ij}^{(0)} + \left(\frac{\alpha_S}{\pi}\right) \eta_{ij}^{(1)} + \left(\frac{\alpha_S}{\pi}\right)^2 \left[\eta_{ij}^{(2)} + \delta_{ig} \delta_{jg} \delta(1-z) \frac{2\text{Re}(C_{LO})}{|C_{LO}|^2} (C_H^{(2)} - C_{HH}^{(2)}) \right] \right\}$$

Known coefficients from single Higgs XS

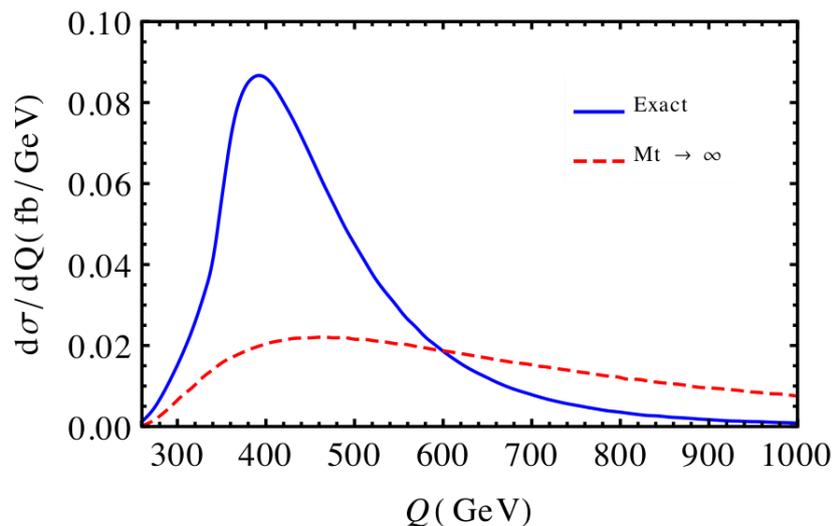
Differences in the second order corrections to the effective vertex
(Grigo et al., 1408.2422)



- The “new” contributions are essentially a NLO calculation of a $2 \rightarrow 2$ process

Top mass effects

- Is the EFT calculation reliable?



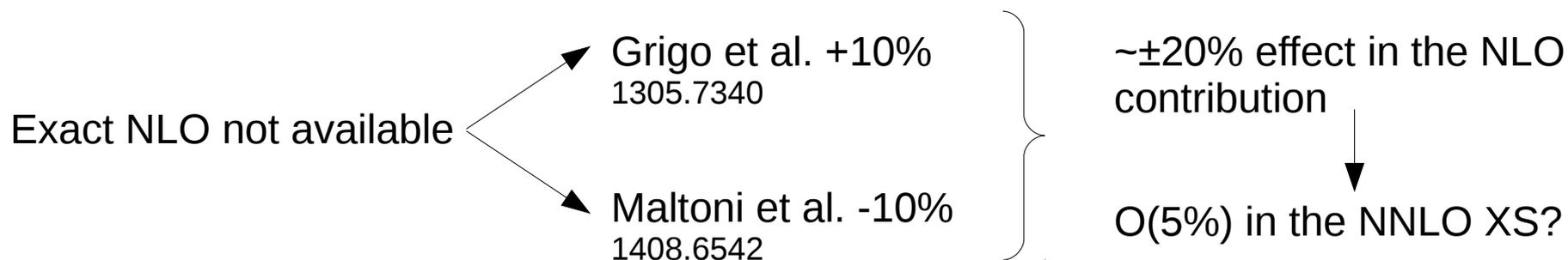
- Worse than in the single Higgs case (larger inv. mass)
- Not reliable for distributions
- Underestimation of the total XS at LO of O(20%)

- Should be more reliable to compute the QCD corrections

- Corrections are dominated by initial state soft radiation, not sensitive to the vertex structure

