

EHDECAY

Michael Spira (PSI)

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III eHDECAY

<http://www.itp.kit.edu/~maggie/eHDECAY/>

IV Other Applications

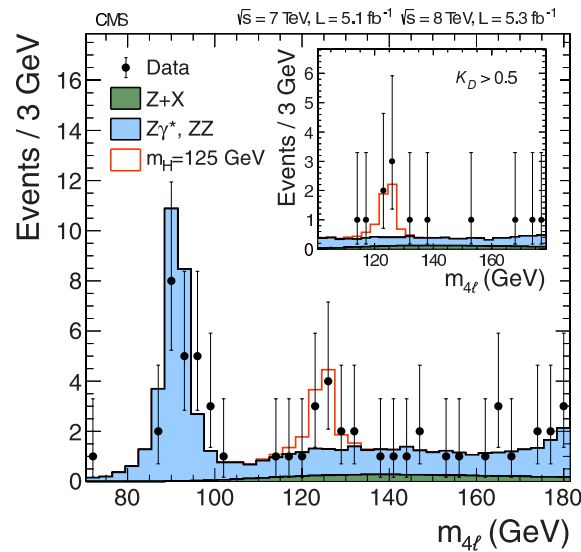
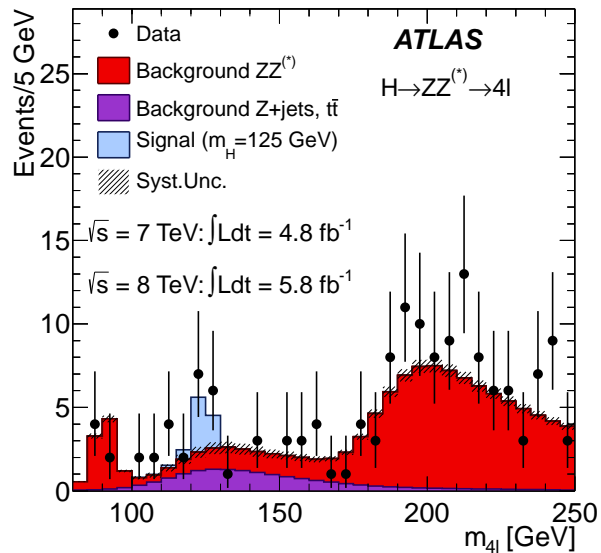
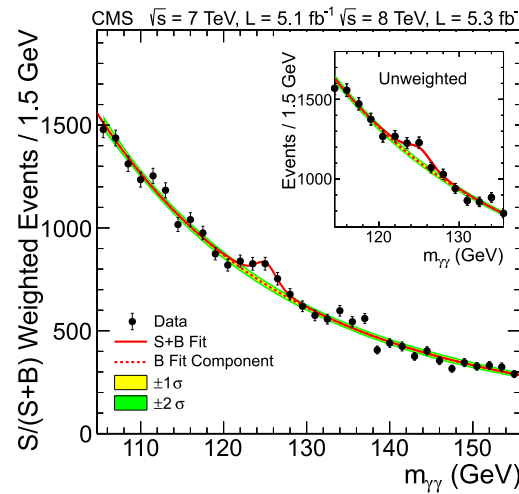
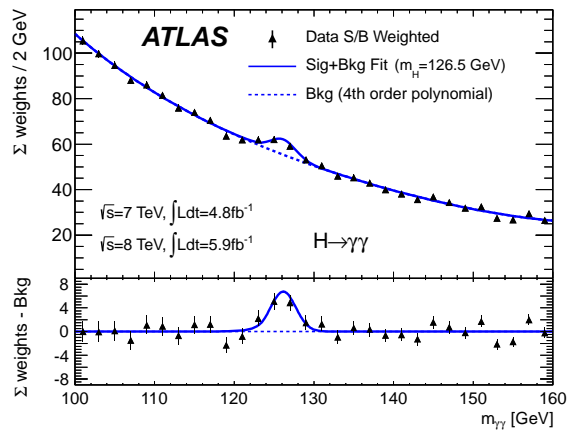
V Conclusions

eHDECAY in collaboration with R. Contino, M. Ghezzi, C. Grojean and M. Mühlleitner

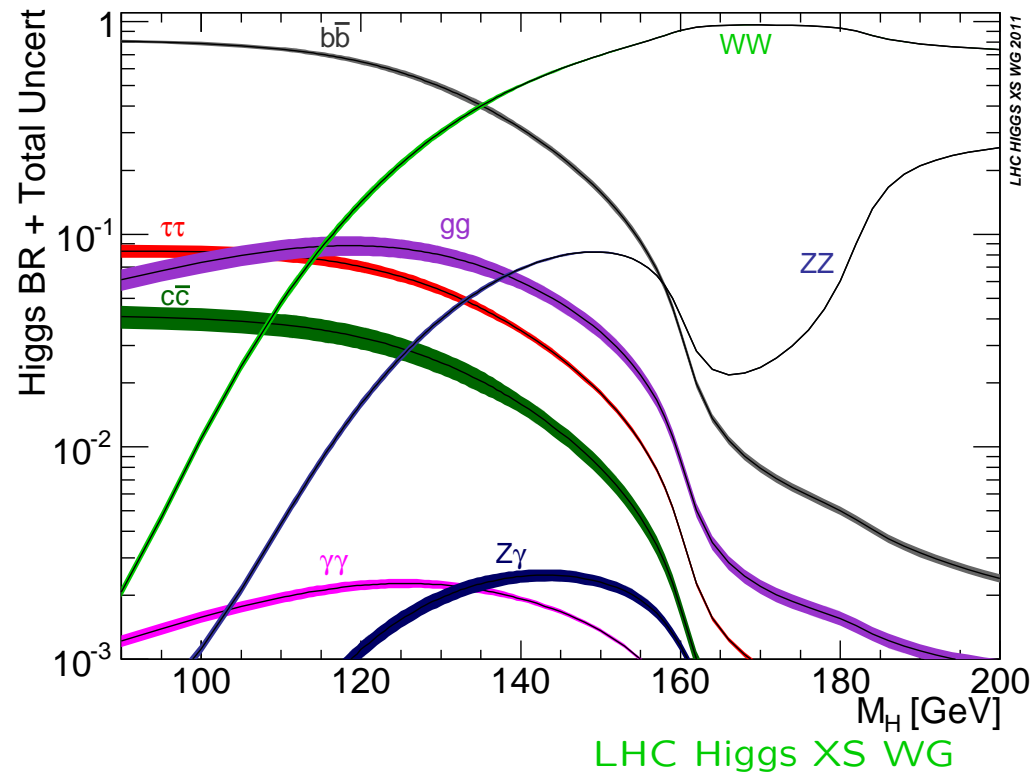
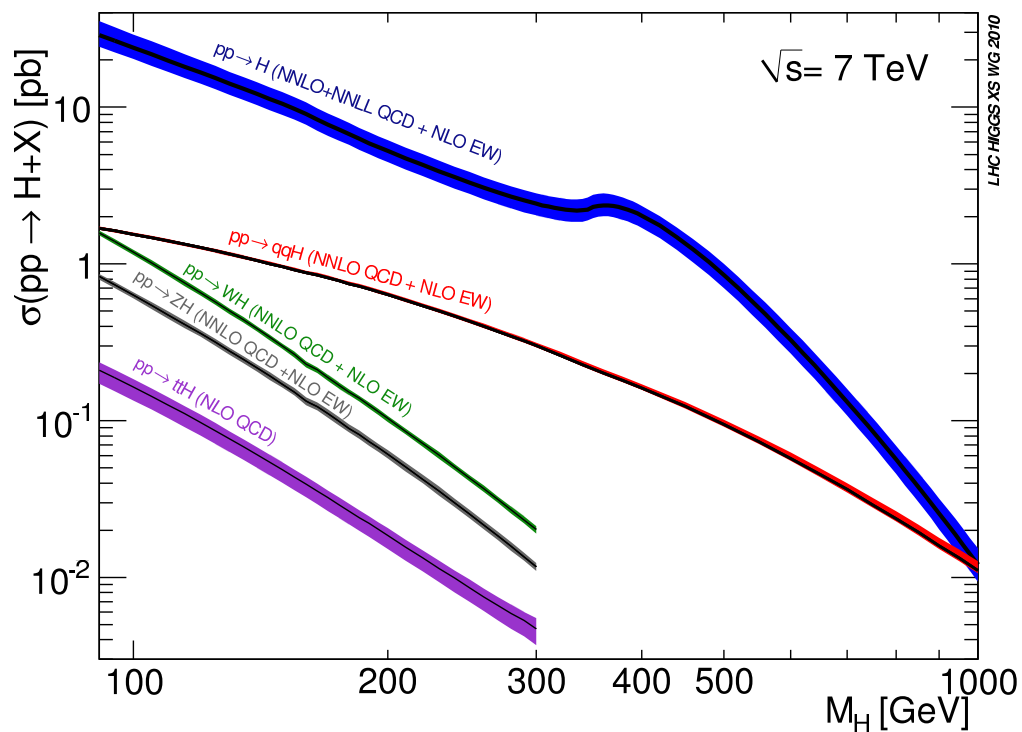
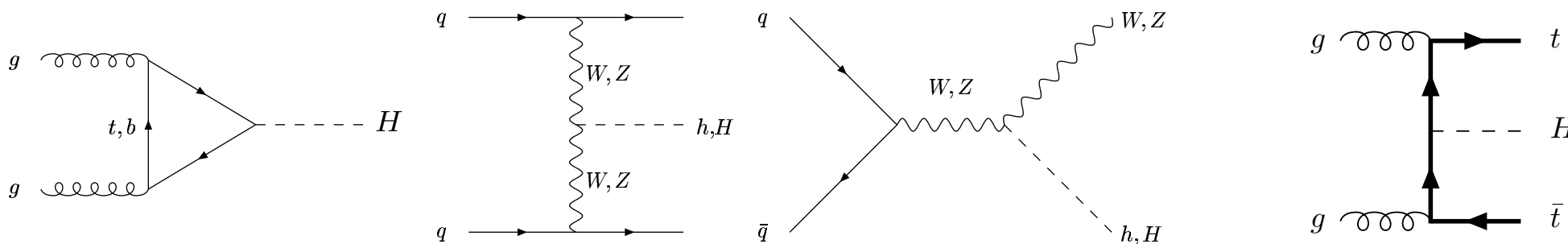
I INTRODUCTION

