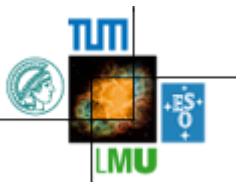


Renaissance of Kaon Flavour Physics

Andrzej J. Buras
(Technical University Munich, TUM-IAS)

**Beauty 2016, Marseille,
May 2016**



Plan for next 28 min

1.

ε'/ε strikes back

2.

$\varepsilon_K \leftrightarrow \Delta M_{s,d}$ tension in SM and CMFV

Intermezzo: $K \rightarrow \pi\nu\bar{\nu}$ in the Standard Model

3.

Implications for ε'/ε , ε_K , $K \rightarrow \pi\nu\bar{\nu}$
(Z, Z'- FCNCs) ΔM_K

4.

Highlights from 331, LHT, Vector-Like Quark Models

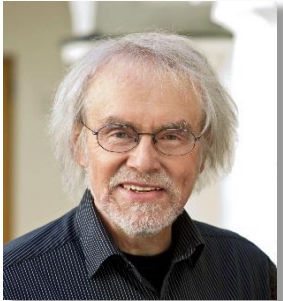
5.

Outlook

Section 1

ε'/ε strikes back

2015 Anatomy of ε'/ε : 1507.06345



AJB



Martin Gorbahn



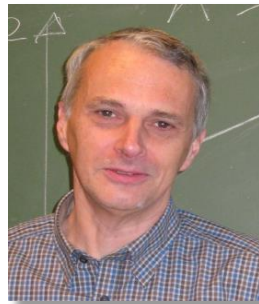
Sebastian Jäger



Matthias
Jamin



AJB



Jean-Marc Gérard

Large N news
1507.06326

FSI
1603.05686

ε'/ε strikes back (CP-Violation in $K_L \rightarrow \pi\pi$)

New results on hadronic matrix elements of QCD penguin (B_6) and electroweak penguin (B_8) operators

Large N approach to QCD

: $B_6 < B_8 < 1$



Upper Bound on ε'/ε in the Standard Model

AJB + Gérard (1507.06326)

Confirmed by Lattice QCD

: $B_6 = 0.57 \pm 0.19$ $B_8 = 0.76 \pm 0.05$

RBC-UKQCD

Anatomy of ε'/ε in the Standard Model

: $(\varepsilon'/\varepsilon) = (1.9 \pm 4.5) \cdot 10^{-4}$

AJB, Gorbahn, Jäger, Jamin (1507.06345)

$(\varepsilon'/\varepsilon) = (6.0 \pm 2.4) \cdot 10^{-4}$ for $B_6 = B_8 = 0.76$

$(\varepsilon'/\varepsilon)_{\text{exp}} = (16.6 \pm 2.3) \cdot 10^{-4}$

Possible New Physics

$(8.6 \pm 3.2) \cdot 10^{-4}$ for $B_6 = B_8 = 1.0$

Implications for $K \rightarrow \pi\nu\bar{\nu}$

- Z' general (AJB, Buttazzo, Knecht, 1507.08672)
- Littlest Higgs Model (Blanke, AJB, Recksiegel, 1507.06316)
- 331 Models (AJB, De Fazio, 1512.02869, 1604.02344)
- New Strategy (AJB, 1601.00005)
- Vector-like Quarks (Bobeth, AJB, Celis, Jung, 1605.xxxx)

Four dominant contributions to ε'/ε in the SM

AJB, Jamin, Lautenbacher (1993); AJB, Gorbahn, Jäger, Jamin (2015)

$$\text{Re}(\varepsilon'/\varepsilon) = \left[\frac{\text{Im}(V_{td} V_{ts}^*)}{1.4 \cdot 10^{-4}} \right] 10^{-4} \left[-3.6 + 21.4 \cdot B_6^{(1/2)} + 1.2 - 10.4 \cdot B_8^{(3/2)} \right]$$

From $\text{Re}A_0$
From $\text{Re}A_2$

(Q₄)

$(V-A) \otimes (V-A)$
QCD Penguins

$(V-A) \otimes (V+A)$
QCD Penguins

$(V-A) \otimes (V-A)$
EW Penguins

$(V-A) \otimes (V+A)$
EW Penguins

Assumes that $\text{Re}A_0$ and $\text{Re}A_2$ ($\Delta I=1/2$ Rule) fully described by SM (includes isospin breaking corrections)

Extracted from



RBC-UKQCD

$B_6^{(1/2)} = B_8^{(3/2)} = 1$ in the large N limit

$B_6^{(1/2)} = 0.57 \pm 0.19$

$B_8^{(3/2)} = 0.76 \pm 0.05$

Why $B_6^{(1/2)} < B_8^{(3/2)} < 1$?

and not $B_6^{(1/2)} > 1$, $B_8^{(3/2)} < 1$ (Pallante, Pich... FSI
2000)

Answer in Large N (Dual QCD) Approach

AJB + Gérard (1507.06326)

Before 2015 it was wrongly assumed that

$$B_6^{(1/2)} = B_8^{(3/2)} = 1 \quad \text{at } \mu \approx 0(1 \text{ GeV})$$

But $B_6^{(1/2)} = B_8^{(3/2)} = 1$ is large N prediction
for $\mu = m_\pi$ not $\mu = 0(1 \text{ GeV})$

Meson evolution $m_\pi \rightarrow \mu = 0(1 \text{ GeV})$ suppresses
 $B_6^{(1/2)}$ and $B_8^{(3/2)}$ below 1 and $B_6^{(1/2)}$ stronger than $B_8^{(3/2)}$
in accordance with quark evolution for $\mu > 1 \text{ GeV}$

FSI in $K \rightarrow \pi\pi$

AJB, Gérard 1603.05686

**Relevant for $\Delta I=1/2$ Rule
(in agreement with Pallante, Pich,...)**

**Less important for ε'/ε
(in variance with Pallante, Pich,...)**

**New application of dual QCD to $K \rightarrow \pi^+\pi^-$
(Caluccio-Leskow, D'Ambrosio, Greynat, Nath, 1604.09721)**

ε'/ε within SM

$$\varepsilon'/\varepsilon \sim \left[\frac{\text{Re } A_2}{\text{Re } A_0} \text{Im } C_6 \langle Q_6 \rangle_0 - \text{Im } C_8 \langle Q_8 \rangle_2 + \text{smaller contributions} \right]$$
$$\left\{ \frac{\text{Re } A_2}{\text{Re } A_0} \approx \frac{1}{22} \quad \frac{\text{Im } C_6}{\text{Im } C_8} \approx 90 \quad \frac{\langle Q_8 \rangle_2}{\langle Q_6 \rangle_0} \approx 2 \right\} \Rightarrow \text{strong cancellations}$$

ε'/ε beyond SM

(Q_6, Q_8, Q'_6, Q'_8)

1. Generally Q_8 wins over Q_6 because $\left(\frac{\text{Im } C_6}{\text{Im } C_8} \right)^{\text{NP}} \approx 0(1)$ but can provide $\Delta(\varepsilon'/\varepsilon) > 0$
2. Q_6 wins over Q_8 in the presence of a flavour symmetry forbidding Q_8
3. Chromomagnetic operators (not in this talk)

Section 2

$\varepsilon_K \leftrightarrow \Delta M_{s,d}$ tension in SM and CMFV

(1602.04020)



Monika Blanke



AJB

Universal Unitarity Triangle 2016

(CMFV)

AJB, Gambino, Gorbahn, Jäger, Silvestrini 0007085

New Results
from
Fermilab Lattice
+ MILC
1602.03560

$$\left\{ \begin{array}{l} F_{B_s} \sqrt{\hat{B}_{B_s}} = (276.0 \pm 8.5 \text{ MeV}) \\ F_{B_d} \sqrt{\hat{B}_{B_d}} = (229.4 \pm 9.3 \text{ MeV}) \end{array} \right\}$$

Their ratio

$$\xi = 1.203 \pm 0.019$$



Blanke + AJB

(1602.04020)

Similar but less precise
results from ETM
(1308.1851)

$$\left\{ \frac{\Delta M_d}{\Delta M_s}, s_{\psi K_s} \right\} \Rightarrow \left\{ \begin{array}{l} \left| \frac{V_{ub}}{V_{cb}} \right| = 0.0864 \pm 0.0025 \\ \gamma = (62.7 \pm 2.1)^\circ \end{array} \right\}$$

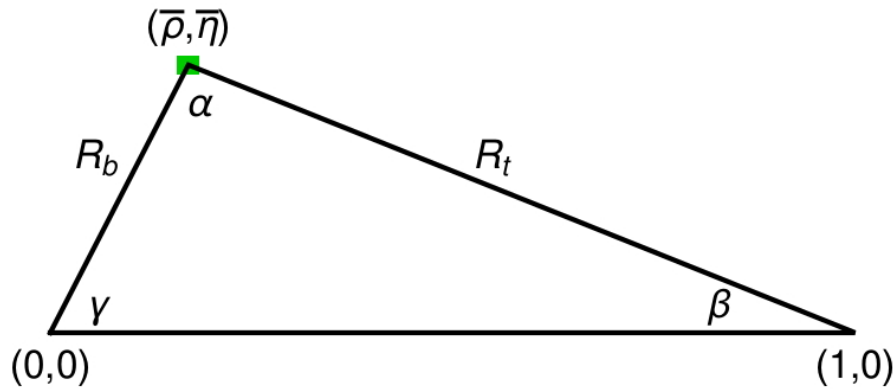
S₁: Strategy 1: ΔM_s determines $V_{cb} \Rightarrow |V_{cb}| = (39.5 \pm 1.3) \cdot 10^{-3}$

S₂: Strategy 1: ε_K determines $V_{cb} \Rightarrow |V_{cb}| = (43.4 \pm 1.2) \cdot 10^{-3}$



Tension between ΔM_K and ε_K

Universal Unitarity Triangle 2016



CMFV :

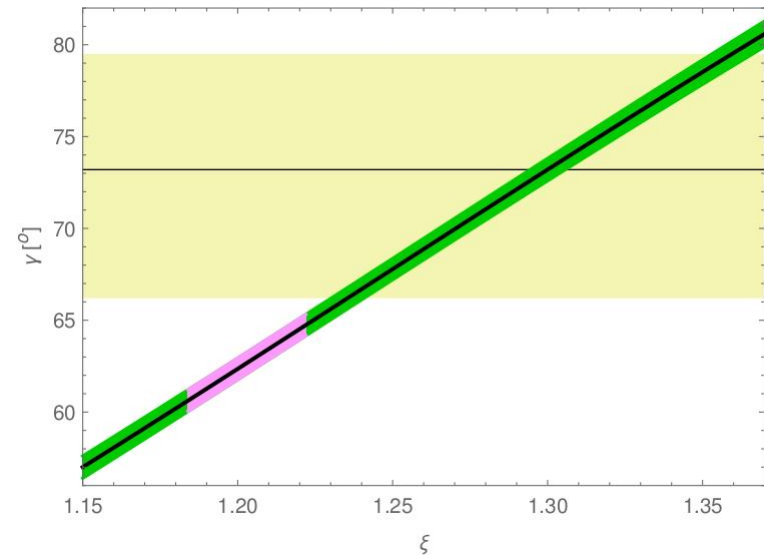
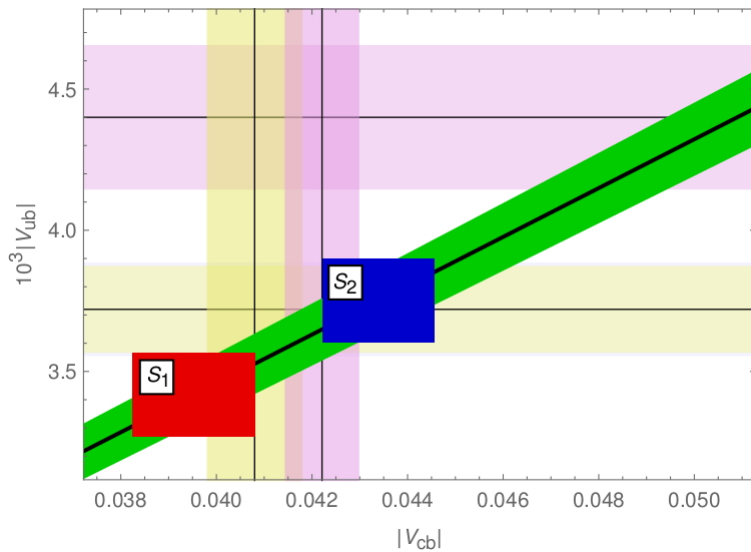
$$\bar{\rho} = 0.172 \pm 0.013$$

