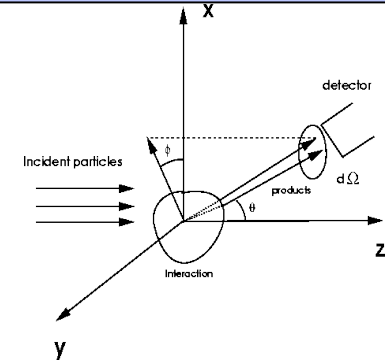


# Theory confronts experiment - Cross sections and decay rates

## Scattering in Quantum Mechanics



- Prepare state at  $t = -\infty$
- Time evolution (possibly scattering)
- Observe resulting system in state

$$|\psi_{in}(t = -\infty)\rangle = |i\rangle$$

$$|\psi_{in}(t = +\infty)\rangle = S |\psi_{in}(t = -\infty)\rangle$$

$$|\psi_{out}(t = +\infty)\rangle = |f\rangle$$

QM : probability amplitude :

$$\begin{aligned} \langle \psi_{out}(t = +\infty) | \psi_{in}(t = +\infty) \rangle &= \langle \psi_{out}(t = +\infty) | S | \psi_{in}(t = -\infty) \rangle \\ &= \langle f | S | i \rangle = S_{fi} \end{aligned}$$

$$S_{fi} = \delta_{fi} + iT_{fi}$$

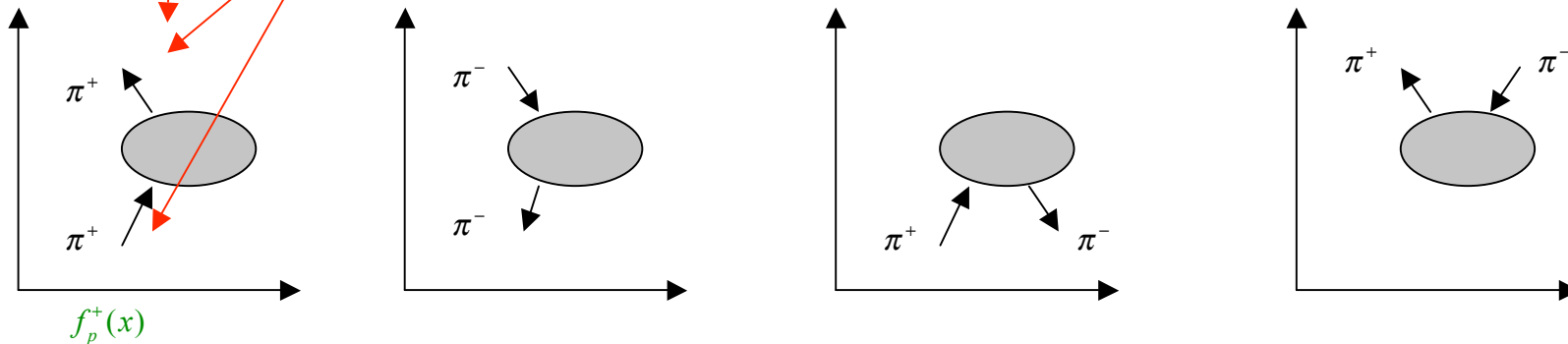
# S matrix for Klein Gordon scattering

$$S_{\mathbf{p}'_+, \mathbf{p}_+} = \langle \psi_{out}(t = +\infty) | \psi_{in}(t = +\infty) \rangle = \lim_{t \rightarrow \infty} \int d^3x f_{\mathbf{p}'_+}^{+*} i \partial_0 \psi(x)$$

$$\psi(x) = \phi(x) - \int d^4x' \Delta_F(x' - x) V(x') \phi(x') + \dots$$

$$\Delta_F(x' - x) = -i \int d^3p f_p^+(x') f_p^{+*}(x) \theta(t' - t) - i \int d^3p f_p^-(x') f_p^{-*}(x) \theta(t - t')$$

$$S_{\mathbf{p}'_+, \mathbf{p}_+} = \delta^3(\mathbf{p}'_+ - \mathbf{p}_+) - i \int d^4x' f_{\mathbf{p}'_+}^{+*}(x') V(x') f_{\mathbf{p}_+}^+(x') + \dots$$



