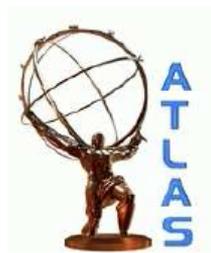


Probing quartic gauge boson couplings at the LHC with the vector boson fusion, vector boson scattering and triboson processes

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on behalf of ATLAS and CMS collaborations



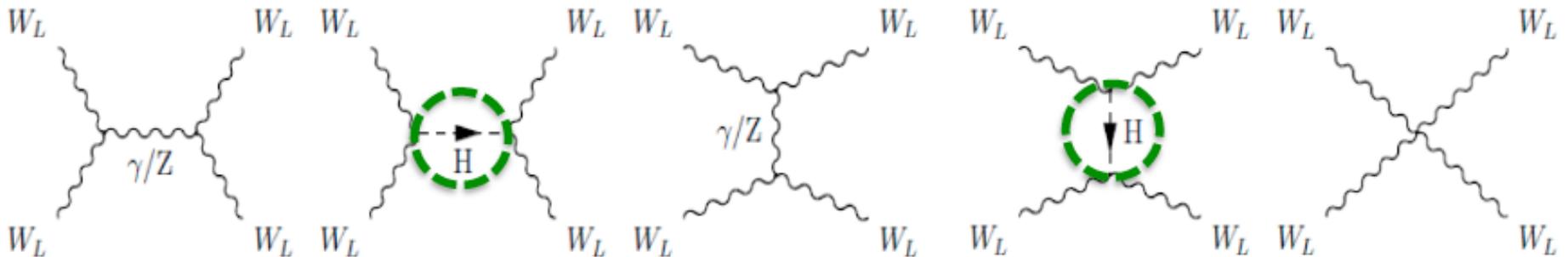
Presented in SM@LHC Conference 2015 at Galileo Galilei Institute

Outline

- Vector boson fusion (VBF):
 - VBF Z+2jets
 - ATLAS: JHEP 04 (2014) 031
 - CMS: Eur. Phys. J. C 75 (2015) 66
- Vector boson scattering (VBS):
 - $W^\pm W^\pm jj$
 - ATLAS: Phys. Rev. Lett. 113, 141803 (2014)
 - CMS: Phys. Rev. Lett. 114, 051801 (2015)
- Triboson measurements
 - $WW\gamma/WZ\gamma$ (CMS: Phys. Rev. D 90, 032008 (2014))
 - $W\gamma\gamma$ (ATLAS: arXiv:1503.03243)
- Exclusive WW
 - CMS: JHEP 1307 (2013) 116

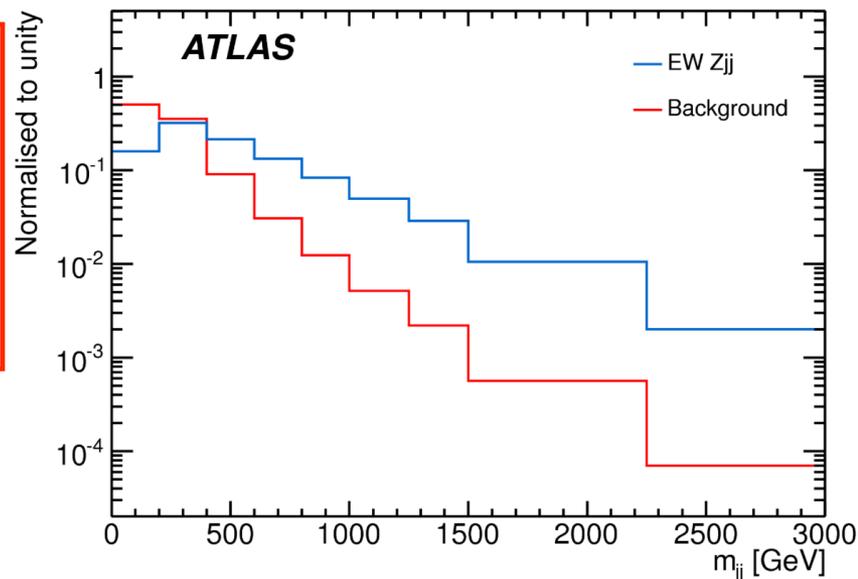
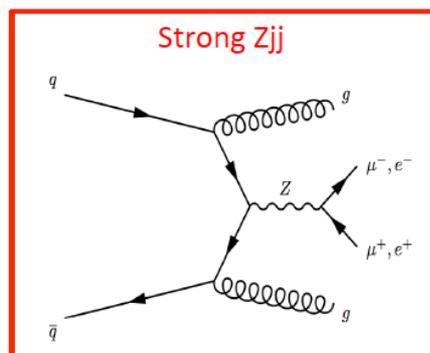
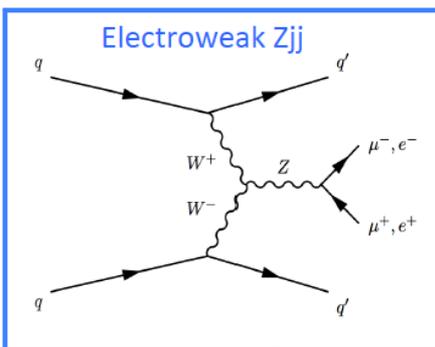
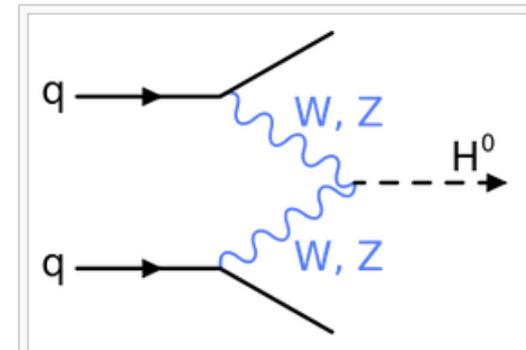
Introduction

- **The Higgs mechanism \neq a Higgs boson !**
- Vector boson self-coupling is a fundamental prediction of the Electroweak Sector of SM
- Its study is important to understanding electroweak symmetry breaking mechanism



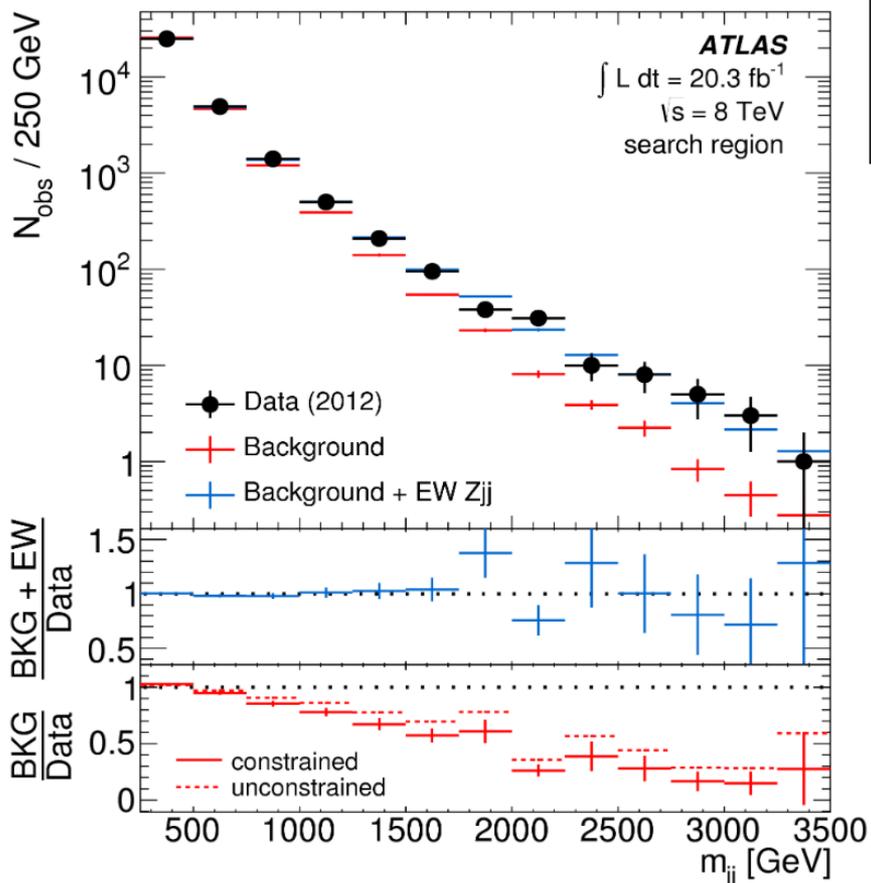
VBF Z

- Two main Higgs production mechanism
 - Gluon-gluon fusion 4
 - Vector boson fusion (VBF)
- Higgs VBF production mechanism has not been confirmed by experimental measurements with 5 sigma significance
- Important to confirm VBF production mechanism in VBF Z+2jets channel