$$\bar{\eta} = 0.332 \pm 0.011$$

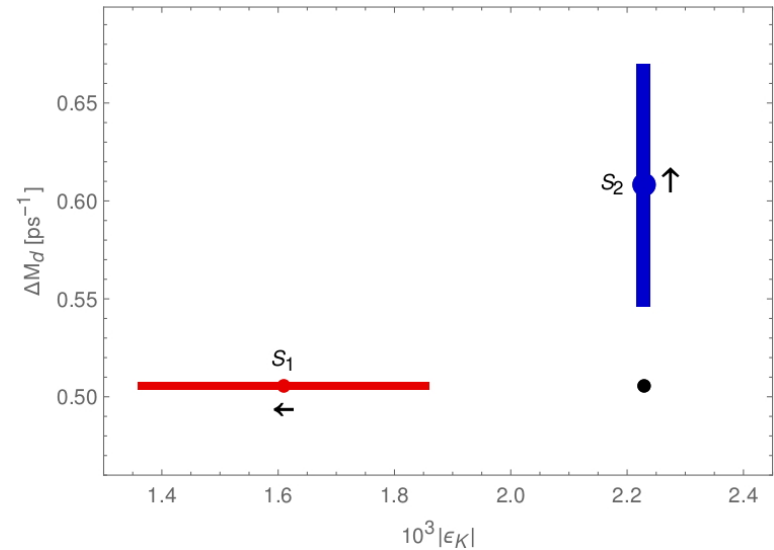
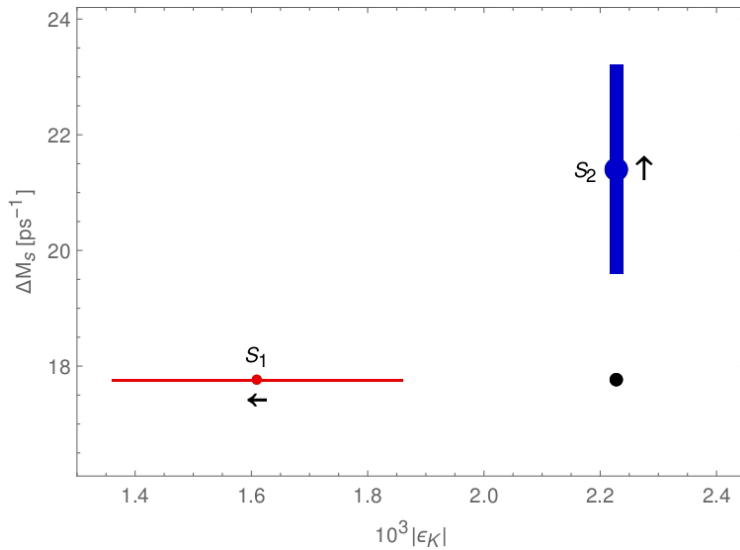
UT fit :

$$\bar{\rho} = 0.137 \pm 0.022$$

$$\bar{\eta} = 0.349 \pm 0.014$$



Tensions between $\Delta M_{d,s}$ and ε_K



**Constrained
MFV
CKM +
SM Operators**
 $S_{\text{Box}} > S_{\text{Box}}^{\text{SM}}$
**(Blanke, AJB
0610037)**

$$S_1: |\varepsilon_K| \leq (1.61 \pm 0.25) \cdot 10^{-3} \quad |\varepsilon_K^{\text{exp}}| = 2.23 \cdot 10^{-3}$$

$$S_2: \Delta M_s \geq (21.4 \pm 1.8) \text{ps}^{-1} \quad (\Delta M_s)^{\text{exp}} = 17.56 / \text{ps}$$

$$\Delta M_d \geq (0.608 \pm 0.062) \text{ps}^{-1} \quad (\Delta M_d)^{\text{exp}} = 0.506 / \text{ps}$$

Intermezzo

$$\mathbf{K}^+ \rightarrow \pi^+ \nu \bar{\nu} \text{ and } \mathbf{K}_L \rightarrow \pi^0 \nu \bar{\nu}$$

in the Standard Model

1503.02693



AJB



D. Buttazzo



J. Girrbach-Noe



R. Kneijens

$K^+ \rightarrow \pi^+ \nu \bar{\nu}$ and $K_L \rightarrow \pi^0 \nu \bar{\nu}$ in the SM

QCD Corrections:

NLO Buchalla, AJB; Misiak, Urban (93, 98)
 NNLO AJB, Gorbahn, Haisch, Nierste (2005)

NLO EW Corrections:

Large m_t : Buchalla, AJB (1997)
 Exact NLO (m_t): Brod, Gorbahn, Stamou (2010)
 " " (m_c): Brod, Gorbahn (2008)

LD Effects:

Isidori, Mescia, Smith (2005)
 Mescia, Smith (2007)

+ Isospin breaking corrections



TH uncertainties at the level of 2% in BR

Unique in Flavour Physics !!

But significant parametric uncertainties

due to $|V_{ub}|, |V_{cb}|, \gamma$

Data

$$\begin{aligned} \text{Br}(K^+ \rightarrow \pi^+ \nu \bar{\nu}) &= (17.3 \pm 11) \cdot 10^{-11} \\ \text{Br}(K_L \rightarrow \pi^0 \nu \bar{\nu}) &\leq 2.6 \cdot 10^{-8} \end{aligned}$$

CKM Uncertainties

AJB, Buttazzo,
Girrbach-Noe,
Knegjens
1503.02693

$$\text{Br}(\text{K}^+ \rightarrow \pi^+ \nu \bar{\nu}) = (8.39 \pm 0.30) \cdot 10^{-11} \left[\frac{|\mathbf{V}_{cb}|}{0.0407} \right]^{2.8} \left[\frac{\gamma}{73.2^\circ} \right]^{0.74}$$
$$\text{Br}(\text{K}_L \rightarrow \pi^0 \nu \bar{\nu}) = (3.36 \pm 0.05) \cdot 10^{-11} \left[\frac{|\mathbf{V}_{ub}|}{3.88 \cdot 10^{-3}} \right]^2 \left[\frac{|\mathbf{V}_{cb}|}{0.0407} \right]^2 \left[\frac{\sin \gamma}{\sin(73.2)} \right]^2$$

$$\text{Br}(\text{K}^+ \rightarrow \pi^+ \nu \bar{\nu}) = (8.39 \pm 0.58) \cdot 10^{-11} \left[\frac{\gamma}{73.2^\circ} \right]^{0.81} \left[\frac{\bar{\text{Br}}(\text{B}_s \rightarrow \mu^+ \mu^-)}{3.4 \cdot 10^{-9}} \right]^{1.42} \left[\frac{227.7}{F_{\text{B}_s}} \right]^{2.84}$$
$$\text{Br}(\text{K}^+ \rightarrow \pi^+ \nu \bar{\nu}) = (8.39 \pm 1.11) \cdot 10^{-11} \left[\frac{|\epsilon_K|}{2.23 \cdot 10^{-3}} \right]^{1.07} \left[\frac{\gamma}{73.2^\circ} \right]^{-0.11} \left[\frac{|\mathbf{V}_{ub}|}{3.88 \cdot 10^{-3}} \right]^{-0.95}$$

$$\text{Br}(\text{K}^+ \rightarrow \pi^+ \nu \bar{\nu}) = (8.4 \pm 1.0) \cdot 10^{-11}$$
$$\text{Br}(\text{K}_L \rightarrow \pi^0 \nu \bar{\nu}) = (3.4 \pm 0.6) \cdot 10^{-11}$$

Section 3

$$\varepsilon' / \varepsilon, \varepsilon_K, \mathbf{K} \rightarrow \pi \nu \bar{\nu}, \Delta \mathbf{M}_K$$

beyond SM

AJB (1601.00005)

Section 3

$\varepsilon'/\varepsilon, \varepsilon_K, \mathbf{K} \rightarrow \pi v \bar{v}, \Delta \mathbf{M}_K$

beyond SM

AJB (1601.00005)

**What are the implications
of NP in ε'/ε and ε_K on
 $\mathbf{K} \rightarrow \pi v \bar{v}$ and $\Delta \mathbf{M}_K$?**

Strategy

AJB (1601.00005)

$$\left(\varepsilon'/\varepsilon\right)^{\text{NP}} = \kappa_{\varepsilon'} \cdot 10^{-3}$$

$$0.5 \leq \kappa_{\varepsilon'} \leq 1.5$$

(Im)

$$\varepsilon_{\text{K}}^{\text{NP}} = \kappa_{\varepsilon} \cdot 10^{-3}$$

$$0.1 \leq \kappa_{\varepsilon} \leq 0.4$$

(Im, Re)

In some models
 $\text{K}_L \rightarrow \mu^+ \mu^-$
 more important
 than ε_{K}

Re and Im Parts: \mathbf{Z} and \mathbf{Z}' Couplings

$$\Delta_{\text{L}}^{\text{sd}}(\mathbf{Z}), \Delta_{\text{R}}^{\text{sd}}(\mathbf{Z})$$

$$\Delta_{\text{L}}^{\text{sd}}(\mathbf{Z}'), \Delta_{\text{R}}^{\text{sd}}(\mathbf{Z}')$$

$$\text{K}^+ \rightarrow \pi^+ \nu \bar{\nu}, \text{K}_L \rightarrow \pi^0 \nu \bar{\nu}, \text{K}_L \rightarrow \mu^+ \mu^-, \Delta \mathbf{M}_{\text{K}}$$

$$(\text{Re, Im}) \quad (\text{Im}) \quad (\text{Re}) \quad (\text{Im, Re})$$

Basic Structure of NP Contributions

AJB (1601.00005)

$$\begin{aligned}
 (\varepsilon'/\varepsilon)^{\text{NP}} &\rightarrow \text{Im} & \varepsilon_{\text{K}}^{\text{NP}} &\rightarrow \text{Im} \cdot \text{Re} \\
 (\kappa_{\varepsilon'} \geq 0.5) & & (\kappa_{\varepsilon} \geq 0.1) & \\
 \Delta M_{\text{K}}^{\text{NP}} &\sim \left[(\text{Re})^2 - (\text{Im})^2 \right]
 \end{aligned}$$

Dominance of $Q_6 (Q_6')$ \Rightarrow $\text{Im} \gg \text{Re} \Rightarrow \left\{ \Delta M_{\text{K}}^{\text{NP}} < 0 \right\}$ (Z)
 (large)

Dominance of $Q_8 (Q_8')$ \Rightarrow $\text{Re} \gg \text{Im} \Rightarrow \left\{ \Delta M_{\text{K}}^{\text{NP}} > 0 \right\}$ (Z/Z')
 (small)



Implications for

$$\mathbf{R}_{+}^{\nu\bar{\nu}} = \frac{\text{Br}(\text{K}^+ \rightarrow \pi^+ \nu\bar{\nu})}{\text{Br}(\text{K}^+ \rightarrow \pi^+ \nu\bar{\nu})_{\text{SM}}}$$

(Re, Im)

$$\mathbf{R}_{0}^{\nu\bar{\nu}} = \frac{\text{Br}(\text{K}_L \rightarrow \pi^0 \nu\bar{\nu})}{\text{Br}(\text{K}_L \rightarrow \pi^0 \nu\bar{\nu})_{\text{SM}}}$$

(Im)

Lessons on ε'/ε , ε_K , $K \rightarrow \pi\nu\bar{\nu}$: BSM

AJB: 1601.00005

Lesson 1

**Do not expect much from MFV
(tensions cannot be removed) 1507.08672**

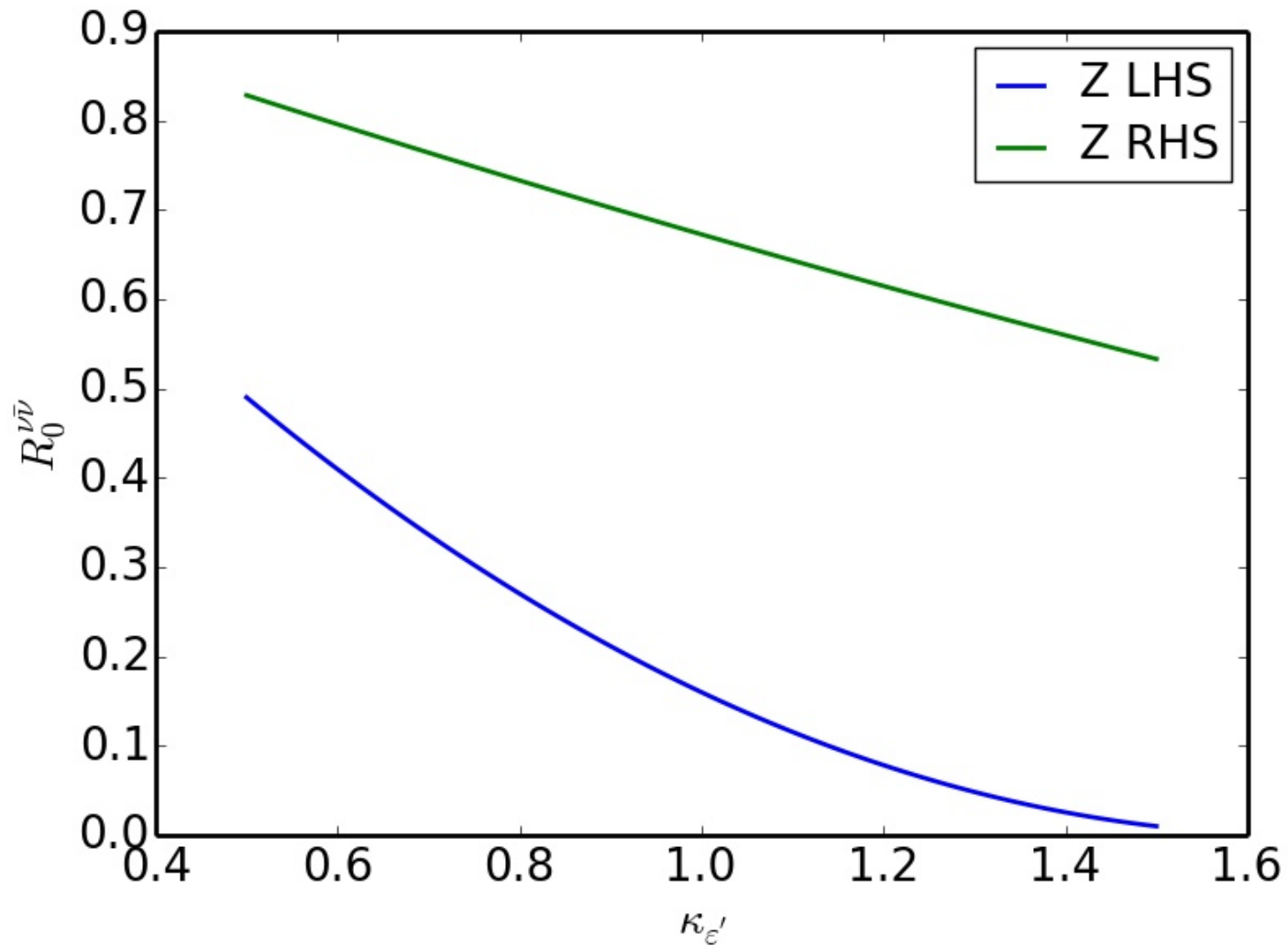
Lesson 2

**Tree-Level Z with LH or RH FCNC currents
(Anticorrelation of ε'/ε and $K_L \rightarrow \pi^0 \nu\bar{\nu}$)
 $K^+ \rightarrow \pi^+ \nu\bar{\nu}$ can be significantly enhanced**