Standard Model

- we have found the Higgs: $M_H \sim 125$ GeV
- $gg \rightarrow H$ dominant



● Higgs Boson Production



- Discovery: LHC [Tevatron]

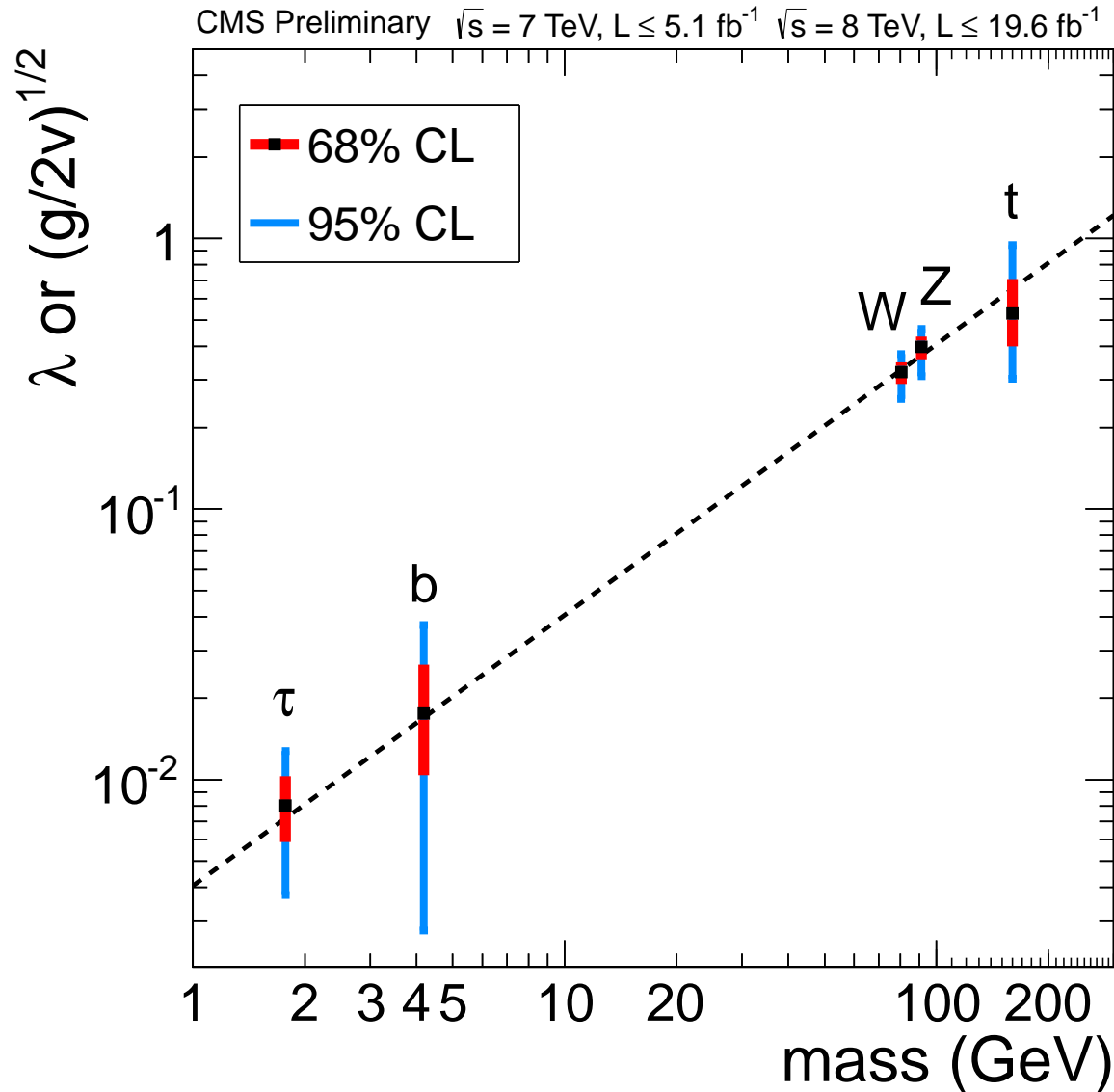
→ Higgs mass

couplings

spin

CP

$\lambda ?$



CMS

- deviations → effective Lagrangians

II EFFECTIVE LAGRANGIANS

(i) weakly interacting theories

- effective higher dimension operators up to dim 6 Buchmüller, Wyler
Grzadkowski, Iskrzynski, Misiak, Rosiek
Giudice, Grojean, Pomarol, Rattazzi

$$\begin{aligned}\mathcal{L} &= \mathcal{L}_{SM} + \frac{1}{\Lambda^2} \sum_i \alpha_i O_i \\ &\equiv \mathcal{L}_{SM} + \sum_i \bar{c}_i O_i \\ &\equiv \mathcal{L}_{SM} + \Delta\mathcal{L}_{SILH} + \Delta\mathcal{L}_{F_1} + \Delta\mathcal{L}_{F_2} + \Delta\mathcal{L}_{bos} + \Delta\mathcal{L}_{4f} + \Delta\mathcal{L}_{CP}\end{aligned}$$

[assume Λ large]

