

Family symmetries.

G.G.Ross, Valencia, December 2008



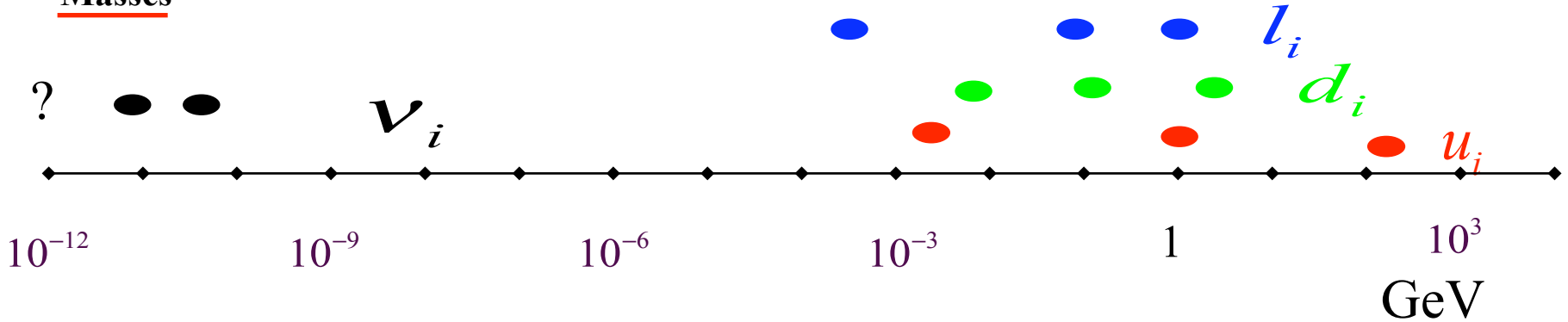
Family symmetries.

G.G.Ross, Valencia, December 2008

- $q \leftrightarrow l$ Neutrino v/s quark masses and mixing?
See-saw with sequential domination
- Symmetry GUT
Family Abelian, Non-Abelian, Discrete, String?
- Tri-bi-maximal mixing Non-Abelian discrete symmetry

DATA :

Masses



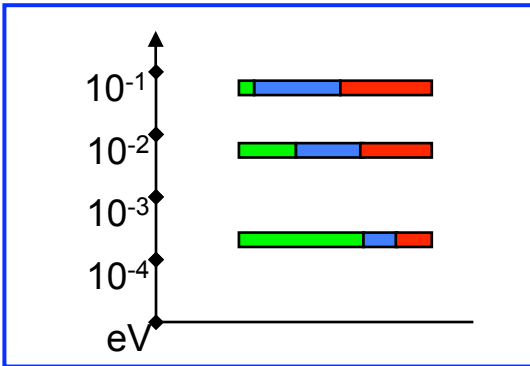
Mixing

Quarks

Leptons

$$V_{CKM} \approx \begin{pmatrix} 1 & 0.218-0.224 & 0.002-0.005 \\ 0.218-0.224 & 1 & 0.032-0.048 \\ 0.004-0.015 & 0.03-0.048 & 1 \end{pmatrix}$$

$$V_{MNS} = \begin{pmatrix} 0.79-0.88 & 0.48-0.61 & < 0.2 \\ 0.27-0.49 & 0.45-0.71 & 0.52-0.82 \\ 0.28-0.5 & 0.51-0.65 & 0.57-0.81 \end{pmatrix}$$



$$\approx \begin{pmatrix} \sqrt{\frac{2}{3}} & \frac{1}{\sqrt{3}} & \sim 0 \\ -\frac{1}{\sqrt{6}} & \frac{1}{\sqrt{3}} & -\frac{1}{\sqrt{2}} \\ -\frac{1}{\sqrt{6}} & \frac{1}{\sqrt{3}} & \frac{1}{\sqrt{2}} \end{pmatrix}$$

Bi-Tri Maximal
Mixing ...
Discrete
Non Abelian
Structure?

DATA :

$$L_{Yukawa} = Y_{ij}^u Q^i u^{c,j} H + Y_{ij}^d Q^i d^{c,j} \bar{H}$$

$$M_{ij}^u = Y_{ij}^u \langle H^0 \rangle \quad M_{ij}^d = Y_{ij}^d \langle \bar{H}^0 \rangle$$

$$M^u = V_L^\dagger \underline{M_{Diag}^u} V_R \quad \underline{V_{CKM}} = V_L^\dagger U_L$$

$$M^d = U_L^\dagger \underline{M_{Diag}^d} U_R$$

The data for quarks is *consistent* with a very symmetric structure :

$$\frac{M^{d,u}}{m^{b,t}} \simeq \begin{pmatrix} \langle \epsilon^4 & \epsilon^3 & \epsilon^3 \\ \epsilon^3 & \epsilon^2 & \epsilon^2 \\ \epsilon^3 & \epsilon^2 & 1 \end{pmatrix} \quad \begin{array}{l} \epsilon^d = 0.15 \\ \epsilon^u = 0.05 \end{array}$$

O(1) coefficients suppressed

$q \leftrightarrow l$ symmetry?