LH	$R_+^{\nu\bar{\nu}} < 2$
RH	$R_+^{\nu\bar{\nu}} < 5.7$

Q_8
Q_8'

Z with LH or RH Flavour Violating Couplings



Lesson 3

Tree-Level Z with LH + RH FCNC currents

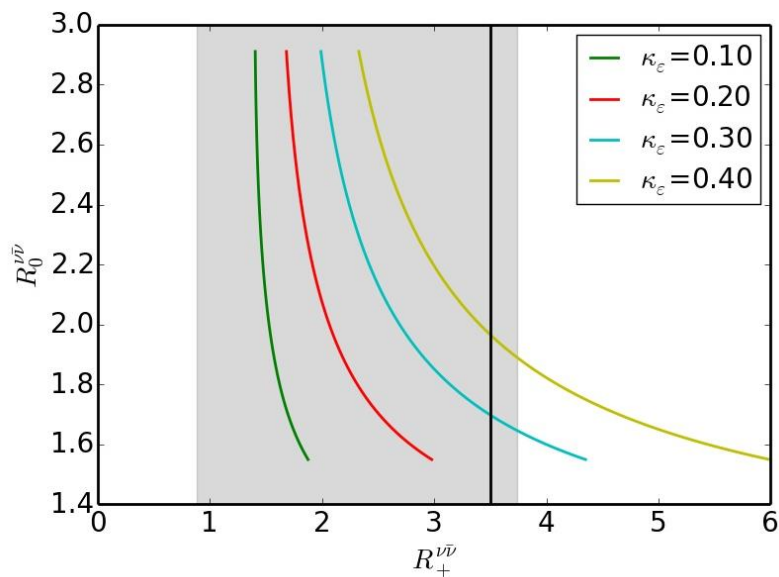
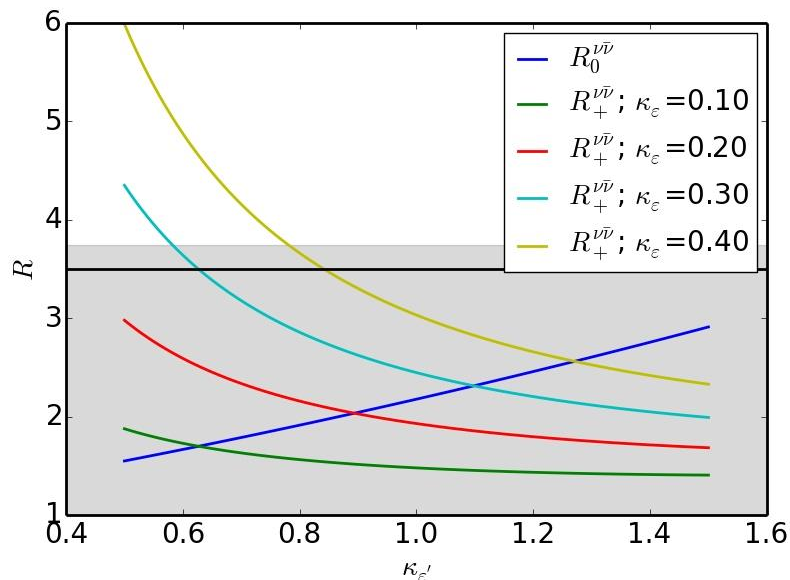
$\varepsilon'/\varepsilon, \varepsilon_K, K^+ \rightarrow \pi^+ \nu \bar{\nu}$ and $K_L \rightarrow \pi^0 \nu \bar{\nu}$
can be simultaneously enhanced

**Correlation depends on hierarchy
between $\text{Re}\Delta_{L,R}$ and $\text{Im}\Delta_{L,R}$**

Z with LH and RH Flavour Violating Couplings

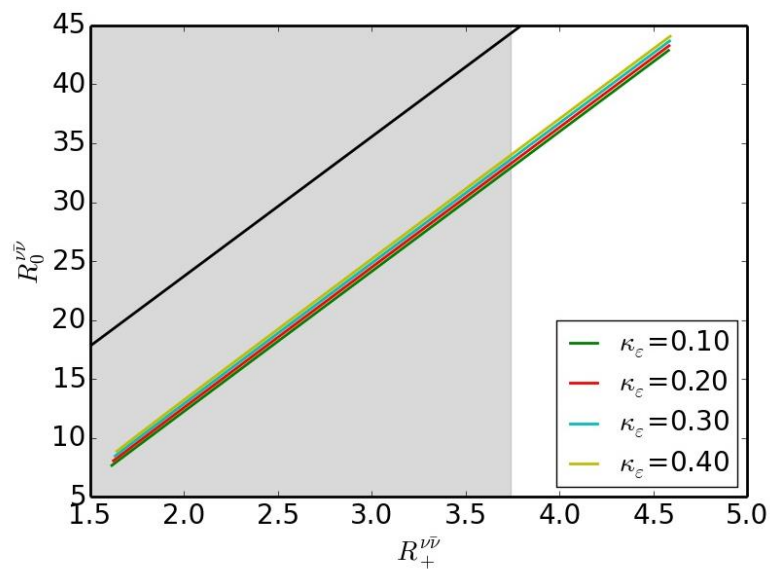
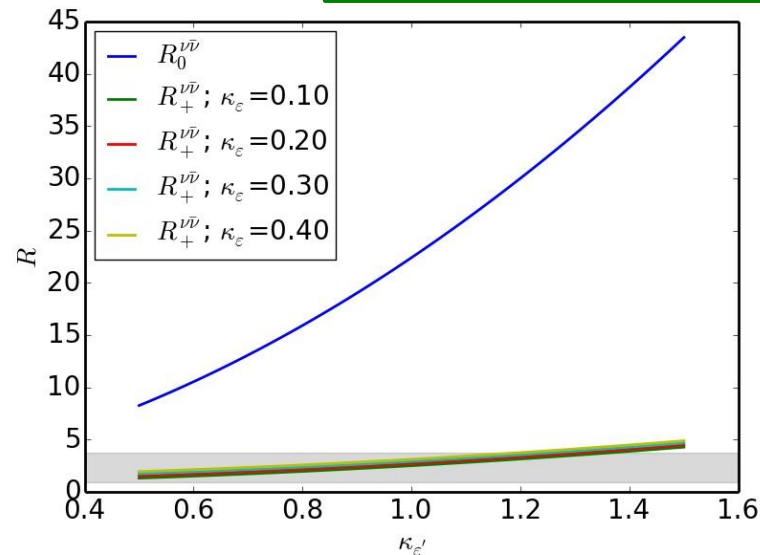
Example 1

$\text{Im} \Delta_{L,R} < \text{Re} \Delta_{L,R}$



Example 2

$\text{Im} \Delta_{L,R} \gg \text{Re} \Delta_{L,R}$



Lesson 4

**Correlation between ε'/ε , $K \rightarrow \pi\nu\bar{\nu}$
in Z' scenarios depends on whether
QCP Penguin (Q_6) or EWP (Q_8) dominates
NP in ε'/ε**

Z' Scenarios with LH Couplings $\Delta_L^{sd}(Z')$

AJB (1601.00005)

Dominance
of QCD
Penguins (Q_6)
in ε'/ε

- Strong correlation between K^+ and K_L on the branch parallel to GN bound
- Very large effects in K_L , moderate in K^+
- $(\Delta M_K)^{NP} < 0$ (could be 20%) ε_K anomaly can be solved

Dominance
of electroweak
Penguins (Q_8)
in ε'/ε

- Both enhanced but anticorrelated

$K_L \uparrow \quad K^+ \downarrow \quad \text{with } \kappa_{\varepsilon'} \uparrow$

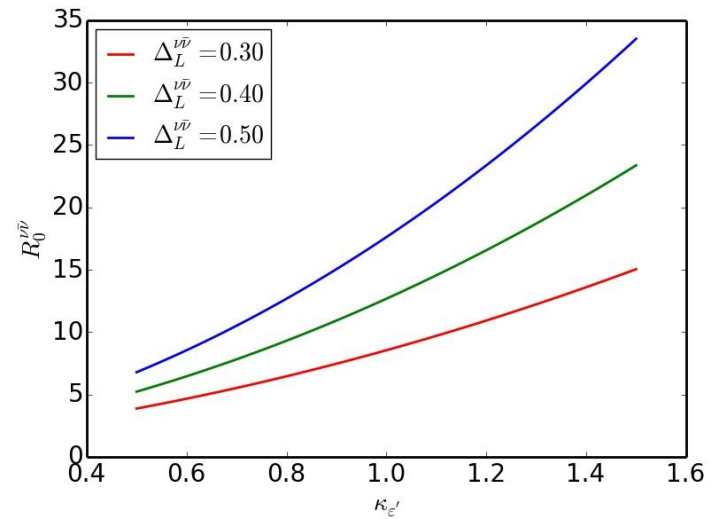
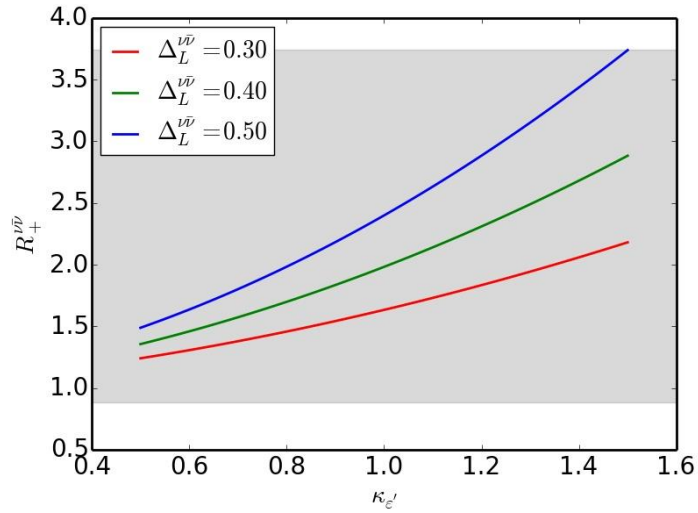
$(K^+ \uparrow \text{ with } \kappa_\varepsilon \uparrow)$ Only (20-40)% effects

- $(\Delta M_K)^{NP} > 0$ (below 10%) ε_K anomaly can be solved

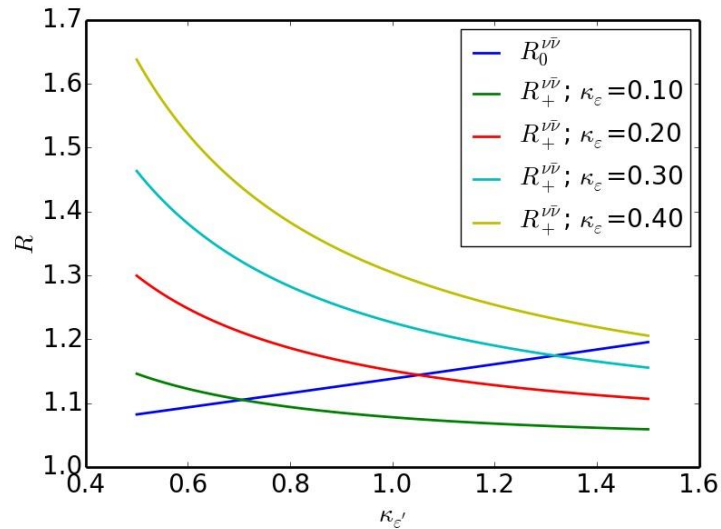
Pattern for
 $\Delta_R^{q\bar{q}}(Z') \approx 0(1)$
in ε'/ε

$M_{Z'} = 3 \text{ TeV}$

QCD Penguin (Q_6)

