

Studies of non-irradiated and irradiated HV-CMOS detectors

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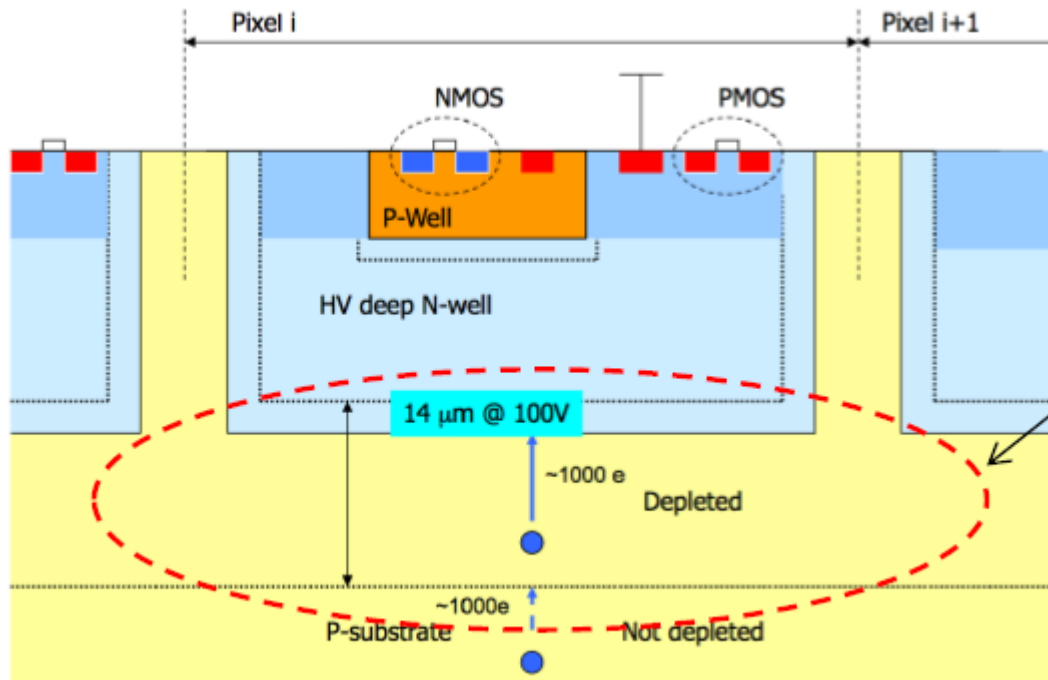
work done in the framework of ATLAS CMOS pixels
and ATLAS CMOS strips collaborations

Motivation

- HV-CMOS technology is a very interesting option for detectors at upgraded LHC
 - reduced cost
 - improved granularity and spatial resolution
- essential to understand charge collection properties especially after irradiation and at LHC readout speed
 - very thin depleted layers compared to usual detectors
 - low resistivity silicon compared to usual detector material
- Edge-TCT and CCE measurements with ^{90}Sr are ideal tools for such studies
- charge collection studies with HV-CMOS sensors (two types of samples CCPDv2, CHES1) before and after irradiation will be presented here

HV-CMOS detectors

- charge collecting electrode n-well in p-type substrate, resistivity $\sim 20 \Omega\text{cm}$
- depletion layer $\sim 14 \mu\text{m}$ thick at 100 V bias
- CMOS circuitry is implemented in the n-well



Interested in charge collection from n-well - substrate junction

Daniel Muenstermann | TWEPP 2012 | Oxford | September 19th, 2012

Samples

1) **CHES1 chip (ATLAS CMOS strips studies):**

350 nm AMS, 20 Ωcm , 120 V max bias

Passive pixels: no amplifier in n-well, n-well connected directly to readout (as in standard diode)

(see. Zhijun's talk earlier today)

E-TCT

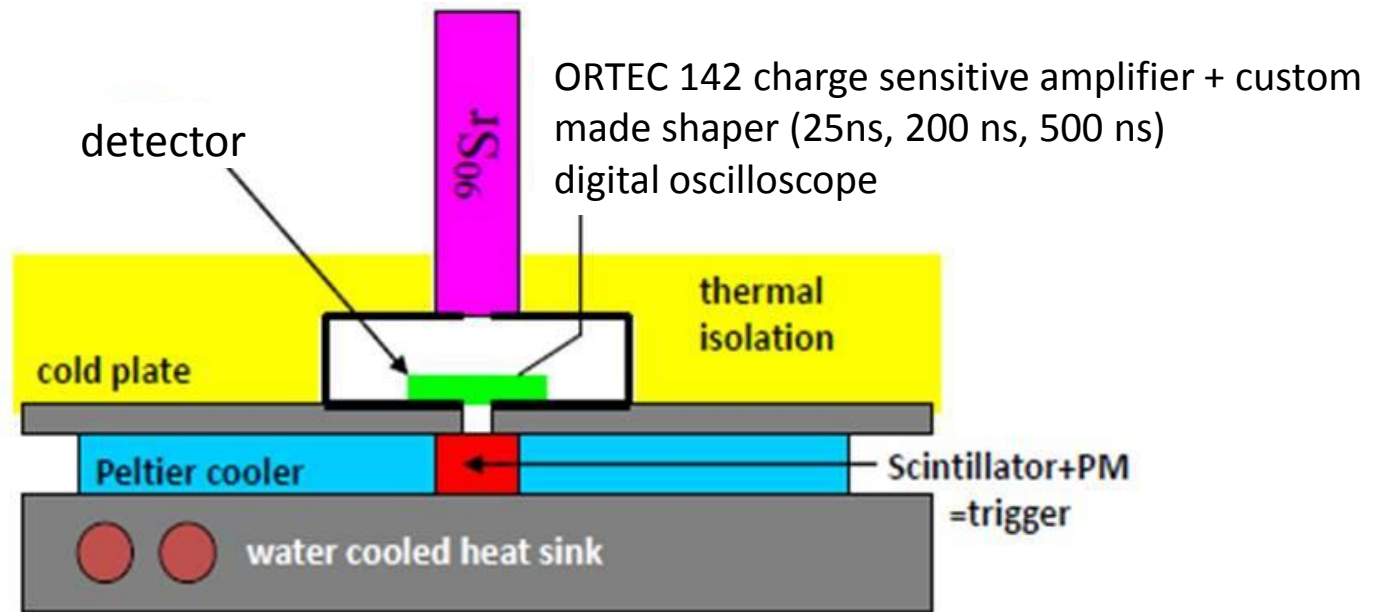
- 100 x 45 μm^2 pixels
- passive pixel
- ➔ induced current directly observed on the scope



CCE measurement with MIPs from Sr^{90} source

- 2 x 2 mm^2 total area
- 880 passive pixels (45 x 200 μm^2) tied together
- ➔ Large structure needed for good collimation and sufficient event rate

MIP CCE measurement setup



- HV-CMOS: thin \rightarrow small signals, large noise \rightarrow S/N bad
 - \rightarrow must have clean sample
 - \rightarrow need large detector for reasonable trigger rate and good collimation, small scintillator

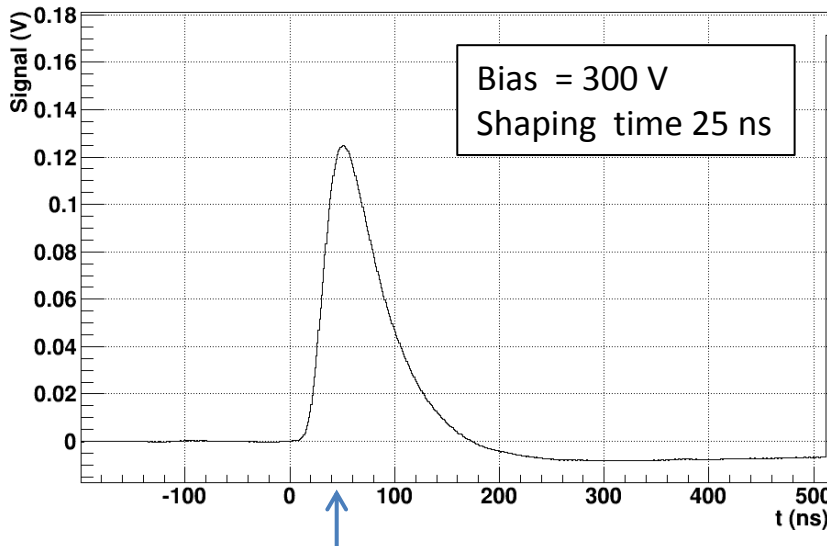
MIP CCE measurement

Calibration

- Sr^{90} with 300 μm thick FZ p-type silicon diode, $V_{fd} = 70 \text{ V}$
→ calibration at 25 ns shaping time : **230 electrons (mean)/mV**
- confirmed with Am241 source, 59.5 keV photon peak in 300 μm thick detector

