

# Radar detection of high-energy neutrino induced particle cascades in ice

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Vrije  
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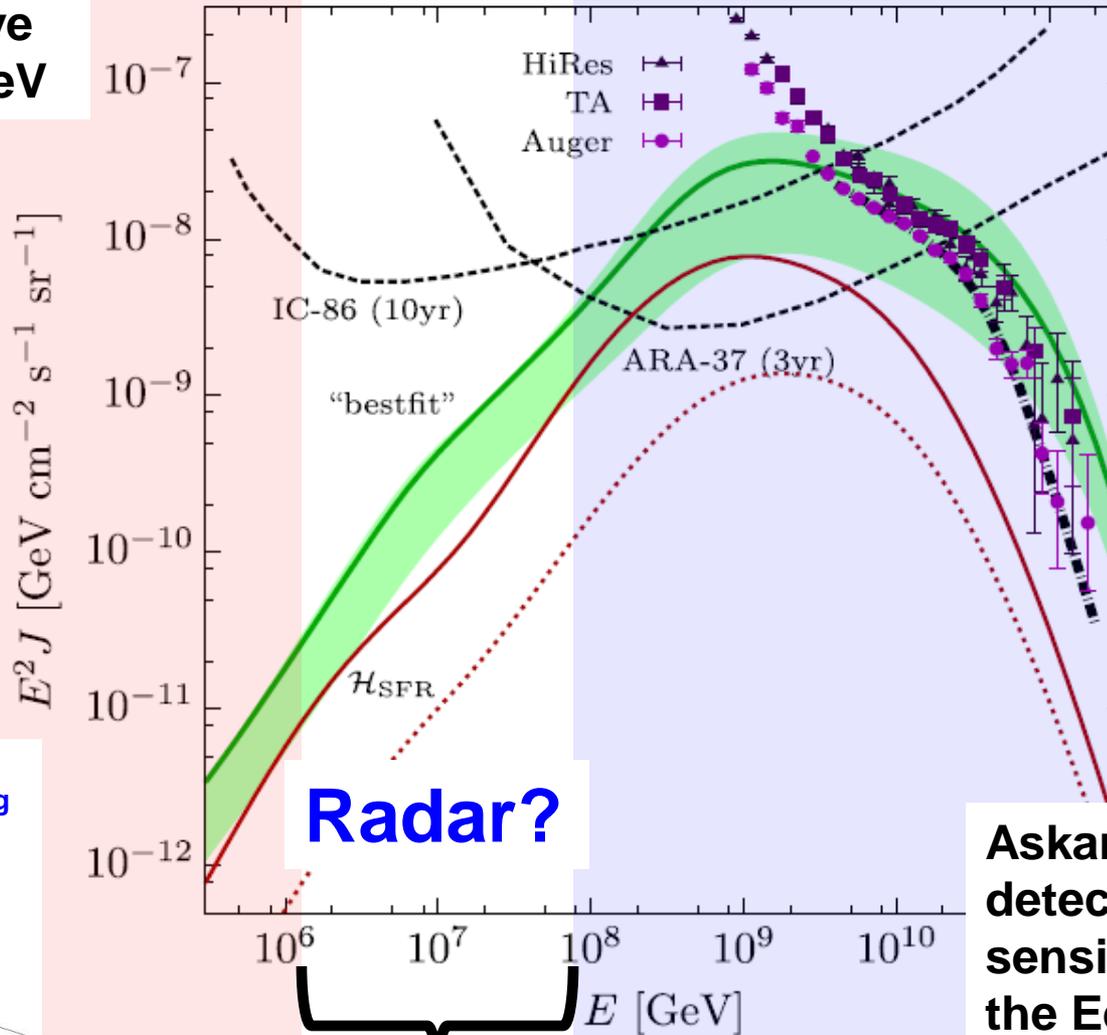
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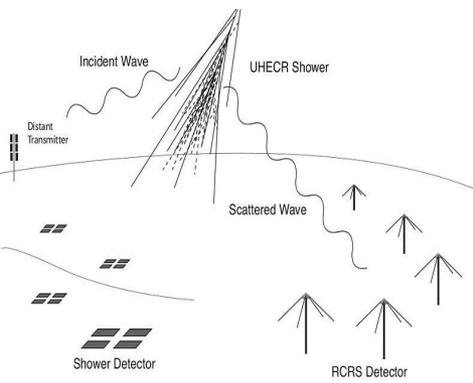


# Motivation

IceCube sensitive below several PeV



M. Abou Bakr Othman et al,  
Proceedings 32nd ICRC, Beijing  
2011



Sensitivity Gap in  
PeV – EeV region

# Radar scattering of a neutrino induced plasma

Leftover electrons from ionization:  
Extension:  $O(30 \text{ cm})$   
Lifetime:  $O(1-20 \text{ ns})$

Shower front electrons:  
Extension:  $R_L = O(10 \text{ cm})$   
Lifetime:  $O(100 \text{ ns})$   
Moving!

Leftover protons from ionization:  
Wide extension:  $O(5 \text{ m})$   
Lifetime:  $O(10-1000 \text{ ns})$

Ionization numbers come from Physical Chemistry research!

Figure from arXiv:1210.5140v2

6. Laws, J. O. & Parsons, D. A. *EOS* 24, 452-460 (1943)

## Proton mobility in ice

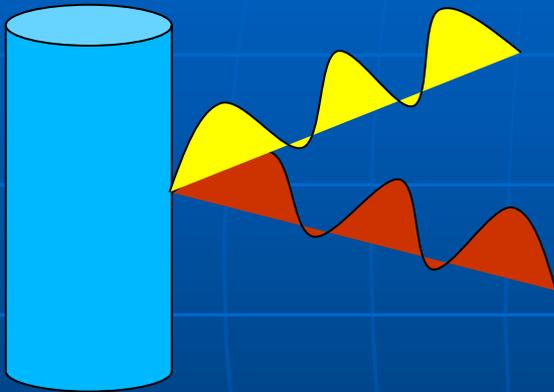
Marinus Kunst & John M. Warman

Interuniversitair Reactor Instituut, Mekelweg 15, 2629 JB Delft, The Netherlands

Ice is frequently taken as a model when factors controlling proton transport in hydrogen-bonded molecular networks are discussed. Such discussions have increased with the acknowledgement that proton transfer across cell membranes may play a significant part in energy conversion and storage in biological systems<sup>1-4</sup> and that this transfer may involve hydrogen-bonded chains spanning the membrane<sup>5,6</sup>. However, there is still much

# RADAR scattering

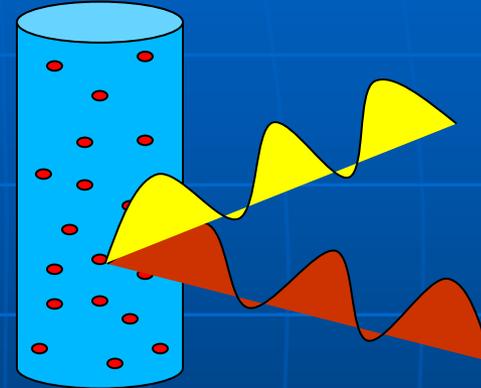
- Over-dense scattering:



Radar frequency  $<$  Plasma Frequency

Reflection from the surface of the plasma tube

- Under-dense scattering:



Radar frequency  $>$  Plasma Frequency

Scattering off of the individual charges in the plasma

# RADAR return power estimation

## Bi-static RADAR configuration

Effective area of receiver:  $A_{eff}$



Transmitted power:  $P_t$



Re-scattering over a sphere:  $1/(4\pi R^2)$

Transmission over  $1/4$  of a sphere:  $1/(\pi R^2)$

Plasma scattering surface:  $\sigma_{eff}$

Attenuation by the medium

$$P_r = P_t \eta \frac{\sigma_{eff}}{\pi R^2} \frac{A_{eff}}{4\pi R^2} e^{-4R/L\alpha}$$