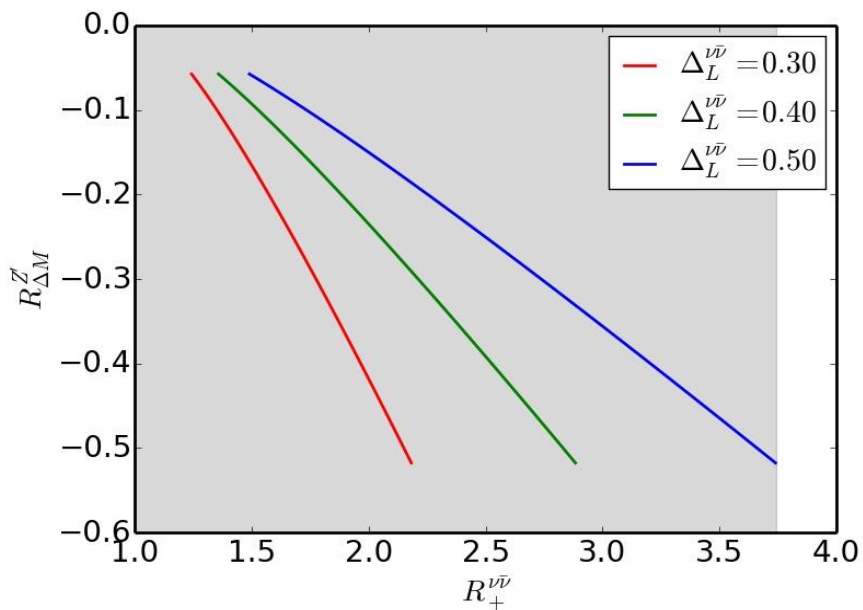
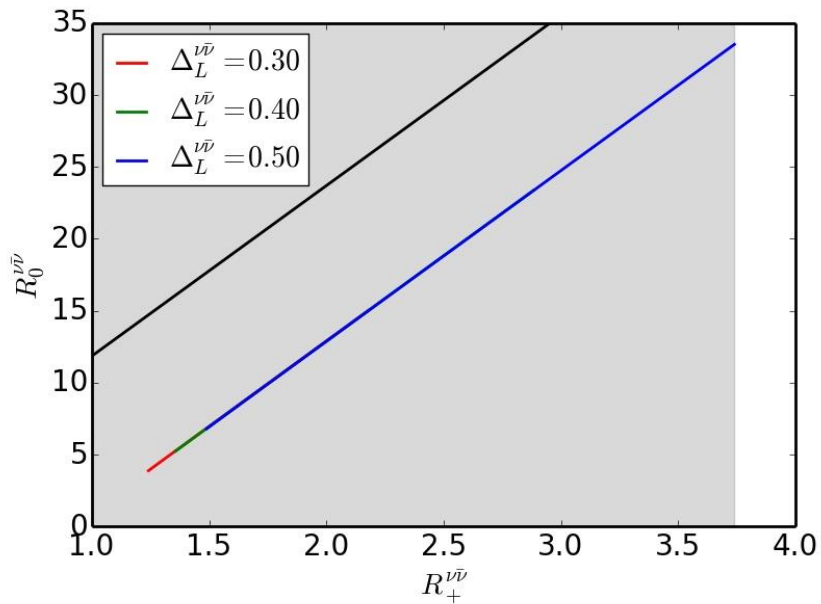


Electroweak Penguin (Q_8)

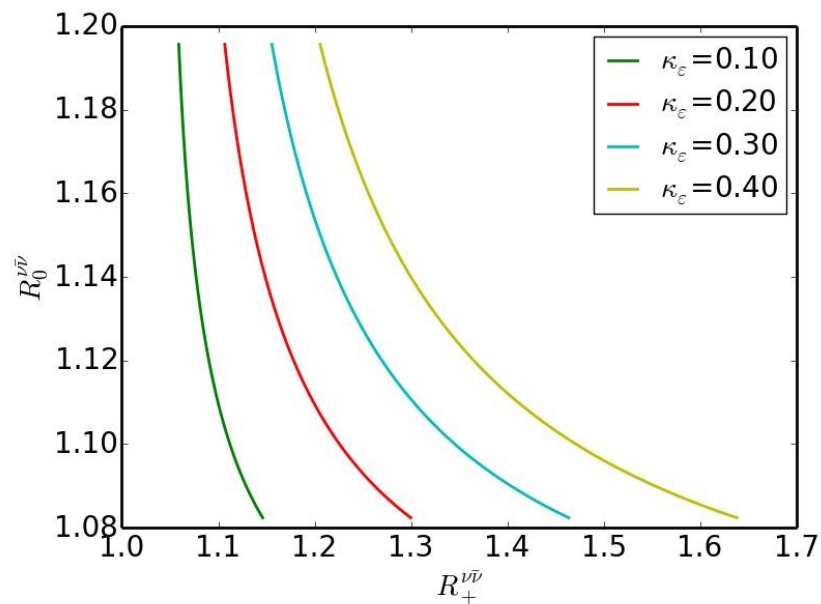


(Z)

QCDP (Q₆)



EWP (Q₈)



($R_{\Delta M}^Z > 0$ but small)

(Z)

Section 4

**Highlights from 331, LHT,
Vector-Like Quark Models**

$\varepsilon'/\varepsilon + K \rightarrow \pi\nu\bar{\nu}$ beyond SM



AJB



Fulvia de Fazio



Jennifer Girrbach-Noe

Z, Z' 331

1404.3824, ...
1311.6729



AJB



Dario Buttazzo



Rob Kneijens

Simplified NP
Models
1507.08672



Monika Blanke



AJB



Stefan Recksiegel

LHT
1507.0631

Most Recent



AJB



Fulvia de Fazio

331 models facing
 $\Delta M_{s,d} \leftrightarrow \varepsilon_K$ tension

$\varepsilon'/\varepsilon, B_s \rightarrow \mu^+ \mu^-$,

$B \rightarrow K^* \mu^+ \mu^-$

Model with Vektor-like Quarks



Christoph Bobeth



AJB



Alejandro Celis



Martin Jung

331 Models Facing ε'/ε Anomaly

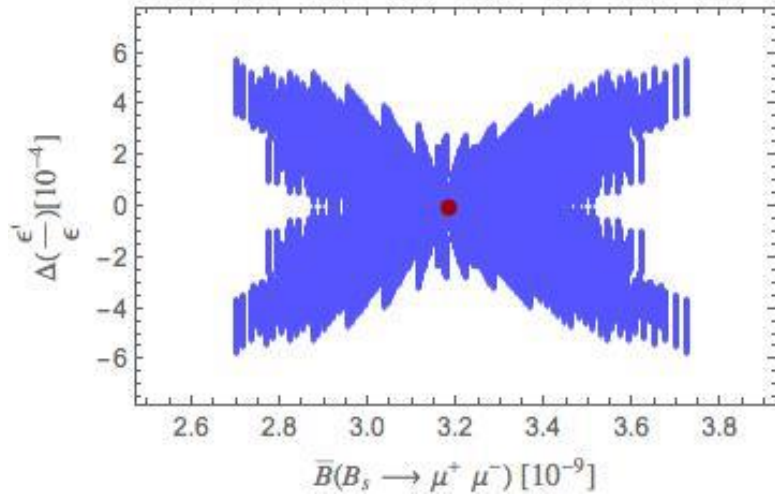
AJB, De Fazio 1512.02869, 1604.02344

- 1.** $\kappa_{\varepsilon'} \leq 0.8$ (only 3 models can reach upper bound)
- 2.** None of them can explain suppressions of C_9 ($B \rightarrow K(K^*)\mu^+\mu^-$) and $B_s \rightarrow \mu^+\mu^-$ simultaneously.
None R_K
- 3.** Small NP effects in $K^+ \rightarrow \pi^+ \nu\bar{\nu}$ and $K_L \rightarrow \pi^0 \nu\bar{\nu}$

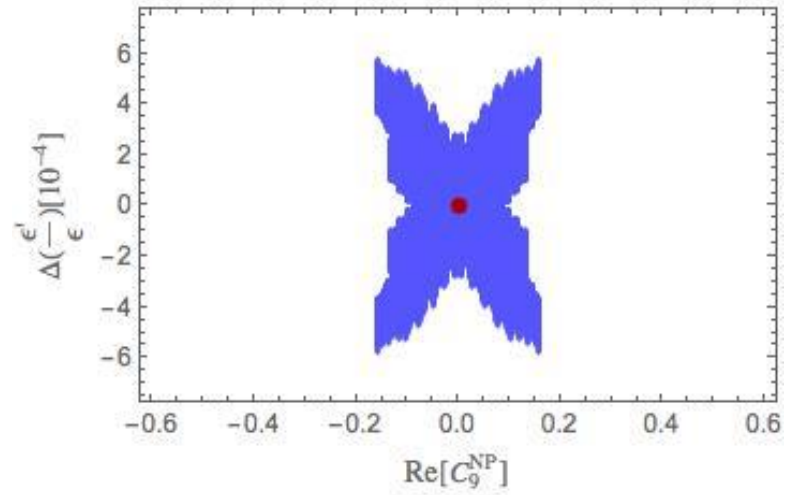
Correlations in Favorite 331 Models

(AJB+De Fazio, 1604.02344)

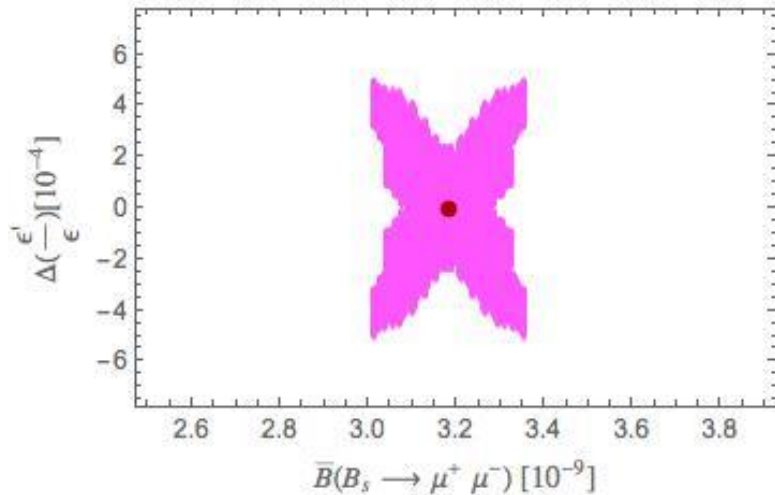
M8



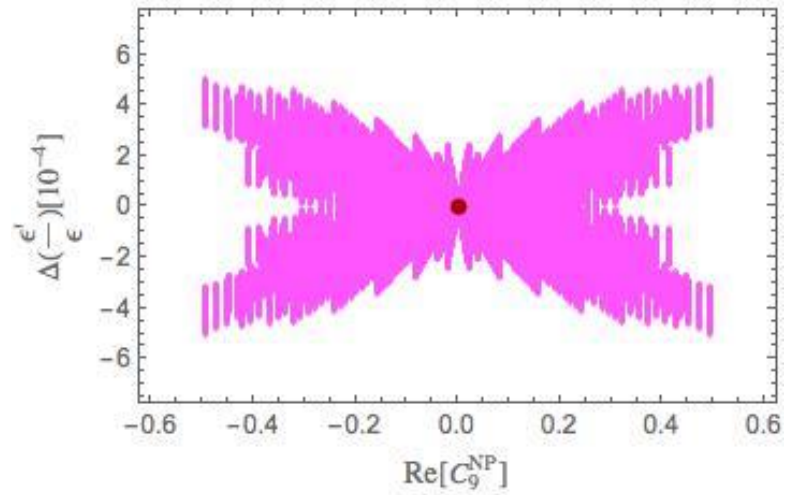
M8



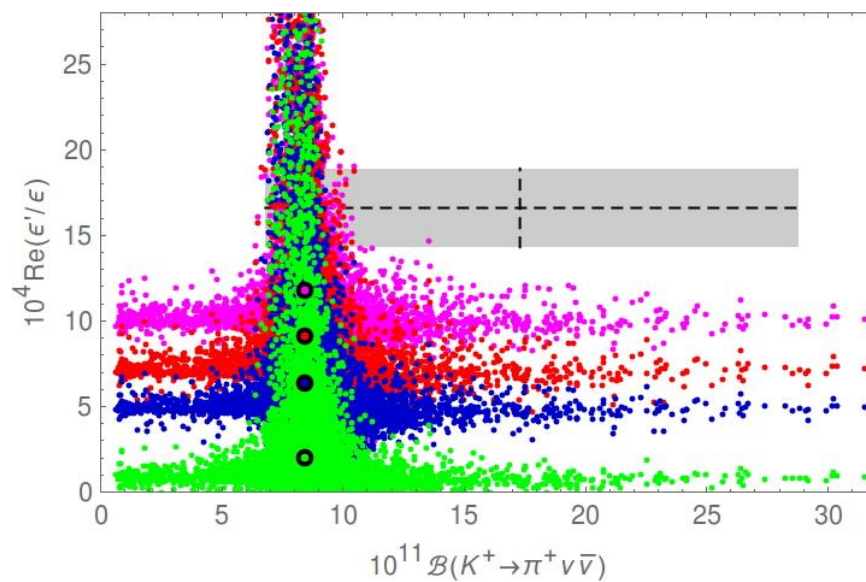
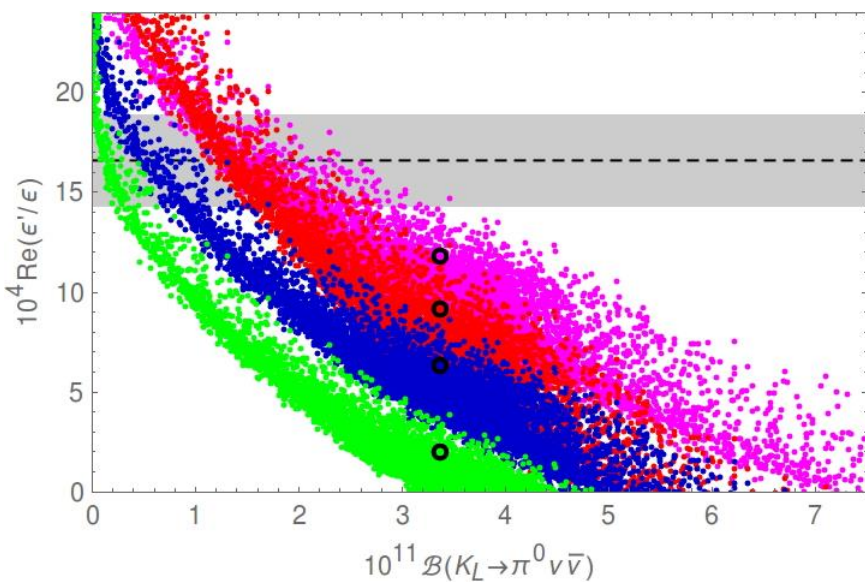
M16