VBF Z (ATLAS)

Background-only hypothesis
Rejected at greater than
 5σ significance



Object selections:

- two electron/muon
- two high- p_T forward jets

Kinematic selections:

- $81 < m_{ll} < 101 \text{ GeV}$
- $p_{T}^{ll} > 20 \text{ GeV}$
- $p_{T}^{\text{balance}} < 0.15$
- $N_{\text{jet}}^{\text{gap}} = 0$
- $m_{jj} > 250 \text{ GeV}$

	Electron+muon
Data	32186
MC predicted N_{bkg}	$32600 \pm 2600^{+3400}_{-4000}$
MC predicted N_{EW}	$1333 \pm 50 \pm 40$
Fitted N_{bkg}	$30530 \pm 216 \pm 40$
Fitted N_{EW}	$1657 \pm 134 \pm 40$

Measured EWK Zjj cross section

$$\sigma_{\text{EW}} = 54.7 \pm 4.6 \text{ (stat)}^{+9.8}_{-10.4} \text{ (syst)} \pm 1.5 \text{ (lumi)} \text{ fb.}$$

Powheg Box predictions at
next-to-leading-order (NLO) accuracy
in perturbative QCD

$$46.1 \pm 0.2 \text{ (stat)}^{+0.3}_{-0.2} \text{ (scale)} \pm 0.8 \text{ (PDF)} \pm 0.5 \text{ (model)} \text{ fb,}$$

VBF Z (CMS)

Eur. Phys. J. C 75 (2015) 66,

- Use boosted decision tree (BDT) technique to improve sensitivity

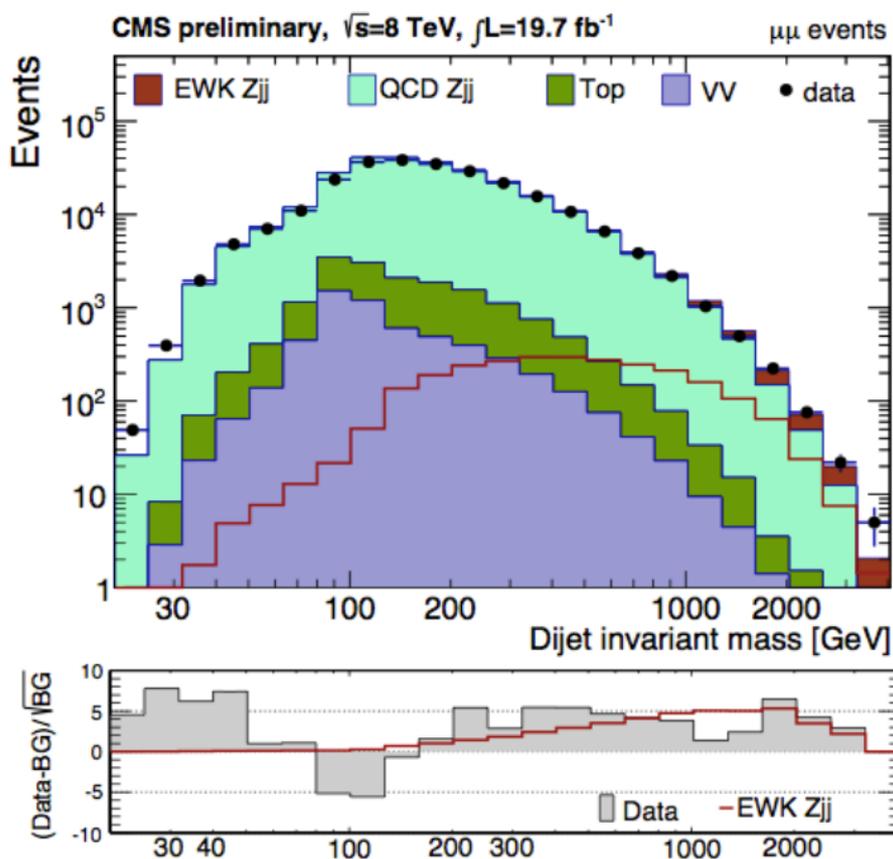
$M_{ll} > 50$ GeV,

$p_{Tj} > 25$ GeV, and $|\eta_j| < 5$, $M_{jj} > 120$ GeV,

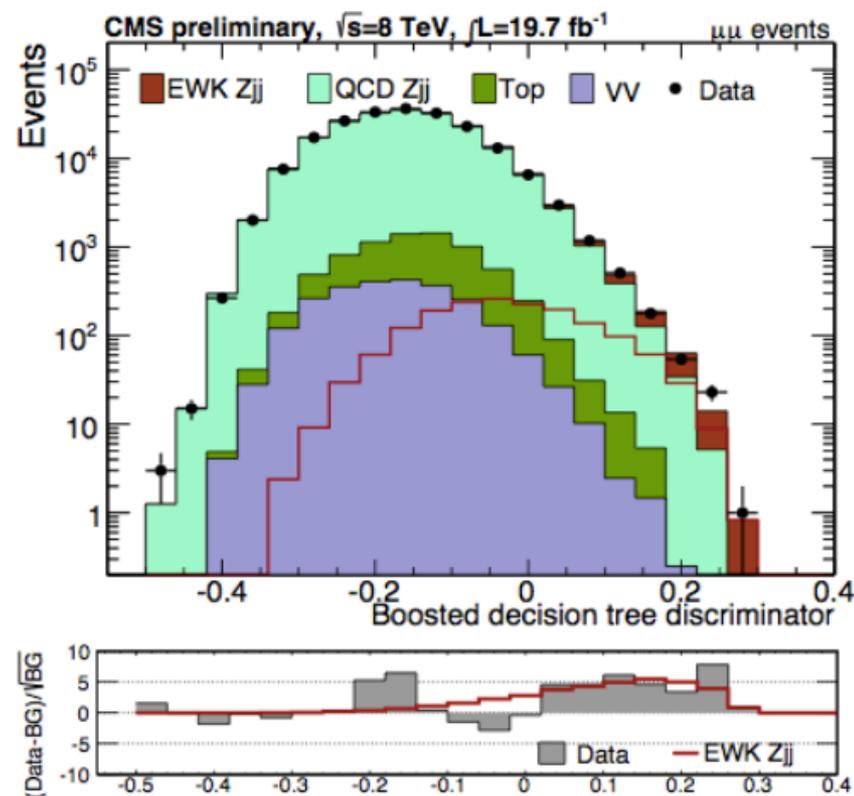
Measured $\sigma_{EW}(lljj) = 174 \pm 15$ (stat) ± 40 (syst) fb

5 σ significance

M_{jj} spectrum

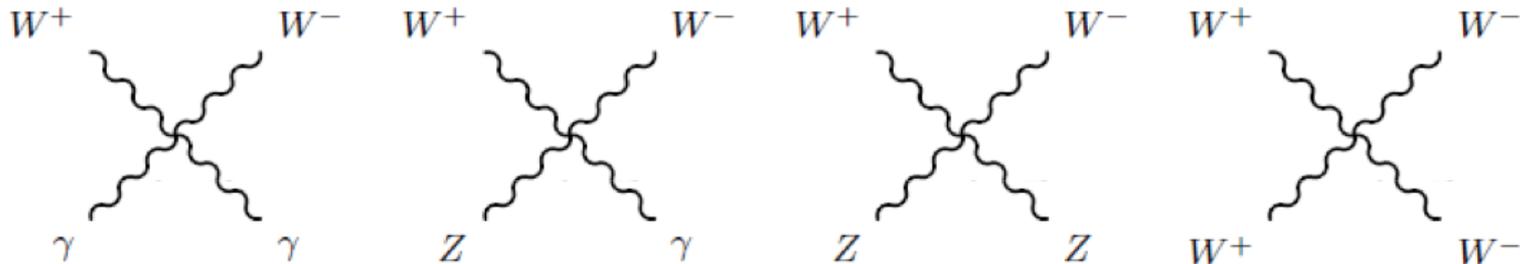


BDT output spectrum



Quartic Gauge Boson Couplings

- Reminder of **Quartic Gauge Boson Couplings (QGCs)**
 - SM model predicts gauge boson self coupling
 - Four gauge boson vertex:
 - $WW\gamma\gamma$, $WWZ\gamma$, $WWWW$, $WWZZ$, $ZZZZ$...



