



# Probing $\gamma\gamma XX$ anomalous gauge couplings with proton tagging at the LHC

New Trends in HEP and QCD workshop @ Natal, Brazil

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C. Royon, O. Kepka, *Phys. Rev. D* **78** (2008)

November, 5th 2014

E. Chapon, C. Royon, O. Kepka, *Phys. Rev. D* **81** (2010)

S. Fichet, G. von Gersdorff, O. Kepka, B. Lenzi, C. Royon, M. Saimpert, *Phys. Rev. D* **89** (2014)

S. Fichet, G. von Gersdorff, B. Lenzi, C. Royon, M. Saimpert, to be submitted



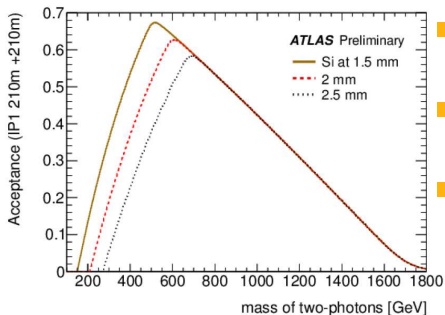
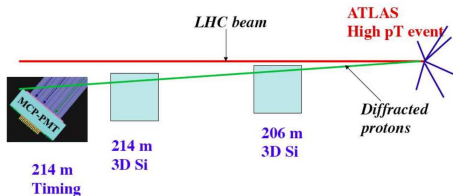
## 1 Motivations and description of the proposed measurements

*What and why are we measuring?*

- 2 The  $WW\gamma\gamma$  and  $ZZ\gamma\gamma$  couplings cases in short
- 3 The  $\gamma\gamma\gamma\gamma$  couplings case in details
- 4 Conclusion and plans

# Forward proton detectors at the LHC

- The ATLAS Forward Physics (AFP) and the CMS-TOTEM Precision Proton Spectrometer (CT-PPS) upgrade projects



- $\xi_{1,2}$  proton momentum loss

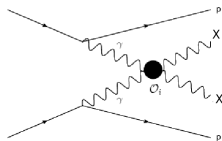
- **missing proton mass**

$$m_{pp}^{miss} = \sqrt{\xi_1 \xi_2 s}$$

- **mass acceptance**

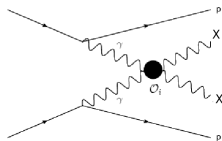
limited by the LHC beam and optics (low  $\beta^*$  case displayed)

# Exclusive production via photon induced processes



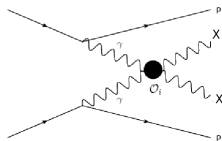
- **All particles at the final state are detected:** two protons in the forward detectors and two high energy particles in the central detector → **strong kinematics constraints**
- Full reconstructed kinematics by comparing the forward to the central particles → **strong background reduction**

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- aQGC important for various physics topics: **electroweak symmetry breaking, extra-dimension models, ...**

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- aQGC important for various physics topics: **electroweak symmetry breaking, extra-dimension models, ...**
- **Drawback:** smaller cross-sections  
(intact protons must be in the acceptance of the forward detectors)

# Summary of the presentation

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- 1 Motivations and description of the proposed measurements
- 2 **The  $WW\gamma\gamma$  and  $ZZ\gamma\gamma$  couplings cases in short**  
*Brief overview of the two studies published in 2008 and 2010.*  
*Comparison on the obtained sensitivities with the latest measurements.*
- 3 The  $\gamma\gamma\gamma\gamma$  couplings case in details
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# WW $\gamma\gamma$ and ZZ $\gamma\gamma$ anomalous couplings

C. Royon, O. Kepka, *Phys. Rev. D* **78** (2008)

E. Chapon, C. Royon, O. Kepka, *Phys. Rev. D* **81** (2010)

- **Effective Field Theory (EFT): dimension 6 operators** parametrized with 4 different parameters

$$\mathcal{L}_6^0 \sim \frac{-e^2 a_0^W}{8 \Lambda^2} F_{\mu\nu} F^{\mu\nu} W^{+\alpha} W_{\alpha}^{-} - \frac{e^2}{16 \cos^2(\theta_W)} \frac{a_0^Z}{\Lambda^2} F_{\mu\nu} F^{\mu\nu} Z^{\alpha} Z_{\alpha}$$

$$\mathcal{L}_6^C \sim \frac{-e^2 a_C^W}{16 \Lambda^2} F_{\mu\alpha} F^{\mu\beta} (W^{+\alpha} W_{\beta}^{-} + W^{-\alpha} W_{\beta}^{+}) \\ - \frac{e^2}{16 \cos^2(\theta_W)} \frac{a_C^Z}{\Lambda^2} F_{\mu\alpha} F^{\mu\beta} Z^{\alpha} Z_{\beta}$$

- **Only the leptonic decays** of the heavy bosons are considered as final states (clean signal experimentally)
- **Background considered:** ND WW/ZZ production, dilepton photoproduction, DPE dilepton, DPE WW/ZZ
- Generation and simulation performed with the **Forward Physics MC generator (FPMC)** interfaced with the fast simulation of the ATLAS detector (ATLFast++ package)

