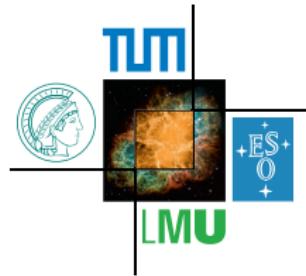


Constraining composite Higgs models with direct and indirect searches

Presented by Peter Stangl

Excellence Cluster Universe
Technische Universität München



Outline

- 1 Composite Higgs models
- 2 Analysis method
- 3 Analysis of a specific model
- 4 Conclusions

Based on:

Christoph Niehoff, PS, David M. Straub [arXiv:1508.00569]

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Composite Higgs models

Solving the naturalness problem

- ▶ Higgs not elementary but bound state of new strong interaction.
- ▶ Lightness of Higgs compared to new physics scale:
Higgs as pseudo-Nambu-Goldstone boson (pNGB)
of spontaneously broken global symmetry.

$$\mathcal{L} = \mathcal{L}_{\text{elementary}} + \mathcal{L}_{\text{composite}}$$

Kaplan, Georgi,
Phys.Lett. B136 (1984) 183
Dugan, Georgi, Kaplan,
Nucl.Phys. B254 (1985) 299
see also talks by
Brian Batell and Kiel Howe

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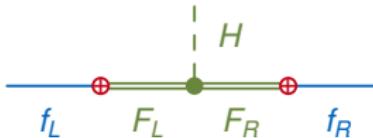
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$$\mathcal{L} = \mathcal{L}_{\text{elementary}} + \mathcal{L}_{\text{composite}} + \mathcal{L}_{\text{mixing}}$$

Avoiding flavour constraints

- Elementary fermions couple linearly to composite fermions.
- Mass eigenstates are mixture of both: **partial compositeness**.

Kaplan, Nucl. Phys. B365 (1991) 259–278



see talk by Francesco Sannino
for problems of UV completions
and alternatives

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The Goal and the challenge

The goal

Use a single framework to constrain the parameter space of a CHM by

- ▶ Radiative electroweak symmetry breaking
- ▶ Higgs physics
- ▶ Electroweak precision tests
- ▶ Flavour physics
- ▶ Direct LHC bounds on fermion and vector boson resonances

The Goal and the challenge

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Use a single framework to constrain the parameter space of a CHM by

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The challenge

Parameter space does not “factorize” into Standard Model and new physics parameters.

- ▶ **Partial compositeness:** fermion masses and CKM parameters are complicated functions of many model parameters.
- ▶ **Dynamically generated Higgs potential:** depends on gauge boson and fermion masses and couplings.

Strategy

The method

- ▶ Brute-force scan of parameter space not applicable.
- ▶ Instead, construct χ^2 -function from all observables:

$$\chi^2(\vec{\theta}) \equiv \sum_{i \in \text{observables}} \left(\frac{\mathcal{O}_i^{\text{th}}(\vec{\theta}) - \mathcal{O}_i^{\text{exp}}}{\sigma_i^{\text{error}}} \right)^2$$

Given a parameter point $\vec{\theta}$, it measures “how close” theoretical predictions $\mathcal{O}_i^{\text{th}}(\vec{\theta})$ are to experimental observations $\mathcal{O}_i^{\text{exp}}$.

- ▶ Numerically minimizing χ^2 -function yields “good” parameter points.

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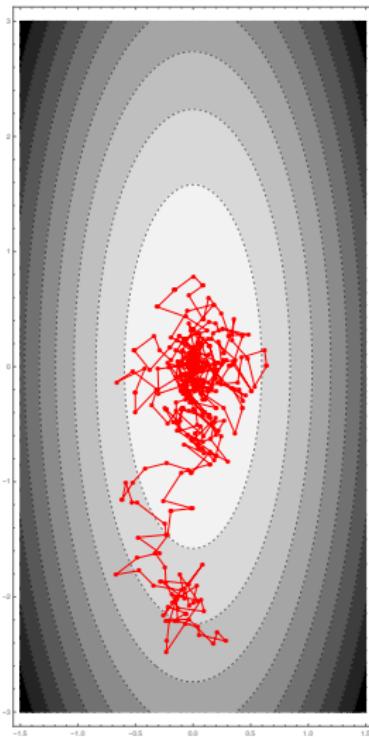
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- ▶ Numerically minimizing χ^2 -function yields “good” parameter points.

