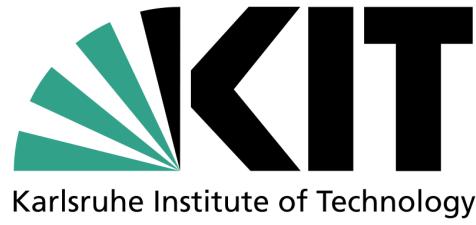


Computational aspects for three-loop DIS calculations

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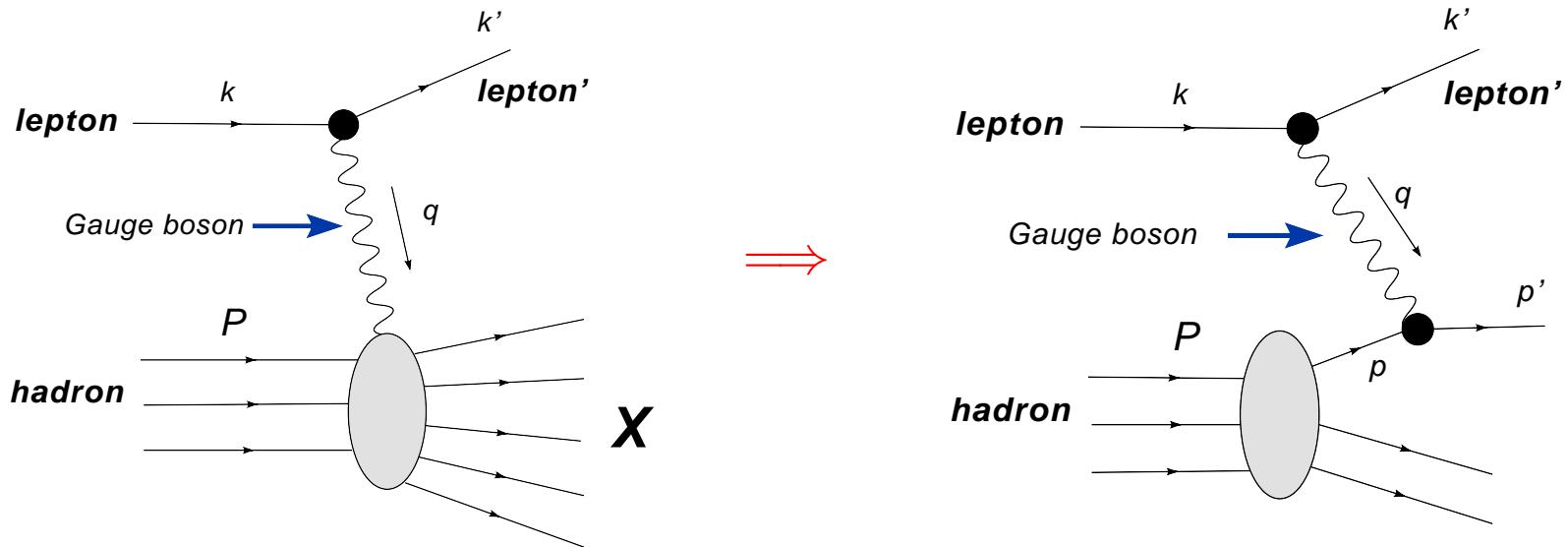
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– ACAT 2008, Erice, Sicily, November 3-7, 2007

Introduction

- Deep-inelastic lepton-hadron scattering ($e^\pm p$, $e^\pm n$, νp , $\bar{\nu} p$, ... - collisions)



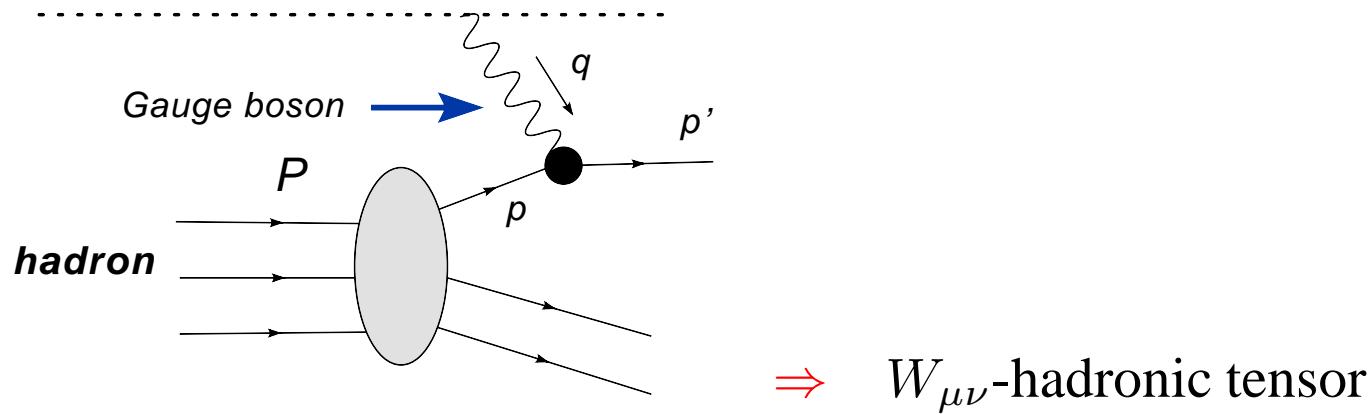
- Gauge boson:
 γ , Z^0 - NC
 W^\pm - CC