Charged leptons are consistent with a similar form

$$\frac{M^{d,l,u}}{m^{b,\tau,t}} \simeq \begin{pmatrix} < \epsilon^4 & \epsilon^3 & \epsilon^3 \\ \epsilon^3 & a\epsilon^2 & a\epsilon^2 \\ \epsilon^3 & a\epsilon^2 & 1 \end{pmatrix} \quad \begin{aligned} \epsilon^d &= 0.15, & a^d &= 1 \\ \epsilon^l &= 0.15, & a^l &= -3 \\ \epsilon^u &= 0.05, & a^u &= 1 \end{aligned}$$

Symmetry 1.

GUT relations

e.g. $SU(4) \subset SO(10)$

$$\Psi_\alpha = \begin{pmatrix} d \\ d \\ d \\ l \end{pmatrix}$$

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$$\text{Det}(M^l) = \text{Det}(M^d) |_{M_X}$$

$$\frac{M^{d,l}}{m_3} = \begin{pmatrix} < \varepsilon^4 & \varepsilon^3 & \varepsilon^3 \\ \varepsilon^3 & a\varepsilon^2 & a\varepsilon^2 \\ \varepsilon^3 & a\varepsilon^2 & 1 \end{pmatrix}$$

$\bar{\Psi} \quad \Psi_\alpha$

$$\frac{m_b}{m_\tau}(M_X) = 1$$

$$\varepsilon^d = 0.15, \quad a^d = 1$$

$$\varepsilon^l = 0.15, \quad a^l = -3$$

Symmetry 1.

GUT relations

e.g. $SU(4) \subset SO(10)$

$$\psi_\alpha = \begin{pmatrix} d \\ d \\ d \\ l \end{pmatrix}$$

$$\langle \Sigma_{45} \rangle$$

||

$$\bar{\psi}^{-\alpha} \begin{pmatrix} 1 & & & \\ & 1 & & \\ & & 1 & \\ & & & -3 \end{pmatrix} \psi_\alpha$$

$$\frac{m_s}{m_\mu}(M_X) = \frac{1}{3}$$

Georgi Jarlskog

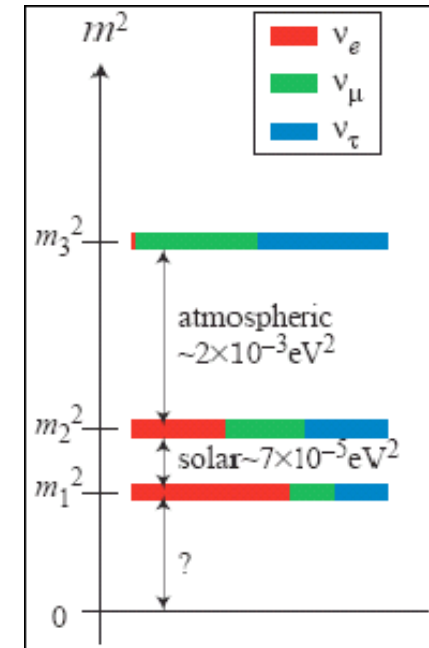
$$\frac{M^{d,l}}{m_3} = \begin{pmatrix} < \epsilon^4 & \epsilon^3 & \epsilon^3 \\ \epsilon^3 & a\epsilon^2 & a\epsilon^2 \\ \epsilon^3 & a\epsilon^2 & 1 \end{pmatrix}$$

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Neutrinos ???

$$L_{eff}^{\nu} = m_3 \bar{\phi}_{23}^i \nu_i \bar{\phi}_{23}^j \nu_j + m_2 \bar{\phi}_{123}^i \nu_i \bar{\phi}_{123}^j \nu_j$$

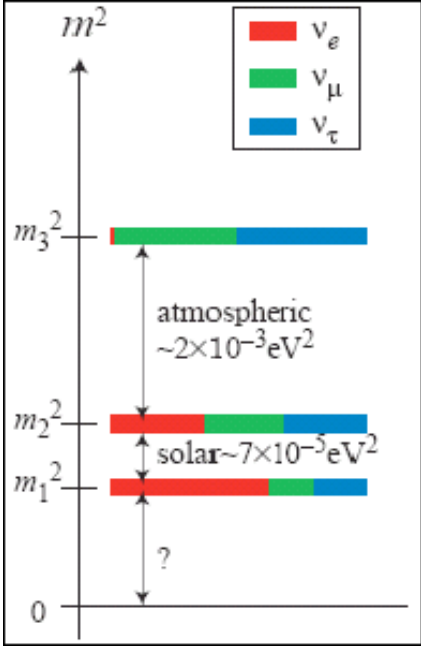
$$\langle \bar{\phi}_{23} \rangle^i = (0, -1, 1), \quad \langle \bar{\phi}_{123} \rangle^i = (1, 1, 1)$$



