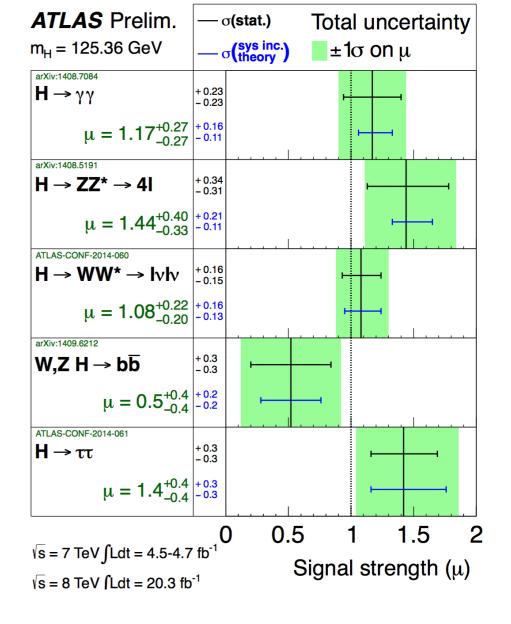


What we're looking for...

- SM H → bb, in Higgs-strahlung
 - One of the most elusive Higgs decay channel at LHC
 - But the one the Higgs boson (supposedly) most frequently decays into
- $H/X \rightarrow HH \rightarrow 4 \text{ b-jets}$
 - In the future hope to access Higgs self-coupling
 - Now: look for BSM heavy X resonances to HH
- (b)bH, H \rightarrow bb in (x)MSSM or 2HDM models
 - First analysis in ATLAS



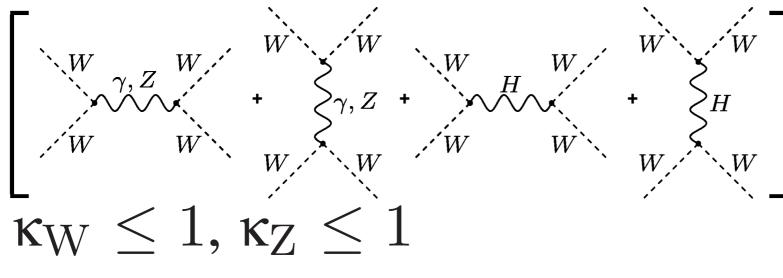
- Quite a few searches involving b-jets reconstruction, where we have >4 people very active also on the performance side (not covered in this talk).
- SM H \rightarrow bb search now published, remaining ones still in ATLAS review.

One motivation for H to bb

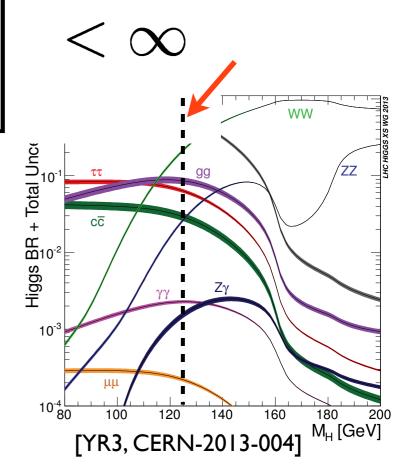
At the LHC the Higgs Higgs couplings can't be measured directly

$$\sigma_{YY \to H} \dot{BR}(H \to XX) \approx \Gamma_Y \frac{\Gamma_X}{\Gamma_H}$$

- A measurement of absolute couplings is possible if the total width is bound
 - Upper limit from fulfilling unitarity in WW scattering (valid for SM and a large class of BSM models)



- Lower limit from sum of all "visible" decay modes $\Gamma_H \ge \Gamma_W + \Gamma_Z + \Gamma_a + \Gamma_\tau + \Gamma_b$
- At ~125 GeV Higgs boson width is expected to be dominated by H to bb (BR ~ 60%)
 - Precise determination of H to bb important for extracting absolute couplings!



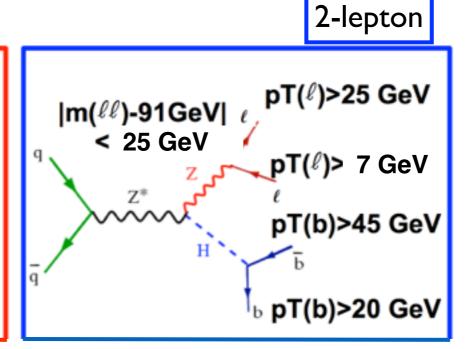
The VH analysis

- Three leptonic signatures:
 - Missing ET
- $ZH \rightarrow \nu \bar{\nu} b \bar{b}$
- I-lepton + Missing ET
- $WH \to \ell \nu b\bar{b}$

- 2-leptons
- $ZH \rightarrow \ell^+\ell^-b\bar{b}$
- Main analysis selection criteria:

- Jets reconstructed with an AntiKt4 jet algorithm.
- Leading two jets in pT used to form Higgs candidate.

PT(ℓ)>25 GeV
pT(b)>45 GeV
pT(b)>20 GeV

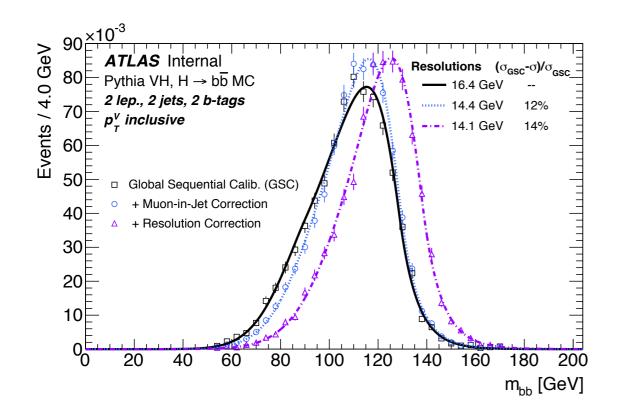


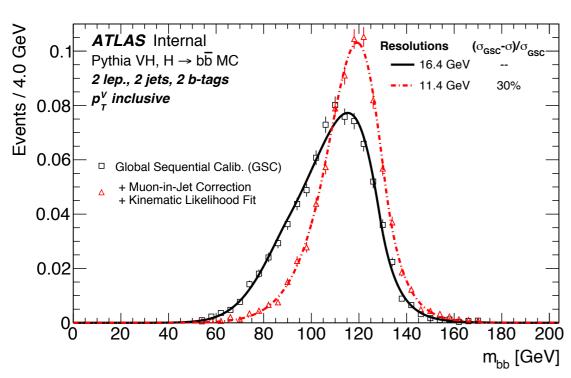
Missing ET >100 GeV

The point of the point

Categories with no or one additional jets.