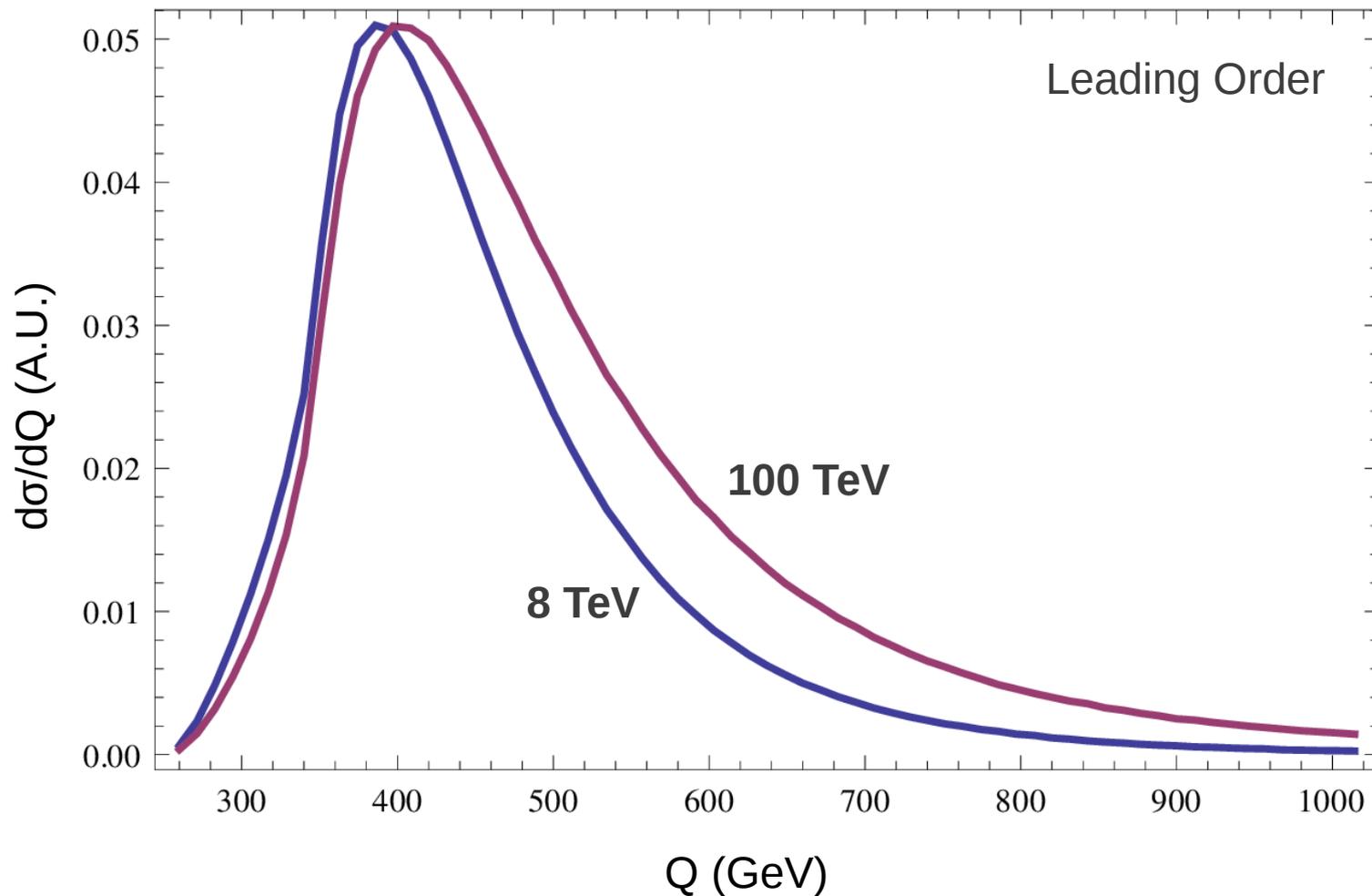
- Usual procedure: compute the corrections in the EFT and normalize by the exact LO