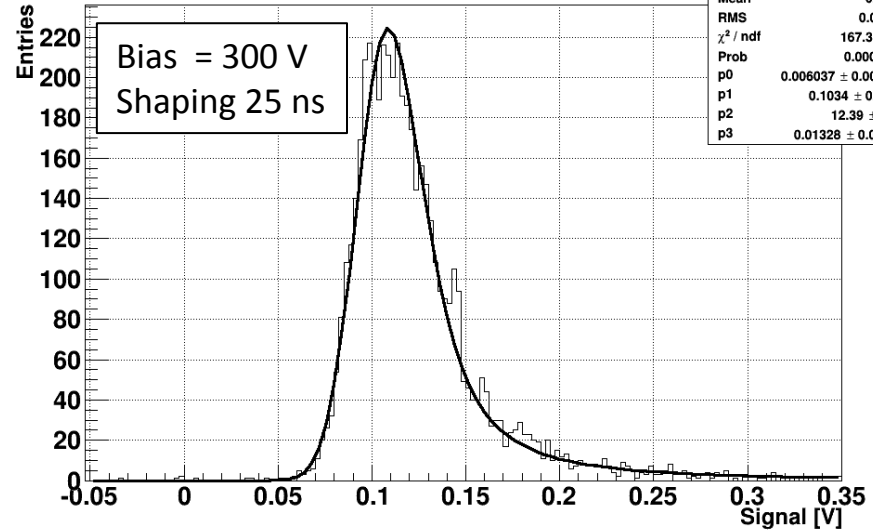
FZ p-type silicon diode, 300 μm thick , $V_{fd} = 70 \text{ V}$:

Averaged 5000 waveforms



Sampling time

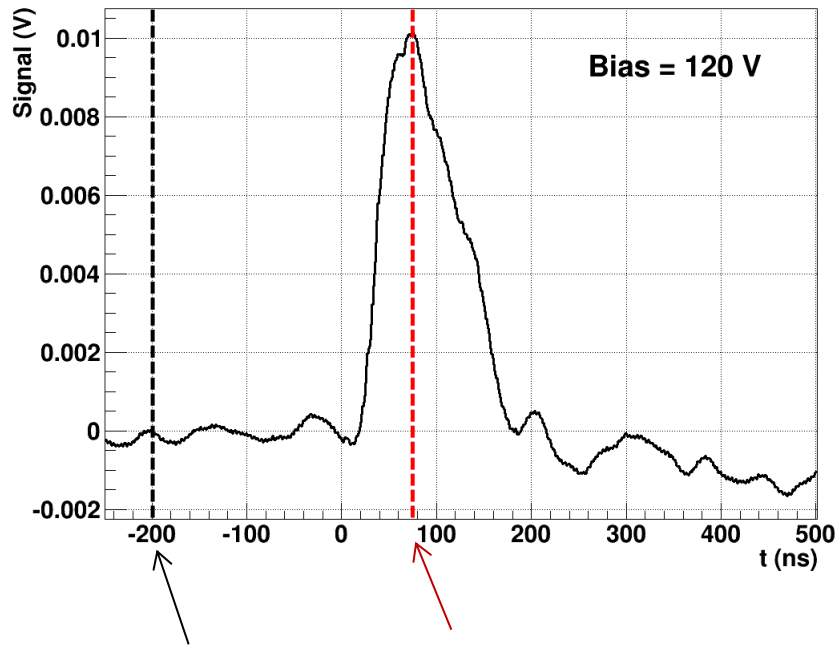
Spectrum of amplitudes at sampling time



MIP CCE measurement

CHES1 large passive HV-CMOS array

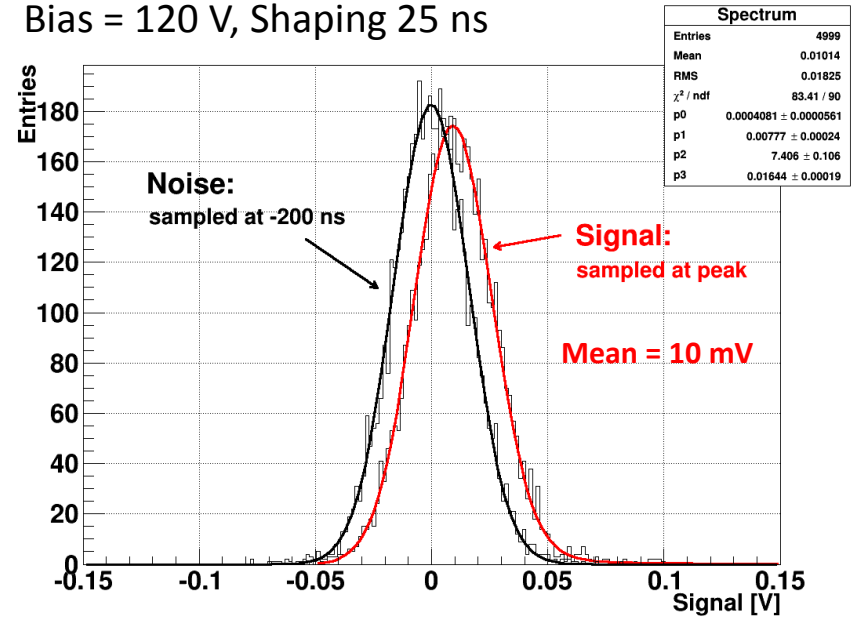
Averaged waveform (5000 samples)