*ATLAS full simulation also performed to probe pile-up effects and gave similar results*



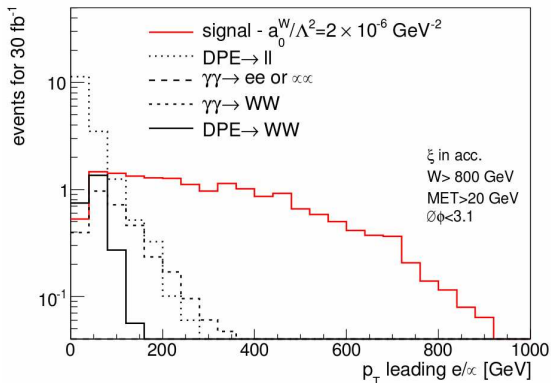
# $WW_{\gamma\gamma}$ and $ZZ_{\gamma\gamma}$ anomalous signal

→ *ATLAS fast simulation study*

E. Chapon, C. Royon, O. Kepka, *Phys. Rev. D* **81** (2010)

## ■ Anomalous signal appears at high energy

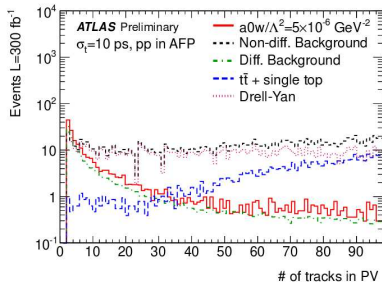
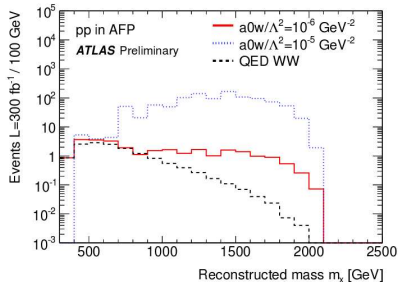
*$WW_{\gamma\gamma}$  signal and background composition passing all the selection but the  $p_T$  cut on the leading lepton*



# Dealing with pile-up at the LHC

→ *ATLAS full simulation study*

- The LHC is operated at **very high luminosity** → **high event multiplicities** in a single bunch-crossing (pile-up)
- Use of the **forward timing detectors** to constrain the vertex z-position of the interaction  
*dependance on the timing detectors resolution*
- Cut on **the number of tracks fitted to the primary vertex**  
*very efficient to remove remaining pile up after requesting a high mass object to be produced*



# WW $\gamma\gamma$ and ZZ $\gamma\gamma$ sensitivities

E. Chapon, C. Royon, O. Kepka, *Phys. Rev. D* **81** (2010)

- **Limits from LEP (OPAL)**,  $\simeq 0.1\text{-}0.2 \text{ GeV}^{-2}$  (for each coupling)  
*Phys. Rev. D* **70** (2004) 032005
- **Recent papers from DØ** for WW $\gamma\gamma$  with reach of the order of  $10^{-2}\text{-}10^{-3} \text{ GeV}^{-2}$   
*Phys. Rev. D* **88** (2013) 012005
- **Recent papers from CMS** for WW $\gamma\gamma$  with reach of the order of  $10^{-4}\text{-}10^{-5} \text{ GeV}^{-2}$   
*Phys. Rev. D* **90** (2014) 032008
- **Sensitivities predictions with AFP** ( $30$  and  $200 \text{ fb}^{-1}$ )  
reach up to  $10^{-6} \text{ GeV}^{-2}$ , *improvement up to a factor  $\simeq 100$*

form factor		limits [ $10^{-6} \text{ GeV}^{-2}$ ]				limits [ $10^{-6} \text{ GeV}^{-2}$ ]			
		$ a_0^W/\Lambda^2 $	$ a_C^W/\Lambda^2 $	$ a_0^Z/\Lambda^2 $	$ a_C^Z/\Lambda^2 $	$ a_0^W/\Lambda^2 $	$ a_C^W/\Lambda^2 $	$ a_0^Z/\Lambda^2 $	$ a_C^Z/\Lambda^2 $
95% c.l.	$\Lambda_{cut} = \infty$	1.2	4.2	2.8	10	0.7	2.4	1.1	4.1
	$\Lambda_{cut} = 2 \text{ TeV}$	2.6	9.4	6.4	24	1.4	5.2	2.5	9.2
3 $\sigma$ evidence	$\Lambda_{cut} = \infty$	1.6	5.8	4.0	14	0.85	3.0	1.6	5.7
	$\Lambda_{cut} = 2 \text{ TeV}$	3.6	13	9.0	34	1.8	6.7	3.5	13
5 $\sigma$ discovery	$\Lambda_{cut} = \infty$	2.3	9.7	6.2	23	1.2	4.3	4.1	8.9
	$\Lambda_{cut} = 2 \text{ TeV}$	5.4	20	14	52	2.7	9.6	5.5	20

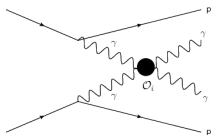
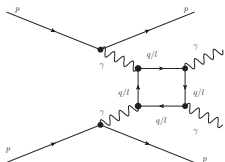
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- 3 **The  $\gamma\gamma\gamma\gamma$  couplings case in details**  
*Review of the study published in 2014 and latest developments not published yet. Ideas for further developments.*
- 4 Conclusion and plans