Still technically challenging

- ▶ Large dimensionality (ca. 30-50 parameters depending on specific model).
- ▶ Complicated dependence on all parameters.

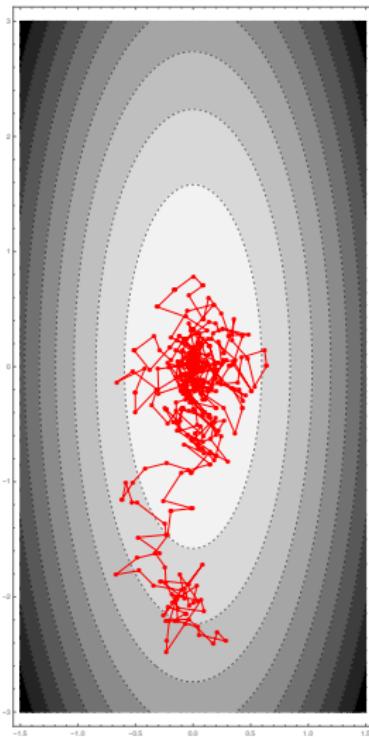
Implementation



Numerical minimization

1. Generate (random) starting points.
2. Use global minimizer [\[NLopt\]](#) to find viable region.
3. Use Markov Chain Monte Carlo [\[pypmc\]](#) to sample viable parameter space.
4. Keep only points satisfying each individual constraint at 3σ level.

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Computations

performed on the [C2PAP](#) computing cluster in Munich.

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The first model we analysed

- Calculable Higgs potential
- Minimal particle content



Deconstructed model
4DCHM with two sites.

De Curtis, Redi, Tesi,
arXiv:1110.1613

The first model we analysed

- Calculable Higgs potential
 - Minimal particle content
 - Only one Higgs doublet
 - Custodial symmetry protecting T parameter
- ⇒ Deconstructed model
4DCHM with two sites.
- ⇒ Minimal $\mathbf{SO(5)/SO(4)}$ coset.

De Curtis, Redi, Tesi,
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The first model we analysed

- Calculable Higgs potential \Rightarrow Deconstructed model
- Minimal particle content \Rightarrow **4DCHM with two sites.**
- Only one Higgs doublet \Rightarrow Minimal **$SO(5)/SO(4)$** coset.
- Custodial symmetry protecting T parameter \Rightarrow Composite quarks in fundamental **5-representations**, no composite leptons.
- Custodial protection of $Z b_L \bar{b}_L$ coupling \Rightarrow Composite quarks in fundamental **5-representations**, no composite leptons.
- Minimal particle content \Rightarrow Composite quarks in fundamental **5-representations**, no composite leptons.

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Agashe et al., arXiv:hep-ph/0605341
Contino et al., arXiv:hep-ph/0612048

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Minimal **$SO(5)/SO(4)$** coset.

- Custodial protection of $Z b_L \bar{b}_L$ coupling
- Minimal particle content



Composite quarks in fundamental **5**-representations, no composite leptons.

Agashe et al., arXiv:hep-ph/0605341
Contino et al., arXiv:hep-ph/0612048

- Suppression of contributions to $\Delta F = 2$ observables



Flavour symmetry only broken by \cancel{L}_{mixing} .
We consider four cases:
 $U(3)_{LC}^3$, $U(3)_{RC}^3$, $U(2)_{LC}^3$ or $U(2)_{RC}^3$.

Cacciapaglia et al., arXiv:0709.1714, Redi, Weiler, arXiv:1106.6357
Barbieri, Buttazzo, Sala, Straub, arXiv:1203.4218

Results

No viable parameter points for $U(3)_{LC}^3$

- ▶ Very strong constraints from electroweak precision tests and CKM unitarity.

cf. Barbieri, Buttazzo, Sala, Straub, arXiv:1211.5085

$U(2)_{LC}^3$, $U(2)_{RC}^3$ and $U(3)_{RC}^3$ flavour structures

- ▶ All constraints can be passed.
- ▶ Fine tuning $\Delta_{BG} < 100$ possible.
- ▶ Tensions in rare B decays can be explained in some part of parameter space.
- ▶ Tension in $\Delta F = 2$ observables could be reduced in $U(2)^3$ models.
- ▶ LHC run 2 can test many of viable parameter points.

$$\Delta_{BG} = \max_{i \in \text{parameters}} \left| \frac{\partial \ln m_Z}{\partial \ln \theta_i} \right|$$

Barbieri, Giudice,
Nucl. Phys. B306 (1988) 63

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Conclusions

New numerical method

Employing a new numerical method, it is possible to constrain composite Higgs models by both direct and indirect constraints in a single framework.