- QGCs can be studied in
 - Tri-boson processes
 - Vector boson scattering processes
 - Exclusive $\gamma\gamma \rightarrow WW$ process

Modeling of aQGCs: Dim 8 EFT models

- Extension of the effective SM-Lagrangian by introducing additional dimension-8 operators for QGCs
 - no effect on TGCs.

Higgs field

$$\mathcal{L}_{S,0} = [(D_\mu \Phi)^\dagger D_\nu \Phi] \times [(D^\mu \Phi)^\dagger D^\nu \Phi]$$

$$\mathcal{L}_{S,1} = [(D_\mu \Phi)^\dagger D^\mu \Phi] \times [(D_\nu \Phi)^\dagger D^\nu \Phi]$$

Higgs - Gauge boson field (L_M)

$$\mathcal{L}_{M,0} = \text{Tr} [\hat{W}_{\mu\nu} \hat{W}^{\mu\nu}] \times [(D_\beta \Phi)^\dagger D^\beta \Phi]$$

$$\mathcal{L}_{M,1} = \text{Tr} [\hat{W}_{\mu\nu} \hat{W}^{\nu\beta}] \times [(D_\beta \Phi)^\dagger D^\mu \Phi]$$

$$\mathcal{L}_{M,2} = [B_{\mu\nu} B^{\mu\nu}] \times [(D_\beta \Phi)^\dagger D^\beta \Phi]$$

$$\mathcal{L}_{M,3} = [B_{\mu\nu} B^{\nu\beta}] \times [(D_\beta \Phi)^\dagger D^\mu \Phi]$$

$$\mathcal{L}_{M,4} = [(D_\mu \Phi)^\dagger \hat{W}_{\beta\nu} D^\mu \Phi] \times B^{\beta\nu}$$

$$\mathcal{L}_{M,5} = [(D_\mu \Phi)^\dagger \hat{W}_{\beta\nu} D^\nu \Phi] \times B^{\beta\mu}$$

$$\mathcal{L}_{M,6} = [(D_\mu \Phi)^\dagger \hat{W}_{\beta\nu} \hat{W}^{\beta\mu} D^\nu \Phi]$$

$$\mathcal{L}_{M,7} = [(D_\mu \Phi)^\dagger \hat{W}_{\beta\nu} \hat{W}^{\beta\mu} D^\nu \Phi]$$

Gauge boson field (L_T)

$$\mathcal{L}_{T,0} = \text{Tr} [\hat{W}_{\mu\nu} \hat{W}^{\mu\nu}] \times \text{Tr} [\hat{W}_{\alpha\beta} \hat{W}^{\alpha\beta}]$$

$$\mathcal{L}_{T,1} = \text{Tr} [\hat{W}_{\alpha\nu} \hat{W}^{\mu\beta}] \times \text{Tr} [\hat{W}_{\mu\beta} \hat{W}^{\alpha\nu}]$$

$$\mathcal{L}_{T,2} = \text{Tr} [\hat{W}_{\alpha\mu} \hat{W}^{\mu\beta}] \times \text{Tr} [\hat{W}_{\beta\nu} \hat{W}^{\nu\alpha}]$$

$$\mathcal{L}_{T,5} = \text{Tr} [\hat{W}_{\mu\nu} \hat{W}^{\mu\nu}] \times B_{\alpha\beta} B^{\alpha\beta}$$

$$\mathcal{L}_{T,6} = \text{Tr} [\hat{W}_{\alpha\nu} \hat{W}^{\mu\beta}] \times B_{\mu\beta} B^{\alpha\nu}$$

$$\mathcal{L}_{T,7} = \text{Tr} [\hat{W}_{\alpha\mu} \hat{W}^{\mu\beta}] \times B_{\beta\nu} B^{\nu\alpha}$$

$$\mathcal{L}_{T,8} = B_{\mu\nu} B^{\mu\nu} B_{\alpha\beta} B^{\alpha\beta}$$

$$\mathcal{L}_{T,9} = B_{\alpha\mu} B^{\mu\beta} B_{\beta\nu} B^{\nu\alpha}$$

VVjj final state	ZZ	Z γ $\gamma\gamma$	W ⁺ W ⁻ WZ	W [±] W [±]	W γ	
VVV final state	ZZZ	ZZ γ Z $\gamma\gamma$	WWZ WZZ	WWW	WV γ	$\gamma\gamma\gamma$
f _{S,0} , f _{S,1}	✓		✓	✓		
f _{M,0} , f _{M,1} , f _{M,6} , f _{M,7}	✓	✓	✓	✓	✓	
f _{M,2} , f _{M,3} , f _{M,4} , f _{M,5}	✓	✓	✓		✓	
f _{T,0} , f _{T,1} , f _{T,2}	✓	✓	✓	✓	✓	✓
f _{T,5} , f _{T,6} , f _{T,7}	✓	✓	✓		✓	✓
f _{T,8} , f _{T,9}	✓	✓				✓

The measurement of Vector boson scattering processes

- Vector boson scattering (VBS) is one of the most promising process to study QGCs.

- Diboson + two forward jets in event topology

- The scattering of longitudinally polarized vector bosons

- violates unitarity at $\sim 1\text{TeV}$ without higgs

- Important to check

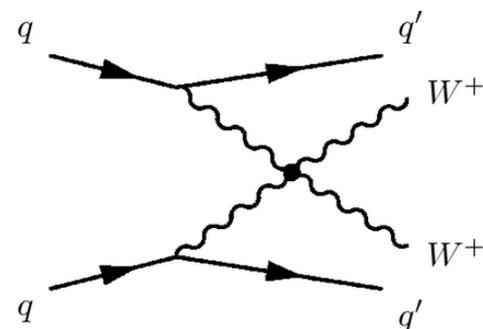
- whether Higgs boson unitarizes it fully or only partially

- The first VBS analysis :Same sign WW

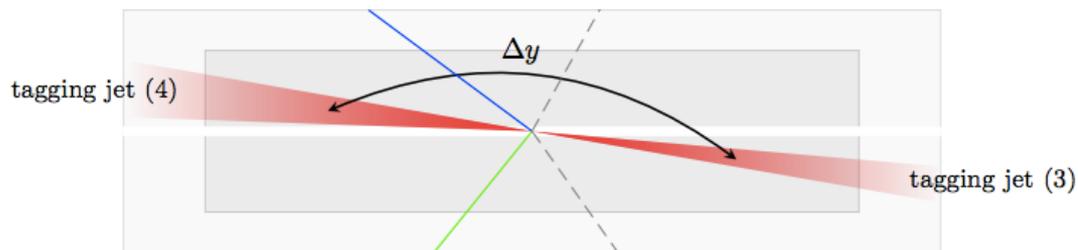
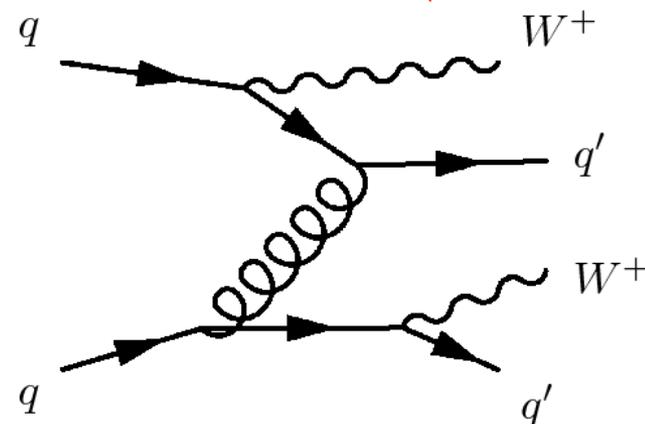
- Sensitive to WWWW vertex

