

Rare K and B Decays in a Warped Extra Dimension with Custodial Protection

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Outline

1. Motivations for WED:

- ◆ Addressing Gauge Hierarchy Problem
- ◆ Natural Generation of Hierarchies in Masses and Mixings
- ◆ ...

2. Randall-Sundrum Scenario:

- ◆ The Model analyzed
- ◆ New Features in the Flavour Sector
- ◆ Rare Decays of B and K mesons: Theoretical Analysis
- ◆ Rare Decays of B and K mesons: Numerical Analysis

3. Conclusions

M. Blanke, A. J. Buras, B. Duling, K. Gemmler, S. Gori,
Rare K and B decays in a Warped Extra Dimension with Custodial Protection, in preparation

Gauge Hierarchy Problem & its Solution

♦ Gauge Hierarchy Problem in 3 sentences:

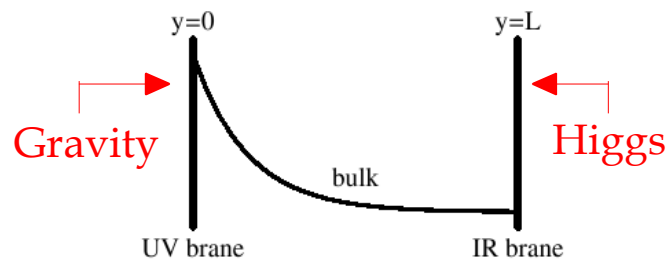
- I. **Huge hierarchy** between the fundamental gravity scale M_{pl} & the EW scale Λ_{EWSB}
- II. **Tremendous fine-tuning** required to keep $\Lambda_{EWSB} \sim 1 \text{ TeV}$
- III. Even if $\frac{\Lambda_{EWSB}}{M_{pl}} \approx 10^{-16}$ is imposed at tree-level, loop corrections push $\Lambda_{EWSB} \sim M_{pl}$

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◆ How to solve it in WED Contexts?



Solution to the 5D Einstein equations

$$ds^2 = e^{-2ky} \eta_{\mu\nu} dx^\mu dx^\nu - dy^2, \quad 0 \leq y \leq L$$

Warped-Factor

- I. Warped-factor along extra dimension leads to: $\Lambda_{eff}(y) = e^{-ky} \Lambda_{fund}$
- II. With $\Lambda_{fund} \sim O(M_{pl})$ only a moderate hierarchy is required to obtain $\Lambda_{eff}(IR \text{ brane}) \approx O(1 \text{ TeV})$

$$kL \approx 30$$

III. fundamental gravity scale however still given by M_{pl}

If Higgs lives on the IR brane, gauge hierarchy problem does not arise!

Flavour Problem & its Solution (1)

◆ Experiments tell us:

I. quarks and charged leptons have

$$m_e \approx 0.5 \text{ MeV} , m_\tau \approx 1800 \text{ MeV}, \dots$$

$$m_u \approx 2.5 \text{ MeV} , m_t \approx 170 \text{ GeV}, \dots$$

◆ ... and the theory:

III. at the same time CKM picture describes data surprisingly well

hierarchies

II. also CKM mixing between quark

$$|V_{ud}| \approx 1 , |V_{us}| \approx 0.226$$

$$|V_{cb}| \approx 0.041 , |V_{ub}| \approx 0.0038$$

SM Yukawa couplings have to exhibit an extremely hierarchical structure, **why?**

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→ How to solve it in WED Contexts?

◆ Preliminaries

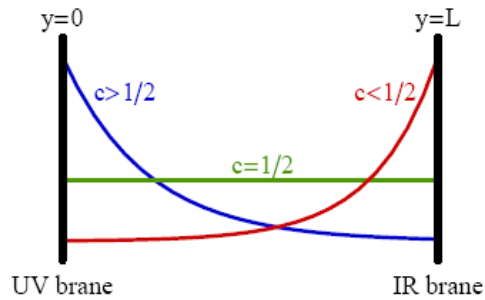
- Gauge fields and matter fields can propagate into the 5th dimension
- For each particle species, there is an infinite number of solutions:

Kaluza-Klein tower of particles

- Zero mode solutions (if existent) are identified with the SM particles (with BC (++)

Flavour Problem & its Solution (2)

◆ Zero Modes of Fermions:

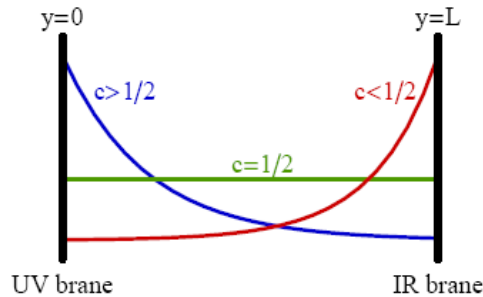


$$f^{(0)}(y, c) = \sqrt{\frac{(1-2c)kL}{e^{(1-2c)kL} - 1}} e^{(\frac{1}{2}-c)ky}$$

Strong dependence on bulk masses

Flavour Problem & its Solution (2)

Zero Modes of Fermions:



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Strong dependence on bulk masses

The Solution of the Flavour Problem:

I. 4D Yukawas in terms of shape functions:

$$Y_{ij} = \int_0^L \frac{dy}{L^{3/2}} \lambda_{ij} h(y) f_L^{(0)}(y, c^i) f_R^{(0)}(y, c^j)$$

5D Yukawas

λ_{ij} assumed to be **anarchical** and O(1)

Higgs localized on the IR brane: $h(y) = \sqrt{2(\beta-1)kL} e^{kL} e^{\beta k(y-L)}$, $\beta > 1$

II. Result: slightly different c parameters of O(1) lead to a large hierarchy in Y_{ij}

Hierarchy of quark masses and mixings explained by a **purely geometrical approach!** 😊

BUT 😞
Still missing a theory for the bulk masses

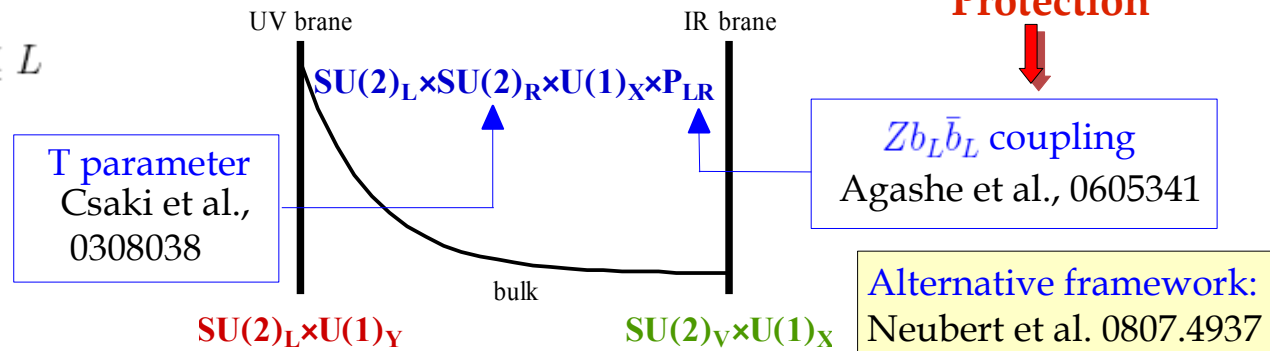
Numerical example:
 $c_1 = 0.66$, $c_2 = 0.59$, $c_3 = 0.41$
 $Y_1 = 0.0017$, $Y_2 = 0.017$, $Y_3 = 0.42$

Definition of the Model

1. Symmetry group and geometric structure:

$$ds^2 = e^{-2ky} \eta_{\mu\nu} dx^\mu dx^\nu - dy^2, \quad 0 \leq y \leq L$$

$$\begin{cases} e^{-kL} \approx 10^{-16} \\ M_{KK} \approx 2.45ke^{-kL} \approx \mathbf{2.45 TeV} \end{cases}$$

