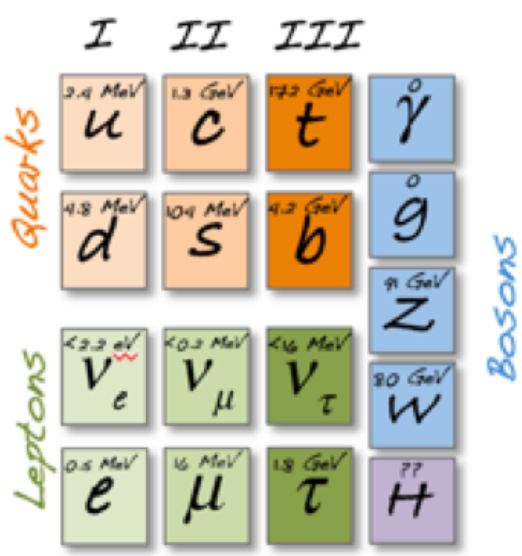
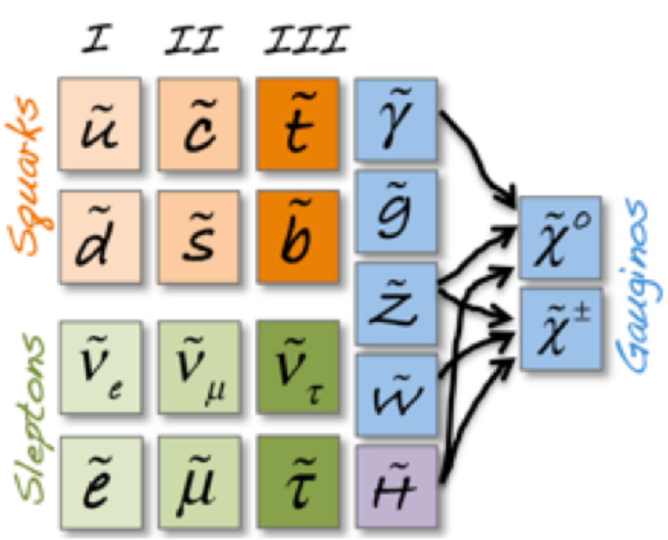
Searches for



Supersymmetric

Supersymmetric

particles in ATLAS



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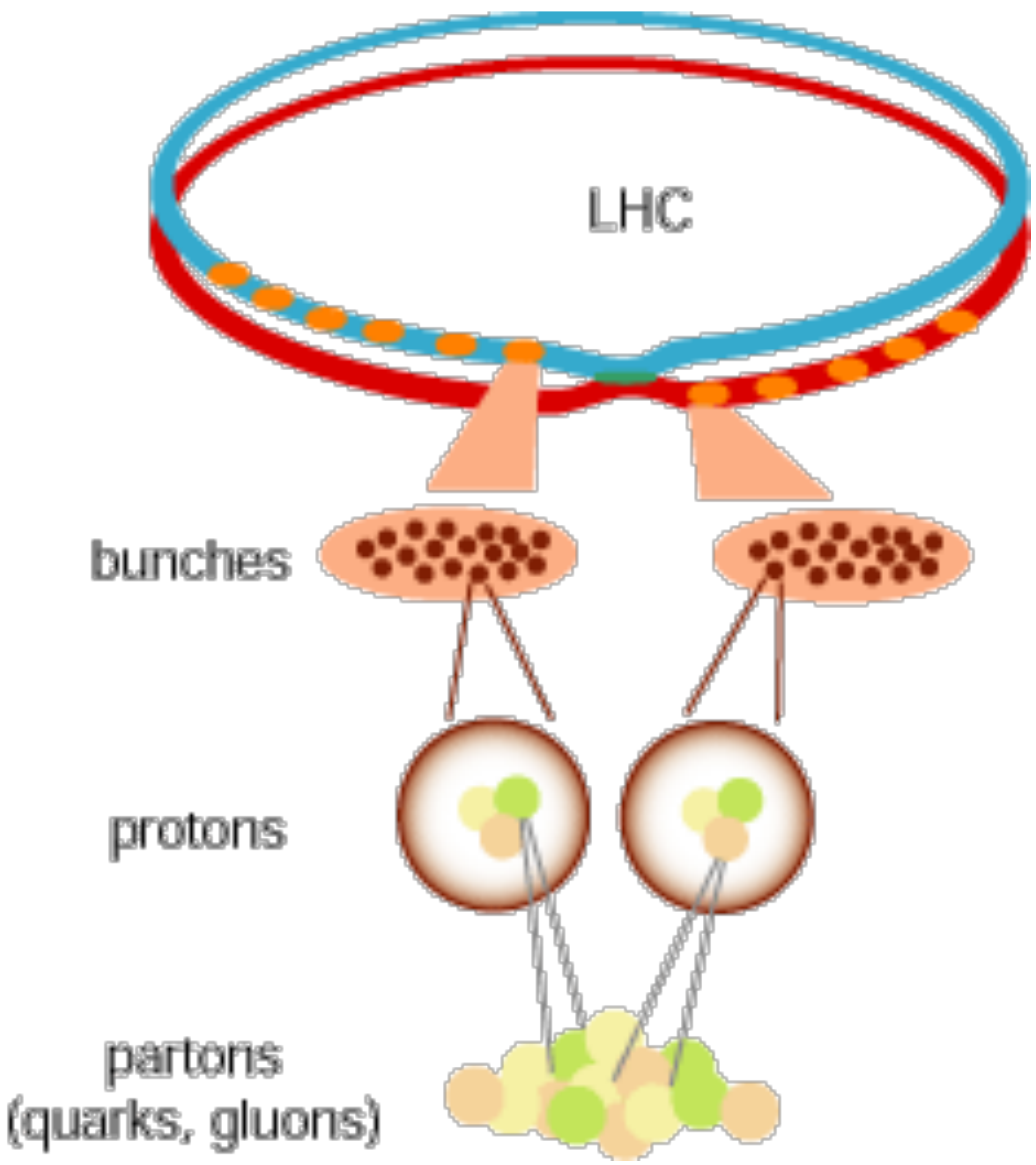
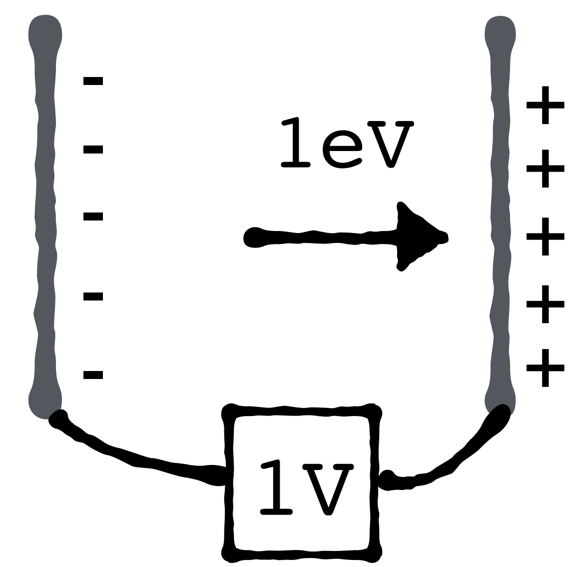
LHC SUMMARY



- 27Km tunnel, 100m underground
- 8.3T superconducting magnets
- proton-proton collisions
  - energy 7-8 TeV
  - bunches/beam 2808
  - protons/bunch 10^{11}
  - luminosity 10^{34} cm^{-2} s^{-1}
  - total crossing rate 40 MHz
  - collision rate $10^{7/8}$ Hz

$E_{LHC} = e\Delta V =$

7.000.000.000.000 eV



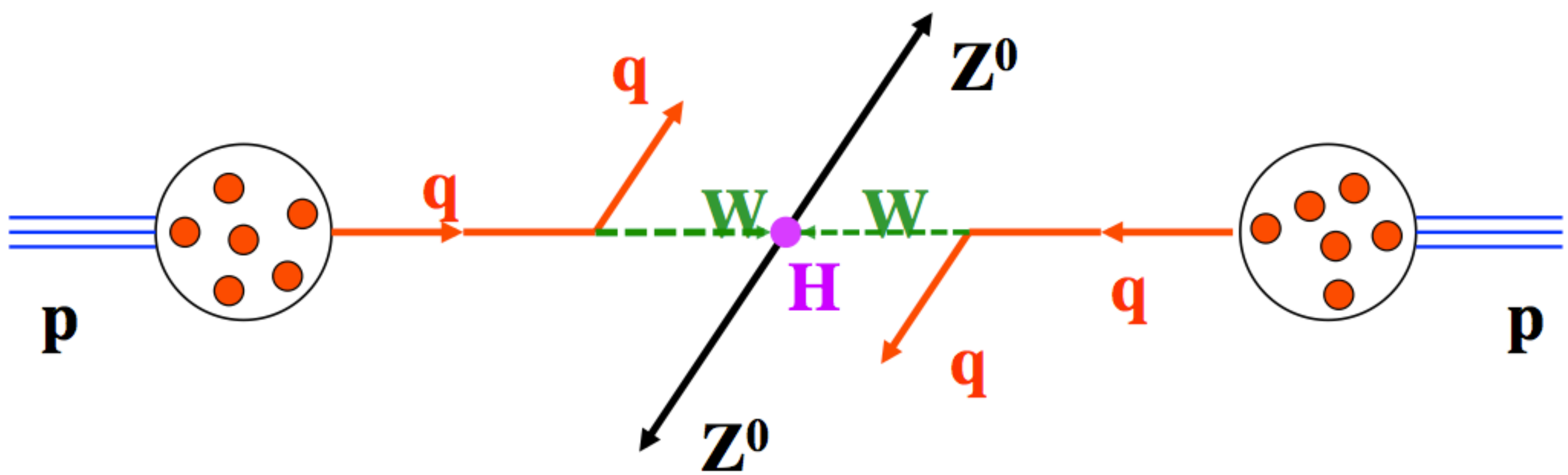
But the rate of New Physics is only ~10^{-5}

Need to select one particle in 10.000.000.000.000. !!