M16



LHT : Blanke, AJB, Recksiegel (1507.06316)



~~$$\left(B_6^{(1/2)} = 1.0, B_8^{(3/2)} = 0.76 \right)$$~~

$$\left(B_6^{(1/2)} = 1.0, B_8^{(3/2)} = 1.0 \right)$$

$$\left(B_6^{(1/2)} = 0.75, B_8^{(3/2)} = 0.76 \right)$$

$$\left(B_6^{(1/2)} = 0.57, B_8^{(3/2)} = 0.76 \right)$$

(Violates
Large N bound)

Supersymmetric Explanation of ε'/ε and ε_K

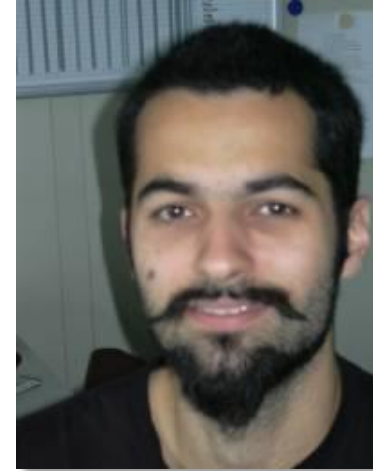
Teppei Kitahara



Ulrich Nierste



Paul Tremper 1604.07400



ε'/ε anomaly can be explained in the MSSM with squark masses above 3 TeV being consistent with ε_K without fine-tuning of CP phases or other parameters.

2018 Vision

$$\text{Br}(\text{K}^+ \rightarrow \pi^+ \nu \bar{\nu}) = (18.0 \pm 2.0) \cdot 10^{-11} \quad (\text{NA62})$$

$$\kappa_{\varepsilon'} \approx 1.0$$

Would point
towards :

Z with LH + RH couplings
Z'(QCDP) with $Z' q \bar{q} \approx 0(1)$
Z'(EWP) with $Z' q \bar{q} \approx 10^{-2}$

Open Questions to be answered hopefully in this Decade

- 1.** What is $\text{Br}(\text{K}^+ \rightarrow \pi^+ \nu \bar{\nu})$ from NA62?
- 2.** What is $\text{Br}(\text{K}_L \rightarrow \pi^0 \nu \bar{\nu})$ from KOPIO?
- 3.** What is the value of $\kappa_{\varepsilon'}$? (Lattice, CKM, NNLO)
- 4.** What is the value of κ_{ε} ? (CKM, η_1)
- 5.** What is $(\Delta M_K)^{\text{SM}}$? (Lattice)
- 6.** Does NP contribute to $\text{Re}A_0$ at 10-20% level? (Lattice)
(see AJB, De Fazio, Girrbach: 1404.3824)
- 7.** Do Z' , G' or other new particles exist?

**Exciting Times are just
ahead of us !!!**

**Exciting Times are just
ahead of us !!!**

Thank You !

Backup

RBC-UK QCD

$$\varepsilon'/\varepsilon = (1.4 \pm 7.0) \cdot 10^{-4}$$

$$\left(\frac{\text{Re } A_0}{\text{Re } A_2} \right) = 31.0 \pm 6.6$$

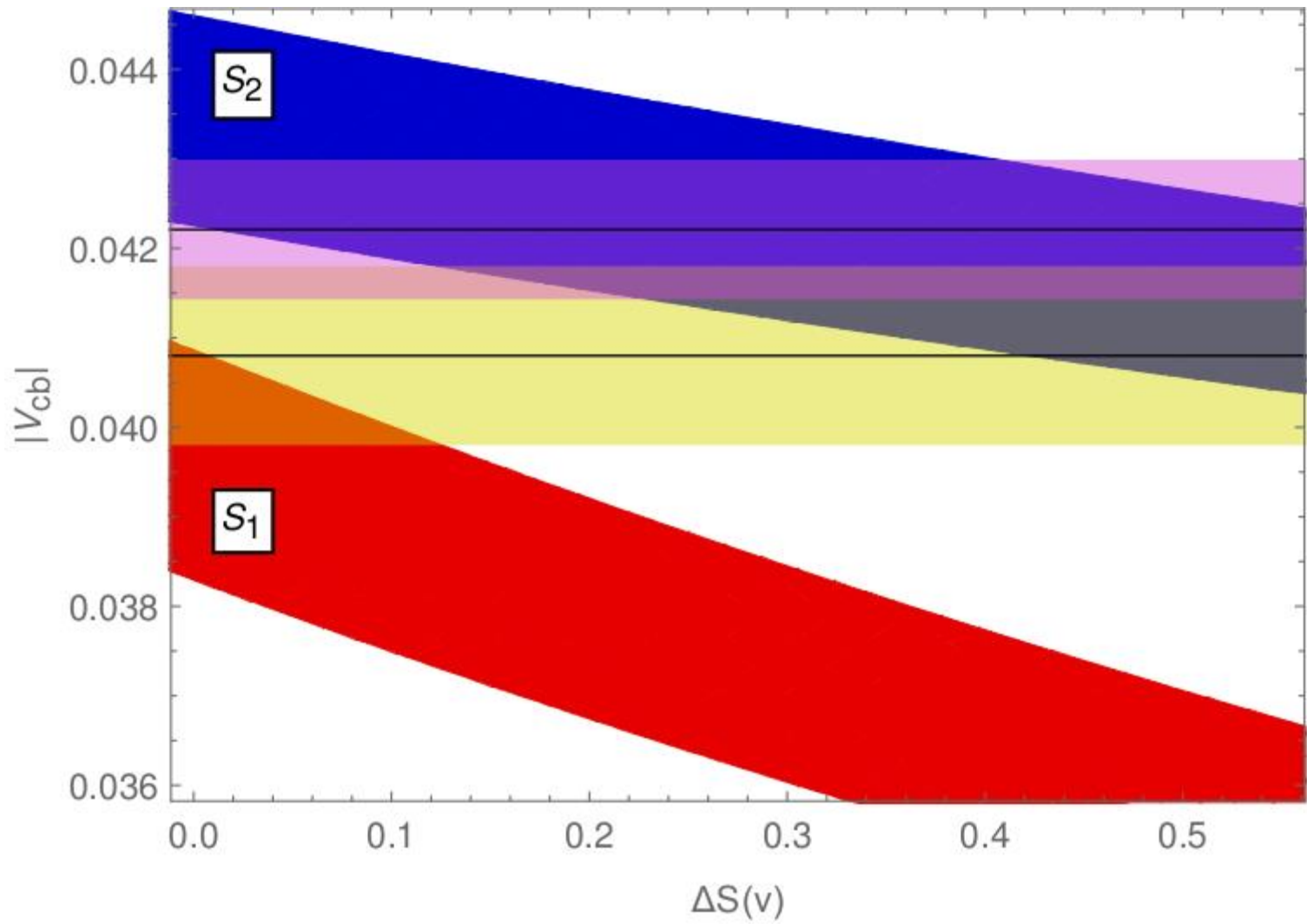
$$(\varepsilon'/\varepsilon)_{\text{exp}} = (16.6 \pm 2.3) \cdot 10^{-4}$$

$$\left(\frac{\text{Re } A_0}{\text{Re } A_2} \right)_{\text{exp}} = 22.4$$

Large N

$$(\varepsilon'/\varepsilon) < (8.6 \pm 3.2) \cdot 10^{-4}$$

$$\left(\frac{\text{Re } A_0}{\text{Re } A_2} \right) = 16.0 \pm 1.5$$



Large N Approach

AJB, Gérard (2015)

vs

Lattice

$$\hat{B}_K = 0.73 \pm 0.02$$

$$(\hat{B}_K \leq 0.75)$$

$$B_6^{(1/2)} = 1 - 0(1/N)$$

$$B_8^{(3/2)} = 1 - 0(1/N)$$

$$\frac{\text{Re } A_0}{\text{Re } A_2} = 16.0 \pm 1.5$$

$$B_8^{(1/2)} = 1 - 0(1/N^2)$$

Exp
22.4

$$\hat{B}_K = 0.766 \pm 0.010 \text{ (FLAG)}$$

(will go down with new results)

$$B_6^{(1/2)} = 0.57 \pm 0.19$$

$$B_8^{(3/2)} = 0.76 \pm 0.05$$

$$\frac{\text{Re } A_0}{\text{Re } A_2} = 31.0 \pm 6.6$$

$$B_8^{(1/2)} = 1.0 \pm 0.2$$

RBC-UKQCD

$\Delta I = 1/2$ Rule

Large N Approach

AJB, Gérard (2015)

vs

Lattice

$$\hat{B}_K = 0.73 \pm 0.02$$
$$(\hat{B}_K \leq 0.75)$$

$$B_6^{(1/2)} \leq B_8^{(3/2)}$$

$$B_8^{(3/2)} = 0.80 \pm 0.10$$

$$\frac{\text{Re } A_0}{\text{Re } A_2} = 16.0 \pm 1.5$$

$$B_8^{(1/2)} = 1 - 0(1/N^2)$$

Exp
22.4



$\Delta I = 1/2$ Rule

$$\hat{B}_K = 0.766 \pm 0.010 \text{ (FLAG)}$$

(will go down with new results)

$$B_6^{(1/2)} = 0.57 \pm 0.19$$

$$B_8^{(3/2)} = 0.76 \pm 0.05$$

$$\frac{\text{Re } A_0}{\text{Re } A_2} = 31.0 \pm 6.6$$

$$B_8^{(1/2)} = 1.0 \pm 0.2$$

RBC-UKQCD

Motivations for New Analysis

1. NA62 in progress: 10% measurement of $K^+ \rightarrow \pi^+ \nu \bar{\nu}$ in 2018.

2. Stress CKM uncertainties in $\text{Br}(K^+ \rightarrow \pi^+ \nu \bar{\nu})$, $\text{Br}(K_L \rightarrow \pi^0 \nu \bar{\nu})$

3. Point out correlation between

$K^+ \rightarrow \pi^+ \nu \bar{\nu}$, $B_s \rightarrow \mu^+ \mu^-$ and γ
(NA62) (LHCb+CMS) (LHCb)

Basically
no CKM
uncertainties

4. Update correlation between

$K^+ \rightarrow \pi^+ \nu \bar{\nu}$, $K_L \rightarrow \pi^0 \nu \bar{\nu}$ and β

(Buchalla, AJP, 94)
(AJP, Fleischer, 00)

5. Use most recent lattice input for CKM

6. Provide the present best value in SM

$K^+ \rightarrow \pi^+ \nu \bar{\nu}$ and $K_L \rightarrow \pi^0 \nu \bar{\nu}$ in simplified NP Models

Review Mod. Phys.: AJB, Schwab, Uhlig (2008) (0405132)
AJB, Buttazzo, Kneijens: hep-ph-1507.08672

MFV : 20-30% effects, strong correlation between K^+ and K_L (Z, Z')

$U(2)^3$: Larger effects in the absence of $B_s \rightarrow \mu^+ \mu^-$ constraint

No MFV : Correlation depends on the presence or absence of ε_K constraint, size on ε'/ε , $K_L \rightarrow \mu^+ \mu^-$

FCNCs Z : Enhancements by factors 2-3 over SM still possible (ε'/ε constraint important)

FCNCs Z' : Still larger enhancements possible as ε'/ε constraint can be eliminated in a model independent analysis but not in specific models with known flavour diagonal quark couplings.