# $\gamma\gamma\gamma\gamma$ SM and anomalous couplings

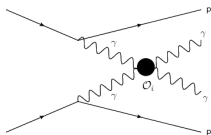
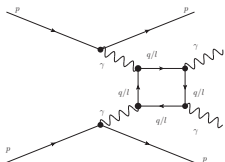


- **Direct coupling absent from the SM**

*Loop induced production strongly suppressed in the SM, measurable at the LHC in Pb-Pb (d'Enterria et al. Phys. Rev. Lett. **111** (2013) 080405)*

- **Never measured in collider experiments**

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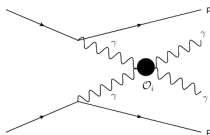
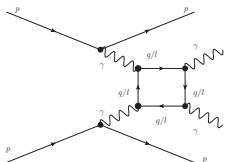
- New Physics search → high integrated luminosity required (**so high pile-up!**)

*300 fb<sup>-1</sup> of data expected at the LHC at  $\sqrt{s} = 14$  TeV with  $\mu > 50$*

- **Huge background** if only 2 high energy  $\gamma$  required

*(SM  $\gamma\gamma$  production + fakes from electrons and jets)*

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- **Additional requirement of two intact protons** with forward detectors highly suppresses the background

*(All particles at the final state are detected)*

# Operators of the $\gamma\gamma\gamma\gamma$ couplings

R.S. Gupta, *Phys. Rev. D* **85** (2012) 014006

S. Fichtel and G. von Gersdorff, *arXiv:1311.6815*

- $\sqrt{\hat{s}_{\gamma\gamma}} \ll \Lambda$ , effective field theory assumption

$$L_{4\gamma} = \zeta_1^\gamma F_{\mu\nu} F^{\mu\nu} F_{\rho\sigma} F^{\rho\sigma} + \zeta_2^\gamma F_{\mu\nu} F^{\nu\rho} F_{\rho\sigma} F^{\sigma\mu} \quad (\text{dimension } 8)$$



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- For low new physics masses, threshold effect to be taken into account → **use of a form factor (f.f.) at the amplitude level**

$$\text{We use } f.f. = \frac{1}{1 + (\frac{\hat{s}_{\gamma\gamma}}{\Lambda'^2})^2} \text{ with } \Lambda' = 1 \text{ TeV} \simeq \sqrt{\hat{s}_{\gamma\gamma, \text{max}}}/2$$

(Unitarity requires  $\zeta_i < 10^{-10} \text{ GeV}^{-4}$ ,  $\simeq 10^4$  higher than our sensitivity limit, so we are safe on this side)

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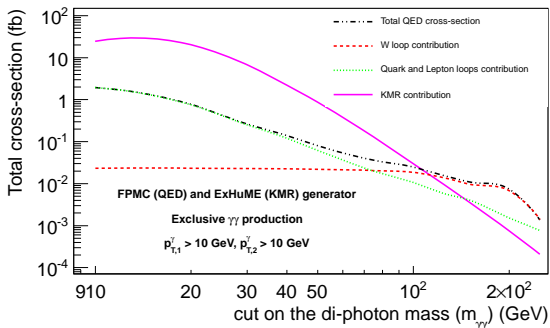
- The signal showed in the plots of this presentation are for a signal with  $\zeta_1 \geq 0$  and  $\zeta_2 = 0$  and with f.f.  
 *$\zeta_1$  and  $\zeta_2$  have a very similar angular behaviour*
- A table of final sensitivities for both  $\zeta_1$  and  $\zeta_2$ , **with and without f.f.** are given at the end of the presentation

- Evaluate the LHC potential to probe  $4\gamma$  couplings using proton tagging and **effective field theory**
  - $4\gamma$  aQGC operators implemented in the **FPMC generator**
  - Simulation of the detector effects
  - Pile-up simulation with Pythia8 minimum bias events
  - Background estimation (expected to be very small thanks to the fully constrained kinematics)
  - Sensitivities calculation: significance  $(\sigma) = S/\sqrt{B}$
  - 2 scenarios were considered
    - LHC full stat (ATLAS or CMS) :  $300 \text{ fb}^{-1}$ ,  $\langle \mu \rangle = 50$
    - HL-LHC (ATLAS) :  $3000 \text{ fb}^{-1}$ ,  $\langle \mu \rangle = 200$

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- **Extra:** update of the exclusive  $\gamma\gamma$  SM production in FPMC ***W loop contribution and the fermion masses included***

# SM exclusive $\gamma\gamma$ production (preliminary)



- Mass of the fermions, W loop contribution and related interference taken into account in the simulation
- W loop non negligible for  $m_{\gamma\gamma} > 70$  GeV
- **Irreducible background for  $\gamma\gamma\gamma\gamma$  new physics searches**  
*Needs to be simulated accurately*

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# Modelisation of the detector effects