- First step: validate it at NLO



Top mass effects

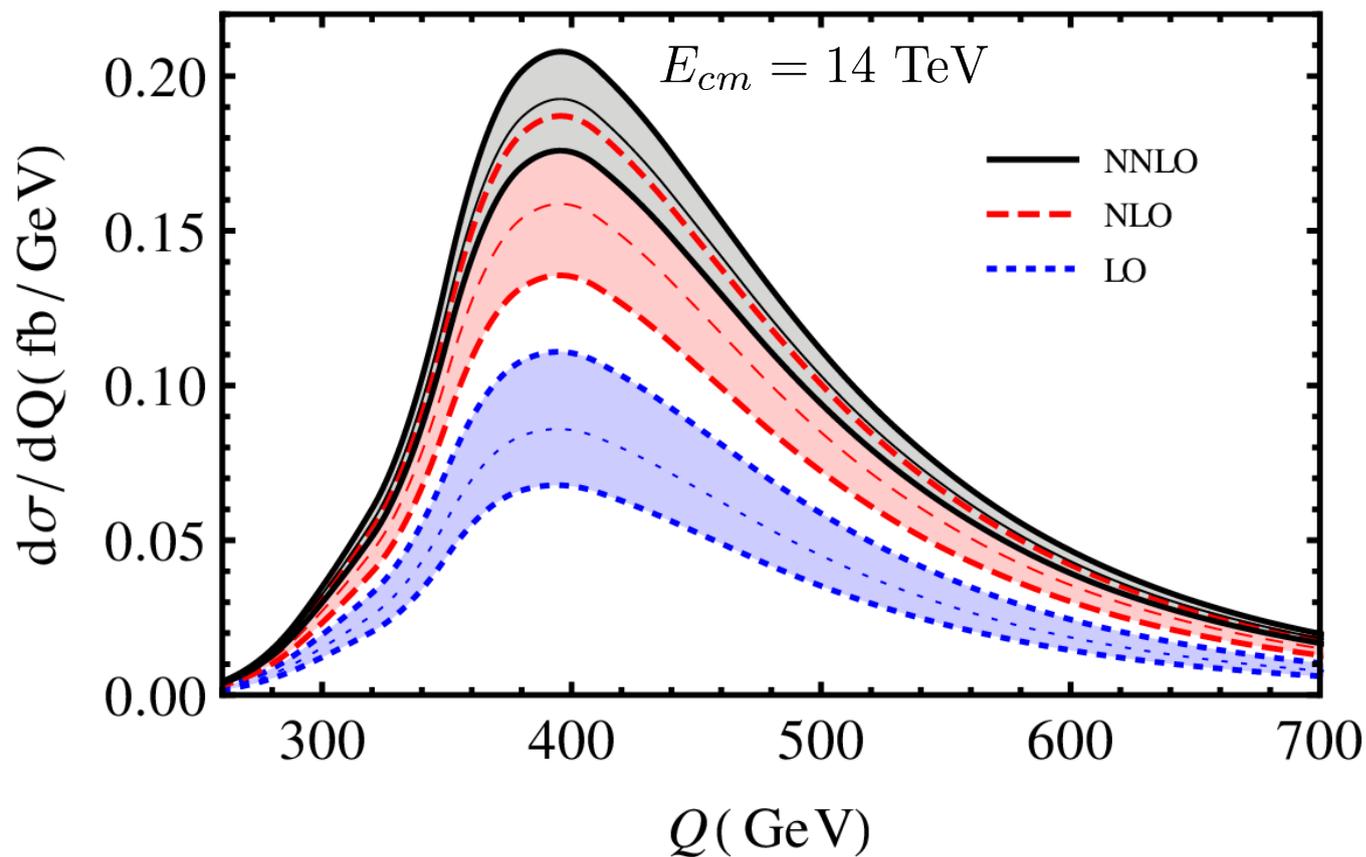
- The previous estimation is for $E_{\text{cm}}=14$ TeV
- The quality of the approximation should decrease with the collider energy



- There is a (moderate) shift towards higher invariant masses as we increase E_{cm}