Noise sampling point

Signal sampling point

Bias = 120 V, Shaping 25 ns

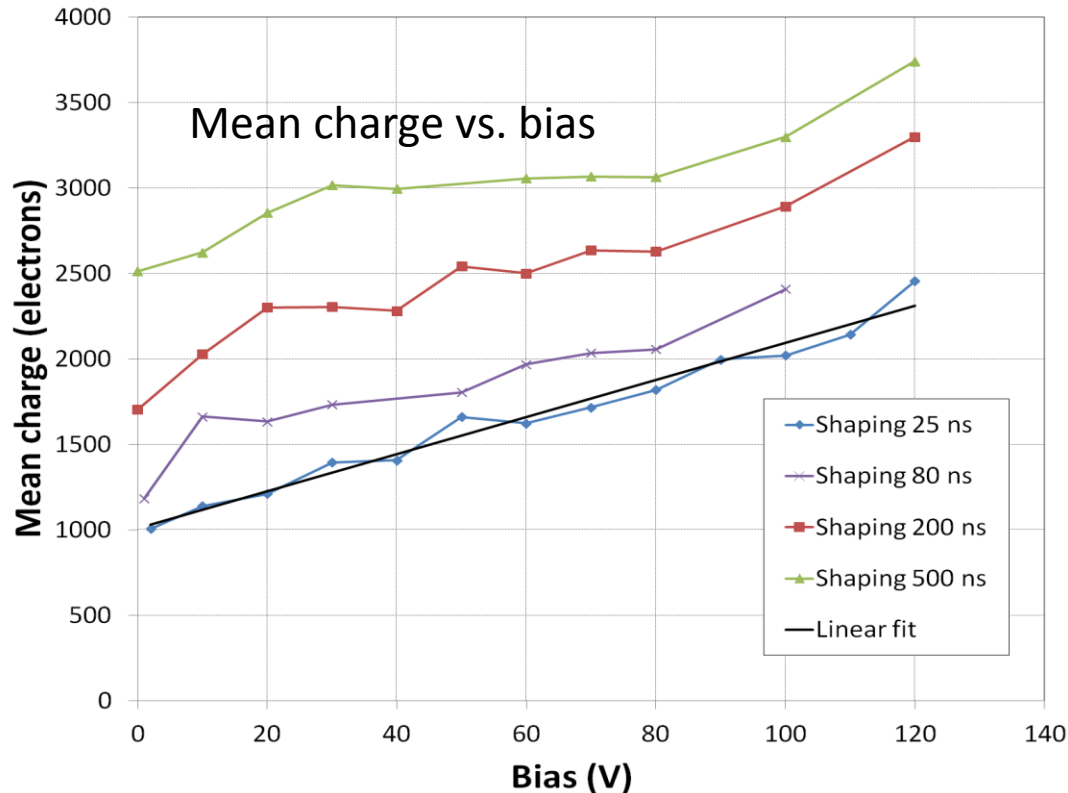


Signal spectrum mean: 10 mV

→ mean charge: 2300 electrons

MIP CCE measurement

CHES1 large passive HV-CMOS array



25 ns shaping: **Mean charge = 1010 el + 11 el/V**

• charge increases with shaping time → diffusion

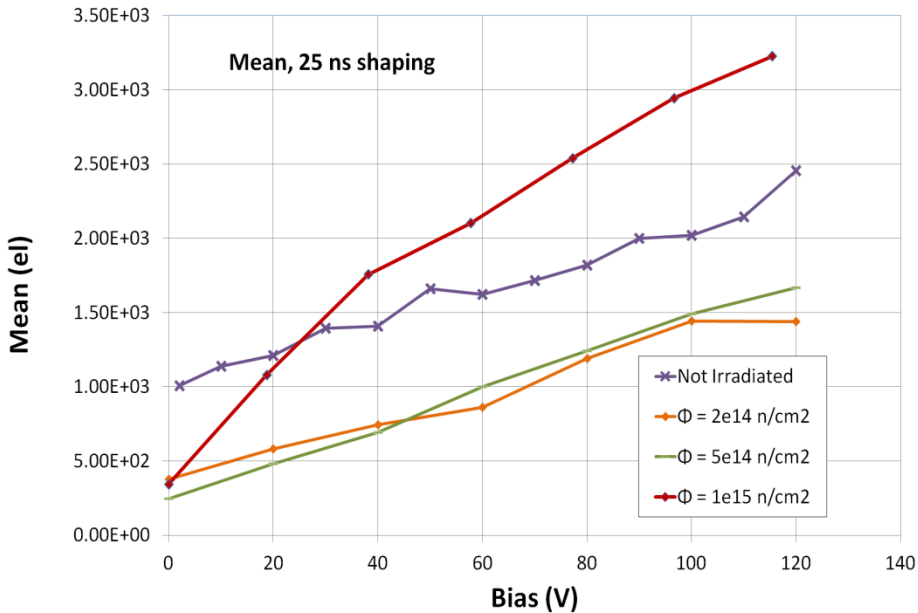
• charge ~ 1000 el at 0 V due to diffusion

→ if at 100 V depleted layer 14 μm thick: ~ **1400 el** expected from depleted layer:
measured: **2100 el** → ~ **700 el** from diffusion (~ consistent)

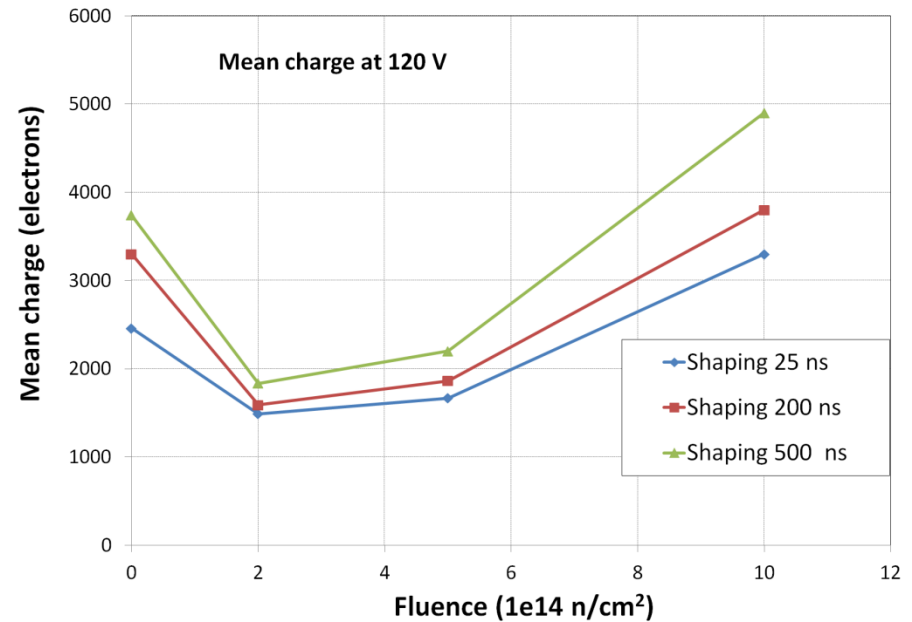
MIP CCE measurement

CHES1 large passive HV-CMOS array – irradiated with neutrons in reactor in Ljubljana