- Final state: W^+W^+jj or W^-W^-jj

Electroweak
 $W^\pm W^\pm jj$ $O(\alpha_{EW}^6)$



QCD $W^\pm W^\pm jj$
 $O(\alpha_{EW}^4 \alpha_{QCD}^2)$



Electroweak $W^\pm W^\pm jj$ (ATLAS)

Object selection

- same sign di-leptons
- $p_{T,l} > 25$ GeV
- Two high pT jets

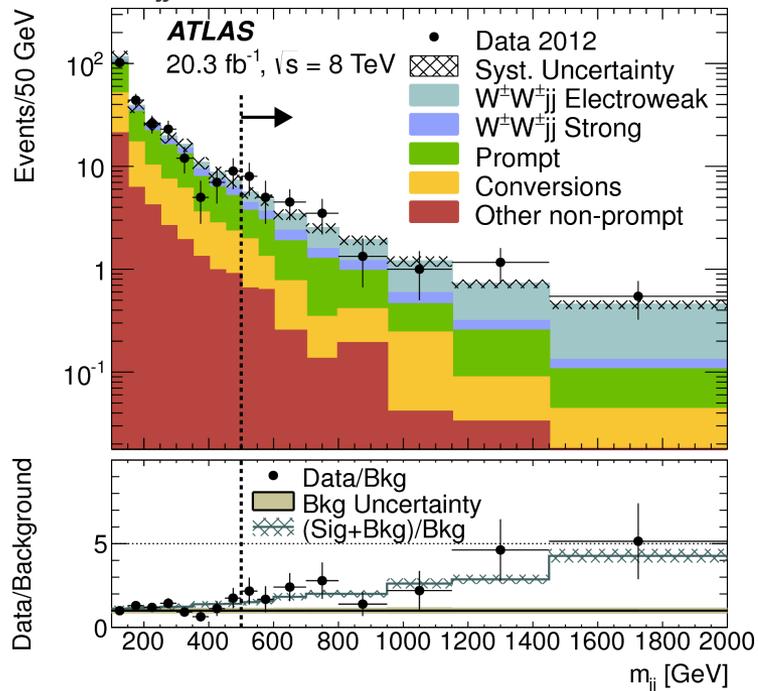
EWK signal only measurement

- $M_{ll} > 20$ GeV
- $|M_{ee} - M_Z| > 10$ GeV
- $M_{jj} > 500$ GeV
- $|\Delta\eta_{jj}| > 2.4$

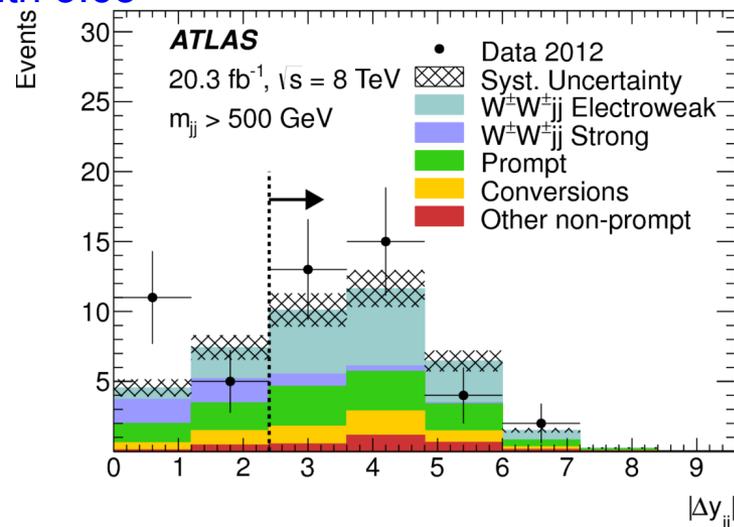
Measured $\sigma_{EW}(W^\pm W^\pm jj) = 1.3 \pm 0.4(\text{stat}) \pm 0.2(\text{syst})$ fb
 Predicted $\sigma_{EW}(W^\pm W^\pm jj) = 0.95 \pm 0.06$ fb
 Observed with 3.6σ

Inclusive QCD+EWK $W^\pm W^\pm jj$ measurement

- $M_{ll} > 20$ GeV
- $|M_{ee} - M_Z| > 10$ GeV
- $M_{jj} > 150$ GeV



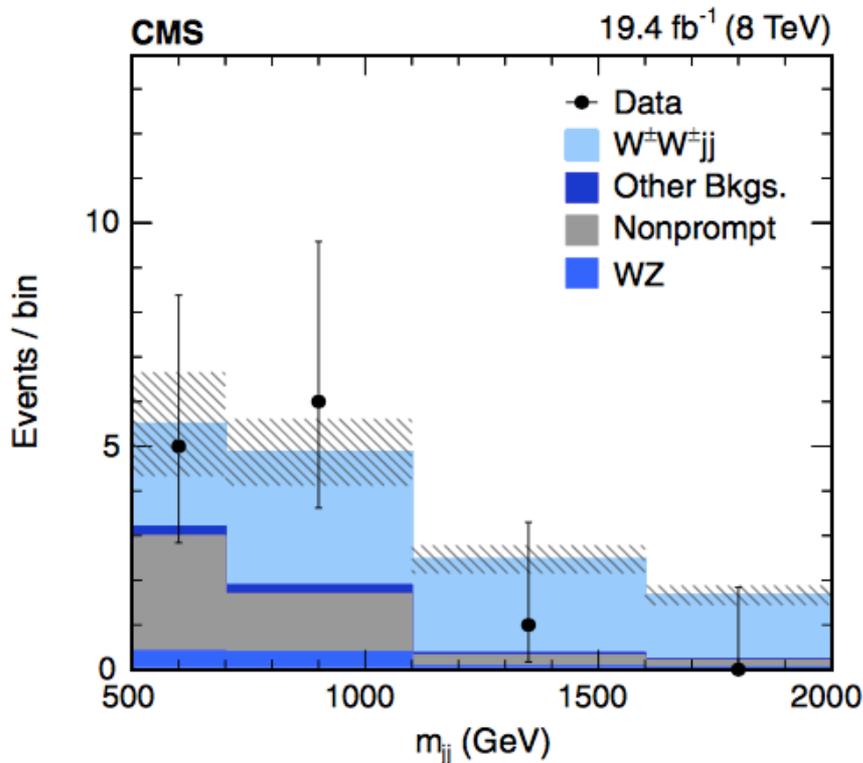
Measured $\sigma_{EW+QCD}(W^\pm W^\pm jj) = 2.1 \pm 0.5(\text{stat}) \pm 0.3(\text{syst})$ fb
 Predicted $\sigma_{EW+QCD}(W^\pm W^\pm jj) = 1.52 \pm 0.11$ fb
 Observed with 4.5σ



SM predictions are calculated using PowhegBox at NLO accuracy in perturbative QCD

Electroweak $W^\pm W^\pm jj$ (CMS)

Phys. Rev. Lett. 114, 051801 (2015)



Object selection

- same sign di-leptons
- $p_{Tj} > 25$ GeV, and $|\eta_j| < 5$
- Two high p_T forward jets
- $E_{T,miss} > 40$ GeV

EWK signal region

- $M_{ll} > 20$ GeV
- $|M_{ee} - M_Z| > 10$ GeV
- $M_{jj} > 500$ GeV
- $|\Delta\eta_{jj}| > 2.4$

$$\text{Measured } \sigma_{EW}(W^\pm W^\pm jj) = 4.0^{+2.4}_{-2.0} \text{ (stat)}^{+1.1}_{-1.0} \text{ (syst)} \text{ fb}$$