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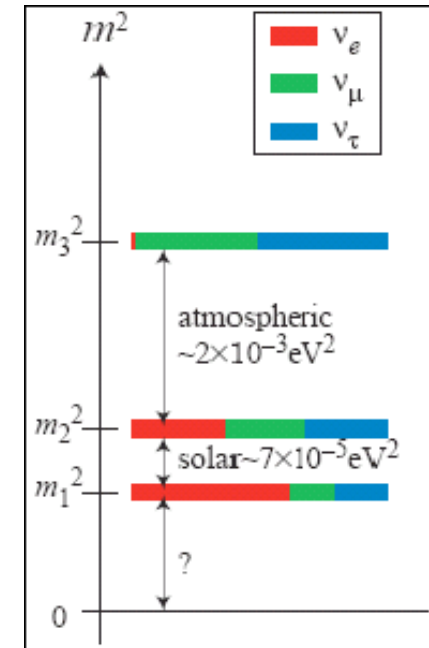
Can one have a unified description of quark, charged lepton **and** neutrinos?

c.f. $L_{Dirac}^{q,l} = m_3 \bar{\phi}_3^i \psi_i \bar{\phi}_3^j \psi_j^c + \dots$ $\langle \bar{\phi}_3 \rangle^i = (0, 0, 1) \quad ???$

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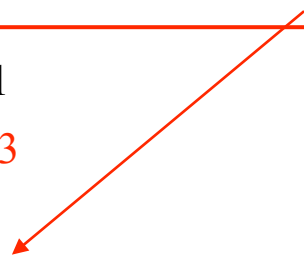
See-Saw

Quarks, charged leptons, neutrinos **can** have similar Dirac mass :

$$L_{Dirac}^{q,l,\nu} = \alpha \psi_i \bar{\phi}_3^i \psi_j^c \bar{\phi}_3^j + \beta \left(\psi_i \bar{\phi}_{123}^i \psi_j^c \bar{\phi}_{23}^j + \psi_i \bar{\phi}_{23}^i \psi_j^c \bar{\phi}_{123}^j \right) + \gamma \psi_i \bar{\phi}_{23}^i \psi_j^c \bar{\phi}_{23}^j \Sigma_{45} \quad \alpha > \beta$$

$$\frac{M^{Dirac}}{m_3} = \begin{pmatrix} \langle \epsilon^4 \rangle & \epsilon^3 + \epsilon^4 & -\epsilon^3 + \epsilon^4 \\ \epsilon^3 + \epsilon^4 & a\epsilon^2 + \epsilon^3 & -a\epsilon^2 + \epsilon^3 \\ -\epsilon^3 + \epsilon^4 & -a\epsilon^2 + \epsilon^3 & 1 \end{pmatrix}$$

$$\begin{aligned} \epsilon^d &= 0.15, & a^d &= 1 \\ \epsilon^l &= 0.15, & a^e &= -3 \\ \epsilon^u &= 0.05, & a^u &= 1 \\ \epsilon^v &= 0.05, & a^v &= 0 \end{aligned}$$

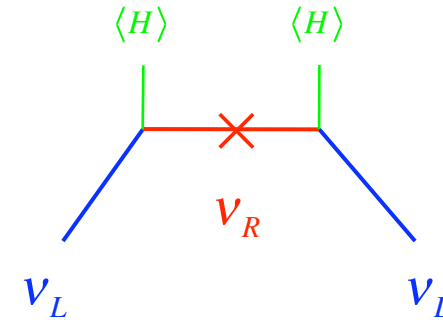


- “See-saw” with sequential domination

$$M_\nu = M_D^\nu M_M^{-1} M_D^{\nu T}$$

Minkowski
Gell-Mann,
Ramond,
Slansky;
Yanagida,
King

$$M_M \approx \begin{pmatrix} M_1 & & \\ & M_2 & \\ & & M_3 \end{pmatrix} \quad M_1 < M_2 \ll M_3$$



$$L_{Dirac}^\nu = \alpha \psi_i \phi_{i3}^{\bar{i}} \psi_j^c \phi_{j3}^{\bar{j}} + \beta \left(\psi_i \phi_{i23}^{\bar{i}} \psi_j^c \phi_{j23}^{\bar{j}} + \psi_i \phi_{i23}^{\bar{i}} \psi_j^c \phi_{j123}^{\bar{j}} \right) \quad \alpha > \beta$$

$$L_{eff}^\nu = \frac{\beta^2}{M_1} \psi_i \phi_{i123}^i \psi_j \phi_{j123}^j + \frac{\beta^2}{M_2} \psi_i \phi_{i23}^i \psi_j \phi_{j23}^j + \frac{(\alpha + \beta)^2}{M_3} \psi_i \phi_{i3}^i \psi_j \phi_{j3}^j$$

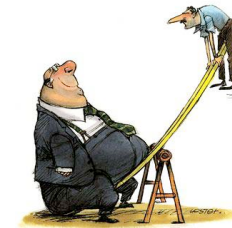
● “See-saw” with sequential domination

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$$M_M \approx \begin{pmatrix} M_1 & & \\ & M_2 & \\ & & M_3 \end{pmatrix}$$

$$M_1 \ll M_2 \ll M_3$$



small

$$L_{Dirac}^\nu = \alpha \psi_i \phi_{i3}^{\bar{i}} \psi_j^c \phi_{j3}^{\bar{j}} + \beta \left(\psi_i \phi_{i23}^{\bar{i}} \psi_j^c \phi_{j23}^{\bar{j}} + \psi_i \phi_{i23}^{\bar{i}} \psi_j^c \phi_{j123}^{\bar{j}} \right) \quad \alpha > \beta$$

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Symmetry 2.