Charge vs. bias

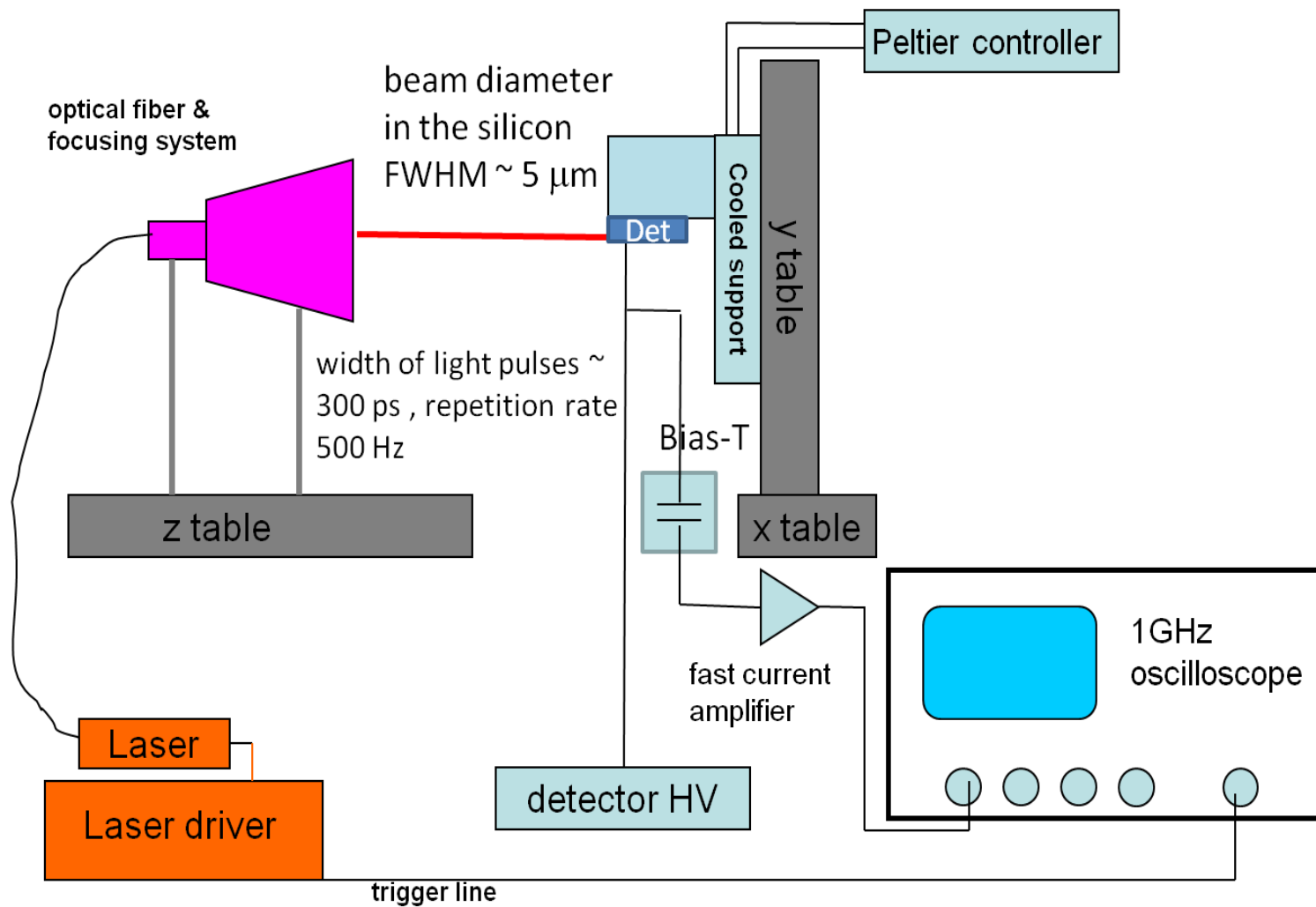


Charge at 120 V vs. fluence



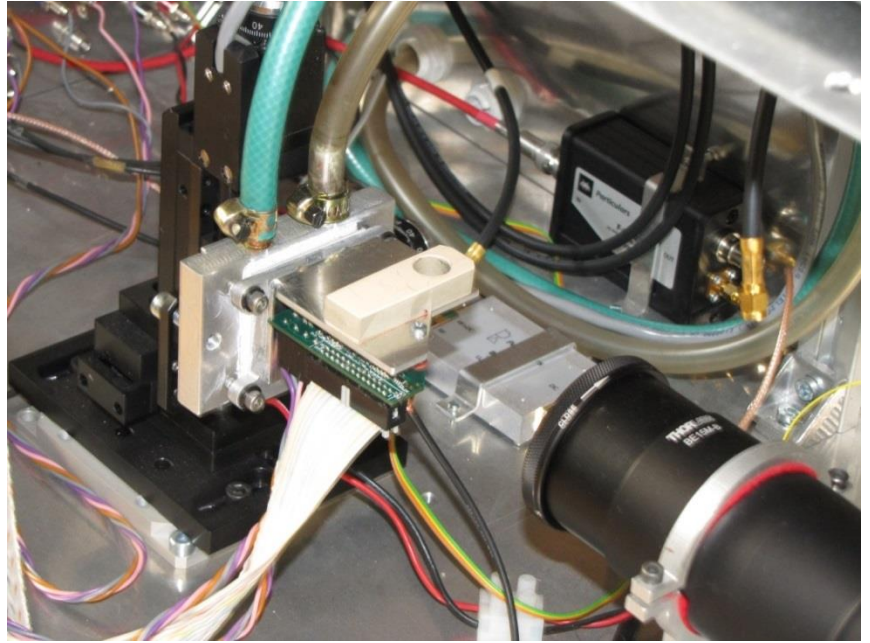
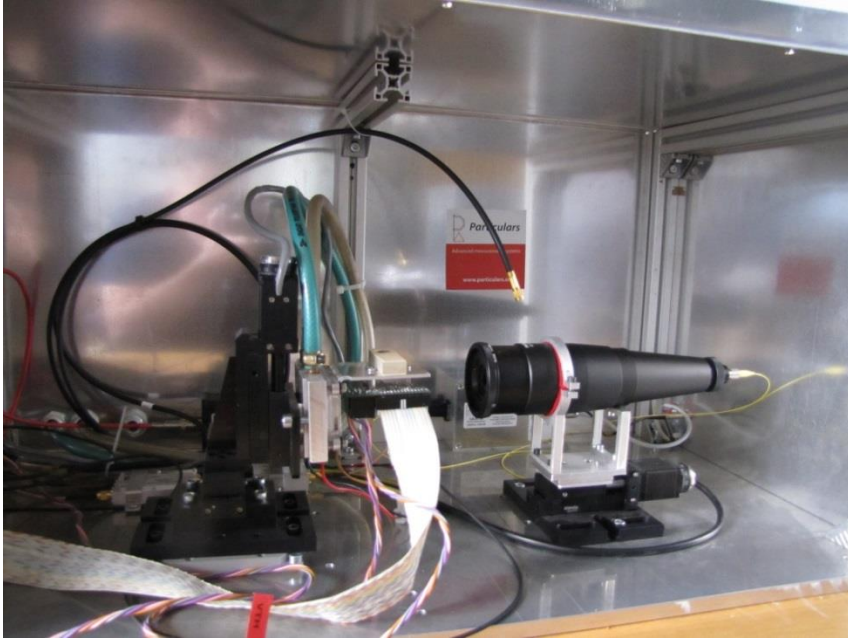
- charge drops after first irradiation step, small influence of shaping time: **small diffusion contribution**
- charge increases with more irradiation → **depleted region increases due to acceptor removal**
→ see talk by G. Kramberger
- larger charge at 500 ns shaping at 1e15 n/cm² : detrapping?
→ $\tau_{d1} \sim 500$ ns for holes measured (however for 20% of the trapped holes),
see: [G. Kramberger et al., 2012 JINST 7 P04006](#)

Edge TCT



Edge-TCT

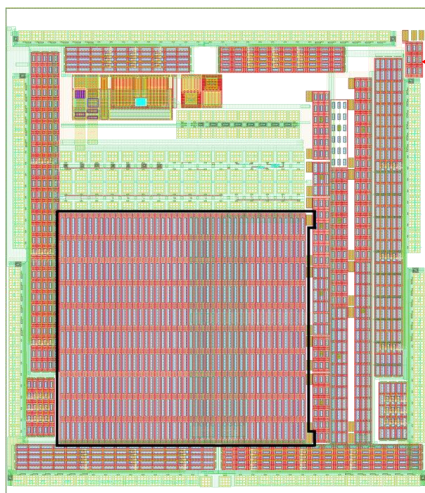
Photos of the setup (more details: www.particlars.si)



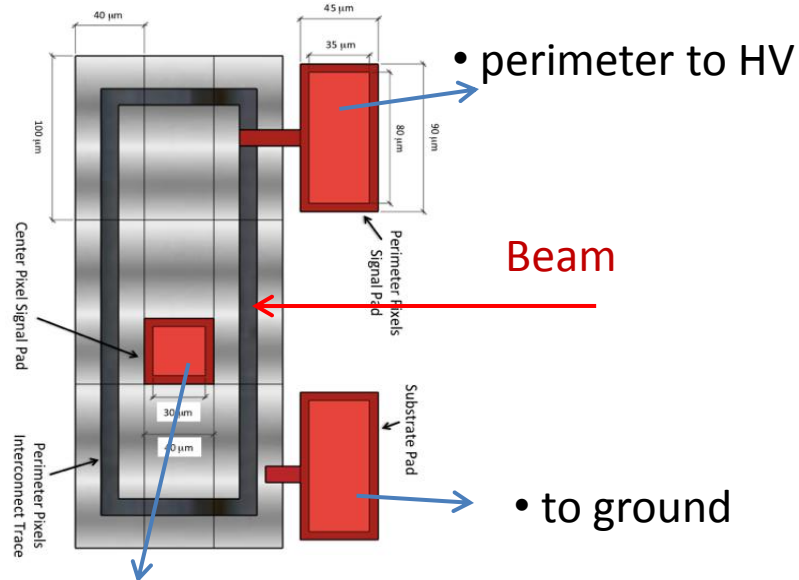
Edge-TCT

Chess1

- passive pixel in the corner

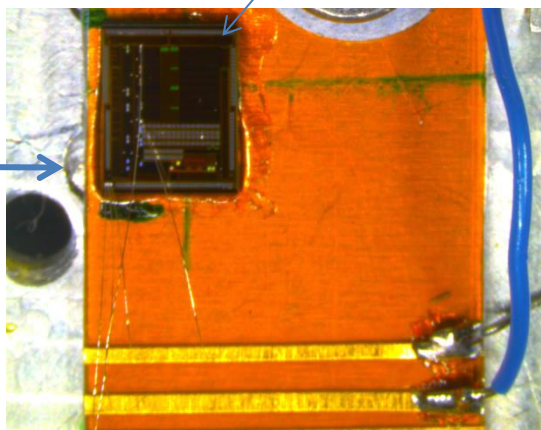


Beam direction



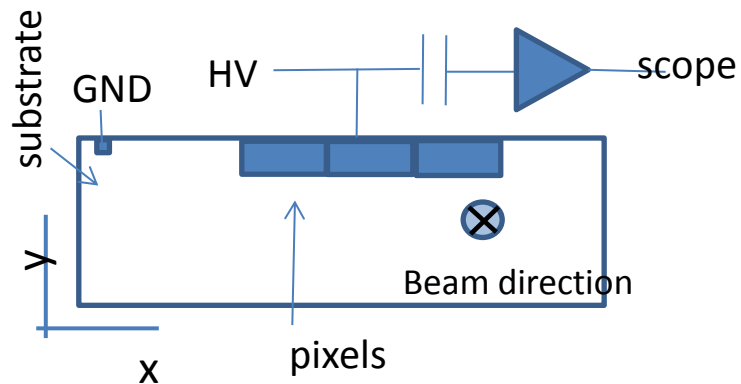
- signal to high voltage and readout (via Bias-T)