Findings of first analysis

- ▶ $U(3)_{\text{LC}}^3$ flavour structure disfavoured.
- ▶ Fine tuning $\Delta_{\text{BG}} < 100$ possible.
- ▶ Tensions in rare B decays could be explained.
- ▶ Tension in $\Delta F = 2$ observables could be reduced in $U(2)^3$ models.
- ▶ LHC run 2 can test many of viable parameter points.

More details

Christoph Niehoff, PS, David M. Straub [arXiv:1508.00569]

Backup slides

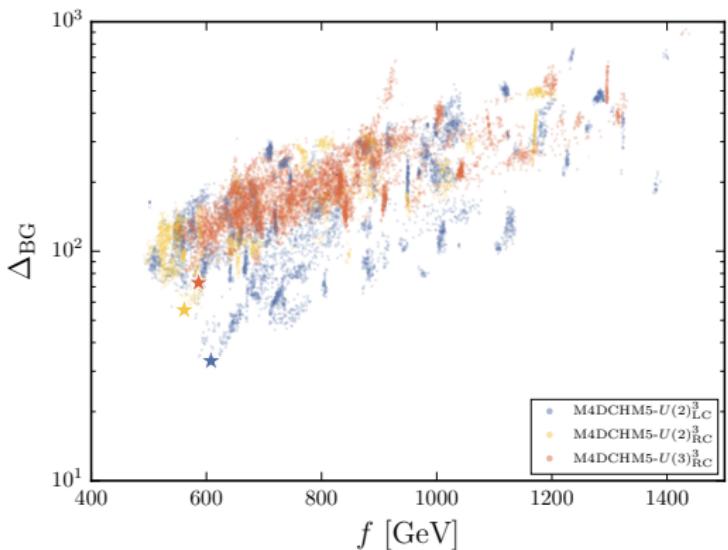
Fine-tuning

Barbieri-Giudice measure

$$\Delta_{\text{BG}} = \max_{i \in \text{parameters}} \left| \frac{\partial \ln m_Z}{\partial \ln \theta_i} \right|$$

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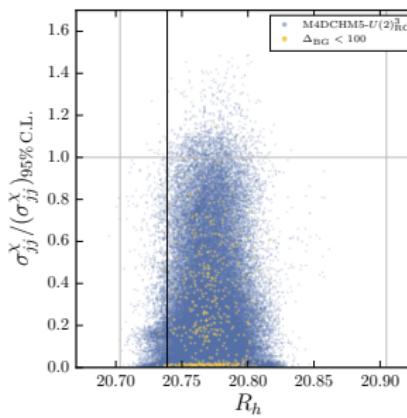
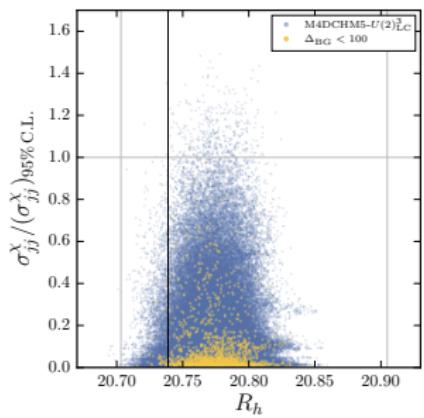
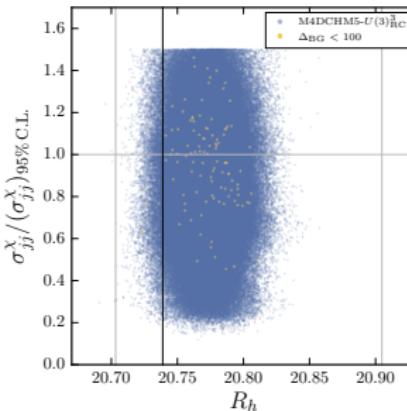
quantifies sensitivity of weak scale to variations in model parameters θ_i .