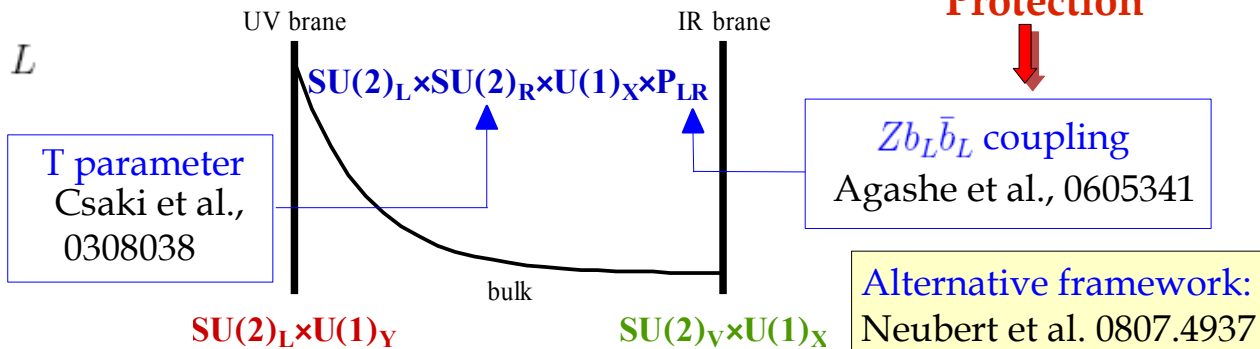


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2. Field content:

◆ Gauge bosons:

I. Gauge eigenstates:

$$\begin{aligned} &W_{L\mu}^a(++), \quad B_\mu(++), \quad G_\mu^c(++), \\ &W_{R\mu}^b(-+), \quad Z_{X\mu}(-+) \\ &a = 1, 2, 3; \quad b = 1, 2; \quad c = 1, \dots, 8 \end{aligned}$$

II. Mass eigenstates:

$$\begin{aligned} &W_\mu^\pm, \quad W_{H\mu}^\pm, \quad \tilde{W}_\mu^\pm \\ &A_\mu, \quad A_\mu^{(1)} \\ &Z_\mu, \quad Z_{H\mu}, \quad Z'_\mu \\ &G_\mu^{(0)}, \quad G_\mu^{(1)} \end{aligned}$$

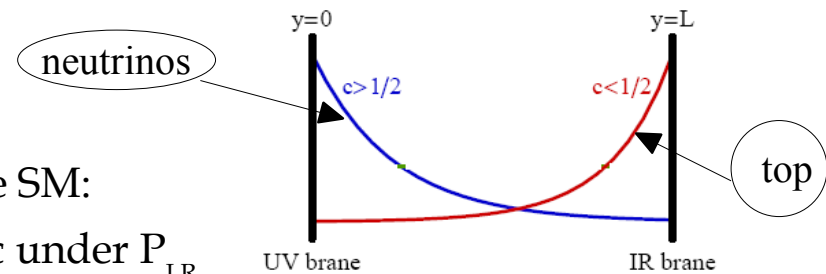
Gauge bosons of the SM

◆ Higgs boson:

- I. Bidoublet of $SU(2)_L \times SU(2)_R$
- II. EWSB mechanism is not specified

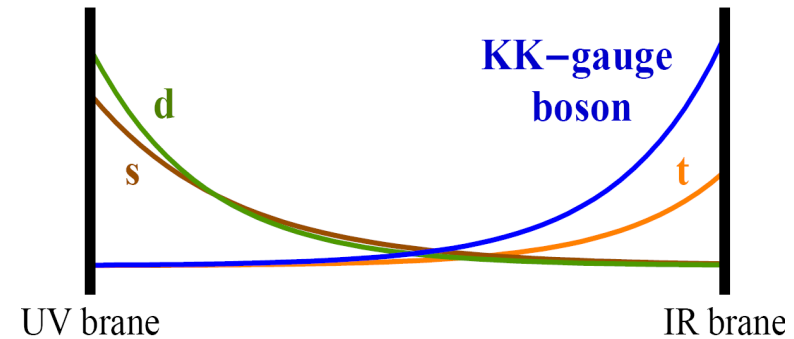
◆ Fermions:

- I. Three different representations of $SU(2)_L \times SU(2)_R$
- II. Different localizations in the bulk of the fermions of the SM:
- III. Left down quarks (all three generations) are symmetric under P_{LR}



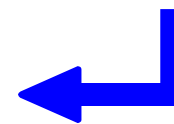
Non Universality & FCNC at Tree Level

- ◆ KK tower of heavy gauge bosons
...that are all localized towards the IR brane



- ◆ Their couplings to SM fermions are **non-universal**
...because couplings to SM fermions depend on their localization

$$\Delta_{L,R} \propto \int_0^L dy e^{ky} \left[f_{L,R}^{(0)}(y, c_\Psi^i) \right]^2 g(y)$$



Rotation to mass eigenstates:

non universalities



off-diagonal terms

Flavour Changing Neutral Currents at Tree Level

$$\Delta_{L,R} \sim U^\dagger \begin{pmatrix} \clubsuit & & \\ & \spadesuit & \\ & & \heartsuit \end{pmatrix} U$$

Non universalities

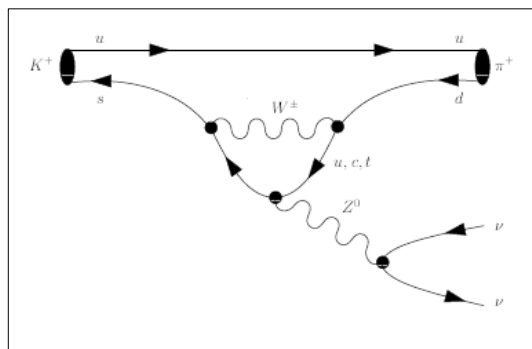
- ◆ New sources of flavour and CP violation beyond CKM: **model is non-MFV**

Rare Decays: some Theoretical Aspects

Example:

$$K^+ \rightarrow \pi^+ \nu \bar{\nu}$$

Standard Model



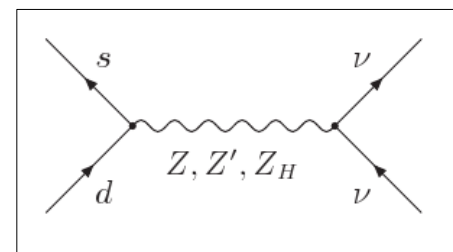
first at one loop level

I. The effective Hamiltonian:

$$\mathcal{H}_{eff}^{SM} \propto V_{ts}^* V_{td} X_{SM} (\bar{s} \gamma_\mu (1 - \gamma_5) d) (\bar{\nu} \gamma_\mu (1 - \gamma_5) \nu)$$

Only operator involved

Warped Extra Dimensions



additional diagrams at tree level

I. Modification of the coefficient of the SM operator

II. New operator is induced:

$$\mathcal{H}_{eff}^{new} \propto V_{ts}^* V_{td} X^V (\bar{s} \gamma_\mu d) (\bar{\nu} \gamma_\mu (1 - \gamma_5) \nu)$$

III. Main contributions from the coupling of Z to right handed down quarks

Rare Decays: K physics vs B physics

$$s \rightarrow d \bar{\nu} \nu \quad vs \quad (b \rightarrow d \bar{\nu} \nu \quad \vee \quad b \rightarrow s \bar{\nu} \nu)$$

Effective Hamiltonian:

$$\mathcal{H}_{eff}^{tot} \propto V_{tq_1}^* V_{tq_2} (X_{SM} + X_{q_1, q_2}^{V-A}) (\bar{q}_1 \gamma_\mu (1 - \gamma_5) q_2) (\bar{\nu} \gamma_\mu (1 - \gamma_5) \nu) + \\ + V_{tq_1}^* V_{tq_2} X_{q_1, q_2}^V (\bar{q}_1 \gamma_\mu q_2) (\bar{\nu} \gamma_\mu (1 - \gamma_5) \nu)$$

$$q_1 \rightarrow q_2 \bar{\nu} \nu$$

where the new functions:

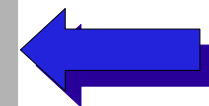
$$X_{q_1, q_2}^{V-A, V} \propto \frac{1}{\lambda_t^{(q)}} F(\Delta_L^{\nu\nu}, \Delta_{L,R}^{q_1, q_2})$$

$$\text{K meson: } \lambda_t^{(q)} = V_{ts}^* V_{td} \approx 4 \cdot 10^{-4}$$

$$\text{B mesons: } \lambda_t^{(q)} = V_{tb}^* V_{tq} \approx 10^{-2}, \quad q = d, s$$

Main Messages:

- I. Non universalities
- II. Expected: bigger contributions of the new physics in the K sector



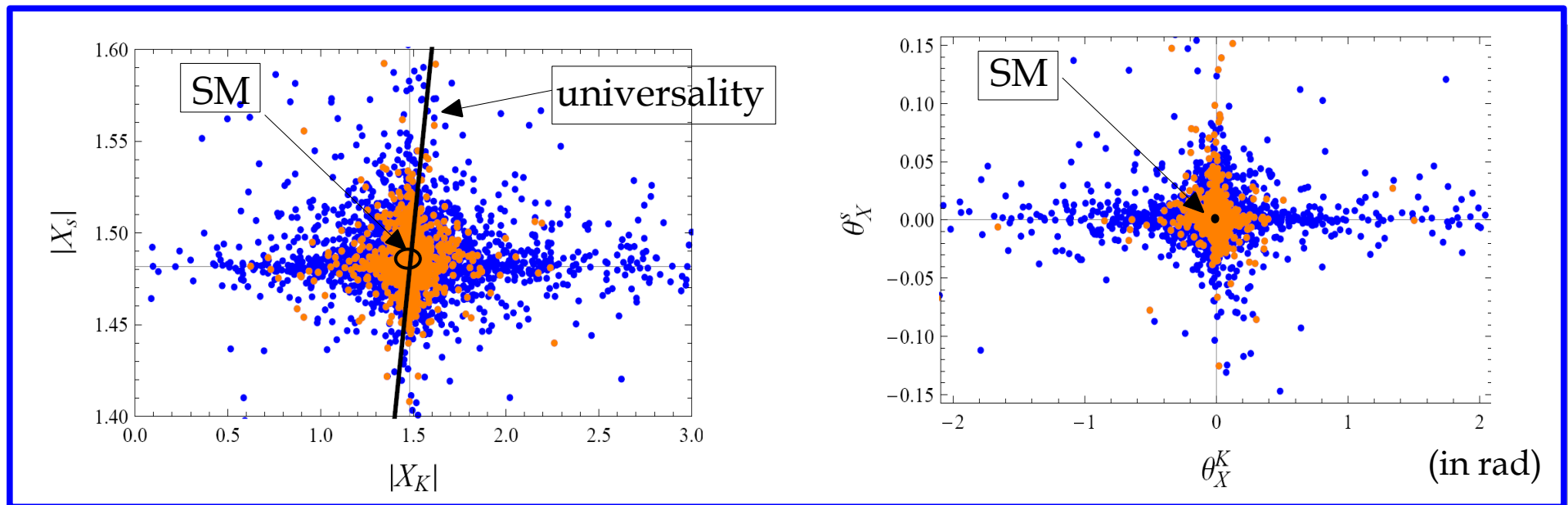
Non universality & New Sources of CP Violation

- 1) Points which satisfy all the $\Delta F=2$ constraints;
- 2) Also small fine tuning is required