# RADAR return power estimation (single antenna)

$$P_r = P_t \eta \frac{\sigma_{eff}(\lambda)}{\pi R^2} \frac{A_{eff}(\lambda)}{4\pi R^2} e^{-4R/L_\alpha}$$

$$\lambda = 0.18 \text{ m}$$

$$\sigma_{eff}^{max} = 0.11 \text{ m}^2$$

$$\sigma_{eff}(\theta = 60^\circ, \phi = 60^\circ) = 1.6 \cdot 10^{-4} \text{ m}^2$$

$$L_\alpha = 1 \text{ km}$$

$$P_{noise} = k_b T_{sys} \Delta \nu$$

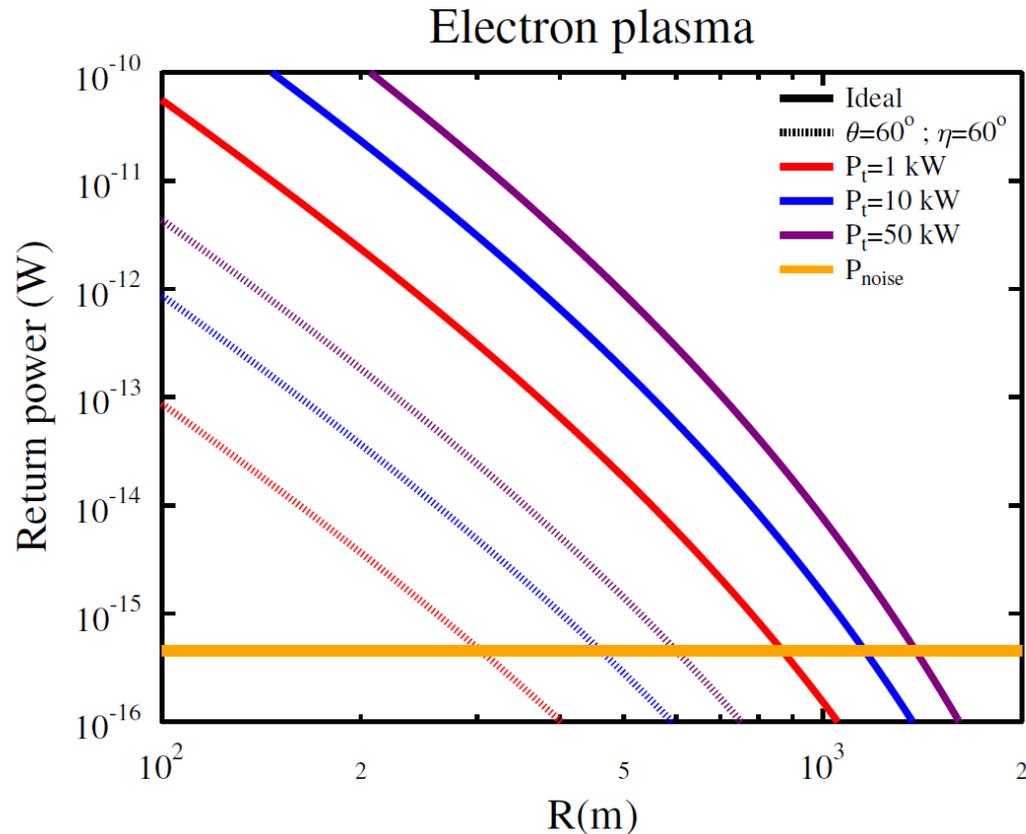
$$T_{sys} = 325 \text{ K}$$

$$\Delta \nu = 100 \text{ kHz}$$

N antennas :

$$P_{Noise}(N) = N \cdot P(N = 1)$$

$$P_{Signal}(N) = N^2 \cdot P(N = 1)$$



$E > 4 \text{ PeV}$

# RADAR return power estimation (single antenna)

$$P_r = P_t \eta \frac{\sigma_{eff}(\lambda)}{\pi R^2} \frac{A_{eff}(\lambda)}{4\pi R^2} e^{-4R/L_\alpha}$$

$$\lambda = 3.6 \text{ m}$$

$$\sigma_{eff}^{\max} = 5.5 \text{ m}^2$$

$$\sigma_{eff}(\theta = 60^\circ, \phi = 60^\circ) = 1.2 \cdot 10^{-2} \text{ m}^2$$

$$L_\alpha = 1.4 \text{ km}$$

$$P_{\text{noise}} = k_b T_{\text{sys}} \Delta \nu$$

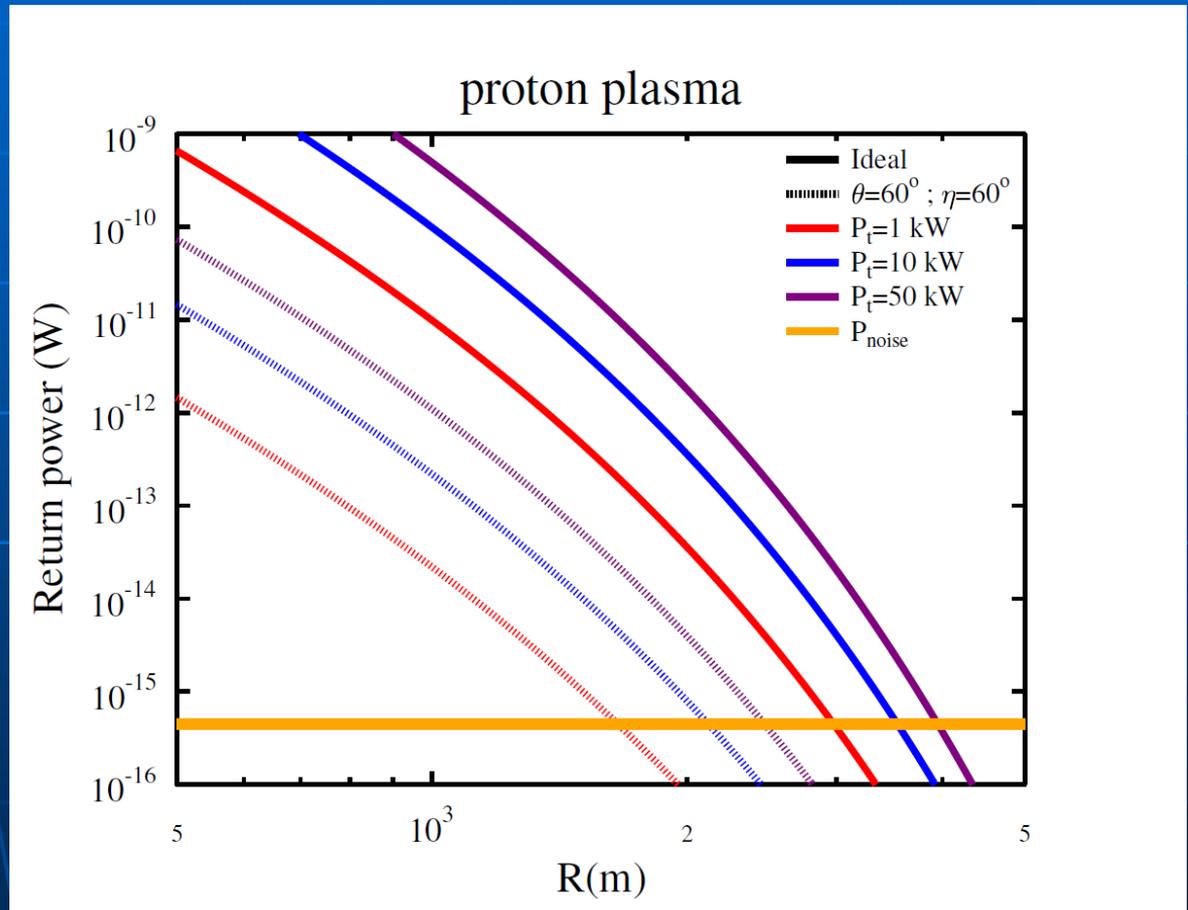
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$E > 20 \text{ PeV}$

# RADAR return power estimation (single antenna)

$$P_r = P_t \eta \frac{\sigma_{eff}(\lambda)}{\pi R^2} \frac{A_{eff}(\lambda)}{4\pi R^2} e^{-4R/L_\alpha}$$

$$\lambda = 2.6 \text{ m}$$

$$\sigma_{eff}^{max} = 5.5 \text{ m}^2$$

$$\sigma_{eff}(\theta = 60^\circ, \phi = 60^\circ) = 1.2 \cdot 10^{-2} \text{ m}^2$$