# Feynman rules

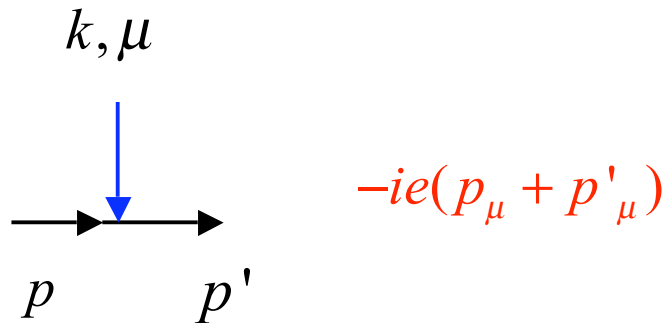
$$iT_{fi} = -i \int d^4 y f_{p+}^{+*}(y) V(y) f_{p+}^+(y) = i \int d^4 y f_{p+}^{+*}(y) ie(A^\mu \partial_\mu + \partial_\mu A^\mu) f_{p+}^+(y)$$

$$= -i \int d^4 y j_\mu^{fi} A^\mu$$

$$f_p^\pm = e^{\mp ip \cdot x} \frac{1}{\sqrt{2p^0 V}}$$

$$j_\mu^{fi} = -ie \left( f_{p+}^{+*}(y) [\partial_\mu f_{p+}^+(y)] - [\partial_\mu f_{p+}^{+*}(y)] f_{p+}^+(y) \right)$$

$$= -e(p_f + p_i)_\mu e^{i(p_f - p_i) \cdot y}$$



Feynman rule associated  
with Feynman diagram



## Fundamental experimental objects

$$\Gamma(a_1 \rightarrow b_1 b_2 \dots b_n)$$

**Decay width = 1/lifetime**

(Dimension  $1/T=M$ )

$$\sigma(a_1 a_2 \rightarrow b_1 b_2 \dots b_n)$$

**Cross section**

(Dimension  $L^2=M^{-2}$ )

Units "barn"

$$1 \text{ barn} = 10^2 \text{ fm}^2$$

$$1 \text{ mb} = 10^{-1} \text{ fm}^2 \quad \text{"milli"}$$

$$1 \text{ } \mu\text{b} = 10^{-4} \text{ fm}^2 \quad \text{"micro"}$$

$$1 \text{ nb} = 10^{-7} \text{ fm}^2 \quad \text{"nano"}$$

$$1 \text{ pb} = 10^{-10} \text{ fm}^2 \quad \text{"pico"}$$

$$1 \text{ fb} = 10^{-13} \text{ fm}^2 \quad \text{"fempto"}$$

(Natural Units  $1 \text{ GeV}^{-2} = 0.39 \text{ mb}$ )

# Fundamental experimental objects

$$\Gamma(a_1 \rightarrow b_1 b_2 \dots b_n)$$

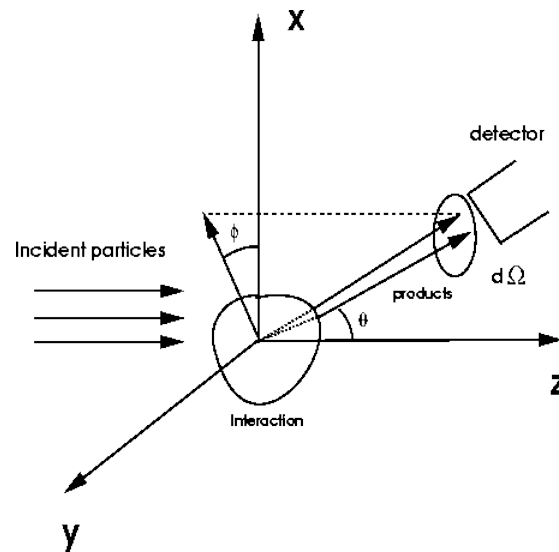
**Decay width = 1/lifetime**

(Dimension  $1/T=M$ )

$$\sigma(a_1 a_2 \rightarrow b_1 b_2 \dots b_n)$$

**Cross section**

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$$\text{Cross section} = \frac{\text{Transition rate} \times \text{Number of final states}}{\text{Initial flux}}$$