Numerical results for the LHC

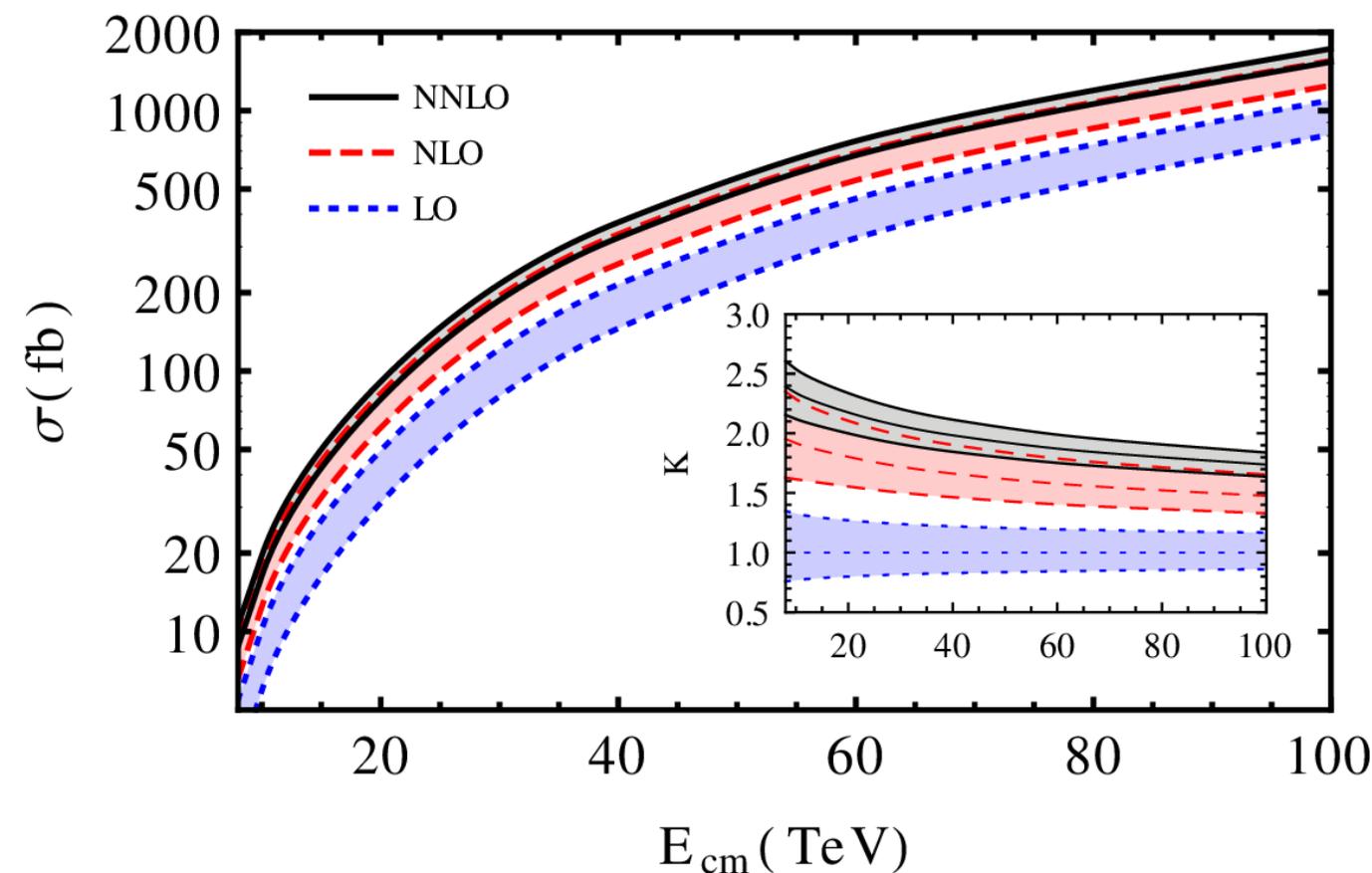
- Central value $\mu_F = \mu_R = Q$ (Higgs pair inv. Mass), and $M_H = 126$ GeV
- Scale variation: $0.5Q \leq \mu_F, \mu_R \leq 2Q$ with the constraint $0.5 \leq \mu_F/\mu_R \leq 2$
- MSTW08 parton distributions and QCD coupling
- Exact LO top and bottom-mass dependence normalization



- Overlap between NLO and NNLO bands
- K factor approx. constant
 $K_{\text{NNLO}} = 2.3$
 $\sim 20\%$ increase w.r.t. NLO
- $\pm 8\%$ scale uncertainty
 (2 times smaller than NLO)