Kinematic variables

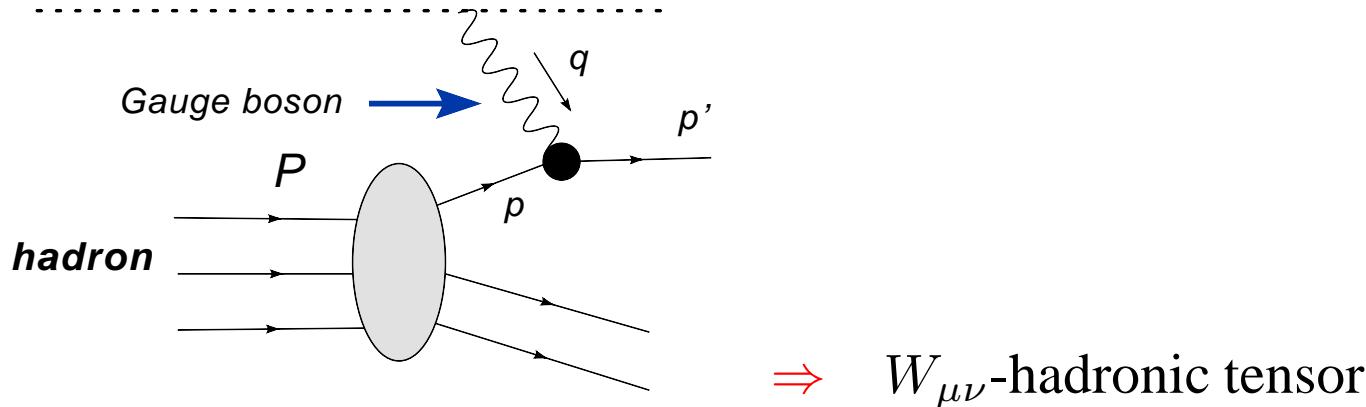
- momentum transfer $Q^2 = -q^2 > 0$
- Bjorken variable $x = Q^2/(2P \cdot q)$
- Inelasticity $y = (P \cdot q)/(P \cdot k)$

■ Particles: **POLARIZED** or **UNPOLARIZED**

Cross section: $d\sigma \sim L^{\mu\nu} W_{\mu\nu}$



Cross section: $d\sigma \sim L^{\mu\nu} W_{\mu\nu}$

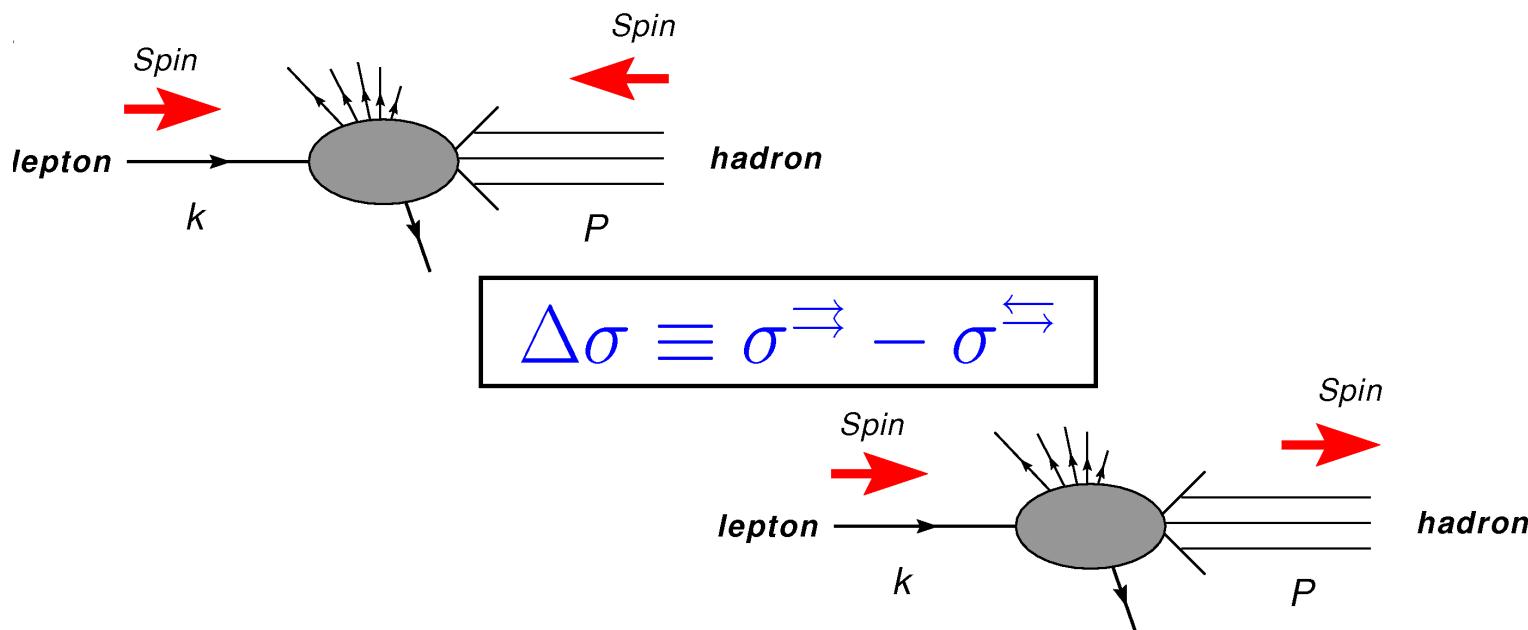


For unpolarized DIS:

$$W_{\mu\nu} = e_{\mu\nu} \frac{1}{2x} F_L(x, Q^2) + d_{\mu\nu} \frac{1}{2x} F_2(x, Q^2) + i\epsilon_{\mu\nu\alpha\beta} \frac{P^\alpha q^\beta}{P \cdot q} F_3(x, Q^2)$$

- $e_{\mu\nu}, d_{\mu\nu}$ - tensors, depend on P, q . symmetric under $\mu \leftrightarrow \nu$
- main objects are $F_{2,L,3}$ structure functions