PP COLLISIONS

- QCD theory is expressed in terms of ***quarks*** and ***gluons***
  - But we are colliding composite objects, protons (uud)!
  - The useful energy (of the partons) is only a fraction of the total (of the protons)

$$\sqrt{\hat{s}} << \sqrt{s} = E_{cm} = 2E_p$$



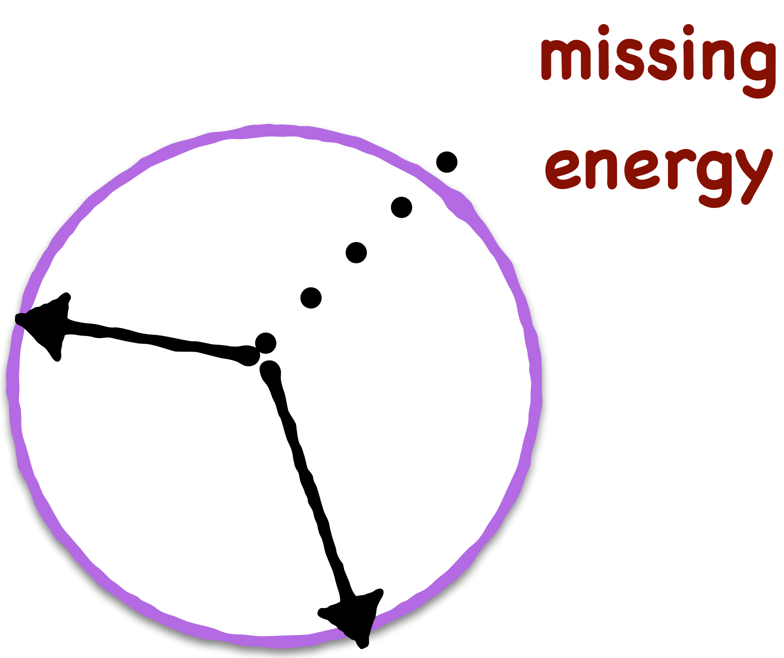
- Since the two processes happens at very different energies (length scales) we assume we can treat them independently: **factorization theorem**
  - The hard scattering is calculable in perturbative QCD and is independent on the long distance effects
  - The physics at larger (non-perturbative) scales is absorbed in some functions (PDFs) which are universal (only dependent on the incoming particles) and are measured by experiments

Kinematic Variables

- The natural variables to use to describe the collisions would be p, θ, Φ
  - But we don't know the longitudinal momentum of the initial patrons scattering
  - p_Z and E are not so useful (we cannot apply any conservation law)

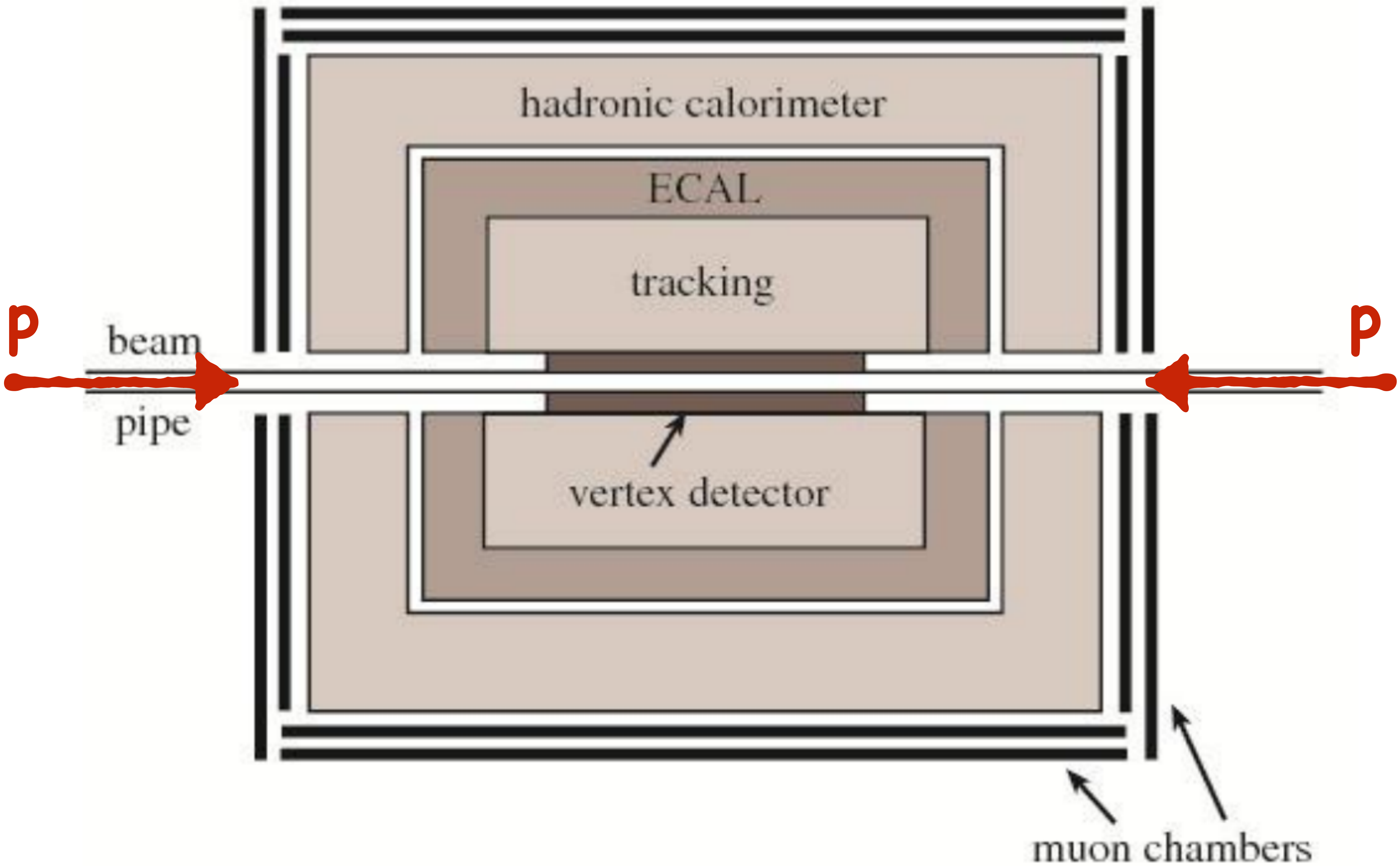
- What is still conserved are the quantities in the ***transverse plane*** : p_T, E_T
  - Before the collisions the transverse momentum/energy is zero

- Sometimes we can have an imbalance if we add up the momenta of all the particles that we see

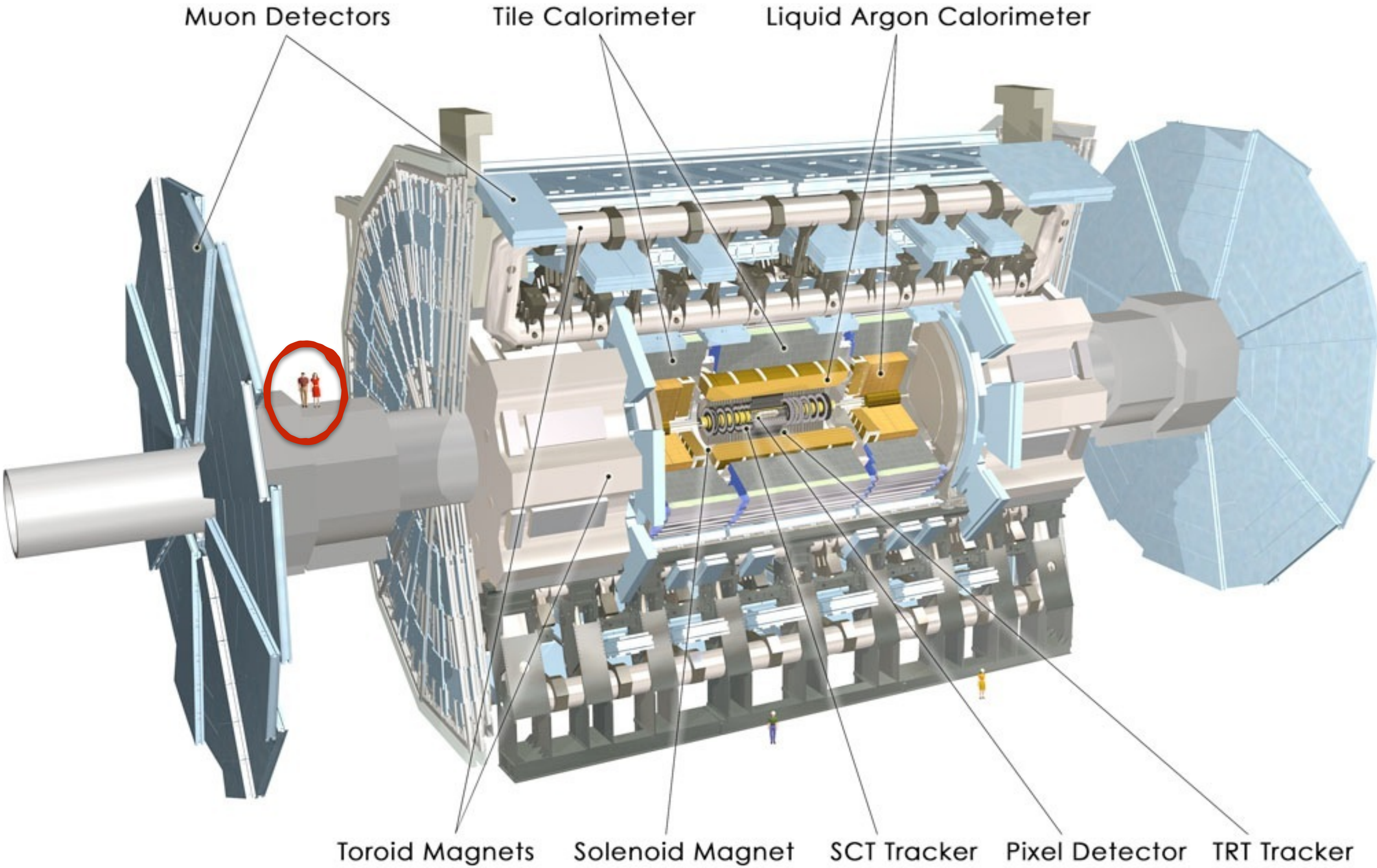


- This can happen if we have particles escaping the detectors undetected
- The *missing energy* quantity is defined to keep track of those