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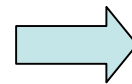
(Dimension  $L^2=M^{-2}$ )

Momenta of final state forms phase space

$$\text{Cross section} = \frac{\text{Transition rate} \times \text{Number of final states}}{\text{Initial flux}}$$

For a single particle the number of final states in volume  $V$  with momenta

in element  $d^3 p$  is  $\frac{V d^3 p}{(2\pi)^3}$



$$\prod_{i=1}^n \frac{V d^3 p}{(2\pi)^3}$$

# Fundamental experimental objects

$$\Gamma(a_1 \rightarrow b_1 b_2 \dots b_n)$$

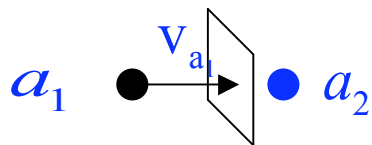
**Decay width = 1/lifetime**

$$\sigma(a_1 a_2 \rightarrow b_1 b_2 \dots b_n)$$

**Cross section**

$$\prod_{i=1}^n \frac{V d^3 p}{(2\pi)^3}$$

$$\text{Cross section} = \frac{\text{Transition rate} \times \text{Number of final states}}{\text{Initial flux}}$$



(Lab frame)

$$\frac{|v_{a1}|}{V} \times \frac{1}{V}$$

# particles passing through unit area in unit time

# target particles per unit volume

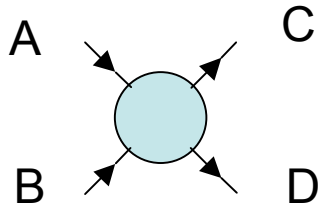


## The transition rate

$$T_{fi} = -\int d^4x \phi_f^*(x) V(x) \phi_i(x) + \dots$$

$$\phi_{i,f} \rightarrow f_p^\pm = e^{\mp ip \cdot x} \frac{1}{\sqrt{2p^0 V}} \equiv \frac{N}{\sqrt{V}} e^{\mp ip \cdot x}$$

e.g.



Transition rate per unit volume

$$W_{fi} = \frac{|T_{fi}|^2}{TV}$$

$$\phi_{f,i} = e^{\mp ip \cdot x}$$

$$T_{fi} = -\frac{N_A N_B N_C N_D}{V^2} (2\pi)^4 \delta^4(p_C + p_D - p_A - p_B) \mathfrak{M}_{fi}$$

$$W_{fi} = (2\pi)^4 \frac{\delta^4(p_C + p_D - p_A - p_B) |\mathfrak{M}|^2}{V^4} \left( \frac{1}{2E_A} \right) \left( \frac{1}{2E_B} \right) \left( \frac{1}{2E_C} \right) \left( \frac{1}{2E_D} \right)$$

## The cross section

Cross section =

Transition rate x Number of final states

Initial flux

$$d\sigma = \frac{V^2}{|\mathbf{v}_A| 2E_A 2E_B} \frac{1}{V^4} |\mathfrak{M}|^2 \frac{(2\pi)^4}{(2\pi)^6} \delta^4(p_C + p_D - p_A - p_B) \frac{d^3 p_C}{2E_C} \frac{d^3 p_D}{2E_D} V^2$$

$$d\sigma = \frac{|\mathfrak{M}|^2}{F} dQ$$

$$dQ = (2\pi)^4 \delta^4(p_C + p_D - p_A - p_B) \frac{d^3 p_C}{(2\pi)^3 2E_C} \frac{d^3 p_D}{(2\pi)^3 2E_D}$$