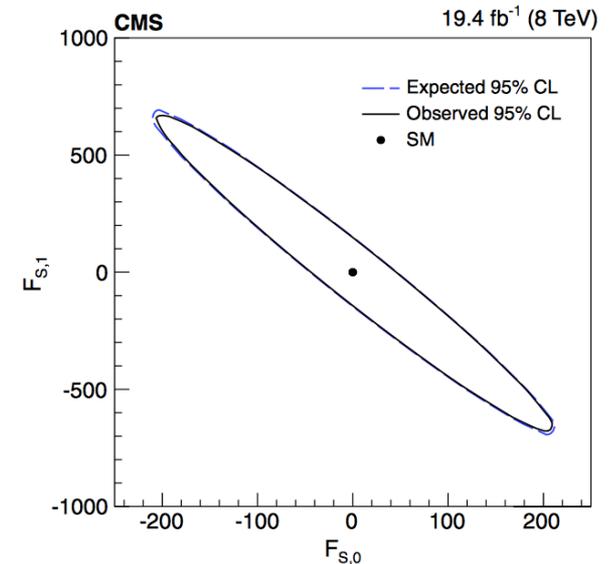
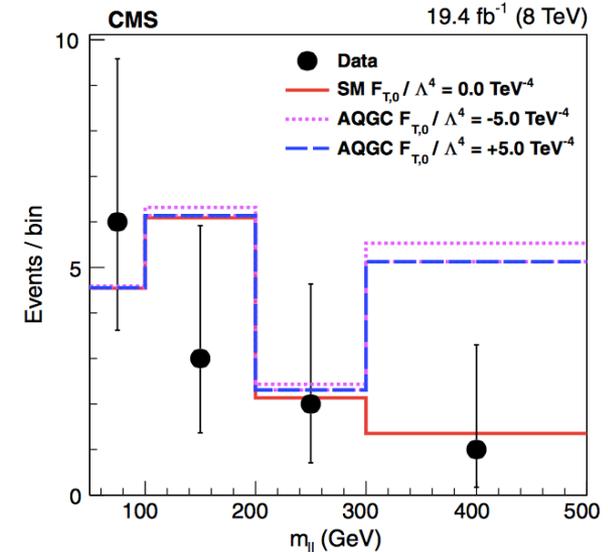
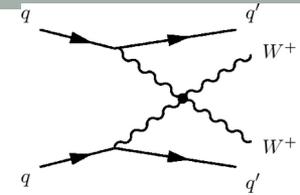
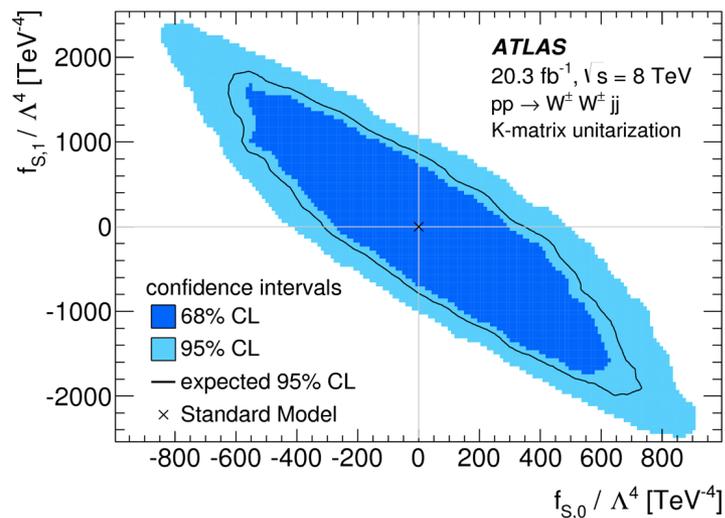
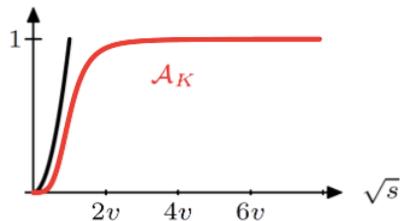
$$\text{Predicted } \sigma_{EW}(W^\pm W^\pm jj) = 5.8 \pm 1.2 \text{ fb.}$$

Observed with 2.0σ

aQGCs limits from VBS process

- FS0 and FS1 in Dim 8 EFT model is related to the Higgs field
- Naive EFT-predicted aTGC/aQGCs amplitudes
 - disrespect the gauge symmetry
 - and violate the unitarity once the \sqrt{s} goes sufficiently
- CMS result is not unitarized
- ATLAS use K matrix unitarization
 - Preserve unitarity in high \sqrt{s}
 - Unitarization with the k-matrix approach (arxiv: 0806.4145)
 - K-matrix amplitude

$$|\mathcal{A}_K(s)|^2 \xrightarrow{s \rightarrow \infty} 1$$



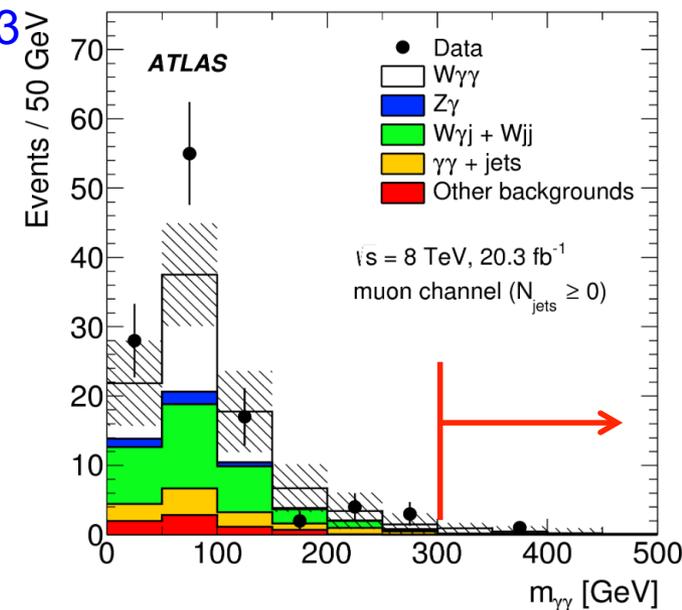
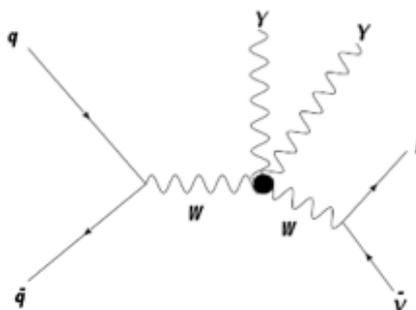
$W\gamma\gamma$ (ATLAS)

arXiv:1503.03243

- $W\gamma\gamma$ process is sensitive to $WW\gamma\gamma$ QGCs Vertex.

- Other contributions from:

- ISR photons
- FSR photons
- Photon from TGC vertex
- Photon from jet fragmentations



- MCFM is used for NLO $W\gamma\gamma$ SM predicted cross section.

- The measured cross section in inclusive case is 1.9 σ higher than predictions
- Better agreement in exclusive case

*veto events with hard jets with $p_T > 30 \text{ GeV}$ in exclusive measurement

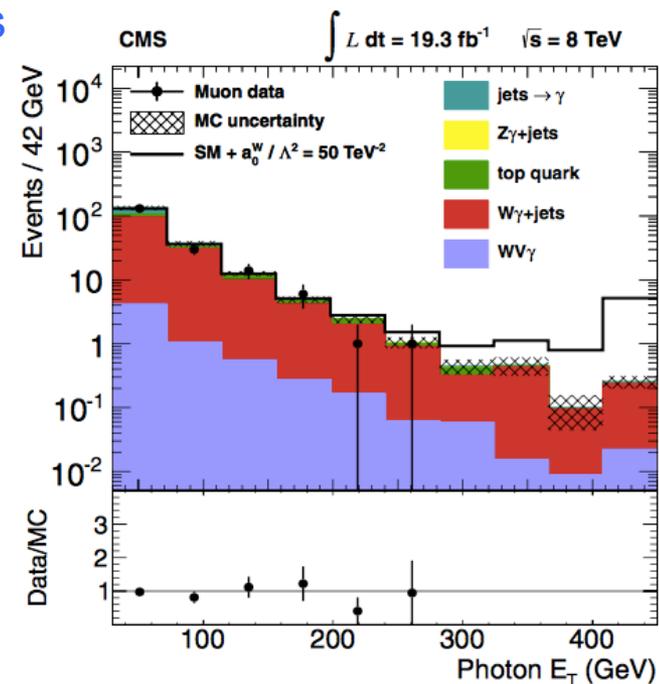
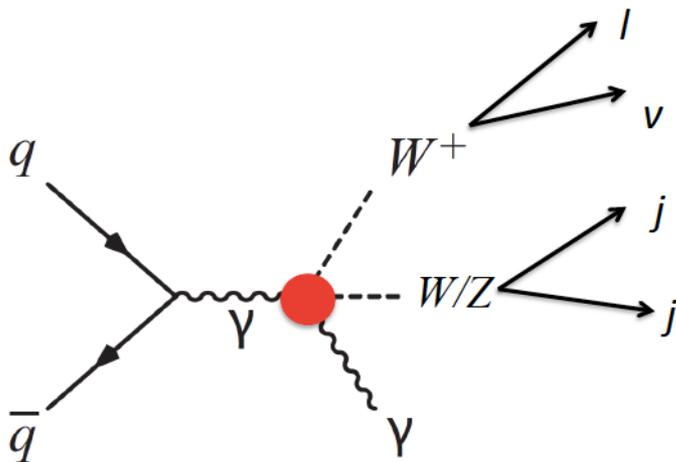
	Measured cross section (fb)	SM prediction (fb)
Inclusive $W\gamma\gamma$	$6.1 \pm 1.1(\text{stat}) \pm 1.2(\text{syst}) \pm 0.2(\text{lumi})$	2.90 ± 0.16
Exclusive $W\gamma\gamma$ With hard jet veto	$2.9 \pm 0.8(\text{stat}) \pm 1.0(\text{syst}) \pm 0.1(\text{lumi})$	1.88 ± 0.2