Basic concept of a general purpose detector



THE ATLAS DETECTOR



What can we "see"

- *Stable particles*, which are directly seen:

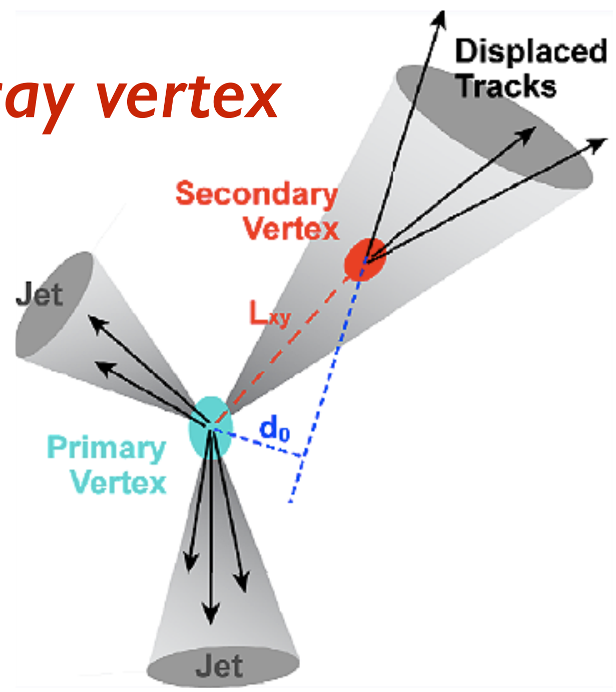
$$p, \bar{p}, e^{\pm}, \gamma$$

- *Quasi-stable particles*, with a long lifetime are also directly seen

$$n, \Lambda, K_L^0, \ldots, \mu^{\pm}, \pi^{\pm}, K^{\pm}, \ldots$$

- Particles with a short lifetime may display a *secondary decay vertex*

$$\tau \sim 10^{-12} \qquad B^{0,\pm}, D^{0,\pm}, \tau, \ldots$$



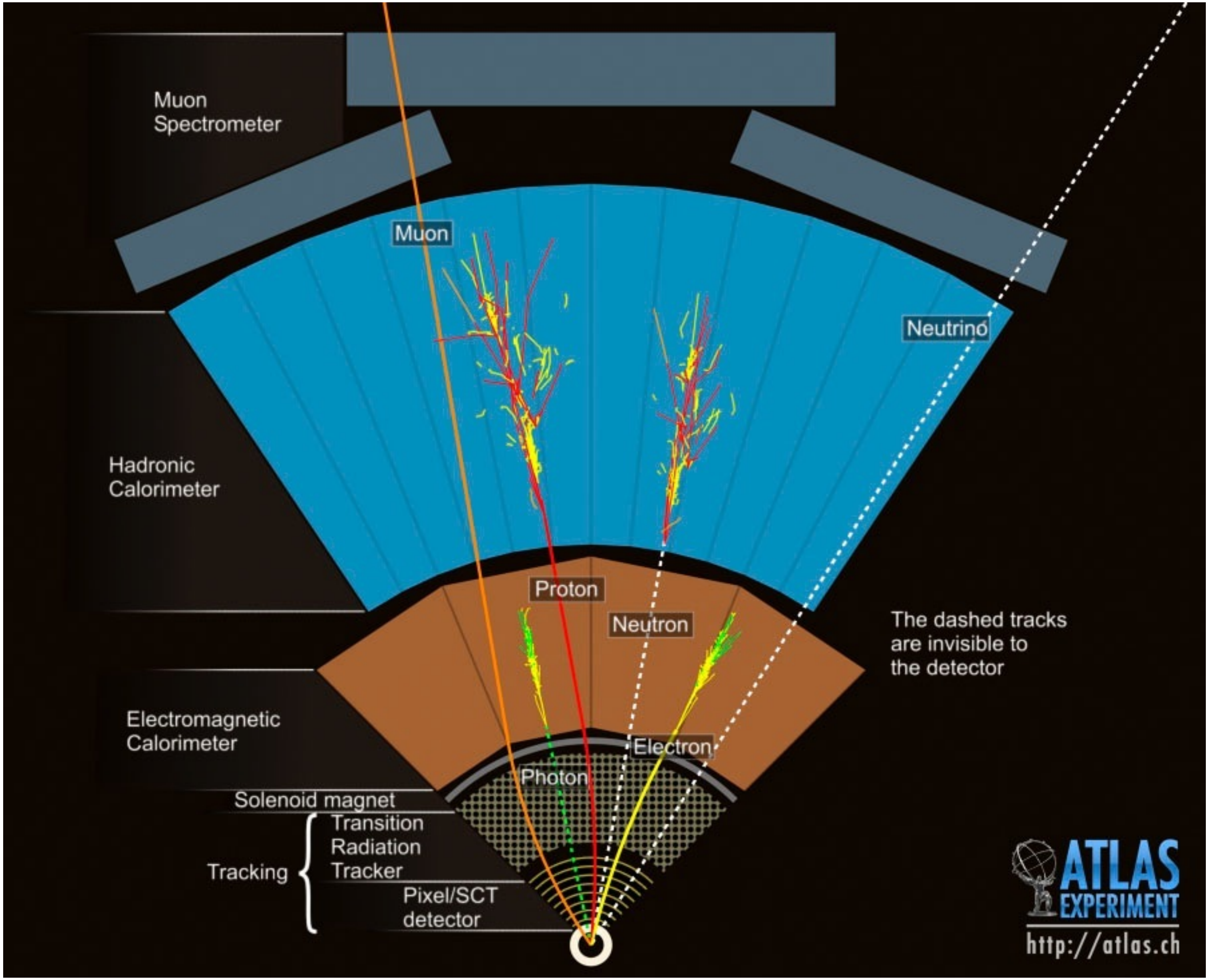
- *Short lived particles* that are not directly seen, but can be *reconstructed* from their decay products

$$\pi^0, \rho^{0,\pm}, \ldots, Z, W^{\pm}, t, H, \ldots$$

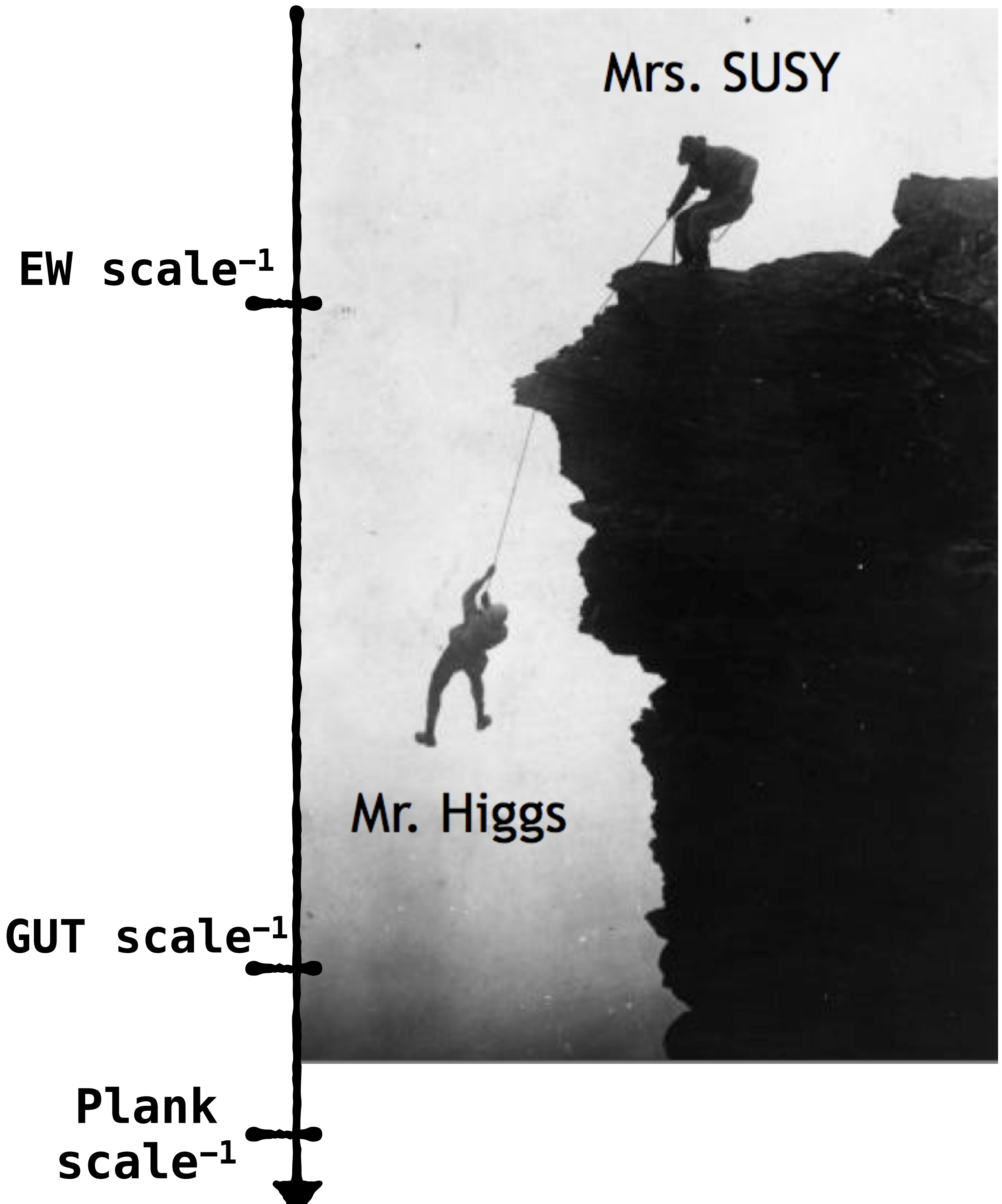
- *Missing particles*, neutral and weakly interacting, that escape the detector:

$$\nu, \tilde{\chi}^0, G_{KK}, \ldots$$

Particles Interactions



SUSY RECAP

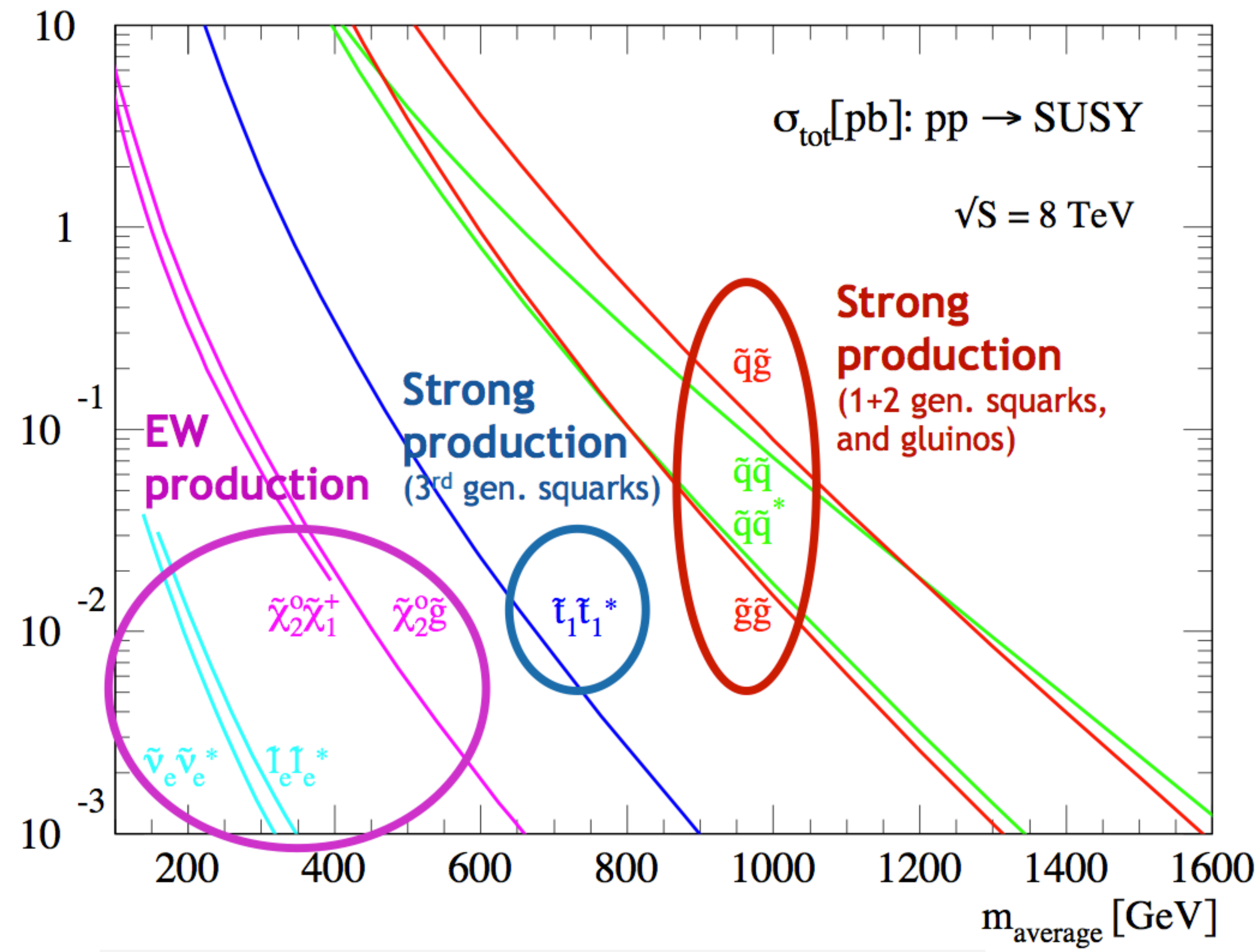


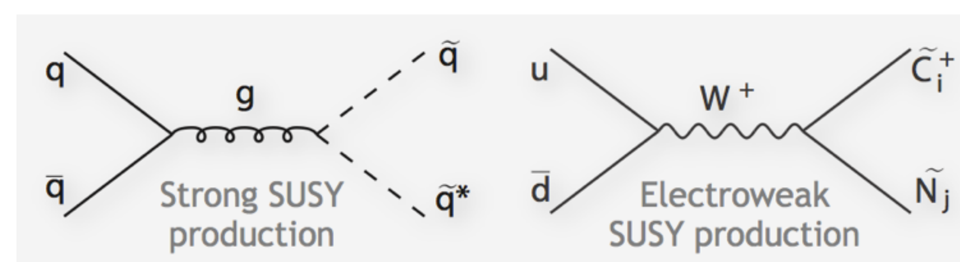
- Maximal possible extension of the space-time symmetry group
- *Moderates the hierarchy problem by cancelling the divergent quadratic corrections to the Higgs boson mass*
- Realise unification of all the known forces (but gravity) at a high scale
- *Provides a suitable dark matter candidate*

It has a rich and complex phenomenology making it interesting and useful to look for even if not realised in nature

SUSY AT COLLIDERS

- *Squarks* and *gluinos* production through strong interaction are prime candidates for discovery due to highest production cross sections
- For fine-tuning arguments *3rd generation* sparticles should have lower masses, but also have lower cross-sections
- *Charginos and Neutralinos* as well as *sneutrinos and staus* have both low mass and are produced with weak processes, very difficult to observe



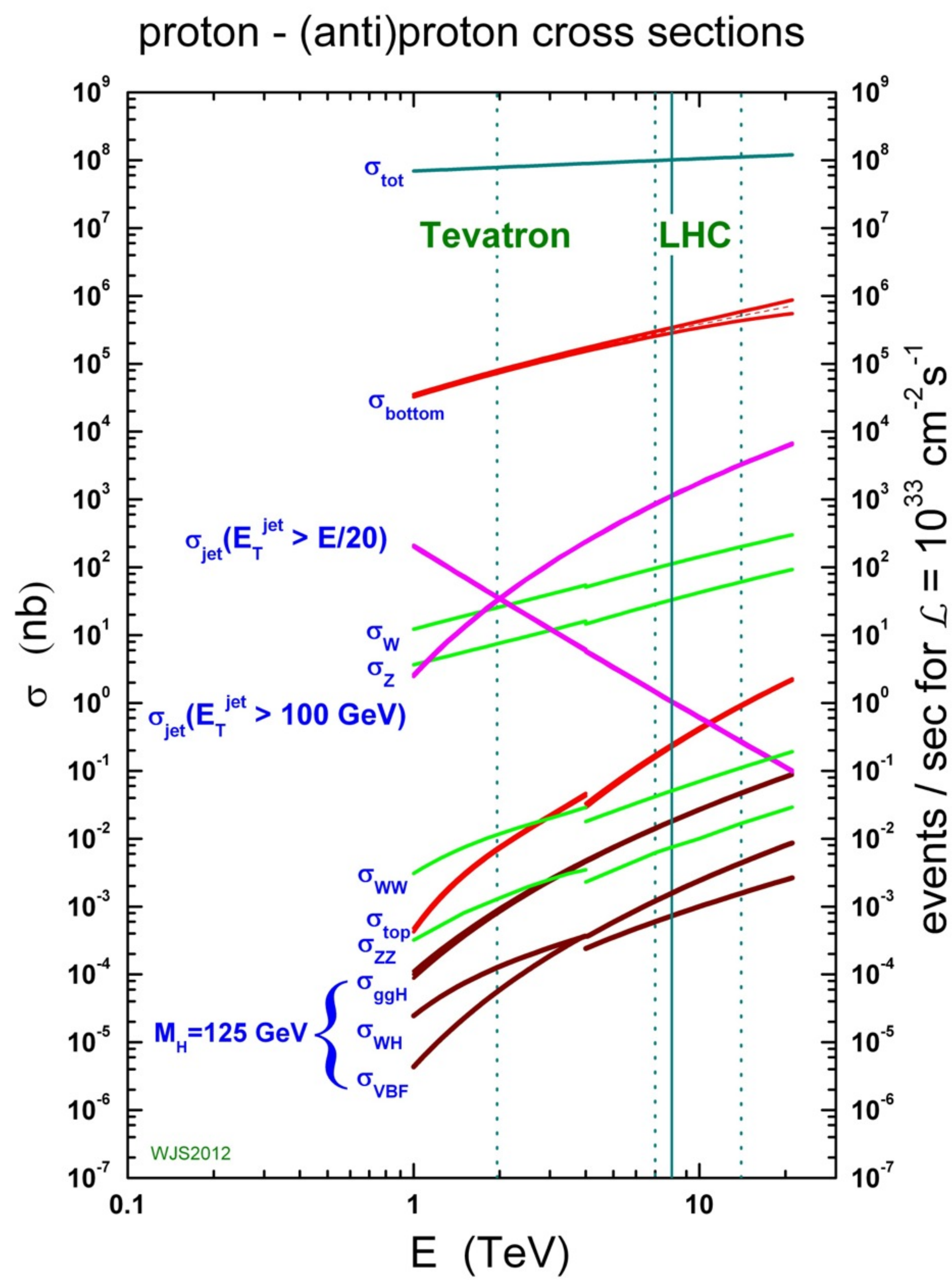


- R-parity conserving (*RPC*) signatures
  - Sparticles are produced in pairs, with each decay ending with an LSP, mostly the lightest neutralino or the gravitino → **high-E_T^{miss}**
- R-parity violating (*RPV*) signatures
  - Resonances or multi-jets/leptons, from production of single sparticles and LSP decays
  - Displaced vertices from late LSP decays