Family symmetry

Non-Abelian family symmetry

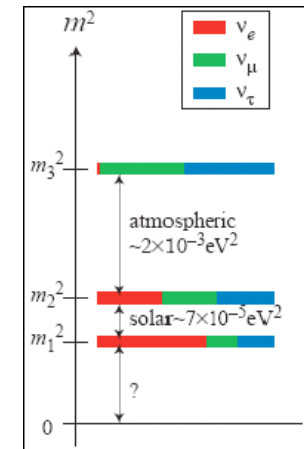
Promote ϕ_i to fields transforming under $SU(3)_{\text{family}}$

$$\frac{\phi_3}{M} = \begin{pmatrix} 0 \\ 0 \\ 1 \end{pmatrix}$$

$$\frac{\phi_{23}}{M} = \begin{pmatrix} 0 \\ 1 \\ -1 \end{pmatrix} \epsilon$$

$$\frac{\phi_{123}}{M} = \begin{pmatrix} 1 \\ 1 \\ 1 \end{pmatrix} \epsilon^2$$

Messenger mass



Symmetry 2.

Family symmetry

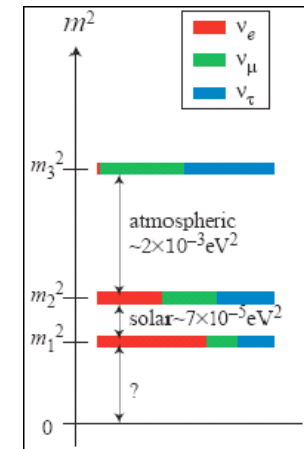
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Vacuum alignment ???

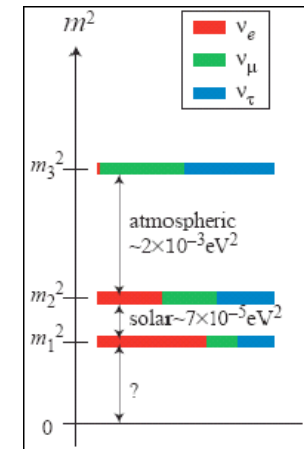
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Vacuum alignment ??? \Rightarrow **Discrete** non Abelian symmetry

List of models with discrete flavour symmetry

(incomplete, by symmetry)

S_3 : Pakvasa et al.(1978), Derman(1979), Ma(2000), Kubo et al.(2003), Chen et al.(2004), Grimus et al.(2005), Dermisek et al (2005), Mohapatra et al.(2006), Morisi(2006), Caravaglios et al.(2006), Haba et al(2006),...Mondragon

S_4 : Pakvasa et al.(1979), Derman(1979), Lee et al.(1994), Mohapatra et al.(2004),Ma(2006), Hagedorn et al.(2006), Caravaglios et al.(2006), Lampe(2007), Sawanaka(2007), ...

A_4 : Wyler(1979), Ma et al.(2001), Babu et al.(2003), Altarelli et al.(2005-8), He et al.(2006), Bazzocchi, Morisi, et al.(2007/8), King et al.(2007),...Luca

D_4 : Seidl(2003), Grimus et al.(2003/4), Kobayashi et al.(2005), ...

D_5 : Ma(2004), Hagedorn et al.(2006), ...

D_n : Chen et al.(2005), Kajiyama et al.(2006), Frampton et al.(1995/6,2000), Frigerio et al (2005), Babu et al.(2005), Kubo(2005),...

T' : Frampton et al.(1994,2007), Aranda et al.(1999,2000),Feruglio et al.(2007), Chen et al.(2007), ...

Δ_n : Kaplan et al.(1994), Schmaltz(1994), Chou et al.(1997), de Madeiros Varzielas et al(2006/7).

T_7 : Luhn et al.(2007).

Non Abelian discrete symmetries

e.g. $Z_3 \rtimes Z_n$

ϕ_i	$Z_3\phi_i$	$Z_n\phi_i$	
ϕ_1	$\rightarrow \phi_2$	$\rightarrow \alpha\phi_1$	$\alpha^n = 1$
ϕ_2	$\rightarrow \phi_3$	$\rightarrow \alpha^2\phi_2$	
ϕ_3	$\rightarrow \phi_1$	$\rightarrow \alpha^{-3}\phi_3$	

$$\Delta(3n^2) \quad n = 2, \quad \Delta(12) \equiv A_4$$

$$n = 3, \quad \Delta(27)$$

Non Abelian discrete symmetries

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ϕ_3	$\rightarrow \phi_1$	$\rightarrow \alpha^{-3}\phi_3$	

Choice of discrete symmetry

- Vacuum structure : $Z_3 \ltimes Z_n \rightarrow \begin{cases} Z_3, & \langle \phi \rangle = (1,1,1) / \sqrt{3} \quad \lambda > 0 \\ Z_n, & \langle \phi \rangle = (0,0,1) \quad \lambda < 0 \end{cases}$

$$V(\phi) = -m^2 \phi^{\dagger i} \phi_i + \dots + \lambda m^2 \phi^{\dagger i} \phi_i \phi^{\dagger i} \phi_i$$