- ▶ $U(2)_{\text{LC}}$: $\Delta_{\text{BG}}^{(\min)} = 33$
- ▶ $U(2)_{\text{RC}}$: $\Delta_{\text{BG}}^{(\min)} = 55$
- ▶ $U(3)_{\text{RC}}$: $\Delta_{\text{BG}}^{(\min)} = 73$

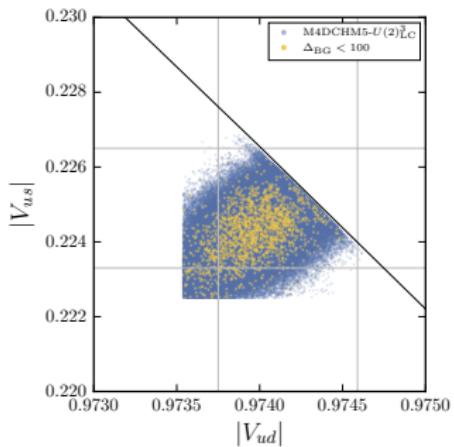
Light quark compositeness

Results for hadronic Z width and dijet angular distributions



Light quark compositeness

Results for CKM unitarity

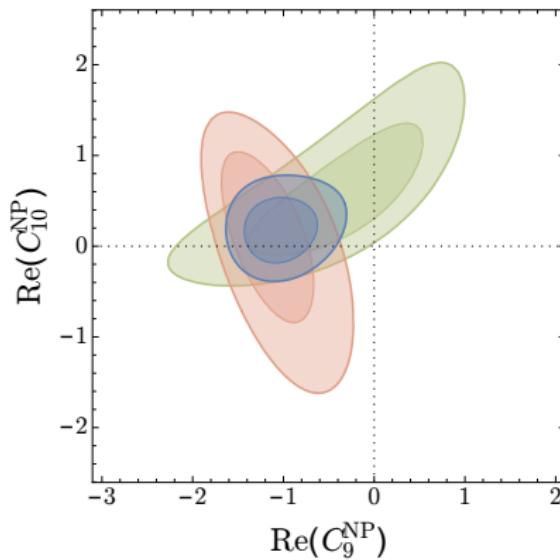


No relevant constraint from first-row CKM unitarity on *RC* models.

Rare B decays

Tensions between SM and data

Recent fits prefer a negative shift in Wilson coefficient C_9 (and possibly positive in C_{10}):

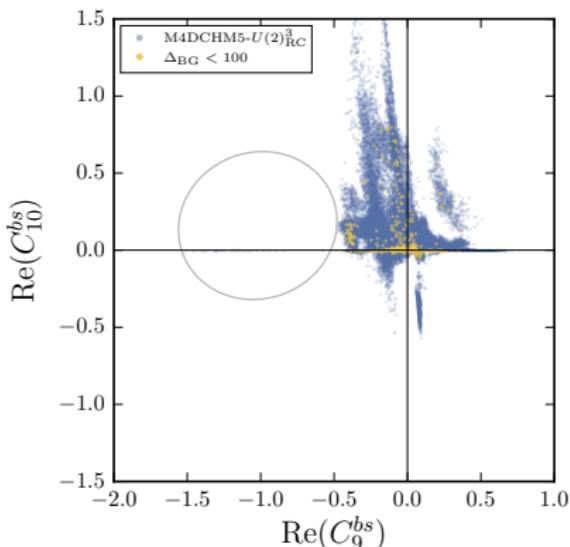
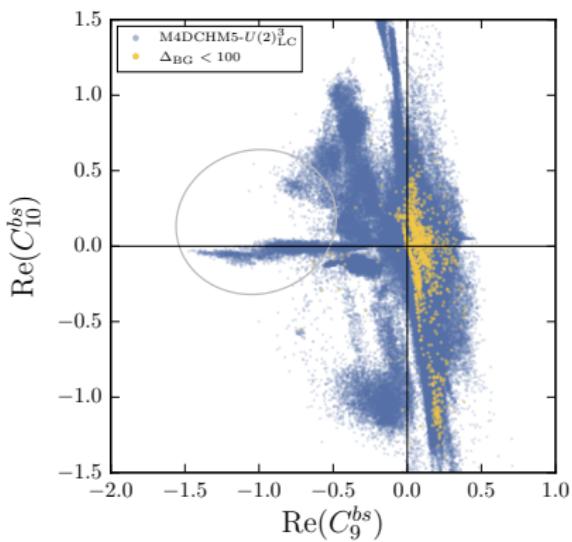