The problem: Backgrounds

The problem: Backgrounds

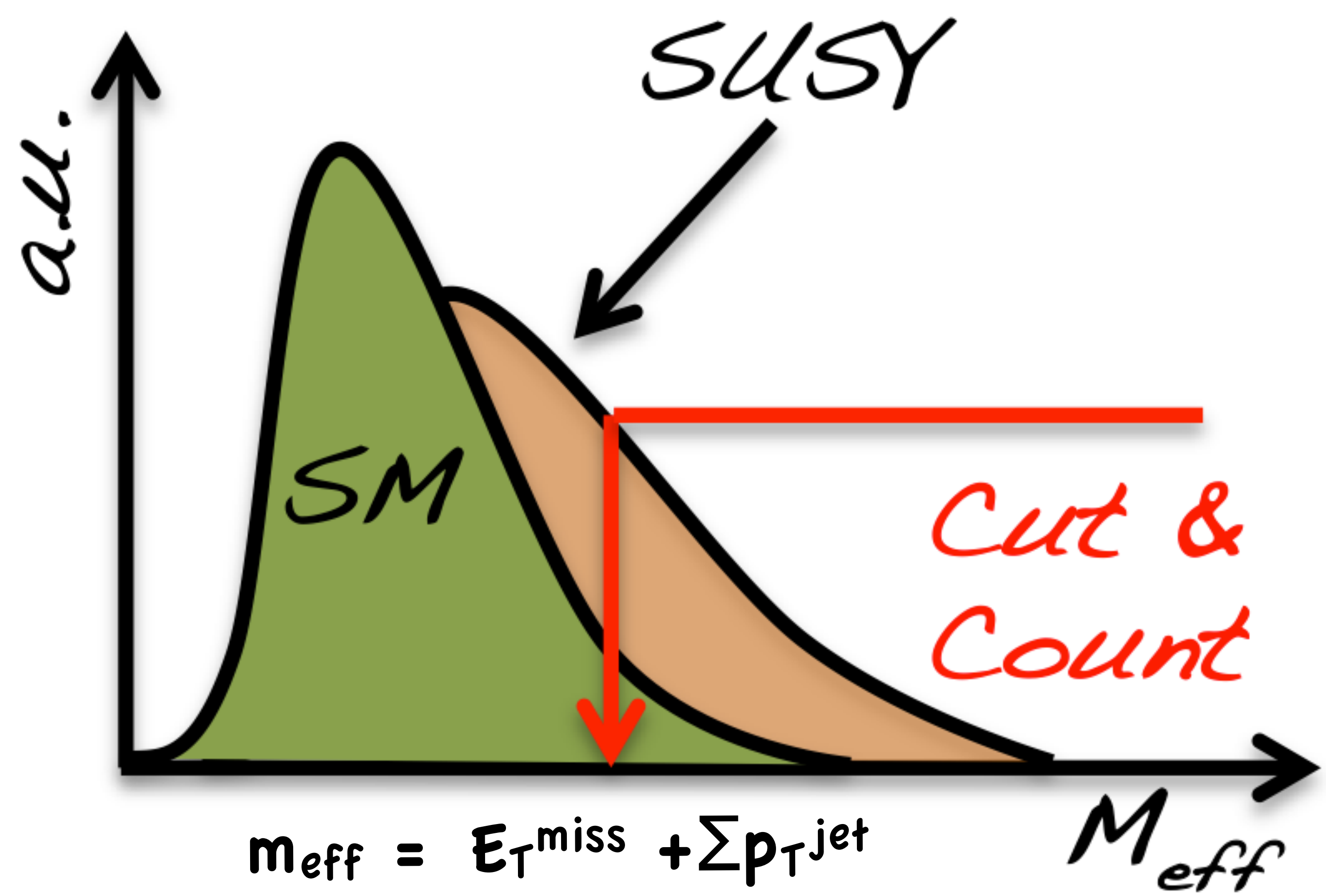
- Many other SM process can mimic the signatures of the signal we look for:
  - QCD jets
  - EWK processes (W/Z+jets)
  - Top quarks
- Their cross-sections are huge, the signal is many (>10) order of magnitudes smaller than the backgrounds, need to have an extremely good understanding of the data
- To predict how much of background will enter our search we make heavy use of simulations (*Monte Carlo*)



How to make a search

- As we cannot identify with certainty a single event to be from a SUSY reaction we must rely on *statistical analysis* to understand if some unusual events are hidden in the data
- A set of *selections* is applied to the events trying to reduce the amount of background events trying to retain as much of the signal as possible
- Given the level of precision that we want we cannot rely on the Monte Carlo simulation only: use lots of *data-driven methods*

- **Control Regions:**
  Regions dominated by a backgrounds, to validate or rescale the prediction from the simulation
- **Fake-estimates:**
  Sometimes we reconstruct a particle for another (*fakes*); since this misidentification probability can depend on a lot of details not in the simulation, it is usually measured in data



$m_{eff} = E_T^{miss} + \Sigma p_T^{jet}$

Discovery Statistics

- To assess if there is a possible signal in the data, we compute the probability for our observation if we assume there is no signal: ***p-value***
- Suppose we observe n events in the signal region, these can consist of:
  - $\mathbf{n_b}$ events from known processes (background)
  - $\mathbf{n_s}$ events from a new process (signal)
- If $\mathbf{n_s}$, $\mathbf{n_b}$ are Poisson with means s, b, then $\mathbf{n=n_s+n_b}$ is also a Poisson with mean **s+b**:

$$P(n; s, b) = \frac{(s+b)^n}{n!} e^{-(s+b)}$$

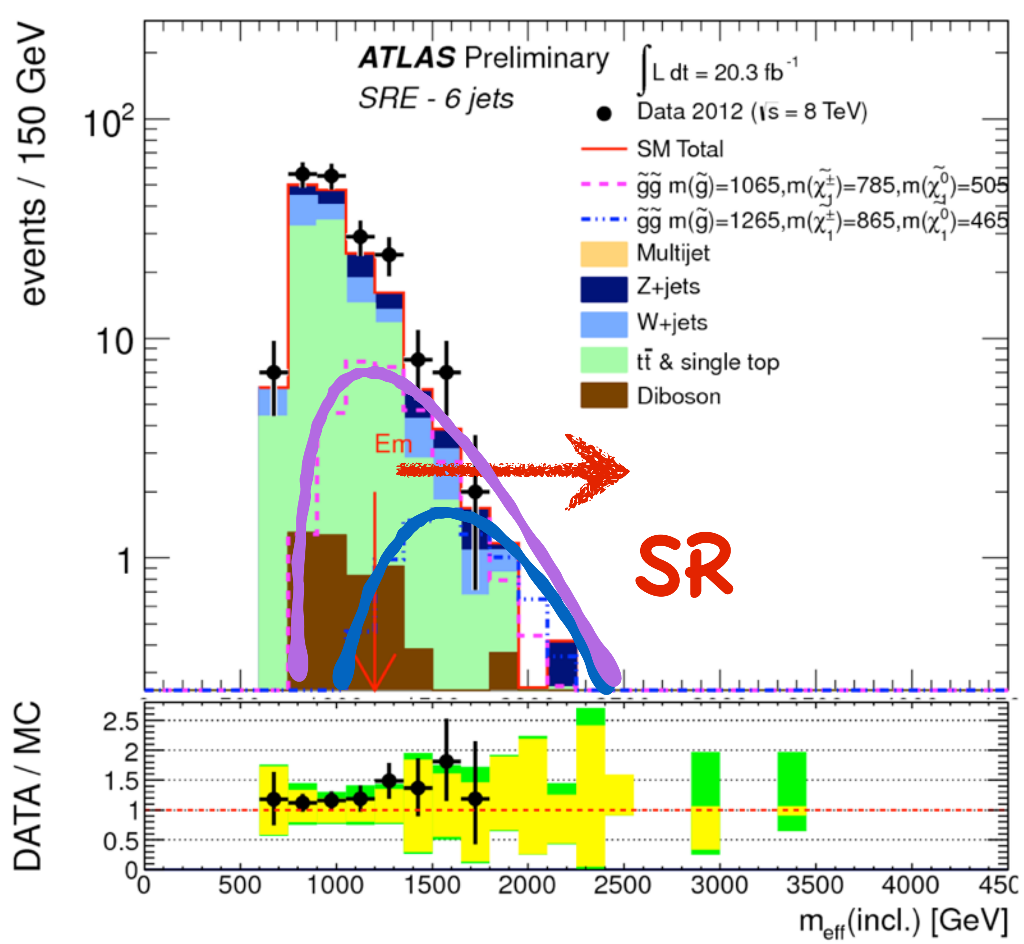
- Suppose we estimate b=0.5 and observe n=5.
Should we claim we have made a discovery?

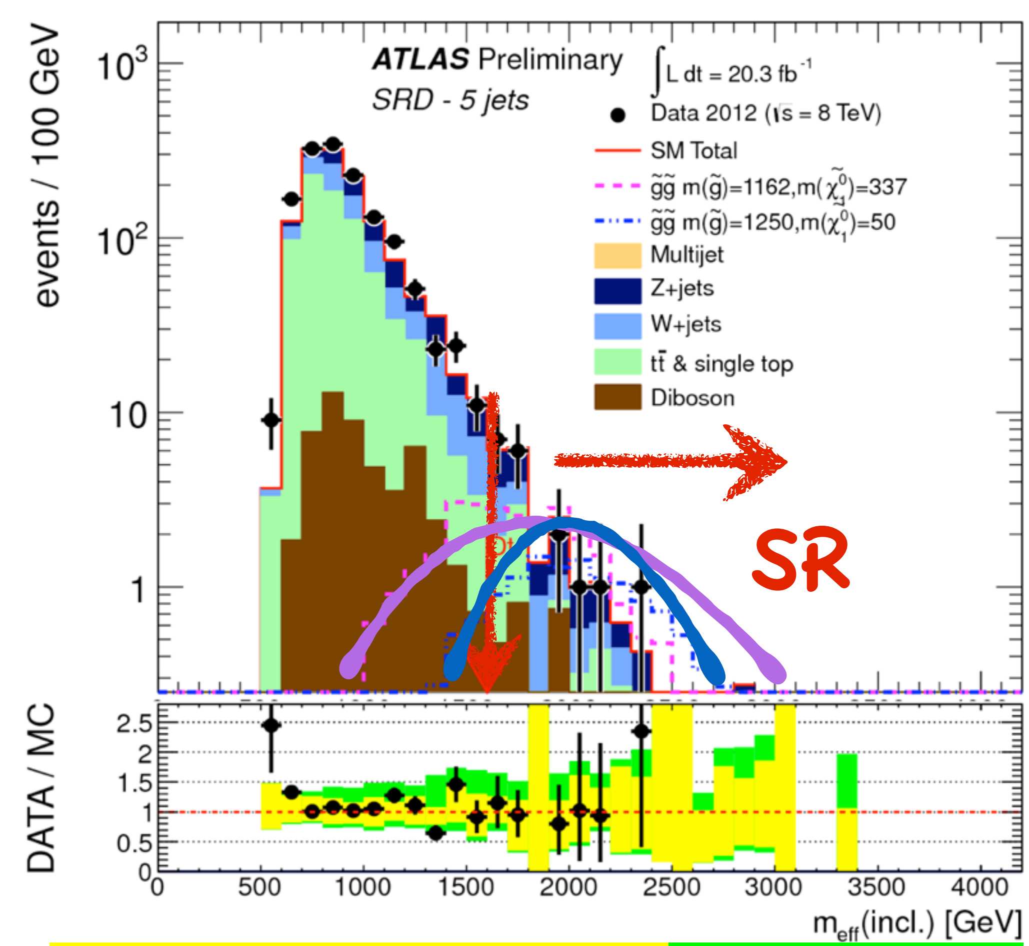
$$\text{p-value} = P(n \geq 5; b = 0.5, s = 0) = 1.7 \cdot 10^{-4}$$

- Conventionally in HEP we claim a discovery if $p = 2.9 \cdot 10^{-7}$

A search example

- SUSY inclusive search for squarks and gluiness in final states with jets and missing energy
  - The "flagship" analysis, sensitive to many scenarios
  - Select events with 2 to 6 jets and high missing energy
  - Looks at the $\mathbf{M_{eff}}$ distribution to discriminate between SUSY and the other SM backgrounds





experimental + MC stat uncertainty +theory uncertainty

Limits

- The results are compatible with the SM expectation, but we still want to extract some useful informations: which SUSY models can we exclude?

