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Charged Higgs effects on top spin correlations

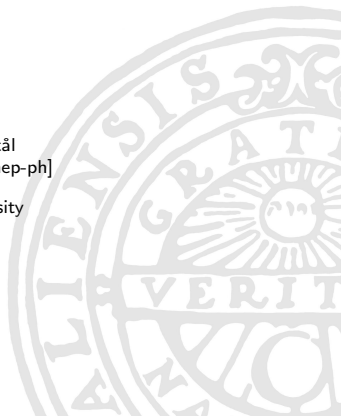
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Work done with

G. Ingelman, J. Rathsman and O. Stål
JHEP 0801:024,2008, arXiv:0710.5906 [hep-ph]

High Energy Physics, Uppsala University

Charged Higgs 2008
2008-09-18





Top spin correlations

Top quark pairs at hadron colliders

- Spin configurations
 - Singlet, $S = 0$: $t_{\uparrow}\bar{t}_{\downarrow}$
 - Triplet, $S = 1$: $t_{\uparrow}\bar{t}_{\uparrow}, t_{\uparrow}\bar{t}_{\downarrow}, t_{\downarrow}\bar{t}_{\downarrow}$
- Different production modes (qq, gg) gives different configurations
- Measurement of spin projection of one top gives, statistically, the spin of the other top
- Parton level correlations:

$$\hat{C}_{ij}(M_{t\bar{t}}^2) = \frac{\hat{\sigma}_{ij}(t_{\uparrow}\bar{t}_{\uparrow} + t_{\downarrow}\bar{t}_{\downarrow}) - \hat{\sigma}_{ij}(t_{\downarrow}\bar{t}_{\uparrow} + t_{\uparrow}\bar{t}_{\downarrow})}{\hat{\sigma}_{ij}(t_{\uparrow}\bar{t}_{\uparrow} + t_{\downarrow}\bar{t}_{\downarrow}) + \hat{\sigma}_{ij}(t_{\downarrow}\bar{t}_{\uparrow} + t_{\uparrow}\bar{t}_{\downarrow})}$$

- Helicity basis, spin quantized along momentum of t (\bar{t}) in $t\bar{t}$ rest frame



Top spin correlations continued

- Parton level correlations

$$\text{Threshold: } \hat{C}_{q\bar{q}}(4m_t^2) = -\frac{1}{3} \qquad \hat{C}_{gg}(4m_t^2) = 1$$

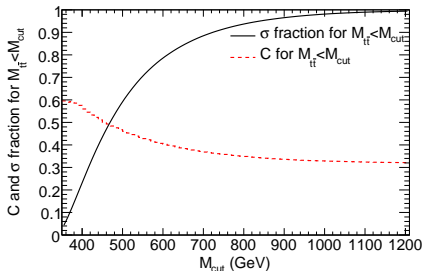
$$\text{Relativistic limit: } \hat{C}_{ij}(M_{t\bar{t}}^2 \gg 4m_t^2) \rightarrow -1$$

- Integrated correlations, NLO calculation [Bernreuter et al, Nucl.Phys. B690 (2004) 81-137]

Tevatron: $C = -0.352$ $q\bar{q}$ dominated **Still unmeasured**

LHC: $C = 0.326$ gg dominated

- Upper cut on $M_{t\bar{t}}^2$ gives more threshold like events
 - Higher correlations at LHC
 - Reduced statistics no problem since cross-section high





Measurement of spin and correlations

Top quark spin analyzing coefficients

- Fully polarized top quark at rest
- Spin along z-axis
- Weak decay, spin \rightarrow angular distributions of decay products

$$\frac{1}{\Gamma} \frac{d\Gamma}{d \cos \theta_i} = \frac{1 + \alpha_i \cos \theta_i}{2}$$