$$L_\alpha = 1.4 \text{ km}$$

$$P_{noise} = k_b T_{sys} \Delta \nu$$

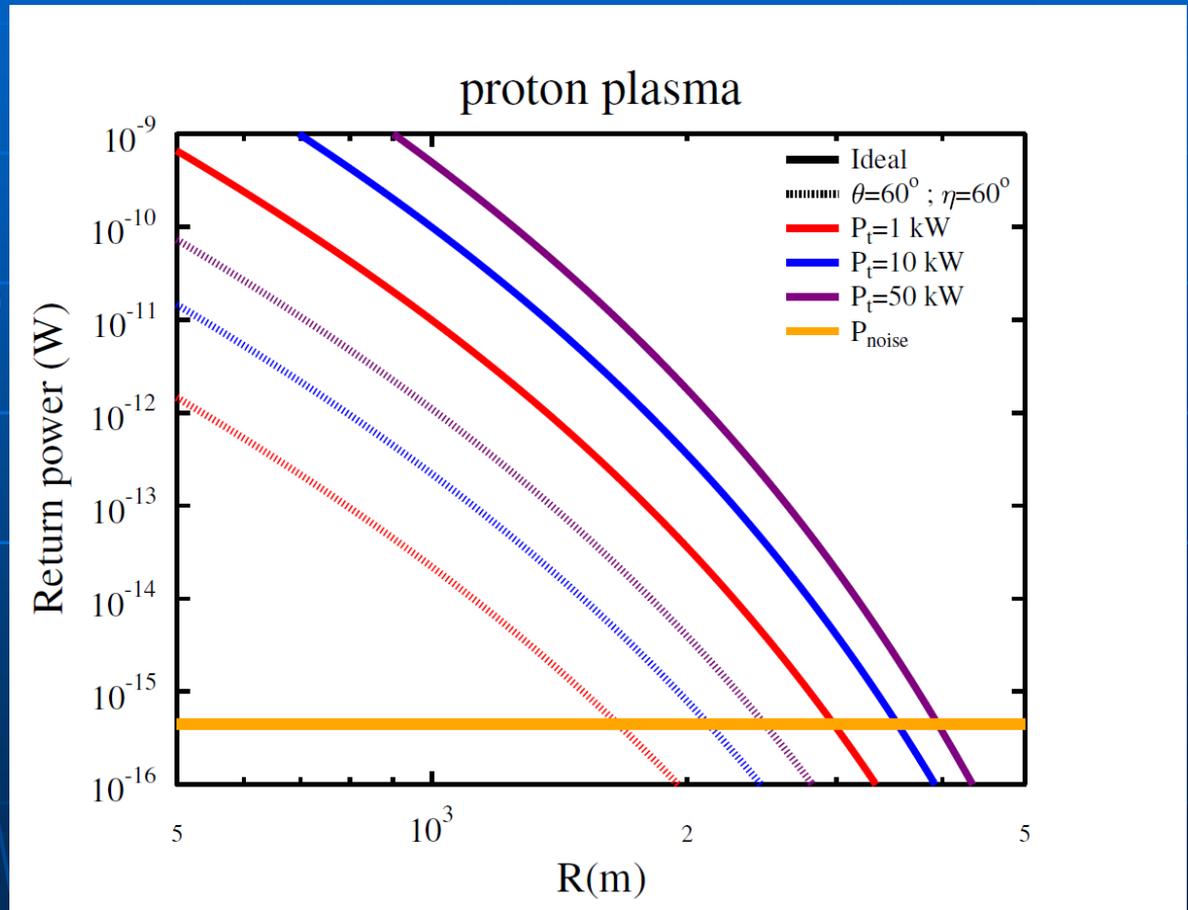
$$T_{sys} = 325 \text{ K}$$

$$\Delta \nu = 100 \text{ kHz}$$

N antennas :

$$P_{Noise}(N) = N \cdot P(N = 1)$$

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$E > 20 \text{ PeV}$

# Open questions: The Plasma

- How large is the over-dense plasma?
- What is the influence of skin-effects?
- What is the lifetime of the plasma?
- Is the plasma collision frequency low enough?