Lorentz  
Invariant  
Phase  
space

$$F = |\mathbf{v}_A| 2E_A 2E_B \\ = 4((p_A \cdot p_B)^2 - m_A^2 m_B^2)^{1/2}$$

The decay rate

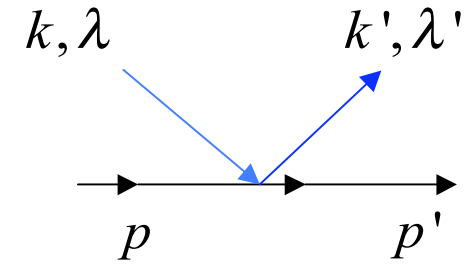
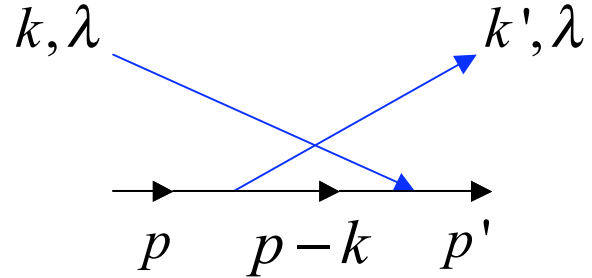
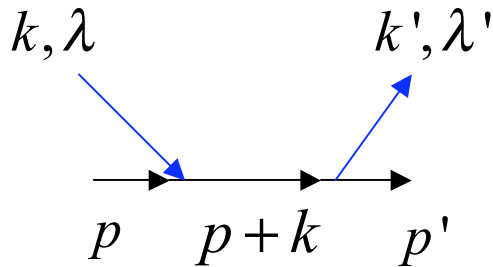
$$d\Gamma = \frac{1}{2E_A} |\mathfrak{M}|^2 dQ$$

$$dQ = (2\pi)^4 \delta^4(p_A - p_{B_1} \dots - p_{B_n}) \frac{d^3 p_{B_1}}{(2\pi)^3 2E_{B_1}} \dots \frac{d^3 p_{B_n}}{(2\pi)^3 2E_{B_n}}$$