Non Abelian discrete symmetries

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Choice of discrete symmetry

- Vacuum structure : $Z_3 \rtimes Z_n \rightarrow \left\{ \begin{array}{l} Z_3, \quad \langle \phi \rangle = (1, 1, 1) \\ Z_n, \quad \langle \phi \rangle = (0, 0, 1) \end{array} \right.$
- Quasi-degenerate neutrinos $D \subset SO(3)$ e.g. A_4 $m\psi^i \psi^i + ..$

GGR, Serna,
Varzielas

Non Abelian discrete symmetries

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Choice of discrete symmetry

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- Quasi-degenerate neutrinos $D \subset SO(3)$ e.g. A_4 $m\psi^i\psi^i + ..$
- GUT : $A_4 \times SU(5), S_3 \times E_6, ...$ $q \leftrightarrow l$ ✗ Altarelli et al, Caravaglios et al
 $\Delta(27) \times SO(10)$ ✓ GGR, Varzielas

A complete model

$$\Delta(27) \otimes SO(10) \otimes G \quad (G = R \otimes U(1))$$

Varzielas, GGR

- $\psi_i^c, \psi_i \in (16, 3) \Rightarrow$ No mass while SU(3) unbroken
- Spontaneous symmetry breaking

$$\begin{array}{cccc} \bar{\phi}_3^i & \bar{\phi}_{23}^i & \bar{\phi}_{123}^i & H_{45} \\ (1, \bar{3}) & (1, \bar{3}) & (1, \bar{3}) & (45, 1) \end{array}$$

c.f. Georgi-Jarskog

$$\begin{pmatrix} 0 \\ 0 \\ 1 \end{pmatrix}, \quad \begin{pmatrix} 0 \\ -1 \\ 1 \end{pmatrix} \varepsilon M, \quad \begin{pmatrix} 1 \\ 1 \\ 1 \end{pmatrix} \varepsilon^2 M, \quad M$$

$$P_Y = \frac{1}{M^2} \bar{\phi}_3^i \psi_i \bar{\phi}_3^j \psi_j^c H + \frac{1}{M^3} \bar{\phi}_{23}^i \psi_i \bar{\phi}_{23}^j \psi_j^c H H_{45} + \frac{1}{M^2} \bar{\phi}_{23}^i \psi_i \bar{\phi}_{123}^j \psi_j^c H + \frac{1}{M^2} \bar{\phi}_{123}^i \psi_i \bar{\phi}_{23}^j \psi_j^c H$$

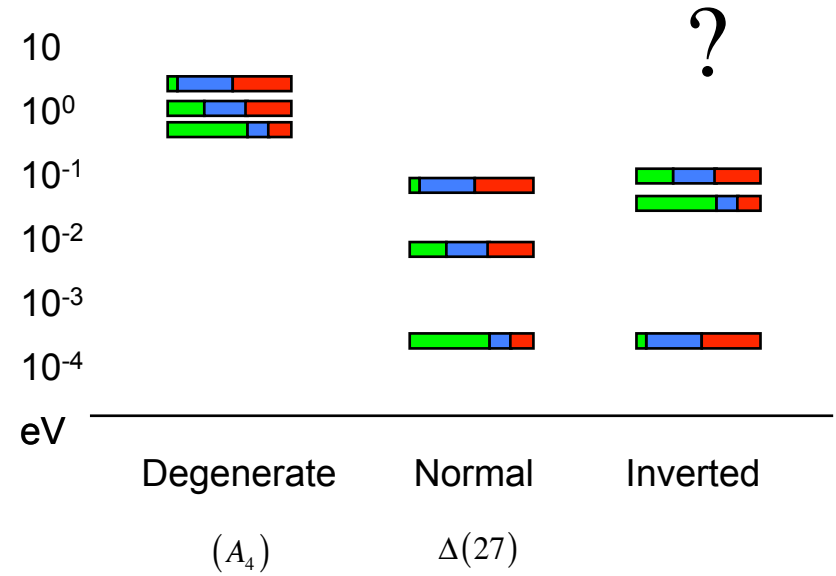
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Neutrino Parameters

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- Tri-bi-maximal mixing

Degenerate and normal spectra favoured



Neutrino Parameters

● Tri-bi-maximal mixing

Degenerate and normal spectra favoured

● Mixing angles

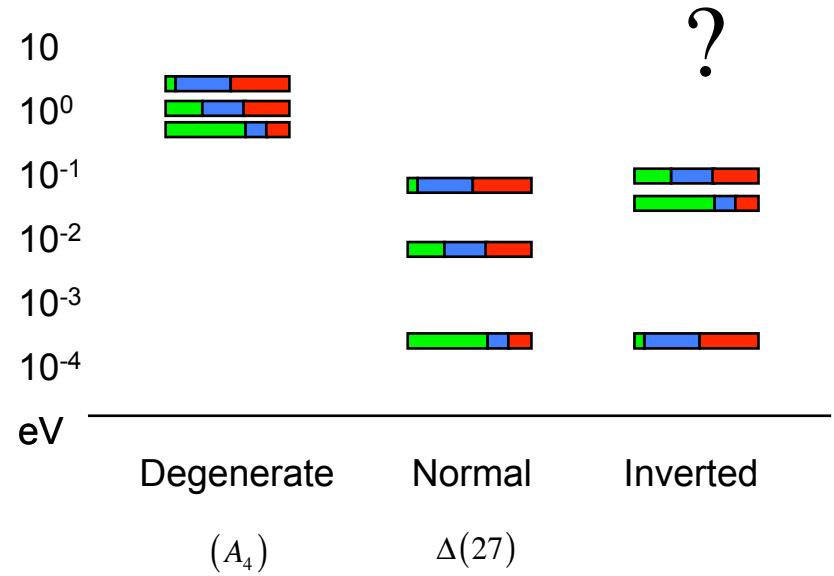
$$\sin^2 \theta_{12} \approx \frac{1}{3} \pm 0.03$$