Analyzing particle	Analyzing coefficients α_i $W^+ (\omega = m_W^2/m_t^2)$	
b	$-\frac{1 - 2\omega}{1 + 2\omega}$	≈ -0.4
W^+	$\frac{1 - 2\omega}{1 + 2\omega}$	≈ 0.4
$l^+ (\bar{d})$	1	1
$\nu_l (u)$	$\frac{(1 - \omega)(1 - 11\omega - 2\omega^2) - 12\omega^2 \ln \omega}{(1 - \omega)^2(1 + 2\omega)}$	≈ -0.35



Measurement of spin and correlations

Spin correlations

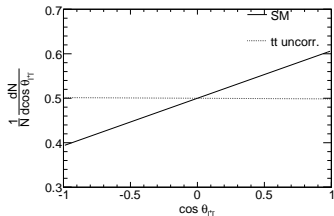
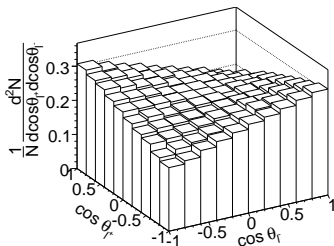
- No net polarization of single top
- Use correlation

$$\frac{1}{N} \frac{d^2 N}{d \cos \theta_i d \cos \theta_j} = \frac{1}{4} \left(1 + \mathcal{C} \alpha_i \alpha_j \cos \theta_i \cos \theta_j \right)$$

- θ_i, θ_j , angles in respective top rest frame
- Alternative one parameter correlation

$$\frac{1}{N} \frac{dN}{d \cos \theta_{ij}} = \frac{1}{2} \left(1 + \mathcal{D} \alpha_i \alpha_j \cos \theta_{ij} \right)$$

- $\cos \theta_{ij} = \hat{p}_i \cdot \hat{p}_j$ is “opening angle”
- Less sensitive to phase-space cuts
- \mathcal{D} is related to \mathcal{C}
- NLO calculation gives $\mathcal{D} = -0.24$ at LHC





Charged Higgs in top decays

- Light charged Higgs mediate top decay
- Spin 0 \rightarrow different spin analyzing coefficients
- Lagrangian for fermion part in 2HDM

$$\mathcal{L}_H = \frac{g_W}{2\sqrt{2}m_W} \sum_{\substack{\{u,c,t\} \\ \{d,s,b\}}} V_{ud} \left\{ H^+ \bar{u} \left[A(1 - \gamma_5) + B(1 + \gamma_5) \right] d + h.c. \right.$$
$$\left. + \frac{g_W}{2\sqrt{2}m_W} \sum_{\{e,\mu,\tau\}} \left[H^+ C \bar{\nu}_l (1 + \gamma_5) l + H^- C^* \bar{l} (1 - \gamma_5) \nu_l \right]. \right.$$

Coupling	2HDM (I)	2HDM (II)
A	$m_u \cot \beta$	$m_u \cot \beta$
B	$-m_d \cot \beta$	$m_d \tan \beta$
C	$m_l \cot \beta$	$m_l \tan \beta$

- With couplings



Charged Higgs in top decays

Spin analyzing coefficients

Analyzing particle	Analyzing coefficients α_j	
	W^+ ($\omega = m_W^2/m_t^2$)	H^+ ($\xi = m_{H^+}^2/m_t^2$)
b	$-\frac{1-2\omega}{1+2\omega}$	$-\frac{A^2-B^2}{A^2+B^2} f(\xi, A, B)$
W^+/H^+	$\frac{1-2\omega}{1+2\omega}$	$\frac{A^2-B^2}{A^2+B^2} f(\xi, A, B)$
l^+ (\bar{d})	1	$\frac{1-\xi^2+2\xi \ln \xi}{(1-\xi)^2} \frac{A^2-B^2}{A^2+B^2} f(\xi, A, B)$
ν_l (u)	$\frac{(1-\omega)(1-11\omega-2\omega^2)-12\omega^2 \ln \omega}{(1-\omega)^2(1+2\omega)}$	$-\frac{1-\xi^2+2\xi \ln \xi}{(1-\xi)^2} \frac{A^2-B^2}{A^2+B^2} f(\xi, A, B)$