- Polarized DIS : lepton and hadron are polarized \mapsto



- Example: lepton-proton scattering, exchange via virtual photon γ
..., Zijlstra, van Neerven '93

$$\frac{d^2\Delta\sigma}{dx dy} = \frac{8\pi\alpha^2}{q^2} \left[\{2 - y - Mxy\}g_1(x, Q^2) - \frac{2Mx}{E}g_2(x, Q^2) \right]$$

M - mass of the proton, E - energy of the lepton

$M = 0 \Rightarrow$ ONLY $g_1(x, Q^2)$ polarized structure function

Structure functions in QCD improved parton model

- Wilson coefficient functions $C_{i,\text{parton}}$

$$F_i(x, Q^2) = \sum_{\text{parton}} \int_x^1 \frac{d\xi}{\xi} PDF_{\text{parton}}(\xi, \mu^2) C_{i,\text{parton}}\left(\frac{x}{\xi}, \alpha_s(\mu^2), \frac{\mu^2}{Q^2}\right)$$

- Parton distribution functions PDF_{parton}
 - NOT calculable in perturbative QCD, extracted from experiment data
 - evolution of $PDF's$ via $DGLAP$ evolution equation with help of splitting functions $P_{\text{parton},\text{parton}'}$

$$\frac{Q^2}{\partial Q^2} \begin{bmatrix} q \\ g \end{bmatrix} = \begin{bmatrix} P_{qq} & P_{qg} \\ P_{gq} & P_{gg} \end{bmatrix} \otimes \begin{bmatrix} q \\ g \end{bmatrix}$$

- The same features for polarized structure function $g_1(x, Q^2)$

Chain of calculations

- $P_{\text{parton},\text{parton}'}^{}, C_{i,\text{parton}}^{}:$ are in hadronic tensor $W_{\mu\nu}^{}$
and also !!! in its analogue – partonic tensor $w_{\mu\nu}^{}$

Chain of calculations

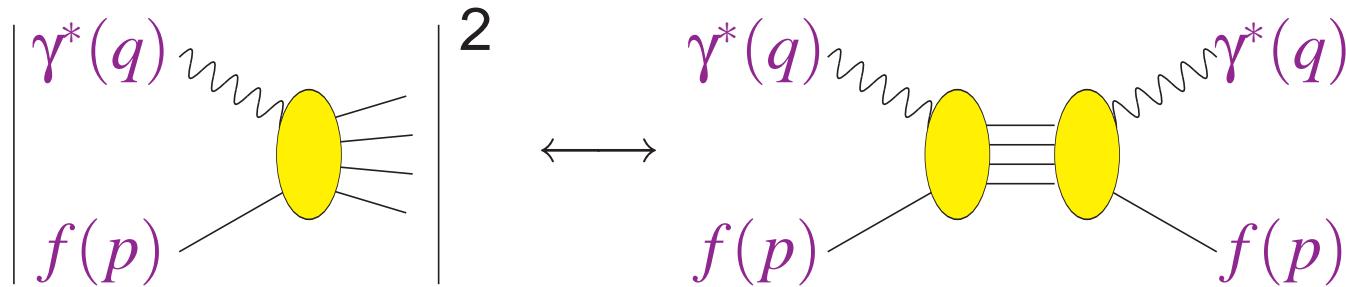
- $P_{\text{parton},\text{parton}'}, C_{i,\text{parton}}$: are in hadronic tensor $W_{\mu\nu}$ and also !!! in its analogue – partonic tensor $w_{\mu\nu}$



- For $w_{\mu\nu}$ use of *OPTICAL THEOREM*: $w_{\mu\nu}(p, q) \propto \text{Im } t_{\mu\nu}(p, q)$

Optical theorem

- The partonic tensor $w_{\mu\nu}(p, q)$ is related to the imaginary part of the partonic forward Compton scattering amplitude $t_{\mu\nu}(p, q)$
 - α_s^3 calculation in DIS with help of **loop technology**



$$w_{\mu\nu}(p, q) = \frac{1}{2\pi} \text{Im} t_{\mu\nu}(p, q)$$

▲ Dispersion relation in x : coefficient of $(2p \cdot q)^N \leftrightarrow N\text{-th moment}$

