

ATLAS SEARCHES FOR NEW HEAVY FLAVOUR STATES

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1. Observation of an excited B_c^\pm meson state with the ATLAS Detector
 - [Phys. Rev. Lett. 113, no. 21, 212004 \(2014\)](#), [arXiv:1407.1032 \[hep-ex\]](#)
 - 4.9 fb⁻¹ of $\sqrt{s} = 7$ TeV and 19.2 fb⁻¹ of $\sqrt{s} = 8$ TeV pp collision data
2. Search for the X_b and other hidden-beauty states in the $\pi^+\pi^-\Upsilon(1S)$ channel at ATLAS
 - [Phys. Lett. B 740, 199 \(2015\)](#), [arXiv:1410.4409 \[hep-ex\]](#)
 - 16.2 fb⁻¹ of $\sqrt{s} = 8$ TeV pp collision data

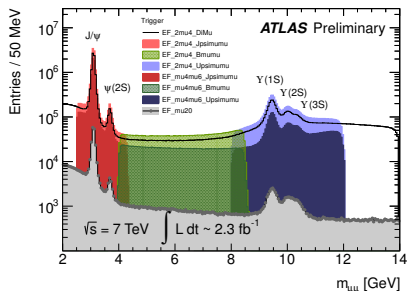
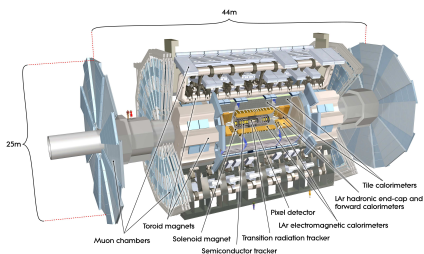
ATLAS DETECTOR AND TRIGGER SYSTEM

ATLAS: general purpose detector:

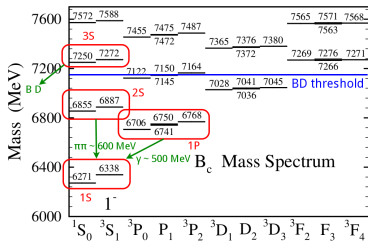
- Subsystems essential for B-physics:
Inner detector and Muon spectrometer,
- Inner detector: tracking, momentum and vertexing, $|\eta| < 2.5$, d_0 resolution $\sim 10\mu m$,
- Muon spectrometer: trigger and muon identification,
- J/ψ mass resolution ~ 60 MeV, $\Upsilon(1S) \sim 120$ MeV,
- ATLAS trigger system: hardware *Level-1 trigger* and two-level software *High-Level Trigger*

Trigger selection for heavy flavour studies is mostly based on di-muon signature:

- muon p_T threshold (4 or 6 GeV)
- di-muon vertex reconstruction
- invariant mass window

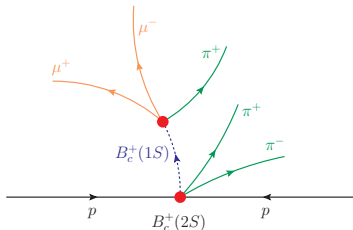


B_c STATES



The B_c mass spectrum, taken from Phys. Rev. D 70, 054017 (2004).

- Excited states of B_c^\pm are predicted by nonrelativistic potential models, perturbative QCD and lattice calculations,
- mass of $B_c^\pm(2S)$ in range 6835–6917 MeV,
- both the $1S$ and $2S$ states have pseudoscalar (0^-) and vector (1^-) spin states that are predicted to differ in mass by about 20–50 MeV,
- ATLAS not sensitive enough to distinguish 0^- and 1^- states: missing soft gamma, mass resolution.



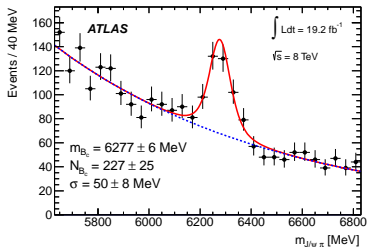
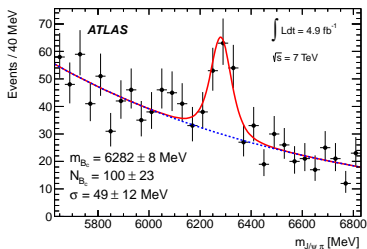
$B_c^\pm \rightarrow J/\psi(\mu^+\mu^-)\pi^\pm$ SELECTION AND FIT