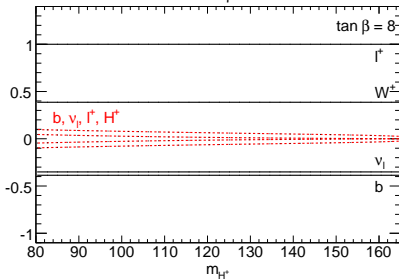
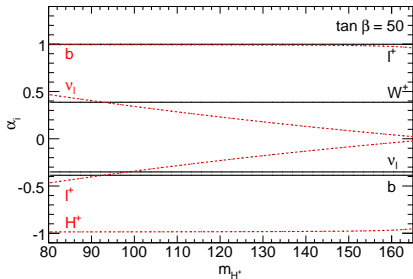
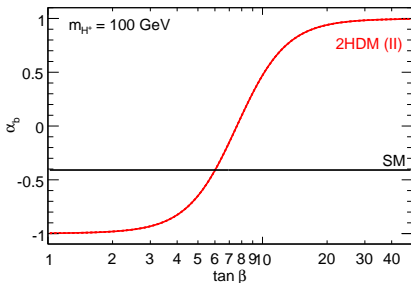
- Charged Higgs coefficients depend on $-\frac{A^2-B^2}{A^2+B^2}$
- Threshold factor $f(\xi, A, B) \simeq 1$ except for $m_{H^+} \rightarrow m_t$
- Remember in 2HDM(II): $A = m_u \cot \beta$ and $B = m_d \tan \beta$



Charged Higgs in top decays

Spin analyzing coefficients, continued

- l^+ most efficient particle in SM
- b and H^+ most efficient particles in 2HDM (II)
- Strong dependence on $\tan \beta$
- l^+ also depend on m_{H^+}



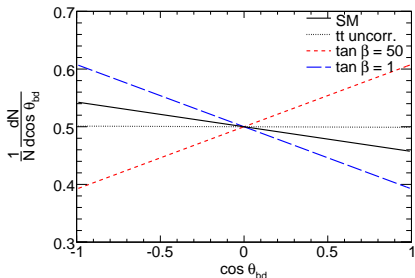
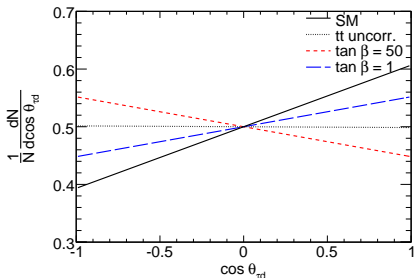


Parton level correlations

Best case

- $t \rightarrow b \quad H^+ / W^+ \rightarrow b \tau^+ \nu_\tau$
 $\bar{t} \rightarrow \bar{b} \quad W^- \rightarrow \bar{b} \bar{u} d$ + c.c. and α_j corresponding to $m_{H^+} = 80$ GeV
- \mathcal{D} -type correlations, $\mathcal{D} = -0.216$ at LHC

$$\frac{1}{N} \frac{dN}{d \cos \theta_{ij}} = \frac{1}{2} (1 + \mathcal{D} \alpha_i \alpha_j \cos \theta_{ij})$$



- Stable τ^+ , fully know final state and CM frame



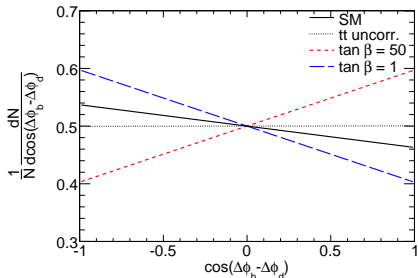
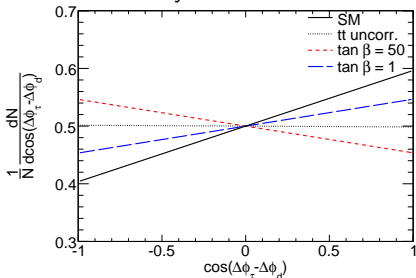
Parton level correlations