- assume Higgs $SU(2)$ -doublet

$$H = \begin{pmatrix} \phi^+ \\ \phi^0 \end{pmatrix}$$

$$\begin{aligned}
\Delta\mathcal{L}_{SILH} &= \frac{\bar{c}_H}{2v^2} \partial^\mu (H^\dagger H) \partial_\mu (H^\dagger H) + \frac{\bar{c}_T}{2v^2} \left(H^\dagger \overleftrightarrow{D}^\mu H \right) \left(H^\dagger \overleftrightarrow{D}_\mu H \right) - \frac{\bar{c}_6 \lambda}{v^2} (H^\dagger H)^3 \\
&+ \left(\frac{\bar{c}_u}{v^2} y_u H^\dagger H \bar{q}_L H^c u_R + \frac{\bar{c}_d}{v^2} y_d H^\dagger H \bar{q}_L H d_R + \frac{\bar{c}_l}{v^2} y_l H^\dagger H \bar{L}_L H l_R + h.c. \right) \\
&+ \frac{i\bar{c}_W g}{2m_W^2} \left(H^\dagger \sigma^i \overleftrightarrow{D}^\mu H \right) (D^\nu W_{\mu\nu})^i + \frac{i\bar{c}_B g'}{2m_W^2} \left(H^\dagger \overleftrightarrow{D}^\mu H \right) (\partial^\nu B_{\mu\nu}) \\
&+ \frac{i\bar{c}_{HW} g}{m_W^2} (D^\mu H)^\dagger \sigma^i (D^\nu H) W_{\mu\nu}^i + \frac{i\bar{c}_{HB} g'}{m_W^2} (D^\mu H)^\dagger (D^\nu H) B_{\mu\nu} \\
&+ \frac{\bar{c}_\gamma g'^2}{m_W^2} H^\dagger H B_{\mu\nu} B^{\mu\nu} + \frac{\bar{c}_g g_S^2}{m_W^2} H^\dagger H G_{\mu\nu}^a G^{a\mu\nu} \\
\Delta\mathcal{L}_{F_1} &= \frac{i\bar{c}_{Hq}}{v^2} (\bar{q}_L \gamma^\mu q_L) \left(H^\dagger \overleftrightarrow{D}_\mu H \right) + \frac{i\bar{c}'_{Hq}}{v^2} (\bar{q}_L \gamma^\mu \sigma^i q_L) \left(H^\dagger \sigma^i \overleftrightarrow{D}_\mu H \right) \\
&+ \frac{i\bar{c}_{Hu}}{v^2} (\bar{u}_R \gamma^\mu u_R) \left(H^\dagger \overleftrightarrow{D}_\mu H \right) + \frac{i\bar{c}_{Hd}}{v^2} (\bar{d}_R \gamma^\mu d_R) \left(H^\dagger \overleftrightarrow{D}_\mu H \right) \\
&+ \left(\frac{i\bar{c}_{Hud}}{v^2} (\bar{u}_R \gamma^\mu d_R) \left(H^{c\dagger} \overleftrightarrow{D}_\mu H \right) + h.c. \right) \\
&+ \frac{i\bar{c}_{HL}}{v^2} (\bar{L}_L \gamma^\mu L_L) \left(H^\dagger \overleftrightarrow{D}_\mu H \right) + \frac{i\bar{c}'_{HL}}{v^2} (\bar{L}_L \gamma^\mu \sigma^i L_L) \left(H^\dagger \sigma^i \overleftrightarrow{D}_\mu H \right) \\
&+ \frac{i\bar{c}_{Hl}}{v^2} (\bar{l}_R \gamma^\mu l_R) \left(H^\dagger \overleftrightarrow{D}_\mu H \right) \\
\Delta\mathcal{L}_{F_2} &= \frac{\bar{c}_{uB} g'}{m_W^2} y_u \bar{q}_L H^c \sigma^{\mu\nu} u_R B_{\mu\nu} + \frac{\bar{c}_{uW} g}{m_W^2} y_u \bar{q}_L \sigma^i H^c \sigma^{\mu\nu} u_R W_{\mu\nu}^i + \frac{\bar{c}_{uG} g_S}{m_W^2} y_u \bar{q}_L H^c \sigma^{\mu\nu} \lambda^a u_R G_{\mu\nu}^a \\
&+ \frac{\bar{c}_{dB} g'}{m_W^2} y_d \bar{q}_L H \sigma^{\mu\nu} d_R B_{\mu\nu} + \frac{\bar{c}_{dW} g}{m_W^2} y_d \bar{q}_L \sigma^i H \sigma^{\mu\nu} d_R W_{\mu\nu}^i + \frac{\bar{c}_{dG} g_S}{m_W^2} y_d \bar{q}_L H \sigma^{\mu\nu} \lambda^a d_R G_{\mu\nu}^a \\
&+ \frac{\bar{c}_{lB} g'}{m_W^2} y_l \bar{L}_L H \sigma^{\mu\nu} l_R B_{\mu\nu} + \frac{\bar{c}_{lW} g}{m_W^2} y_l \bar{L}_L \sigma^i H \sigma^{\mu\nu} l_R W_{\mu\nu}^i + h.c.
\end{aligned}$$

$$\begin{aligned}
\Delta\mathcal{L}_{bos} &= \frac{\bar{c}_{3W} g^3}{m_W^2} \epsilon^{ijk} W_\mu^{i\nu} W_\nu^{j\rho} W_\rho^{k\mu} + \frac{\bar{c}_{3G} g_S^3}{m_W^2} f^{abc} G_\mu^{a\nu} G_\nu^{b\rho} G_\rho^{c\mu} \\
&+ \frac{\bar{c}_{2W}}{m_W^2} (D^\mu W_{\mu\nu})^i (D_\rho W^{\rho\nu})^i + \frac{\bar{c}_{2B}}{m_W^2} (\partial^\mu B_{\mu\nu}) (\partial_\rho B^{\rho\nu}) + \frac{\bar{c}_{2G}}{m_W^2} (D^\mu G_{\mu\nu})^a (D_\rho G^{\rho\nu})^a \\
\Delta\mathcal{L}_{4f} &= \sum_{\psi, L/R, T^a} \bar{\psi}_i \gamma^\mu T^a \psi_j \bar{\psi}_k \gamma_\mu T^a \psi_l + \bar{\psi}_i T^a \psi_j \bar{\psi}_k T^a \psi_l \\
\Delta\mathcal{L}_{\mathcal{CP}} &= \frac{i\tilde{c}_{HW} g}{m_W^2} (D^\mu H)^\dagger \sigma^i (D^\nu H) \tilde{W}_{\mu\nu}^i + \frac{i\tilde{c}_{HB} g'}{m_W^2} (D^\mu H)^\dagger (D^\nu H) \tilde{B}_{\mu\nu} \\
&+ \frac{\tilde{c}_\gamma g'^2}{m_W^2} H^\dagger H B_{\mu\nu} \tilde{B}^{\mu\nu} + \frac{\tilde{c}_g g_S^2}{m_W^2} H^\dagger H G_{\mu\nu}^a \tilde{G}^{a\mu\nu} \\
&+ \frac{\tilde{c}_{3W} g^3}{m_W^2} \epsilon^{ijk} W_\mu^{i\nu} W_\nu^{j\rho} \tilde{W}_\rho^{k\mu} + \frac{\tilde{c}_{3G} g_S^3}{m_W^2} f^{abc} G_\mu^{a\nu} G_\nu^{b\rho} \tilde{G}_\rho^{c\mu}
\end{aligned}$$

$$\tilde{V}_{\mu\nu} = \frac{1}{2} \epsilon_{\mu\nu\alpha\beta} V^{\alpha\beta}$$