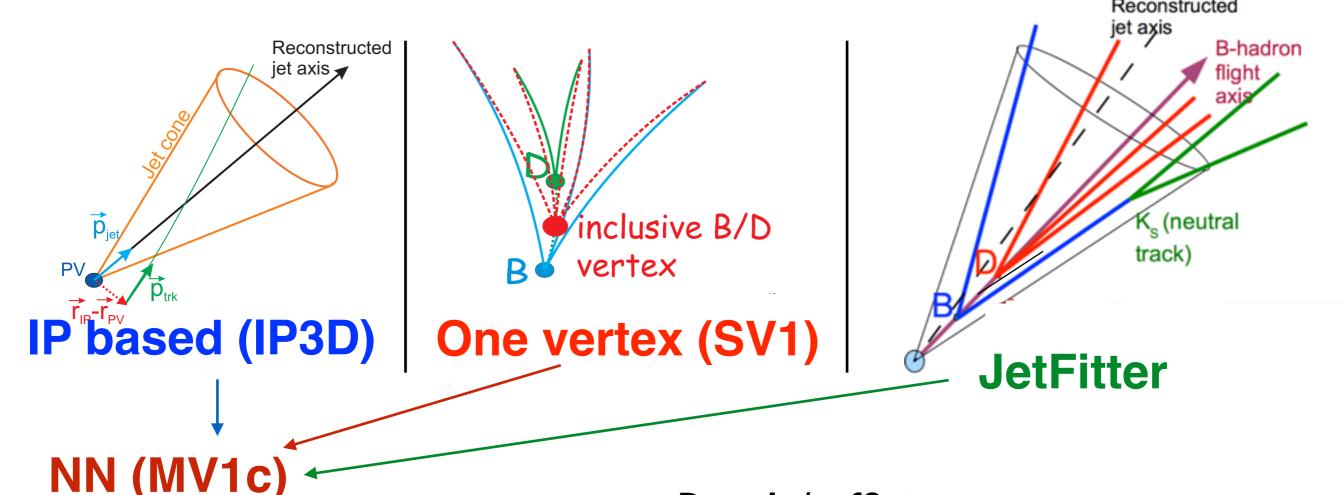
Ingredients: (1) mass





- Std antiKt R=0.4 jets
- Use Global Sequential
 Calibration as basis (already ~5% improvement w.r.t.
 EM jets)
- Then add muon-in-jet
- For 2-lepton channel, run full kinematic fit
 - Exploit closed kinematics of Ilbb system (maximize likelihood)

Ingredients: (2) b-tagging

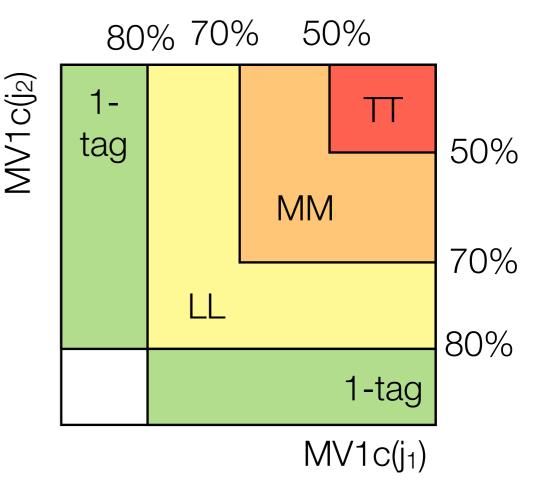


ε (B)	R(c)	R(light)
80%	~3	~29
70%	~5.3	~136
60%	~10.5	~450
50%	~26	~1400

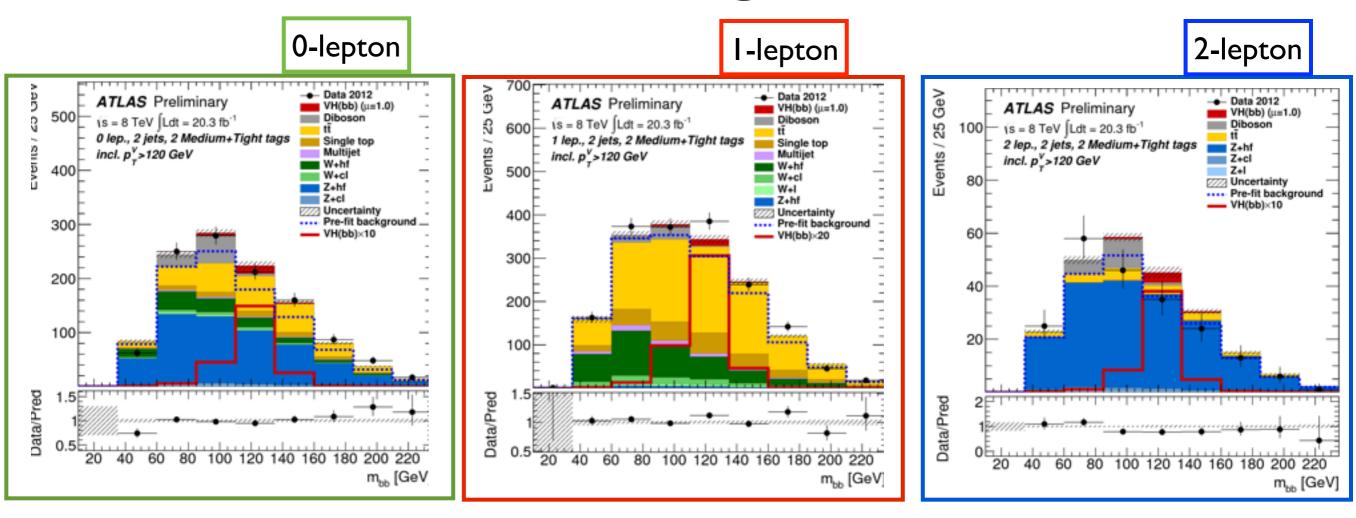
- R = I / efficiency
- Optimized to reject c-jets
- Simultaneous use of several working points.
- Dedicated calibration with reduced uncertainties (~2-3%).

Analysis categories

- Three channels: 0-lepton, 1-lepton and 2-lepton
- **Two** pT(W/Z) regions
 - <120, >120 GeV
- Four b-tag regions (1-tag, LL,MM,TT)
- Two jet bins (2 and 3 jets)
- Discriminating variables in fit
 - 1-tag: **MV1c**
 - 2-tag: **BDT**



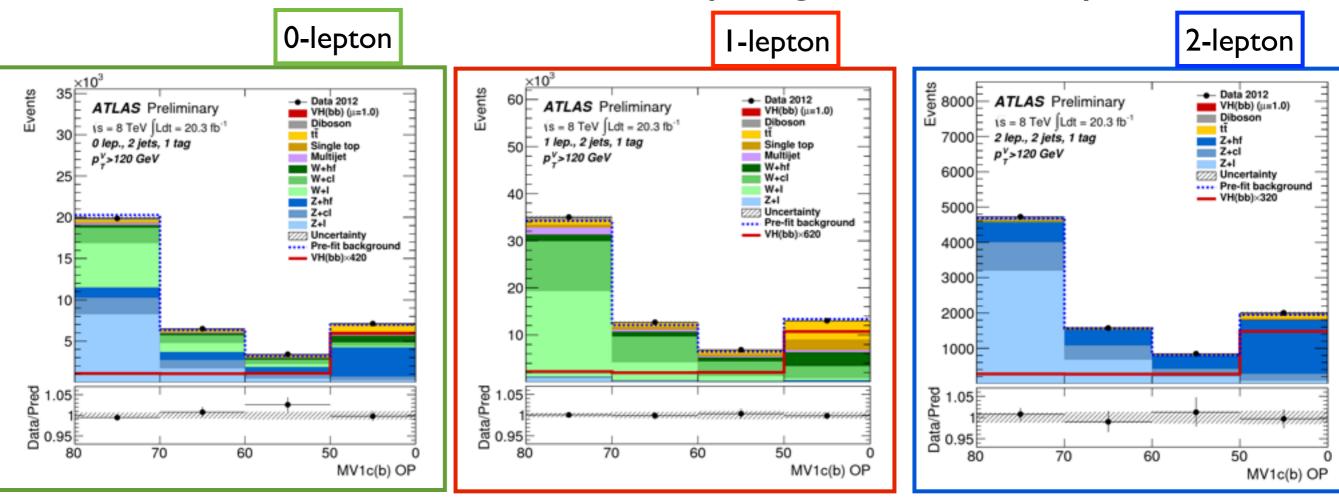
Main backgrounds



- W/Z+jets (mainly +bb/cc/cl)
- Top background (especially I- and 0-lepton)
- Diboson (W/Z+Z to bb)
- QCD (data-driven estimates)