Altmannshofer, Straub,
arXiv:1411.3161

Rare B decays

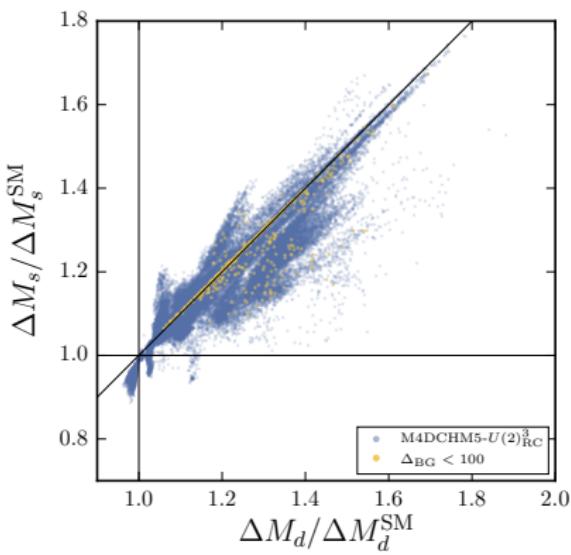
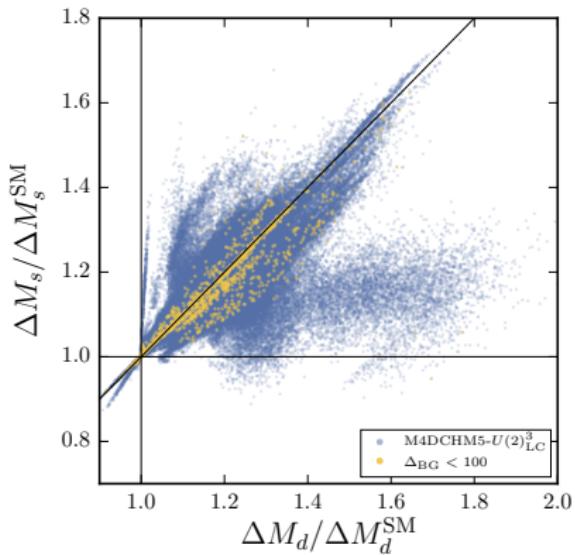
Tensions could be explained

- We find points with negative shift in C_9 .
- Then, we predict light neutral vector resonance (with large BR to $t\bar{t}$).



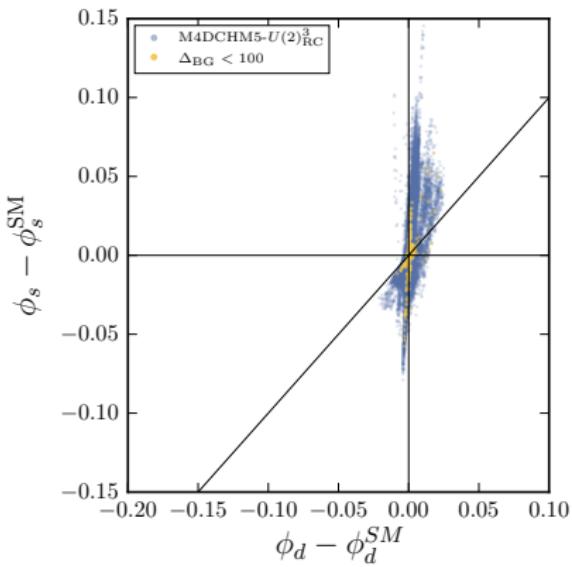
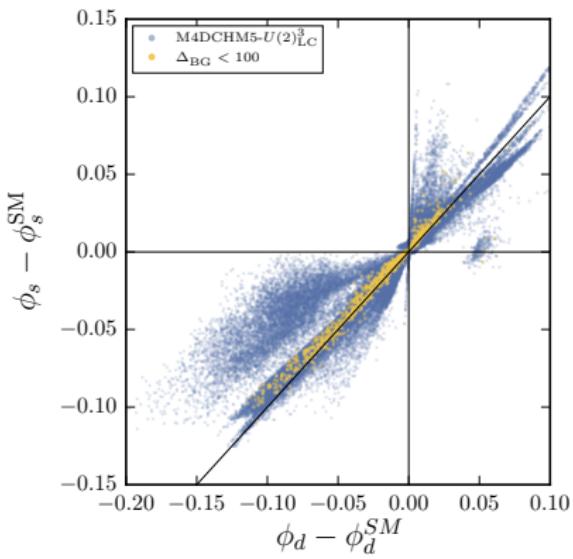
Meson-Antimeson mixing