Compton scattering of a  $\pi$  meson

$$\gamma\pi \rightarrow \gamma\pi$$

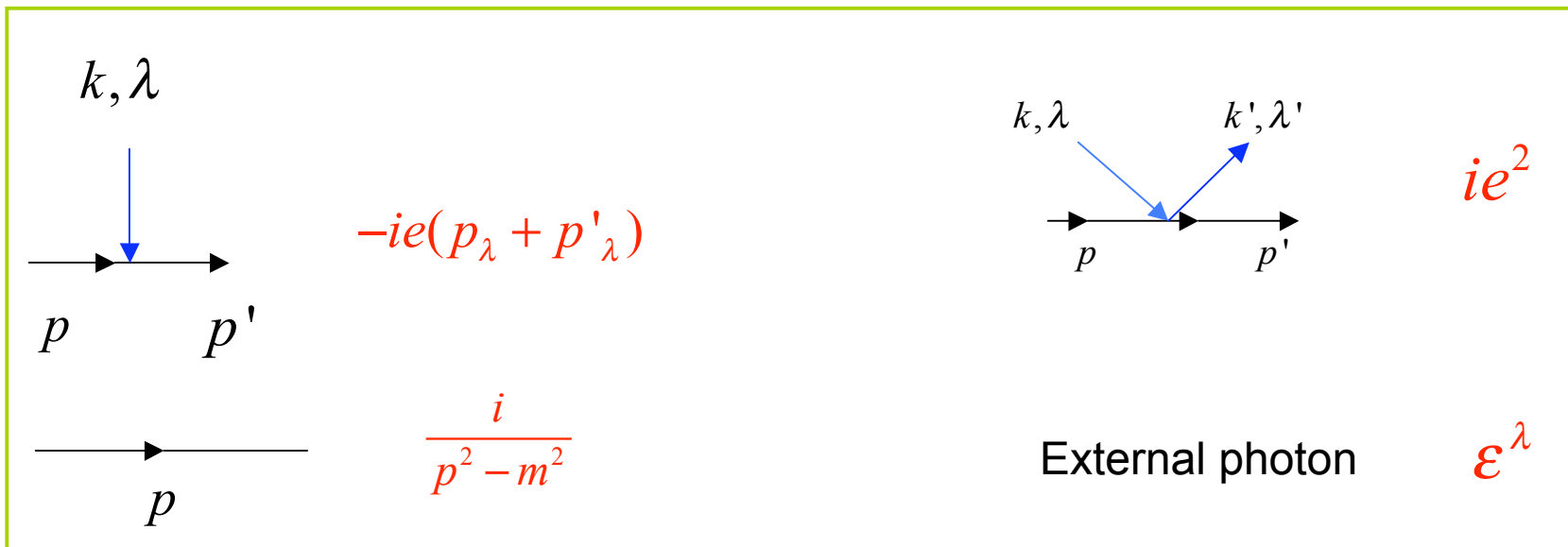


Feynman rules

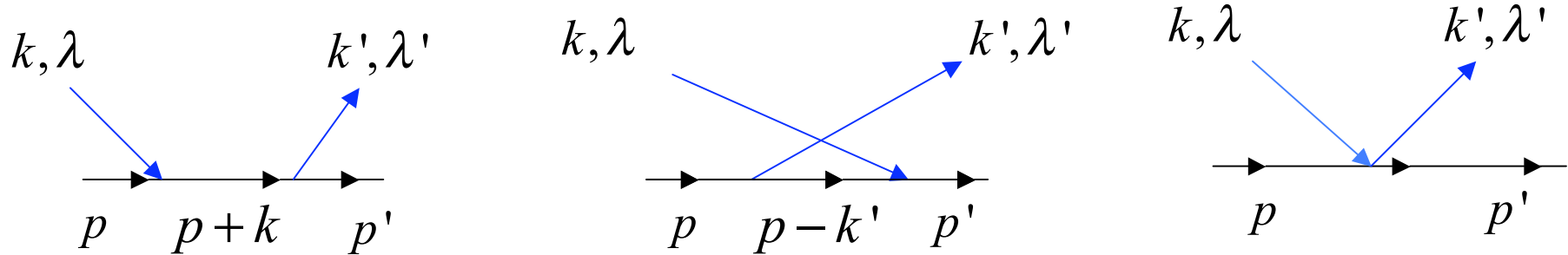
Klein Gordon

$$(\partial_\mu \partial^\mu + m^2)\psi = -V\psi$$

$$V = -ie(\partial_\mu A^\mu + A^\mu \partial_\mu) - e^2 A^2$$



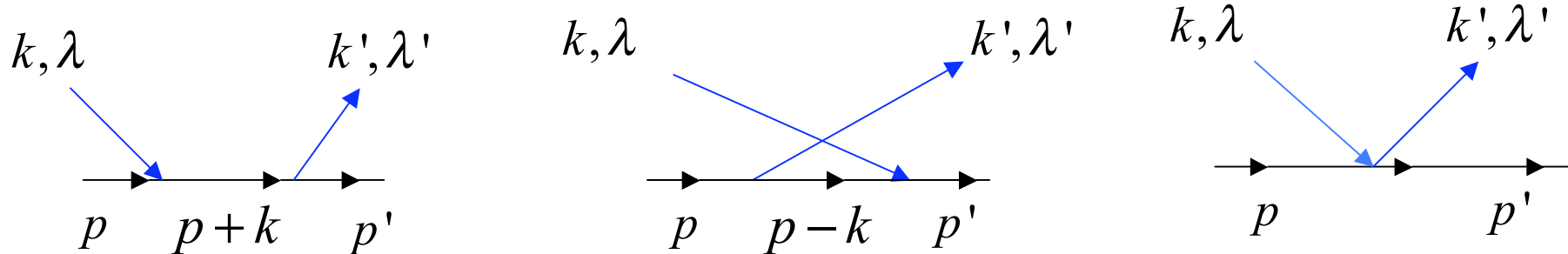
## Compton scattering of a $\pi$ meson



$$i\mathfrak{M}_{fi} = (-ie)^2 \left[ \varepsilon \cdot (2p+k) \frac{i}{(p+k)^2 - m^2} \varepsilon' \cdot (2p'+k') \right. \\ \left. + \varepsilon \cdot (2p'-k) \frac{i}{(p-k')^2 - m^2} \varepsilon' \cdot (2p-k') - 2i\varepsilon \cdot \varepsilon' \right]$$