W+jet, Z+jet backgrounds

- High statistics control regions are the 0- and 1-tag regions
- Residual contamination from top bkg in 0- and 1-lepton

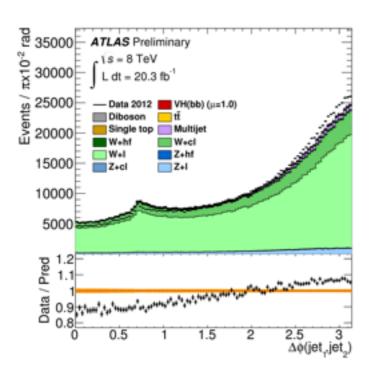


• The MVIc distribution in the I-tag regions is used to obtain the fractions of W+ and Z+cl, cc/cb/bb from data.

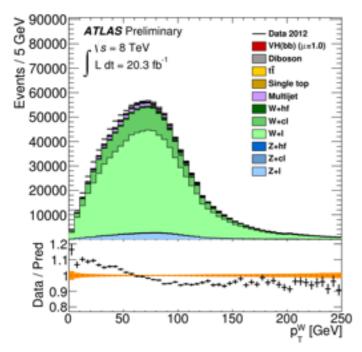
Process	Scale factor
\overline{Wbb}	0.83 ± 0.15
Wcl	1.14 ± 0.10
Zbb	1.09 ± 0.05
Zcl	0.88 ± 0.12

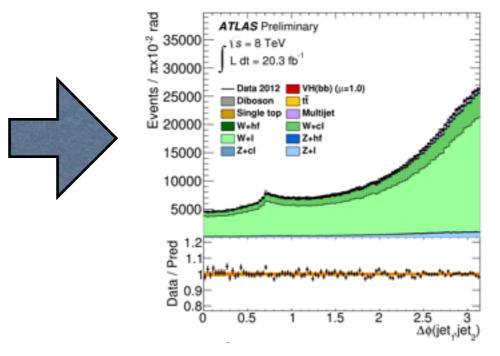
Data driven corrections

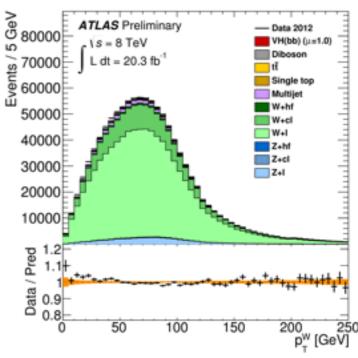
 In both 0- and I-tag regions, a significant discrepancy between data and MC is seen in both the DeltaPhi(j,j) and pT(V) distributions.



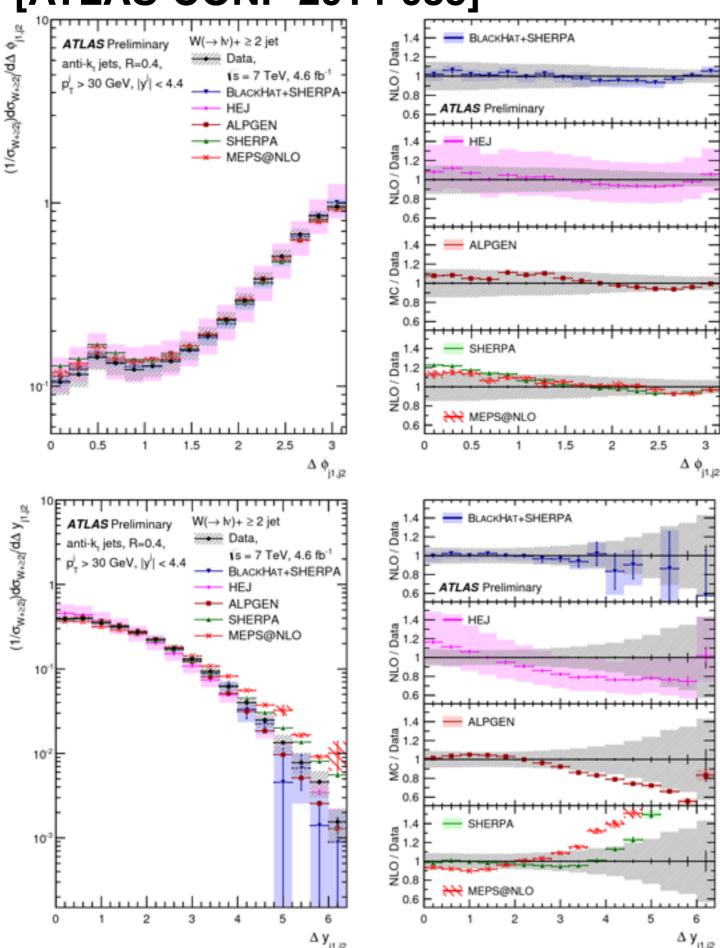
- A data/MC correction is derived using 0-tag events based on DeltaPhi, and applied to MC. The low and high pT(V) bins are separately corrected.
- The correction works well also in I-tag events, which means W+cl needs the same correction. But hard to conclude on W+cc/bb: no correction applied there.
- Half of the correction as systematic uncertainty. The full W+light/cl correction as systematics on W+bb/cc.







[ATLAS-CONF-2014-035]

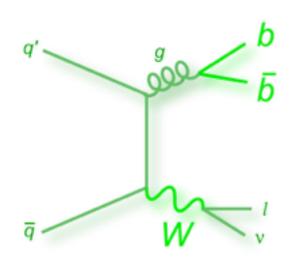


SM measurement

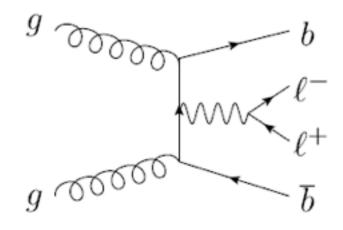
- Despite the different cuts, the same effect is seen in the unfolded W+jet measurement, for N(jets)≥2.
- NLO for W+2 jets (BlackHat +Sherpa) gets closest to data.
- For Run-II, hopefully we can establish a good modelling for W+jets first, before the fun with heavy flavor starts!
 - First indications based on Sherpa 2.1.x are not so promising.
 - Some NLO codes give slightly improved modelling, is the matching to the parton shower the problem?

W+bb/Z+bb processes

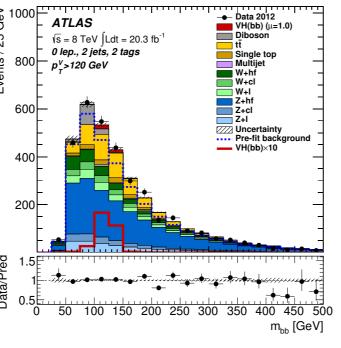
• For heavy flavor processes additional systematic uncertainties are considered.