CHES1



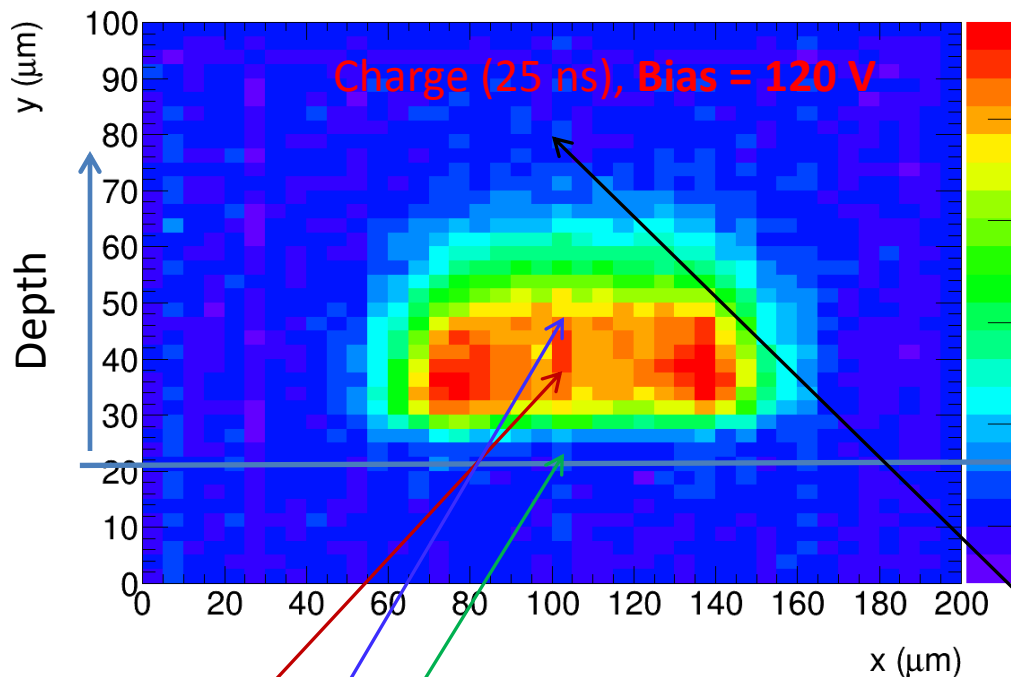
Beam direction

Detector connection scheme:



Edge-TCT

Chess1, not irradiated, pixel 100 μm x 45 μm



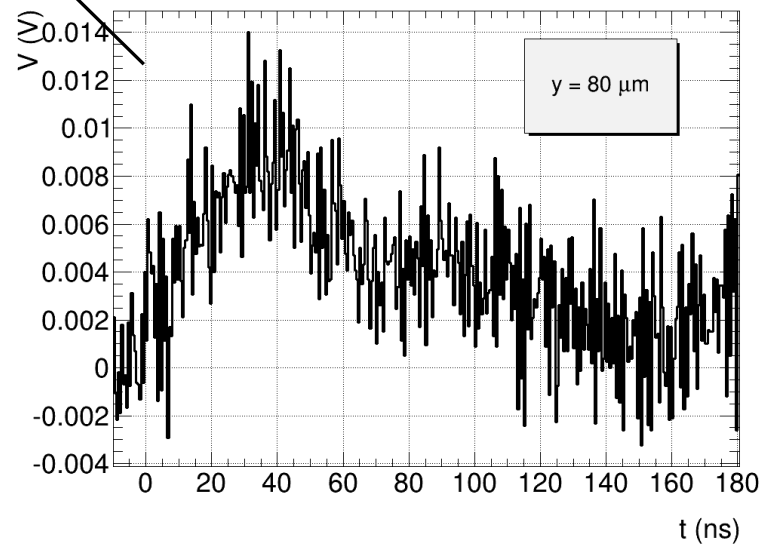
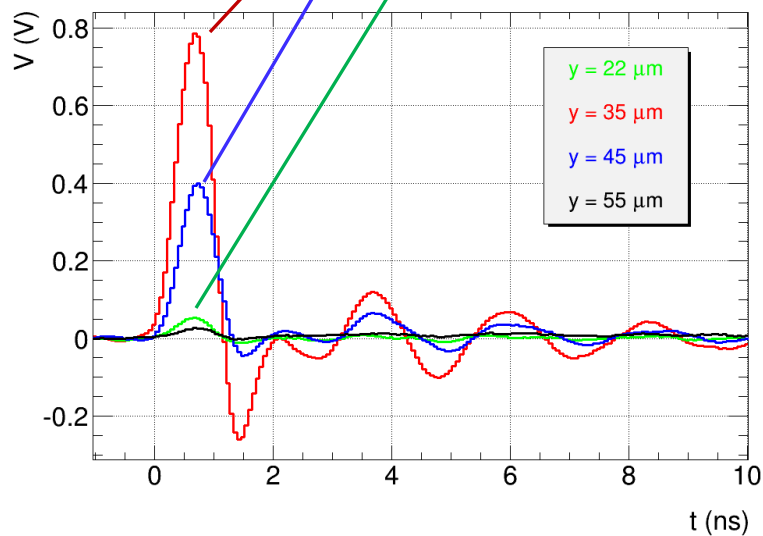
Scan across pixel:

- 2.5 μm steps in y
- 5 μm steps in x



Beam direction

Chip surface

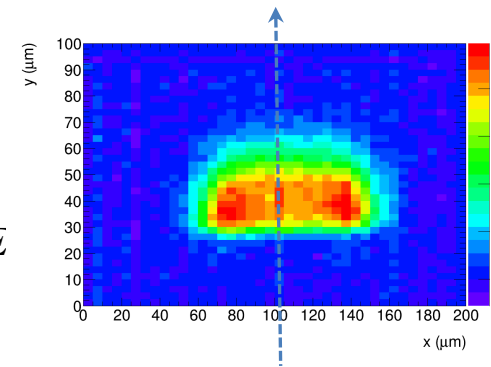
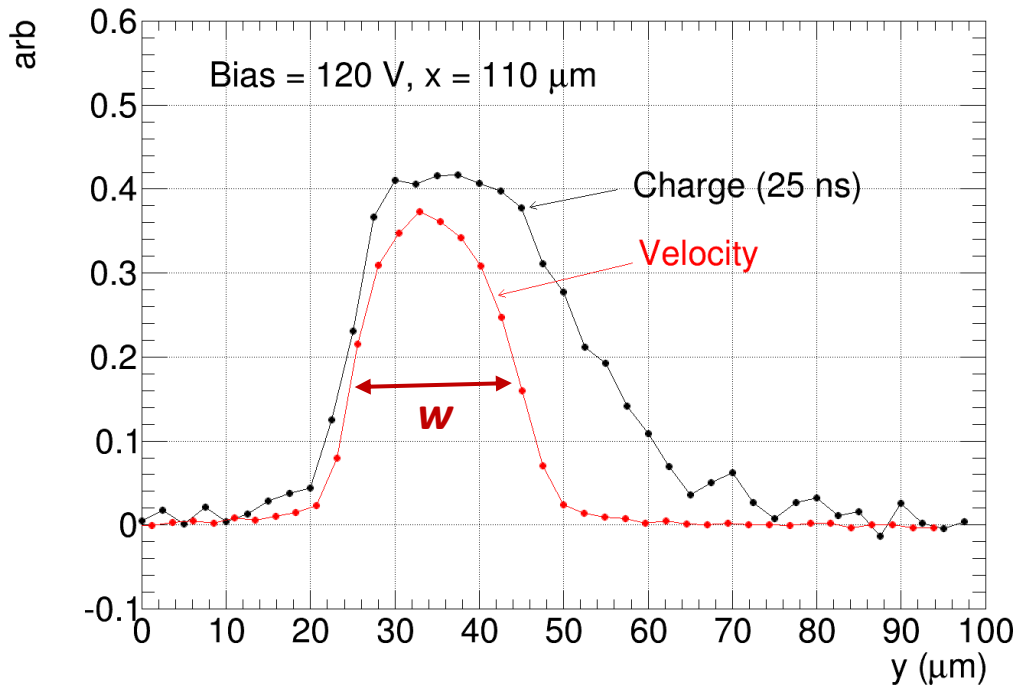


Edge-TCT

Chess1, not irradiated

- 1) **charge**: integral of induced current pulse
- 2) **velocity** (in E-TCT): induced current immediately after the laser pulse

$$I(x, y, t \sim 0) \approx qE_w(x, y) [\bar{v}_e(x, y) + \bar{v}_h(x, y)]; \quad \bar{v}_e(x, y) + \bar{v}_h(x, y) \propto E$$