- We consider *grids* in the SUSY parameter space, and for each point we test if our observed events are compatible with it, if not the point is *excluded*

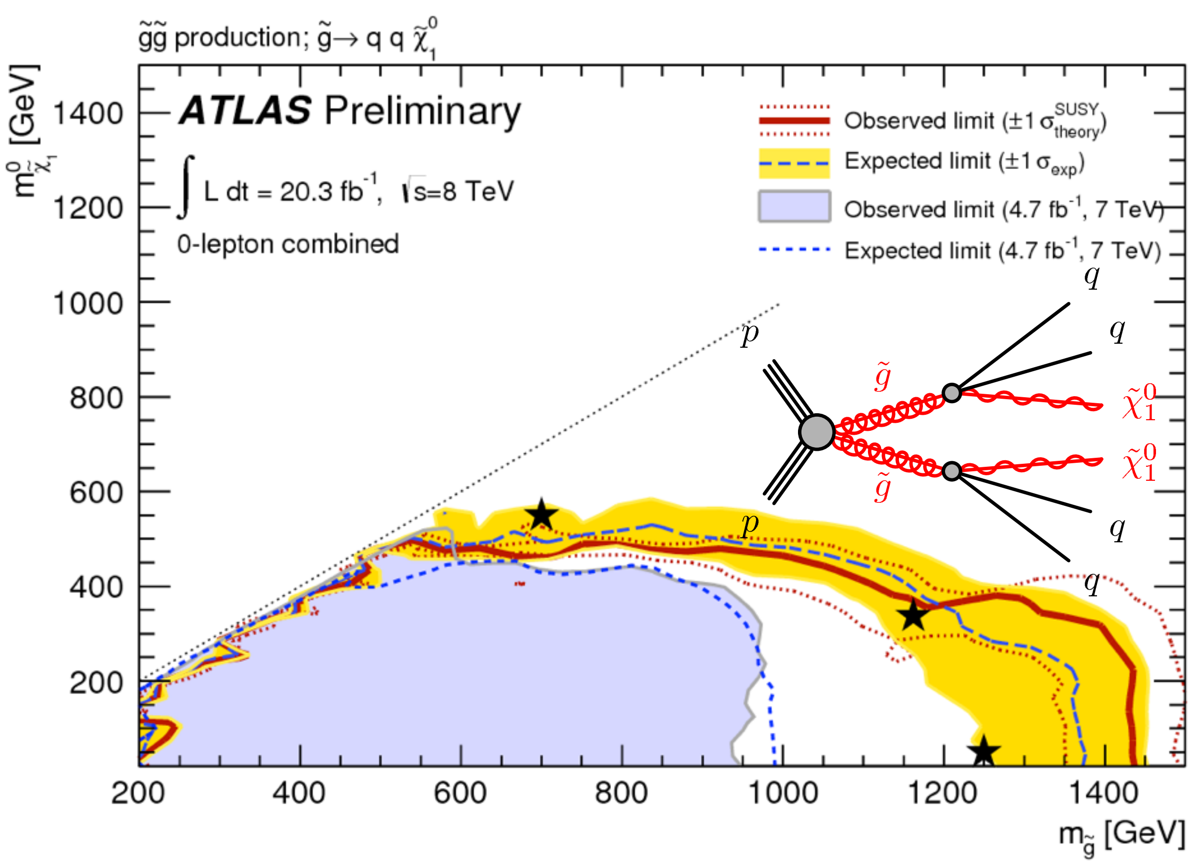
Squark/Gluino pair production, with direct decays to the LSP; all other SUSY particles decoupled