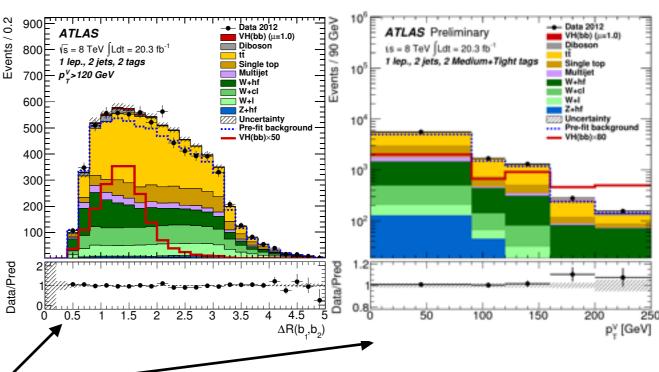


- W+bb is hard to control with data (top/single-top backgrounds).
 Different generators / models are considered:
 - I. Powheg+Pythia/Herwig
 - 2. aMC@NLO
 - 3. Sherpa
- Variations of O(20%) on m(bb), Δ R(bb)! (gluon splitting? NLO corrections?)
- Really hoping in theory improvements in this area...
 - Sherpa 2? Powheg MiNLO Wbb? (needs W+bbj @ NLO)?
 aMC@NLO (needs extension of FxFx to b-jets)?
- Z+bb can be controlled a bit easier from m(bb) sidebands, but signal region lacks statistics
 - Systematic variation on m(bb) derived in 0- and I-tag region, and cross-checked in 2-tag region.



BDT





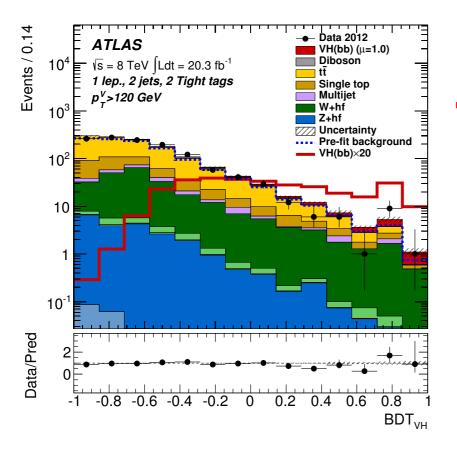
	0-Lepton	1-Lepton	2-Lepton
$\overline{p_{ m T}^V}$		×	×
$E_{ m T}^{ m miss}$	×	×	×
$p_{ m T}^{b_1}$	×	×	×
$p_{ m T}^{b_2}$	×	×	×
m_{bb}	×	×	× /
$\Delta R(b_1,b_2)$	×	×	×
$ \Delta\eta(b_1,b_2) $	×		×
$\Delta\phi(V,bb)$	×	×	×
$ \Delta\eta(V,bb) $			×
$H_{ m T}$	×		
$\min[\Delta\phi(\ell,b)]$		×	
$m_{ m T}^W$		×	
$m_{\ell\ell}$			×
$MV1c(b_1)$	×	×	×
$MV1c(b_2)$	×	×	×
	Only in 3-jet events		
$p_{ m T}^{ m jet_3}$	×	×	×
m_{11} .		¥	~

- Good description of input variables & correlations.
- Main improvements from EPS 2013:
 - MVA based on BDT (+15%)
 - "Continuous" b-tagging (+15%)
 - GSC for jets / kinematic fit in 2-lepton
 - Re-optimization of cuts

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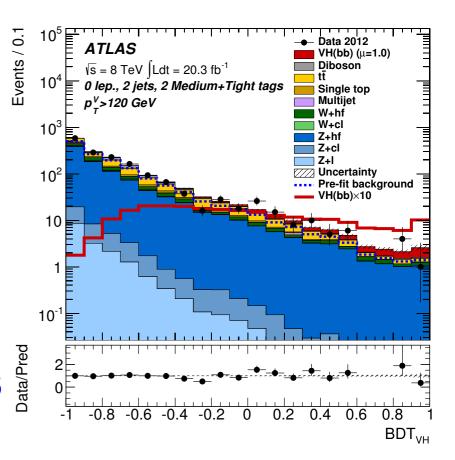
- >50% improvement in sensitivity
 - SLAC goal was to improve H to bb by >20%. Exceeded!!

Fit model



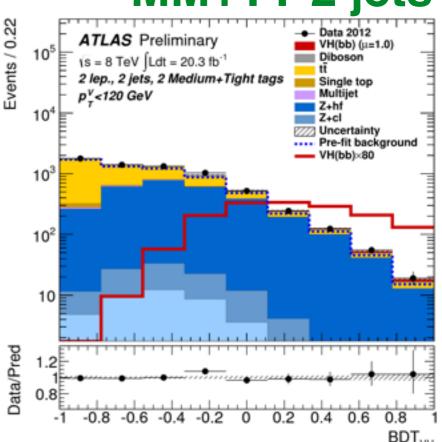
1-lepton TT 2 jets

0-lepton MM+TT 2 jets

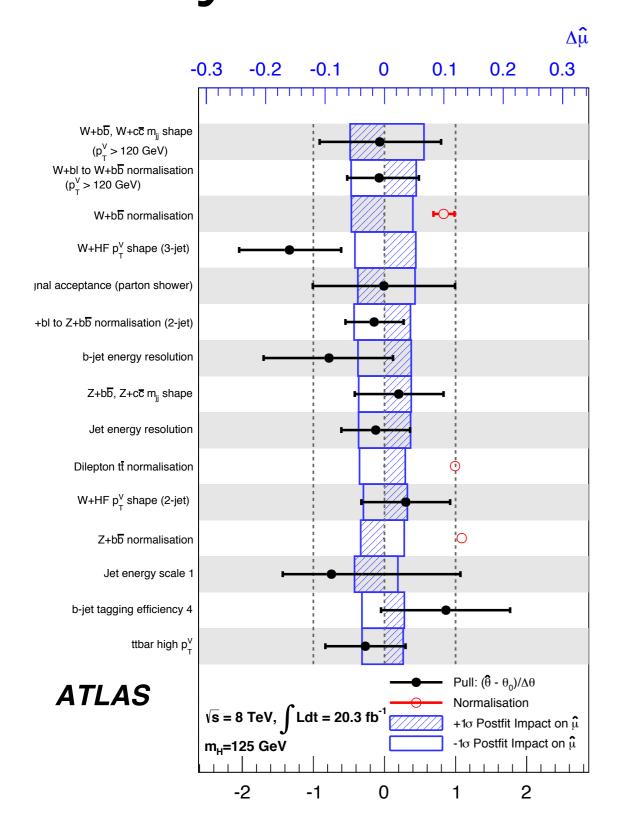


2-lepton MM+TT 2 jets

- Simultaneous profile likelihood fit to:
 - 2-tag: 27 signal regions
 - 1-tag: 11 control regions
- constraining backgrounds & extracting μ
- ~170 NPs to account for systematic effects