More info
in BBK

see Rob Kneijens (Moriond) 1505.04928

Different Patterns of Flavour Violation

Z with LH couplings: $\Delta_L^{sd}(Z)$

Q₈ EWP

AJB (1601.00005)

- Anticorrelation of ε'/ε and $K_L \rightarrow \pi^0 \nu \bar{\nu}$
- Strong suppression of $\text{Br}(K_L \rightarrow \pi^0 \nu \bar{\nu})$
- $\text{Br}(K^+ \rightarrow \pi^+ \nu \bar{\nu}) \leq 2 \text{Br}(K^+ \rightarrow \pi^+ \nu \bar{\nu})^{\text{SM}}$
- NP effects in ΔM_K and ε_K very small

} No specific correlation

($K_L \rightarrow \mu^+ \mu^-$ constraint more important)

Z with RH couplings: $\Delta_R^{sd}(Z)$

- Anticorrelation of ε'/ε and $K_L \rightarrow \pi^0 \nu \bar{\nu}$
- Moderate suppression of $\text{Br}(K_L \rightarrow \pi^0 \nu \bar{\nu})$
- $\text{Br}(K^+ \rightarrow \pi^+ \nu \bar{\nu}) \leq 6 \text{Br}(K^+ \rightarrow \pi^+ \nu \bar{\nu})^{\text{SM}}$
- NP effects in ΔM_K and ε_K very small

Unless Loop effects important

Q₈ EWP

Z with LH and RH Couplings $\Delta_{L,R}^{sd}(\mathbf{Z})$

AJB (1601.00005)

New Features

ε_K constraint dominates over $K_L \rightarrow \mu^+ \mu^-$
 because of LR operators \rightarrow " ε_K anomaly"
 can be resolved.

Possibility of simultaneous enhancements of

$$\varepsilon'/\varepsilon, \varepsilon_K, K_L \rightarrow \pi^0 \nu \bar{\nu}, K^+ \rightarrow \pi^+ \nu \bar{\nu}$$

Example 1

$$\text{Im} \Delta_{L,R} < \text{Re} \Delta_{L,R}$$

Both $K_L \rightarrow \pi^0 \nu \bar{\nu}$ and $K^+ \rightarrow \pi^+ \nu \bar{\nu}$ enhanced
 but anticorrelated

$$K_L \uparrow \quad K^+ \downarrow \quad \text{with } \kappa_{\varepsilon'} \uparrow$$

$$(K^+ \uparrow \text{ with } \kappa_{\varepsilon} \uparrow)$$

NP Effects
 in ΔM_K
 small

Example 2

$$\text{Im} \Delta_{L,R} \gg \text{Re} \Delta_{L,R}$$

$$K_L \uparrow \quad K^+ \uparrow \quad \text{with } \kappa_{\varepsilon'} \uparrow$$

(no dependence on κ_{ε})

Correlation between K_L and K^+

On the branch parallel to Grossmann-Nir Bound

What about $\Delta I = 1/2$ Rule?

$$\frac{\text{Re } A_0}{\text{Re } A_2} \approx 22.4$$

Since 1955

Gell-Mann
Pais

1986, 2014

Large N including
I/N corrections

Quark Evolution $1 \text{ GeV} \leq \mu \leq M_W$
Meson Evolution $0 \leq \mu \leq 1 \text{ GeV}$

Correct value
of $\text{Re}A_2$

$$\left(\frac{\text{Re } A_0}{\text{Re } A_2} \right)_{\text{I/N}} \approx 16.0 \pm 1.5^*)$$

Dominance
of current-
current
operators

Correct value
of $\text{Re}A_2$

$$\left(\frac{\text{Re } A_0}{\text{Re } A_2} \right)_{\text{Lattice}} \approx 31 \pm 7$$

RBC-QCD
(2013, 2015)

*) G' with particular couplings ($M_{G'} \approx 3.5 \text{ TeV}$)
could be responsible for the missing piece

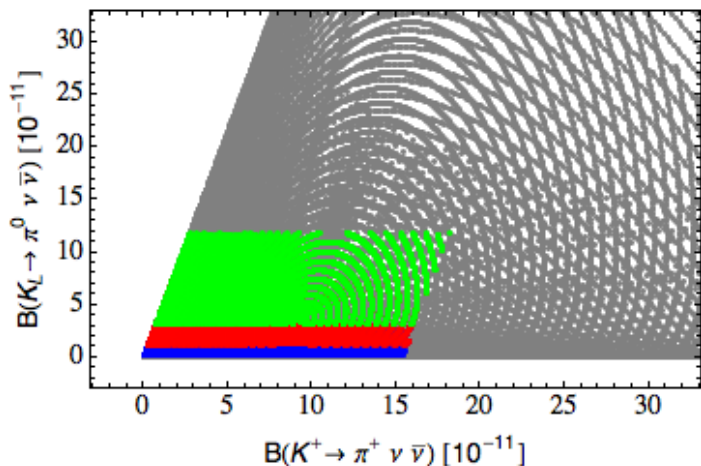
AJB
De Fazio
Girrbach-Noe
1404.3824

Z with FCNCs at Work

AJB, de Fazio,
Girrbach-Noe
1404.3824

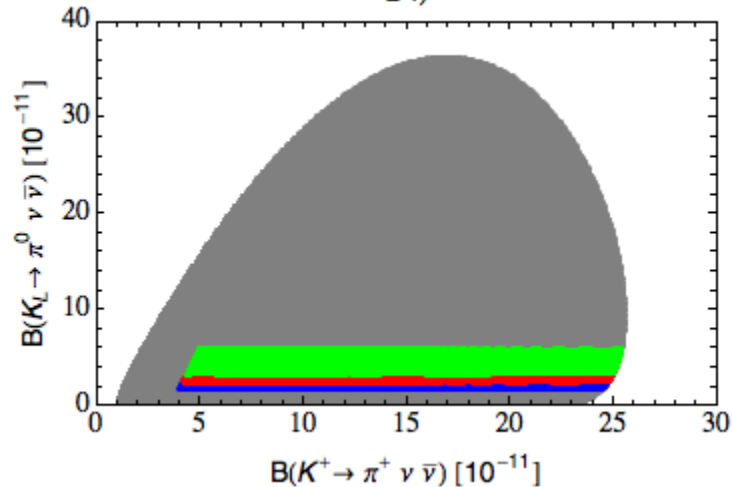
LHS

B f)



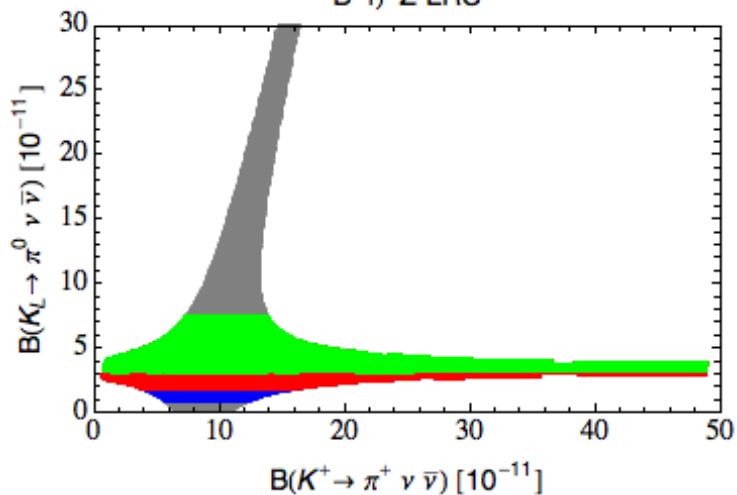
RHS

B f)



LRS

B f) Z LRS



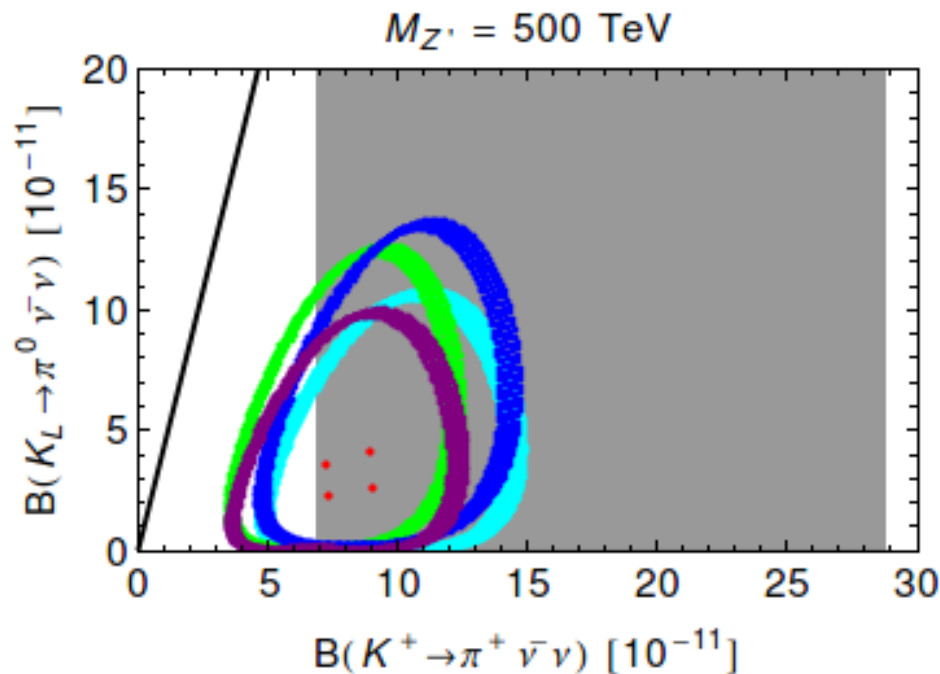
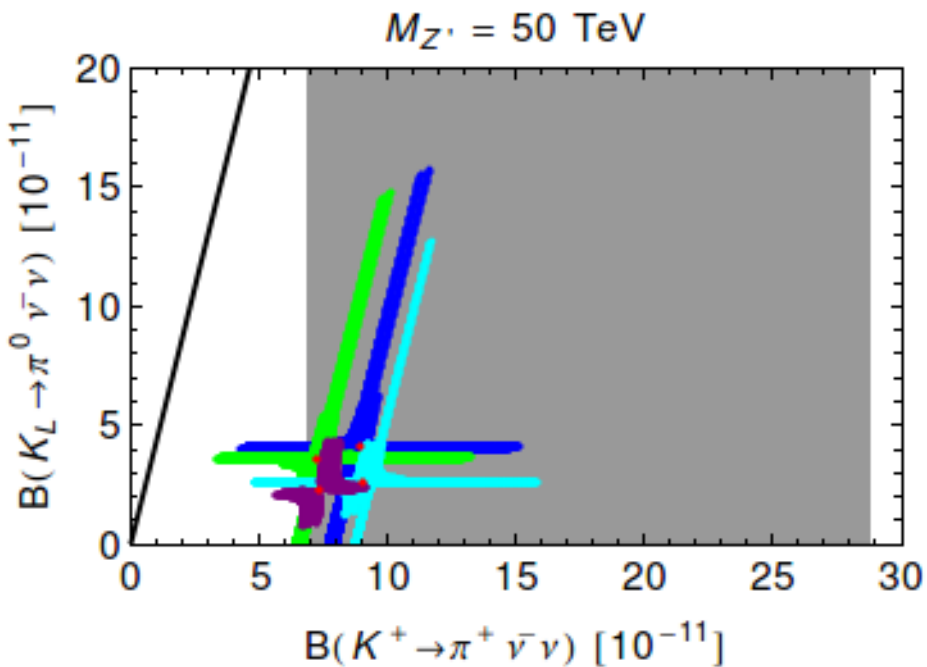
$\epsilon_K, \Delta M_K$ constraint

ϵ'/ϵ
+
 $K_L \rightarrow \mu^+ \mu^-$

$B_6 = 1.25$
 $B_6 = 1.00$
 $B_6 = 0.75$

Heavy Z' at Work

AJB, Buttazzo, Girrbach-Noe, Kneijens, 1407.0728



ϵ_K constraint

General discussion:
Blanke 0904.2528

No ϵ_K constraint

Can we reach Zeptouniverse through Quark Flavour Physics ?

(Z')

AJB, Buttazzo, Girschbach-Noe, Kneijens, 1407.0728

If only left-handed
or only right-handed
couplings present in NP

:

Only with K rare Decays
 $B_s \sim 15 \text{ TeV}, B_d \sim 15 \text{ TeV}$

If both LH and RH
present but
 $g_L^{ij} \ll g_R^{ij}$ or $g_L^{ij} \gg g_R^{ij}$

:

$K \rightarrow \pi \nu \bar{\nu}$: $\Lambda_{\text{NP}}^{\text{max}} \simeq 2000 \text{ TeV}$
 B_d : $\Lambda_{\text{NP}}^{\text{max}} \simeq 160 \text{ TeV}$
 B_s : $\Lambda_{\text{NP}}^{\text{max}} \simeq 160 \text{ TeV}$

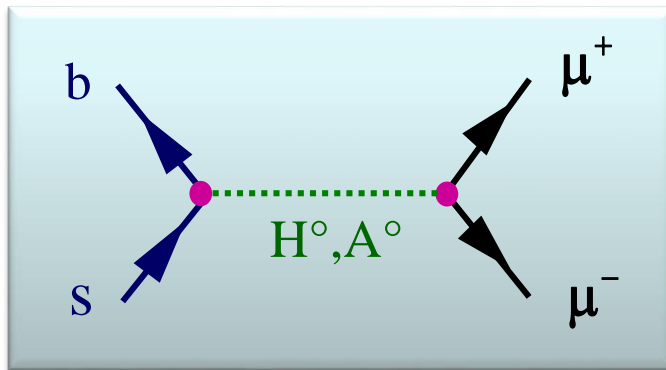
Yes we can !!