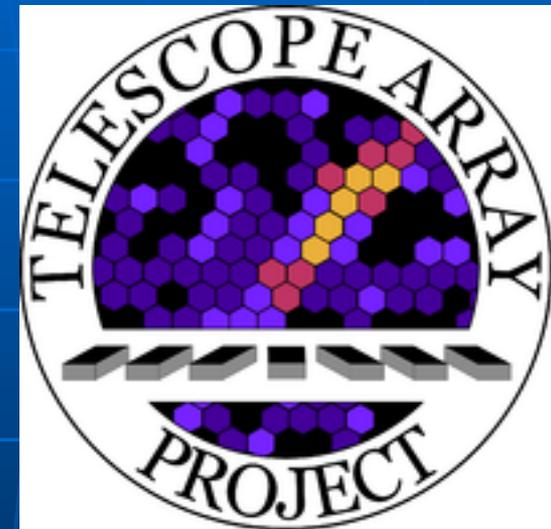


**Experimental verification  
needed!**

# Radar scattering experiment at TA-ELS

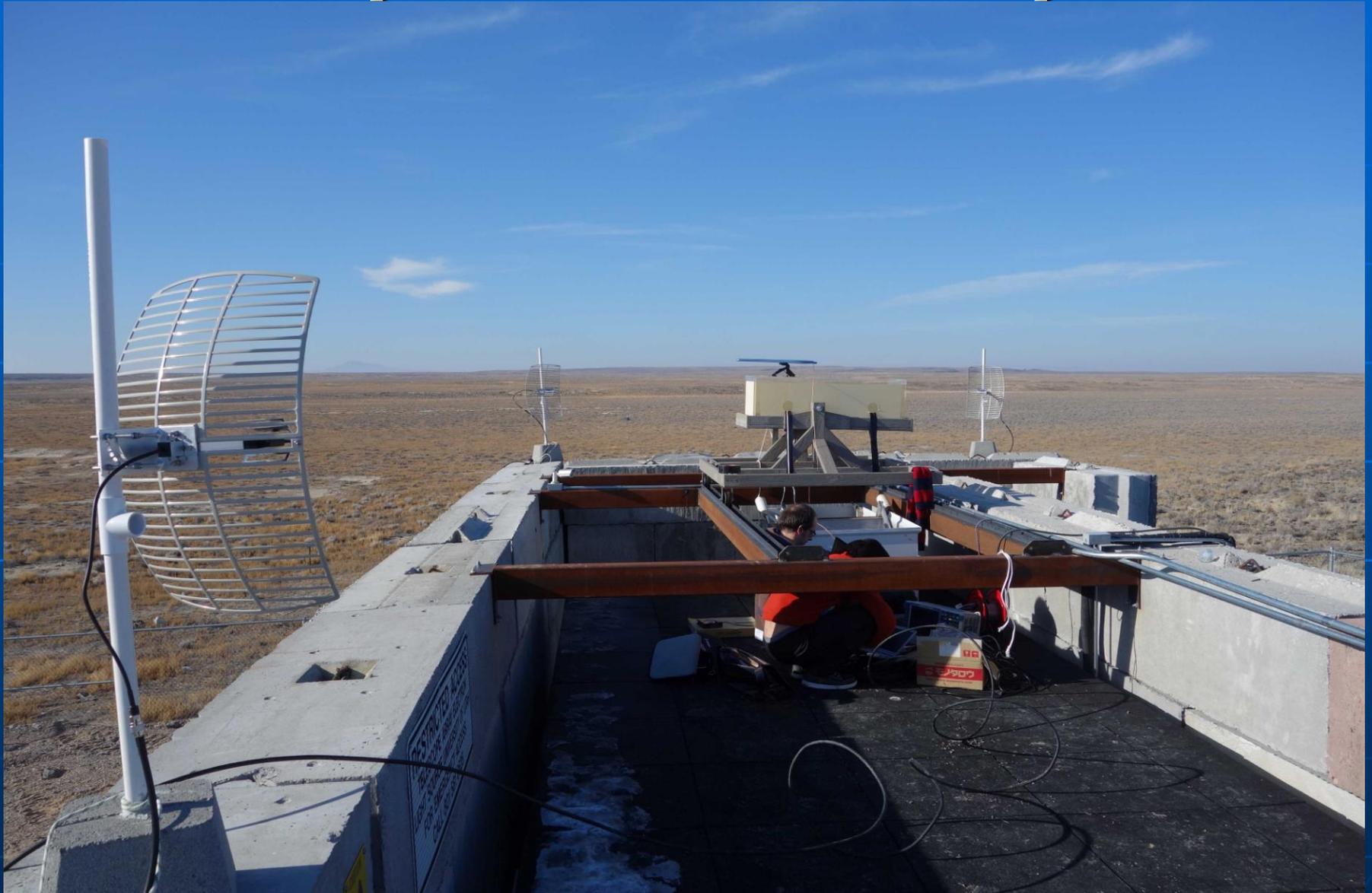


Aya  
Matt  
Kael

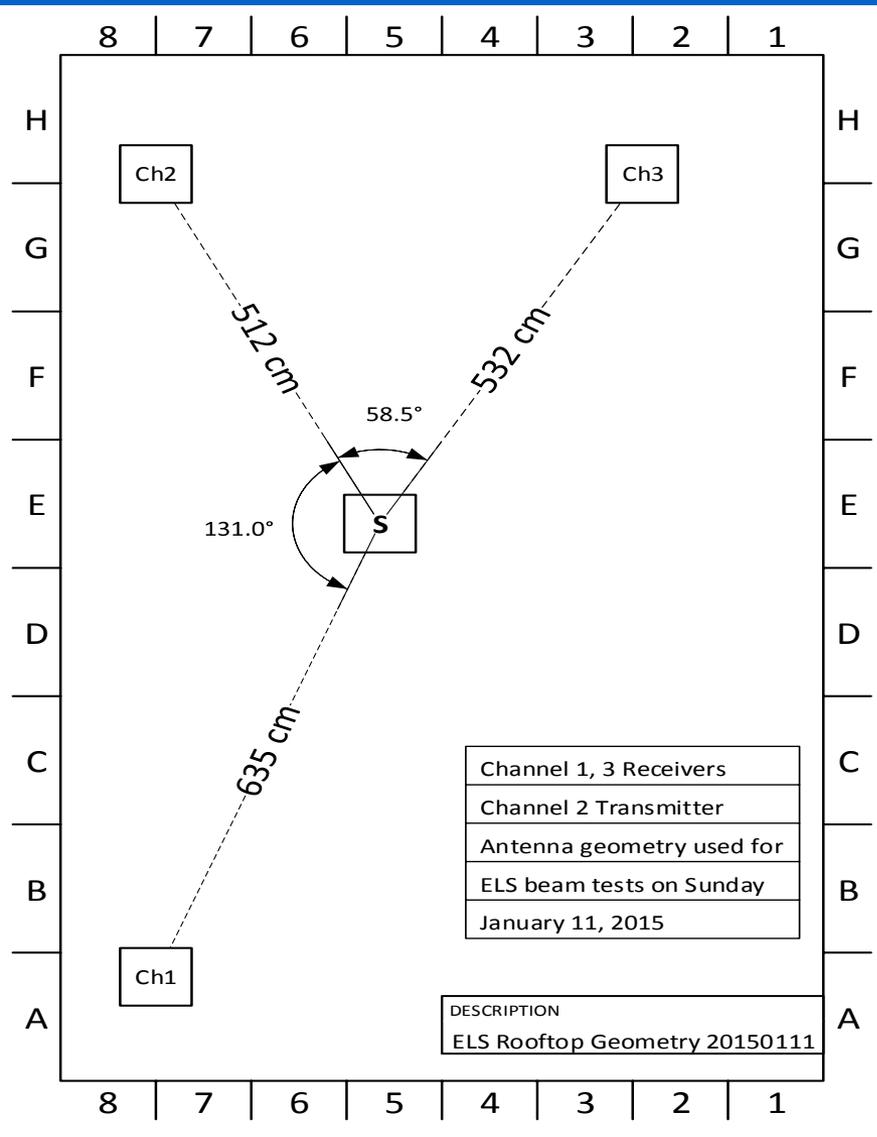


Many thanks to the **Chiba group** and  
the **Telescope Array Collaboration** !