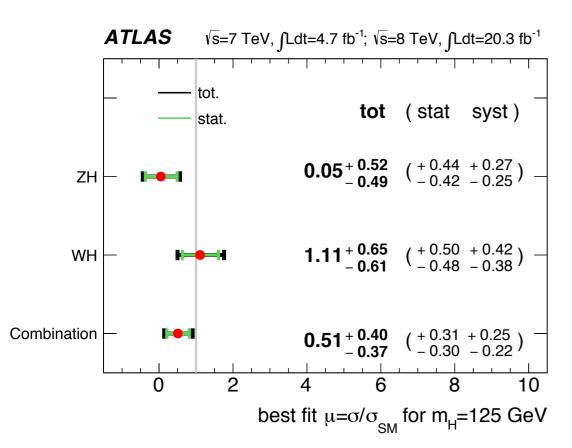


Systematic uncertainties



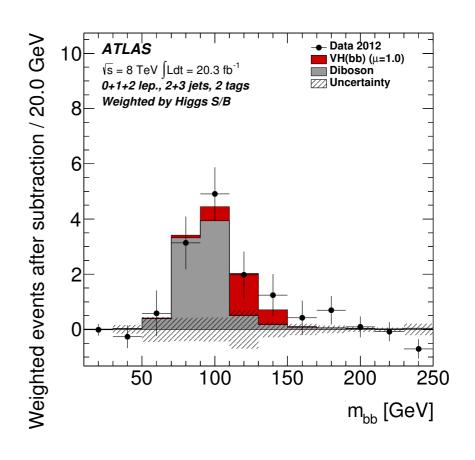
- Leading uncertainties:
 - W+b/c theory (shapes + flavor composition)
 - signal theory (parton shower)
 - jet energy resolution
- Important message for Run-II
 - Need to urgently improve our MC modeling (gluon splitting, heavy flavor fractions,...)

Events / 0.5 **ATLAS** Data 2012 10⁷ $VH(bb) (\mu=1.0)$ $\sqrt{s} = 8 \text{ TeV } \int Ldt = 20.3 \text{ fb}^{-1}$ Diboson 10⁶ Single top Multijet 10⁵ W+cl 10^{4} 10³ 10^{2} 10 Pull (stat.) $\log_{10}(S/B)$



Main result

- Uncertainty on µ decreased from ~0.65 to ~0.4!
- Expected sensitivity 2.6σ, observ. 1.4σ
 - $\mu = 0.5 \pm 0.4$
- No evidence for H to bb yet, but no incompatibility either.....



H to hh to 4b

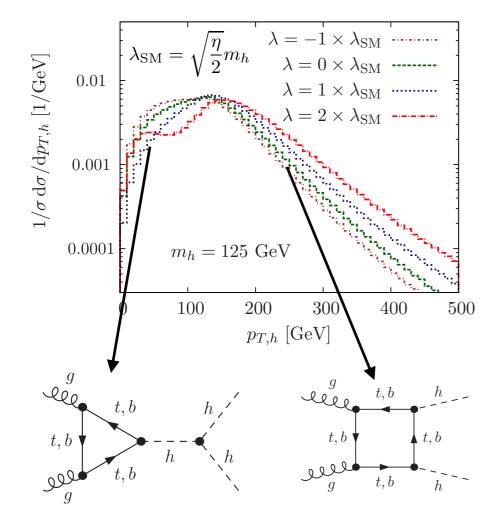
 Measurement of Higgs self-coupling crucial for full characterization of EW sector

$$V(H^{\dagger}H) = \mu^{2}H^{\dagger}H + \eta(H^{\dagger}H)^{2}$$

$$\supset \frac{1}{2}m_{h}^{2}h^{2} + \sqrt{\frac{\eta}{2}}m_{h}h^{3} + \frac{\eta}{4}h^{4}$$

 Very challenging at the LHC! (even with L=3000 inv. fb.)

$$2\eta = g^2 m_h^2 / m_W^2$$
 ?



- Di-Higgs production does not only involve triple Higgs vertex.
- Variation of coupling result in variation of differential pT(H) cross section.
- Important to finalize sensitivity studies now, since we already need to finalize the main detector design choices for Phase-II!!

X to hh to 4b

 Well before, we can start to investigate di-Higgs production in several BSM scenarios:

Spin-2 KK graviton in the Randall-Sundrum model with

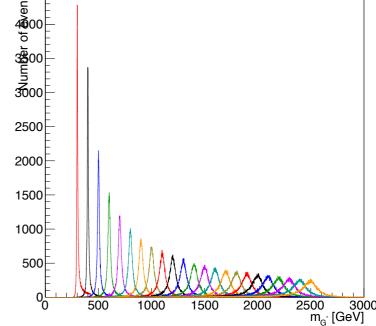
warped extra-dimensions

Search for masses up to ~2 TeV

- SUSY H to hh
- ... and many other models

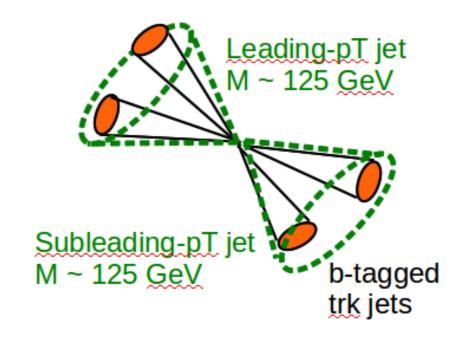


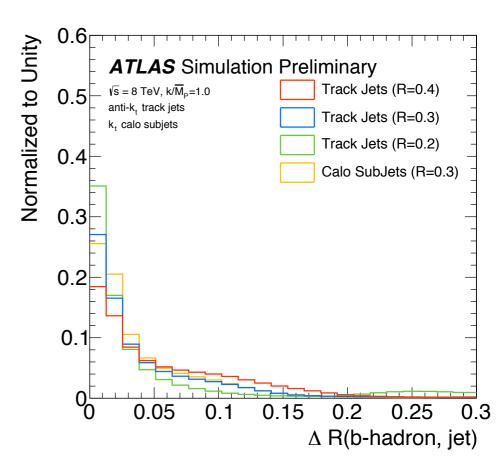
Challenging for reconstruction algorithms!



Reconstruction strategy

- Trigger based on multi-jet and multi-jet + b-jets signatures.
- Reconstruct two Anti-Kt R=1.0 calorimeter jets and trim them
 - Leading pT>350 GeV, sub-leading pT>250 GeV
- Then match these two jets to four Anti-Kt R=0.3 track-jets.
- Allows to resolve jets into sub-structure.
- Results in better alignment of "subjets" along true b-hadron direction (crucial for b-tagging!).

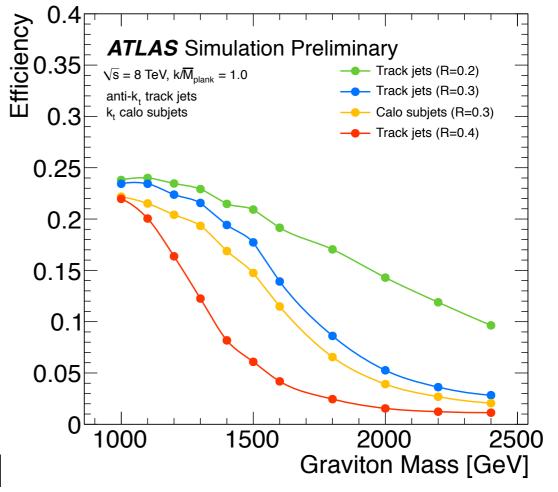




B-tagging the 4 b-jets...