- after using EOM: 53 (59) independent dim6 operators
- canonical normalization

$$\begin{aligned}
\mathcal{L} = & \frac{1}{2} \partial_\mu h \partial^\mu h - \frac{1}{2} m_h^2 h^2 - c_3 \frac{1}{6} \left(\frac{3m_h^2}{v} \right) h^3 - \sum_{\psi=u,d,l} m_{\psi^{(i)}} \bar{\psi}^{(i)} \psi^{(i)} \left(1 + c_\psi \frac{h}{v} + \dots \right) \\
& + m_W^2 W_\mu W^\mu \left(1 + 2c_W \frac{h}{v} + \dots \right) + \frac{1}{2} m_Z^2 Z_\mu Z^\mu \left(1 + 2c_Z \frac{h}{v} + \dots \right) + \dots \\
& + \left(c_{WW} W_{\mu\nu}^+ W^{-\mu\nu} + \frac{c_{ZZ}}{2} Z_{\mu\nu} Z^{\mu\nu} + c_{Z\gamma} Z_{\mu\nu} \gamma^{\mu\nu} + \frac{c_{\gamma\gamma}}{2} \gamma_{\mu\nu} \gamma^{\mu\nu} + \frac{c_{gg}}{2} G_{\mu\nu}^a G^{a\mu\nu} \right) \frac{h}{v} \\
& + \left(c_{W\partial W} \left(W_\nu^- D_\mu W^{+\mu\nu} + h.c. \right) + c_{Z\partial Z} Z_\nu \partial_\mu Z^{\mu\nu} + c_{Z\partial\gamma} Z_\nu \partial_\mu \gamma^{\mu\nu} \right) \frac{h}{v} + \dots
\end{aligned}$$

Higgs couplings	$\Delta\mathcal{L}_{SILH}$	MCHM4	MCHM5
c_W	$1 - \bar{c}_H/2$	$\sqrt{1 - \xi}$	$\sqrt{1 - \xi}$
c_Z	$1 - \bar{c}_H/2 - \bar{c}_T$	$\sqrt{1 - \xi}$	$\sqrt{1 - \xi}$
c_ψ ($\psi = u, d, l$)	$1 - (\bar{c}_H/2 + \bar{c}_\psi)$	$\sqrt{1 - \xi}$	$\frac{1 - 2\xi}{\sqrt{1 - \xi}}$
c_3	$1 + \bar{c}_6 - 3\bar{c}_H/2$	$\sqrt{1 - \xi}$	$\frac{1 - 2\xi}{\sqrt{1 - \xi}}$
c_{gg}	$8(\alpha_s/\alpha_2) \bar{c}_g$	0	0
$c_{\gamma\gamma}$	$8 \sin^2 \theta_W \bar{c}_\gamma$	0	0
$c_{Z\gamma}$	$(\bar{c}_{HB} - \bar{c}_{HW} - 8 \bar{c}_\gamma \sin^2 \theta_W) \tan \theta_W$	0	0
c_{WW}	$-2 \bar{c}_{HW}$	0	0
c_{ZZ}	$-2 (\bar{c}_{HW} + \bar{c}_{HB} \tan^2 \theta_W - 4 \bar{c}_\gamma \tan^2 \theta_W \sin^2 \theta_W)$	0	0
$c_{W\partial W}$	$-2(\bar{c}_W + \bar{c}_{HW})$	0	0
$c_{Z\partial Z}$	$-2(\bar{c}_W + \bar{c}_{HW}) - 2(\bar{c}_B + \bar{c}_{HB}) \tan^2 \theta_W$	0	0
$c_{Z\partial\gamma}$	$2(\bar{c}_B + \bar{c}_{HB} - \bar{c}_W - \bar{c}_{HW}) \tan \theta_W$	0	0

small deviations from SM couplings

Contino, Ghezzi, Grojean, Mühlleitner, S.

$$\begin{aligned}
\mathcal{L} = & \frac{1}{2} \partial_\mu h \partial^\mu h - \frac{1}{2} m_h^2 h^2 - c_3 \frac{1}{6} \left(\frac{3m_h^2}{v} \right) h^3 - \sum_{\psi=u,d,l} m_{\psi^{(i)}} \bar{\psi}^{(i)} \psi^{(i)} \left(1 + c_\psi \frac{h}{v} + \dots \right) \\
& + m_W^2 W_\mu W^\mu \left(1 + 2c_W \frac{h}{v} + \dots \right) + \frac{1}{2} m_Z^2 Z_\mu Z^\mu \left(1 + 2c_Z \frac{h}{v} + \dots \right) + \dots \\
& + \left(c_{WW} W_{\mu\nu}^+ W^{-\mu\nu} + \frac{c_{ZZ}}{2} Z_{\mu\nu} Z^{\mu\nu} + c_{Z\gamma} Z_{\mu\nu} \gamma^{\mu\nu} + \frac{c_{\gamma\gamma}}{2} \gamma_{\mu\nu} \gamma^{\mu\nu} + \frac{c_{gg}}{2} G_{\mu\nu}^a G^{a\mu\nu} \right) \frac{h}{v} \\
& + \left(c_{W\partial W} (W_\nu^- D_\mu W^{+\mu\nu} + h.c.) + c_{Z\partial Z} Z_\nu \partial_\mu Z^{\mu\nu} + c_{Z\partial\gamma} Z_\nu \partial_\mu \gamma^{\mu\nu} \right) \frac{h}{v} + \dots
\end{aligned}$$

- also valid in case of a non-linear Lagrangian for a light Higgs-like scalar [h generic \mathcal{CP} -even scalar]

\Rightarrow expansion in E/M (derivatives) only, large deviations from SM couplings allowed

SILH: expansion in $v^2/f^2, E^2/M^2, \alpha_s/\pi, \alpha/\pi$
non-lin.: expansion in $E^2/M^2, \alpha_s/\pi$

Partial Width	QCD	Electroweak	Total	
$H \rightarrow b\bar{b}/c\bar{c}$	$\sim 0.1\%$	$\sim 1\text{--}2\%$ for $M_H \lesssim 135\text{GeV}$	$\sim 2\%$	NNNNLO / NLO
$H \rightarrow \tau^+\tau^-/\mu^+\mu^-$		$\sim 1\text{--}2\%$ for $M_H \lesssim 135\text{GeV}$	$\sim 2\%$	NLO
$H \rightarrow t\bar{t}$	$\lesssim 5\%$	$\lesssim 2\text{--}5\%$ for $M_H < 500\text{GeV}$ $\sim 0.1(\frac{M_H}{1\text{TeV}})^4$ for $M_H > 500\text{GeV}$	$\sim 5\%$ $\sim 5\text{--}10\%$	(NNN)NLO / LO
$H \rightarrow gg$	$\sim 3\%$	$\sim 1\%$	$\sim 3\%$	NNNLO approx. / NLO
$H \rightarrow \gamma\gamma$	$< 1\%$	$< 1\%$	$\sim 1\%$	NLO / NLO
$H \rightarrow Z\gamma$	$< 1\%$	$\sim 5\%$	$\sim 5\%$	(N)LO / LO
$H \rightarrow WW/ZZ \rightarrow 4f$	$< 0.5\%$	$\sim 0.5\%$ for $M_H < 500\text{GeV}$ $\sim 0.17(\frac{M_H}{1\text{TeV}})^4$ for $M_H > 500\text{GeV}$	$\sim 0.5\%$ $\sim 0.5\text{--}15\%$	(N)NLO

- QCD: variation of Higgs widths for scale by factor 2 and 1/2
elw: missing HO estimated from known structure at NLO
 $M_H \gtrsim 500$ GeV: Higgs self-interactions dominate error
different uncertainties added linearly for each channel
- parametric uncertainties:

$m_t = 172.5 \pm 2.5$ GeV	$\alpha_s(M_Z) = 0.119 \pm 0.002$
$m_b(m_b) = 4.16 \pm 0.06$ GeV	$m_c(m_c) = 1.28 \pm 0.03$ GeV

 different uncertainties added quadratically for each channel
- total uncertainties: parametric & theor. uncertainties added linearly