$$A^N \equiv \int_0^1 dx x^{N-1} A(x)$$

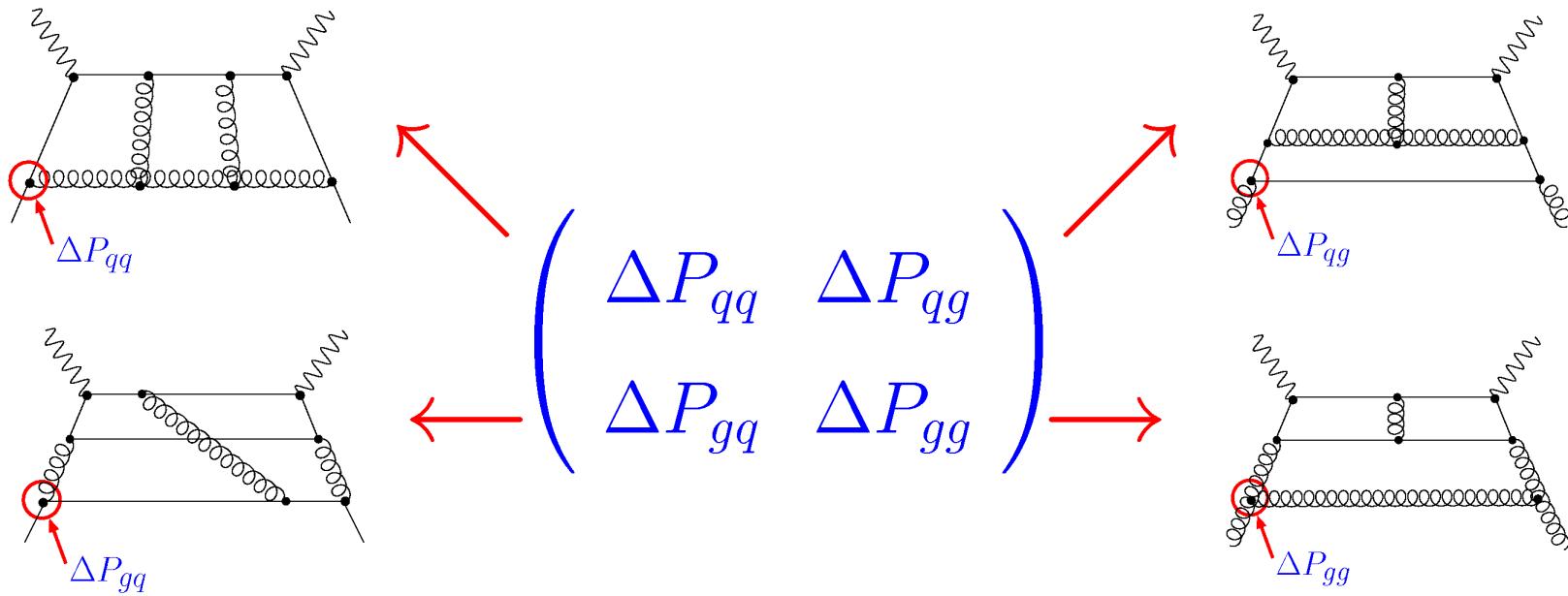
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Nogueira '93

- Example: diagrams for the polarized splitting functions



Chain of calculations

- $P_{\text{parton},\text{parton}'}^{}, C_{i,\text{parton}}^{}:$ are in hadronic tensor $W_{\mu\nu}^{} \downarrow$ and also !!! in its analogue – partonic tensor $w_{\mu\nu}^{} \downarrow$
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Chain of calculations

- $P_{\text{parton},\text{parton}'}, C_{i,\text{parton}}$: are in hadronic tensor $W_{\mu\nu}$ and also !!! in its analogue – partonic tensor $w_{\mu\nu}$
↓
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↓
- $t_{\mu\nu}(p, q)$ - generation of all possible diagrams up to 3 loops → QGRAF
Nogueira '93
↓
- Depending on NC/CC, polarized/unpolarized use of symmetries to reduce number of diagrams: *procedure of* Moch, Vermaseren Vogt '05
- ▲ Symbolic manipulations → FORM and TFORM Tentyukov, Vermaseren '07

Use of symmetries

- Number of diagrams for the polarized DIS after treatment by *procedure of Moch, Vermaseren Vogt '05*

legs	tree	1-loop	2-loop	3-loop
$q\gamma q\gamma$	1	3	25	359
$g\gamma g\gamma$	0	4	46	900
$q\varphi q\psi$	0	4	131	3890
$g\varphi g\psi$	2	31	924	29383 !!!

Original QGRAF output has up to **4 times more!**

Chain of calculations

- $P_{\text{parton},\text{parton}'}, C_{i,\text{parton}}$: are in hadronic tensor $W_{\mu\nu}$ and also !!! in its analogue – partonic tensor $w_{\mu\nu}$
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Chain of calculations

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↓
- Calculation of diagrams → done in Mellin- N space
 - Method of projection Gorishnii, Larin, Tkachev '83: Gorishnii, Larin '87
 - fixed Mellin-moments → MINCER in FORM
Larin, Tkachev, Vermaseren '91
 - symbolic N calc. and inverse Mellin transf. to Bjorken- x space → *recipe of* Moch, Vermaseren, Vogt '04

Method of projection & MINCER

Method of projection in pictures

- Identify scalar topologies
- Scalar diagram with external momenta P and Q
- N -th moment:
→ coefficient of $(2P \cdot Q)^N$
- Taylor expansion
- Feed scalar two-point functions in MINCER

$$\text{Scalar loop diagram} = \int \prod_n^3 d^D l_n \frac{1}{(P - l_1)^2} \frac{1}{l_1^2 \dots l_8^2}$$

$$\text{Scalar loop diagram} = \frac{(2P \cdot Q)^N}{(Q^2)^{N+\alpha}} C_N$$

$$\frac{1}{(P - l_1)^2} = \sum_i \frac{(2P \cdot l_1)^i}{(l_1^2)^{i+1}} \rightarrow \frac{(2P \cdot l_1)^N}{(l_1^2)^N}$$

⇒

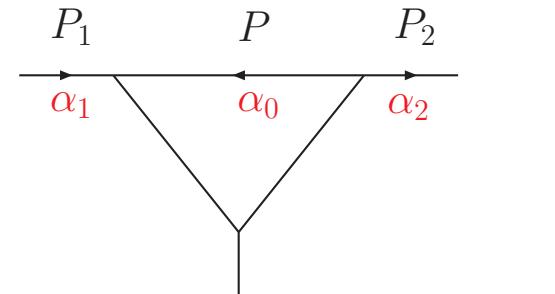


Mincer

- $\int dP \frac{\partial}{\partial P^\mu} [(P - l_j)^\mu \times I(l_1, \dots, P, \dots)] = 0$ - integration by part identities
t'Hooft, Veltman '72; Chetyrkin , Tkachov '81
Leibniz, Newton :-)