real function

$$X_K = X_{SM} + X_{sd}^{V-A} + X_{sd}^V = |X_K| e^{i\theta_X^K}$$

$$X_s = X_{SM} + X_{bs}^{V-A} + X_{bs}^V = |X_s| e^{i\theta_X^s}$$



I. Deviation from the universality:

$$|X_K| \neq |X_s|$$

II. Bigger new physics contribution in X_K

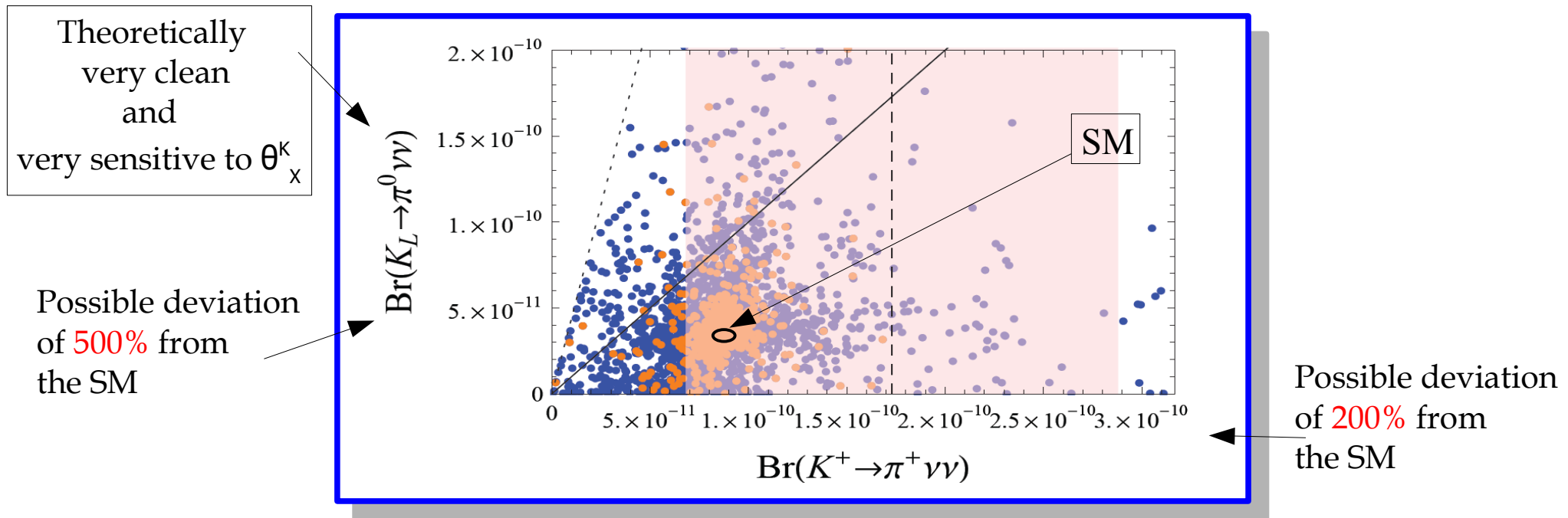
I. New phases:



new sources of CP violation

II. Bigger contribution in θ_X^K

Rare Decays of K mesons...



◆ Values predicted by the SM:

$$Br(K^+ \rightarrow \pi^+ \nu \bar{\nu}) = (8.4 \pm 0.8) 10^{-11}$$

$$Br(K_L \rightarrow \pi^0 \nu \bar{\nu}) = (2.9 \pm 0.4) 10^{-11}$$

◆ Experimental bounds:

$$Br(K^+ \rightarrow \pi^+ \nu \bar{\nu}) = (17.3^{+11.5}_{-10.5}) 10^{-11}$$

$$Br(K_L \rightarrow \pi^0 \nu \bar{\nu}) < 2.1 \cdot 10^{-7} \quad (90\% CL)$$

◆ Some Observations:

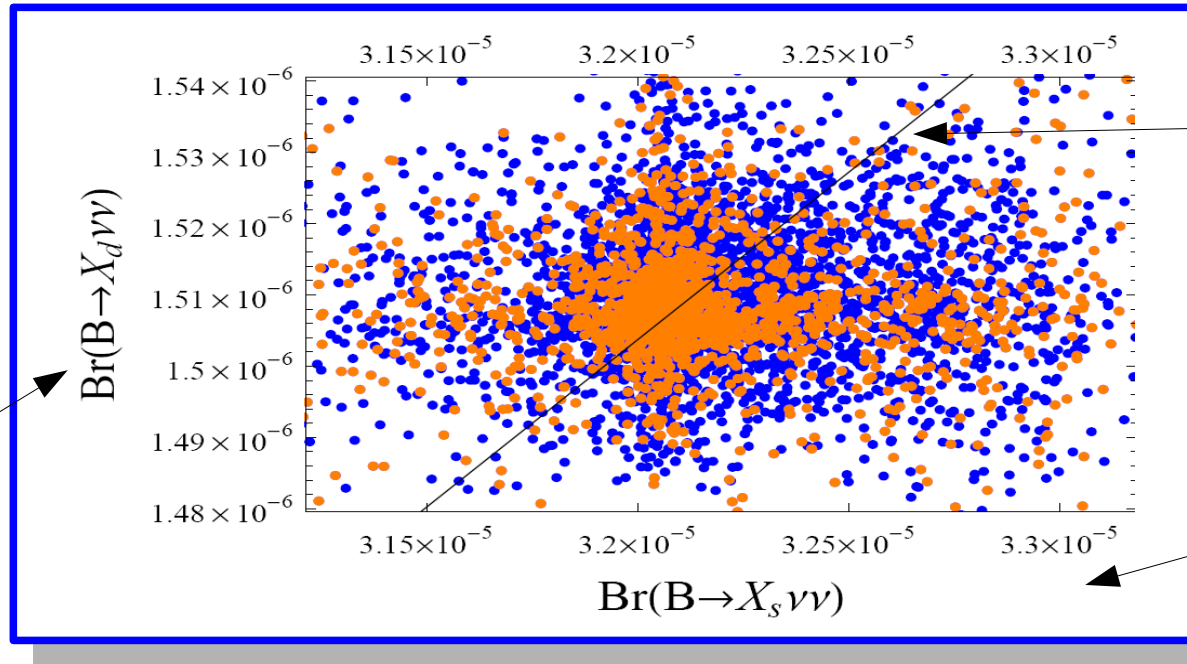
- I. It's possible to have **simultaneously big contributions** for both these branching ratios
- II. The most part of the points stays **in the experimental range** for $K^+ \rightarrow \pi^+ \nu \bar{\nu}$

...and Rare Decays of B mesons

Inclusive Decay

Possible deviation of **3%** from the SM

B decays not so sensitive to new physics



Model with **MFV**

Possible deviation of **5%** from the SM

◆ Values predicted by the SM:

$$Br(B \rightarrow X_s \nu \bar{\nu}) \approx 3.2 \cdot 10^{-5}$$

$$Br(B \rightarrow X_d \nu \bar{\nu}) \approx 1.5 \cdot 10^{-6}$$

◆ Experimental bounds:

$$Br(B \rightarrow X_s \nu \bar{\nu}) < 64 \cdot 10^{-5}$$

$$Br(B \rightarrow X_d \nu \bar{\nu}) < ??$$

◆ Some Observations:

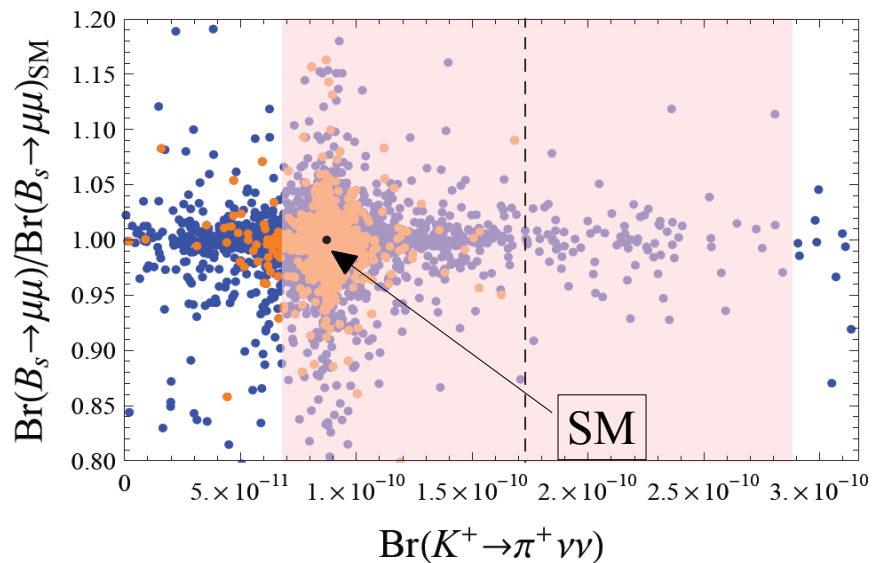
I. Very **clean correlation** between the two observables in models with MFV