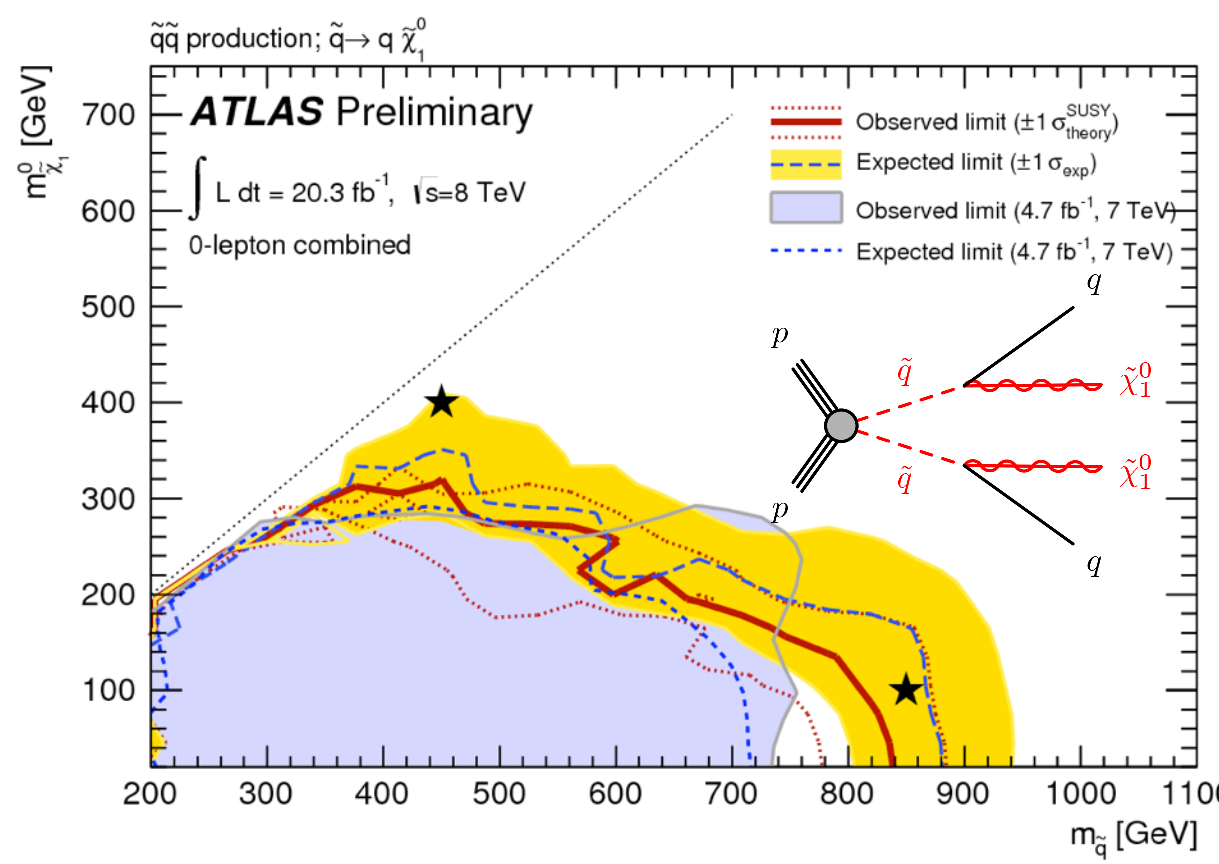
Expected limit

Yellow band ±1σ experimental uncertainties

Observed limit

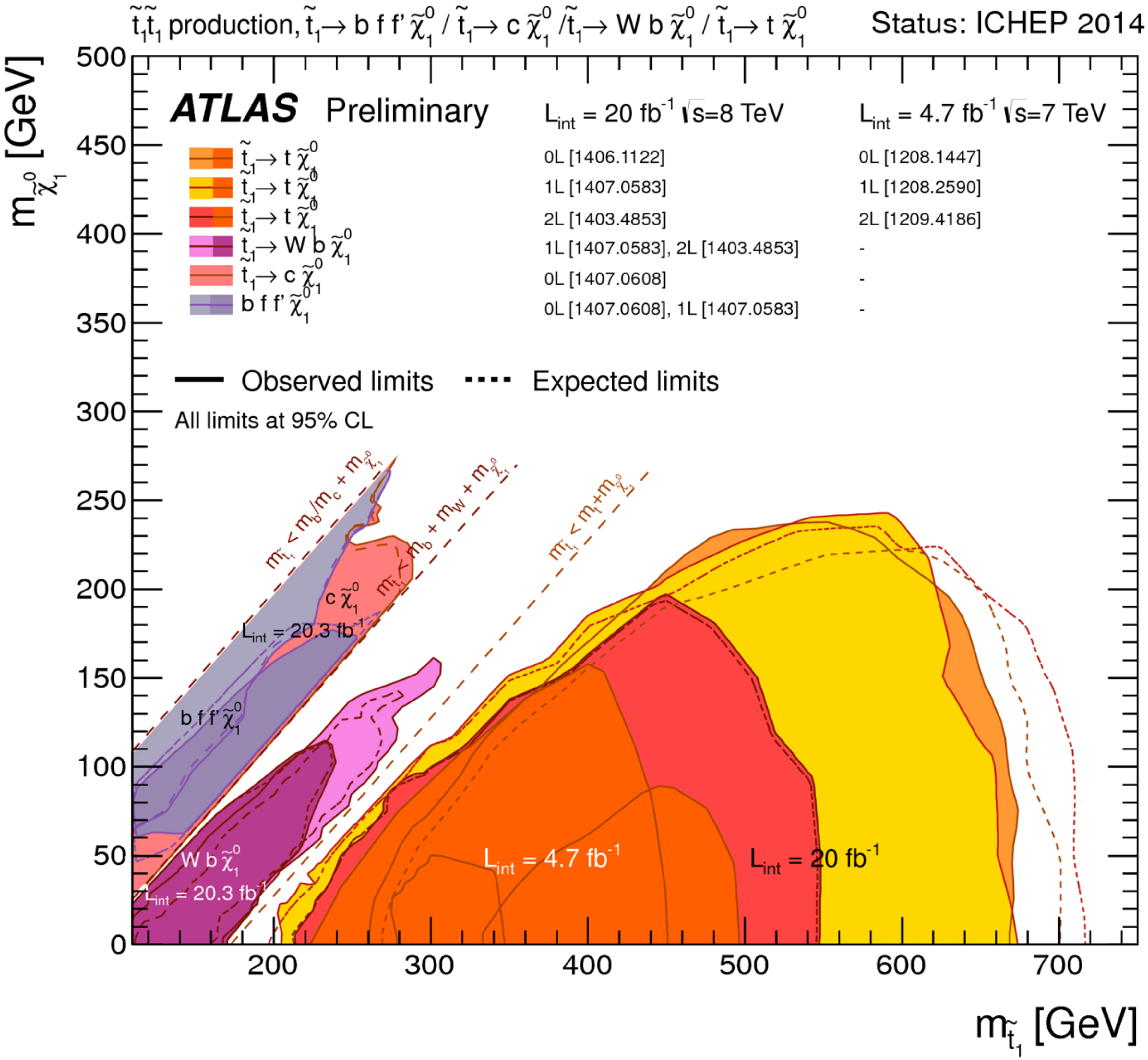
±1σ signal theory uncertainty





3rd Generation Limits

Lots of analyses have looked for signals of 3rd generation squarks, that would need to be light to solve the fine-tuning problem



We haven't found anything yet, and the limits are now close to the masses where the theory is not so appealing anymore