III eHDECAY

<http://www.itp.kit.edu/~maggie/eHDECAY/>

- $h \rightarrow f \bar{f}$:

$$\Gamma(\bar{\psi}\psi)|_{SILH} = \Gamma_0^{SM}(\bar{\psi}\psi) \left[1 - \bar{c}_H - 2\bar{c}_\psi + \frac{2}{|A_0^{SM}|^2} \text{Re}(A_0^{*SM} A_{1,ew}^{SM}) \right] [1 + \delta_\psi \kappa^{QCD}]$$

$$\Gamma(\bar{\psi}\psi)|_{NL} = c_\psi^2 \Gamma_0^{SM}(\bar{\psi}\psi) [1 + \delta_\psi \kappa^{QCD}]$$

A_0^{SM} : SM tree-level amplitude

$A_{1,ew}^{SM}$: SM elw. amplitude [real corrections treated analogously]

- factorization of QCD \leftrightarrow elw. [limit small m_h]
- NL: no elw. corrections!
- other decay modes analogous

- approximate formulae [w/o elw. corrections]: $\alpha_2 = \sqrt{2}G_F m_W^2 / \pi$

$$\frac{\Gamma(\bar{\psi}\psi)}{\Gamma(\bar{\psi}\psi)_{SM}} \simeq 1 - \bar{c}_H - 2\bar{c}_\psi \quad (\text{small } \bar{c}_t \text{ contamination for } s\bar{s}, c\bar{c}, b\bar{b})$$

$$\frac{\Gamma(h \rightarrow W^{(*)}W^*)}{\Gamma(h \rightarrow W^{(*)}W^*)_{SM}} \simeq 1 - \bar{c}_H + 2.2\bar{c}_W + 3.7\bar{c}_{HW}$$

$$\begin{aligned} \frac{\Gamma(h \rightarrow Z^{(*)}Z^*)}{\Gamma(h \rightarrow Z^{(*)}Z^*)_{SM}} &\simeq 1 - \bar{c}_H - 2\bar{c}_T + 2.0(\bar{c}_W + \tan^2\theta_W \bar{c}_B) \\ &+ 3.0(\bar{c}_{HW} + \tan^2\theta_W \bar{c}_{HB}) - 0.26\bar{c}_\gamma \end{aligned}$$

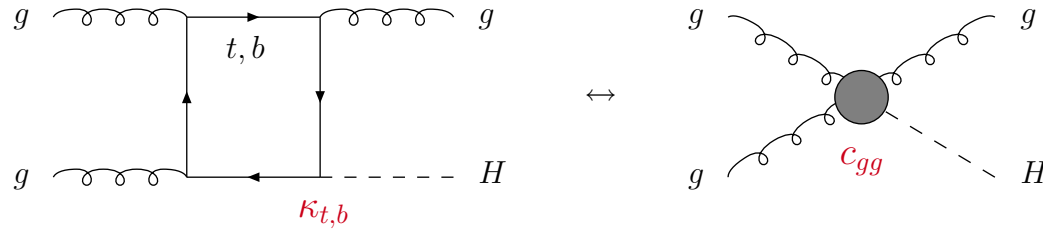
$$\begin{aligned} \frac{\Gamma(h \rightarrow Z\gamma)}{\Gamma(h \rightarrow Z\gamma)_{SM}} &\simeq 1 - \bar{c}_H + 0.12\bar{c}_t - 5 \cdot 10^{-4}\bar{c}_c - 0.003\bar{c}_b - 9 \cdot 10^{-5}\bar{c}_\tau \\ &+ 4.2\bar{c}_W + 0.19(\bar{c}_{HW} - \bar{c}_{HB} + 8\bar{c}_\gamma \sin^2\theta_W) \frac{4\pi}{\sqrt{\alpha_2\alpha_{em}}} \end{aligned}$$

$$\begin{aligned} \frac{\Gamma(h \rightarrow \gamma\gamma)}{\Gamma(h \rightarrow \gamma\gamma)_{SM}} &\simeq 1 - \bar{c}_H + 0.54\bar{c}_t - 0.003\bar{c}_c - 0.007\bar{c}_b - 0.007\bar{c}_\tau \\ &+ 5.04\bar{c}_W - 0.54\bar{c}_\gamma \frac{4\pi}{\alpha_{em}} \end{aligned}$$

$$\frac{\Gamma(h \rightarrow gg)}{\Gamma(h \rightarrow gg)_{SM}} \simeq 1 - \bar{c}_H - 2.12\bar{c}_t + 0.024\bar{c}_c + 0.1\bar{c}_b + 22.2\bar{c}_g \frac{4\pi}{\alpha_2}$$

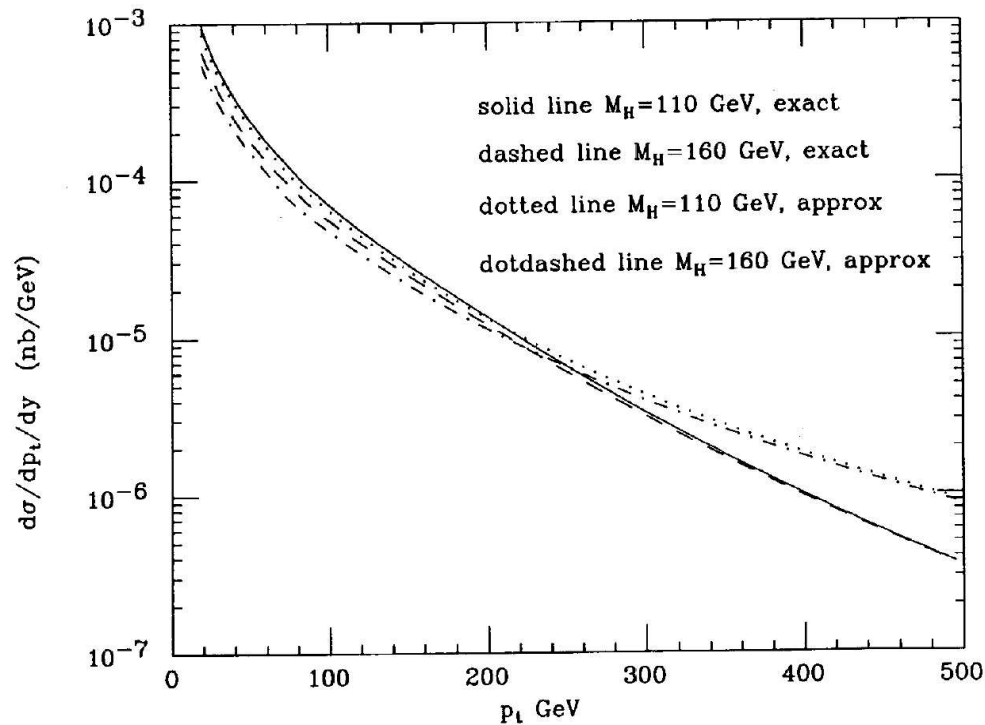
IV OTHER APPLICATIONS

Higgs p_T (or how to prove that ggF is loop-mediated)



- distinction dim4 \leftrightarrow dim5

Harlander, Neumann



$m_t = 160 \text{ GeV}$

Ellis, Hinchliffe, Soldate, van der Bij

- implementation of dim6 $G^{a\mu\nu}G_{\mu\nu}^a H$ in HIGLU (σ_{tot}, p_{TH}) planned

V CONCLUSIONS

- Higgs decays: eHDECAY
- inclusion of SILH (dim 6) and non-linear Lagrangians
- systematic extension of SM \rightarrow well-defined expansions
- SILH: expansion in $v^2/f^2, E^2/M^2, \alpha_s/\pi, \alpha/\pi$
- non-lin.: expansion in $E^2/M^2, \alpha_s/\pi$
- eHDECAY provides consistent tool for BSM Higgs decays
<http://www.itp.kit.edu/~maggie/eHDECAY/>
[extension to subleading operators planned]
- important impact on Higgs production modes