$$d\sigma = \frac{V^2}{|\mathbf{v}_A| 2E_A 2E_B} \frac{1}{V^4} |\mathfrak{M}|^2 \frac{(2\pi)^4}{(2\pi)^6} \delta^4(p_C + p_D - p_A - p_B) \frac{d^3 p_C}{2E_C} \frac{d^3 p_D}{2E_D} V^2$$

# Compton scattering of a $\pi$ meson



$$\mathfrak{M}_{fi} = \varepsilon \cdot (2p + k) \frac{i}{(p + k)^2 - m^2} \varepsilon' \cdot (2p' + k')$$

$$+ \varepsilon \cdot (2p' - k') \frac{i}{(p - k)^2 - m^2} \varepsilon' \cdot (2p - k) - 2i\varepsilon \cdot \varepsilon'$$

$$\left( \frac{d\sigma}{d\Omega} \right)_{lab} = \frac{\alpha^2}{m^2} \frac{(\varepsilon \cdot \varepsilon')^2}{\left[ 1 + \frac{k}{m} (1 - \cos \theta) \right]^2}$$

( $\varepsilon \cdot p = \varepsilon' \cdot p = 0$  gauge)

$$\sigma_{total} |_{k=0} = \int \frac{d\sigma}{d\Omega} d\Omega = \frac{8\pi\alpha^2}{3m_\pi^2} \approx 8 \cdot 10^{-2} \text{ GeV}^{-2} = 3 \cdot 10^{-2} \text{ mb}$$

$$\sigma_{total} |_{k/m \gg 1} \approx \frac{2\pi\alpha^2}{mk}$$





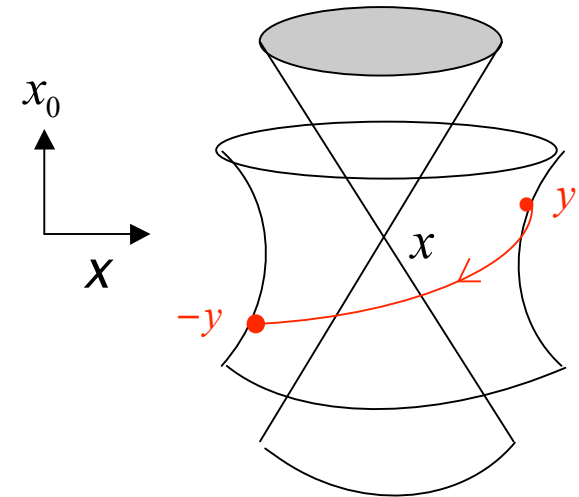
# Causality?

QM :  $U(x'-x) \propto e^{-m\sqrt{(x'-x)^2 - (t'-t)^2}}$  ✗

See Peskin & Schroeder  
"Quantum Field Theory" p28

Field theory :

$$\begin{aligned} \Delta_F(x'-x) &= -i \int \frac{d^3 p}{(2\pi)^3 2\omega_p} e^{-i\omega_p |t'-t| - i\mathbf{p}\cdot(\mathbf{x}'-\mathbf{x})} \\ &= D(x-y) - D(y-x) \end{aligned}$$



When  $(x-y)^2 < 0$ , we can perform a Lorentz transformation taking  $(x-y) \rightarrow -(x-y)$   
...causality preserved  $(e^{-m|r|} - e^{-m|r|})$

No (continuous) transformation possible for  $(x-y)^2 > 0$   
...and amplitude nonvanishing  $(e^{-imt} - e^{imt})$