B_c^+ (1S) selection for 7 TeV (8 TeV) data:

- $p_T(\mu_1) > 4$ GeV and $p_T(\mu_2) > 6$ GeV,
- J/ψ vertex fit $\chi^2/\text{NDF} < 15$,
- $m(J/\psi)$ within $\pm 3\sigma$ of the nominal (σ depending on the rapidity range),
- B_c^+ vertex fit $\chi^2/\text{NDF} < 2.0$ (1.5) with dimuon mass constrained to mass J/ψ ,
- $p_T(B_c^+) > 15$ GeV (18 GeV),
- pion candidate from $B_c^+ p_T(\pi^+) > 4$ GeV and its impact $d^0/\sigma(d^0)(\pi^+) > 5$ (4.5)

Extended unbinned fit of the mass distribution

- *Signal*: Gaussian with per-candidate errors
- *Background*: exponential
- the stability of B_c^\pm yield was checked through its normalization to $B^\pm \rightarrow J/\psi K^\pm$ that were reconstructed with similar requirements.



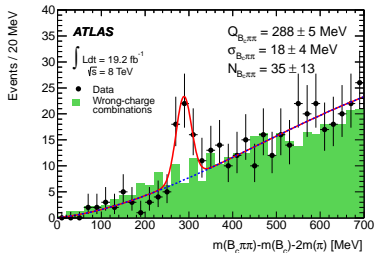
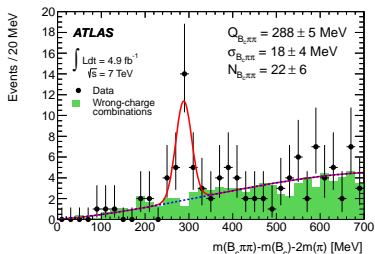
$B_c^\pm(2S) \rightarrow B_c^\pm(1S)\pi^+\pi^-$ SELECTION AND FIT

Selection of $B_c^\pm(2S)$ candidates:

- $B_c^+(1S)$ candidates within $\pm 3\sigma$ of the fitted mass
- pion tracks from primary vertex with $p_T(\pi^+, \pi^-) > 400$ MeV,
- cascade vertex fitter was applied to build decay topology: the refitted triplet of $B_c^+(1S)$ tracks and the pair of PV pion tracks must intersect in two separate vertices,
- for several candidates in event, the one with the best cascade fit χ^2 is kept

Extended unbinned fit of Q -value distribution:

- use $Q = m(B_c^+\pi^+\pi^-) - m(B_c^+) - 2m(\pi^+)$ to improve resolution,
- *Signal*: Gaussian,
- *Background*: 3rd order polynomial,
- *Wrong charge combination* (same-sign π) used for background control



SUMMARY OF $B_c^\pm(2S) \rightarrow B_c^\pm(1S)\pi^+\pi^-$ ANALYSIS

- Significance of the observed signal calculated with toy studies accounting for a “look elsewhere effect” ,
- combined significance is 5.2σ (3.7σ and 4.5σ in 7 and 8 TeV data, respectively),
- dominant source of systematic of the Q -value is the *fitting procedure*,
- a new state observed at $Q = 288.3 \pm 3.5$ (stat.) ± 4.1 (syst.) MeV (error-weighted mean of 7 and 8 TeV values),
- this new state corresponds to a mass 6842 ± 4 (stat.) ± 5 (syst.) MeV, and is consistent with the predicted mass of $B_c(2S)$.

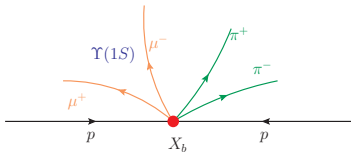
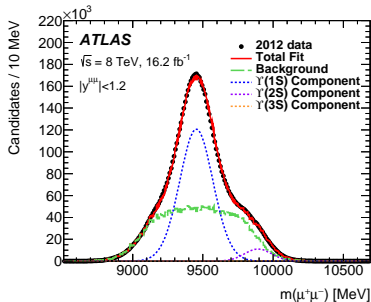
- $X(3872)$ is the best-studied new hidden-charm state, observed by many experiments
 - mass, narrow width, $J^{PC} = 1^{++} \rightarrow$ unlikely a conventional quarkonium,
 - weakly bound $D^0\bar{D}^{*0}$ molecule or $[qc][\bar{q}\bar{c}]$ tetraquark
 - the relative production rate R for $X(3872)$ (CMS, [JHEP 1304, 154 \(2013\)](#))

$$R = \frac{\sigma_{pp \rightarrow X(3872)} \cdot \mathcal{B}_{X(3872) \rightarrow J/\psi \pi^+ \pi^-}}{\sigma_{pp \rightarrow \psi(2S)} \cdot \mathcal{B}_{\psi(2S) \rightarrow J/\psi \pi^+ \pi^-}} = (6.56 \pm 0.29 \pm 0.65)\%$$