Mass differences in B_d^0 and B_s^0 system



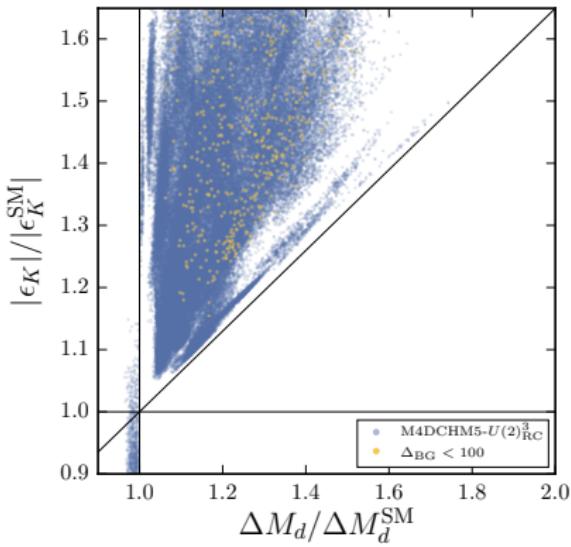
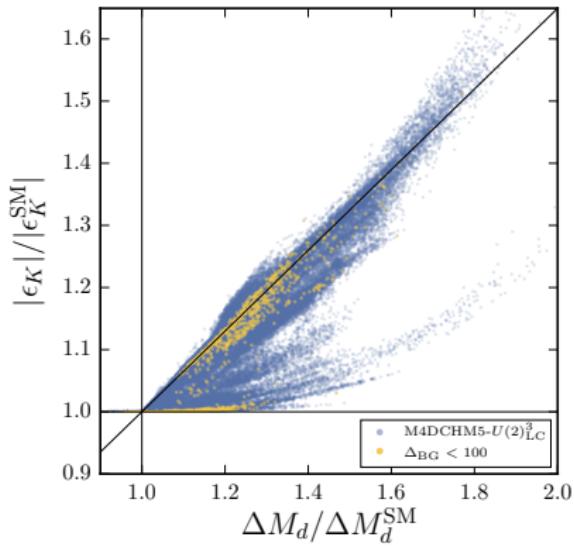
Meson-Antimeson mixing

Mixing phases in B_d^0 and B_s^0 system



Meson-Antimeson mixing

Indirect CP violation in K^0 mixing



Meson-Antimeson mixing

Facing new lattice results

Tension between $\Delta M_d/\Delta M_s$ and ϵ_K increased.