$$\sin^2 \theta_{23} \approx \frac{1}{2} \pm 0.03$$

$$\sin \theta_{13} \approx \sqrt{\frac{m_e}{2m_\mu}} = 0.053 \pm 0.05 \quad (3 \pm 3^\circ)$$

From charged lepton mixing

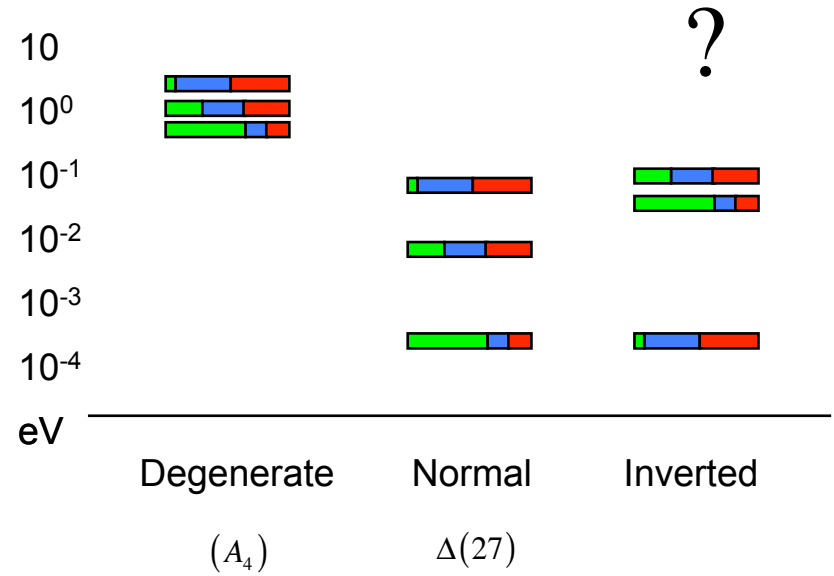
King; GGR, Varzielas



Neutrino Parameters

● Tri-bi-maximal mixing

Degenerate and normal spectra favoured

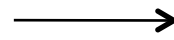


● Mixing angles

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$$\sin \theta_{13} \approx \sqrt{\frac{m_e}{2m_\mu}} = 0.053 \pm 0.05 \quad (3 \pm 3^\circ)$$



$$\theta_{12} + \frac{1}{\sqrt{2}} \frac{\theta_c}{3} \cos(\delta - \pi) \approx 35.26 \pm 2^\circ$$

