

# On distinguishing the direct and spontaneous CP violation in 2HDM

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## Motivation

- ▶ Sources of CP violation – a fundamental question in particle physics
- ▶ **CP conservation** – the Lagrangian **and** vacuum conserve CP
- ▶ **CP violation** – explicit or spontaneous
- ▶ 2HDM allows both the explicit and spontaneous violation of CP symmetry
- ▶ CP violating weak-basis invariants
- ▶ To distinguish spontaneous and explicit violation – comparison of various invariants

## 2 Higgs Doublet Model

- We consider the general 2 Higgs Doublet Model (without fermions):

$$\begin{aligned}\mathcal{L} &= \mathcal{L}_{gauge} + (D_\mu \Phi_1)^\dagger (D^\mu \Phi_1) + (D_\mu \Phi_2)^\dagger (D^\mu \Phi_2) - V, \\ V &= \frac{\lambda_1}{2} (\Phi_1^\dagger \Phi_1)^2 + \frac{\lambda_2}{2} (\Phi_2^\dagger \Phi_2)^2 + \lambda_3 (\Phi_1^\dagger \Phi_1) (\Phi_2^\dagger \Phi_2) + \lambda_4 (\Phi_1^\dagger \Phi_2) (\Phi_2^\dagger \Phi_1) \\ &\quad + \left[ \frac{1}{2} \lambda_5 (\Phi_1^\dagger \Phi_2) (\Phi_1^\dagger \Phi_2) + \lambda_6 (\Phi_1^\dagger \Phi_1) (\Phi_1^\dagger \Phi_2) + \lambda_7 (\Phi_2^\dagger \Phi_2) (\Phi_1^\dagger \Phi_2) + H.c. \right] \\ &\quad - \frac{1}{2} m_{11}^2 (\Phi_1^\dagger \Phi_1) - \frac{1}{2} m_{22}^2 (\Phi_2^\dagger \Phi_2) - \frac{1}{2} m_{12}^2 (\Phi_1^\dagger \Phi_2) - \frac{1}{2} (m_{12}^2)^* (\Phi_2^\dagger \Phi_1)\end{aligned}$$

- Scalar doublets in **the Higgs (Georgi) Basis** (ie.  $\langle \Phi_2 \rangle = 0$ ):

$$\Phi_1 = \begin{pmatrix} G^+ \\ \frac{v + \eta_1 + iG^0}{\sqrt{2}} \end{pmatrix} \quad \Phi_2 = \begin{pmatrix} H^+ \\ \frac{\eta_2 + iA}{\sqrt{2}} \end{pmatrix}$$

- There are 2 charged physical Higgs fields:  $H^+$ ,  $H^-$  and 3 neutral:
- in case of CP conservation:  $h, H, A$
  - in case of CP violation:  $h_1, h_2, h_3$

## CP transformation

- ▶ In Lagrangian we must identify CP-invariant part in order to determine the CP transformation

[Branco, Lee '66]

- ▶ Let us consider the simplest CP transformation in 2HDM:

$$\Phi_1(\vec{x}, t) \rightarrow \Phi_1^\dagger(-\vec{x}, t), \quad \Phi_2(\vec{x}, t) \rightarrow \Phi_2^\dagger(-\vec{x}, t)$$

- ▶ Under this transformation the fields transform as follows:

$$\text{CP-even: } \eta_{1,2} \rightarrow \eta_{1,2}, \quad \text{CP-odd: } A \rightarrow -A$$

## Conditions for CP violation

CP-odd weak-basis invariants – invariant under weak basis transformation but change sign under CP transformation.

► Jarlskog invariant for quarks in SM ( $m_i$  - mass,  $V_{ij}$  - CKM matrix):

$$J \propto (m_t^2 - m_c^2)(m_t^2 - m_u^2)(m_c^2 - m_u^2)(m_b^2 - m_s^2)(m_b^2 - m_d^2)(m_s^2 - m_d^2) \\ \times \text{Im}(V_{ud}V_{cs}V_{us}^*V_{cd}^*)$$

If  $J = 0$  then CP is conserved.