BACKUP SLIDES

II EFFECTIVE LAGRANGIANS

- $WW \rightarrow WW$ @ high energies

(a)

(b)

$$\mathcal{A} = \frac{s}{v^2} \left\{ 1 - \frac{\kappa_V^2 s}{s - M_H^2} \right\} \Rightarrow \kappa_V = 1$$

- $f\bar{f} \rightarrow WW$ @ high energies

(a)

(b)

$$\mathcal{A} = \frac{m_f \sqrt{s}}{v^2} \left\{ 1 - \frac{\kappa_f \kappa_V s}{s - M_H^2} \right\} \Rightarrow \kappa_f = \kappa_V = 1$$

- analogously for κ_H

- constraints from precision measurements:

$$\Delta\epsilon_1 \equiv \Delta\rho = \bar{c}_T(m_Z), \quad -1.5 \times 10^{-3} < \bar{c}_T(m_Z) < 2.2 \times 10^{-3}$$

$$\Delta\epsilon_3 = \bar{c}_W(m_Z) + \bar{c}_B(m_Z), \quad -1.4 \times 10^{-3} < \bar{c}_W(m_Z) + \bar{c}_B(m_Z) < 1.9 \times 10^{-3}$$

- Z -pole measurements:

$$\frac{\delta g_{L\psi}}{g_{L\psi}} = \frac{1}{2} \frac{\bar{c}_{H\psi} + 2 T_{3L} \bar{c}'_{H\psi}}{T_{3L} - Q \sin^2 \theta_W}, \quad \frac{\delta g_{R\psi}}{g_{R\psi}} = \frac{1}{2} \frac{\bar{c}_{H\psi}}{Q \sin^2 \theta_W}$$

$$-0.03 < \bar{c}_{Hq1} < 0.02, \quad -0.002 < \bar{c}'_{Hq1} < 0.003,$$

$$-0.005 < \bar{c}_{Hq2} < 0.003, \quad -0.003 < \bar{c}'_{Hq2} < 0.005,$$

$$-0.008 < \bar{c}_{Hu} < 0.02, \quad -0.03 < \bar{c}_{Hd} < 0.02, \quad -0.03 < \bar{c}_{Hs} < 0.02$$

$$-0.004 < \bar{c}_{HL} + \bar{c}'_{HL} < 0.002, \quad -0.003 < \bar{c}_{HL} - \bar{c}'_{HL} < 0.0002, \quad -0.0007 < \bar{c}_{Hl} < 0.003,$$

$$-0.02 < \bar{c}_{Hq2} + \bar{c}'_{Hq2} < 0.005, \quad -0.02 < \bar{c}_{Hc} < 0.03,$$

$$-0.003 < \bar{c}_{Hq3} - \bar{c}'_{Hq3} < 0.009, \quad -0.07 < \bar{c}_{Hb} < -0.005$$

- EDMs: neutron & mercury:

$$-7.01 \times 10^{-6} < \text{Im}(\bar{c}_{uB} + \bar{c}_{uW}) < 7.86 \times 10^{-6},$$

$$-9.42 \times 10^{-7} < \text{Im}(\bar{c}_{dB} - \bar{c}_{dW}) < 8.40 \times 10^{-7},$$

$$-1.62 \times 10^{-6} < \text{Im}(\bar{c}_{uG}) < 2.01 \times 10^{-6},$$

$$-7.71 \times 10^{-7} < \text{Im}(\bar{c}_{dG}) < 5.70 \times 10^{-7},$$

- top quark: nEDM, $b \rightarrow s\gamma, s\ell^+\ell^-$:

$$-1.39 \times 10^{-4} < \text{Im}(\bar{c}_{tG}) < 1.21 \times 10^{-4}$$

$$-0.057 < \text{Re}(\bar{c}_{tW} + \bar{c}_{tB}) - 2.65 \text{Im}(\bar{c}_{tW} + \bar{c}_{tB}) < 0.20$$

$t\bar{t}$ cxns @ Tevatron & LHC:

$$-6.12 \times 10^{-3} < \text{Re}(\bar{c}_{tG}) < 1.94 \times 10^{-3}$$

$$-1.2 < \text{Re}(\bar{c}_{bW}) < 1.1, \quad -0.01 < \text{Re}(\bar{c}_{tW}) < 0.02$$

- leptons: EDMs & anomalous magnetic moments:

$$-1.64 \times 10^{-2} < \text{Re}(\bar{c}_{eB} - \bar{c}_{eW}) < 3.37 \times 10^{-3},$$

$$1.88 \times 10^{-4} < \text{Re}(\bar{c}_{\mu B} - \bar{c}_{\mu W}) < 6.43 \times 10^{-4},$$

$$-2.97 \times 10^{-7} < \text{Im}(\bar{c}_{eB} - \bar{c}_{eW}) < 4.51 \times 10^{-7},$$

$$-0.26 < \text{Im}(\bar{c}_{\mu B} - \bar{c}_{\mu W}) < 0.29,$$

• $h \rightarrow gg$:

$$\begin{aligned}
\Gamma(gg)|_{SILH} &= \frac{G_F \alpha_s^2 m_h^3}{4\sqrt{2}\pi^3} \left[\frac{1}{9} \sum_{q,q'=t,b,c} (1 - \bar{c}_H - \bar{c}_q - \bar{c}_{q'}) A_{1/2}^*(\tau_{q'}) A_{1/2}(\tau_q) c_{eff}^2 \kappa_{soft} \right. \\
&\quad + 2 \operatorname{Re} \left(\sum_{q=t,b,c} \frac{1}{3} A_{1/2}^*(\tau_q) \frac{16\pi \bar{c}_g}{\alpha_2} \right) c_{eff} \kappa_{soft} \\
&\quad + \left| \sum_{q=t,b,c} \frac{1}{3} A_{1/2}(\tau_q) \right|^2 c_{eff}^2 \kappa_{ew} \kappa_{soft} \\
&\quad \left. + \frac{1}{9} \sum_{q,q'=t,b} (1 - \bar{c}_H - \bar{c}_q - \bar{c}_{q'}) A_{1/2}^*(\tau_q) A_{1/2}(\tau_{q'}) \kappa^{NLO}(\tau_q, \tau_{q'}) \right] \\
\Gamma(gg)|_{NL} &= \frac{G_F \alpha_s^2 m_h^3}{4\sqrt{2}\pi^3} \left[\left| \sum_{q=t,b,c} \frac{c_q}{3} A_{1/2}(\tau_q) \right|^2 c_{eff}^2 \kappa_{soft} \right. \\
&\quad + 2 \operatorname{Re} \left(\sum_{q=t,b,c} \frac{c_q}{3} A_{1/2}^*(\tau_q) \frac{2\pi c_{gg}}{\alpha_s} \right) c_{eff} \kappa_{soft} + \left| \frac{2\pi c_{gg}}{\alpha_s} \right|^2 \kappa_{soft} \\
&\quad \left. + \frac{1}{9} \sum_{q,q'=t,b} c_q A_{1/2}^*(\tau_q) c_{q'} A_{1/2}(\tau_{q'}) \kappa^{NLO}(\tau_q, \tau_{q'}) \right]
\end{aligned}$$

$$A_{1/2}(\tau) = \frac{3}{2}\tau [1 + (1 - \tau) f(\tau)]$$

$$f(\tau) = \begin{cases} \arcsin^2 \frac{1}{\sqrt{\tau}} & \tau \geq 1 \\ -\frac{1}{4} \left[\ln \frac{1 + \sqrt{1 - \tau}}{1 - \sqrt{1 - \tau}} - i\pi \right]^2 & \tau < 1. \end{cases}$$

$$\kappa_{soft}^{NLO} = 1 + \frac{\alpha_s^{NLO}}{\pi} \left(\frac{73}{4} - \frac{7}{6} N_F \right), \quad c_{eff}^{NLO} = 1 + \frac{\alpha_s^{NLO}}{\pi} \frac{11}{4}$$