Antusch, King

From charged lepton mixing

King; GGR, Varzielas

Neutrino Parameters

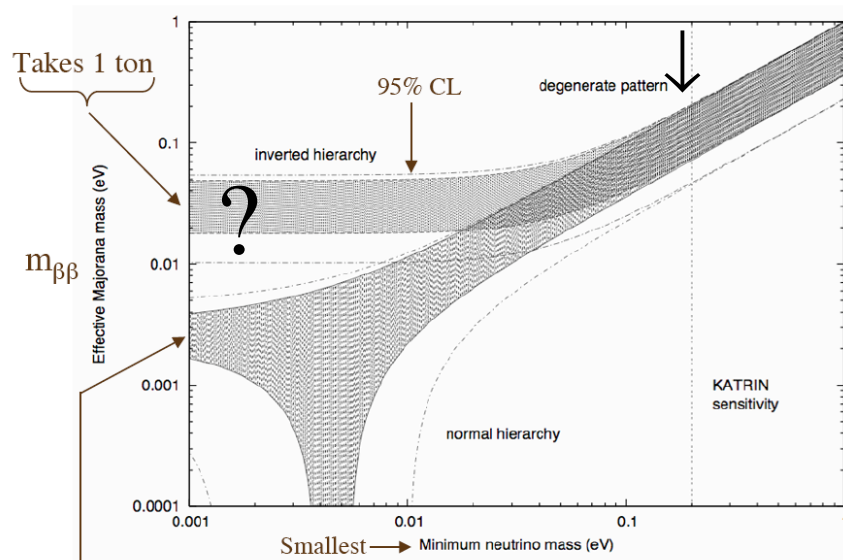
- Tri-bi-maximal mixing

Degenerate and normal spectra possible

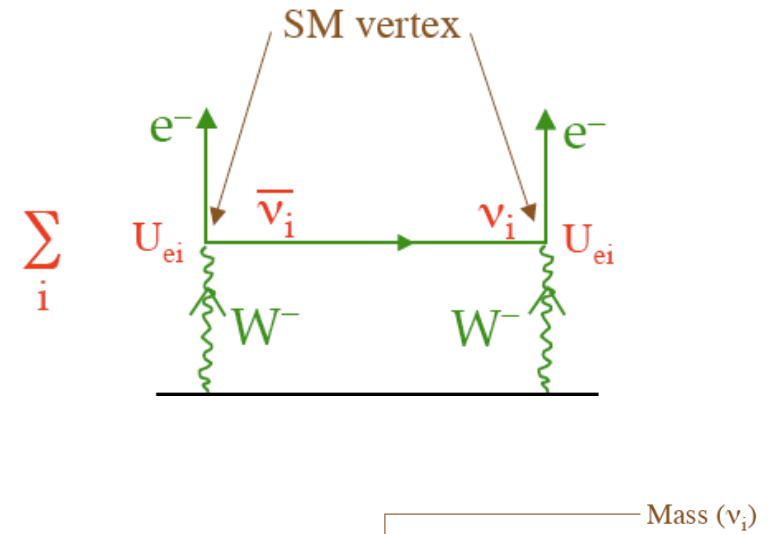
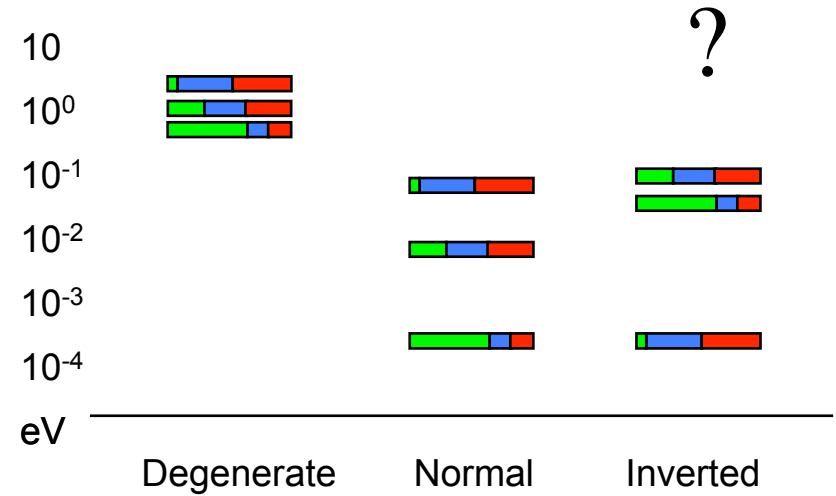
- Neutrinoless double β decay

(3 light neutrinos)

... Ivo's talk



$m_{\beta\beta}$ For Each Hierarchy



$$\text{Amp}[0\nu\beta\beta] \propto \left| \sum m_i U_{ei}^2 \right| \equiv m_{\beta\beta}$$

Sequential see saw and (1,1) texture zero

$$\frac{M^{Dirac}}{m_3} = \begin{pmatrix} < \varepsilon^4 & \varepsilon^3 + \varepsilon^4 & -\varepsilon^3 + \varepsilon^4 \\ \varepsilon^3 + \varepsilon^4 & a\varepsilon^2 + \varepsilon^3 & -a\varepsilon^2 + \varepsilon^3 \\ -\varepsilon^3 + \varepsilon^4 & -a\varepsilon^2 + \varepsilon^3 & 1 \end{pmatrix}$$

(1,1) texture zero

Sequential see saw and (1,1) texture zero

Quark sector zero $\Rightarrow |V_{us}| = \left| \sqrt{\frac{m_d}{m_s}} - e^{i\delta} \sqrt{\frac{m_u}{m_c}} \right|$

(1,1) texture zero

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Sequential see saw and (1,1) texture zero

Lepton sector zero $\Rightarrow \sin^2 \theta_{23} \simeq \frac{1}{2}$, $\sin \theta_{13} \simeq \sqrt{\frac{m_e}{2m_\mu}}$

(1,1) texture zero

$$\frac{M^{Dirac}}{m_3} = \begin{pmatrix} < \varepsilon^4 & \varepsilon^3 + \varepsilon^4 & -\varepsilon^3 + \varepsilon^4 \\ \varepsilon^3 + \varepsilon^4 & a\varepsilon^2 + \varepsilon^3 & -a\varepsilon^2 + \varepsilon^3 \\ -\varepsilon^3 + \varepsilon^4 & -a\varepsilon^2 + \varepsilon^3 & 1 \end{pmatrix}$$

CP ?

$$\phi_A = \text{Arg}(Am_{1/2}^*), \quad \phi_B = \text{Arg}(Bm_{1/2}^*) < 10^{-2} ?$$

SUSY CP problem has a simple solution in familion models:

- CP invariant in underlying theory (string compactification)
- Spontaneously broken in **flavour changing** familion sector

Dine, Leigh, MacIntyre

GGR, Velasco Sevilla, Vives

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GGR, Velasco Sevilla, Vives

Soft masses

$$m_{3/2}^2 (\tilde{f}^{i\dagger} \tilde{f}_i + \tilde{f}^{i\dagger} \phi_{3i}^\dagger \tilde{f}_j \phi_3^j + \dots)$$

Characteristic pattern

CP ?