Mincer

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Triangle rule

Define

$$I(\alpha_0, \beta_1, \beta_2, \alpha_1, \alpha_2) = \int d^D P \frac{1}{(P^2)^{\alpha_0} ((P + P_1)^2)^{\beta_1} (P_1^2)^{\alpha_1} ((P + P_2)^2)^{\beta_2} (P_2^2)^{\alpha_2}}$$

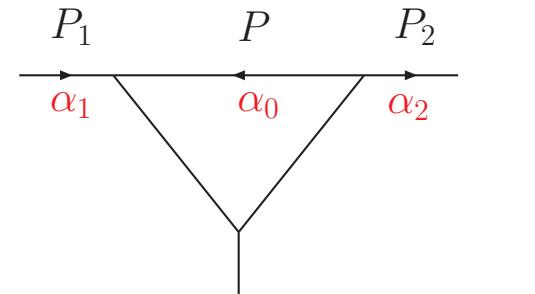
and act the integrand with $\frac{\partial}{\partial P_\mu} P_\mu = D + P_\mu \frac{\partial}{\partial P_\mu}$. Result \Rightarrow

Recursion relation:

$$\begin{aligned} & I(\alpha_0, \beta_1, \beta_2, \alpha_1, \alpha_2) \times (D - 2\alpha_0 - \beta_1 - \beta_2) = \\ & \beta_1 (I(\alpha_0 - 1, \beta_1 + 1, \beta_2, \alpha_1, \alpha_2) - I(\alpha_0, \beta_1 + 1, \beta_2, \alpha_1 - 1, \alpha_2)) \\ & \beta_2 (I(\alpha_0 - 1, \beta_1, \beta_2 + 1, \alpha_1, \alpha_2) - I(\alpha_0, \beta_1, \beta_2 + 1, \alpha_1, \alpha_2 - 1)) \end{aligned}$$

Mincer

- $\int dP \frac{\partial}{\partial P^\mu} [(P - l_j)^\mu \times I(l_1, \dots, P, \dots)] = 0$ - integration by part identities
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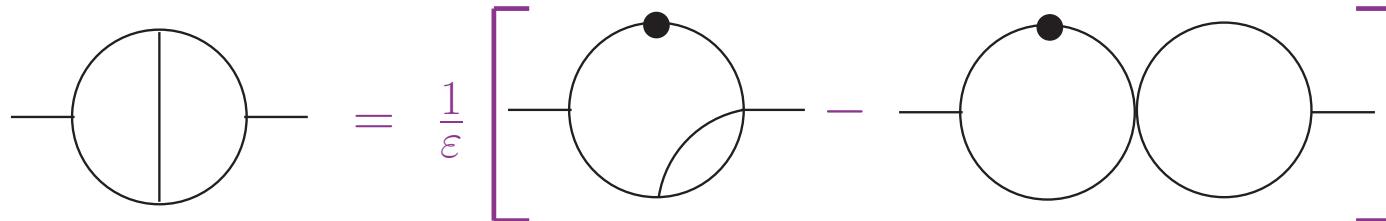
$$I(\alpha_0, \beta_1, \beta_2, \alpha_1, \alpha_2) = \int d^D P \frac{1}{(P^2)^{\alpha_0} ((P + P_1)^2)^{\beta_1} (P_1^2)^{\alpha_1} ((P + P_2)^2)^{\beta_2} (P_2^2)^{\alpha_2}}$$

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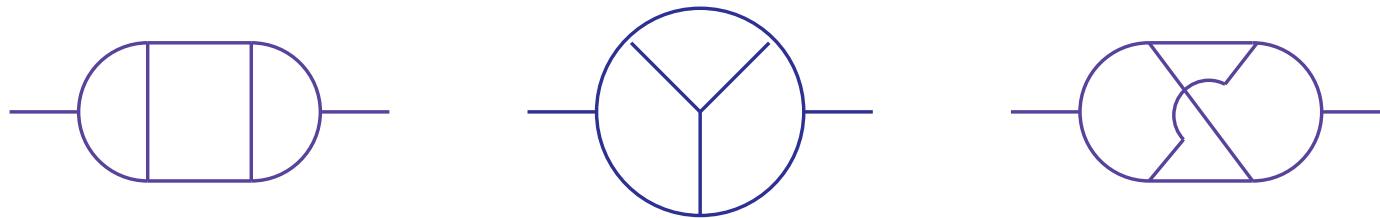
$$\begin{aligned} I(\alpha_0, \beta_1, \beta_2, \alpha_1, \alpha_2) \times (D - 2\alpha_0 - \beta_1 - \beta_2) = \\ \beta_1 (I(\alpha_0 - 1, \beta_1 + 1, \beta_2, \alpha_1, \alpha_2) - I(\alpha_0, \beta_1 + 1, \beta_2, \alpha_1 - 1, \alpha_2)) \\ \beta_2 (I(\alpha_0 - 1, \beta_1, \beta_2 + 1, \alpha_1, \alpha_2) - I(\alpha_0, \beta_1, \beta_2 + 1, \alpha_1, \alpha_2 - 1)) \end{aligned}$$

In pictures



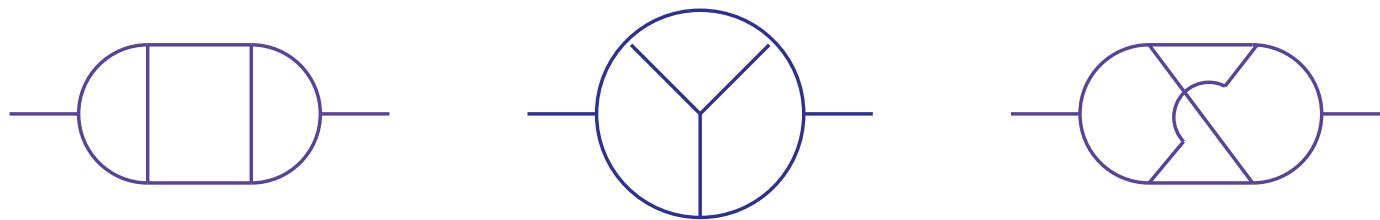
Classification of loop integrals

- Classify according to topology of underlying two-point function
 - top-level topology types ladder, benz, non-planar \Rightarrow