Fermilab Lattice, MILC, arXiv:1602.03560

Blanke, Buras, arXiv:1602.04020

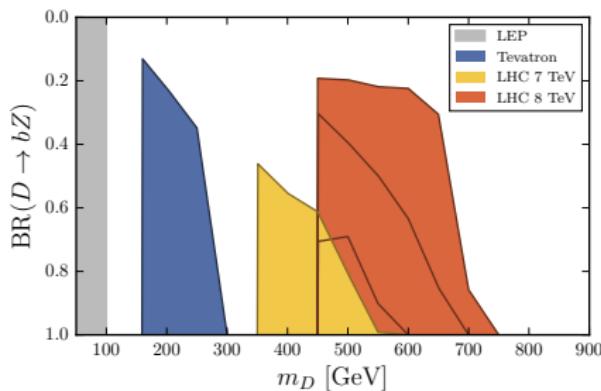
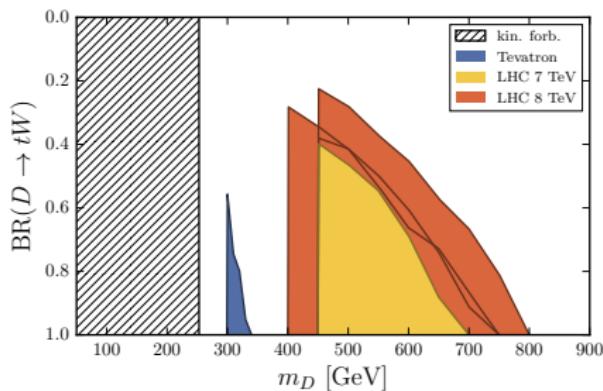
Reducing the tension	$U(2)^3_{LC}$	$U(2)^3_{RC}$	$U(3)^3_{RC}$
$\Delta M_d/\Delta M_d^{\text{SM}} < \Delta M_s/\Delta M_s^{\text{SM}}$	✓	✓	✗
$\phi_d - \phi_d^{\text{SM}} \neq 0$	✓	✗	✗
$ \epsilon_K / \epsilon_K^{\text{SM}} > \Delta M_d/\Delta M_d^{\text{SM}}$	✗	✓	✗

Quark partner searches

Included experimental analyses

20 different analyses from **ATLAS**, **CMS** and **CDF** included into numerics.

Gaps in experimental analyses

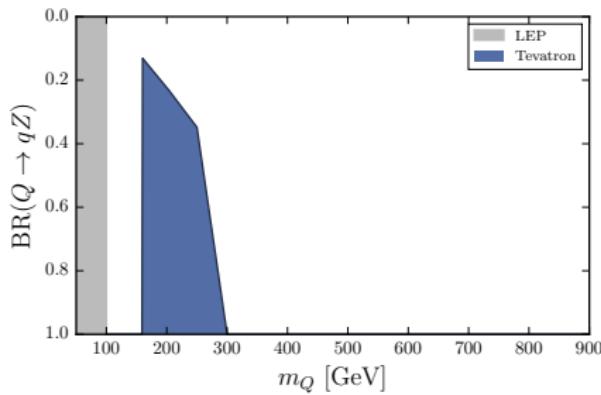
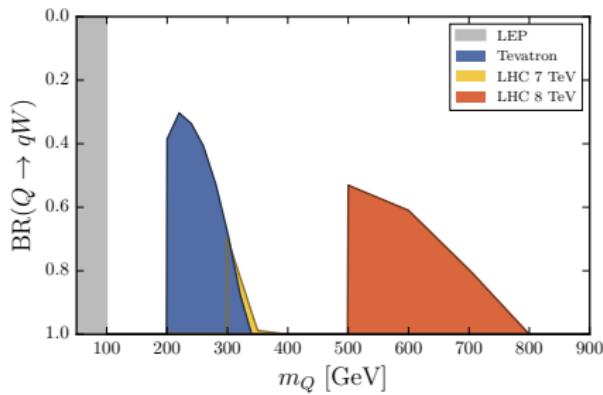