Azimuthal correlations

- Charged Higgs decays to $\tau^+ \nu_\tau$ and τ^+ to $X + \bar{\nu}_\tau$
- Center of mass frame not know
- Use hadronic W^+ and τ^+ to get transverse rest frame
- Use azimuthal angles and the correlation

$$\frac{1}{N} \frac{dN}{d \cos(\Delta\phi_i - \Delta\phi_j)} = \frac{1}{2} \left[1 + \mathcal{D}' \alpha_i \alpha_j \cos(\Delta\phi_i - \Delta\phi_j) \right].$$

- Numerically $\mathcal{D}' = 0.9\mathcal{D}$ at LHC

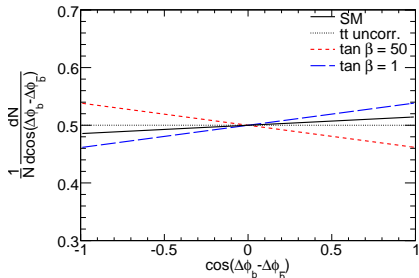
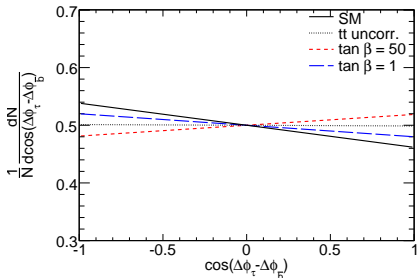




Parton level correlations

Azimuthal correlations, realistic variables

- Hard to identify d -jets
- Use b -jet (or least energetic light jet) on W^+ side
- Use τ^+ -jet or b -jet on H^+ side
- Realistic best case, all quantities measurable

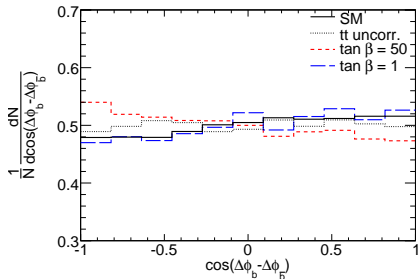
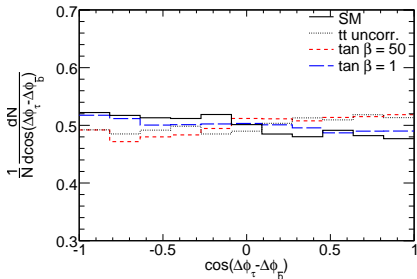




Hadron level correlations

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- Full $2 \rightarrow 6$ ME with MadGraph/MadEvent
 $p p \rightarrow (t \bar{t}) \rightarrow (b H^+ / W^+ \bar{b} W^-) \rightarrow b \tau^+ \nu_\tau \bar{b} \bar{u} d$
- τ^+ decay with Tauola
- Parton shower, hadronization and underlying event with Pythia
- k_\perp jet finding in $|\eta| < 5$ with $d_{cut} = 20$ GeV
- “Flavor tag” $\Delta R(\text{jet}, \text{truth}) < 0.4$ in $|\eta| < 2.5$
- W^+ candidate $|m_{jj} - m_{W^+}| < 10$ GeV, t candidate $|m_{jjb} - m_t| < 15$ GeV
- No background or detector effects





Summary and Conclusions

- Top quark spin correlations should be measurable at LHC
- In SM I^+ is the most efficient analyzer
- Charged Higgs in top decays alter angular distributions
- In 2HDM(II) b quark is the most efficient analyzer
- Charged Higgs decays to τ^+ and neutrinos so full reconstruction not possible
- Correlations with azimuthal angles can be constructed, \mathcal{D}'
- Realistic hadron-level correlations are small in both SM and 2HDM(II)