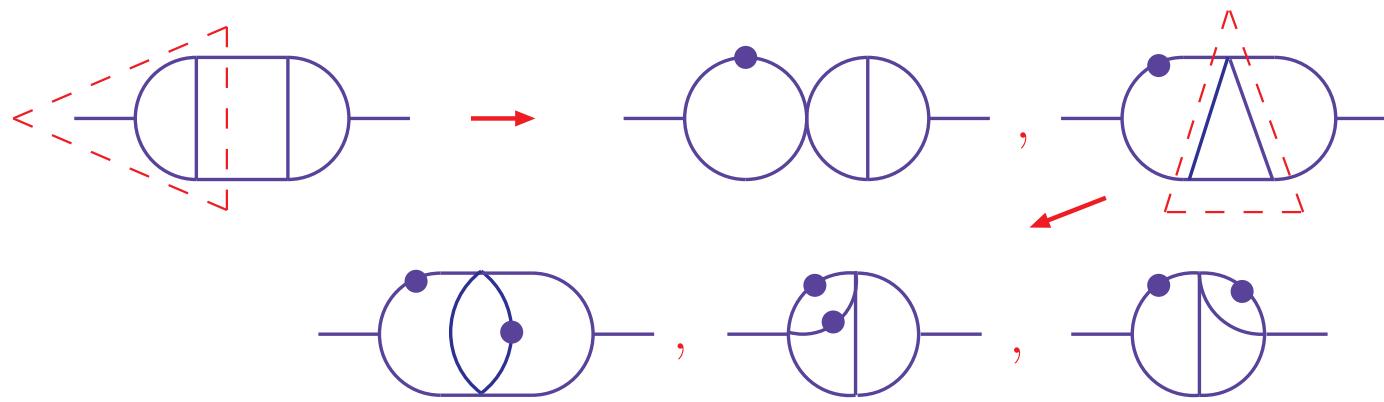


Classification of loop integrals

- Classify according to topology of underlying two-point function
 - top-level topology types ladder, benz, non-planar \Rightarrow



- Using **IBP** identities more complicated topologies are reduced to simpler topologies



Symbolic N calculations

- Combine identities: integration by parts, scaling, Passarino-Veltman type

⇒ Difference equations for $I(N)$ [recall: coefficient of $(2p \cdot q)^N$]

$$a_0(N)I(N) - \dots - a_n(N)I(N-n) = I_0(N)$$

Simple scalar example [red line: flow of massless parton momentum p]

$$\begin{array}{c} \text{Diagram 1: } \text{Two green ovals connected by a horizontal bar. The top bar has two red arcs above it. Vertices are labeled '1' at all four corners.} \\ + \frac{N+3+3\epsilon}{N+2} \frac{2p \cdot q}{q^2} \begin{array}{c} \text{Diagram 2: } \text{Two green ovals connected by a horizontal bar. The top bar has three red arcs above it. Vertices are labeled '1' at all four corners.} \\ = \frac{2}{N+2} \begin{array}{c} \text{Diagram 3: } \text{Two green ovals connected by a horizontal bar. The top bar has one red arc above it. The rightmost oval has a '2' at its top vertex. Vertices are labeled '1' at all four corners.} \end{array} \end{array} \end{array}$$