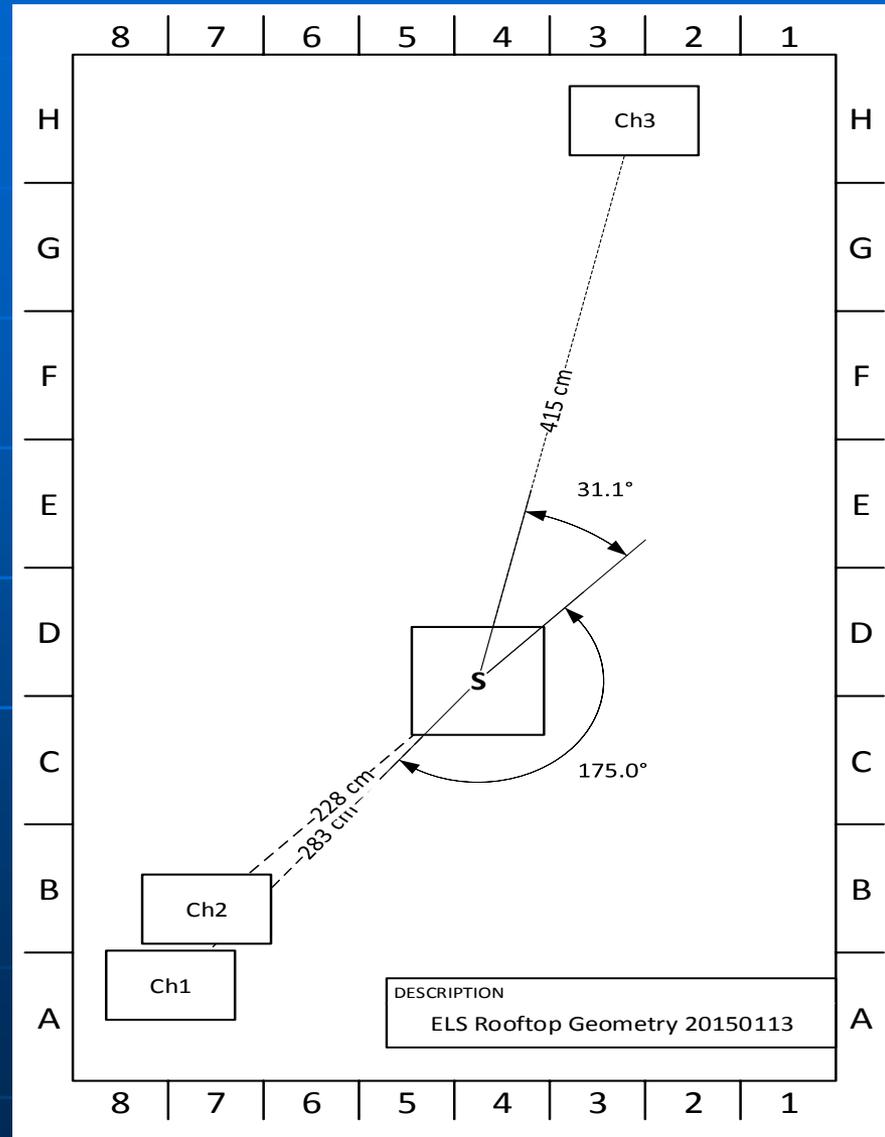
# Experimental setup



# Experimental setup

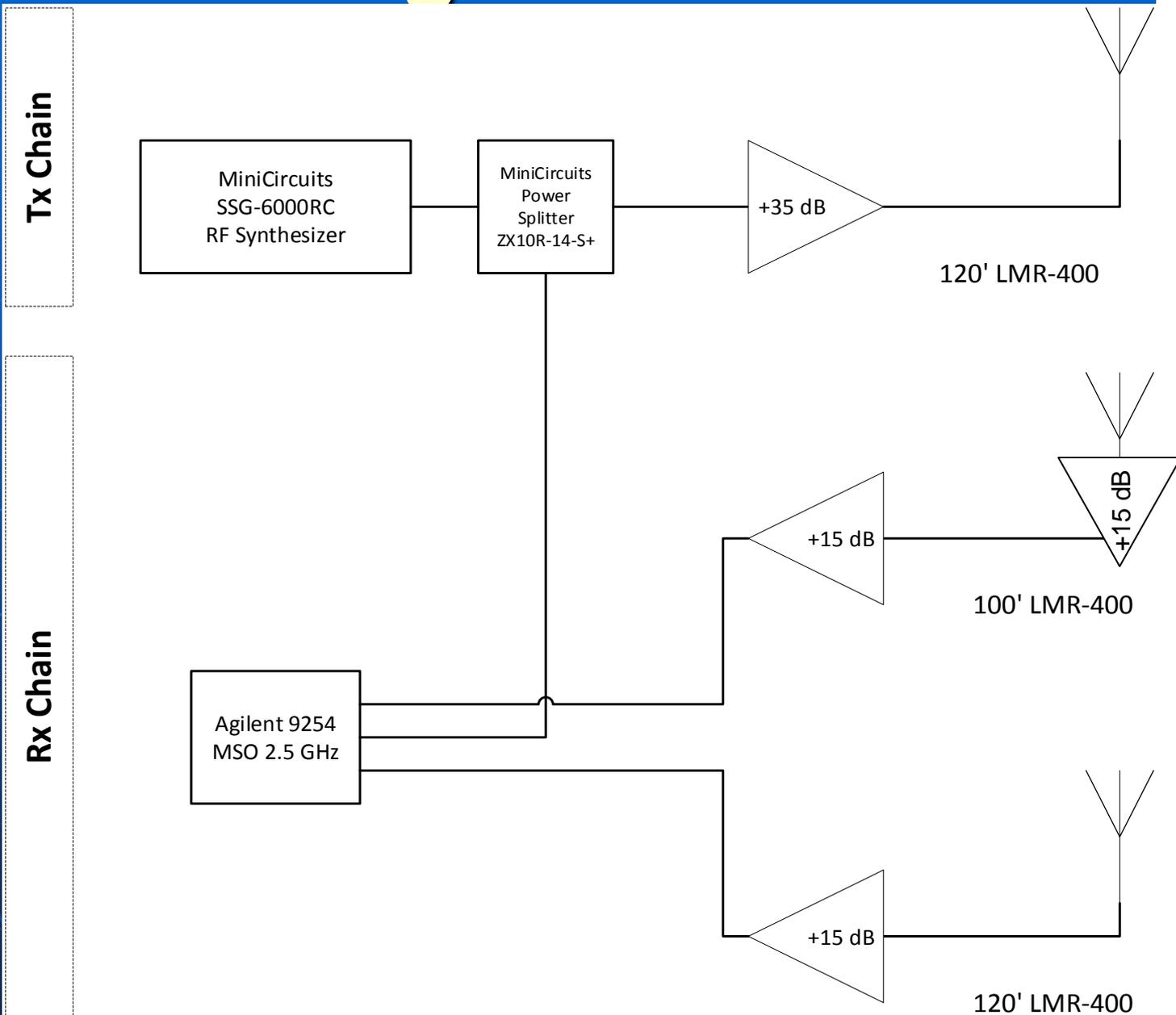


Early Configuration



Later Configuration

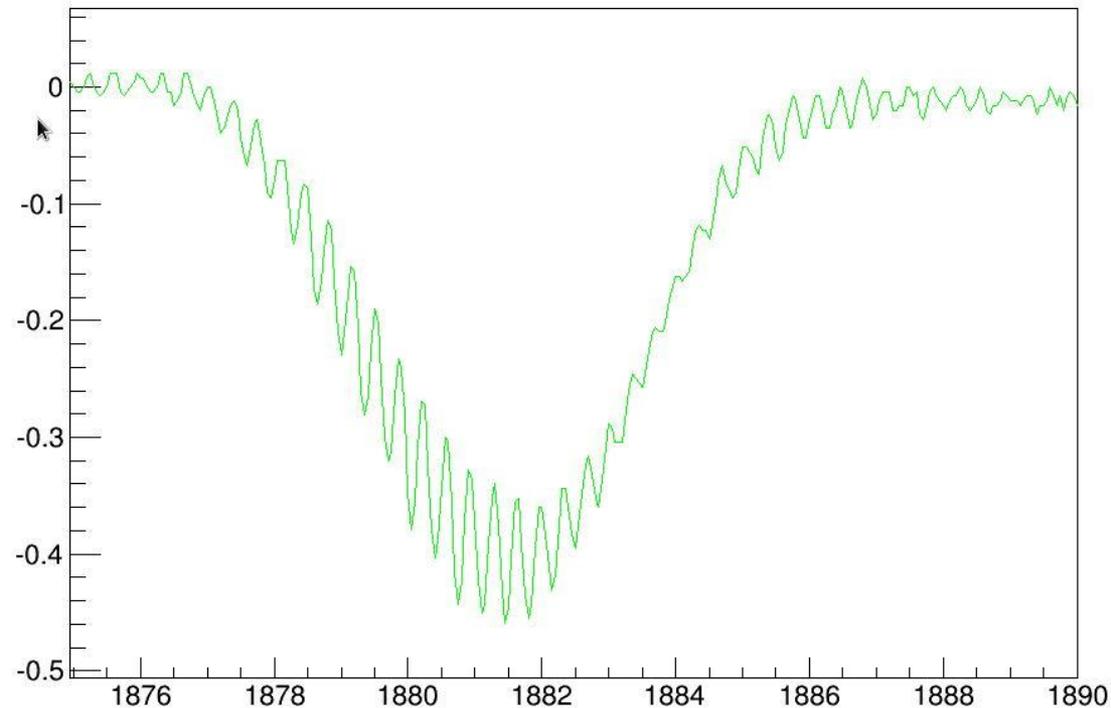
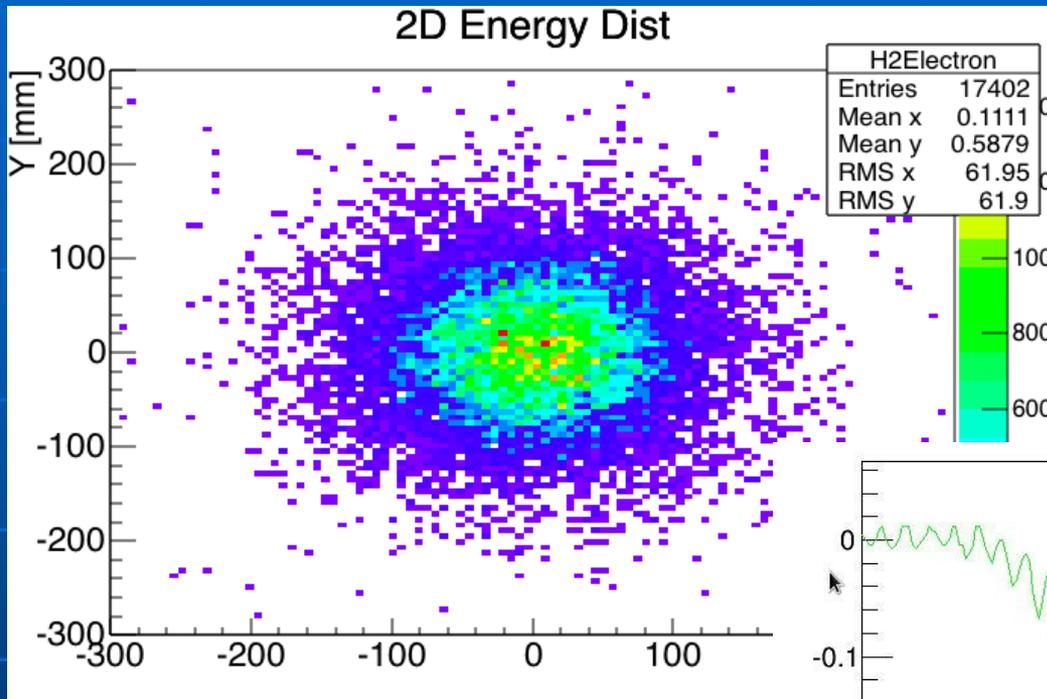
# Signal chain



# Radar scattering

## Beam characteristics

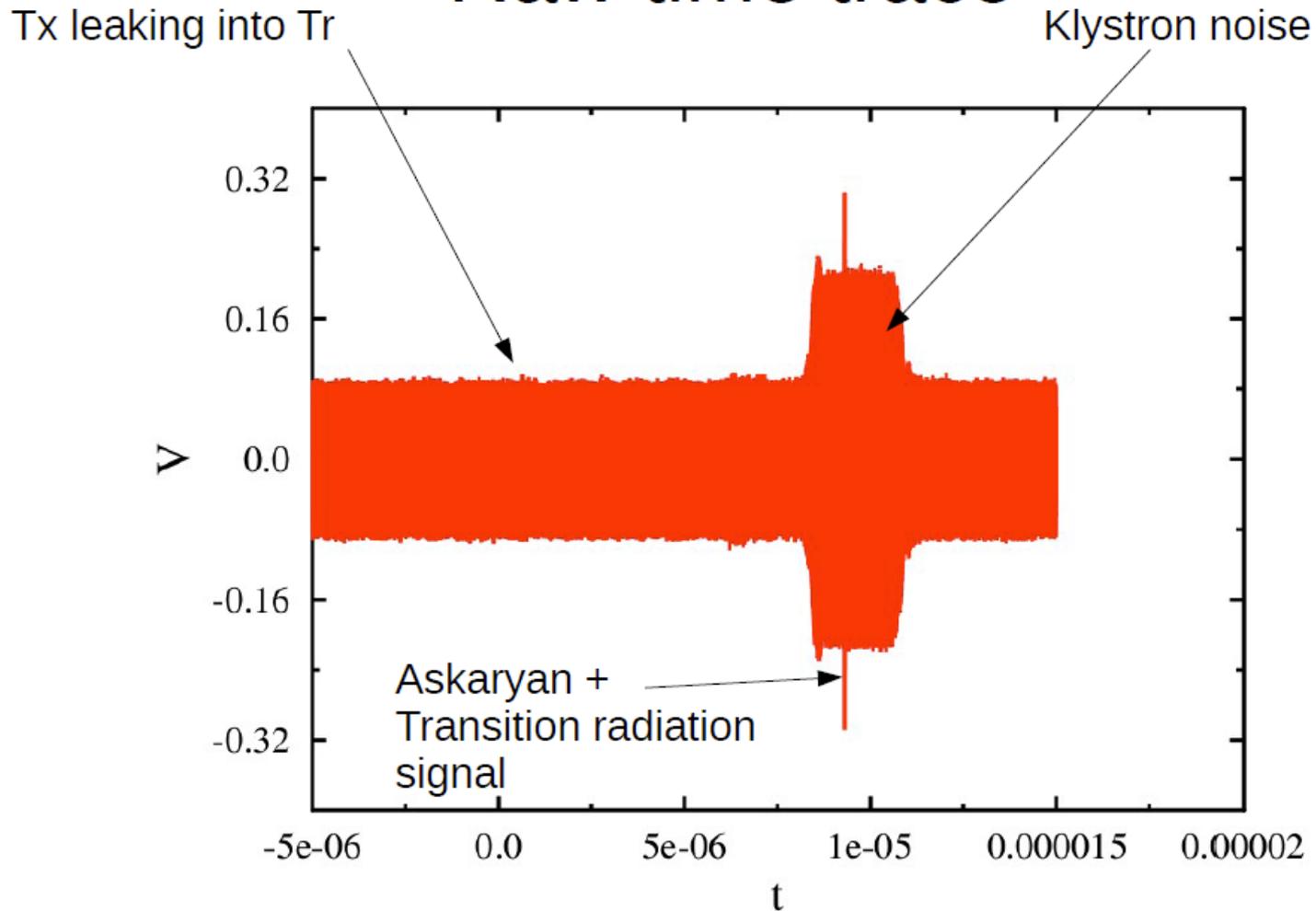
$\sim 10^9$  (40 MeV) electrons  
 $\sim 40$  PeV