All ATLAS SUSY results

ATLAS SUSY Searches* - 95% CL Lower Limits

Status: ICHEP 2014

ATLAS Preliminary

$\sqrt{s}$ = 7, 8 TeV

Strong production
Natural SUSY
EWinos
RPV

	Model	e, μ, τ, γ	Jets	E_T^{miss}	$\int \mathcal{L}\, dt[fb^{-1}]$	Mass limit		Reference
Inclusive Searches	MSUGRA/CMSSM	0	2-6 jets	Yes	20.3	$\tilde{q}, \tilde{g}$ 1.7 TeV	$m(\tilde{q})=m(\tilde{g})$	1405.7875
	MSUGRA/CMSSM	1 e, μ	3-6 jets	Yes	20.3	$\tilde{g}$ 1.2 TeV	any $m(\tilde{q})$	ATLAS-CONF-2013-062
	MSUGRA/CMSSM	0	7-10 jets	Yes	20.3	$\tilde{g}$ 1.1 TeV	any $m(\tilde{q})$	1308.1841
	$\tilde{q}\tilde{q}, \tilde{q}\to q\tilde{\chi}_1^0$	0	2-6 jets	Yes	20.3	$\tilde{q}$ 850 GeV	$m(\tilde{\chi}_1^0)$=0 GeV, $m(1^{st}$ gen. $\tilde{q})=m(2^{nd}$ gen. $\tilde{q})$	1405.7875
	$\tilde{g}\tilde{g}, \tilde{g}\to q\bar{q}\tilde{\chi}_1^0$	0	2-6 jets	Yes	20.3	$\tilde{g}$ 1.[illegible]3 TeV	$m(\tilde{\chi}_1^0)$=0 GeV	1405.7875
	$\tilde{g}\tilde{g}, \tilde{g}\to qq\tilde{\chi}_1^\pm\to qqW^\pm\tilde{\chi}_1^0$	1 e, μ	3-6 jets	Yes	20.3	$\tilde{g}$ 1.18 TeV	$m(\tilde{\chi}_1^0)$<200 GeV, $m(\tilde{\chi}^\pm)=0.5(m(\tilde{\chi}_1^0)+m(\tilde{g}))$	ATLAS-CONF-2013-062
	$\tilde{g}\tilde{g}, \tilde{g}\to qq(\ell\ell/\ell\nu/\nu\nu)\tilde{\chi}_1^0$	2 e, μ	0-3 jets	-	20.3	$\tilde{g}$ 1.12 TeV	$m(\tilde{\chi}_1^0)$=0 GeV	ATLAS-CONF-2013-089
	GMSB ($\tilde{\ell}$ NLSP)	2 e, μ	2-4 jets	Yes	4.7	$\tilde{g}$ 1.24 TeV	$\tan\beta$<15	1208.4688
	GMSB ($\tilde{\ell}$ NLSP)	1-2 τ + 0-1 ℓ	0-2 jets	Yes	20.3	$\tilde{g}$ 1.6 TeV	$\tan\beta$ >20	1407.0603
	GGM (bino NLSP)	2 γ	-	Yes	20.3	$\tilde{g}$ 1.2[illegible] TeV	$m(\tilde{\chi}_1^0)$>50 GeV	ATLAS-CONF-2014-001
	GGM (wino NLSP)	1 $e, \mu + \gamma$	-	Yes	4.8	$\tilde{g}$ 619 GeV	$m(\tilde{\chi}_1^0)$>50 GeV	ATLAS-CONF-2012-144
	GGM (higgsino-bino NLSP)	γ	1 b	Yes	4.8	$\tilde{g}$ 900 GeV	$m(\tilde{\chi}_1^0)$>220 GeV	1211.1167
	GGM (higgsino NLSP)	2 e, μ (Z)	0-3 jets	Yes	5.8	$\tilde{g}$ 690 GeV	m(NLSP)>200 GeV	ATLAS-CONF-2012-152
	Gravitino LSP	0	mono-jet	Yes	10.5	$F^{1/2}$ scale 645 GeV	$m(\tilde{G})>10^{-4}$ eV	ATLAS-CONF-2012-147
3^{rd} gen. $\tilde{g}$ med.	$\tilde{g}\to b\bar{b}\tilde{\chi}_1^0$	0	3 b	Yes	20.1	$\tilde{g}$ 1.2[illegible] TeV	$m(\tilde{\chi}_1^0)$<400 GeV	1407.0600
	$\tilde{g}\to t\bar{t}\tilde{\chi}_1^0$	0	7-10 jets	Yes	20.3	$\tilde{g}$ 1.1 TeV	$m(\tilde{\chi}_1^0)$ <350 GeV	1308.1841
	$\tilde{g}\to t\bar{t}\tilde{\chi}_1^0$	0-1 e, μ	3 b	Yes	20.1	$\tilde{g}$ 1.[illegible]4 TeV	$m(\tilde{\chi}_1^0)$<400 GeV	1407.0600
	$\tilde{g}\to b\bar{t}\tilde{\chi}_1^+$	0-1 e, μ	3 b	Yes	20.1	$\tilde{g}$ 1.[illegible] TeV	$m(\tilde{\chi}_1^0)$<300 GeV	1407.0600
3^{rd} gen. squarks direct production	$\tilde{b}_1\tilde{b}_1, \tilde{b}_1\to b\tilde{\chi}_1^0$	0	2 b	Yes	20.1	$\tilde{b}_1$ 100-620 GeV	$m(\tilde{\chi}_1^0)$<90 GeV	1308.2631
	$\tilde{b}_1\tilde{b}_1, \tilde{b}_1\to t\tilde{\chi}_1^\pm$	2 e, μ (SS)	0-3 b	Yes	20.3	$\tilde{b}_1$ 275-440 GeV	$m(\tilde{\chi}_1^\pm)=2\, m(\tilde{\chi}_1^0)$	1404.2500
	$\tilde{t}_1\tilde{t}_1$(light), $\tilde{t}_1\to b\tilde{\chi}_1^\pm$	1-2 e, μ	1-2 b	Yes	4.7	$\tilde{t}_1$ 110-167 GeV	$m(\tilde{\chi}_1^0)$=55 GeV	1208.4305, 1209.2102
	$\tilde{t}_1\tilde{t}_1$(light), $\tilde{t}_1\to Wb\tilde{\chi}_1^0$	2 e, μ	0-2 jets	Yes	20.3	$\tilde{t}_1$ 130-210 GeV	$m(\tilde{\chi}_1^0)=m(\tilde{t}_1)-m(W)$-50 GeV, $m(\tilde{t}_1)<<m(\tilde{\chi}_1^\pm)$	1403.4853
	$\tilde{t}_1\tilde{t}_1$(medium), $\tilde{t}_1\to t\tilde{\chi}_1^0$	2 e, μ	2 jets	Yes	20.3	$\tilde{t}_1$ 215-530 GeV	$m(\tilde{\chi}_1^0)$=1 GeV	1403.4853
	$\tilde{t}_1\tilde{t}_1$(medium), $\tilde{t}_1\to b\tilde{\chi}_1^\pm$	0	2 b	Yes	20.1	$\tilde{t}_1$ 150-580 GeV	$m(\tilde{\chi}_1^0)$<200 GeV, $m(\tilde{\chi}_1^\pm)-m(\tilde{\chi}_1^0)$=5 GeV	1308.2631
	$\tilde{t}_1\tilde{t}_1$(heavy), $\tilde{t}_1\to t\tilde{\chi}_1^0$	1 e, μ	1 b	Yes	20	$\tilde{t}_1$ 210-640 GeV	$m(\tilde{\chi}_1^0)$=0 GeV	1407.0583
	$\tilde{t}_1\tilde{t}_1$(heavy), $\tilde{t}_1\to t\tilde{\chi}_1^0$	0	2 b	Yes	20.1	$\tilde{t}_1$ 260-640 GeV	$m(\tilde{\chi}_1^0)$=0 GeV	1406.1122
	$\tilde{t}_1\tilde{t}_1, \tilde{t}_1\to c\tilde{\chi}_1^0$	0	mono-jet/c-tag	Yes	20.3	$\tilde{t}_1$ 90-240 GeV	$m(\tilde{t}_1)-m(\tilde{\chi}_1^0)$<85 GeV	1407.0608
	$\tilde{t}_1\tilde{t}_1$(natural GMSB)	2 e, μ (Z)	1 b	Yes	20.3	$\tilde{t}_1$ 150-580 GeV	$m(\tilde{\chi}_1^0)$>150 GeV	1403.5222
	$\tilde{t}_2\tilde{t}_2, \tilde{t}_2\to\tilde{t}_1+Z$	3 e, μ (Z)	1 b	Yes	20.3	$\tilde{t}_2$ 290-600 GeV	$m(\tilde{\chi}_1^0)$<200 GeV	1403.5222
EW direct	$\tilde{\ell}_{L,R}\tilde{\ell}_{L,R}, \tilde{\ell}\to\ell\tilde{\chi}_1^0$	2 e, μ	0	Yes	20.3	$\tilde{\ell}$ 90-325 GeV	$m(\tilde{\chi}_1^0)$=0 GeV	1403.5294
	$\tilde{\chi}_1^+\tilde{\chi}_1^-, \tilde{\chi}_1^+\to\tilde{\ell}\nu(\ell\tilde{\nu})$	2 e, μ	0	Yes	20.3	$\tilde{\chi}_1^\pm$ 140-465 GeV	$m(\tilde{\chi}_1^0)$=0 GeV, $m(\tilde{\ell},\tilde{\nu})=0.5(m(\tilde{\chi}_1^\pm)+m(\tilde{\chi}_1^0))$	1403.5294
	$\tilde{\chi}_1^+\tilde{\chi}_1^-, \tilde{\chi}_1^+\to\tilde{\tau}\nu(\tau\tilde{\nu})$	2 τ	-	Yes	20.3	$\tilde{\chi}_1^\pm$ 100-350 GeV	$m(\tilde{\chi}_1^0)$=0 GeV, $m(\tilde{\tau},\tilde{\nu})=0.5(m(\tilde{\chi}_1^\pm)+m(\tilde{\chi}_1^0))$	1407.0350
	$\tilde{\chi}_1^\pm\tilde{\chi}_2^0\to\tilde{\ell}_L\nu\tilde{\ell}_L\ell(\tilde{\nu}\nu), \ell\tilde{\nu}\tilde{\ell}_L\ell(\tilde{\nu}\nu)$	3 e, μ	0	Yes	20.3	$\tilde{\chi}_1^\pm, \tilde{\chi}_2^0$ 700 GeV	$m(\tilde{\chi}_1^\pm)=m(\tilde{\chi}_2^0)$, $m(\tilde{\chi}_1^0)$=0, $m(\tilde{\ell},\tilde{\nu})=0.5(m(\tilde{\chi}_1^\pm)+m(\tilde{\chi}_1^0))$	1402.7029
	$\tilde{\chi}_1^\pm\tilde{\chi}_2^0\to W\tilde{\chi}_1^0Z\tilde{\chi}_1^0$	2-3 e, μ	0	Yes	20.3	$\tilde{\chi}_1^\pm, \tilde{\chi}_2^0$ 420 GeV	$m(\tilde{\chi}_1^\pm)=m(\tilde{\chi}_2^0)$, $m(\tilde{\chi}_1^0)$=0, sleptons decoupled	1403.5294, 1402.7029
	$\tilde{\chi}_1^\pm\tilde{\chi}_2^0\to W\tilde{\chi}_1^0h\tilde{\chi}_1^0$	1 e, μ	2 b	Yes	20.3	$\tilde{\chi}_1^\pm, \tilde{\chi}_2^0$ 285 GeV	$m(\tilde{\chi}_1^\pm)=m(\tilde{\chi}_2^0)$, $m(\tilde{\chi}_1^0)$=0, sleptons decoupled	ATLAS-CONF-2013-093
	$\tilde{\chi}_2^0\tilde{\chi}_3^0, \tilde{\chi}_{2,3}^0\to\tilde{\ell}_R\ell$	4 e, μ	0	Yes	20.3	$\tilde{\chi}_{2,3}^0$ 620 GeV	$m(\tilde{\chi}_2^0)=m(\tilde{\chi}_3^0)$, $m(\tilde{\chi}_1^0)$=0, $m(\tilde{\ell},\tilde{\nu})=0.5(m(\tilde{\chi}_2^0)+m(\tilde{\chi}_1^0))$	1405.5086
Long-lived particles	Direct $\tilde{\chi}_1^+\tilde{\chi}_1^-$ prod., long-lived $\tilde{\chi}_1^\pm$	Disapp. trk	1 jet	Yes	20.3	$\tilde{\chi}_1^\pm$ 270 GeV	$m(\tilde{\chi}_1^\pm)-m(\tilde{\chi}_1^0)$=160 MeV, $\tau(\tilde{\chi}_1^\pm)$=0.2 ns	ATLAS-CONF-2013-069
	Stable, stopped $\tilde{g}$ R-hadron	0	1-5 jets	Yes	27.9	$\tilde{g}$ 832 GeV	$m(\tilde{\chi}_1^0)$=100 GeV, 10 μs<$\tau(\tilde{g})$<1000 s	1310.6584
	GMSB, stable $\tilde{\tau}$, $\tilde{\chi}_1^0\to\tilde{\tau}(\tilde{e},\tilde{\mu})+\tau(e,\mu)$	1-2 μ	-	-	15.9	$\tilde{\chi}_1^0$ 475 GeV	10<$\tan\beta$<50	ATLAS-CONF-2013-058
	GMSB, $\tilde{\chi}_1^0\to\gamma\tilde{G}$, long-lived $\tilde{\chi}_1^0$	2 γ	-	Yes	4.7	$\tilde{\chi}_1^0$ 230 GeV	0.4<$\tau(\tilde{\chi}_1^0)$<2 ns	1304.6310
	$\tilde{q}\tilde{q}, \tilde{\chi}_1^0\to qq\mu$ (RPV)	1 μ, displ. vtx	-	-	20.3	$\tilde{q}$ 1.0 TeV	1.5 <$c\tau$<156 mm, BR(μ)=1, $m(\tilde{\chi}_1^0)$=108 GeV	ATLAS-CONF-2013-092
RPV	LFV $pp\to\tilde{\nu}_\tau+X, \tilde{\nu}_\tau\to e+\mu$	2 e, μ	-	-	4.6	$\tilde{\nu}_\tau$ 1.61 TeV	λ'_{311}=0.10, λ_{132}=0.05	1212.1272
	LFV $pp\to\tilde{\nu}_\tau+X, \tilde{\nu}_\tau\to e(\mu)+\tau$	1 $e, \mu + \tau$	-	-	4.6	$\tilde{\nu}_\tau$ 1.1 TeV	λ'_{311}=0.10, $\lambda_{1(2)33}$=0.05	1212.1272
	Bilinear RPV CMSSM	2 e, μ (SS)	0-3 b	Yes	20.3	$\tilde{q}, \tilde{g}$ 1.[illegible]5 TeV	$m(\tilde{q})=m(\tilde{g})$, $c\tau_{LSP}$<1 mm	1404.2500
	$\tilde{\chi}_1^+\tilde{\chi}_1^-, \tilde{\chi}_1^+\to W\tilde{\chi}_1^0, \tilde{\chi}_1^0\to ee\tilde{\nu}_\mu, e\mu\tilde{\nu}_e$	4 e, μ	-	Yes	20.3	$\tilde{\chi}_1^\pm$ 750 GeV	$m(\tilde{\chi}_1^0)>0.2\times m(\tilde{\chi}_1^\pm)$, $\lambda_{121}\neq0$	1405.5086
	$\tilde{\chi}_1^+\tilde{\chi}_1^-, \tilde{\chi}_1^+\to W\tilde{\chi}_1^0, \tilde{\chi}_1^0\to\tau\tau\tilde{\nu}_e, e\tau\tilde{\nu}_\tau$	3 $e, \mu + \tau$	-	Yes	20.3	$\tilde{\chi}_1^\pm$ 450 GeV	$m(\tilde{\chi}_1^0)>0.2\times m(\tilde{\chi}_1^\pm)$, $\lambda_{133}\neq0$	1405.5086
	$\tilde{g}\to qqq$	0	6-7 jets	-	20.3	$\tilde{g}$ 916 GeV	BR(t)=BR(b)=BR(c)=0%	ATLAS-CONF-2013-091
	$\tilde{g}\to\tilde{t}_1t, \tilde{t}_1\to bs$	2 e, μ (SS)	0-3 b	Yes	20.3	$\tilde{g}$ 850 GeV		1404.250
Other	Scalar gluon pair, sgluon$\to q\bar{q}$	0	4 jets	-	4.6	sgluon 100-287 GeV	incl. limit from 1110.2693	1210.4826
	Scalar gluon pair, sgluon$\to t\bar{t}$	2 e, μ (SS)	2 b	Yes	14.3	sgluon 350-800 GeV		ATLAS-CONF-2013-051
	WIMP interaction (D5, Dirac χ)	0	mono-jet	Yes	10.5	M* scale 704 GeV	$m(\chi)$<80 GeV, limit of<687 GeV for D8	ATLAS-CONF-2012-147

$\sqrt{s}$ = 7 TeV full data | $\sqrt{s}$ = 8 TeV partial data | $\sqrt{s}$ = 8 TeV full data

10^{-1} — Mass scale [TeV]

1TeV

*Only a selection of the available mass limits on new states or phenomena is shown. All limits quoted are observed minus 1σ theoretical signal cross section uncertainty.

Summary

- ATLAS has performed a broad search for SUSY
- Inclusive analysis have been extended to cover challenging scenarios

But no excess observed so far!

- Mass limits pushed further in the TeV region where the theory suggested to look
  - Gluino masses up to 1.3 TeV
  - Squark masses up to 740 GeV
  - Stop above 700 GeV

- The upcoming Run2 of the LHC at an even higher energy of 13 TeV will bring new exciting opportunities to discover (or finally exclude) Supersymmetry

Backup