

Tree-level Vaccum stability in THDM R. Santos

NExT

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CHARGE BREAKING IN THE SM

Can the SM potential have a Normal and a CB minima simultaneously?





CHARGE BREAKING IN THE SM



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CHARGE BREAKING IN THE SM

$$m_1^2 = m_2^2 = \frac{g^2 v^2}{4}$$
$$m_3^2 = \frac{v^2}{4} (g^2 + g'^2 Y^2)$$
$$m_4^2 = 0$$

Gauge boson masses in the most general case

U(1) always survives in the SM!

Look, it's the PHOTON!

Electric charge is conserved!

$$v = \sqrt{v_1^2 + v_2^2 + v_3^2 + v_4^2}$$

The Weinberg angle is the same!

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CHARGE BREAKING IN THE THDM

<u>Conclusion</u>: No Charge Breaking in the SM!

What about CB in THDM?



Now we have two doublets



8 possible vevs









CHARGE BREAKING IN THE THDM

CB is possible in THDM! SUPPOSE WE LIVE IN A THDM! Are we in DANGER?!...

Three minimum field configurations

► NORMAL

$$\Phi_1 = \begin{pmatrix} 0 \\ v_1 \end{pmatrix} e \Phi_2 = \begin{pmatrix} 0 \\ v_2 \end{pmatrix}$$

CHARGE BREAKING

$$\Phi_1 = \begin{pmatrix} 0 \\ v'_1 \end{pmatrix} e \Phi_2 = \begin{pmatrix} \alpha \\ v'_2 \end{pmatrix}$$

CP BREAKING

$$\Phi_1 = \begin{pmatrix} 0 \\ v'_1 + i\delta \end{pmatrix} e \Phi_2 = \begin{pmatrix} 0 \\ v'_2 \end{pmatrix}$$

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FOR A SAFER THDM

Some algebra

 M^2

2v

H[±]

$$V_{CB} - V_{N} = \frac{M_{H^{\pm}}^{2}}{2v^{2}} \left[\left(v'_{1} v_{2} - v'_{2} v_{1} \right)^{2} + \alpha^{2} v_{1}^{2} \right] = \frac{1}{2} \frac{M_{H^{\pm}}^{2}}{2v^{2}} \left[\frac{\left(v'_{1} v_{2} - v'_{2} v_{1} \right)^{2} + \alpha^{2} v_{1}^{2}}{2v^{2}} \right]$$

Calculated at the Normal stationary point.

$$V_N \text{ is a MINIMUM} \implies \frac{M_{H^{\pm}}^2}{2v^2} > 0 \implies V_N < V_{CB}$$

Bonus V_{CB} is a Saddle Point No! V_N is below V_{CB} and tunnelling will not happen!



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However

If V_{CP} is a MINIMUM



Any competing N SP is a saddle point above it

Any competing CB SP is a saddle point above it

The THDM build on V_{CP} is also safe!







For charge breaking

$$V_{CB} - V_N = \frac{M_{H^{\pm}}^2}{2v^2} \left[(v'_1 v_2 - v'_2 v_1)^2 + \alpha^2 v_1^2 \right]$$

For CP breaking

$$V_{CP} - V_N = \frac{M_A^2}{2v^2} \left[\left(v''_1 v_2 - v''_2 v_1 \right)^2 + \delta^2 v_1^2 \right]$$

For 2 competing normal minima

$$V_{N_{2}} - V_{N_{1}} = \frac{1}{2} \left\{ \left(\frac{M_{H^{\pm}}^{2}}{v^{2}} \right)_{N_{1}} - \left(\frac{M_{H^{\pm}}^{2}}{v^{2}} \right)_{N_{2}} \right\} \left[\left(v''_{1} v_{2} - v''_{2} v_{1} \right)^{2} + \delta^{2} v_{1}^{2} \right]$$

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- 1. THDM bounded from below have no maxima.
- 2. THDM have at most two minima.
- 3. Minima of different nature never coexist.
- 4. CB and CP minima are uniquely determined.
- 5. If a THDM has only one normal minimum than this is the absolute minimum all other if they exist are saddle points and above it.