# Radar scattering

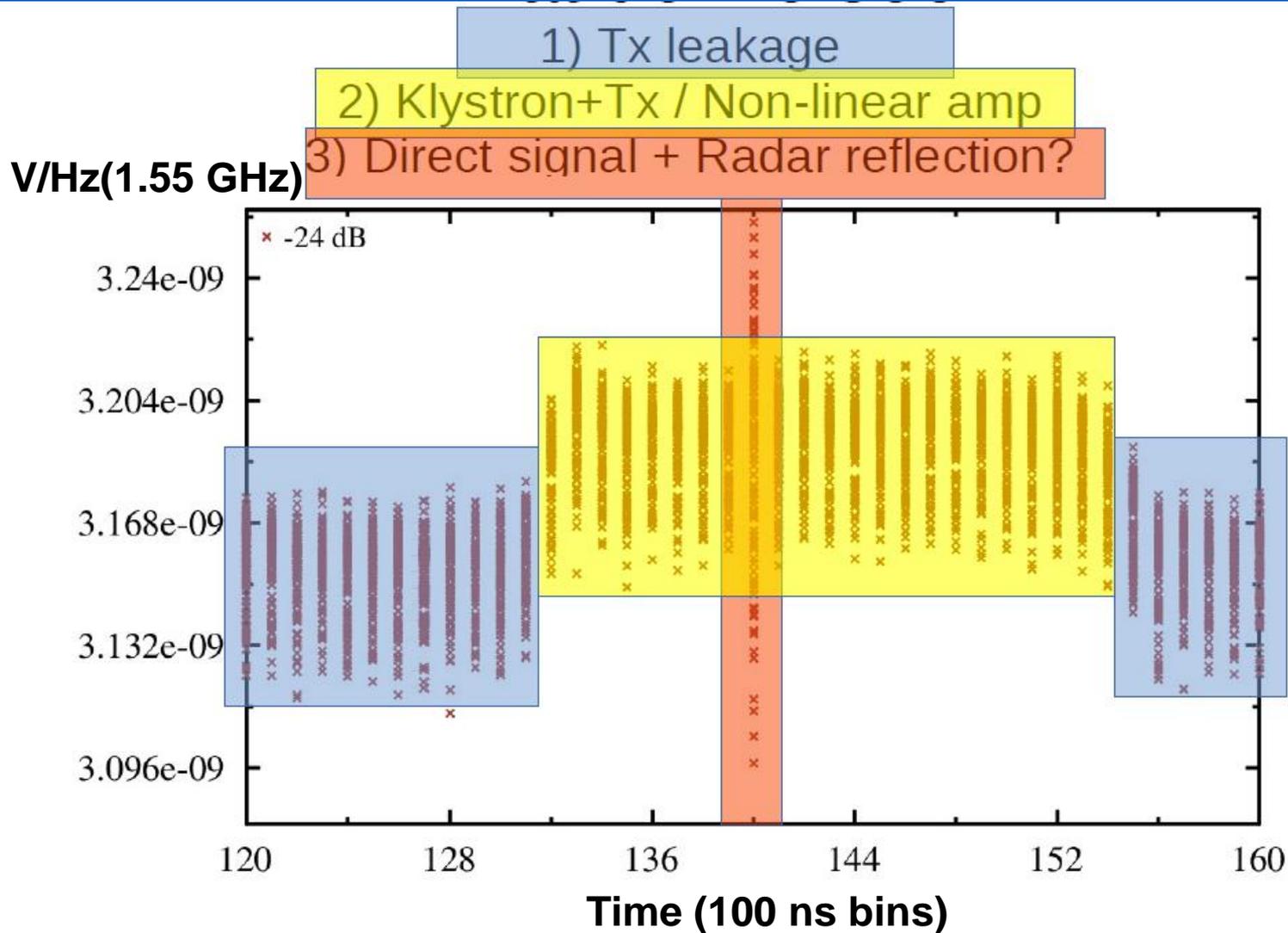
## What do we see?

### Raw time trace



# Radar scattering

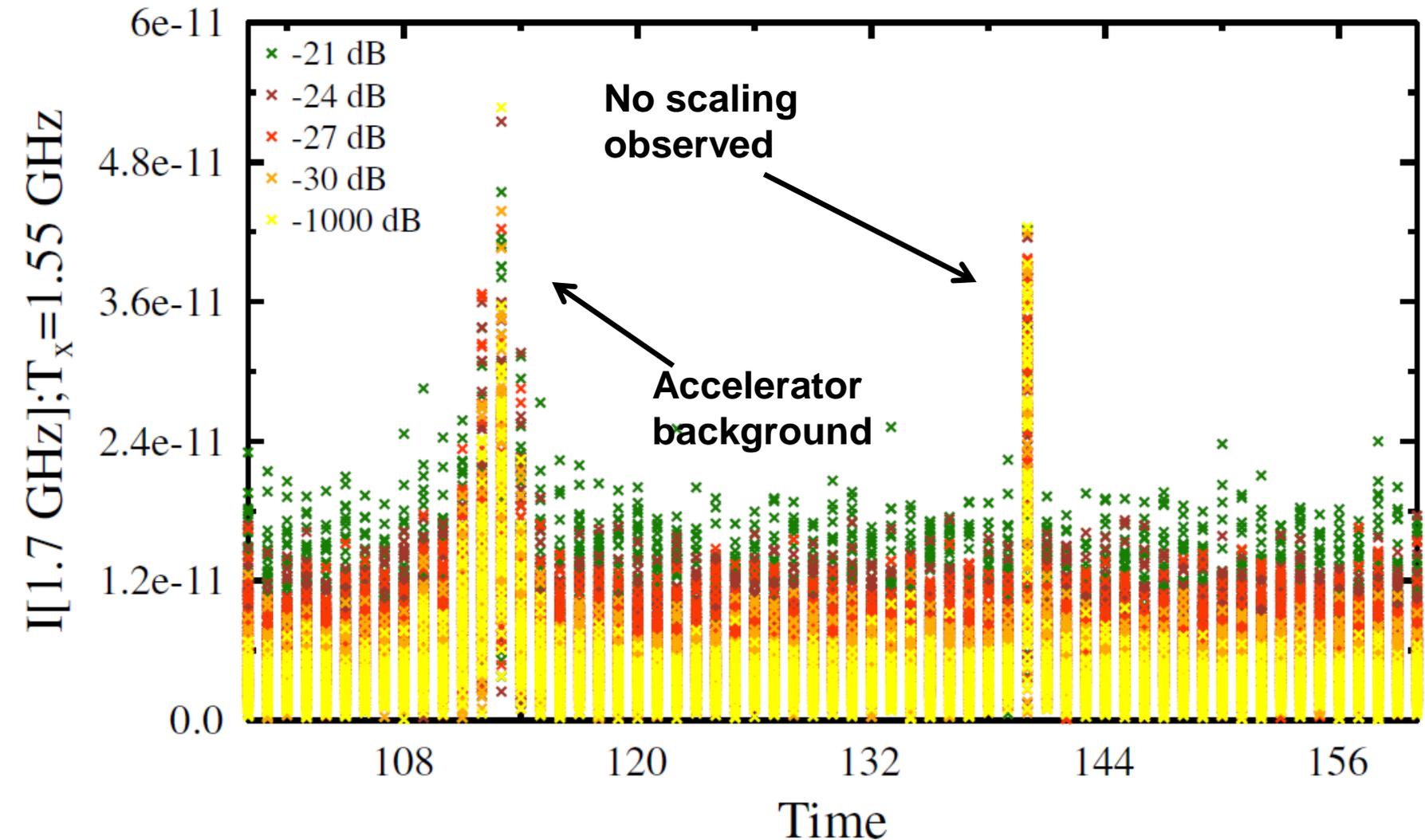
## What do we see?