- Heavy-quark symmetry suggests a hidden-beauty partner X_b ;
 - mass predictions vary (e.g. 10561 MeV for the molecular model of Swanson, while tetraquark predicts 10492, 10593, or 10682 MeV, depending on the flavour of the light quarks)
 - Decay $X_b \rightarrow \pi^+ \pi^- \Upsilon(1S)$ is a straightforward way to reconstruct X_b ,
 - as by-product, $\Upsilon(1^3D_J)$, $\Upsilon(10860)$, $\Upsilon(11020)$ can be studied with the same final state.

$\Upsilon(1S)$ reconstruction

- Pairs of oppositely charged muons with $p_T(\mu) > 4$ GeV and $|\eta(\mu)| < 2.3$ are fitted to the common vertex,
- require the matching between the reconstructed and trigger-level muons,
- dimuons in mass range ± 350 MeV around $\Upsilon(1S)$ mass are retained.

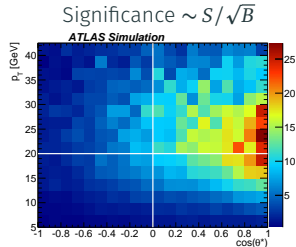
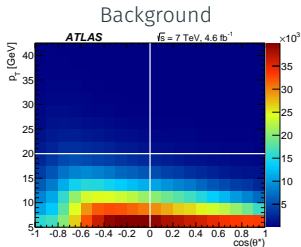
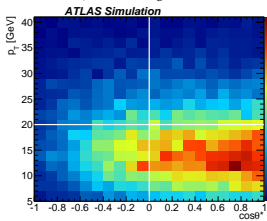
Build X_b candidate

- Dipion candidate are formed from oppositely charged pions with $p_T(\pi) > 400$ MeV and $|\eta(\pi)| < 2.5$
- $\Upsilon(1S) \rightarrow \mu^+\mu^-$ candidate and dipion system are combined by performing a four-track common-vertex fit ($\chi^2 < 20$),
- the $\mu^+\mu^-$ mass constrained to mass $\Upsilon(1S)$; this significantly improve the mass resolution,
- candidates with mass < 11.2 GeV are retained.

The analysis is performed in 8 bins

- $|y(X_b)|$: barrel ($|y| < 1.2$) and endcap ($1.2 < |y| < 2.4$) due to different mass resolution,
- $(p_T(X_b), \cos \theta^*)$: split into 4 quadrants – different S/B ratio, (θ^* is an angle between $\pi^+\pi^-$ momentum in the parent rest frame and the parent momentum in lab frame),
- fraction of the signal in each bins are defined by *splitting functions* derived from the simulation,
- optimize the expected significance for a weak signal at 10561 MeV,
- binned, extended maximum-likelihood fit is performed simultaneously to the 8 bins ($|y|, p_T, \cos \theta^*$) in a local region around the mass of interest.

MC sample $m(X_b) = 10.561$ GeV

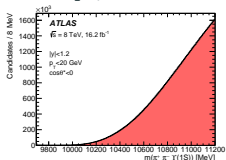


$\Upsilon(2S)$ AND $\Upsilon(3S)$ FITS

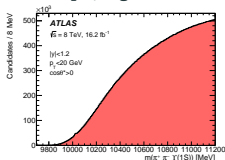
- Clear peaks of $\Upsilon(2S)$ and $\Upsilon(3S)$ are observed; no other visible signals,
- kinematic bin most sensitive to an X_b signal: $|y| < 1.2$, $p_T > 20$ GeV and $\cos\theta^*$,
- normalized to $\Upsilon(2S)$ yields and validated on $\Upsilon(3S)$.