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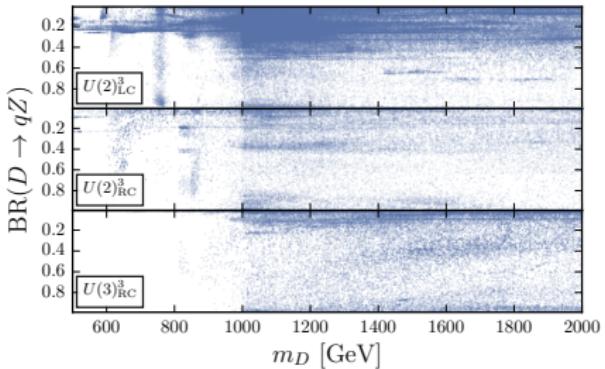
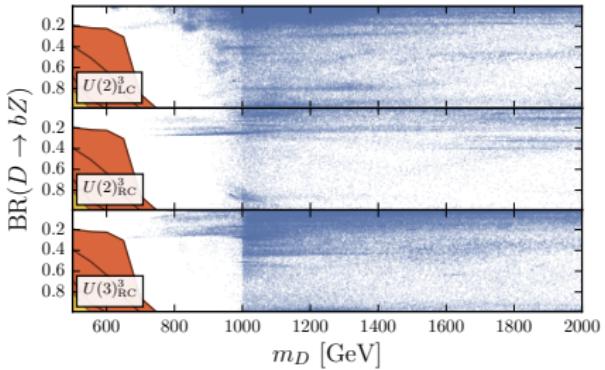
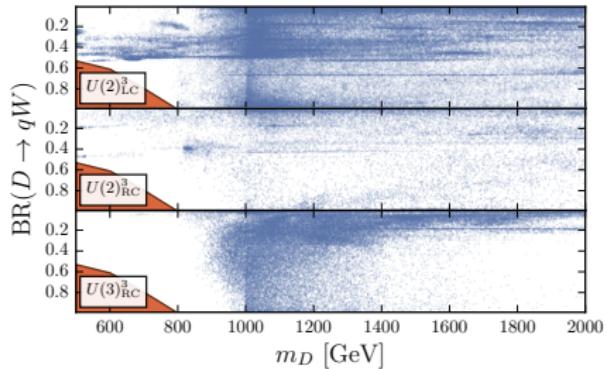
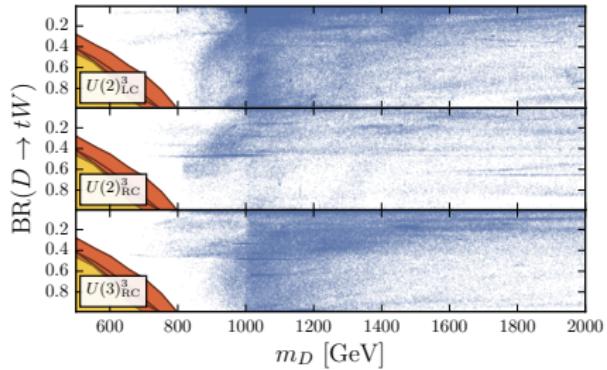
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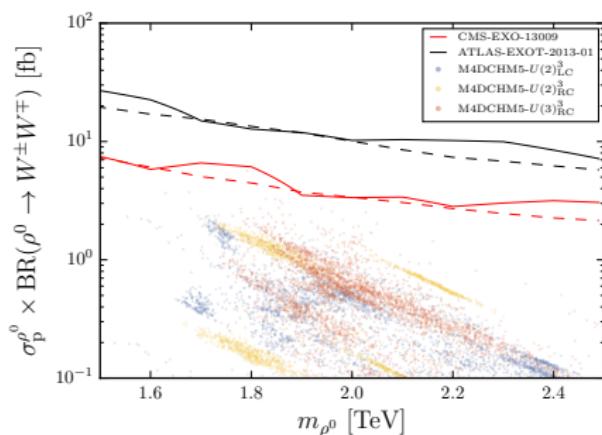
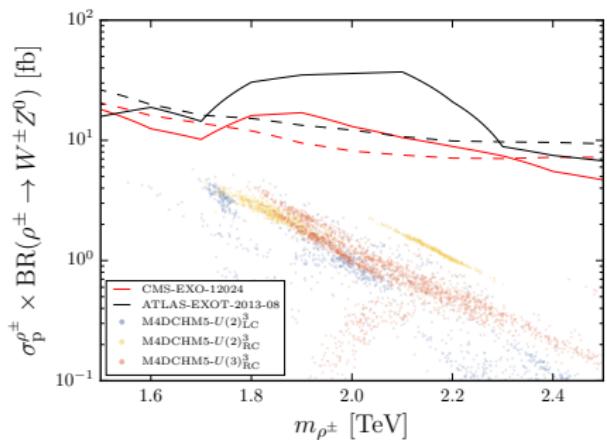


Vector partner searches

Included experimental analyses

14 different analyses from **ATLAS** and **CMS** included into numerics.

Results



Additional constraints

In addition to the ones discussed in these slides, our analysis includes constraints from

- ▶ Higgs production and decay
- ▶ Oblique parameters
- ▶ Top decays