- **Detector effects (acceptance/efficiency) must be taken into account to get realistic predictions**
- **Analysis performed at particle level** but taking into account main detector effects
  - Estimation of  $\gamma$  **conversion rates** ( $\eta$  function), **fake photon rates**, **reconstruction efficiency** ( $p_T$  functions) from ECFA ATLAS studies
  - **Smearing** of 1% in  $\gamma$  energies, 0.001 in  $\eta$  and  $\phi$  (absolute), 2% for  $\xi$  to mimic detector resolution
  - Requirement of **at least one converted photon**  $\rightarrow$  **constraint on the  $\gamma$  vertex**, possibility to combine with forward proton timing measurement
  - Selection on high  $p_T^\gamma$ , high diphoton mass,  $\Delta\Phi^{\gamma\gamma}$ , match proton missing/ $\gamma\gamma$  mass (summary S17)

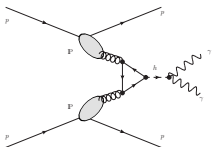
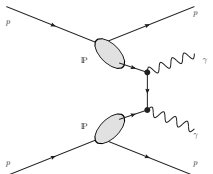
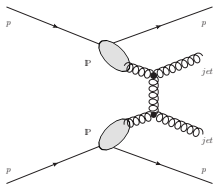


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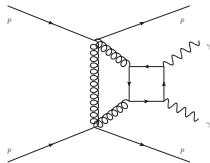
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- **Final outputs**
  - $5\sigma$  and 95% C.L sensitivities on the  $\gamma\gamma\gamma\gamma$  couplings at the LHC *EFT, valid for  $m > 2(1)$  TeV for tree-level (pair) production*
  - M-Q exclusion plane for generic new fermions/vectors at the LHC *full amplitude calculation, valid for all masses*

# Backgrounds (FPMC, ExHuME)

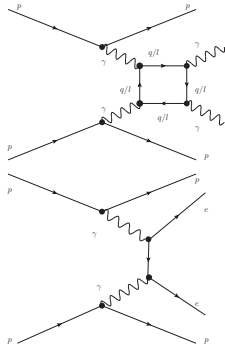
## IP backgrounds (FPMC)



## Exclusive QCD (ExHuME)

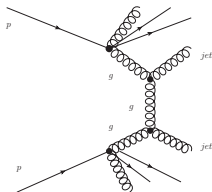


## Exclusive QED (FPMC)

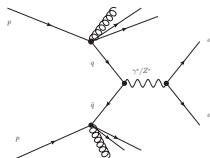


# Pile-up backgrounds (HERWIG 6.5)

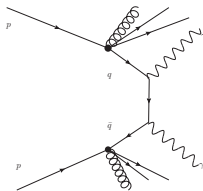
Dijet



Drell-Yan



Diphoton



**+ intact protons generated from minimum bias events (Pythia 8)**

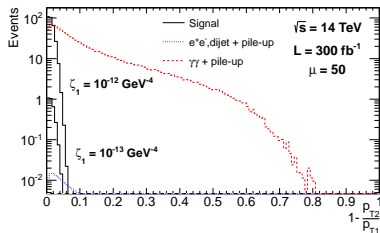
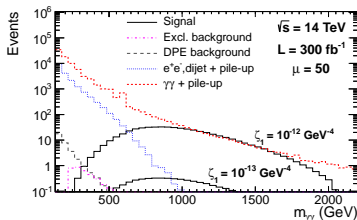
transported to the forward detectors through the LHC magnets with FPTracker/MADX (used by the LHC machine people)

# Mass and $p_T$ balance distribution of signal and backgrounds

smearing, fakes and reconstruction factors,  $\geq 1$  converted  $\gamma$  required

$0.015 < \xi < 0.15$ ,  $|\eta| < 2.37$ ,  
 $p_{T1,2}^\gamma > 50$  GeV **ONLY**

if we request also  $m_{\gamma\gamma} > 600$  GeV  
 and  $p_{T1,2} > 200, 100$  GeV

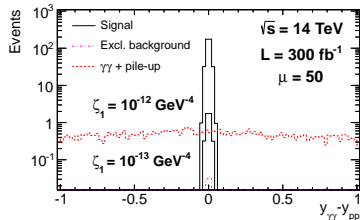
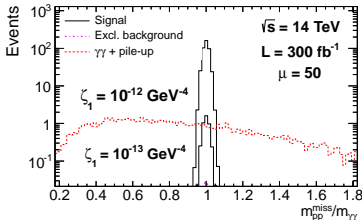


- By requesting  $m_{\gamma\gamma} > 600$  GeV (left), **Only pile-up backgrounds remain**
- $p_T$  ratio distribution after  $p_T$  and  $m_{\gamma\gamma}$  cuts (right) provides another efficient cut (**exclusive process**)
- $\Delta\phi > \Pi - 0.01$  also applied in the final selection

# Use of the forward detector $\xi$ measurement

smearing, fakes and reconstruction factors,  $\geq 1$  converted  $\gamma$  required

$0.015 < \xi < 0.15$ ,  $|\eta| < 2.37$ ,  $m_{\gamma\gamma} > 600$  GeV,  $p_{T1,2} > 200, 100$  GeV,  
 $p_T$  ratio  $> 0.95$ ,  $\Delta\phi > \Pi-0.01$