IVAN IVANOV HAS HELPED TO COMPLETE THE PICTURE! PROVED 2. AND THE PART OF 3. WE DIDN'T PROVE - THAT WHEN THE CP SP WAS ABOVE THE NORMAL MINIMUM IT WAS ALWAYS A SADDLE POINT — WE ONLY PROVED THE REVERSE STATEMENT WHICH IS OBVIOUS FROM THE MASS EXPRESSIONS

But how many normal minima?





NORMAL MINIMA

Two competing normal minima

$$V_{N_{2}} - V_{N_{1}} = \frac{1}{2} \left\{ \left(\frac{M_{H^{\pm}}^{2}}{v^{2}} \right)_{N_{1}} - \left(\frac{M_{H^{\pm}}^{2}}{v^{2}} \right)_{N_{2}} \right\} \left[\left(v''_{1} v_{2} - v''_{2} v_{1} \right)^{2} + \delta^{2} v_{1}^{2} \right]$$

The potential with the largest charged Higgs mass wins!







EQUAL BUT DIFFERENT?



<u>No!</u>

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 $\left(M_{h}^{2}\right)_{N_{1}}\neq\left(M_{h}^{2}\right)_{N_{2}}$

 $=V_{N_{1}}$

 V_{N_2}



THE POTENTIAL WE LOVE

The potential has 8 parameters:

It has a softly broken Z2 symmetry

It can break CP

It can break charge

But once we are at a Normal minimum, CB and CP will not occur. What about two normal minima for the same parameter set?





BAD NEWS?





Note however that this potential is fine for CP-violating minima- in this case the minimum is uniquely determined!

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EXACT Z₂ - THE PERFECT THDM FOR NORMAL MINIMA

$$m_{12}^2 = 0$$

$$\begin{cases} v_1^2 = \frac{2\lambda_2 m_1^2 - (\lambda_3 + \lambda_5)m_2^2}{(\lambda_3 + \lambda_5)^2 - 4\lambda_1\lambda_2} \\ v_2^2 = \frac{2\lambda_1 m_2^2 - (\lambda_3 + \lambda_5)m_1^2}{(\lambda_3 + \lambda_5)^2 - 4\lambda_1\lambda_2} \end{cases}$$



$$m_h^2 \propto \frac{2\lambda_1 m_2^2 - (\lambda_3 + \lambda_5)m_1^2}{2\lambda_1}$$

Just one! Normal minimum

 $m_{h}^{2} <$

charged

Using the condition for the potential to be bounded from below we show that it is always above and it is a saddle point

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Force the CP minima not to exist (minimum equations have no solution).

2 easy ("natural") ways of avoiding that the CP minimum equations have a solution:

- $m_{12}=0$ and $\lambda_5 \neq 0$ ($\Phi_1 \rightarrow -\Phi_1$; $\Phi_2 \rightarrow \Phi_2$)
- $\lambda_5 = 0$ and $m_{12} \neq 0$ ($\Phi_1 \rightarrow e^{i\theta} \Phi_1; \Phi_2 \rightarrow \Phi_2$) softy broken (otherwise axion)

Two 7-parameters models Z2 and global U(1)

Extending the symmetry to the fermions Implies flavour conservation at tree level.





CONCLUSIONS I





The conclusions are valid for all renormalizable THDM at tree level.

The soft broken (Z₂) potential is fine for competing minima of different nature and <u>for CP violating minima</u> (which are unique) but shows "problems" for two competing Normal minima. Exact Z₂ potential is doing great for normal minima but it doesn't allow for CP violation.

The stability of the Normal minimum is not guaranteed. At one-loop the CB minimum could be below the Normal one.





Thank you









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BASED ON

▶ PLB322 (1994) Tree Level Vacuum Stability in THDM

- PLB603 (2004) Stability of the tree level vacuum in THDM against charge or CP spontaneous violation
- ► PLB632 (2006) Charge and CP breaking in THDM
- ▶ PRD74 (2006) Stability of the normal vacuum in MHDM
- ▶ PLB652 (2007) Charge and CP breaking in THDM