[Jarlskog '85]

► Invariant for scalars in 2HDM ( $m_i$  - mass,  $T_{ij}$  - elements of rotation matrix):

$$J_1 \propto (m_1^2 - m_2^2)(m_1^2 - m_3^2)(m_2^2 - m_3^2)T_{11}T_{21}T_{31}$$

$J_1 = 0 \not\Rightarrow$  CP conservation

[Lavoura '94]

## J-invariants

► The squared mass matrix for neutral scalars  $(\eta_1, \eta_2, A)$  in the Higgs basis has following form:

In the Higgs basis:  $\lambda_i \rightarrow \Lambda_i, m_{ij}^2 \rightarrow \mu_{ij}^2$

$$\begin{pmatrix} M_{11} & M_{12} & M_{13} \\ M_{12} & M_{22} & M_{23} \\ M_{13} & M_{23} & M_{33} \end{pmatrix} = \begin{pmatrix} v^2 \Lambda_1 & v^2 \operatorname{Re} \Lambda_6 & -v^2 \operatorname{Im} \Lambda_6 \\ v^2 \operatorname{Re} \Lambda_6 & [v^2 \Lambda_{345} - \mu_{22}^2]/2 & -v^2 \operatorname{Im} \Lambda_5/2 \\ -v^2 \operatorname{Im} \Lambda_6 & -v^2 \operatorname{Im} \Lambda_5/2 & [v^2 \tilde{\Lambda}_{345} - \mu_{22}^2]/2 \end{pmatrix}$$

► The invariant constructed from mass matrix:

$$J_1 = M_{12}M_{13}(M_{22} - M_{33}) + M_{23}(M_{13}^2 - M_{12}^2) = -8v^6 \operatorname{Im}(\Lambda_5^* \Lambda_6^2)$$

- $M_{13}$  or  $M_{23}$  non-zero  $\Rightarrow$  CP-mixing between states, CP violation
- $M_{13} = M_{23} = 0 \Rightarrow$  **no CP-mixing, yet possible CP violation**

► Note that  $\Lambda_7$  does not appear in mass matrix and it can appear **only in interactions**  $\Rightarrow J_2, J_3$ :

$$J_2 = -4v^4 \operatorname{Im}(\Lambda_5^* \Lambda_7^2), \quad J_3 = 2\sqrt{2}v^3 \operatorname{Im}(\Lambda_7^* \Lambda_6)$$

## $J$ -invariants - summary

From mass matrix:

$$J_1 \propto \text{Im}(\Lambda_5^* \Lambda_6^2)$$

From interaction:

$$J_2 \propto \text{Im}(\Lambda_5^* \Lambda_7^2), \quad J_3 \propto \text{Im}(\Lambda_7^* \Lambda_6)$$

- ▶ To have CP conserving model all  $J$ -invariants must be 0.
- ▶ CP violation:
  - If  $J_1 \neq 0$  then there is mixing between neutral states of different CP properties and physical states  $h_1, h_2, h_3$  have no defined CP.
  - If  $J_1 = 0$  and  $J_{2,3} \neq 0$  we have CP violation from interactions even if there is no mixing between states!
  - Note that  $J_i$  do not distinguish between spontaneous and explicit CP violation.
- ▶ Note that all that is valid not only in the Higgs Basis!

[Lavoura, Silva, Botella '94]



## I-invariants

- ▶ Other invariants formed from the parameters of the potential:

$$V = Y_{ab} \bar{\Phi}_a^\dagger \Phi_b + \frac{1}{2} Z_{abcd} (\Phi_a^\dagger \Phi_b) (\Phi_c^\dagger \Phi_d)$$

$$I_1 = \text{Im}(Z_{a\bar{c}}^{(1)} Z_{e\bar{b}}^{(1)} Z_{b\bar{e}c\bar{d}} Y_{d\bar{a}}) \quad I_2 = \text{Im}(Y_{a\bar{b}} Y_{c\bar{d}} Z_{b\bar{a}d\bar{f}} Z_{f\bar{c}}^{(1)})$$

$$I_3 = \text{Im}(Z_{a\bar{c}b\bar{d}} Z_{b\bar{f}}^{(1)} Z_{d\bar{h}}^{(1)} Z_{f\bar{a}j\bar{k}} Z_{k\bar{j}m\bar{n}} Z_{n\bar{m}h\bar{c}}), \quad I_4 = \text{Im}(Z_{a\bar{c}b\bar{d}} Z_{c\bar{e}d\bar{g}} Z_{e\bar{h}f\bar{q}} Y_{g\bar{a}} Y_{h\bar{b}} Y_{q\bar{f}})$$