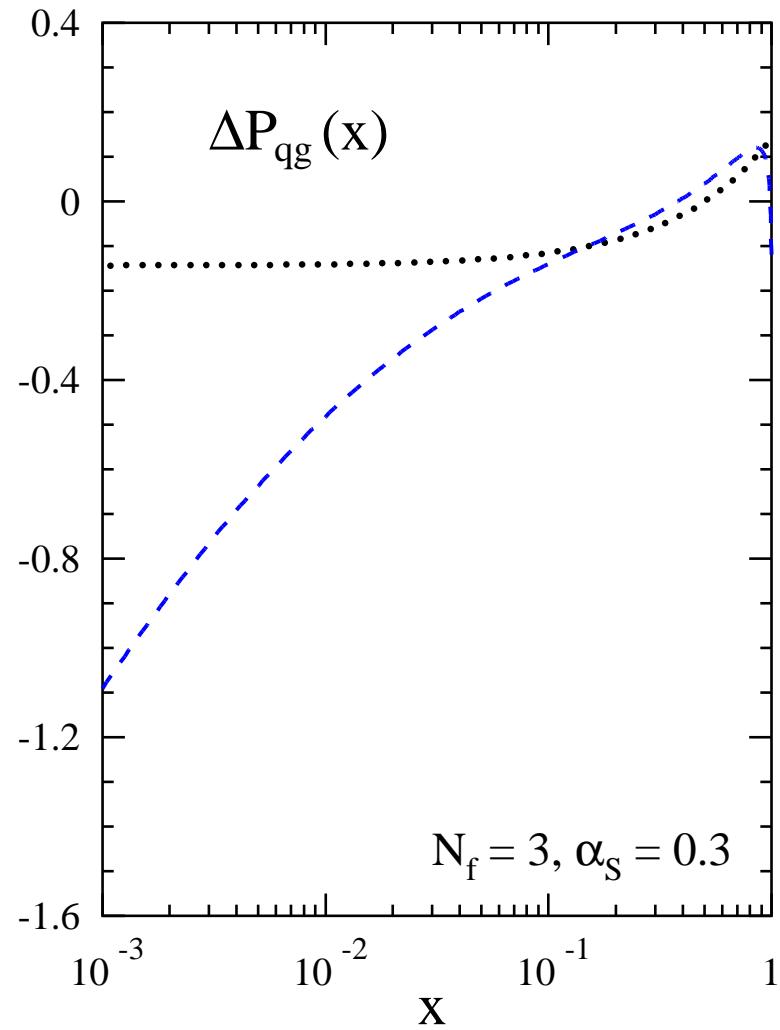
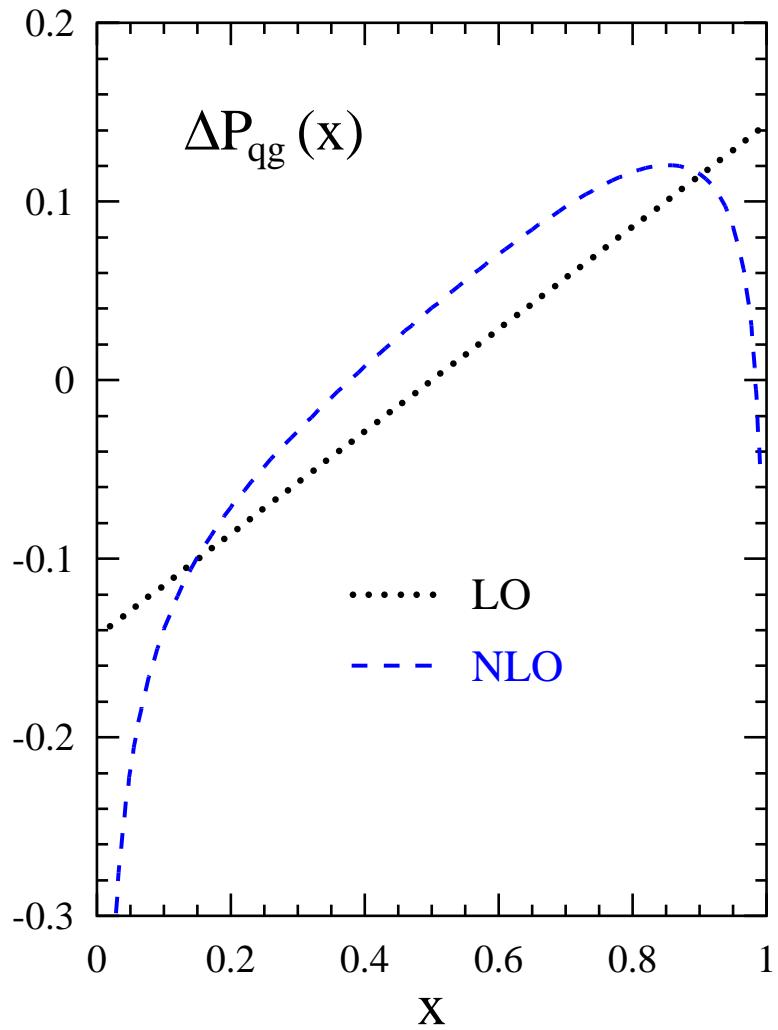
- Successive reduction to simpler (lower topologies or ‘less red’) integrals

Examples of results

$$\begin{aligned}
C_{3,10}^{\text{ns}} = & 1 + a_s C_F \frac{1953379}{138600} + a_s^2 C_F n_f \left(-\frac{537659500957277}{15975002736000} \right) + a_s^2 C_F^2 \left(\frac{597399446375524589}{14760902528064000} \right. \\
& \left. + \frac{7202}{105} \zeta_3 \right) + a_s^2 C_A C_F \left(\frac{5832602058122267}{29045459520000} - \frac{99886}{1155} \zeta_3 \right) \\
& + a_s^3 C_F n_f^2 \left(\frac{51339756673194617191}{996360920644320000} + \frac{48220}{18711} \zeta_3 \right) \\
& + a_s^3 C_F^2 n_f \left(-\frac{125483817946055121351353}{209235793335307200000} - \frac{59829376}{3274425} \zeta_3 + \frac{24110}{693} \zeta_4 \right) \\
& + a_s^3 C_F^3 \left(-\frac{744474223606695878525401307}{7088908678200207936000000} + \frac{28630985464358}{24960941775} \zeta_3 \right. \\
& \left. + \frac{151796299}{8004150} \zeta_4 - \frac{53708}{99} \zeta_5 \right) \\
& + a_s^3 C_A C_F n_f \left(-\frac{185221350045507487753}{226445663782800000} + \frac{8071097}{39690} \zeta_3 - \frac{24110}{693} \zeta_4 \right) \\
& + a_s^3 C_A C_F^2 \left(\frac{19770078729338607732075449}{8369431733412288000000} - \frac{619383700181}{5546875950} \zeta_3 \right. \\
& \left. - \frac{151796299}{5336100} \zeta_4 - \frac{37322}{99} \zeta_5 \right) \\
& + a_s^3 C_A^2 C_F \left(\frac{93798719639056648125143}{36231306205248000000} - \frac{43202630363}{20582100} \zeta_3 \right. \\
& \left. + \frac{151796299}{16008300} \zeta_4 + \frac{195422}{231} \zeta_5 \right).
\end{aligned}$$