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- Overlap between NLO and NNLO bands
- K factor smaller as the center of mass energy increases
- Substantial reduction of the scale uncertainty

Numerical results for the LHC

- To estimate the size of the missing higher orders  Scale variation
- Uncertainties in the PDFs and QCD coupling determination  90% C.L. MSTW08 sets

E_{cm}	8 TeV	14 TeV	33 TeV	100 TeV
σ_{NNLO}	9.76 fb	40.2 fb	243 fb	1638 fb
Scale [%]	+9.0 – 9.8	+8.0 – 8.7	+7.0 – 7.4	+5.9 – 5.8
PDF [%]	+6.0 – 6.1	+4.0 – 4.0	+2.5 – 2.6	+2.3 – 2.6
PDF+ α_S [%]	+9.3 – 8.8	+7.2 – 7.1	+6.0 – 6.0	+5.8 – 6.0

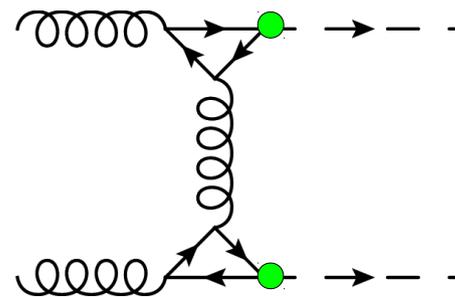
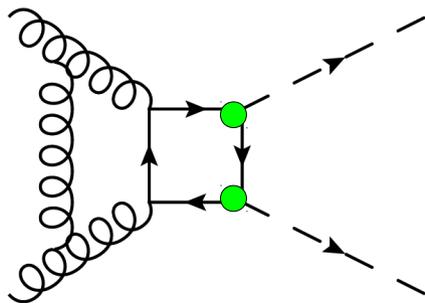
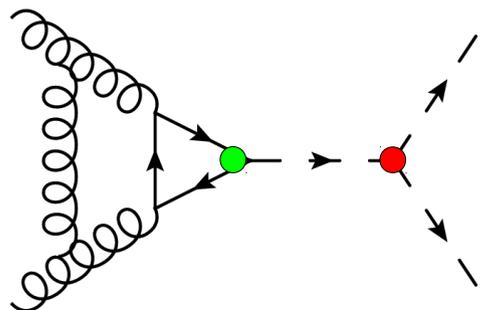
- Total scale uncertainty drops from 19% (8TeV) to 12% (100TeV)
- Also non-perturbative uncertainties drop as E_{cm} increases
- Pert. and non pert. unc. are of the same order
- Uncertainty coming from the use of the EFT is not included!

Numerical results for the LHC

- How does the K factor behave in near-SM scenarii?

• Trilinear coupling λ

• Top yukawa coupling y_t



EFT:

Single Higgs like

"New" contribution

- In principle not trivial because of the different contributions
- We found that dependence is almost negligible for a large range of the parameters

$$\lambda/\lambda_{SM} \in [-10; 20]$$

$$y_t/y_{t,SM} \in [0.5; 2.5]$$

Numerical results for the LHC

- **Soft-virtual approximation:**

Consists in considering only soft radiation in the real corrections



Keep only the plus distributions and delta functions (calculation is much simpler)

- **Why should it work?**

Gluon PDFs prefer low values of x  Threshold production

In the HH production case:

- Overestimates the total NNLO result by less than a 2%
- Better than in the single Higgs case (larger inv. Mass, closer to the threshold)

- Next step in the EFT: N3LO cross section in the SV approx

Conclusions

- We computed the NNLO corrections to the inclusive Higgs pair production XS for the gluon fusion production channel.
- We worked within the large top-mass approximation (normalizing our results with the exact LO).
- We found an increase of $\sim 20\%$ w.r.t. the NLO prediction ($E_{\text{cm}}=14\text{TeV}$).
- The scale uncertainty is substantially reduced ($\sim 17\%$), and we found an overlap with the previous order result.
- Finite top-mass effects still have to be analyzed at NLO.
- If at NLO the effect on the total XS is of $O(10\%)$ at 14 TeV, at NNLO we can expect them to be of $O(5\%)$

Thanks!