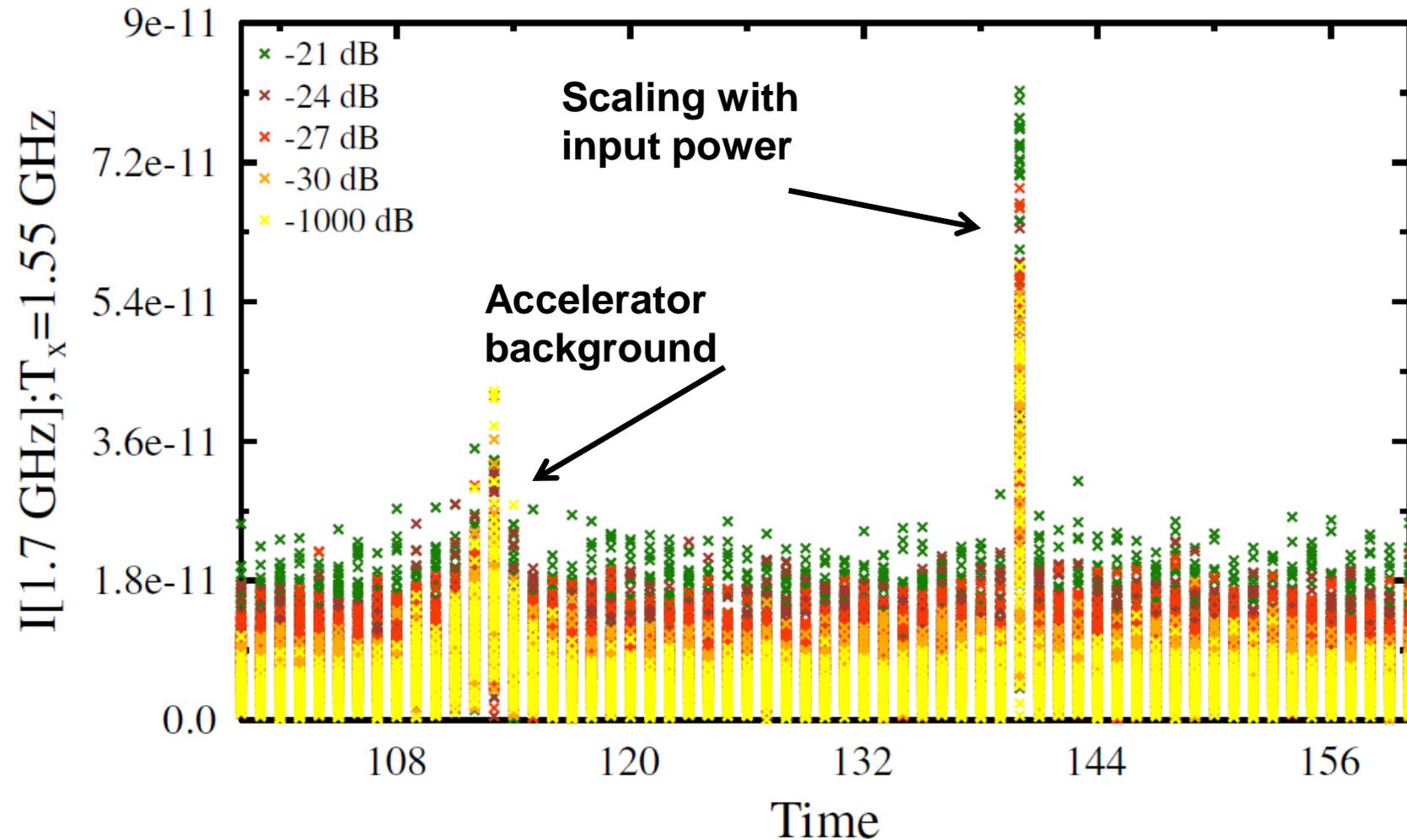
# Radar scattering Interference and instrumental effects

- Accelerator noise interferes with our transmit signal
- Non-linear amplifier response
- **Signal can be mimicked by these effects!**
- What if we look at a different frequency than our transmit frequency?

# Radar scattering Air



# Radar scattering Ice



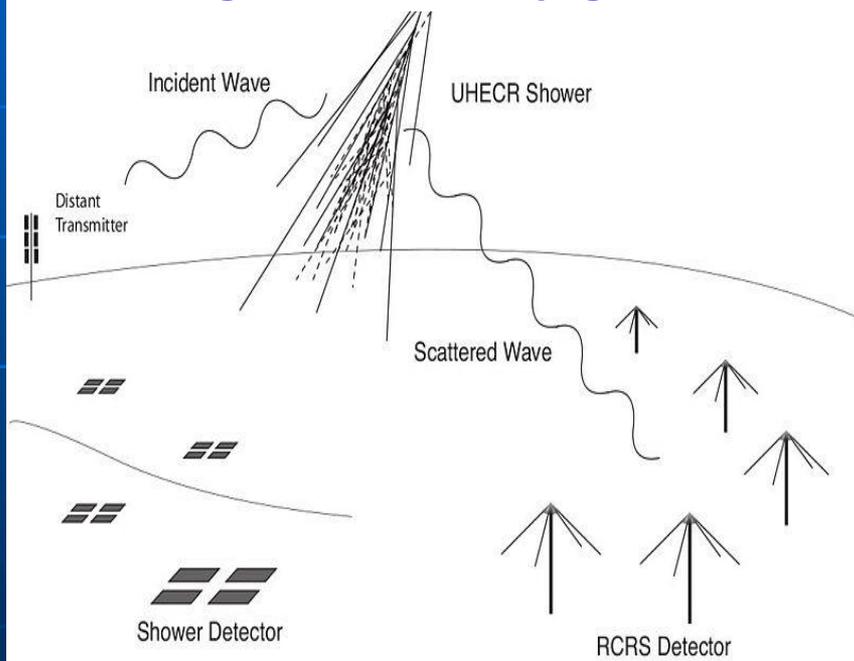
# Conclusions

- Modeling the RADAR scattering of high-energy neutrino induced cascades gives an energy threshold of **several PeV**.
- We performed a measurement to determine the feasibility of this method.
- Obtained data **hints toward a scattered signal, analysis is ongoing.**

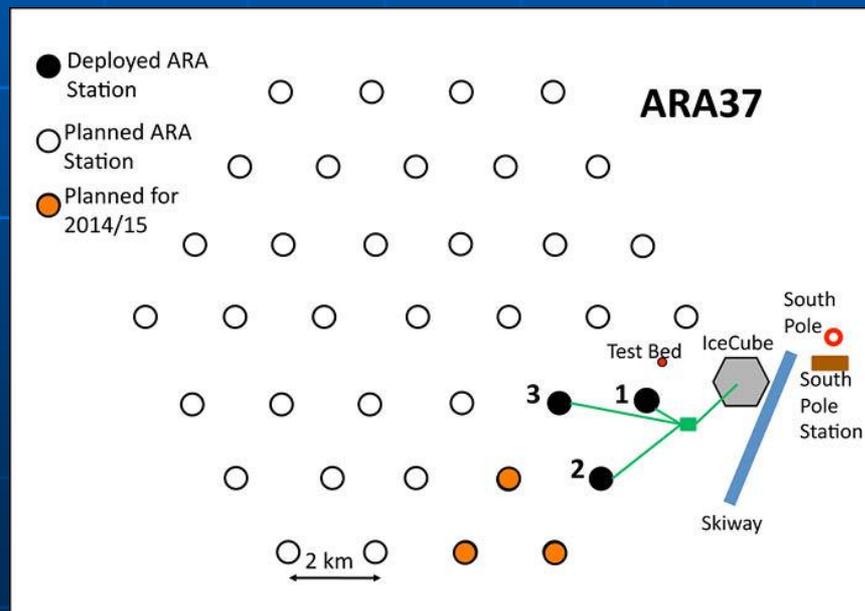
# New detection method

If a RADAR signal can be bounced off of a neutrino induced cascade in ice, we have **control over the signal strength!**