- ▶ 4 invariants formed from quartic and quadratic parameters of  $V$
- ▶ The **potential** is **explicitly CP conserving**  $\Leftrightarrow$  all  $I_i$  are 0  
 $\Leftrightarrow$  there exists „real basis” of fields  $\Phi_1, \Phi_2$  in which all  $\lambda_i, m_{ij}^2$  are real  
 $\Rightarrow$  **CP symmetry can be conserved or violated spontaneously**
- ▶ If  $\exists I_i \neq 0$  then **CP is explicitly violated.**

[Branco; Haber, Gunion '05]

## $J$ - and $I$ -invariants: a comparison

Model properties	J-invariants	I-invariants
CP explicitly violated	$\exists J_i \neq 0$	$\exists I_i \neq 0$
CP spontaneously violated	$\exists J_i \neq 0$	$\forall I_i = 0$
CP conserved	$\forall J_i = 0$	$\forall I_i = 0$

- ▶ if  $\forall J_i = 0$  then we have CP conserving case
- ▶ if  $\exists J_i \neq 0$  then we have CP violation
  - if  $\forall I_i = 0$  then we have spontaneous CP violation
  - if  $\exists I_i \neq 0$  then we have explicit CP violation

## CP violation without mixing: $J_1 = 0$

- ▶ Let's consider a special case when  $\text{Im } \Lambda_5 = \text{Im } \Lambda_6 = 0$ . The neutral mass matrix in basis  $(\eta_1, \eta_2, A)$  is:

$$\begin{pmatrix} v^2 \Lambda_1 & v^2 \text{Re } \Lambda_6 & 0 \\ v^2 \text{Re } \Lambda_6 & [v^2 \Lambda_{345} - \mu_{22}^2]/2 & 0 \\ 0 & 0 & [v^2 \tilde{\Lambda}_{345} - \mu_{22}^2]/2 \end{pmatrix} \Rightarrow J_1 \sim \text{Im}(\Lambda_5^* \Lambda_6^2) = 0$$

- ▶ After diagonalization we get states  $h, H$  and  $A$  with defined CP properties - there is **no mixing between states of different CP**.
- ▶ The interactions with gauge bosons are like **in CP conserving case** ( $\alpha$  - mixing angle between  $\eta_1$  and  $\eta_2$ ,  $\chi_i$  - relative coupling with respect to SM):

$$hW^+W^- \propto \chi_1^V = \cos \alpha, \quad HW^+W^- \propto \chi_2^V = \sin \alpha,$$

$$AW^+W^- \propto \chi_3^V = 0$$

- ▶ Note, that if there is a mixing then  $\chi_3^V \neq 0$

## CP violation without mixing - self-interactions

- ▶  $\text{Im } \Lambda_5 = \text{Im } \Lambda_6 = 0 \Rightarrow J_1 = 0$
- ▶  $\text{Im } \Lambda_7 \neq 0 \Rightarrow J_2, J_3 \neq 0$  - there is a sign of CP violation in the interaction!

$$AAA \propto -\text{Im} \Lambda_7 v, \quad Ahh \propto -\text{Im} \Lambda_7 v \sin^2 \alpha,$$

$$AHH \propto -\text{Im} \Lambda_7 v \cos^2 \alpha, \quad AhH \propto \text{Im} \Lambda_7 v \cos \alpha \sin \alpha,$$

$$AH^+H^- \propto -\text{Im} \Lambda_7 v$$

- ▶ No couplings with odd number of  $A$  in CP conserving case
- ▶ To conclude:
  - $J_1 = 0, J_{2,3} \neq 0$ : there is a CP violation in the model.
  - To determine whenever it is spontaneous or explicit violation we need to use I-invariants.

## Summary

- ▶ CP violation in 2HDM without fermions
- ▶ Study of sources of CP violation with focus on distinguishing explicit and spontaneous violation
- ▶ We found that both  $J_i$  and  $I_i$  are needed to distinguish sources
- ▶ Usual approach: CP violation  $\Leftrightarrow$  mixing between states of different CP properties (true for soft violation of  $Z_2$  symmetry)
- ▶ However, with CP conservation in the gauge interactions of scalars still possible CP violation in self-interactions  
 $\Rightarrow$  CP violation effects in vertices with odd number of A  
(eg.  $A \rightarrow H^+H^-$ )