See Vasiliki's talk on Wednesday electroweak physics sections for more details

WV γ (CMS)

Phys. Rev. D 90, 032008 (2014)

- The selected data events is dominated by W γ +jets
 - not enough signal statistics for measurements
- 95% CL cross section upper limit is set at 311 fb
 - The limit is 3.4 times larger than SM predictions



Event selection highlight

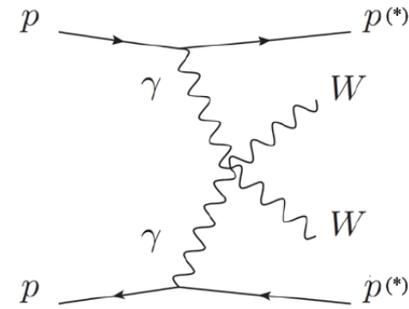
- One good lepton
- One good photon
- Two high pT jets
- $70\text{GeV} < M_{jj} < 100\text{GeV}$

Major BG in WV γ

- W γ +jets
- WV +jet , jet fake as γ
- Multijets

Exclusive $\gamma\gamma \rightarrow WW$ (CMS)

- Exclusive $\gamma\gamma$ production without color exchange
 - Very clean event signature (two tracks vertex)
 - Require No extra track on di-lepton vertex

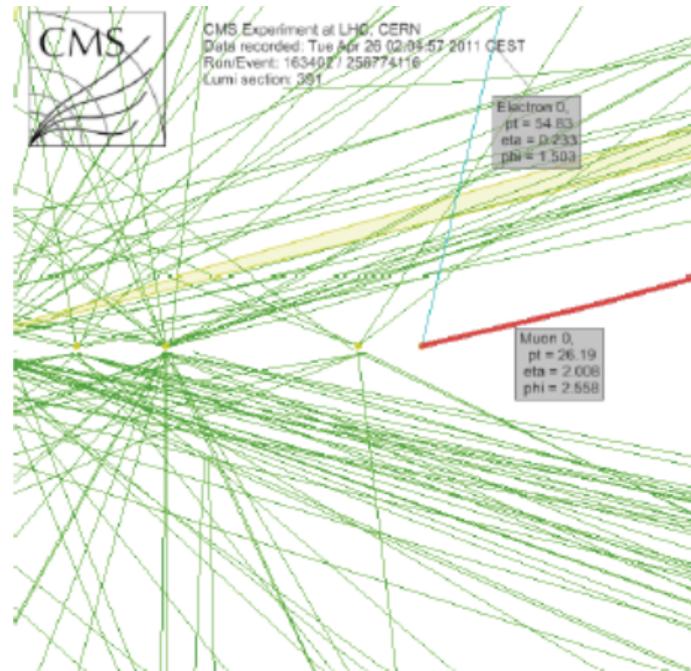
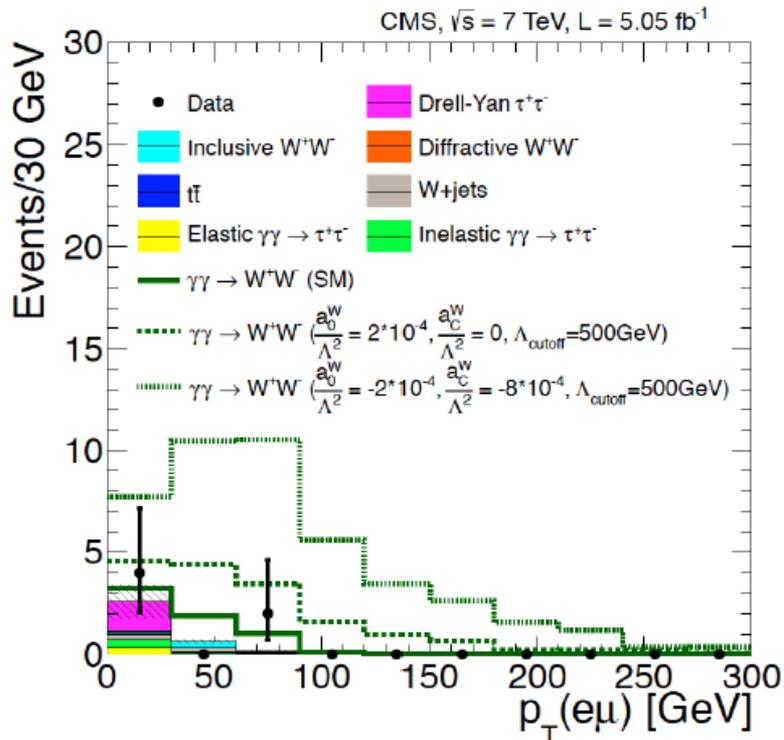


$$\sigma(pp \rightarrow p^{(*)}W^+W^-p^{(*)} \rightarrow p^{(*)}\mu^\pm e^\mp p^{(*)}) = 2.2^{+3.3}_{-2.0} \text{ fb}$$

1σ significance

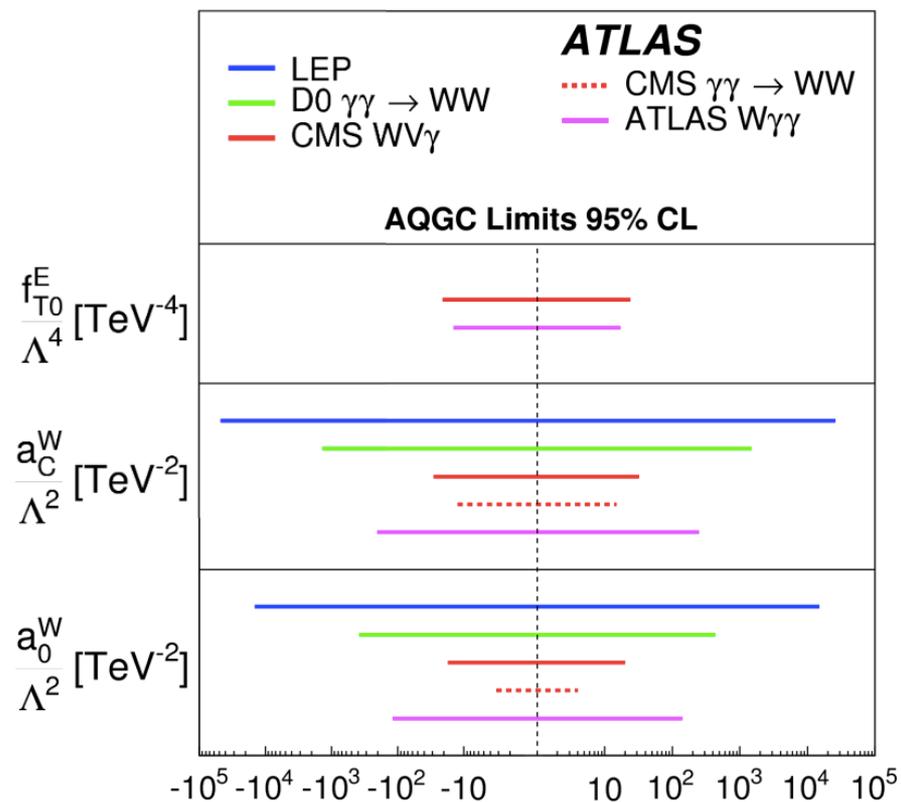
$$-4.0 \times 10^{-6} < a_0^W / \Lambda^2 < 4.0 \times 10^{-6} \text{ GeV}^{-2}$$

$$-1.5 \times 10^{-5} < a_C^W / \Lambda^2 < 1.5 \times 10^{-5} \text{ GeV}^{-2}$$