Can we reach Zeptouniverse through S and P

AJB, Buttazzo, Girrbach-Noe, Kneijens, 1407.0728

Yes :

$$\mathbf{B}_{s,d} \rightarrow \mu^+ \mu^-$$



S : ≈ 350 TeV

P : ≈ 700 TeV

Pseudoscalars more powerful than scalars because of the interference with SM contribution

Similar to $K \rightarrow \pi \nu \bar{\nu}$ (**Z**):
No tuning necessary to reach Zeptouniverse

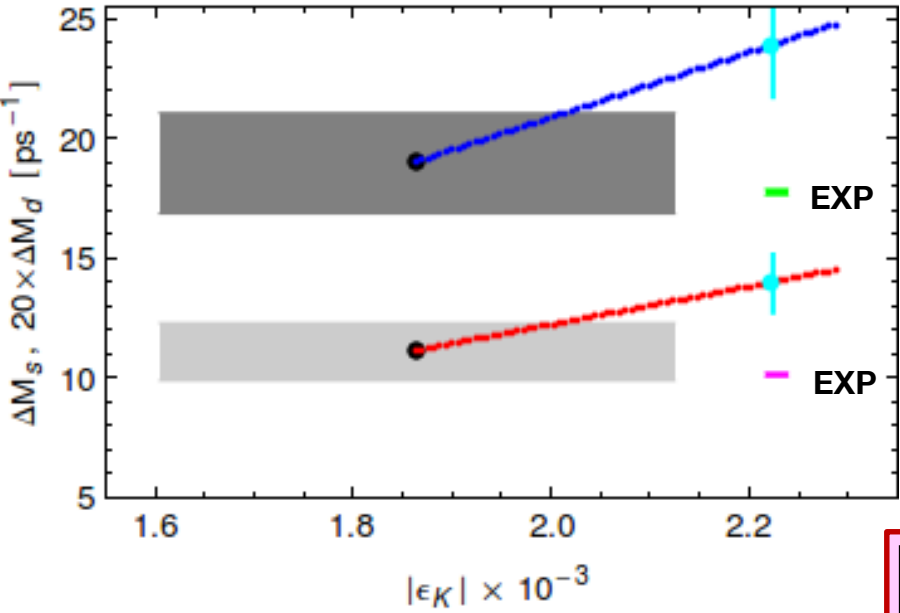
$$S=H^0$$

$$P=A^0$$

2 Tensions in $\Delta F=2$ within MFV

$$\epsilon_K \leftrightarrow \Delta M_{s,d}$$

$$\epsilon_K \leftrightarrow S_{\psi K_s}$$



$$\left\{ |V_{ub}|_{\text{excl}} \right\} \Rightarrow \left\{ \begin{array}{l} \epsilon_K^{\text{SM}} < \epsilon_K^{\text{exp}} \\ S_{\psi K_s}^{\text{SM}} \approx S_{\psi K_s}^{\text{exp}} \end{array} \right\}^* \quad (2\sigma)$$

$$\left\{ |V_{ub}|_{\text{incl}} \right\} \Rightarrow \left\{ \begin{array}{l} \epsilon_K^{\text{SM}} \approx \epsilon_K^{\text{exp}} \\ S_{\psi K_s}^{\text{SM}} > S_{\psi K_s}^{\text{Data}} \end{array} \right\} \quad (3\sigma)$$

$$|V_{cb}|$$

Lunghi + Soni (2008)
 AJB + Guadagnoli (2008)

AJB + Girrbach 1306.3755
 Similar tension in
 Gauged Flavour Models:
 AJB, Merlo, Stamou (2011)

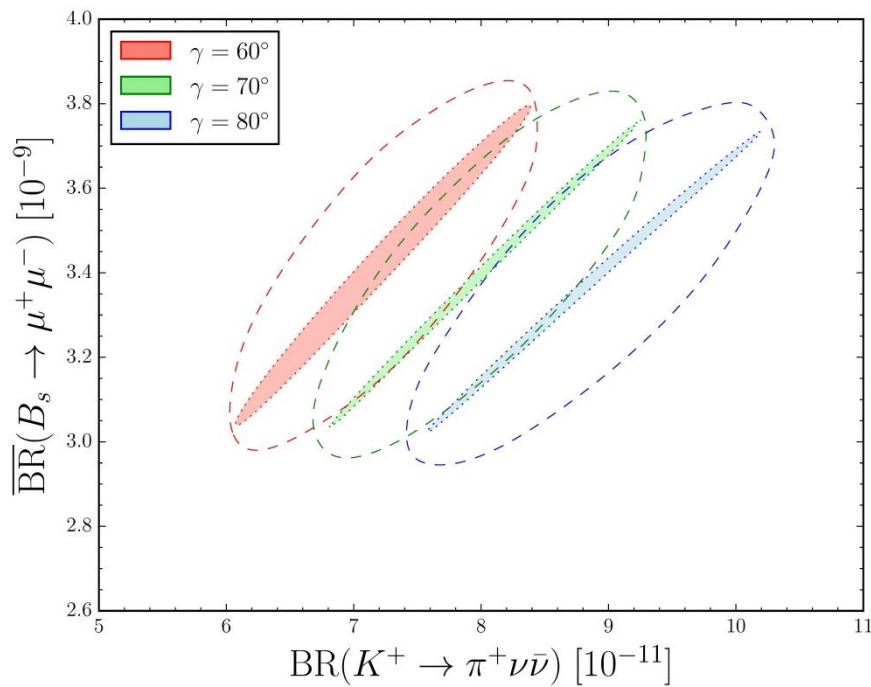
*) Can still work within MFV
 ($\Delta\epsilon_K > 0$ in MFV) Blanke + AJB
 (2006)

Both tensions can only be clarified through improved $|V_{ub}|, |V_{cb}|$ + Lattice Input and improved measurement of $S_{\psi K_s}$

Correlations within SM

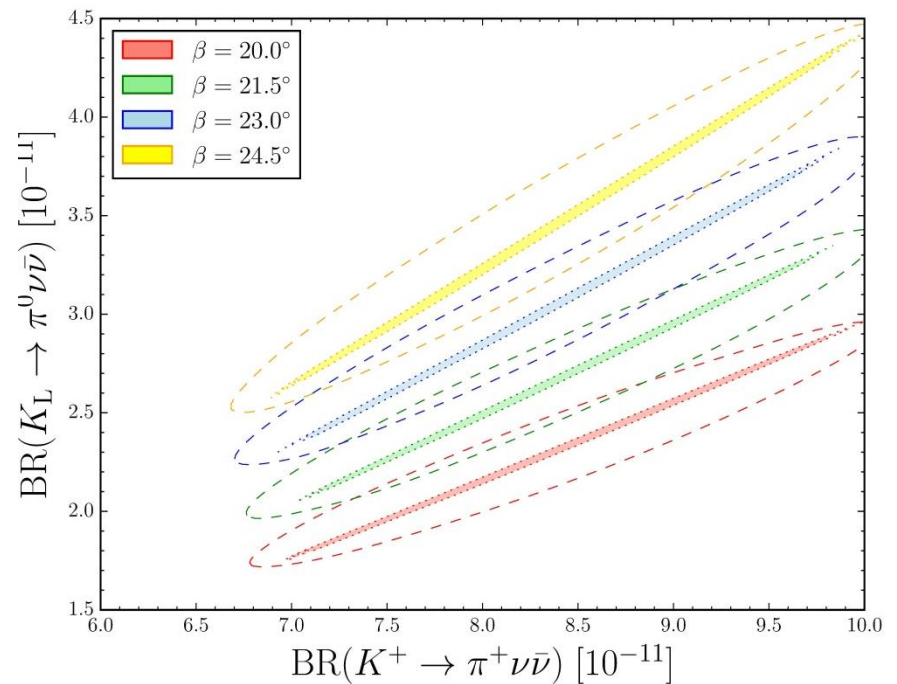
$$B_s \rightarrow \mu^+ \mu^-, K^+ \rightarrow \pi^+ \nu \bar{\nu}, \gamma$$

BBGK (2015)



$$K^+ \rightarrow \pi^+ \nu \bar{\nu}, K_L \rightarrow \pi^0 \nu \bar{\nu}, \beta$$

Buchalla, AJB (94)



General Properties

- 1.** $K^+ \rightarrow \pi^+ \nu \bar{\nu}$ **CP-conserving**
- 2.** $K_L \rightarrow \pi^0 \nu \bar{\nu}$ **CP-violating**
- 3.** **Both sensitive to New Physics (NP)**
 $K^+ \rightarrow \pi^+ \nu \bar{\nu}$ bounded by $K_L \rightarrow \mu^+ \mu^-$
 $K_L \rightarrow \pi^0 \nu \bar{\nu}$ bounded by ε'/ε
- 4.** **The correlation between $K^+ \rightarrow \pi^+ \nu \bar{\nu}$ and $K_L \rightarrow \pi^0 \nu \bar{\nu}$ depends on the ε_K constraint (Blanke 0904.2528)**
- 5.** **Can probe scales far above LHC.**

Strategy B: use ε_K , ΔM_s , ΔM_d , $S_{\psi K_s}$

$$|V_{cb}| = (42.4 \pm 1.0) \cdot 10^{-3}$$

$$|V_{ub}| = (3.61 \pm 0.13) \cdot 10^{-3}$$

$$\gamma = (69.5 \pm 5.0)^\circ \Rightarrow \gamma = (70.8 \pm 2.3)^\circ$$

(after new lattice results for ξ)

$$\begin{aligned} \text{Br}(K^+ \rightarrow \pi^+ \nu \bar{\nu}) &= (9.1 \pm 0.7) \cdot 10^{-11} \\ \text{Br}(K_L \rightarrow \pi^0 \nu \bar{\nu}) &= (3.0 \pm 0.3) \cdot 10^{-11} \end{aligned}$$

$$\text{UTfit} : |V_{cb}| = (41.7 \pm 0.6) \cdot 10^{-3}$$

$$|V_{ub}| = (3.63 \pm 0.12) \cdot 10^{-3}$$

$$\text{CKMfitter} : |V_{cb}| = (41.2 \pm 1.0) \cdot 10^{-3}$$

$$|V_{ub}| = (3.55 \pm 0.16) \cdot 10^{-3}$$

New Bound on $B_6^{(1/2)}$ and $B_8^{(3/2)}$ from Large N

AJB + Gérard 1507.06326

$$B_6^{(1/2)} \leq B_8^{(3/2)} < 1$$



Using BGJJ formula

$$B_6^{(1/2)} = 1.0 \quad B_8^{(3/2)} = 1.0 \quad \Rightarrow \quad (\varepsilon'/\varepsilon)_{SM} = 8.6 \cdot 10^{-4}$$

$$B_6^{(1/2)} = 0.8 \quad B_8^{(3/2)} = 0.8 \quad \Rightarrow \quad (\varepsilon'/\varepsilon)_{SM} = 6.4 \cdot 10^{-4}$$

$$B_6^{(1/2)} = 0.6 \quad B_8^{(3/2)} = 0.8 \quad \Rightarrow \quad (\varepsilon'/\varepsilon)_{SM} = 2.2 \cdot 10^{-4}$$