- **Missing proton mass**  $\sqrt{\xi_1 \xi_2 s}$  matches  $m_{\gamma\gamma}$  for the signal  
*A mass window of 3% (= resolution) is required in the event selection*
- Same effect with **rapidity variables**  
 $|y_{\gamma\gamma} - y_{pp}| < 0.03$  with  $y_{pp} = (0.5 * \ln(\frac{\xi_1}{\xi_2}))$  is applied
- The small width of the signal distributions is due to the smearing applied to simulate detector effects

Very efficient cuts due to very good  $\xi$  resolution, **absolutely needed in order to suppress the pile-up background**

# Event selection: summary

## ■ Kinematic cuts

1  $p_{T1}^\gamma > 200 \text{ GeV}, p_{T2}^\gamma > 100 \text{ GeV}$

2  $m_{\gamma\gamma} > 600 \text{ GeV}$

## ■ Selection of **exclusive events**

1  $\frac{p_{T2}}{p_{T1}} > 0.95$

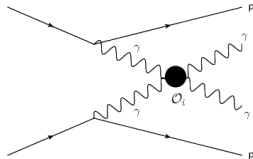
2  $|\Delta\Phi| > \pi - 0.01$

## ■ **Forward detectors** cuts

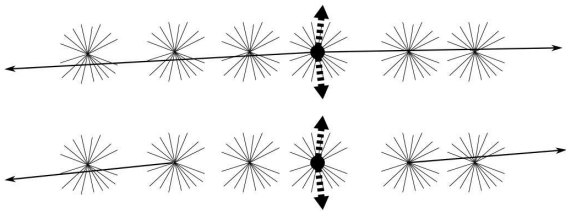
1  $m_{pp}^{miss} = m_{\gamma\gamma} \pm 3\%$

2  $|y_{\gamma\gamma} - y_{pp}| < 0.03$   
with  $y_{pp} = (0.5 * \ln(\frac{\xi_1}{\xi_2}))$

- 3 Possible proton timing measurement with forward detectors **(Not used)**



# Possible extra-cut: proton timing requirement



- Proton timing will be measured by forward detectors
  - 10 ps resolution assumed  $\rightarrow$  proton vertex constrained within 2.1 millimeters
  - Requirement of 1 converted  $\gamma \rightarrow < 1$  mm resolution on the  $\gamma$  vertex
  - **Resolution on the vertex position driven by forward timing detectors**
- Matching the two proton and photon vertices provide an additional background rejection factor of  $\simeq 40$  at  $\mu = 50$
- No need to use for this study, **robustness of the  $\gamma\gamma\gamma\gamma$  analysis**  
 $\rightarrow$  needed for  $\gamma\gamma WW$  and  $\gamma\gamma ZZ$ !
- can be used for unknown backgrounds (beam-induced)

# Expected events for $\zeta_1^\gamma = 2 \cdot 10^{-13} \cdot \text{GeV}^{-4}$

- $\sqrt{s} = 14 \text{ TeV}, L = 300 \text{ fb}^{-1}$ , at least one converted  $\gamma$

Cut / Process	Signal	Excl.	DPE	$e^+e^-$ , dijet + pu	$\gamma\gamma$ + pu
$0.015 < \xi < 0.15, p_{T1,2} > 200, 100\text{GeV}$	37.0	0.25	0.2	1.6	2968
$m_{\gamma\gamma} > 600 \text{ GeV}$	34.9	0.20	0	0.2	1023
$p_{T2}/p_{T1} > 0.95,  \Delta\phi  > \pi - 0.01$	34.0	0.19	0	0	80.2
$\sqrt{\xi_1 \xi_2 s} = m_{\gamma\gamma} \pm 3\%$	33.0	0.18	0	0	2.8
$ y_{\gamma\gamma} - y_{pp}  < 0.03$	31.7	0.18	0	0	0

- **Signal selection efficiency > 70%** (after detector effects)
  - Signal increased by a factor 3-4 when relaxing the >1 conv. photon requirement  
(the di-photon vertex cannot be identified accurately anymore from the central detector)



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- **Signal selection efficiency > 70%** (after detector effects)
  - Signal increased by a factor 3-4 when relaxing the  $>1$  conv. photon requirement  
(the di-photon vertex cannot be identified accurately anymore from the central detector)
- **Background completely suppressed thanks to forward detectors  $\xi$  measurement**
  - Very high significance per observed event
  - $< 5$  background events expected at  $\mu = 200$   
*Robust analysis, good background control*
  - proton time-of-flight **not used**  
*Possible additional rejection factor of 40 at  $\mu = 50$*

# Final discovery ( $5\sigma$ ) and exclusion (95% CL) sensitivities on $\zeta_1$ and $\zeta_2$

**EFT approach:** S. Fichet, G. von Gersdorff, O. Kepka, B. Lenzi, C. Royon, M. Saimpert, Phys. Rev. D **89** (2014)

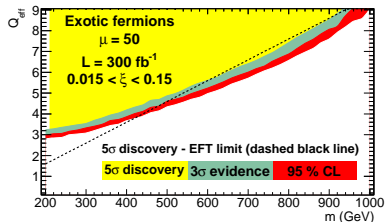
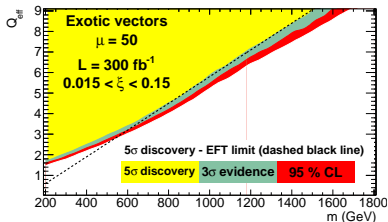
**Update of the EFT + full amplitude calculation:** S. Fichet, G. von Gersdorff, B. Lenzi, C. Royon, M. Saimpert, to be submitted