II. In general: $\frac{Br(B \rightarrow X_d \nu \bar{\nu})}{Br(B \rightarrow X_s \nu \bar{\nu})} = \frac{|V_{td}|^2}{|V_{ts}|^2} P$ where $P \equiv \frac{|X_d^{V-A} + X_d^V/2|^2 + |X_d^V/2|^2}{|X_s^{V-A} + X_s^V/2|^2 + |X_s^V/2|^2}$

III. In models with MFV, $P=1$ (universality); in WED we have deviations

Correlations

B physics vs K physics



I. SM prediction

$$Br(B_s \rightarrow \mu^+ \mu^-) = (3.35 \pm 0.32) \cdot 10^{-9}$$

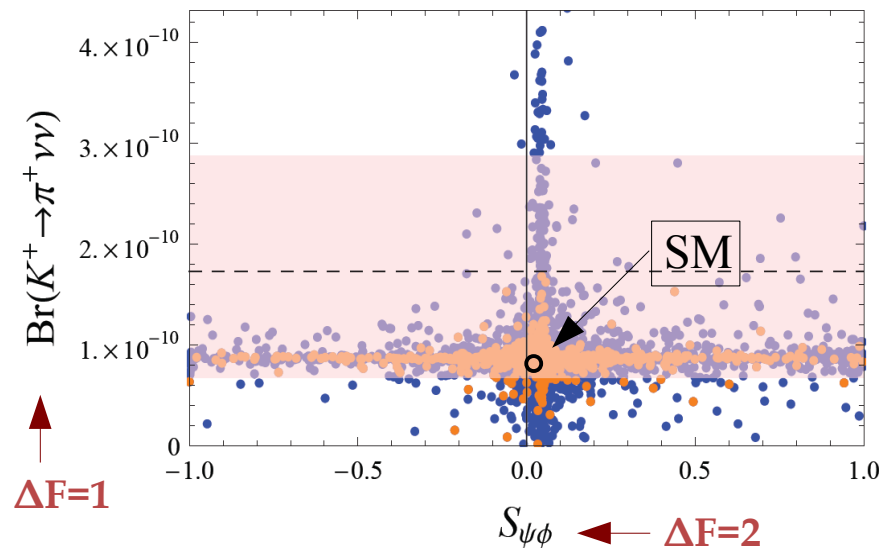
II. Measurements

$$Br(B_s \rightarrow \mu^+ \mu^-) < 4.7 \cdot 10^{-8}$$

For these two decays:

Possible deviations of **15%** in the B system;
Possible deviations of **200%** in the K system

$\Delta F=1$ vs $\Delta F=2$ observables




I. SM prediction

$$S_{\psi\phi} \approx 0.04, \quad Br(K^+ \rightarrow \pi^+ \nu \bar{\nu}) = (8.4 \pm 0.8) 10^{-11}$$

II. Measurements

$$S_{\psi\phi} \approx 0.4, \quad Br(K^+ \rightarrow \pi^+ \nu \bar{\nu}) = (17.3^{+11.5}_{-10.5}) 10^{-11}$$

Difficult to obtain **simultaneously large deviations** from the SM 
for both the observables

See next talk...

Conclusions

Warped Extra Dimension with custodial Protection shows:

♦ Elegant solutions for:

- I. Gauge Hierarchy Problem;
- II. Flavour Problem;
- III. ...



Testability at LHC since
 $M_{KK} \approx (2 - 3) TeV$

♦ In the Flavour Sector:

• In Rare Decays:

- I. Big Effects in **K physics** ← Testable Deviations from the SM
- II. Smaller Effects in **B physics** ← Challenging Testability of Deviations from Models with MFV
- III. Interesting Correlations between different Observables with $\Delta F=1$

• In $\Delta F=2$ Observables:

See next talk

...and Correlations between
 $\Delta F=2$ and $\Delta F=1$ Observables