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 Baikov, Chetyrkin

- $\kappa^{NLO}(\tau_q, \tau_{q'})$: NLO mass effects ($\lesssim 5\%$ in SM)

- $h \rightarrow \gamma\gamma$:

$$\Gamma(\gamma\gamma)|_{SILH} = \frac{G_F \alpha_{em}^2 m_h^3}{128 \sqrt{2} \pi^3} \left\{ |A_{NLO}^{SM}(\gamma\gamma)|^2 + 2 \operatorname{Re} \left(A_{LO}^{SM*}(\gamma\gamma) A_{ew}^{SM}(\gamma\gamma) \right) + 2 \operatorname{Re} \left[A_{NLO}^{SM*}(\gamma\gamma) \left(\Delta A(\gamma\gamma) + \frac{32\pi \sin^2 \theta_W \bar{c}_\gamma}{\alpha_{em}} \right) \right] \right\}$$

$$\Gamma(\gamma\gamma)|_{NL} = \frac{G_F \alpha_{em}^2 m_h^3}{128 \sqrt{2} \pi^3} \left| \sum_{q=t,b,c} \frac{4}{3} c_q 3Q_q^2 A_{1/2}^{NLO}(\tau_q) + \frac{4}{3} c_\tau Q_\tau^2 A_{1/2}(\tau_\tau) + c_V A_1(\tau_W) + \frac{4\pi}{\alpha_{em}} c_{\gamma\gamma} \right|^2$$

$$\Delta A(\gamma\gamma) = - \sum_{q=t,b,c} \frac{4}{3} \left(\frac{\bar{c}_H}{2} + \bar{c}_q \right) 3Q_q^2 A_{1/2}^{NLO}(\tau_q) - \left(\frac{\bar{c}_H}{2} + \bar{c}_\tau \right) \frac{4}{3} Q_\tau^2 A_{1/2}(\tau_\tau) - \left(\frac{\bar{c}_H}{2} - 2\bar{c}_W \right) A_1(\tau_W)$$

$$A_1(\tau) = -[2 + 3\tau + 3\tau(2 - \tau)f(\tau)]$$

$$A_{1/2}^{NLO}(\tau_q) = A_{1/2}(\tau_q)(1 + \kappa_{QCD})$$

- κ_{QCD} : massive QCD corrections

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• $h \rightarrow Z\gamma$:

$$\begin{aligned}
\Gamma(Z\gamma)|_{SILH} &= \frac{G_F^2 \alpha_{em} m_W^2 m_h^3}{64\pi^4} \left(1 - \frac{m_Z^2}{m_h^2}\right)^3 \\
&\times \left\{ |A^{SM}(Z\gamma)|^2 + 2 \operatorname{Re}(A^{SM*}(Z\gamma) \Delta A(Z\gamma)) \right. \\
&\quad \left. + 2 \operatorname{Re} \left[-\frac{4\pi \tan \theta_W}{\sqrt{\alpha_{em} \alpha_2}} (\bar{c}_{HB} - \bar{c}_{HW} - 8\bar{c}_\gamma \sin^2 \theta_W) A^{SM*}(Z\gamma) \right] \right\} \\
\Gamma(Z\gamma)|_{NL} &= \frac{G_F^2 \alpha_{em} m_W^2 m_h^3}{64\pi^4} \left(1 - \frac{m_Z^2}{m_h^2}\right)^3 \\
&\times \left| \sum_\psi \frac{c_\psi N_c Q_\psi \hat{v}_\psi}{\cos \theta_W} A_{1/2}^{Z\gamma}(\tau_\psi, \lambda_\psi) + c_V A_1^{Z\gamma}(\tau_W, \lambda_W) - \frac{4\pi}{\sqrt{\alpha_{em} \alpha_2}} c_{Z\gamma} \right|^2 \\
A_{1/2}^{Z\gamma}(\tau, \lambda) &= [I_1(\tau, \lambda) - I_2(\tau, \lambda)] \\
A_1^{Z\gamma}(\tau, \lambda) &= \cos \theta_W \left\{ 4(3 - \tan^2 \theta_W) I_2(\tau, \lambda) \right. \\
&\quad \left. + \left[\left(1 + \frac{2}{\tau}\right) \tan^2 \theta_W - \left(5 + \frac{2}{\tau}\right) \right] I_1(\tau, \lambda) \right\} \\
\Delta A(Z\gamma) &= - \sum_\psi \left(\frac{\bar{c}_H}{2} + \bar{c}_\psi \right) \frac{N_c Q_\psi \hat{v}_\psi}{\cos \theta_W} A_{1/2}^{Z\gamma}(\tau_\psi, \lambda_\psi) - \left(\frac{\bar{c}_H}{2} - 2\bar{c}_W \right) A_1^{Z\gamma}(\tau_W, \lambda_W) \\
A^{SM}(Z\gamma) &= \sum_\psi \frac{N_c Q_\psi \hat{v}_\psi}{\cos \theta_W} A_{1/2}^{Z\gamma}(\tau_\psi, \lambda_\psi) + A_1^{Z\gamma}(\tau_W, \lambda_W)
\end{aligned}$$

$$I_1(\tau, \lambda) = \frac{\tau\lambda}{2(\tau - \lambda)} + \frac{\tau^2\lambda^2}{2(\tau - \lambda)^2} [f(\tau) - f(\lambda)] + \frac{\tau^2\lambda}{(\tau - \lambda)^2} [g(\tau) - g(\lambda)]$$

$$I_2(\tau, \lambda) = -\frac{\tau\lambda}{2(\tau - \lambda)} [f(\tau) - f(\lambda)]$$

$$g(\tau) = \begin{cases} \sqrt{\tau - 1} \arcsin \frac{1}{\sqrt{\tau}} & \tau \geq 1 \\ \frac{\sqrt{1 - \tau}}{2} \left[\ln \frac{1 + \sqrt{1 - \tau}}{1 - \sqrt{1 - \tau}} - i\pi \right] & \tau < 1. \end{cases}$$

• $h \rightarrow Z^*Z^*, W^*W^*$:

$$\Gamma(V^*V^*) = \frac{1}{\pi^2} \int_0^{m_h^2} \frac{dQ_1^2 m_V \Gamma_V}{(Q_1^2 - m_V^2)^2 + m_V^2 \Gamma_V^2} \int_0^{(m_h - Q_1)^2} \frac{dQ_2^2 m_V \Gamma_V}{(Q_2^2 - m_V^2)^2 + m_V^2 \Gamma_V^2} \Gamma(VV)$$

$$\Gamma(VV)|_{NL} = \Gamma^{SM}(VV) \times \left\{ c_V^2 - 2c_V \left[\frac{a_{VV}}{2} \left(1 - \frac{Q_1^2 + Q_2^2}{m_h^2} \right) + a_{V\partial V} \frac{Q_1^2 + Q_2^2}{m_h^2} \right] \right. \\ \left. + c_V a_{VV} \frac{\lambda(Q_1^2, Q_2^2, m_h^2) (1 - (Q_1^2 + Q_2^2)/m_h^2)}{\lambda(Q_1^2, Q_2^2, m_h^2) + 12 Q_1^2 Q_2^2 / m_h^4} \right\}$$

$$a_{VV} = c_{VV} \frac{m_h^2}{m_V^2}, \quad a_{V\partial V} = \frac{c_{V\partial V}}{2} \frac{m_h^2}{m_V^2}$$

$$\Gamma(VV)|_{SILH} = \Gamma^{SILH}(VV) + \Gamma^{SM}(VV) \frac{2}{|A_0^{SM}|^2} \text{Re} \left(A_0^{*SM} A_{ew}^{SM} \right)$$