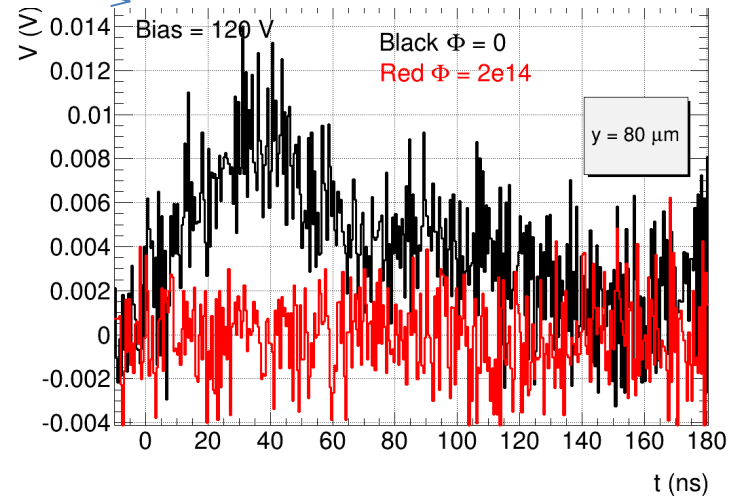
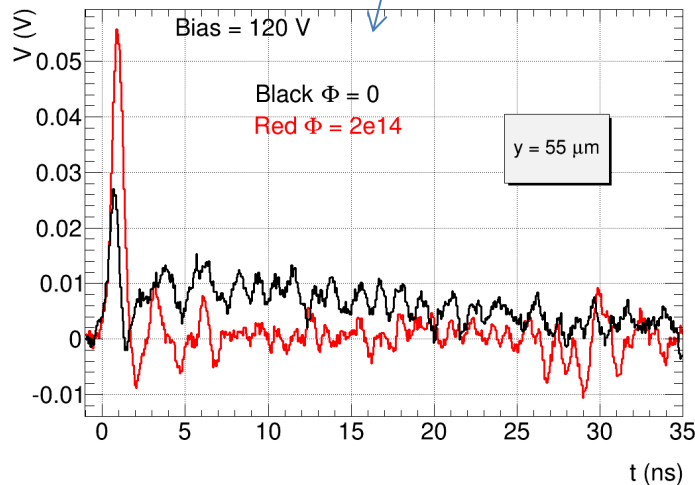
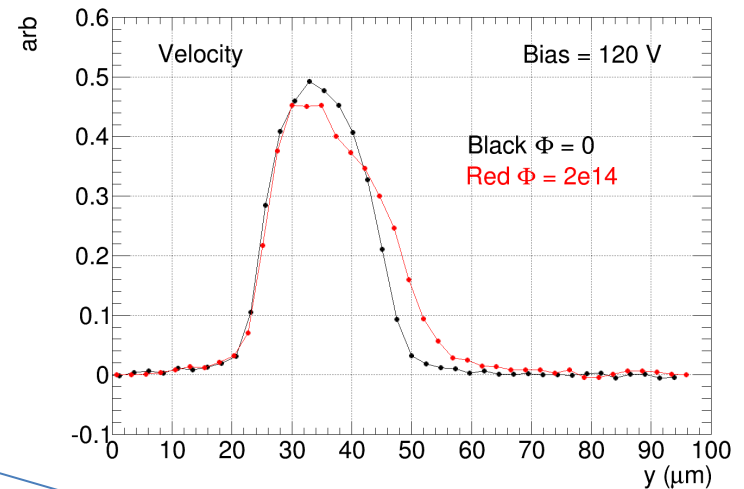
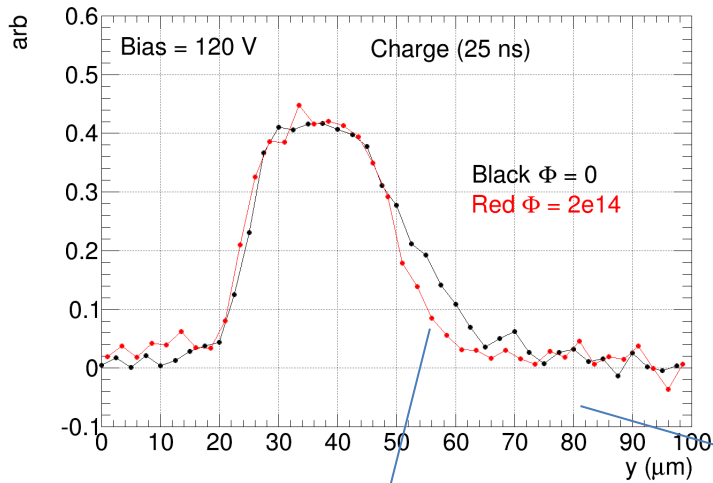


- **high velocity** \sim **depleted region** ~ 20 μm \rightarrow about 60% charge within this region
- total charge collection region wider (diffusion) \sim consistent with 2400 el measured with Sr-90
 - \rightarrow take into account laser beam width

Edge-TCT

Chess1, irradiated with $2e14 \text{ n/cm}^2$

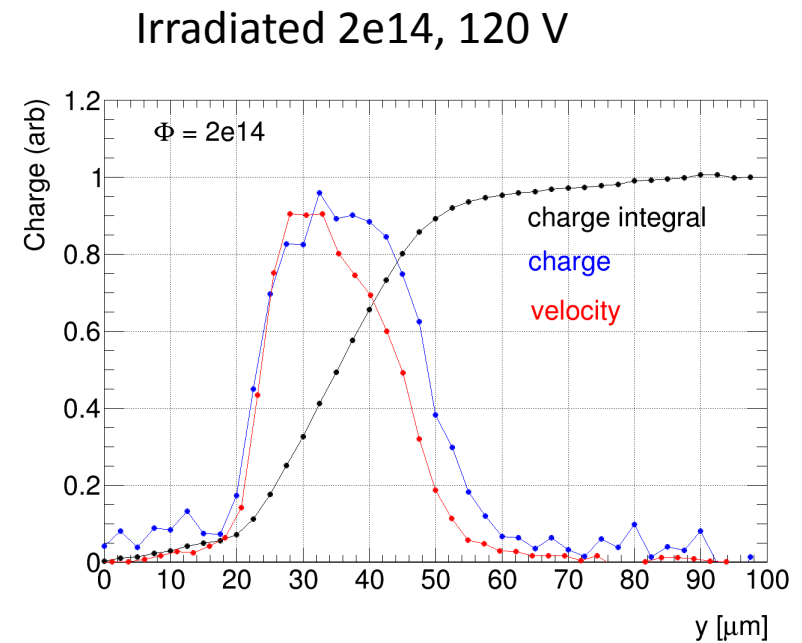
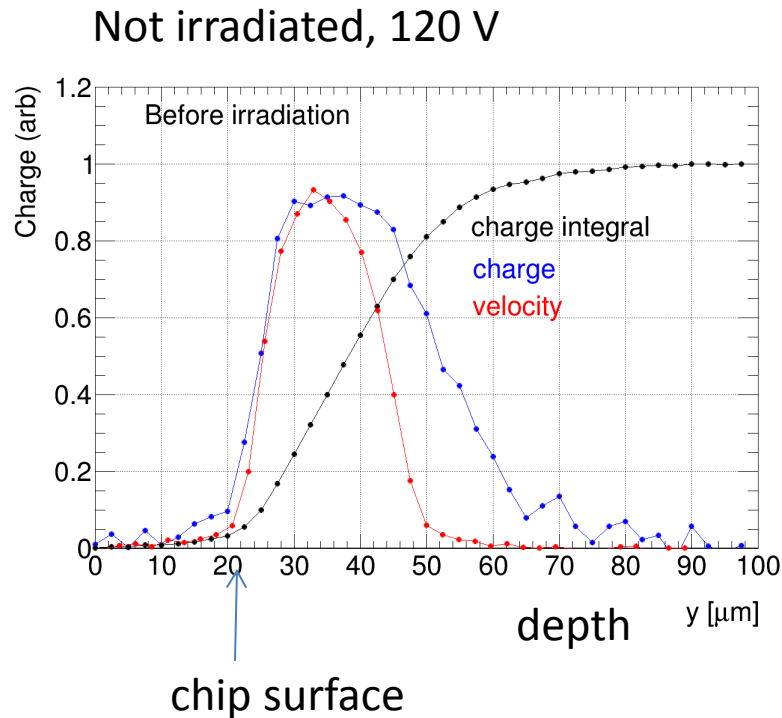
- charge collection region narrower
- field region (velocity) increases \rightarrow acceptor removal
- no long tails of induced current pulses \rightarrow less diffusion



Edge-TCT

Chess1, irradiated with $2e14$ n/cm²

Charge collection and velocity profiles across pixel centre



After irradiation:

→ smaller difference between charge collection and velocity (depleted) region

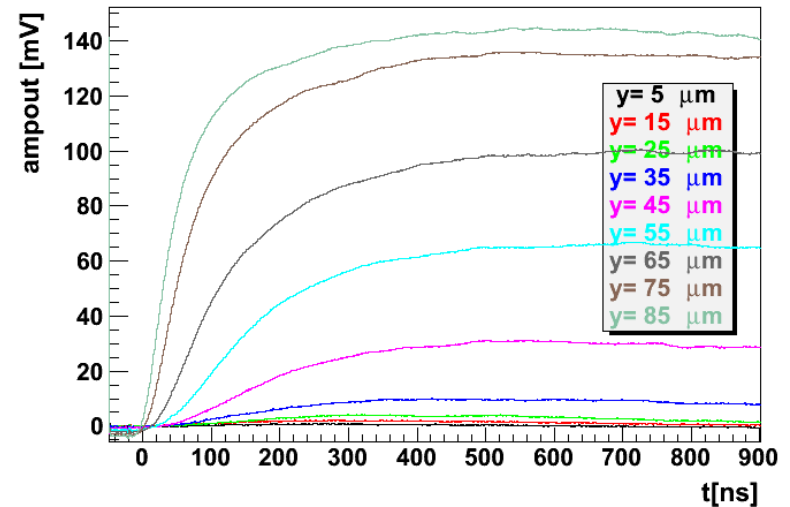
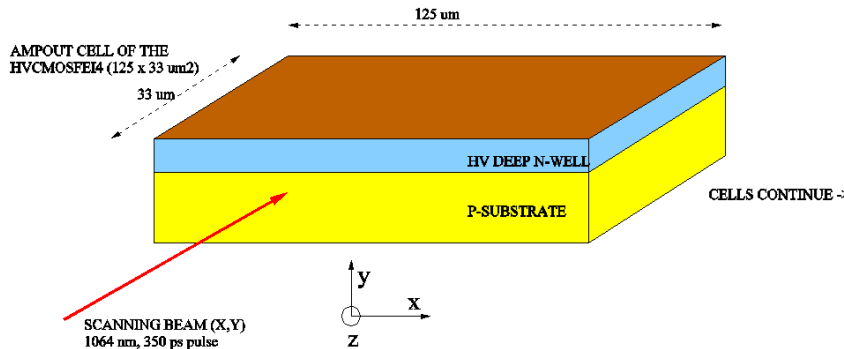
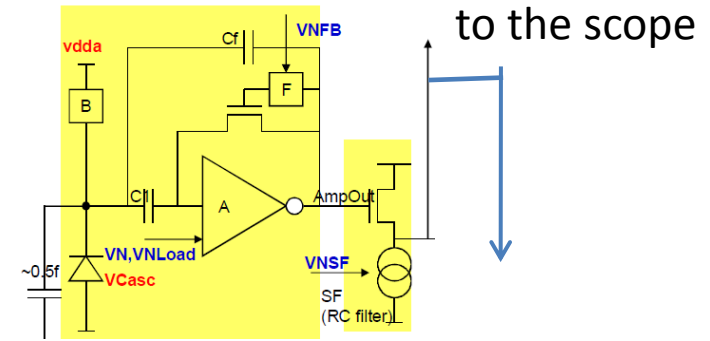
→ charge within depleted region: before irradiation $\sim 60\%$ after irradiation $\sim 90\%$

Samples

2) CCPDv2 chip (ATLAS CMOS pixels studies): 180 nm, AMS, $\sim 20 \Omega\text{cm}$, 60 V max bias,
Active pixels: output of the amplifier monitored on the scope

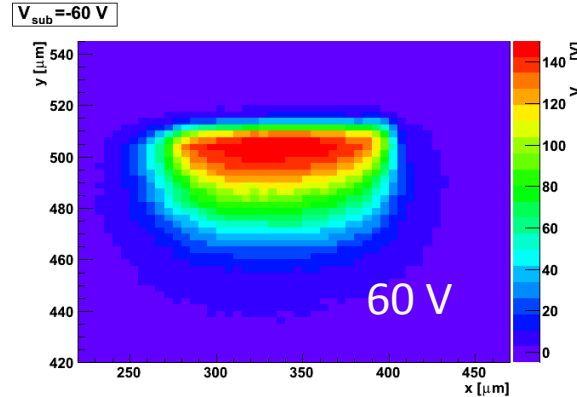
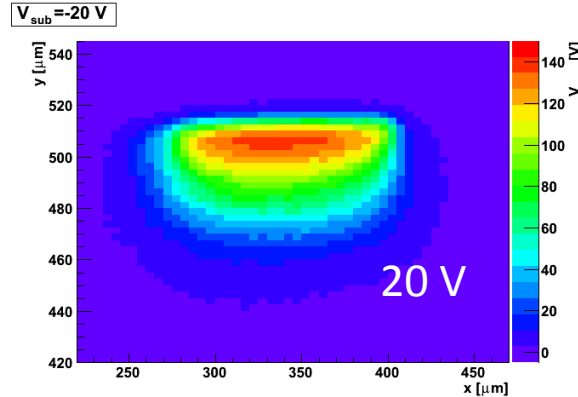
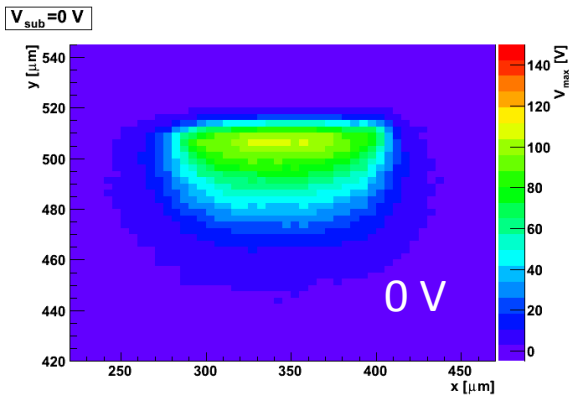
E-TCT on single cell, $125 \times 33 \mu\text{m}^2$
readout after the charge sensitive
amplifier (not observing induced current)

Single cell charge
sensitive amplifier:

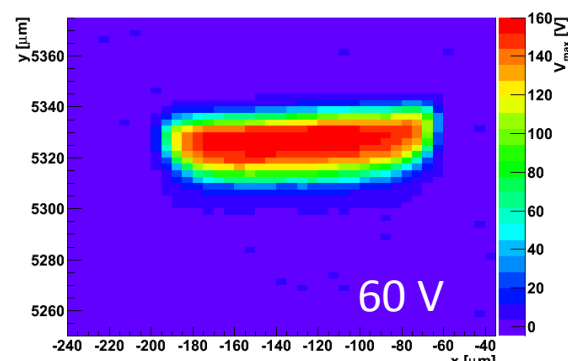
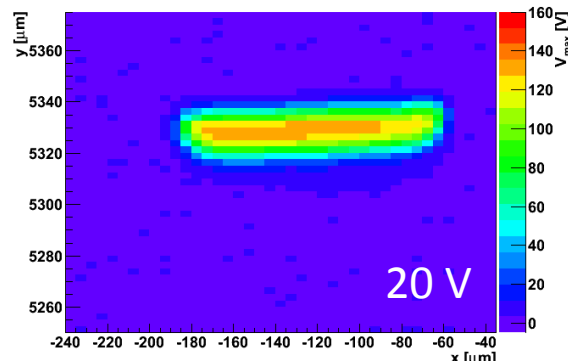
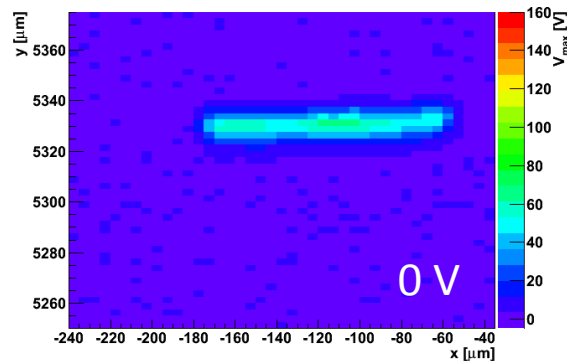


Edge-TCT, HV2FEI4, pixel 125 $\mu\text{m} \times 33 \mu\text{m}$, irradiated with neutrons in Ljubljana

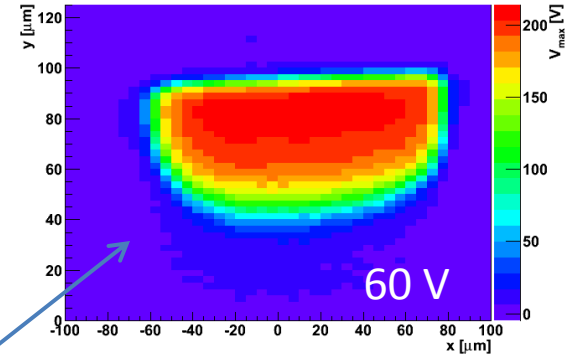
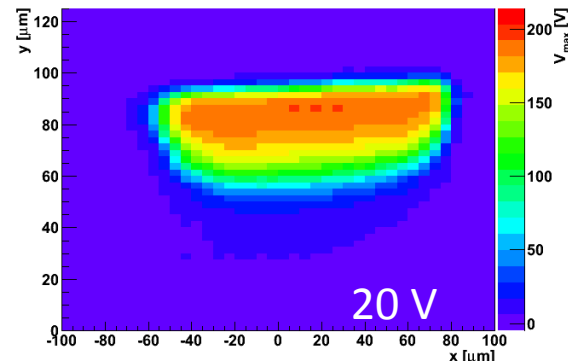
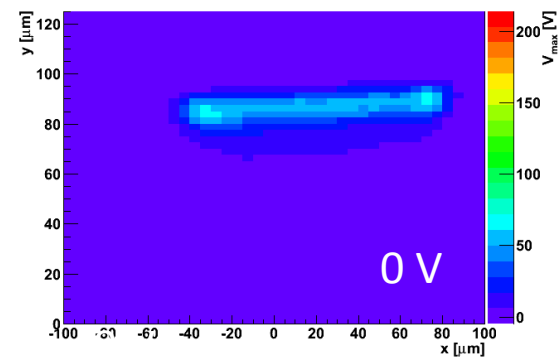
not irradiated