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GGR, Velasco Sevilla, Vives

Soft masses

$$m_{3/2}^2 (\tilde{f}^{i\dagger} \tilde{f}_i + \tilde{f}^{i\dagger} \phi_{3i}^\dagger \tilde{f}_j \phi_{3j}^j + \dots)$$

Characteristic pattern

Soft A-terms

Misalignment of Yukawa and A-terms?

$$P = Y_{ij} \tilde{Q}_i q_j^c H_a \propto \left(\frac{\theta}{M}\right)^{\alpha(i,j)} \tilde{Q}_i q_j^c H_a$$

$$A_{ij} Y_{ij} \tilde{Q}_i \tilde{q}_j^c H_a = (3 + \alpha(i,j)) m_{3/2} Y_{ij} \tilde{Q}_i \tilde{q}_j^c H_a$$

$$F_\phi = m_{3/2} \langle \phi \rangle, \text{ "natural" value}$$

GGR, Vives

Here $\alpha(i, j) = 2$ so Y and A **aligned** (with $F_\Sigma = 0$) - suppresses FCNC

CP ?

$$\phi_A = \text{Arg}(Am_{1/2}^*), \quad \phi_B = \text{Arg}(Bm_{1/2}^*) < 10^{-2} ?$$

SUSY CP problem has a simple solution in familon models:

- CP invariant in underlying theory (string compactification)

Dine, Leigh, MacIntyre

- Spontaneously broken in **flavour changing** familon sector

GGR, Velasco Sevilla, Vives

Soft masses

$$m_{3/2}^2 (\tilde{f}^{i\dagger} \tilde{f}_i + \tilde{f}^{i\dagger} \phi_{3i}^\dagger \tilde{f}_j \phi_3^j + \dots)$$

Characteristic pattern

Soft A-terms: CP violation and FCNC

EDMs

$$|\text{Im}(\delta_{LR}^u)_{11}| \lesssim 2 \times 10^{-8} \frac{A_0}{100 \text{ GeV}} \left(\frac{500 \text{ GeV}}{\langle \tilde{m}_u \rangle_{LR}} \right)^2 \left(\frac{\bar{\epsilon}}{0.15} \right)^3 \left(\frac{\epsilon}{0.05} \right)^3 \sin \phi_1 \leq 10^{-6} |_{\text{Expt}}$$

$$|\text{Im}(\delta_{LR}^d)_{11}| \sim 2 \times 10^{-7} \frac{A_0}{100 \text{ GeV}} \left(\frac{500 \text{ GeV}}{\langle \tilde{m}_d \rangle_{LR}} \right)^2 \left(\frac{\bar{\epsilon}}{0.15} \right)^6 \frac{10}{\tan \beta} \sin \phi_1 \leq 10^{-6} |_{\text{Expt}}$$

$$|\text{Im}(\delta_{LR}^\ell)_{11}| \sim 6 \times 10^{-8} \frac{A_0}{100 \text{ GeV}} \left(\frac{200 \text{ GeV}}{\langle \tilde{m}_e \rangle_{LR}} \right)^2 \left(\frac{\bar{\epsilon}}{0.15} \right)^6 \frac{10}{\tan \beta} \sin \phi_1 \leq 10^{-7} |_{\text{Expt}}$$

$\mu \rightarrow e\gamma$

$$|(\delta_{LR}^e)_{12}| \lesssim |(\delta_{LR}^\ell)_{12}| \sim 3 \times 10^{-5} \frac{A_0}{100 \text{ GeV}} \frac{(200 \text{ GeV})^2}{\langle \tilde{m}_l \rangle_{LR}^2} \frac{10}{\tan \beta} \left(\frac{\bar{\epsilon}}{0.15} \right)^4 \leq 10^{-5} |_{\text{Expt}}$$

Antusch, King, Malinsky, GGR

Summary

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- CP violation and g-2

SUSY models: sparticles have related mass structure

$(m^2 \phi_i^\dagger \phi_i \dots$ degeneracy split by small fermion mass related terms)

FCNC : L_i, B_i violation – close to present bounds

$(\mu \rightarrow e\gamma, \text{mercury EDM within a factor of } 10 \text{ of present limits})$

See saw parameters

Masses + Mixing Angles	6	(15)	(8)
CP violating phases	3	(6)	(3)

$$\kappa = Y_\nu^T M^{-1} Y_\nu$$

Insufficient information

$$P = Y_\nu^\dagger Y_\nu$$

RGE for sleptons (SUSY)

In *principle* all parameters determined by neutrino structure and radiative corrections
(assuming degenerate sleptons at GUT scale)

Davidson, Ibarra

Thermal Leptogenesis (Baryogenesis)

Fukugita, Yanagida

$$M_{\nu_R} > 10^9 \text{ GeV}, \quad T_{\text{Reheat}} > 10^{10} \text{ GeV} \quad (\text{hierarchical spectrum})$$

$$\cancel{CP} = \cancel{CP}(\delta, \phi', \phi_{i=1,2,3})$$

↙ $\beta\beta$ phase
← see saw phases
↘ ν factory phase

Sequential see saw and symmetry (texture zero) simplifies structure...