Luminosity	300 fb <sup>-1</sup>	300 fb <sup>-1</sup>	300 fb <sup>-1</sup>	3000 fb <sup>-1</sup>
pile-up ( $\mu$ )	50	50	50	200
coupling (GeV <sup>-4</sup> )	$\geq 1$ conv. $\gamma$ $5\sigma$	$\geq 1$ conv. $\gamma$ 95% CL	all $\gamma$ 95% CL	all $\gamma$ 95% CL
$\zeta_1$ f.f.	$8 \cdot 10^{-14}$	$5 \cdot 10^{-14}$	$3 \cdot 10^{-14}$	$2.5 \cdot 10^{-14}$
$\zeta_1$ no f.f.	$2.5 \cdot 10^{-14}$	$1.5 \cdot 10^{-14}$	$9 \cdot 10^{-15}$	$7 \cdot 10^{-15}$
$\zeta_2$ f.f.	$2 \cdot 10^{-13}$	$1 \cdot 10^{-13}$	$6 \cdot 10^{-14}$	$4.5 \cdot 10^{-14}$
$\zeta_2$ no f.f.	$5 \cdot 10^{-14}$	$4 \cdot 10^{-14}$	$2 \cdot 10^{-14}$	$1.5 \cdot 10^{-14}$

- A large panel of extra-dimension models can be probed in **the multi-TeV range** (see Sylvain's talk)
- The form factor is not needed for any new physics scale **beyond  $\simeq 2$  TeV**  
*because of the forward detector acceptance (see slide 9)*

# Full amplitude computation for generic heavy charged fermions/vectors contributions (preliminary)

- The existence of new heavy charged particles will enhance the  $\gamma\gamma\gamma$  coupling at high mass via loops
- This enhancement can be parametrized by **only the mass and the effective charge**  $Q_{\text{eff}} = Q \cdot N^{1/4}$ ,  $N$  multiplicity (= values taken by the new degree of freedom)
- Generic full amplitude implementation for fermions and vectors implemented in FPMC
- Paper in preparation, preliminary  $M$ - $Q_{\text{eff}}$  exclusion plane



# Summary of the presentation


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- 1 Motivations and description of the proposed measurements
- 2 The  $WW\gamma\gamma$  and  $ZZ\gamma\gamma$  couplings cases in short
- 3 The  $\gamma\gamma\gamma\gamma$  couplings case in details
- 4 **Conclusion and plans**

# Where does proton tagging do better?

- Proton tagging allows a **very high background rejection** at the cost of a **smaller cross-section**
  - A single observation has **a high significance**
  - **Ideal to probe small deviations from the Standard Model like aQGC**  
*ex: new charged particles via loops, ADD gravity effects, ...*
  - Interesting “subleading” constraints on resonances searches at tree level  
*ex: new neutral particles at tree level*
  - Hard to be more quantitative: no  $\gamma\gamma\gamma$  limits from collider experiments yet (with intact protons or not)
- We reach sensitivities allowing to **probe directly a large class of new models**
  - **Extra-dimensions:** KK gravitons, dilaton, high  $\kappa$  untested domain (Randall-Sundrum model)
  - **Strongly-interacting composite states, monopoles:** generic searches of new heavy charged particles

- 
- **Forward proton tagging at the LHC seems promising to probe anomalous  $\gamma\gamma XX$  Gauge Couplings**
    - proton tagging associated with high energy object in the central detector allow to **highly suppress the background**
    - $WW\gamma\gamma$  and  $ZZ\gamma\gamma$  couplings sensitivity improvement by a factor up to 100 compared to latest the CMS measurement
    - $\gamma\gamma\gamma\gamma$  couplings: sensitivities around  $10^{-13} - 10^{-14} \text{ GeV}^{-4}$ , down to  $7 \cdot 10^{-15} \text{ GeV}^{-4}$  → **allows to probe directly a large panel of new physics models in the multi-TeV range** (KK gravitons, strongly-interacting heavy dilaton, ...) *(no previous constraints from collider experiments)*
  - **$\gamma\gamma\gamma\gamma$  couplings: a way to probe exotic heavy charged vectors/fermions in a completely model-independant way**

→ *very interesting subleading constraints in addition to direct searches at the LHC*

- 
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(no previous constraints from collider experiments)
  - **$\gamma\gamma\gamma\gamma$  couplings: a way to probe exotic heavy charged vectors/fermions in a completely model-independant way**
    - *very interesting subleading constraints in addition to direct searches at the LHC*
  - The cross-section of those potential new contributions seems to be very dependant on the spin value of the new particles
    - **sensitivities to higher spin resonances under study**  
(requires further theoretical development)