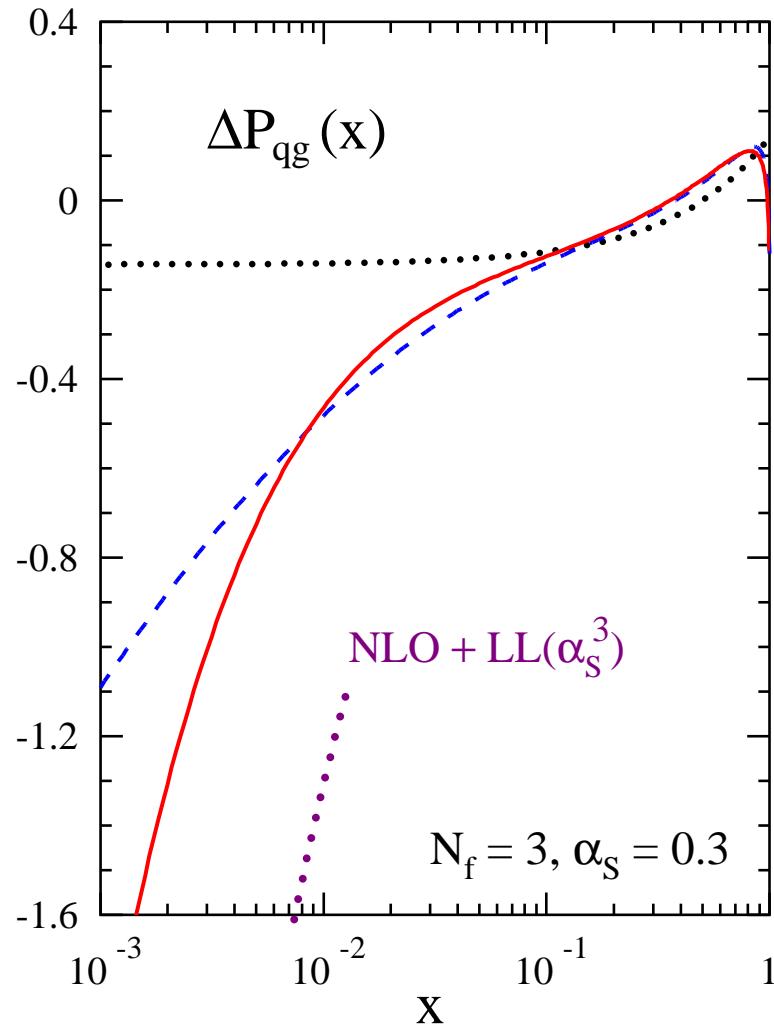
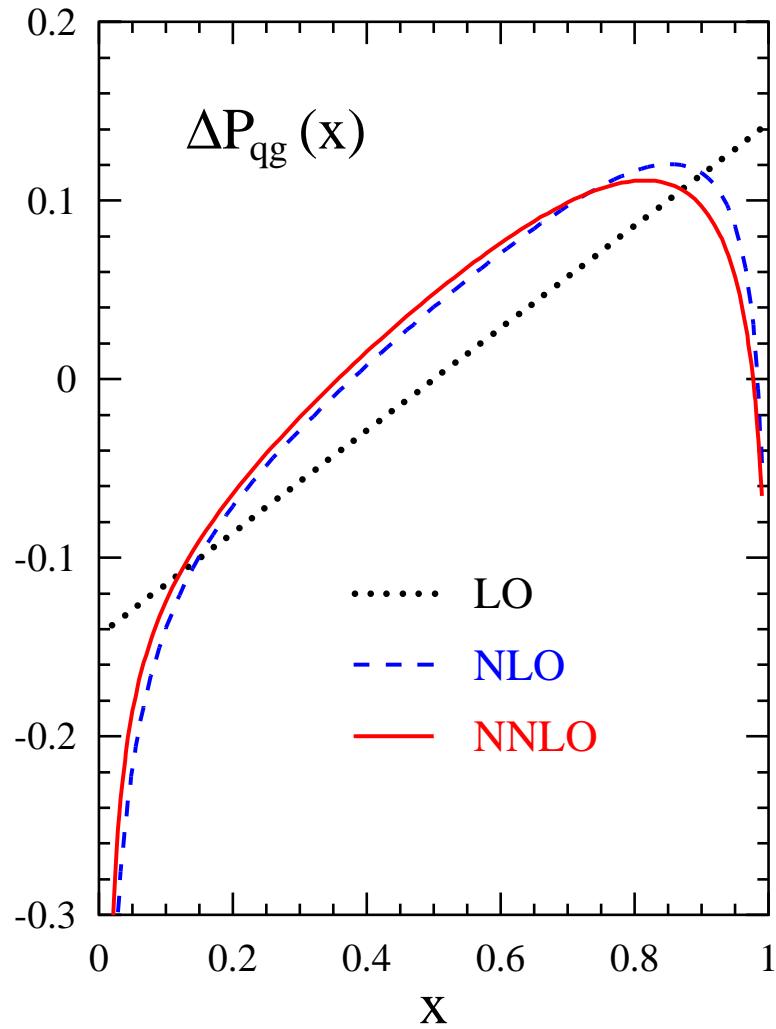
Polarized DIS: The spin splitting function ΔP_{qg} to NNLO

Vogt, Moch, Rogal, Vermaseren '08



Polarized DIS: The spin splitting function ΔP_{qg} to NNLO

Vogt, Moch, Rogal, Vermaseren '08



NNLO corr. $\leq 15\%$ for $0.005 \leq x < 0.9$

Summary

- The *recipe* for higher order DIS calculations has been described
- Application examples:
 - **Charged Current DIS**
S. Moch, M. Rogal, Nucl.Phys.B782 '07
S. Moch, M. Rogal, A. Vogt, Nucl.Phys.B790:317-335, '07
M. Rogal, arXiv:0711.0521 [hep-ph]
 - 3-loop splitting functions for **polarized DIS**
A. Vogt, S. Moch, M. Rogal, J.A.M. Vermaseren, arXiv:0807.1238 [hep-ph] '08