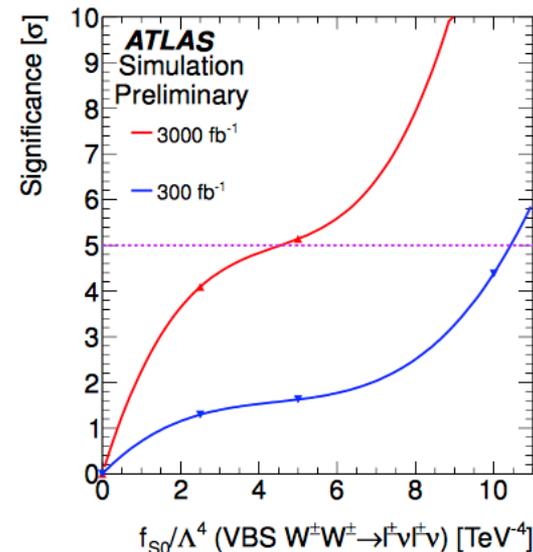
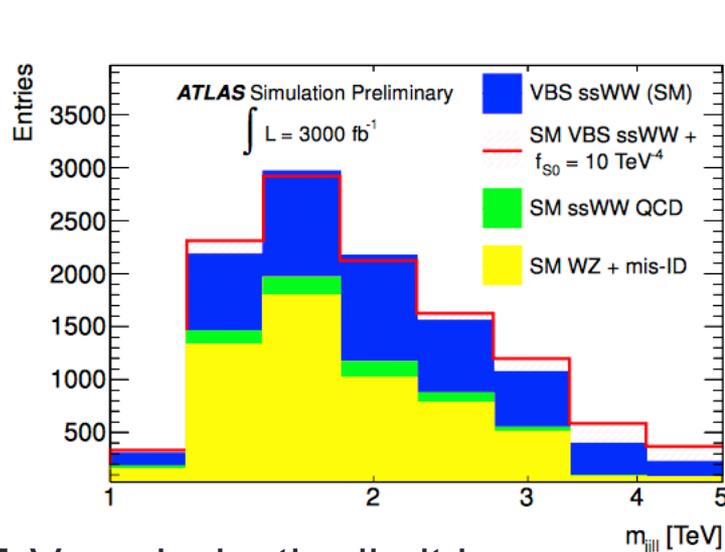
aQGCs limits from tri-boson and Exclusive $\gamma\gamma \rightarrow WW$ processes

- Tri-boson processes and $\gamma\gamma \rightarrow WW$ are sensitive to
 - F_T operators (Gauge boson field)
 - F_M operators (Higgs-Gauge field)
 - $F_{M,2}, F_{M,3}$ can be converted to LEP convention $a_C^W a_0^W$



aQGCs limits obtained by ATLAS/CMS are orders of magnitude more stringent than the best limits obtained at LEP.

Prospects for LHC upgrade



□ At 8TeV analysis, the limit is :

□ $f_S/\Lambda^4 \sim O(1000)$, $f_T/\Lambda^4 \sim O(10)$

□ aQGCs sensitivity is orders of magnitudes better in 13TeV and in HI-LHC upgrade

Channel	Parameter	(95% CL limits) 14TeV, 300 fb^{-1}	(95% CL limits) 14TeV, 3000 fb^{-1}
$W^\pm W^\pm jj$	$f_{S,0}/\Lambda^4 (\text{TeV}^{-4})$	$[-6.8, 6.8]$	$[-0.8, 0.8]$
$WZ jj$	$f_{T,1}/\Lambda^4 (\text{TeV}^{-4})$	$[-0.7, 0.7]$	$[-0.3, 0.3]$
$Z\gamma\gamma$	$f_{T,9}/\Lambda^4 (\text{TeV}^{-4})$	$[-0.9, 0.9]$	$[-0.3, 0.3]$

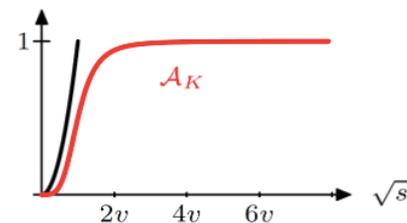
Summary

- Observation of electroweak Zjj production at LHC
 - A benchmark process for future studies of VBF processes at LHC
 - Background only hypothesis rejected at greater than 5σ in both ATLAS and CMS
- First evidence of tri-boson and Vector boson scattering production processes.
 - Contributions of aQGCs not observed
 - limits are set on aQGCs using Dimension 8 EFT models
 - Expect aQGCs limits will improve with the upcoming 13 TeV data by order of magnitudes.
 - limits with 8TeV 20 fb⁻¹ data: $f_S/\Lambda^4 \sim O(1000)$, $f_T/\Lambda^4 \sim O(10)$
 - expected limits with 300 fb⁻¹ data at 14TeV : $f_S/\Lambda^4 \sim O(10)$, $f_T/\Lambda^4 \sim O(1)$
 - expected limits with 3000 fb⁻¹ data at 14TeV: $f_S/\Lambda^4 \sim O(1)$, $f_T/\Lambda^4 \sim O(1)$

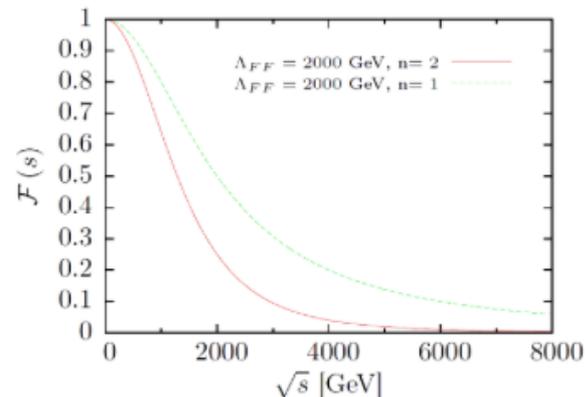
backup

Unitarity violation treatment

- EFT-predicted aTGC/aQGCs amplitudes
 - disrespect the gauge symmetry
 - and violate the unitarity once the \sqrt{s} goes sufficiently high
- Unitarization with the k-matrix approach (arxiv: 0806.4145)
 - Unitarization by infinitely heavy and wide resonance
 - K-matrix amplitude $|\mathcal{A}_K(s)|^2 \xrightarrow{s \rightarrow \infty} 1$
- Form factor approach (arxiv:1205.4231)
 - Unitarity can be preserved by introducing form-factor (FF)
 - energy-dependent form factors



$$\mathcal{F}(s) = \frac{1}{\left(1 + \frac{s}{\Lambda_{FF}^2}\right)^n}$$



W $\gamma\gamma$ (ATLAS)

- MCFM is used for NLO W $\gamma\gamma$ SM predicted cross section.
 - The measured cross section in inclusive case is 1.9 σ higher than predictions
 - Better agreement in exclusive case
- Exclusive measured cross section is used for aQGCs study