# Probing $\gamma\gamma XX$ anomalous gauge couplings with proton tagging at the LHC

Back-up

Matthias Saimpert<sup>1</sup>  
E. Chapon, S. Fichet, G. von Gersdorff,  
O. Kepka, B. Lenzi, C. Royon<sup>1</sup>

<sup>1</sup>CEA Saclay - Irfu/SPP

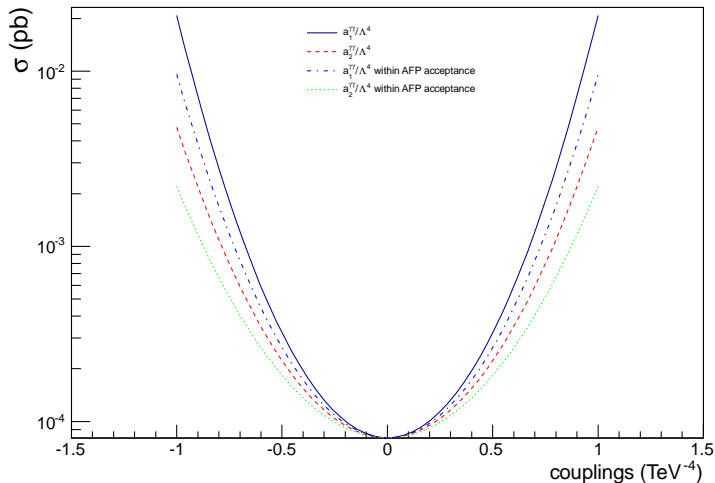
November, 5th 2014



# Integrated total cross-section against couplings for anomalous $\gamma\gamma\gamma$ couplings



Form factor applied



# Effective Field Theory cross-section of the $4\gamma$ couplings (G. Von Gersdorff)



## EFT OF 4 PHOTON INTERACTIONS

- ▶ Focus on **AAAA** (**AAZZ** and **AAWW** see [Chapon et al '12])
- ▶ EFT for 4-photon interaction contains two dim-8 structures

$$\mathcal{L}_{4\gamma} = \zeta_1 (F_{\mu\nu}F^{\mu\nu})^2 + \zeta_2 F_{\mu\nu}F^{\nu\rho}F_{\rho\sigma}F^{\sigma\mu}$$

- ▶ Cross section has a simple form

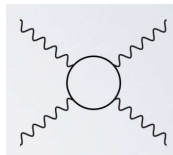
$$\frac{d\sigma}{d\Omega} = \frac{1}{16\pi^2 s} (s^2 + t^2 + st)^2 [48\zeta_1^2 + 40\zeta_1\zeta_2 + 11\zeta_2^2]$$

- ▶ Unitarity breaks down for  $\zeta_i s^2 \gtrsim 2\pi$
- ▶ Demanding unitarity for **LHC energies**  $\Rightarrow \zeta_i \lesssim 10^{-10} \text{GeV}^{-4}$
- ▶ In explicit models EFT breaks down before that!
- ▶ LHC sensitivities to  $\zeta_i$  are  $\sim 10^{4-5}$  **better** than unitarity bound

# New physics contributions to $4\gamma$ couplings

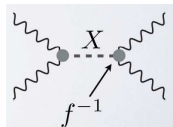
## 1 New charged particles via loops

- Effective coupling only depends on the mass, charge and spin :  
 $\zeta_i^\gamma \propto c_i^s Q^4 m^{-4}$
- Example: top partners

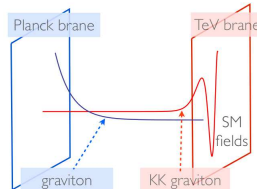


## 2 New neutral particles at tree level

- Effective coupling depends on mass, spin and the non-renormalizable  $\gamma\gamma X$  coupling  $\zeta_i^\gamma \propto b_i^s f^{-2} m^{-2}$
- Example: KK gravitons, dilaton (warped extra-dimension)



if coupling  $\simeq \text{TeV}$  and  $m_{KK} \simeq \text{few TeV}$ ,  
 $\zeta_i^\gamma \simeq 10^{-14}-10^{-13} \text{ GeV}^{-4}$  achievable,  
which we are sensitive



# Conversion, fake and efficiency reconstruction rates



- Inputs from the **ECFA ATLAS studies**
- **Photon conversion factors:** 15% in the barrel, 30% in the end-caps

- **Photon and electron reconstruction efficiency:**

$$Eff(p_T) = 0.76 - 1.98 \exp\left(\frac{-p_T}{16.1(\text{GeV})}\right)$$

- **Photon fake factors:** 1% for electron

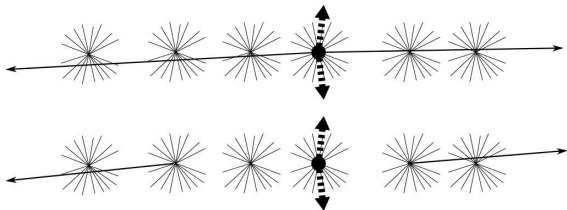
*European Strategy studies*

- **Fake photon  $p_T$  for jets:** gaussian draw (Mean=75%,  $\sigma=13\%$ ) on the jet  $p_T$  and use of

$$Eff_{fake}(p_T) = 0.0093 \exp\left(\frac{-\min(p_T, 200\text{GeV})}{17.5(\text{GeV})}\right)$$