- A proof of the concept is described in the ATLAS pub note ATL-PHYS-PUB-2014-13 and shown at Boost 2014.
- The new technique has the potential to significantly improve the sensitivity to a high mass resonance w.r.t. previous analyses.
- Track-jets are now being calibrated with data.
- X to 4b analysis progressing well





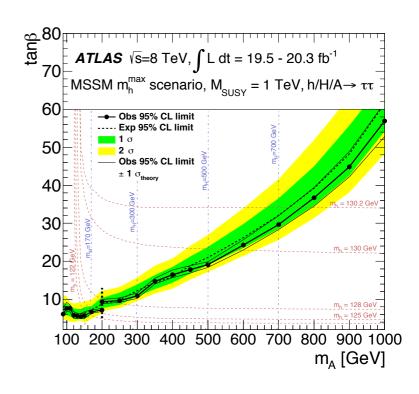
Plan is to publish full Run-I result on a few months timescale.

bH to bbb analysis

- Together with H → TT, main search channel for heavy A/H Higgs boson in MSSM scenarios with large tan(β).
- Very challenging analysis because of pure hadronic environment and relatively soft b-jet from associated production

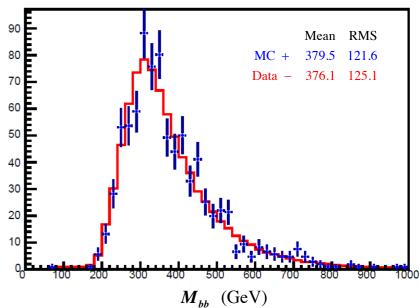


- Look for di b-jet peak over large bbb continuum background.
- Determine background shape from simultaneous likelihood fit to m(bb) in control region (bb + anti-tag).
 - Assumption that m(bb) is nearly independent on flavor of third jet crucial. Need theory and simulations to back this up!



Background simulation

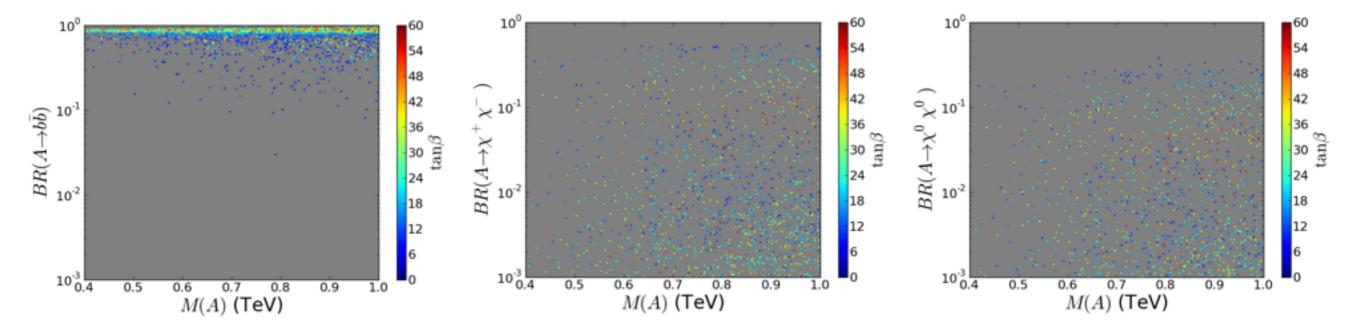
- Background mainly from multi-jet production of ≥ 3 b-jets difficult process to simulate
 - → help from Stephan to produce realistic Sherpa sample (thanks!!)
 - Replaces more simplistic model based on Pythia (2-to-2 scattering ME only)
 - Of general interest for studying multi-b production in ATLAS
- After quite some validation, got the MC request finally approved.
- Plan to finalize analysis soon, and have public results to show!