Definition of the fiducial region

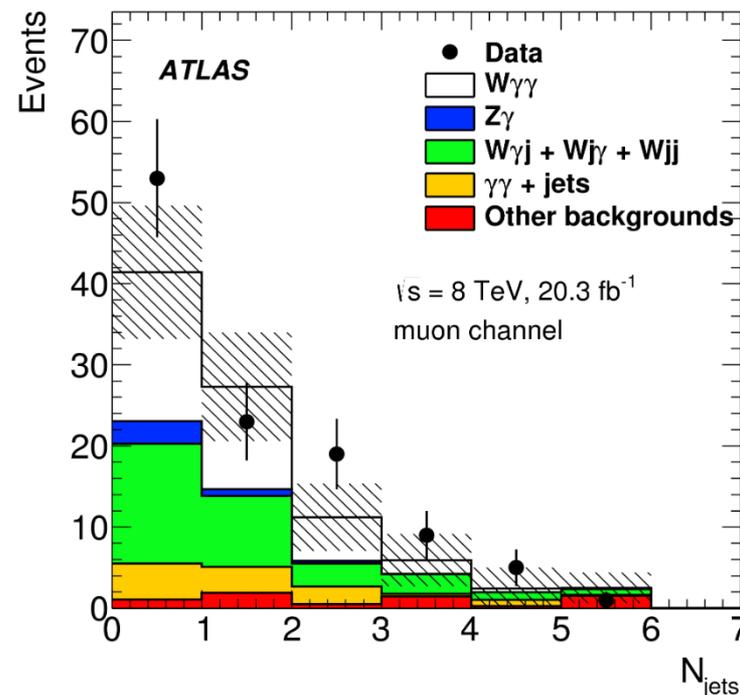
$$p_T^\ell > 20 \text{ GeV}, p_T^\nu > 25 \text{ GeV}, |\eta_\ell| < 2.5$$

$$m_T > 40 \text{ GeV}$$

$$E_T^\gamma > 20 \text{ GeV}, |\eta^\gamma| < 2.37, \text{ iso. fraction } \epsilon_h^p < 0.5$$

$$\Delta R(\ell, \gamma) > 0.7, \Delta R(\gamma, \gamma) > 0.4, \Delta R(\ell/\gamma, \text{jet}) > 0.3$$

$$\text{Exclusive: no anti-}k_t \text{ jets with } p_T^{\text{jet}} > 30 \text{ GeV}, |\eta^{\text{jet}}| < 4.4$$



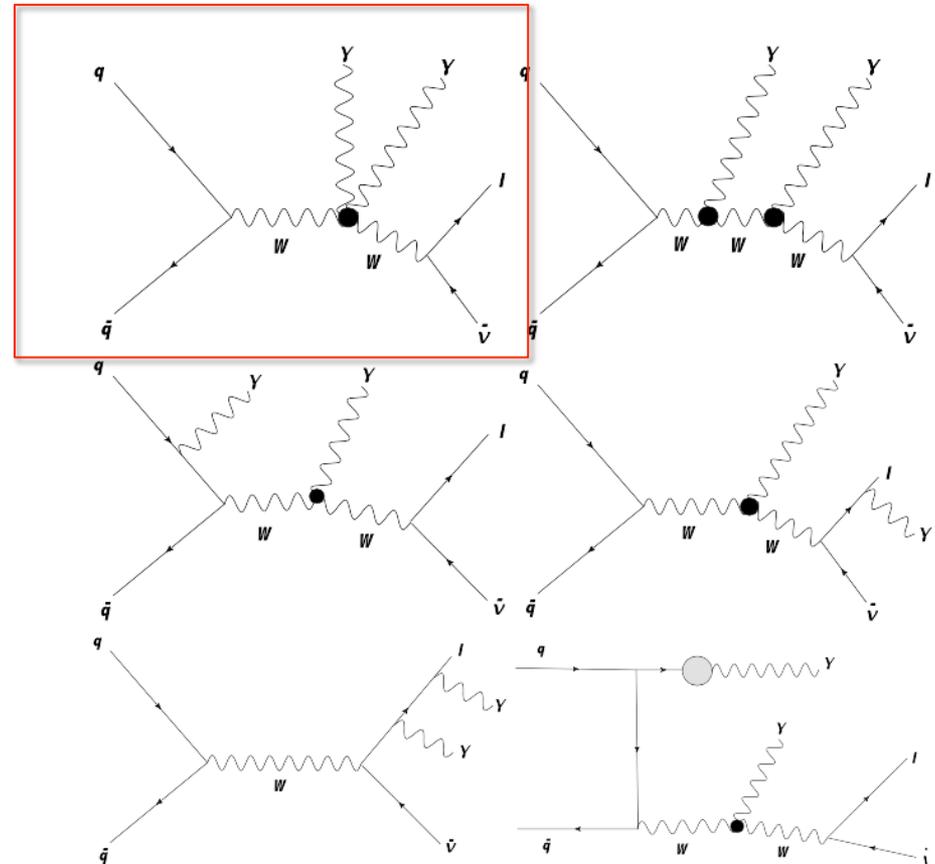
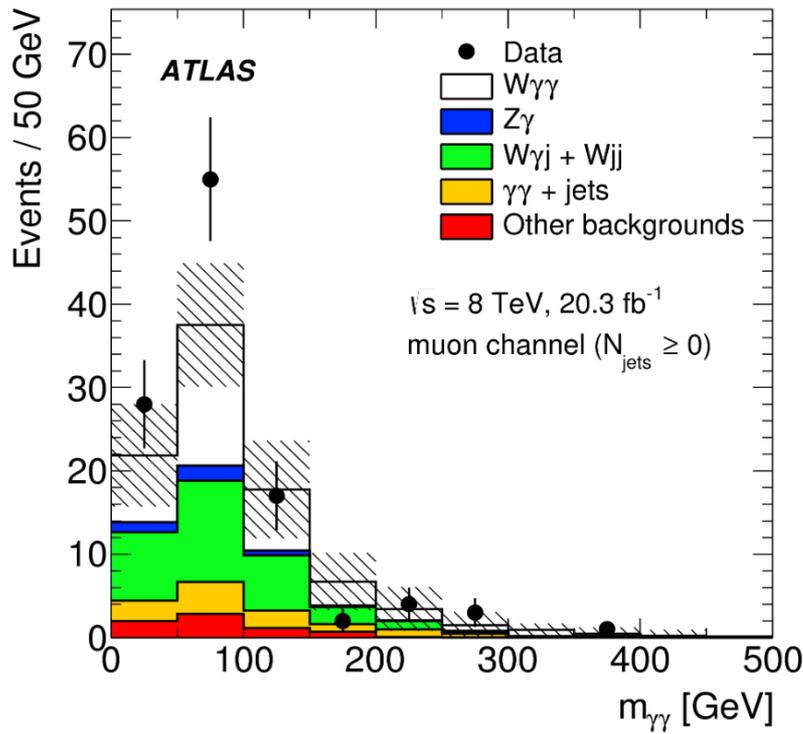
More than 3 σ significance

	Measured cross section (fb)	SM prediction (fb)
Inclusive W $\gamma\gamma$	$6.1 \pm 1.1(\text{stat}) \pm 1.2(\text{syst}) \pm 0.2$ (lumi)	2.90 ± 0.16
Exclusive W $\gamma\gamma$	$2.9 \pm 0.8(\text{stat}) \pm 1.0(\text{syst}) \pm 0.1$ (lumi)	1.88 ± 0.2

$W\gamma\gamma$ (ATLAS)

arXiv:1503.03243

- $W\gamma\gamma$ process is sensitive to $WW\gamma\gamma$ and $WZ\gamma\gamma$ QGCs Vertex.
- Other contributions from:
 - ISR photons
 - FSR photons
 - Photon from TGC vertex
 - Photon from jet fragmentations



Definition of the fiducial region

$$p_{\text{T}}^{\ell} > 20 \text{ GeV}, p_{\text{T}}^{\nu} > 25 \text{ GeV}, |\eta_{\ell}| < 2.5$$

$$m_{\text{T}} > 40 \text{ GeV}$$

$$E_{\text{T}}^{\gamma} > 20 \text{ GeV}, |\eta^{\gamma}| < 2.37, \text{ iso. fraction } \epsilon_{\text{h}}^{\text{P}} < 0.5$$

$$\Delta R(\ell, \gamma) > 0.7, \Delta R(\gamma, \gamma) > 0.4, \Delta R(\ell/\gamma, \text{jet}) > 0.3$$

$$\text{Exclusive: no anti-}k_{\text{t}} \text{ jets with } p_{\text{T}}^{\text{jet}} > 30 \text{ GeV}, |\eta^{\text{jet}}| < 4.4$$