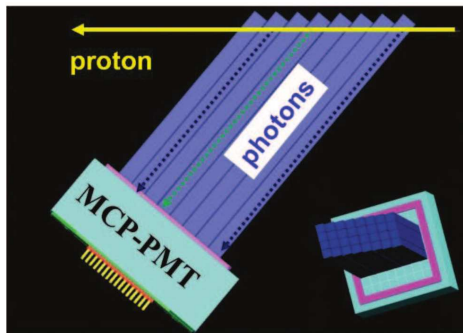
**almost no fake  $\gamma$  from jets at very high  $p_T$**

# Forward detectors measurements



- **Proton missing mass** measurement with 3% resolution in case of double tag
- **It matches the central  $\gamma\gamma$  mass for signal.** Can match as well for pile-up backgrounds as a statistical fluctuation
- **Double tag probability** from pile-up protons on the forward detectors (no missing mass requirement) :  
32% ( $\mu = 50$ )    66% ( $\mu = 100$ )    93% ( $\mu = 200$ )

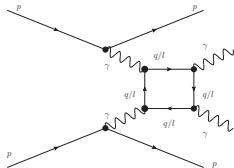
# Forward timing detectors : inefficiencies due to pile-up protons



Inefficiencies - 2mm bar detector										
Bar	1	2	3	4	5	6	7	8	9	10
$\mu = 50$	0.129	0.085	0.067	0.057	0.049	0.046	0.043	0.040	0.036	0.011
$\mu = 100$	0.185	0.122	0.097	0.082	0.071	0.066	0.062	0.057	0.051	0.016

*M. Saimpert. Search for new states of matter with the ATLAS experiment at the LHC, Master Thesis MINES ParisTech (2013)*

# SM QED exclusive $\gamma\gamma$ production



- Different loop contributions: fermions (quarks, leptons), vectors (W)
- **W loop contribution** and **massive fermions** added to the process in FPMC rev.913 (negligible at low mass but not at high mass, usually not included in the MCs)
- Interferences SM/Exotics added for the full amplitude calculation of new heavy charged particles



## The BSM amplitudes



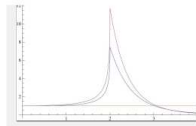
- Loops of spin 0, 1/2, 1 new **electric** particles contribute to  $4\gamma$ . Because all vertices are fixed by gauge invariance, the NP contributions depend **only** on spin, mass and electric charge!  $\Rightarrow$  Very **model-independent**

- For example in the effective theory limit :  $\zeta_i^\gamma = \alpha_{\text{em}}^2 Q^4 m^{-4} N c_{i,s}$

$$c_{1,s} = \begin{cases} \frac{1}{288} & s = 0 \\ -\frac{1}{36} & s = \frac{1}{2} \\ -\frac{5}{32} & s = 1 \end{cases}, \quad c_{2,s} = \begin{cases} \frac{1}{360} & s = 0 \\ \frac{7}{90} & s = \frac{1}{2} \\ \frac{27}{40} & s = 1 \end{cases}$$

Scalar loops are smaller !

- Full amplitudes** for fermions and vectors are now implemented in FPMC.
- Amplitudes get **enhanced** near the threshold







## The SM background



- All electric particles of the SM contribute : **leptons, quarks and W bosons**
- The imaginary part of certain W helicity amplitudes grows with the energy, while the fermion amplitudes are finite. Background is dominated by the **W loop**
- When the new particle is real, it **interferes** with the W loop.
  - ➔ On-shell NP signal **enhanced** by SM interference
- All SM background amplitudes are implemented in FPMC (+ switches to separately turn off them)
- One can check that SM fermions contributions are negligible.
  - ➔ Keeping only the W loop provides a huge gain of CPU time !

# Full amplitude computation for generic heavy charged fermions/vectors contributions (preliminary)



- **Link full amplitude - effective field theory**

$$\zeta_i^\gamma = c_i^s Q_{eff}^4 m^{-4} \alpha_{em}^2, c_i \simeq 0.01 (0.1) \text{ for fermions (vectors)}$$

- Typical sensitivity with the **full amplitude calculation**

$$M = 800 \text{ GeV}, Q_{eff} > 7 (4) \text{ for fermions (vectors)}$$

- Gives a coupling of  $\simeq \mathbf{3.10}^{-15}$  in terms of  $\zeta_i$

- Same order of magnitude than the sensitivity we had using the effective field theory  $\rightarrow$  **successful cross-check of the method**

## Conclusion: additional information

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- Effective field theory:  $5\sigma$  discovery with less luminosity ( $1 \text{ fb}^{-1}$ ,  $10 \text{ fb}^{-1}$ ,  $50 \text{ fb}^{-1}$ ):  $7 \cdot 10^{-13}$ ,  $2 \cdot 10^{-13}$ ,  $9 \cdot 10^{-14} \text{ GeV}^{-4}$
- $\gamma\gamma\gamma$  paper: S. Fichet et al Phys. Rev. D **89** (2014)
- More detailed paper including the full amplitude calculations for loop contributions and SM exclusive production update in preparation