For $\text{Im}(V_{ts} V_{td}^*) = 1.4 \cdot 10^{-4}$

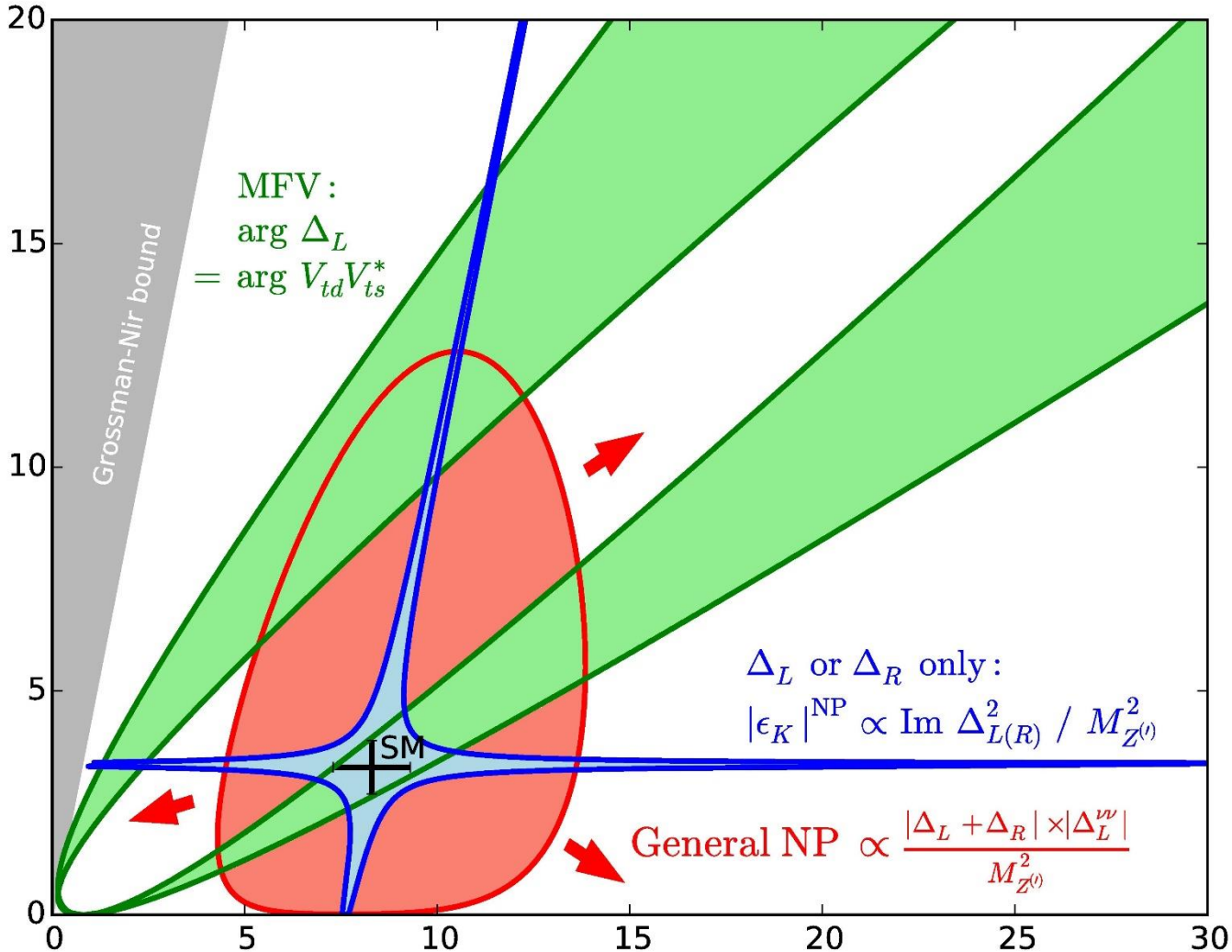
Below data but positive

Yet still large
uncertainties

$K_L \rightarrow \pi^0 \nu \bar{\nu}$ versus $K^+ \rightarrow \pi^+ \nu \bar{\nu}$

AJB, Buttazzo, Knegjens, 1507.08672

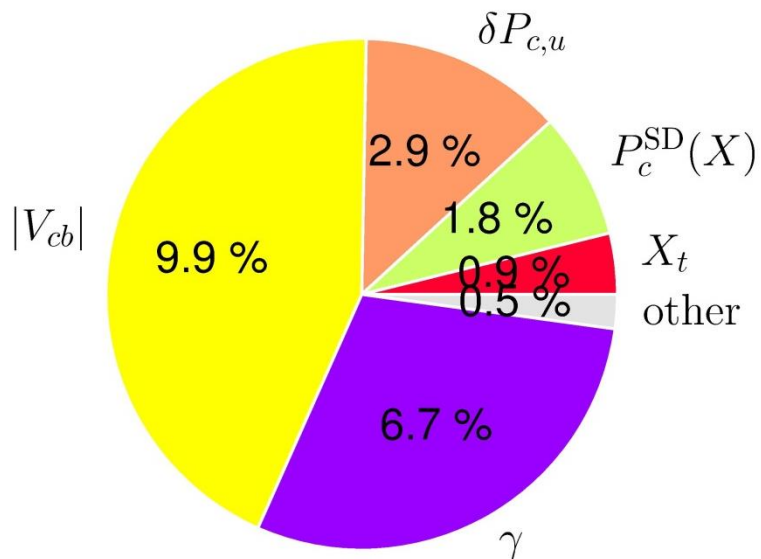
$K_L \rightarrow \pi^0 \nu \bar{\nu}$



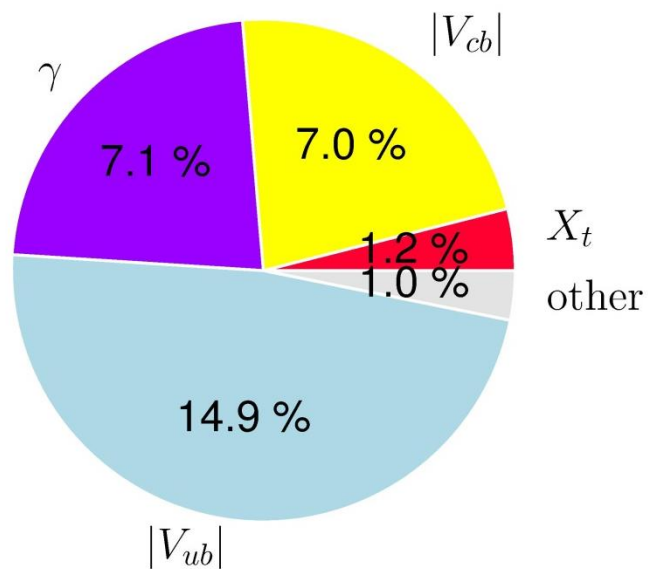
$K^+ \rightarrow \pi^+ \nu \bar{\nu}$

Error Budgets

$$\mathcal{B}(K^+ \rightarrow \pi^+ \nu \bar{\nu})$$



$$\mathcal{B}(K_L \rightarrow \pi^0 \nu \bar{\nu})$$



Update: 1503.02693

$$P_c = 0.404 \pm 0.024$$

$$X_t = 1.481 \pm 0.005_{\text{th}} \pm 0.008_{\text{exp}}$$

Z' outside the reach of the LHC

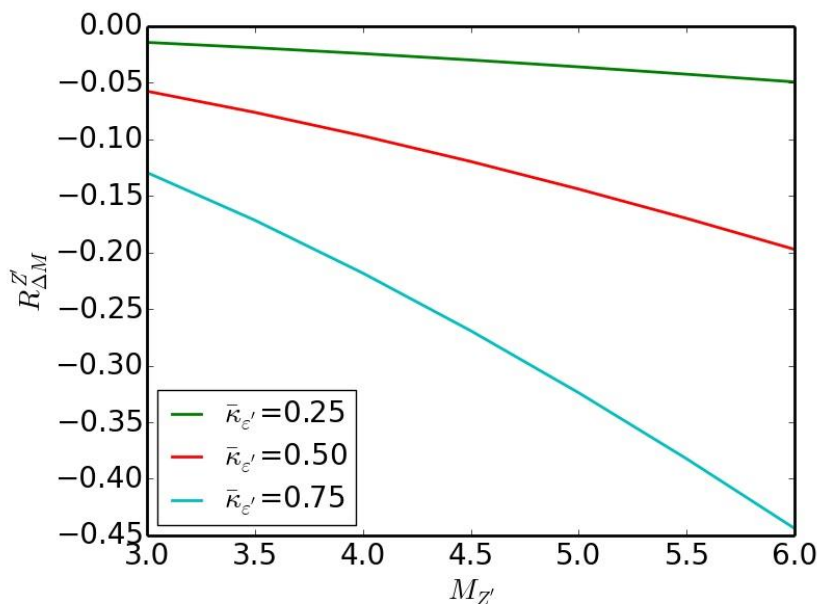
QCD Penguin

For fixed $\bar{\kappa}_{\varepsilon'}$: $\text{Br}(\text{K}_L \rightarrow \pi^0 \nu \bar{\nu})$, $\text{Br}(\text{K}^+ \rightarrow \pi^+ \nu \bar{\nu})$

But constraint
from ΔM_K

Independent of $M_{Z'}$

$Z' q \bar{q} \approx 0(1)$

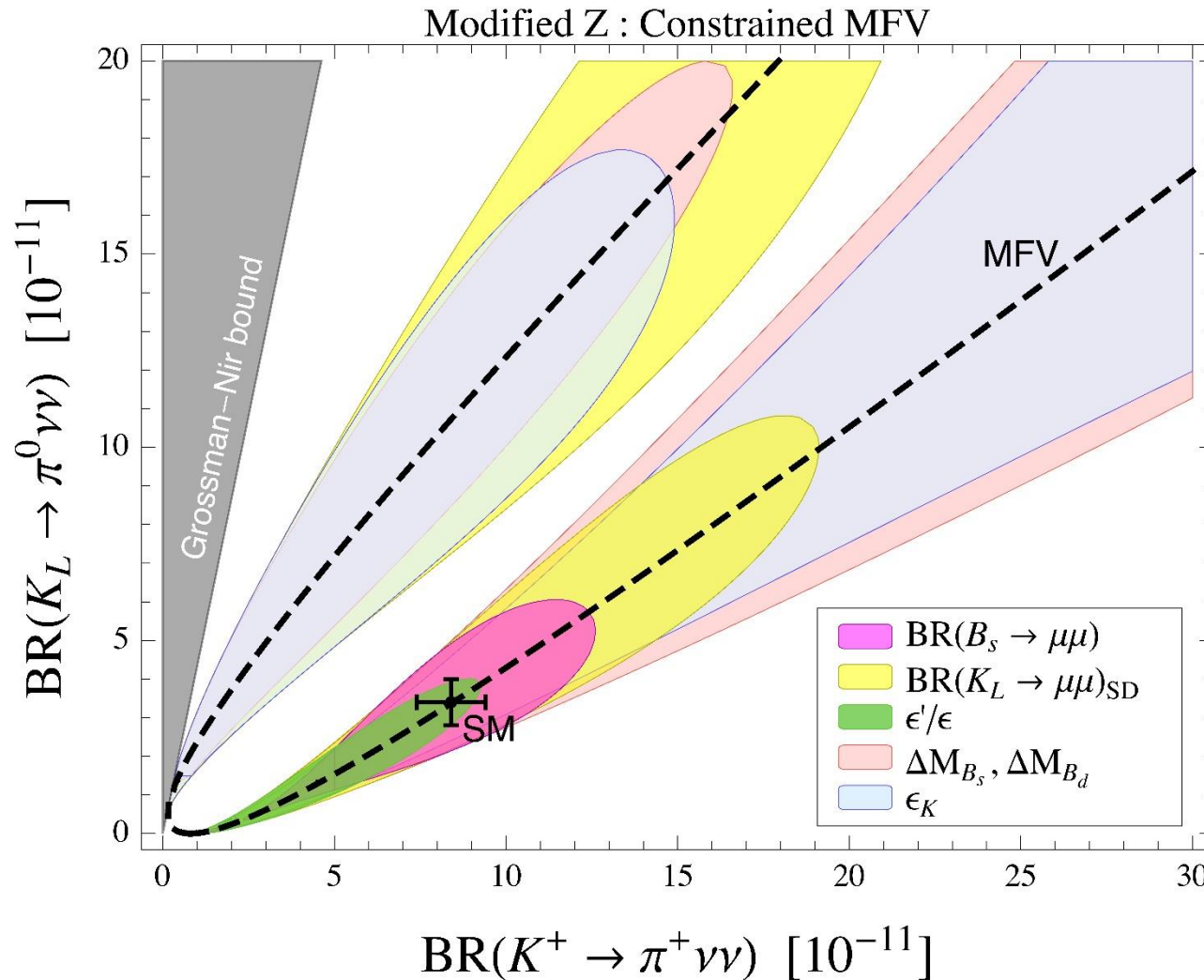


$$\bar{\kappa}_{\varepsilon'} \equiv \left[\frac{\kappa_{\varepsilon'}}{\Delta_R^{\rho\bar{\rho}}(\mathbf{Z}')} \right]$$

EWP Penguin : Significant effects in rare decays
only for $q \bar{q} Z' \approx 0(10^{-2})$

$K^+ \rightarrow \pi^+ \nu \bar{\nu}$ and $K_L \rightarrow \pi^0 \nu \bar{\nu}$ in MFV and $U(2)^3$

AJB + Fleischer (MFV)
0104238



AJB, Buttazzo, Kneijens: hep-ph-1507.08672

Using Tree Level Determination of CKM

$$|V_{ub}|_{\text{excl}} = (3.72 \pm 0.14) \cdot 10^{-3}$$

$$|V_{cb}|_{\text{excl}} = (39.36 \pm 0.75) \cdot 10^{-3}$$

$$|V_{ub}|_{\text{incl}} = (4.40 \pm 0.25) \cdot 10^{-3}$$

$$|V_{cb}|_{\text{incl}} = (42.21 \pm 0.78) \cdot 10^{-3}$$



$$|V_{ub}|_{\text{avg}} = (3.88 \pm 0.29) \cdot 10^{-3}$$

$$|V_{cb}|_{\text{avg}} = (40.7 \pm 1.4) \cdot 10^{-3}$$

$$\gamma = \left(73.2 \begin{matrix} +6.3 \\ -7.0 \end{matrix} \right)^\circ$$

$$\begin{aligned} \overline{\text{Br}}(B_s \rightarrow \mu^+ \mu^-) &= (3.4 \pm 0.3) \cdot 10^{-9} \\ \overline{\text{Br}}(B_s \rightarrow \mu^+ \mu^-)_{\text{exp}} &= (2.8 \pm 0.7) \cdot 10^{-9} \end{aligned}$$

$$\begin{aligned} \text{Br}(K^+ \rightarrow \pi^+ \nu \bar{\nu}) &= (8.4 \pm 1.0) \cdot 10^{-11} \\ \text{Br}(K_L \rightarrow \pi^0 \nu \bar{\nu}) &= (3.4 \pm 0.6) \cdot 10^{-11} \end{aligned}$$



AJB, Buttazzo,
Girrbach-Noe,
Knegjens
1503.02693