subprocesscross section (pb) $jj \rightarrow b\bar{b}jj$ 58531 $jj \rightarrow b\bar{b}j$ 33411 $bj \rightarrow bjj + \bar{b}j \rightarrow \bar{b}jj$ 22147 $bj \rightarrow bjj + \bar{b}j \rightarrow \bar{b}jj$ 16282 $bj \rightarrow bj + \bar{b}j \rightarrow \bar{b}j$ 12135 $jj \rightarrow b\bar{b}$ 1672 $cj \rightarrow cb\bar{b}j + \bar{c}j \rightarrow \bar{c}b\bar{b}j$ 1602 $bj \rightarrow bb\bar{b}j + b\bar{b}j \rightarrow \bar{b}b\bar{b}j$ 997 $jj \rightarrow bb\bar{c}\bar{c}$ 776 $cj \rightarrow cb\bar{b} + \bar{c}j \rightarrow \bar{c}b\bar{b}$ 681 $bj \rightarrow bb\bar{b} + \bar{b}j \rightarrow \bar{b}b\bar{b}$ 387 $jj \rightarrow bb\bar{b}b$ 387 $b\bar{c} \rightarrow b\bar{c}j + \bar{b}c \rightarrow \bar{b}cj$ 206 $bc \rightarrow bcj + \bar{b}\bar{c} \rightarrow \bar{b}\bar{c}j$ 194 $b\bar{c} \rightarrow b\bar{c}jj + \bar{b}\bar{c} \rightarrow \bar{b}\bar{c}jj$ 136 $bc \rightarrow bcj + \bar{b}\bar{c} \rightarrow \bar{b}\bar{c}j$ 136 $bc \rightarrow bc + \bar{b}\bar{c} \rightarrow \bar{b}\bar{c}$ 122 $b\bar{b} \rightarrow b\bar{b}j$ 62 $bb \rightarrow bbj + \bar{b}\bar{b} \rightarrow \bar{b}\bar{b}j$ 53 $b\bar{b} \rightarrow b\bar{b}j$ 39 $b\bar{b} \rightarrow b\bar{b}$ 37 $bb \rightarrow bb + \bar{b}\bar{b} \rightarrow \bar{b}\bar{b} \rightarrow \bar{b}\bar{b}$ 37		
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$cj \rightarrow cb\bar{b}j + \bar{c}j \rightarrow \bar{c}b\bar{b}j$ $bj \rightarrow bb\bar{b}j + \bar{b}j \rightarrow \bar{b}b\bar{b}j$ $jj \rightarrow b\bar{b}c\bar{c}$ $cj \rightarrow cb\bar{b} + \bar{c}j \rightarrow \bar{c}b\bar{b}$ $bj \rightarrow bb\bar{b} + \bar{b}j \rightarrow \bar{b}b\bar{b}$ 387 $jj \rightarrow b\bar{b}b\bar{b}$ $b\bar{c} \rightarrow b\bar{c}j + \bar{b}c \rightarrow \bar{b}cj$ $bc \rightarrow bcj + \bar{b}\bar{c} \rightarrow \bar{b}cj$ $bc \rightarrow bc + \bar{b}\bar{c} \rightarrow \bar{b}\bar{c}j$ $bc \rightarrow bc + \bar{b}\bar{c} \rightarrow \bar{b}\bar{c}$ $b\bar{b} \rightarrow b\bar{b}j$ $bb \rightarrow bbj + \bar{b}\bar{b} \rightarrow \bar{b}\bar{b}j$ 39	$bj \rightarrow bj + \bar{b}j \rightarrow \bar{b}j$	12135
$\begin{array}{c} bj \rightarrow bb\bar{b}j + \bar{b}j \rightarrow \bar{b}b\bar{b}j \\ jj \rightarrow b\bar{b}c\bar{c} \\ cj \rightarrow cb\bar{b} + \bar{c}j \rightarrow \bar{c}b\bar{b} \\ bj \rightarrow bb\bar{b} + \bar{b}j \rightarrow \bar{b}b\bar{b} \\ jj \rightarrow b\bar{b}b\bar{b} \\ b\bar{c} \rightarrow b\bar{c}j + \bar{b}c \rightarrow \bar{b}cj \\ bc \rightarrow bcj + \bar{b}c \rightarrow \bar{b}cj \\ bc \rightarrow bcjj + \bar{b}c \rightarrow \bar{b}cjj \\ bc \rightarrow bcj + \bar{b}c \rightarrow \bar{b}cjj \\ bc \rightarrow bcj + \bar{b}c \rightarrow \bar{b}cjj \\ bc \rightarrow bcj + \bar{b}c \rightarrow \bar{b}cjj \\ bc \rightarrow bc + \bar{b}c \rightarrow \bar{b}c \\ b\bar{c} \rightarrow b\bar{c} + \bar{b}c \rightarrow \bar{b}c \\ b\bar{b} \rightarrow b\bar{b}j \\ bb \rightarrow bbj + \bar{b}b \rightarrow \bar{b}bj \\ b\bar{b} \rightarrow b\bar{b}j \\ b\bar{b} \rightarrow b\bar{b}j \\ \end{array}$	$jj \to b\bar{b}$	1672
$jj \rightarrow b\bar{b}c\bar{c}$ $cj \rightarrow cb\bar{b} + \bar{c}j \rightarrow \bar{c}b\bar{b}$ $bj \rightarrow bb\bar{b} + \bar{b}j \rightarrow \bar{b}b\bar{b}$ $jj \rightarrow b\bar{b}b\bar{b}$ $b\bar{c} \rightarrow b\bar{c}j + \bar{b}c \rightarrow \bar{b}cj$ $bc \rightarrow bcj + \bar{b}c \rightarrow \bar{b}cj$ $bc \rightarrow bcjj + \bar{b}c \rightarrow \bar{b}cjj$ $bc \rightarrow bcjj + \bar{b}c \rightarrow \bar{b}cjj$ $bc \rightarrow bcjj + \bar{b}c \rightarrow \bar{b}cjj$ $bc \rightarrow bc + \bar{b}c \rightarrow \bar{b}c$ $b\bar{c} \rightarrow b\bar{c} + \bar{b}c \rightarrow \bar{b}c$ $b\bar{c} \rightarrow b\bar{c} + \bar{b}c \rightarrow \bar{b}c$ $b\bar{b} \rightarrow b\bar{b}j$ $bb \rightarrow bbj + \bar{b}b \rightarrow \bar{b}bj$ $bb \rightarrow bbjj + \bar{b}b \rightarrow \bar{b}bjj$ $bb \rightarrow bbjj + \bar{b}b \rightarrow \bar{b}bjj$ 39	$cj \to cb\bar{b}j + \bar{c}j \to \bar{c}b\bar{b}j$	1602
$cj \rightarrow cb\bar{b} + \bar{c}j \rightarrow \bar{c}b\bar{b}$ $bj \rightarrow bb\bar{b} + \bar{b}j \rightarrow \bar{b}b\bar{b}$ $jj \rightarrow b\bar{b}b\bar{b}$ $b\bar{c} \rightarrow b\bar{c}j + \bar{b}c \rightarrow \bar{b}cj$ $bc \rightarrow bcj + \bar{b}\bar{c} \rightarrow \bar{b}cj$ $b\bar{c} \rightarrow b\bar{c}jj + \bar{b}c \rightarrow \bar{b}cjj$ $bc \rightarrow bcjj + \bar{b}\bar{c} \rightarrow \bar{b}cjj$ $bc \rightarrow bcjj + \bar{b}\bar{c} \rightarrow \bar{b}\bar{c}jj$ $bc \rightarrow bc + \bar{b}\bar{c} \rightarrow \bar{b}\bar{c}$ $b\bar{c} \rightarrow b\bar{c} + \bar{b}c \rightarrow \bar{b}c$ $b\bar{b} \rightarrow b\bar{b}j$ $bb \rightarrow bbj + \bar{b}\bar{b} \rightarrow \bar{b}\bar{b}j$ $bb \rightarrow bbj + \bar{b}\bar{b} \rightarrow \bar{b}\bar{b}jj$ $bb \rightarrow bbj + \bar{b}\bar{b} \rightarrow \bar{b}\bar{b}jj$ 39	$bj \to bb\bar{b}j + \bar{b}j \to \bar{b}b\bar{b}j$	997
$\begin{array}{cccccccccccccccccccccccccccccccccccc$	$jj \to b\bar{b}c\bar{c}$	776
$jj \rightarrow b\bar{b}b\bar{b}$ $b\bar{c} \rightarrow b\bar{c}j + \bar{b}c \rightarrow \bar{b}cj$ $bc \rightarrow bcj + \bar{b}\bar{c} \rightarrow b\bar{c}j$ $b\bar{c} \rightarrow b\bar{c}jj + \bar{b}c \rightarrow \bar{b}cjj$ $bc \rightarrow bcjj + \bar{b}\bar{c} \rightarrow \bar{b}cjj$ $bc \rightarrow bcjj + \bar{b}\bar{c} \rightarrow \bar{b}\bar{c}jj$ $bc \rightarrow bc + \bar{b}\bar{c} \rightarrow \bar{b}\bar{c}$ $b\bar{c} \rightarrow b\bar{c} + \bar{b}c \rightarrow \bar{b}c$ $b\bar{b} \rightarrow b\bar{b}j$ $bb \rightarrow bbj + \bar{b}\bar{b} \rightarrow \bar{b}\bar{b}j$ $bb \rightarrow bbjj + \bar{b}\bar{b} \rightarrow \bar{b}\bar{b}jj$ $bb \rightarrow bbjj + \bar{b}\bar{b} \rightarrow \bar{b}\bar{b}jj$ 39 $b\bar{b} \rightarrow b\bar{b}$	$cj \to cb\bar{b} + \bar{c}j \to \bar{c}b\bar{b}$	681
$\begin{array}{cccccccccccccccccccccccccccccccccccc$	$bj \to bb\bar{b} + \bar{b}j \to \bar{b}b\bar{b}$	387
$bc \rightarrow bcj + \bar{b}\bar{c} \rightarrow \bar{b}\bar{c}j$ $b\bar{c} \rightarrow b\bar{c}jj + \bar{b}c \rightarrow \bar{b}cjj$ $bc \rightarrow bcjj + \bar{b}\bar{c} \rightarrow \bar{b}\bar{c}jj$ $bc \rightarrow bcjj + \bar{b}\bar{c} \rightarrow \bar{b}\bar{c}jj$ $bc \rightarrow bc + \bar{b}\bar{c} \rightarrow \bar{b}\bar{c}$ $b\bar{c} \rightarrow b\bar{c} + \bar{b}c \rightarrow \bar{b}c$ $b\bar{b} \rightarrow b\bar{b}j$ $bb \rightarrow bbj + \bar{b}\bar{b} \rightarrow \bar{b}\bar{b}j$ $bb \rightarrow bbjj + \bar{b}\bar{b} \rightarrow \bar{b}\bar{b}jj$ $bb \rightarrow bbjj + \bar{b}\bar{b} \rightarrow \bar{b}\bar{b}jj$ $bb \rightarrow bbjj + \bar{b}\bar{b} \rightarrow \bar{b}\bar{b}jj$ 39 $b\bar{b} \rightarrow b\bar{b}$	$jj \to b\bar{b}b\bar{b}$	376
$\begin{array}{cccccccccccccccccccccccccccccccccccc$	$b\bar{c} \to b\bar{c}j + \bar{b}c \to \bar{b}cj$	206
$bc \rightarrow bcjj + \bar{b}\bar{c} \rightarrow \bar{b}\bar{c}jj$ $bc \rightarrow bc + \bar{b}\bar{c} \rightarrow \bar{b}\bar{c}$ $b\bar{c} \rightarrow b\bar{c} + \bar{b}c \rightarrow \bar{b}c$ $b\bar{b} \rightarrow b\bar{b}j$ $bb \rightarrow bbj + \bar{b}b \rightarrow \bar{b}bj$ $bb \rightarrow bbjj$ $bb \rightarrow b\bar{b}j$ $b\bar{b} \rightarrow b\bar{b}j$	$bc \to bcj + \bar{b}\bar{c} \to \bar{b}\bar{c}j$	194
$\begin{array}{cccccccccccccccccccccccccccccccccccc$	$b\bar{c} \to b\bar{c}jj + \bar{b}c \to \bar{b}cjj$	143
$\begin{array}{cccccccccccccccccccccccccccccccccccc$	$bc \to bcjj + \bar{b}\bar{c} \to \bar{b}\bar{c}jj$	136
$\begin{array}{cccccccccccccccccccccccccccccccccccc$	$bc \to bc + \bar{b}\bar{c} \to \bar{b}\bar{c}$	122
$\begin{array}{cccccccccccccccccccccccccccccccccccc$	$b\bar{c} \to b\bar{c} + \bar{b}c \to \bar{b}c$	121
$\begin{array}{cccccccccccccccccccccccccccccccccccc$	$b\bar{b} \to b\bar{b}j$	62
$\begin{array}{cccccccccccccccccccccccccccccccccccc$	$bb \to bbj + \bar{b}\bar{b} \to \bar{b}\bar{b}j$	53
$b\bar{b} \to b\bar{b}$ 37	$b\bar{b} \to b\bar{b}jj$	44
	$bb \to bbjj + \bar{b}\bar{b} \to \bar{b}\bar{b}jj$	39
$bb \to bb + \bar{b}\bar{b} \to \bar{b}\bar{b}$ 30	$b\bar{b} \to b\bar{b}$	37
	$bb \to bb + \bar{b}\bar{b} \to \bar{b}\bar{b}$	30