5e14 cm⁻²

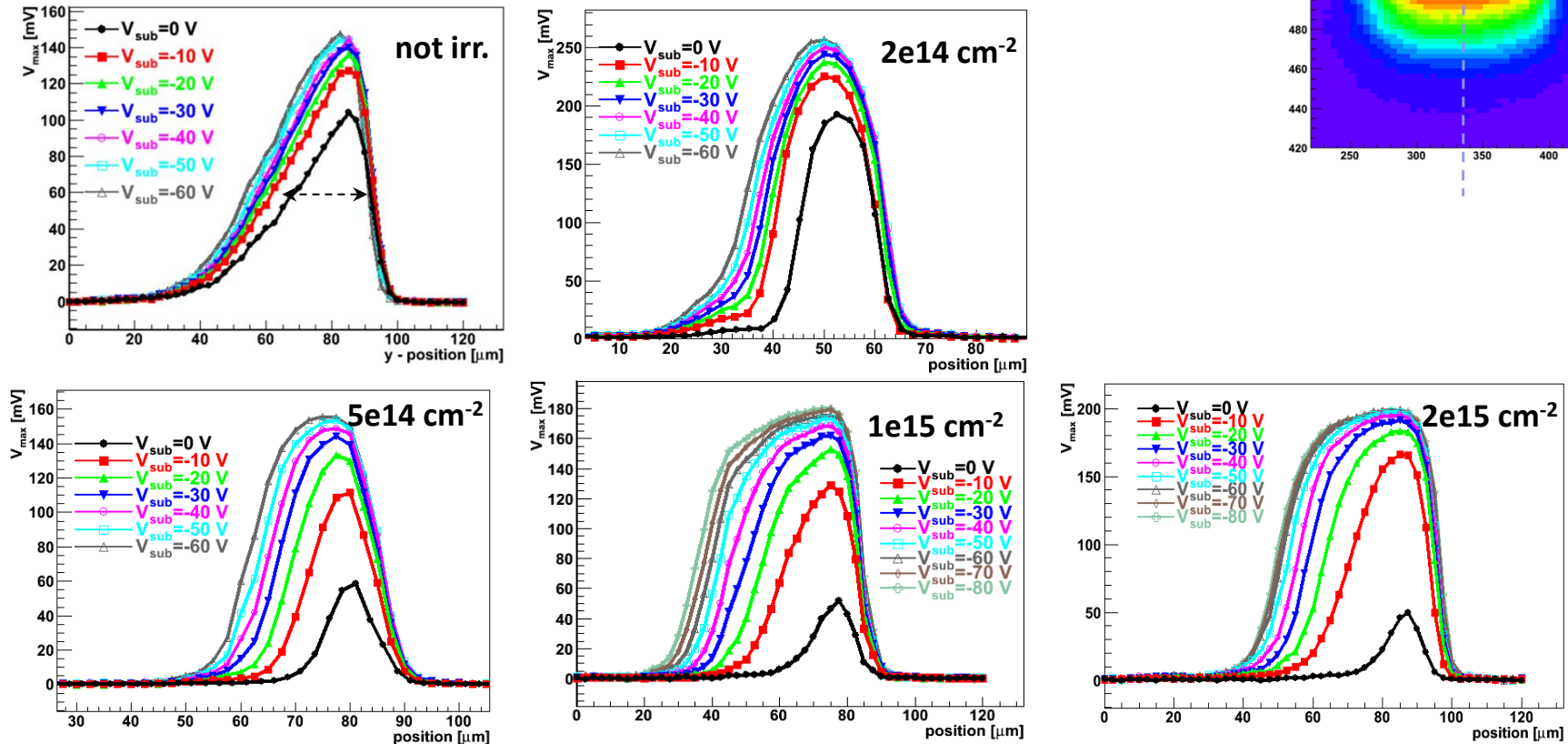


2e15 cm⁻²



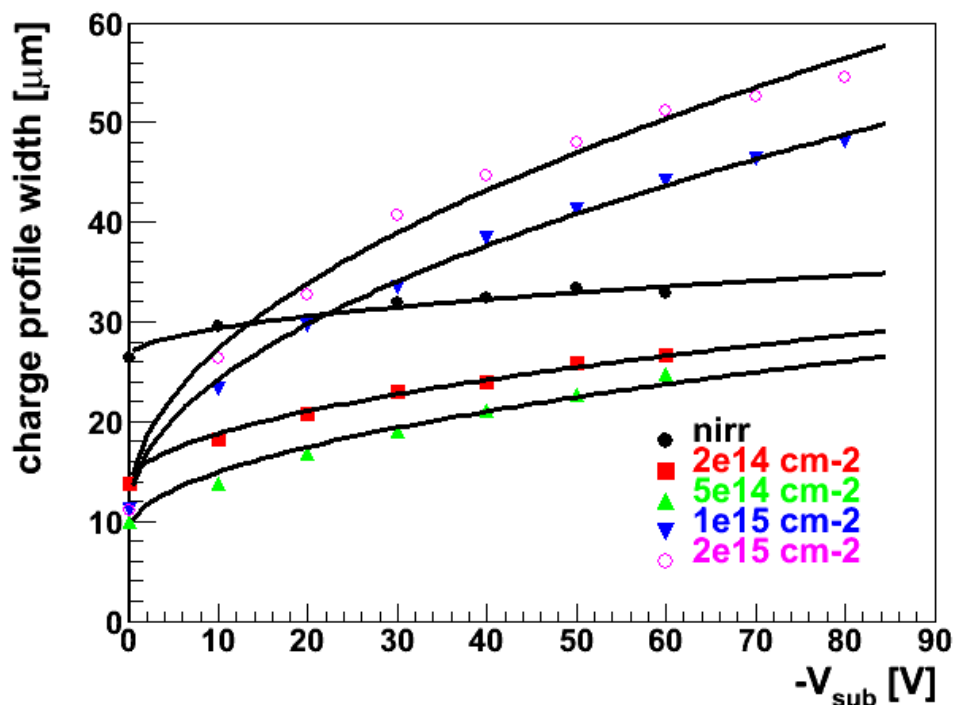
Charge collection region **larger** at high fluence

Charge collection profiles across center of the pixel



- tail (diffusion) seen before irradiation, almost disappears at 5e14 cm⁻²
- profile width (**FWHM**) is a measure of charge collection region (drift + diffusion)
→ the width of the laser beam (~ 8 μm FWHM) should be taken into account

Dependence of charge collection region on bias voltage



- at $V_{sub}=0$ V it is assumed that charge is collected by diffusion (note the FWHM of the beam)
- any additional bias increases depletion layer which adds to the diffusion
- effective doping concentration seems to decrease with fluence
 → **depletion region wider after irradiation!**

Points to effective acceptor removal – see talk by G. Kramberger tomorrow

Conclusions:

CCE with MIPs from Sr-90:

- first measurements with passive HV-CMOS devices on CHESS1 chip:
 - before irradiation: $Q = 1010 \text{ el} + 11 \text{ el/V}$
 - after irradiation with $2e14$ and $5e14 \text{ n/cm}^2$: $Q = 250 \text{ el} + 11 \text{ el/V}$
 - after $1e15 \text{ n/cm}^2$ larger charge than before irradiation: $\sim 3200 \text{ el @ } 120 \text{ V}$

E-TCT:

- passive pixel (CHESS1) : directly probe field region, charge collecting region, observe induced current pulses, long tails (diffusion).
After irradiation: → depleted region increases, no tails of induced pulses
- active pixel (HV2FEI4): map charge collection region, diffusion smaller after irradiation
 - significant widening of charge collecting region after $1e15 \text{ n/cm}^2$

Effective Acceptor removal!

In HVCMOS detectors charge collection increases with irradiation in certain fluence range → **selection of substrate material important to explore the effect**

Consistent!

Edge-TCT

Chess1, not irradiated

- longer integration, more charge collected deeper in the pixel (diffusion)