barrel

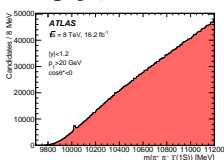
low p_T , low $\cos\theta^*$



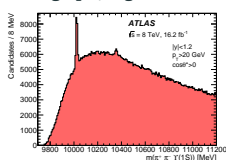
low p_T , high $\cos\theta^*$



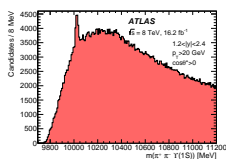
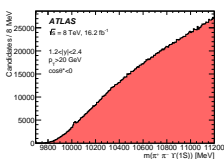
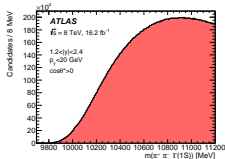
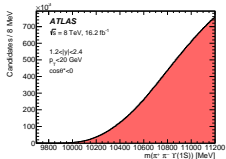
high p_T , low $\cos\theta^*$



high p_T , high $\cos\theta^*$



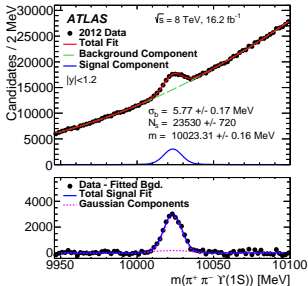
endcap



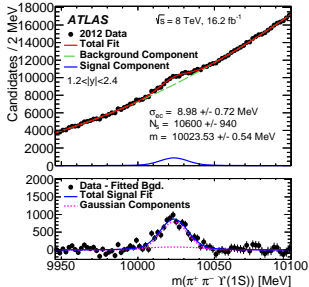
$\Upsilon(2S)$ AND $\Upsilon(3S)$ FITS

- Differential production cross section $d^2\sigma/dydp_T$ for $\Upsilon(2S)$ and $\Upsilon(3S)$ measured by ATLAS/CMS at 7 TeV were used to calculate production weights, which are applied to 8 TeV $\Upsilon(2S)/\Upsilon(3S)$ simulated sample,
- $\Upsilon(2S)$ total fitted yield $N_{2S} = 34300 \pm 800$ is consistent with expected $N_{2S}^{\text{expected}} = (\sigma\mathcal{B})_{2S} \cdot \mathcal{L} \cdot \mathcal{A} \cdot \epsilon = 33300 \pm 2500$,
- $\Upsilon(3S)$ total fitted yield $N_{3S} = 11600 \pm 1300$ agrees with $N_{3S}^{\text{expected}} = 11400 \pm 1500$, statistical significance for simultaneous fit to the analysis bins is $z = 8.7$ while for individual fit to the most sensitive bin $z = 6.5$.

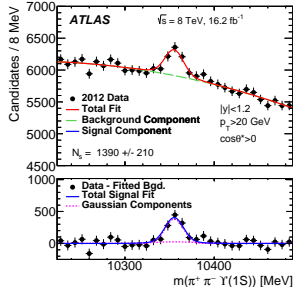
$\Upsilon(2S)$ in barrel



$\Upsilon(2S)$ in endcap



$\Upsilon(3S)$ in the most sensitive bin

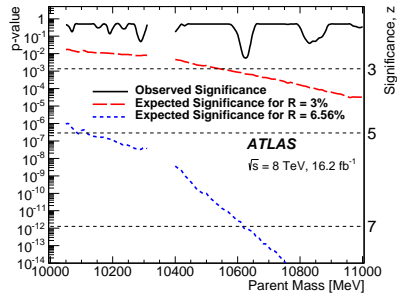


Strategy:

- A hypothesis test for signal presence across 10–11 GeV range every 10 MeV,
- simultaneous fit in 8 analysis bins
- for each mass, extract p -value and significance.

Assumptions:

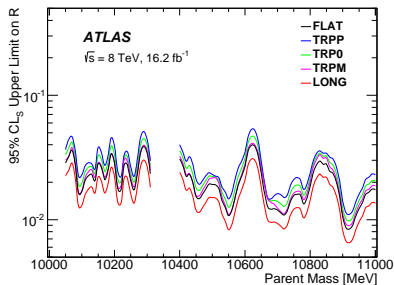
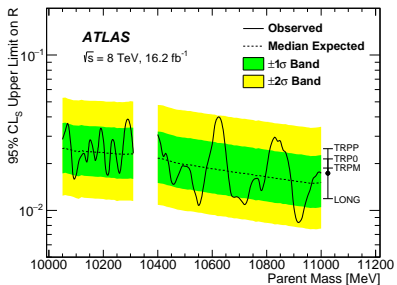
- Look for narrow state,
- resolution dependence on $|y|, p_T$ is $\Upsilon(nS)$ -like,
- phase-space shape of $m(\pi^+\pi^-)$.