MSSM Models

- Benchmark models defined within Higgs cross section working group (https://twiki.cern.ch/twiki/bin/view/LHCPhysics/LHCHXSWGMSSMNeutral).
- Most of these have very low BR(H/A →bb) < 50%.
- We were wondering why and have asked Tom and JoAnne for advice.



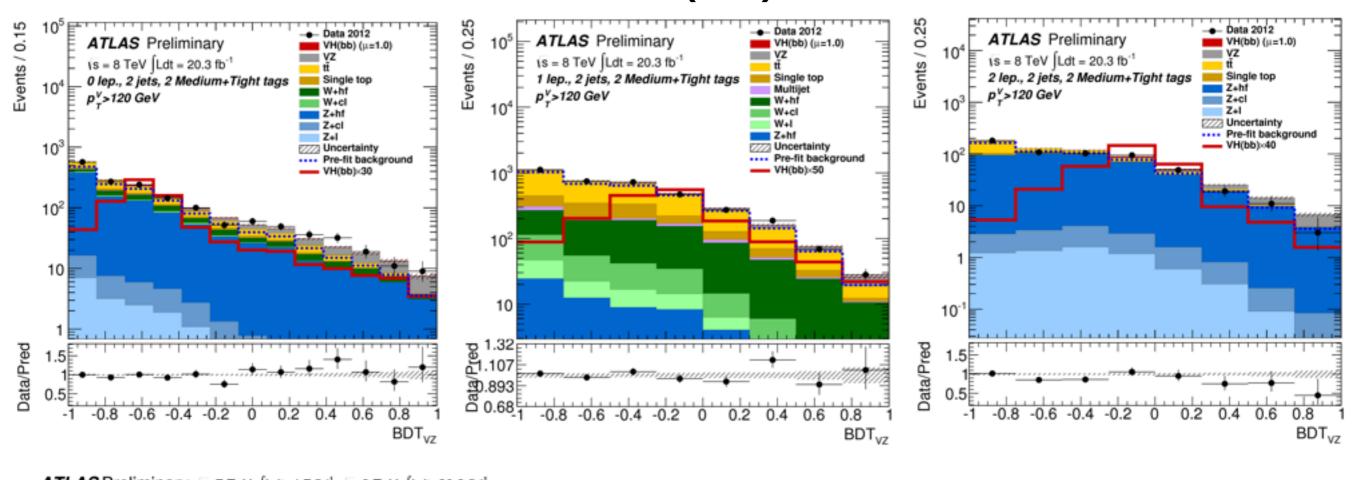
- From their findings (thanks!), most of the allowed A/H decays to SUSY are in charginos and neutralinos, but scenarios with $BR(H/A \rightarrow bb) < 50\%$ remain very unusual.
- Time to add more representative SUSY scenarios?
- An interesting question remains how much BR(bb)/BR(tt) varies from model to model. This determines the relative sensitivity of the TT and bb searches.

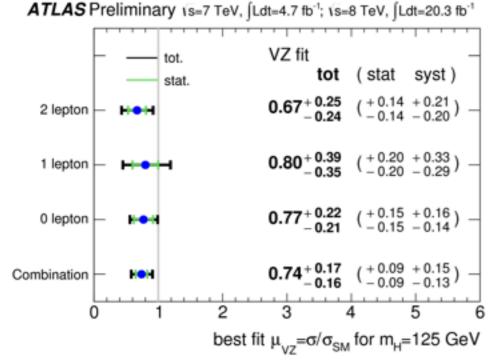
Plans for the future

- Observe a clear H → bb signal
 - But need to reduce uncertainties, including theory:
 - SM measurement of W+bb background, in particular challenging region with g → bb splitting?
- Look in early data for resonances decaying into two or 4 b-jets
 - We are now able to access phase space regions with very close-by b-jets: this will extend our reach in many models!
- In 13 TeV data, triggering on bH will be more challenging, but we plan continuing looking into this channel.

Backup

Cross-check (1): VZ to Vbb



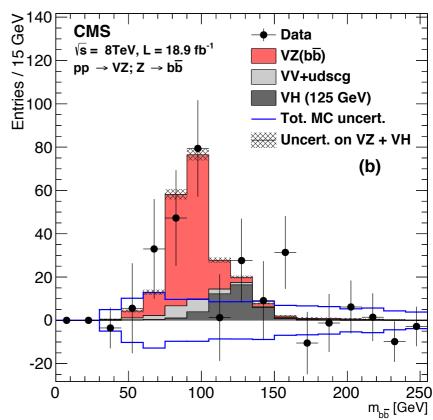


- Train MVA to extract diboson signal
- Clear signal found, all channels well compatible
- $\sim 1.5\sigma$ low

Comparison to CMS

ATLAS		CMS	
σ [exp]	σ [obs]	σ [exp]	σ [obs]
2.6	1.4	2.1*	2.1

* neglects pT dependence of gg to ZH contribution, 10% increase in sensitivity?



- New ATLAS analysis has slightly better sensitivity, but observed signal lower
- Combining ATLAS and CMS won't get us to evidence: will need to wait for some good Run-II data
- In the meantime, hope the coupling fits will profit from increased precision of H to bb signal strength