M. Abou Bakr Othman et al,  
Proceedings 32nd ICRC, Beijing 2011



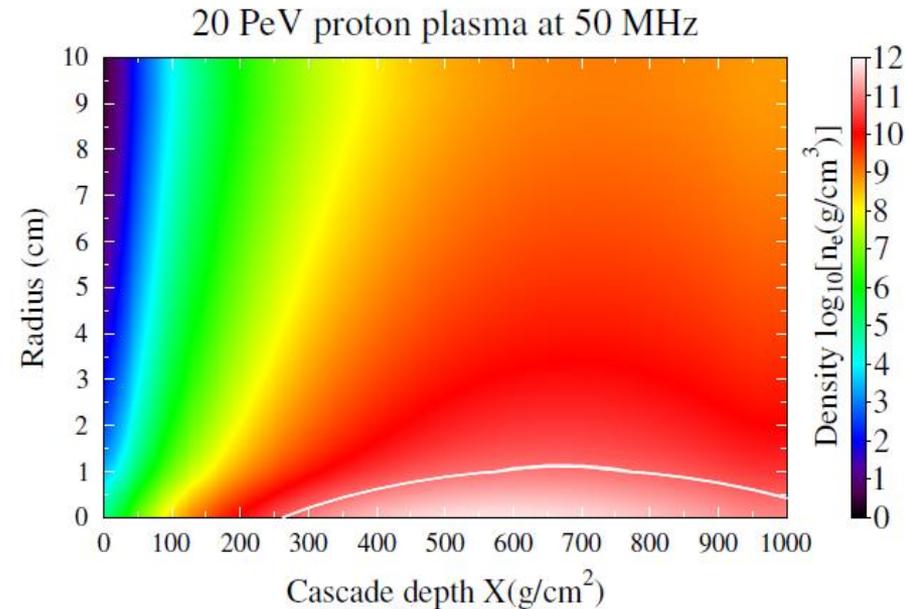
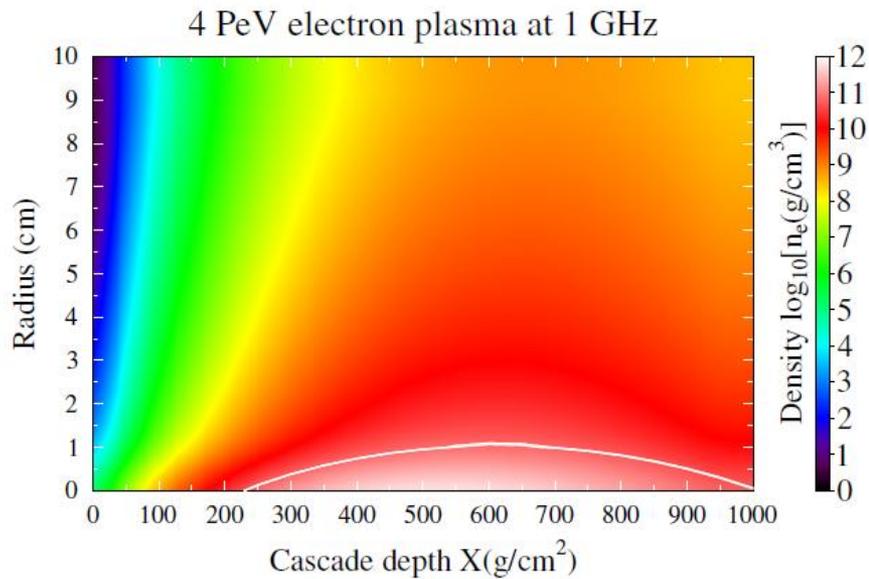
**Infrastructure already available!**



# Over-dense scattering

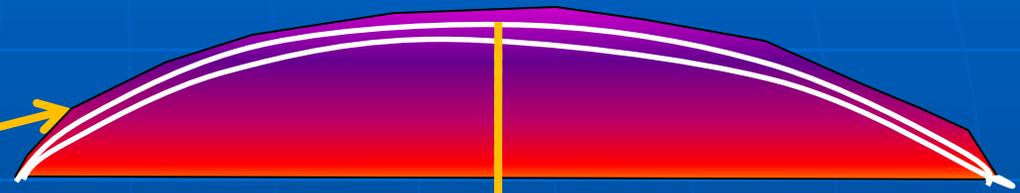
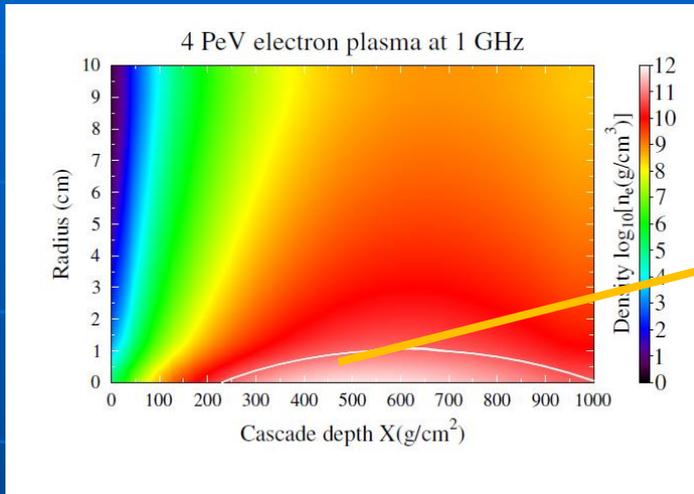
$$v_{Plasma} > v_{Radar} > \begin{cases} 1/\tau_{Plasma} & c_{med} \tau_e < l_c \\ c_{med} / l_c & c_{med} \tau_e > l_c \end{cases}$$

$$v_{Plasma} \propto \sqrt{n_{Plasma}} \propto \sqrt{E_{primary}}$$



# Skin Effects

Model: Consider over-dense cylinders of equal density



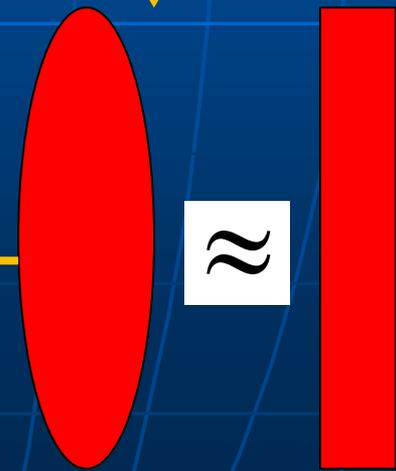
Calculate skin depth for a collision less plasma:

$$\delta = \frac{c}{2\omega_p}$$

Within 1 skin depth the amount of power absorbed and re-scattered equals:

$$f_{skin}^{i+1} = (1 - f_{skin}^i) \left(1 - e^{-\frac{x}{\delta_i}}\right)$$

$$A_{Plasma}^i \approx L_i r_i$$



# The over-dense radar cross-section

This approach:

1. Include skin-effects directly into the radar cross-section.
2. Consider projected area and polarization angles for in/outgoing wave

$$\sigma_{od} = A_{plasma} \times f_{skin} \times f_{geom}$$

$$A_{Plasma}^i \approx L_i r_i$$

$$f_{skin}^{i+1} = (1 - f_{skin}^i)(1 - e^{-x/\delta_i})$$

$$f_{geom} = (\vec{e}_t \cdot \vec{e}_c)(\vec{e}_c \cdot \vec{e}_r)$$

$$\sigma_{od} = \sum_i L_i r_i (1 - f_{skin}^i)(1 - e^{-x/\delta_i})(\vec{e}_t \cdot \vec{e}_c)(\vec{e}_c \cdot \vec{e}_r)$$

# The under-dense radar cross-section

The wave will scatter off of the individual electron given by the Thompson cross-section

$$\sigma_T = \left( \frac{m_e}{m_p} \right)^2 0.665 \cdot 10^{-28} \text{ m}^2$$

We have to take into account for the phase lag of the individual electrons w.r.t. each other:

$$\sigma_{ud} = \sum_{i=1}^N \sigma_T \cos(kx)$$

$$k = \frac{2\pi}{\lambda_d} \quad x = |\vec{x}_1 - \vec{x}_i| + |\vec{x}_2 - \vec{x}_i|$$

# Radar scattering

## What do we see?