mass ranges near $\Upsilon(2S)$ and $\Upsilon(3S)$ are excluded

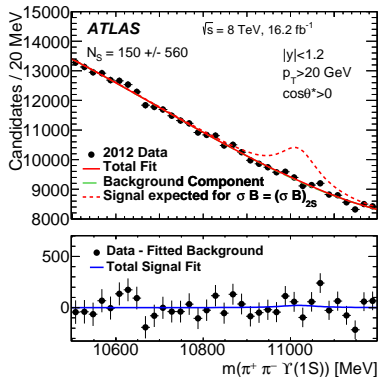
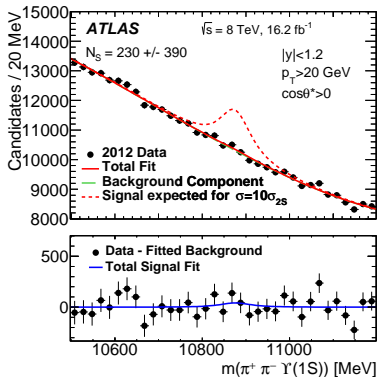
X_b UPPER LIMITS CALCULATION

- CL_s limit on $R = (\sigma\mathcal{B})/(\sigma\mathcal{B})_{2S}$ at 95% CL is evaluated as a function of mass,
- various assumptions on X_b production polarization shift the limits (1 longitudinal and 3 transverse spin-alignment scenarios),
- this is the most sensitive X_b production search for $m > 10.1$ GeV,
- the null result of the search $X_b \rightarrow \pi^+\pi^-\gamma(1S)$ decay does not tell us if the X_b exists, because for an isoscalar with $J^{PC} = 1^{++}$ such a decay is forbidden by G-parity conservation.



RESULTS FOR $\Upsilon(1^3D_J)$, $\Upsilon(10860)$, $\Upsilon(11020)$

- $\Upsilon(1^3D_J)$ triplet fit attempted with additional signal shapes for the three masses: 10156, 10164, and 10170 MeV; a significance of $z = 0.12$ is found,
- upper limit of relative cross-section $\sigma(\Upsilon(1^3D_2))/\sigma(\Upsilon(2S)) < 0.55$ (using known $\mathcal{B}(\Upsilon(1^3D_2) \rightarrow \pi^+ \pi^- \Upsilon(1S)) = (6.6 \pm 1.6) \times 10^{-3}$ from BaBar)
- broad resonances $\Upsilon(10860)$, $\Upsilon(11020)$ searched for in grid of mass and width, (using world-average masses and uncertainties), largest significances of $z = 1.1$ and 0.6
- plots shown with rates $\sigma_{10860} = 10\sigma(2S)$ and $\sigma_{10860} = \sigma(2S) \cdot \mathcal{B}(2S)$



Two analyses performed by ATLAS are reported:

1. Observation of a new state decaying into $B_c^+ \pi^+ \pi^-$ in 7 TeV and 8 TeV data
 - First observation of an excited state in B_c^+ sector
 - Measured mass 6842 ± 4 (stat.) ± 5 (syst.) MeV is consistent with the predicted mass of the $B_c^+(2S)$ state
2. Search for X_b state in $\pi^+ \pi^- \Upsilon(1S)$ channel in 8 TeV data
 - No evidence for new narrow state is found for masses 10.05–10.31 GeV and 10.40–11.00 GeV
 - Upper limits are set on the ratio $R = [\sigma(pp \rightarrow X_b) \mathcal{B}(X_b \rightarrow \pi^+ \pi^- \Upsilon(1S))]/[\sigma(pp \rightarrow \Upsilon(2S)) \mathcal{B}(\Upsilon(2S) \rightarrow \pi^+ \pi^- \Upsilon(1S))]$, with results ranging from 0.8% to 4.0%, depending on the X_b mass
 - No evidences for other searched states $\Upsilon(1^3D_J)$ triplet, $\Upsilon(10860)$ and $\Upsilon(11020)$
 - Upper limit on relative cross section is set $\sigma(\Upsilon(1^3D_2))/\sigma(\Upsilon(2S)) < 0.55$
 - The searches for the isospin-conserving X_b decays are in progress