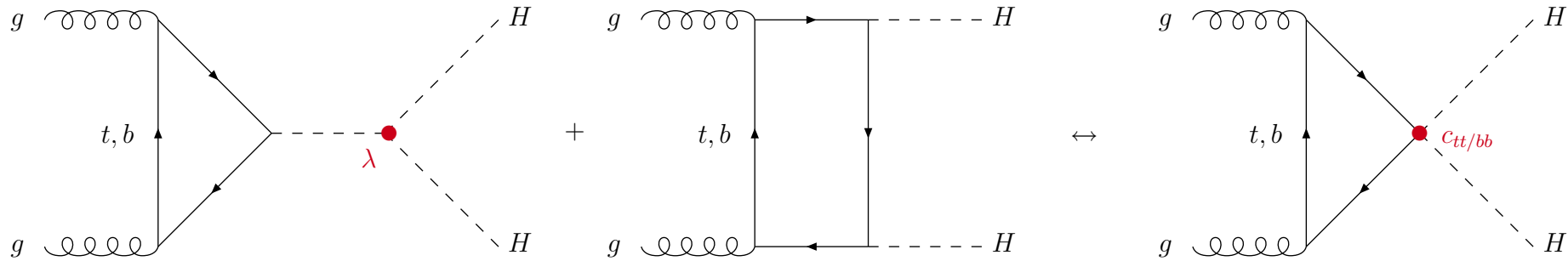
$$\Gamma^{SILH}(VV) = \Gamma^{SM}(VV) \times \left\{ 1 - \bar{c}_H - 2 \left[\frac{\bar{a}_{VV}}{2} \left(1 - \frac{Q_1^2 + Q_2^2}{m_h^2} \right) + \bar{a}_{V\partial V} \frac{Q_1^2 + Q_2^2}{m_h^2} \right] \right. \\ \left. + \bar{a}_{VV} \frac{\lambda(Q_1^2, Q_2^2, m_h^2) (1 - (Q_1^2 + Q_2^2)/m_h^2)}{\lambda(Q_1^2, Q_2^2, m_h^2) + 12 Q_1^2 Q_2^2 / m_h^4} \right\}$$

$$\Gamma^{SM}(VV) = \frac{\delta_V G_F m_h^3}{16\sqrt{2}\pi} \sqrt{\lambda(Q_1^2, Q_2^2, m_h^2)} \left(\lambda(Q_1^2, Q_2^2, m_h^2) + \frac{12 Q_1^2 Q_2^2}{m_h^4} \right)$$

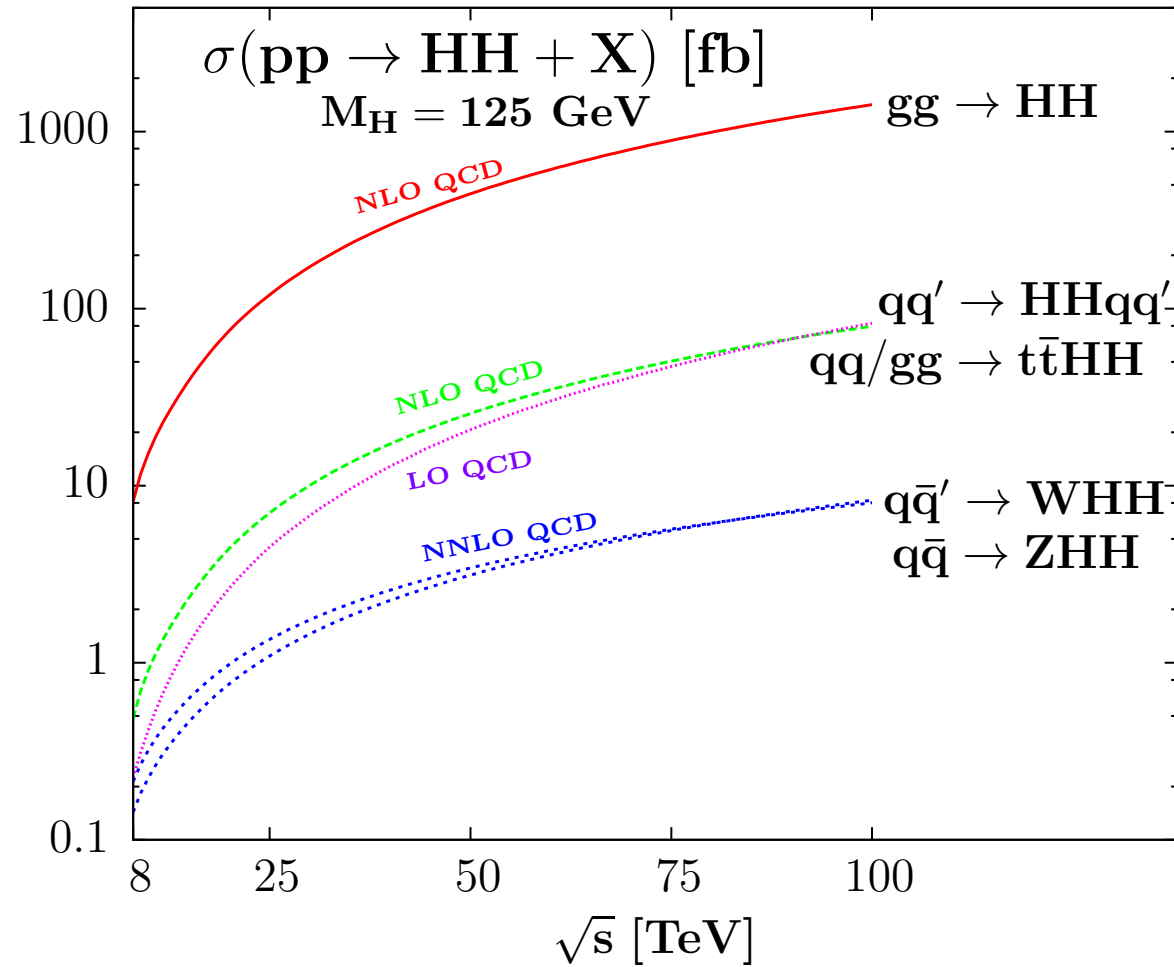
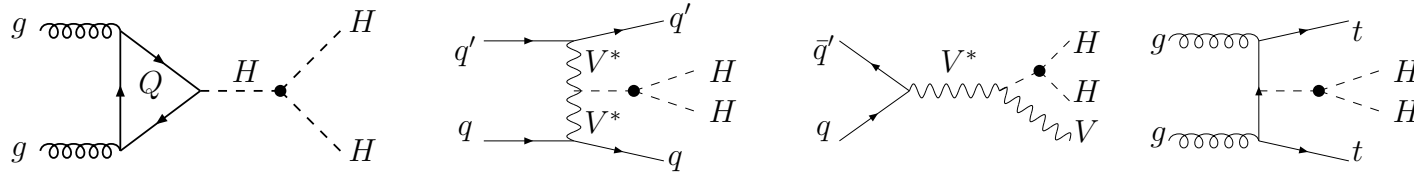
$$\bar{a}_{WW} = -2 \frac{m_h^2}{m_W^2} \bar{c}_{HW}, \quad \bar{a}_{ZZ} = -2 \frac{m_h^2}{m_Z^2} (\bar{c}_{HW} + \bar{c}_{HB} \tan^2 \theta_W - 4\bar{c}_\gamma \tan^2 \theta_W \sin^2 \theta_W)$$

$$\bar{a}_{W\partial W} = -2 \frac{m_h^2}{2m_W^2} (\bar{c}_W + \bar{c}_{HW}), \quad \bar{a}_{Z\partial Z} = -2 \frac{m_h^2}{2m_Z^2} (\bar{c}_W + \bar{c}_{HW} + (\bar{c}_B + \bar{c}_{HB}) \tan^2 \theta_W)$$

(ii) anomalous Higgs couplings [e.g. composite Higgs]



- threshold region: sensitive to λ
large M_{HH} : sensitive to $c_{tt/bb}$ [e.g. boosted Higgs pairs]
- similar effects for other Higgs pair production modes



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