# Interacting relativistic quantum dynamics for multi-time wave functions 

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## Idea of a multi-time wave function

Setting: $N$ particles, coordinates
$x_{k}=\left(t_{k}, \mathbf{x}_{k}\right) \in \mathbb{R}^{1+d}$


Non-relativistic (single-time) Schrödinger picture: $\varphi\left(\mathrm{x}_{1}, \ldots, \mathrm{x}_{\mathrm{N}} ; t\right)$

Multi-time Schrödinger picture: $\psi\left(t_{1}, \mathbf{x}_{1}, \ldots, t_{N}, \mathbf{x}_{N}\right)$

Relation: $\varphi\left(\mathbf{x}_{1}, \ldots, \mathbf{x}_{N} ; t\right)=\psi\left(t, \mathbf{x}_{1}, \ldots, t, \mathbf{x}_{N}\right)$

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## Hamiltonian multi-time equations

Evolution equations:

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\begin{aligned}
& i \frac{\partial}{\partial t_{1}} \psi=H_{1} \psi \\
& \vdots \\
& i \frac{\partial}{\partial t_{N}} \psi=H_{N} \psi
\end{aligned}
$$

## Consistency condition:



## No-go theorem (Petrat/Tumulka 2015)

## Interaction potentials excluded, i.e. if


then this has to be gauge-equivalent to $V_{i} \equiv V_{i}\left(x_{i}\right)$.

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## Interaction via boundary conditions: a $(1+1)$-dim. model

Assumptions: $\quad N=2, \quad d=1, \quad m_{1}=m_{2}=0, \quad \psi=\left(\psi_{1}, \psi_{2}, \psi_{3}, \psi_{4}\right)$

Multi-time Dirac equations on $\Omega$ :

$$
\begin{aligned}
i \frac{\partial}{\partial t_{1}} \psi\left(t_{1}, z_{1}, t_{2}, z_{2}\right) & =-i \sigma_{3} \otimes 1 \frac{\partial}{\partial z_{1}} \psi\left(t_{1}, z_{1}, t_{2}, z_{2}\right) \\
i \frac{\partial}{\partial t_{2}} \psi\left(t_{1}, z_{1}, t_{2}, z_{2}\right) & =-i 1 \otimes \sigma_{3} \frac{\partial}{\partial z_{2}} \psi\left(t_{1}, z_{1}, t_{2}, z_{2}\right)
\end{aligned}
$$

Initial conditions at $t_{1}=t_{2}=0$

Boundary conditions at $\mathcal{C}=\left\{\left(t_{1}, z_{1}, t_{2}, z_{2}\right): t_{1}=t_{2}, z_{1}=z_{2}\right\}$

## Basic idea: multi-time characteristics

Write out the two-time system in matrix-vector form:

$$
\begin{aligned}
i \frac{\partial}{\partial t_{1}}\left(\begin{array}{l}
\psi_{1} \\
\psi_{2} \\
\psi_{3} \\
\psi_{4}
\end{array}\right) & =-i\left(\begin{array}{cccc}
1 & & & \\
& 1 & & \\
& & -1 & \\
& & & -1
\end{array}\right) \frac{\partial}{\partial z_{1}}\left(\begin{array}{l}
\psi_{1} \\
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\text { E.g. for } \psi_{1} \text { : } \quad\left(\frac{\partial}{\partial t_{1}}+\frac{\partial}{\partial z_{1}}\right) \psi_{1}=0, \quad\left(\frac{\partial}{\partial t_{2}}+\frac{\partial}{\partial z_{2}}\right) \psi_{1}=0
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E.g. for $\psi_{1}$ : $\quad\left(\frac{\partial}{\partial t_{1}}+\frac{\partial}{\partial z_{1}}\right) \psi_{1}=0, \quad\left(\frac{\partial}{\partial t_{2}}+\frac{\partial}{\partial z_{2}}\right) \psi_{1}=0$

$$
\Rightarrow \quad \psi_{1}\left(t_{1}, z_{1}, t_{2}, z_{2}\right)=f_{1}\left(z_{1}-t_{1}, z_{2}-t_{2}\right) .
$$

## Existence and uniqueness result

## Theorem

There exists a unique $C^{k}$ solution of the following initial boundary value problem on $\Omega_{1}$ (i.e. $\Omega$ with $z_{1}<z_{2}$ ):

$$
\begin{aligned}
\psi_{i}\left(0, z_{1}, 0, z_{2}\right) & \stackrel{!}{=} g_{i}\left(z_{1}, z_{2}\right), \quad z_{1}<z_{2}, \quad i=1,2,3,4 \\
\psi_{2}(t, z-0, t, z+0) & \stackrel{!}{=} h_{2}(t, z), \quad t<0, \quad z \in \mathbb{R} \\
\psi_{3}(t, z-0, t, z+0) & \stackrel{!}{=} h_{3}(t, z), \quad t>0, \quad z \in \mathbb{R}
\end{aligned}
$$

where $g_{i}, h_{j}$ are $C^{k}$ functions and compatible, i.e. the transitions between initial values and boundary conditions is also $C^{k}$.

## Relativistic probability conservation

Conserved tensor current: $\quad j^{\mu \nu}\left(x_{1}, x_{2}\right)=\bar{\psi}\left(x_{1}, x_{2}\right) \gamma^{\mu} \otimes \gamma^{\nu} \psi\left(x_{1}, x_{2}\right)$
$d=1: \quad \gamma^{0}=\sigma_{1}, \gamma^{1}=\sigma_{1} \sigma_{3}, \quad \partial_{1, \mu} j^{\mu \nu}=\partial_{2, \nu} j^{\mu \nu}=0$,

## Probability conservation on space-like hypersurfaces $\Sigma$ :

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\int_{(\Sigma \times \Sigma) \cap \Omega} d \sigma_{\mu}\left(x_{1}\right) d \sigma_{\nu}\left(x_{2}\right) j^{\mu \mu}\left(x_{1}, x_{2}\right)=1 \forall \Sigma
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Theorem
Probability conservation as well as existence and uniqueness are ensured for boundary conditions


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\psi_{2}(t, z-0, t, z+0) \stackrel{!}{=} e^{-i \theta} \psi_{3}(t, z-0, t, z+0) \forall t, z
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## Lorentz invariance

Apart from the boundary conditions, LI is manifest. Representation of the proper Lorentz group $\mathcal{L}_{+}^{\uparrow}$ in $d=1$ :

- Just one generator (boosts in $z$-direction).

- $S_{i}[\Lambda]=\exp \left(\frac{\beta}{2} \gamma_{i}^{0} \gamma_{i}^{1}\right), \quad \beta \in \mathbb{R}$
- $\psi_{1}\left(x_{1}, x_{2}\right) \stackrel{\wedge}{\longmapsto}\left(\cosh ^{2} \beta+2 \cosh \beta \sinh \beta+\sinh ^{2} \beta\right) \psi_{1}\left(\Lambda^{-1} x_{1}, \Lambda^{-1} x_{2}\right)$ $\psi_{2}\left(x_{1}, x_{2}\right) \stackrel{\wedge}{\longmapsto} \psi_{2}\left(\Lambda^{-1} x_{1}, \Lambda^{-1} x_{2}\right)$


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## Criterion <br> A model is interacting if there exist initial product wave functions that become entangled during time evolution.

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## Questions?



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